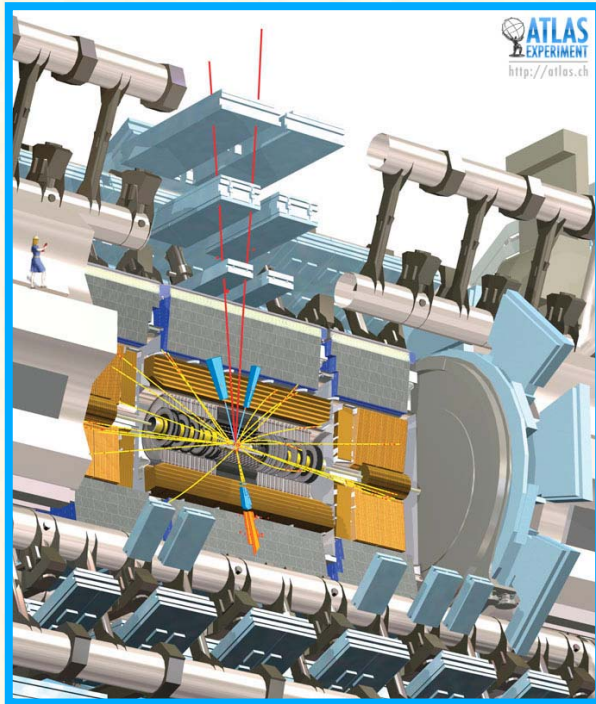


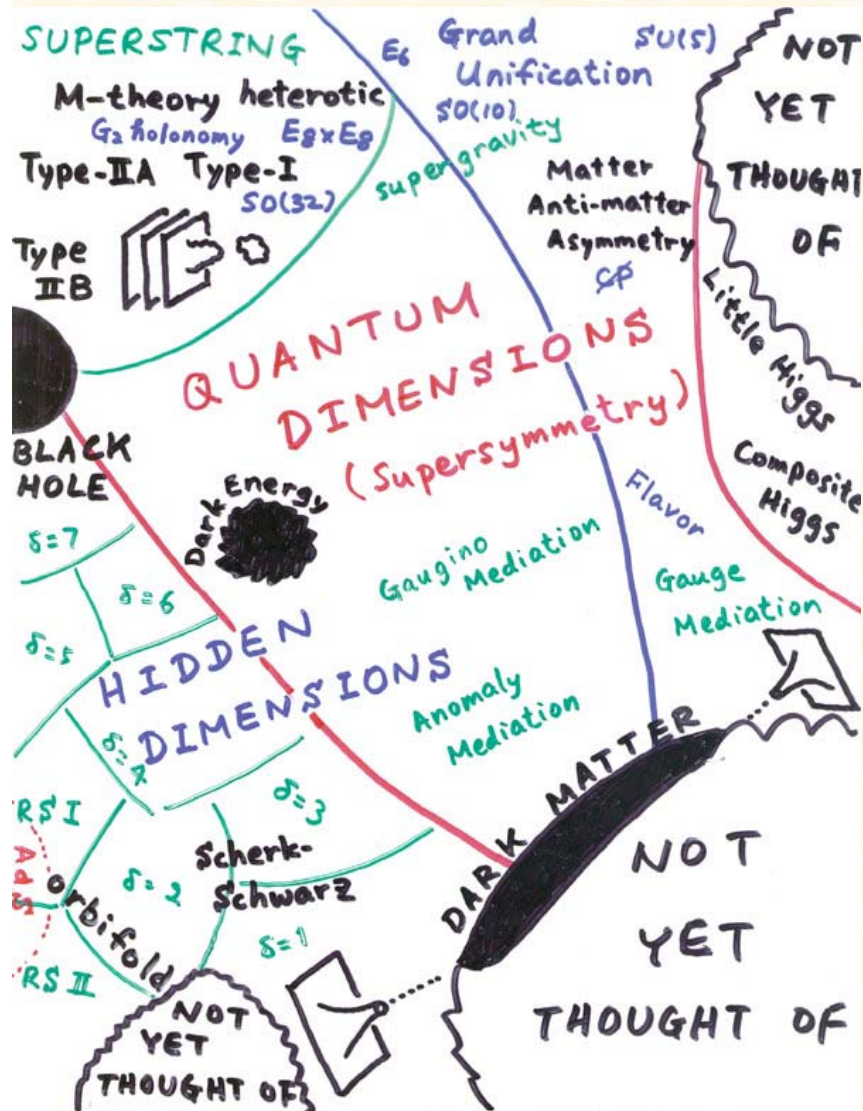
# Higgs boson searches at hadron colliders

## Part 4



- Higgs Bosons in the MSSM
- Invisibly decaying Higgs bosons
- Composite Higgs boson
- WW scattering

# Non-Standard-Model Higgs Scenarios




We really don't know what is going on at the TeV scale....

... be prepared for other scenarios

# Why alternative scenarios?

A. Pomarol, IFAE Barcelona

- „There must be more than a single Higgs.  
Although the Standard Model is *consistent*, it is *not natural*.“
- In quantum field theories it is difficult to protect scalar masses from large radiative corrections



The image shows three Feynman diagrams representing radiative corrections to the Higgs mass. Each diagram consists of a horizontal dashed line representing the Higgs boson. The first diagram has a solid circle loop with two arrows on it, representing a fermion loop. The second diagram has a jagged, sun-like loop representing a scalar loop. The third diagram has a dashed circle loop representing a ghost loop. The diagrams are separated by plus signs and followed by an ellipsis.

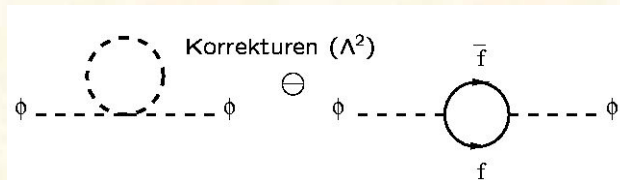
$$\delta m_H^2 = -\frac{\lambda_F^2}{8\pi^2} [\Lambda^2 - m_F^2 \ln \frac{\Lambda^2}{m_F^2}] + \dots$$

Extreme “finetuning” needed, if theory valid up to a high scale  $\Lambda$

# Possible solutions

## (i) An additional symmetry

### - Supersymmetry



$$\Delta m_H = f(m_B^2 - m_f^2)$$

### - The Higgs is not elementary: Composite Higgs

## (ii) Lower the UV scale

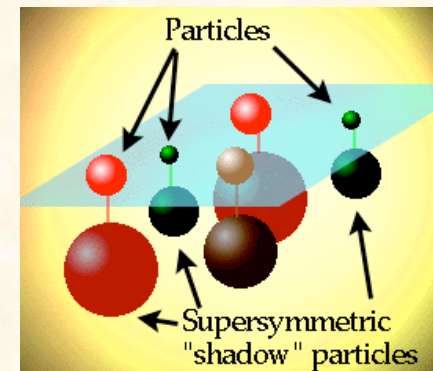
### - Large extra dimensions

## (iii) Remove the Higgs

### - Technicolor models

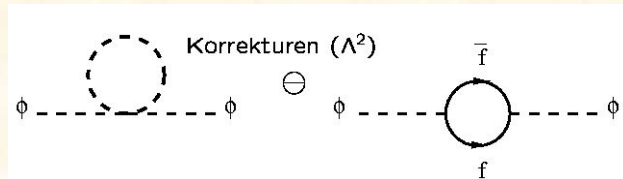
# The Higgs Sector

in the **MSSM**



# Why do we like SUSY so much?

1. Quadratically divergent quantum corrections to the Higgs boson mass are avoided

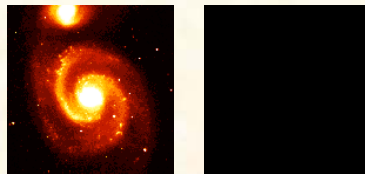


$$\Delta m_H = f(m_B^2 - m_f^2)$$

→  $m_{\text{SUSY}} \sim 1 \text{ TeV}$

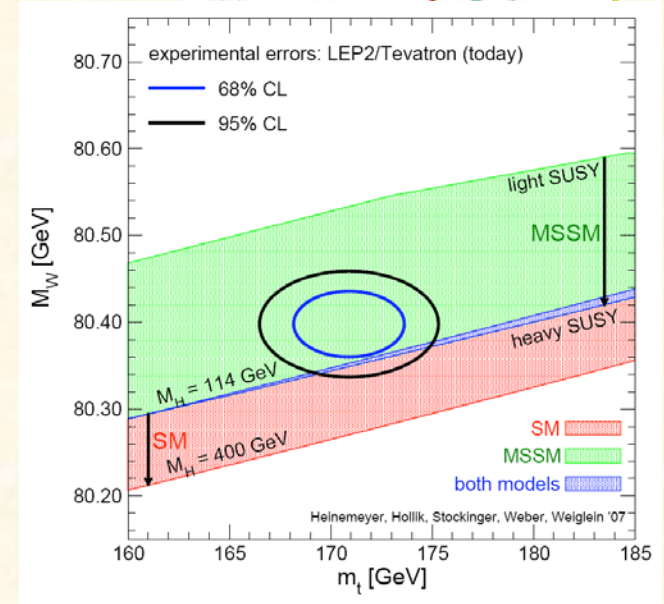
(Hierarchy or naturalness problem)

2. Unification of coupling constants of the three interactions seems possible
3. SUSY provides a candidate for dark matter



**The lightest SUSY particle (LSP)**

4. A SUSY extension is a small perturbation, consistent with the electroweak precision data



# The Higgs Sector in the MSSM

Two Higgs doublets:

5 Higgs particles

H, h, A  
H<sup>+</sup>, H<sup>-</sup>

determined by two parameters:

$m_A, \tan \beta$

fixed mass relations at tree level:

(Higgs self coupling in MSSM fixed by gauge couplings)

$$m_{H,h}^2 = \frac{1}{2} \left( m_A^2 + m_Z^2 \pm \sqrt{(m_A^2 + m_Z^2)^2 - 4m_Z^2 m_A^2 \cos^2 2\beta} \right)$$

$$m_h^2 \leq m_Z^2 \cos^2 2\beta \leq m_H^2$$

Important radiative corrections !! (tree level relations are significantly modified)

→ upper mass bound depends on top mass and mixing in the stop sector

$$m_h^2 \leq m_Z^2 + \frac{3g^2 m_t^4}{8\pi^2 m_W^2} \left[ \ln \left( \frac{M_S^2}{m_t^2} \right) + x_t^2 \left( 1 - \frac{x_t^2}{12} \right) \right]$$

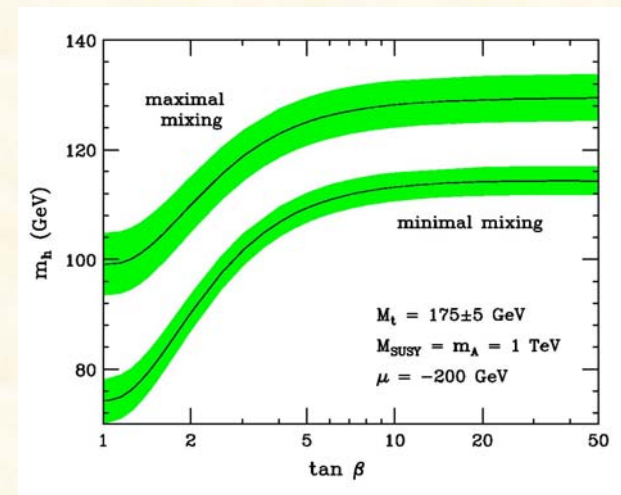
where:  $M_S^2 = \frac{1}{2} (M_{\tilde{t}_1}^2 + M_{\tilde{t}_2}^2)$  and  $x_t = (A_t - \mu \cot \beta) / M_S$

→  $m_h < 115$  GeV for no mixing

→  $m_h < 135$  GeV for maximal mixing

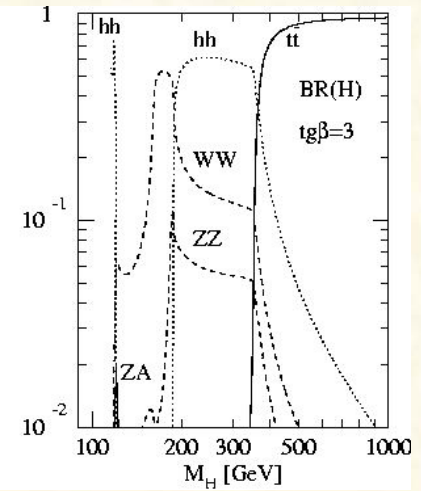
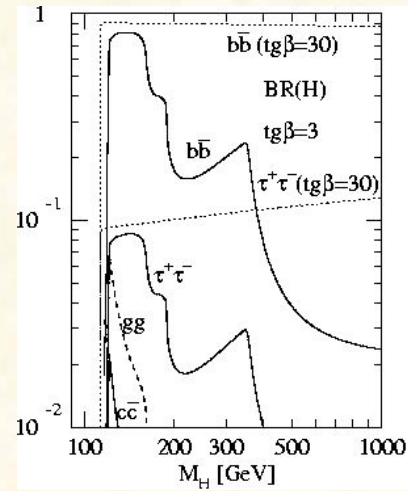
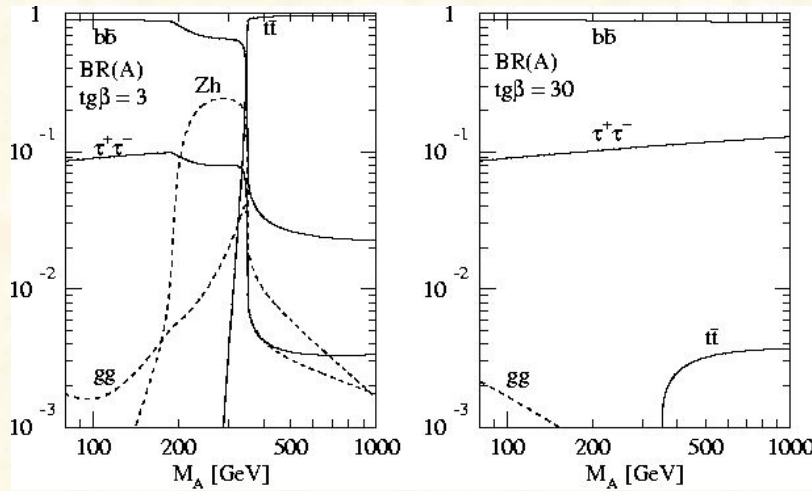
i.e., no mixing scenario: in LEP reach

max. mixing: easier to address at the LHC



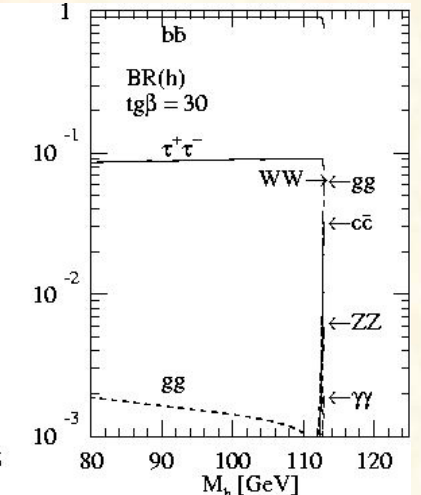
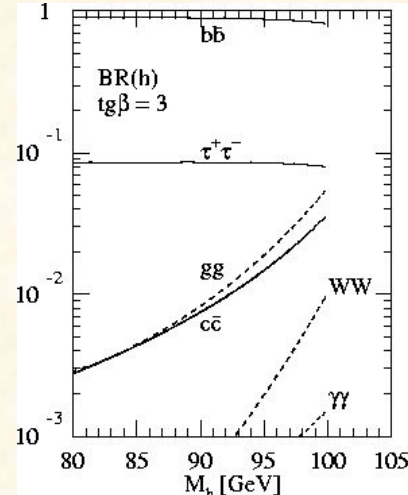
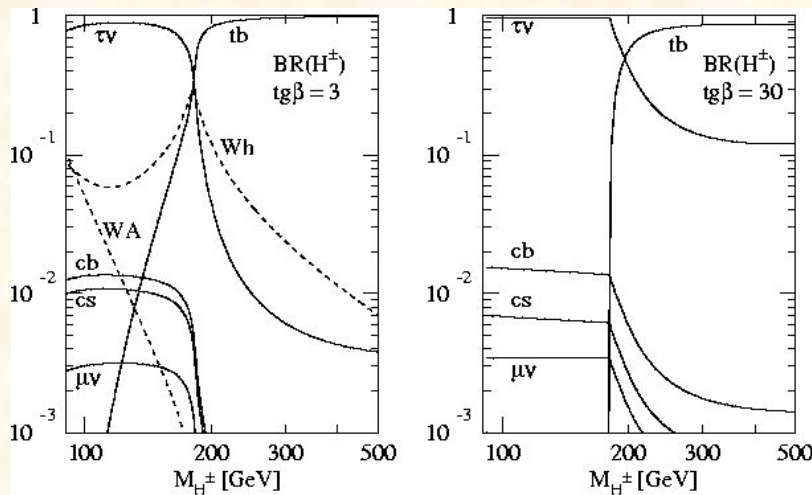
# Branching ratios of MSSM Higgs bosons

A



H

H±



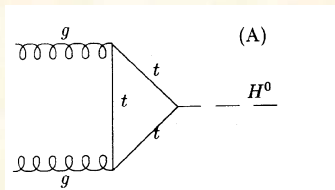
h



# Production of MSSM Higgs bosons

At large  $\tan \beta$ : enhanced couplings of Higgs bosons H and A to down-type fermions

→ important production processes:

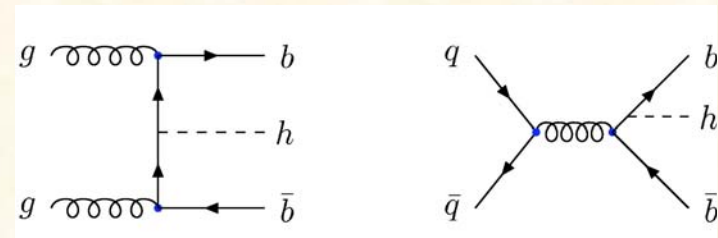


$gg \rightarrow bb \text{ H/A}$

$qq \rightarrow bb \text{ H/A}$

$gg \rightarrow \text{H/A}$

(b, t quarks and **SUSY-particles** in loop)



## Cross section calculation:\*

- associated bbH production becomes dominant process
- NLO calculations are available
  - two approaches –long discussions among theorists-
    - four flavour scheme (bb from gluon splitting)  $K \sim 1.3 - 1.5$  (Tevatron – LHC)
    - five flavour scheme (use b-quark parton distributions,  $bb \rightarrow h$ ,  $gb \rightarrow bh$ )
- Finally: reasonable agreement (within respective scale uncertainties) between the NLO four-<sup>\*\*</sup> and the NNLO five-flavour<sup>\*\*\*</sup> calculation is found for the inclusive (no b-tags) cross section.

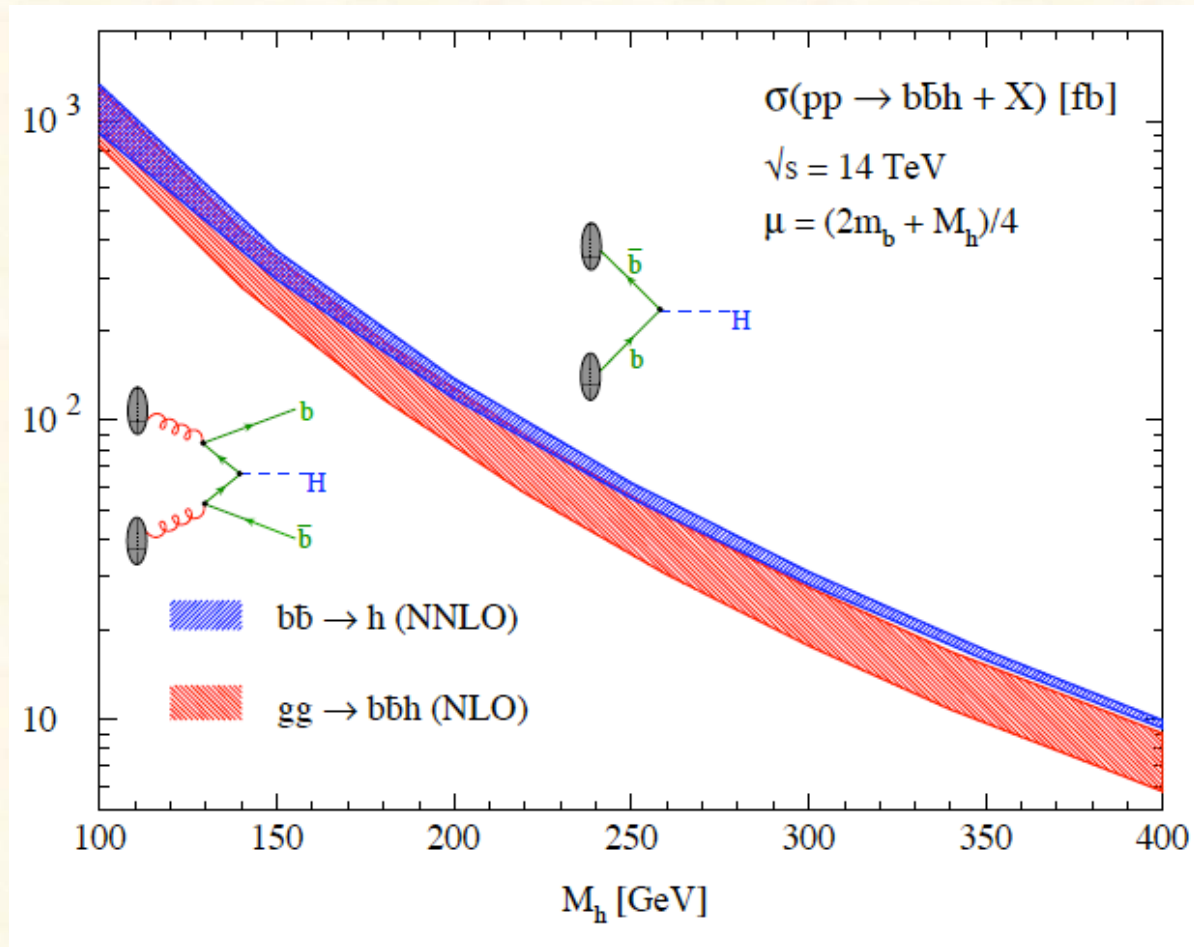
\*) For a review, see: J. Campbell et al. Proc. Les Houches 2003, hep-ph/0405302.

\*\*\*) S. Dittmaier, M. Krämer and M. Spira, Phys. Rev. D70 (2004) 074010; S. Dawson et al., Phys. Rev. D69 (2004) 0740027.

\*\*\*\*) D. Dicus, T. Stelzer, Z. Sullivan and S. Willenbrock, Phys. Rev. D59 (1999) 094016; R. Harlander and W.B. Kilgore Phys. Rev. D68 (2003) 013001.

# Production of MSSM Higgs bosons

Agreement between the four and five flavour scheme calculations:



## MSSM benchmark scenarios

**Masses and couplings** of the Higgs bosons depend –in addition to  $\tan\beta$  and  $m_A$ - on the SUSY parameters through radiative corrections

Most relevant parameters:  $A_t$  = trilinear coupling in the stop sector ( $X_t = A_t - \mu \cot \beta$ )  
 $\mu$  = Higgs mass parameter  
 $M_2$  = gaugino mass term ( $M_1$  from gauge unification)  
 $m_g$  = gluino mass  
 $m_{\text{SUSY}}$  = common scalar mass

**$m_h$ -max:** SUSY parameters chosen such that max mass value for h achieved;  
**No mixing:** vanishing mixing in the stop sector,  $X_t = 0$ ;  
**Gluophobic:** coupling to gluons strongly suppressed, large stop mixing, cancellation between top-quark and stop loop contributions;  
**Small  $\alpha$ :** effective mixing angle between CP-even Higgs bosons is small, reduced BR into bb and  $\tau\tau$  for large  $\tan\beta$  and intermediate values of  $m_A$ .

M. Carena, S. Heinemeyer, C.E.Wagner,  
G. Weiglein, Eur.Phys. J. C26 (2003) 601.

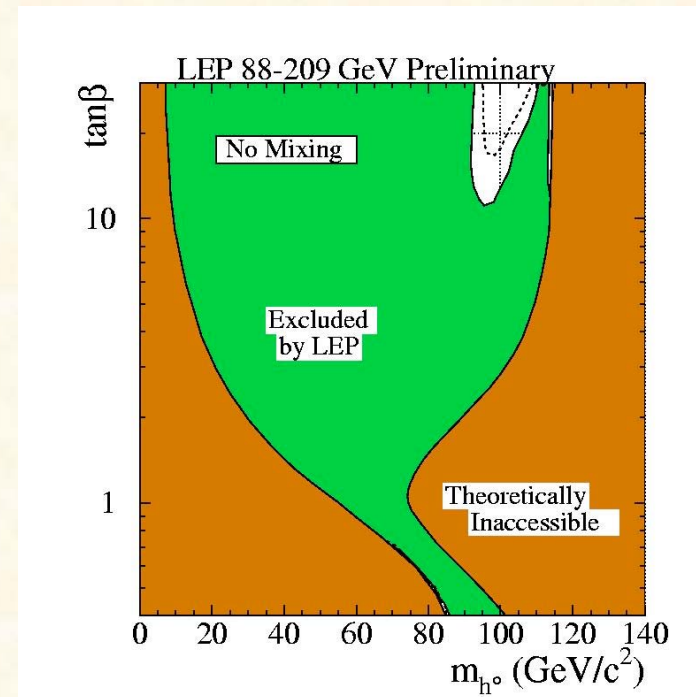
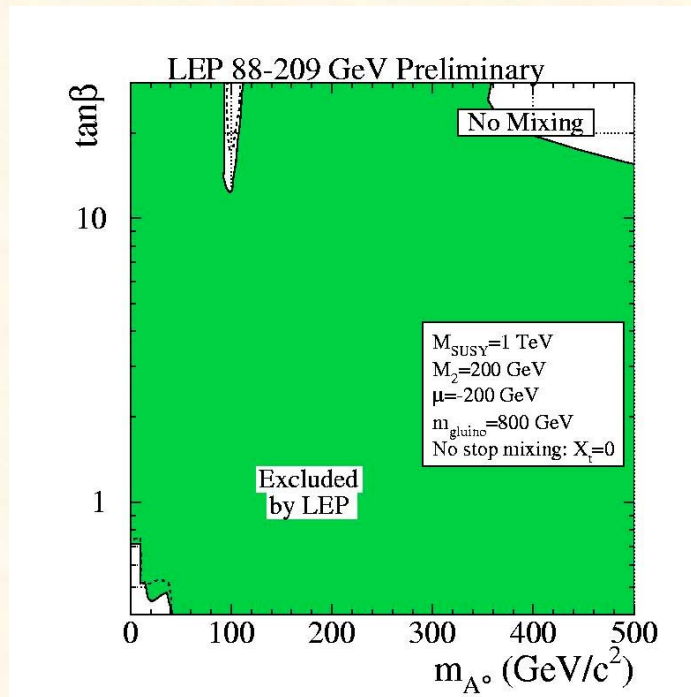
	$m_{\text{SUSY}}$ (GeV/c <sup>2</sup> )	$\mu$ (GeV/c <sup>2</sup> )	$M_2$ (GeV/c <sup>2</sup> )	$X_t$ (GeV/c <sup>2</sup> )	$m_g$ (GeV/c <sup>2</sup> )
$m_h$ -max	1000	200	200	2000	800
No mixing	1000	200	200	0	800
Gluophobic	350	300	300	-750	500
Small $\alpha$	800	2000	500	-1100	500

## LEP results for the no-mixing scenario:

Search for  $e^+e^- \rightarrow h A \rightarrow bb bb$  and  $e^+e^- \rightarrow h Z$   
 $\rightarrow bb \tau\tau$

No significant excess found  $\rightarrow$

set limits in MSSM Higgs boson parameter space ( $M_A$ - $\tan \beta$ )

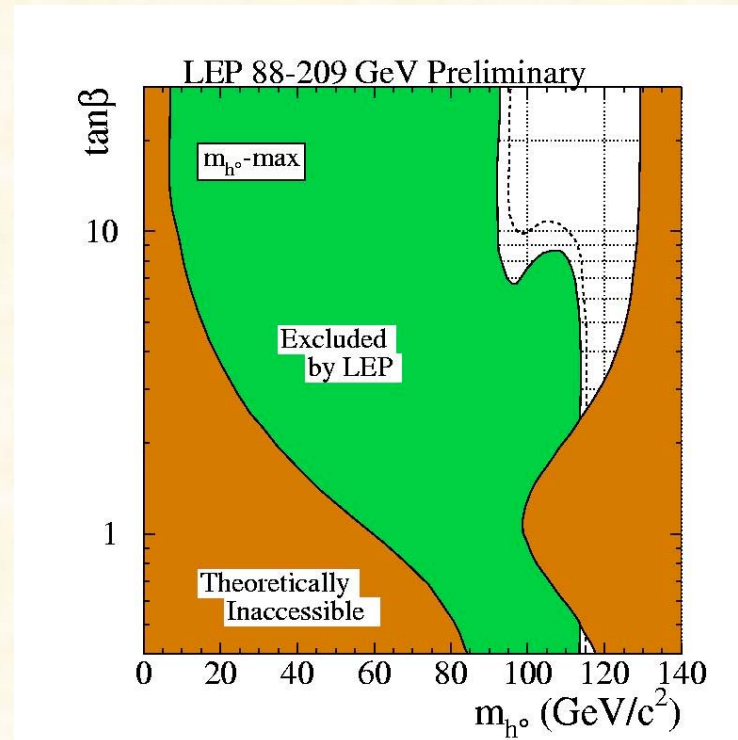
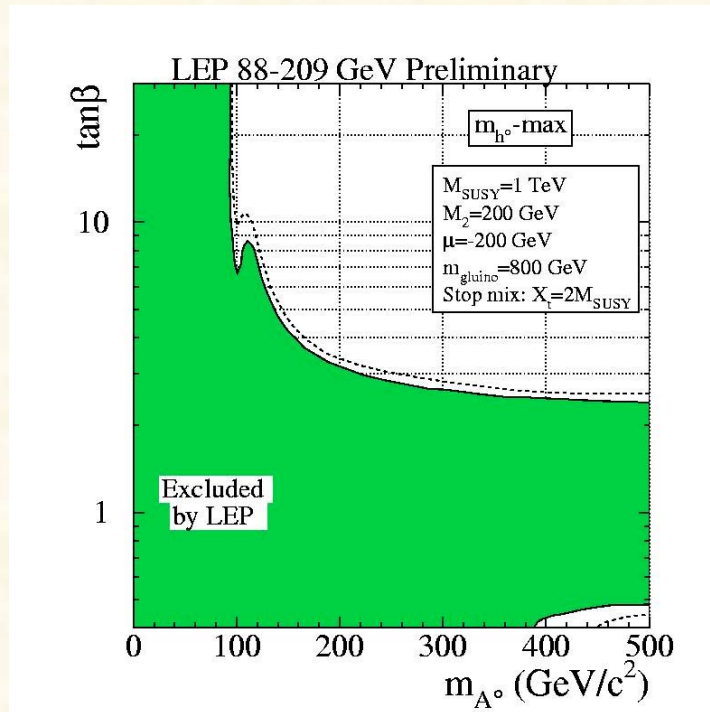


$$M_h > 91.5 \text{ GeV/c}^2$$

$$m_A > 92.2 \text{ GeV/c}^2$$

Excluded  $\tan \beta$  range:  $0.7 < \tan \beta < 10.5$

## Results for the $m_h$ -max scenario:



$$M_h > 91.0 \text{ GeV}/c^2$$

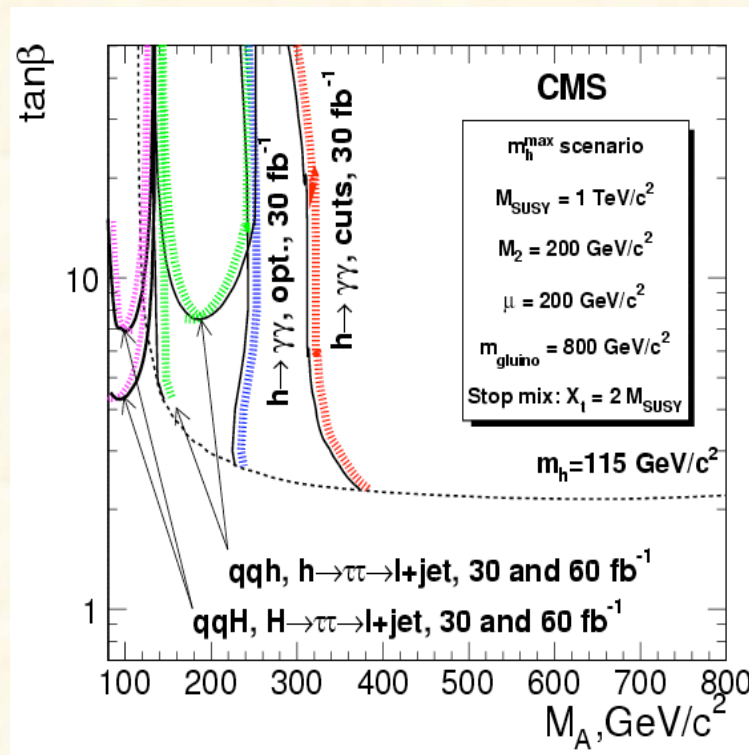
$$m_A > 91.9 \text{ GeV}/c^2$$

Excluded  $\tan\beta$  range:  $0.5 < \tan\beta < 2.4$  ( $m_t = 175 \text{ GeV}/c^2$ )  
 $< 1.9$  ( $m_t = 179 \text{ GeV}/c^2$ )



## Search for the light CP-even Higgs boson $h$

- Standard search channels can be used
- Vector boson fusion channels contribute significantly

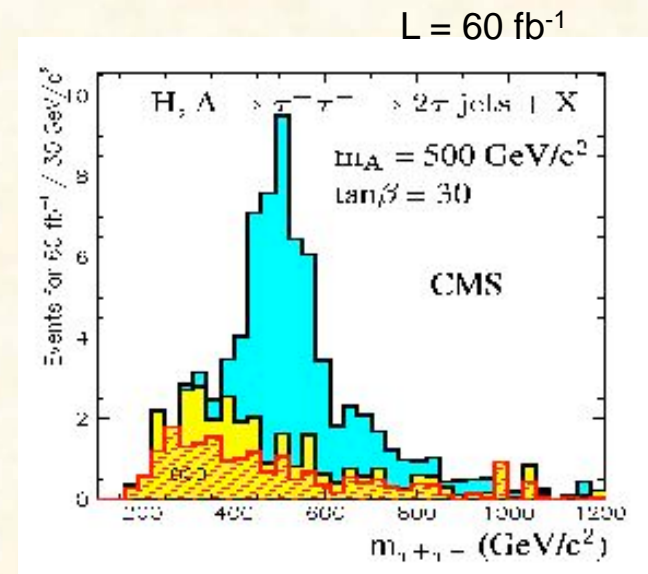


- Uncovered region at small  $m_A$   $\rightarrow$  look for heavier Higgs bosons
- Large integrated luminosities are needed

# Search for the heavy Higgs bosons H and A

at large  $\tan \beta$ : **bb H/A  $\rightarrow$  bb  $\tau\tau$**  plays a key role  
 in addition: bb H/A  $\rightarrow \mu\mu$ ,  
 bb H/A  $\rightarrow bb$  very difficult at the LHC  
 (trigger, backgrounds,.....)

- Selection requires excellent b and  $\tau$  identification
- detailed studies  $\rightarrow$  both leptonic and hadronic tau decays can be used
- $m_{H/A} < 400 \text{ GeV}/c^2$  ( $l - \tau_{had}$ ) dominates
- $> 400 \text{ GeV}/c^2$  ( $\tau_{had} - \tau_{had}$ ) contributes significantly
- H/A mass can be reconstructed, collinear approx.
- Dominant backgrounds: W+jet, tt production

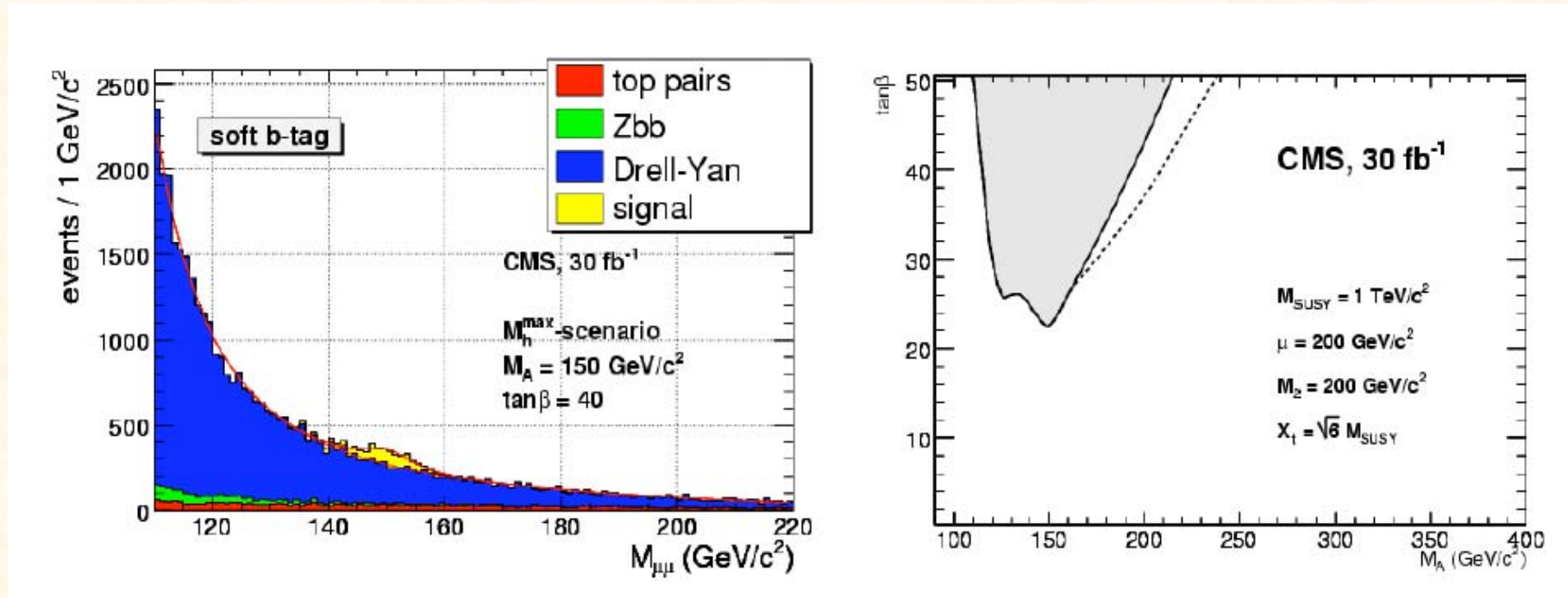


at small  $\tan \beta$ : add. modes: search for H/A  $\rightarrow$  h decays  
 allows for simultaneous observation for two Higgs bosons  
 examples: H  $\rightarrow$  hh, A  $\rightarrow$  Zh



# $A/H \rightarrow \mu\mu$

For MSSM Higgs bosons also the  $\mu\mu$  decay mode can be used (large  $\tan\beta$ )



$5\sigma$  discovery contours for 30 and 100  $\text{fb}^{-1}$

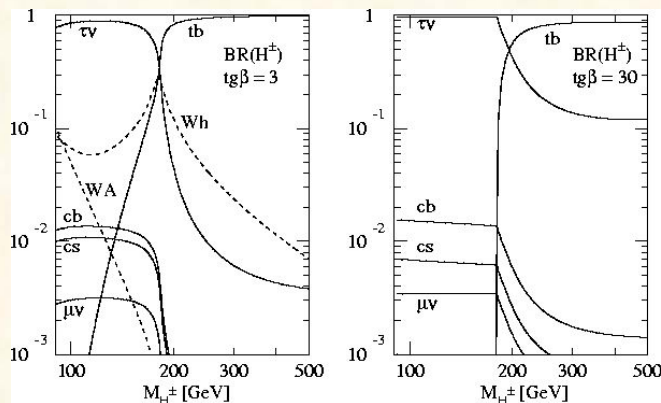
# Search for the Charged Higgs Boson

Detection of a charged Higgs boson → *Physics Beyond the Standard Model*

Production: depends strongly on  $m_{H^\pm}$

- (i) via top decays:  $t \rightarrow H^\pm b$
- (iii)  $gg \rightarrow H^\pm tb$  or  $gb \rightarrow H^\pm t$
- (v)  $gg, qq \rightarrow H^\pm W$
- (iv)  $qq \rightarrow H^+ H^-$

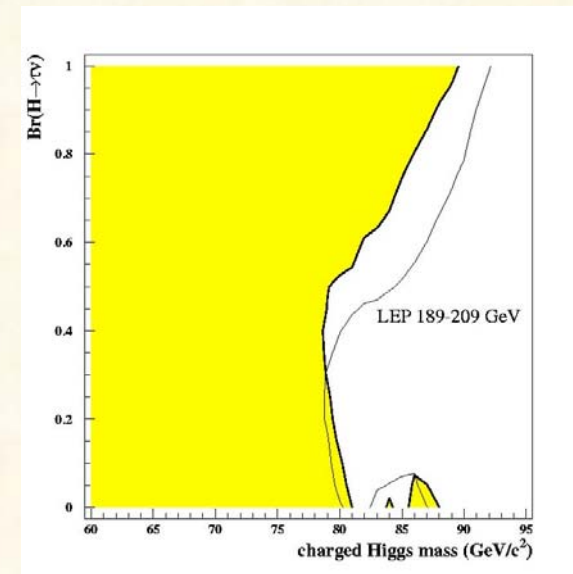
Decays: depend strongly on  $m_{H^\pm}$



$\tau\nu$  decay mode significant at large  $\tan\beta$

LEP results on  $m_{H^\pm}$ :

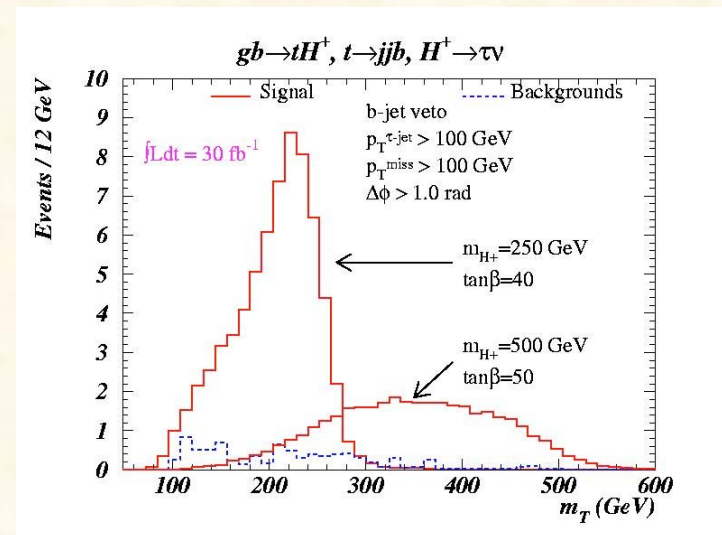
search for  $H^+ \rightarrow cs$  and  $\tau\nu$



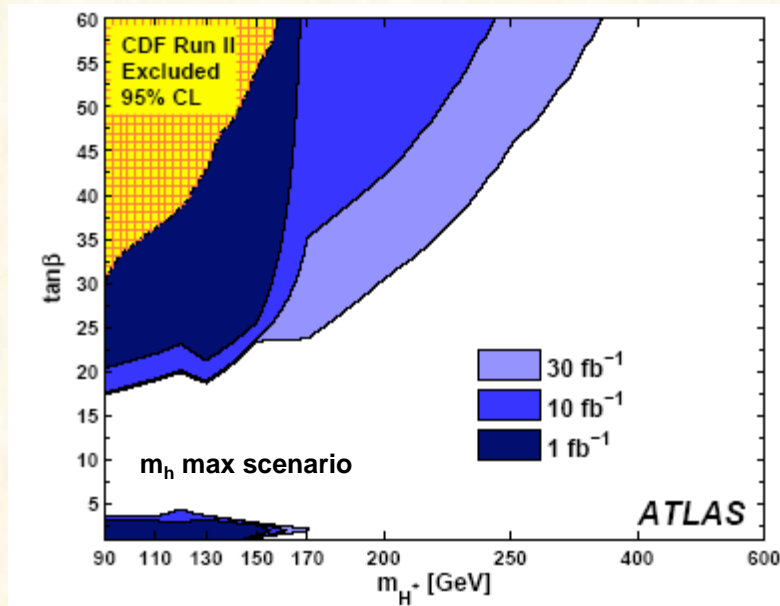
**$M_{H^+} > 78.6 \text{ GeV}/c^2$  (95% CL)**

## Search for the charged Higgs bosons $H^\pm$ at large mass

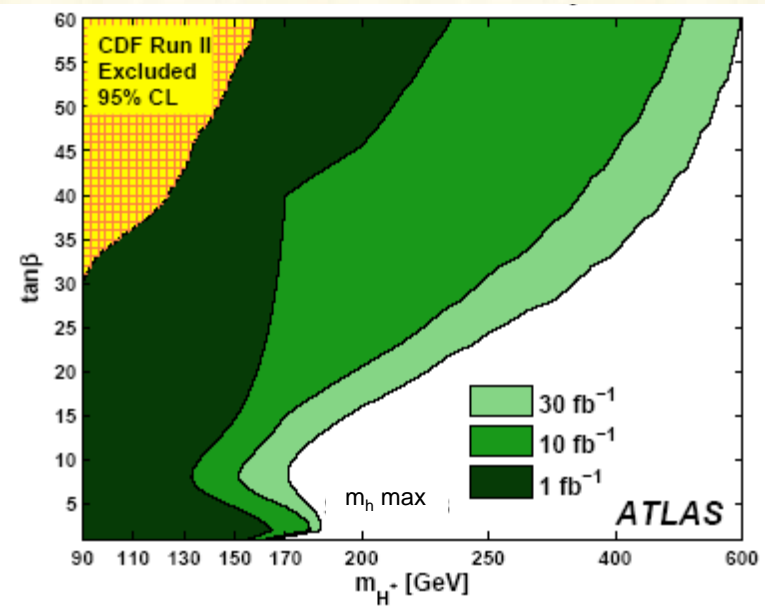
- $H^\pm \rightarrow tb$  decays:
  - promising channel: (i)  $gb \rightarrow H^\pm t \rightarrow tb t \rightarrow \ell\nu bb qqb$   
require **three b-tags** + top-reconstruction  $\rightarrow$  large background suppression
  - more difficult: (ii)  $gg \rightarrow H^\pm tb \rightarrow tb tb \rightarrow \ell\nu bb qqbb$   
require **four b-tags** + top-reconstruction (larger comb. background in rec. of  $H^\pm$ )
- $H^\pm \rightarrow \tau\nu$  decays:
  - promising channel: (i)  $gb \rightarrow H^\pm t \rightarrow \tau\nu t \rightarrow h\nu\nu qqb$   
exploit hadronic decays of the  $\tau$  and  $t$  quark  
 $\rightarrow$  **transverse mass distribution ( $\tau_{\text{had}} + E_T^{\text{miss}}$ )**  
**can be used to reconstruct the  $H^\pm$  mass**
  - additional channel:  
(ii)  $gg \rightarrow H^\pm tb \rightarrow \tau\nu tb \rightarrow h\nu\nu qqbb$   
require **two b-tags**
- Other decay modes ( $H^\pm \rightarrow Wh, WH$ )  
marginal - hopeless in MSSM



## Some examples of updated MSSM studies



$5\sigma$  discovery contours for 1 to 30  $\text{fb}^{-1}$



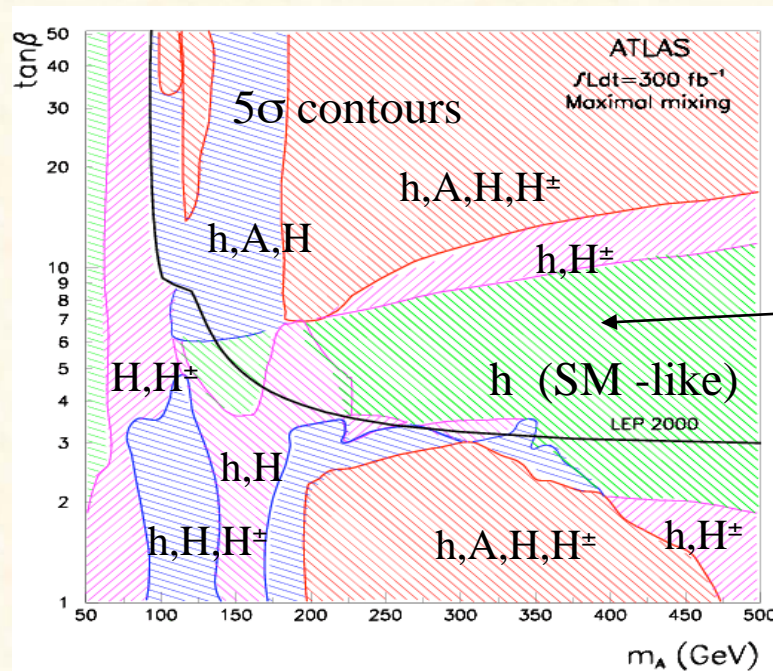
95% CL exclusions for 1 to 30  $\text{fb}^{-1}$

ATLAS: Charged Higgs boson searches

$H^\pm \rightarrow \tau\nu$  and  $tb$  decay modes



# LHC discovery potential for MSSM Higgs bosons



- 4 Higgs observable
- 3 Higgs observable
- 2 Higgs observable
- 1 Higgs observable

Here only SM-like  $h$  observable if SUSY particles neglected.

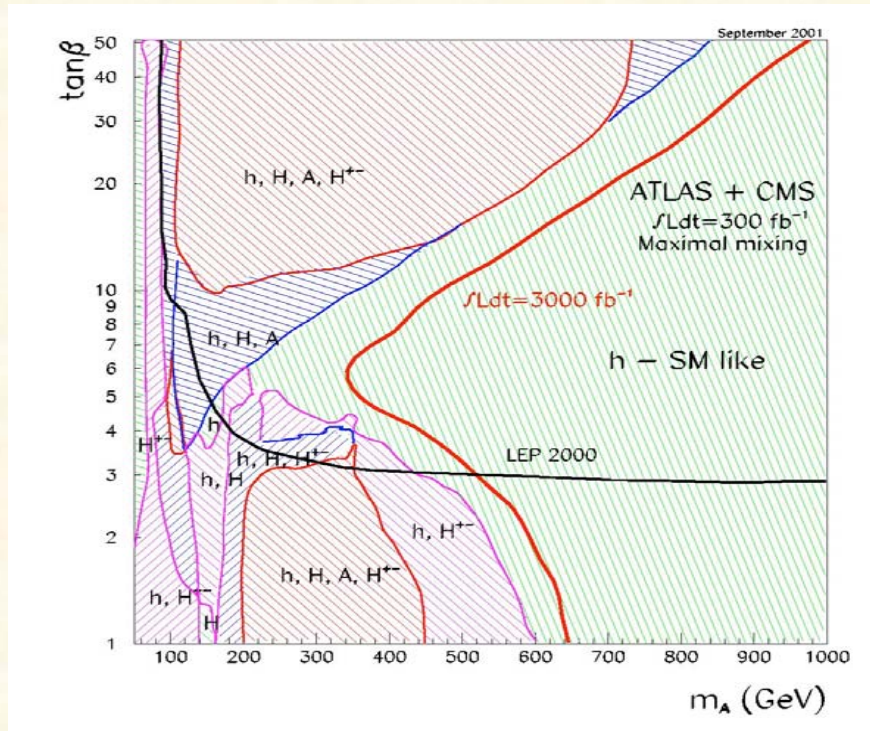
Assuming decays to SM particles only

- Region at large  $m_A$  and moderate  $\tan \beta$  only covered by  $h$ ; difficult to detect other Higgs bosons

Possible coverage: \* via SUSY decays (model dependent, see below)  
 \* luminosity (only moderate improvement)

# MSSM discovery potential for Super-LHC

ATLAS + CMS,  $2 \times 3000 \text{ fb}^{-1}$

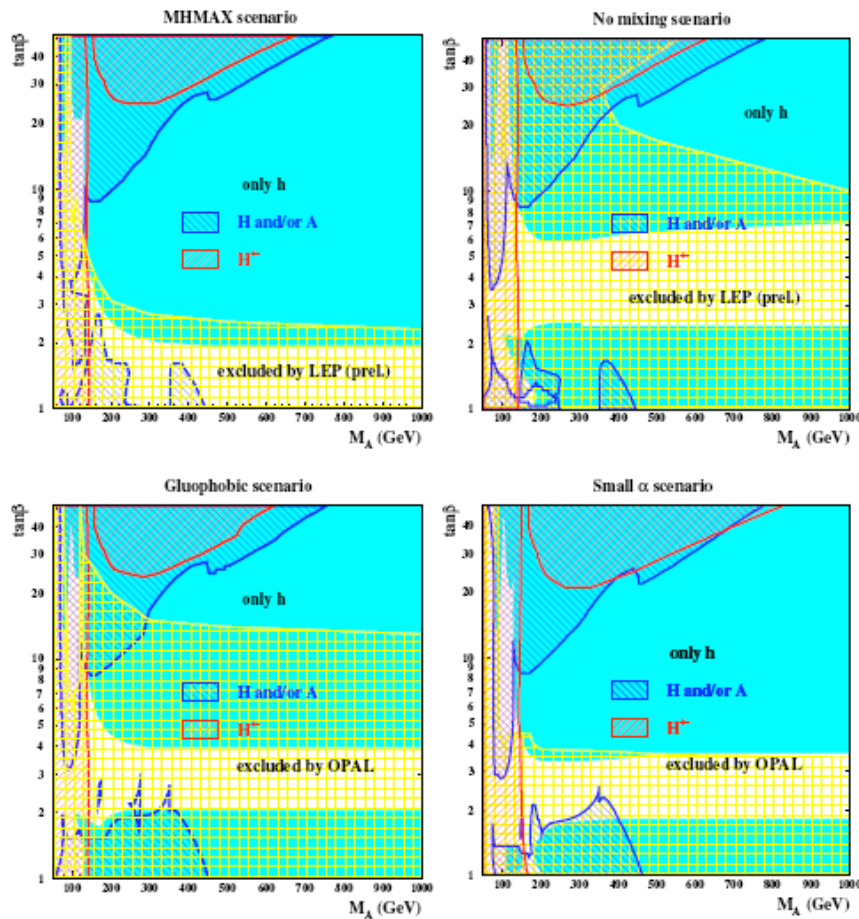


- Situation can be improved, in particular for  $m_A < \sim 400 \text{ GeV}$
- But: (s)LHC can not promise a complete observation of the heavy part of the MSSM Higgs spectrum ....  
.... although the observation of sparticles will clearly indicate that additional Higgs bosons should exist.

# Updated MSSM scan for different benchmark scenarios

Benchmark scenarios as defined by M.Carena et al. (h mainly affected)

**ATLAS preliminary, 30 fb<sup>-1</sup>, 5σ discovery**



**MHMAX scenario** ( $M_{\text{SUSY}} = 1 \text{ TeV}/c^2$ )  
maximal theoretically allowed region for  $m_h$

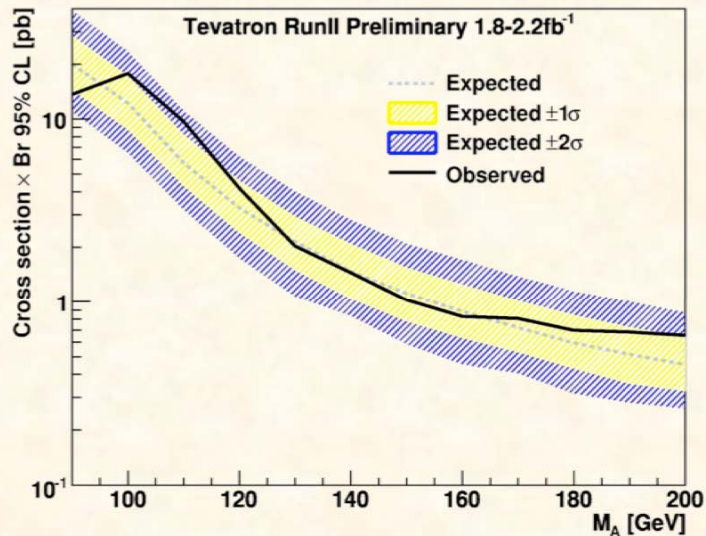
**Nomixing scenario** ( $M_{\text{SUSY}} = 2 \text{ TeV}/c^2$ )  
(1TeV almost excl. by LEP )  
small  $m_h \rightarrow$  difficult for LHC

**Gluophobic scenario** ( $M_{\text{SUSY}} = 350 \text{ GeV}/c^2$ )  
coupling to gluons suppressed  
(cancellation of top + stop loops)  
small rate for  $g g \rightarrow H$ ,  $H \rightarrow \gamma\gamma$  and  $Z \rightarrow 4 \ell$

**Small  $\alpha$  scenario** ( $M_{\text{SUSY}} = 800 \text{ GeV}/c^2$ )  
coupling to b (and t) suppressed  
(cancellation of sbottom, gluino loops) for  
large  $\tan\beta$  and  $M_A$  100 to 500  $\text{GeV}/c^2$

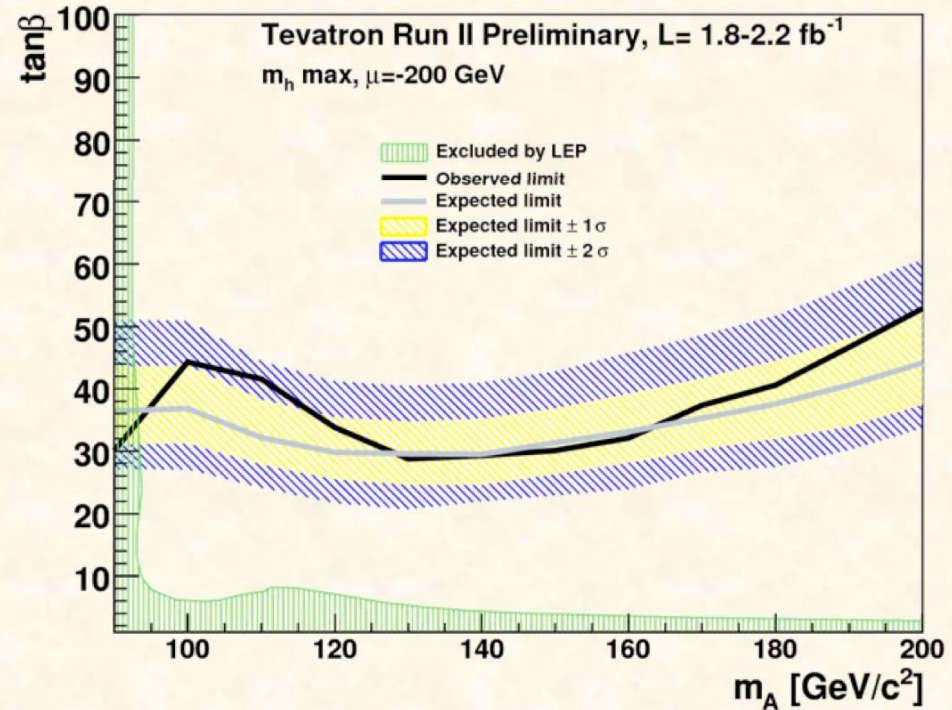


# Tevatron exclusions on MSSM Higgs bosons

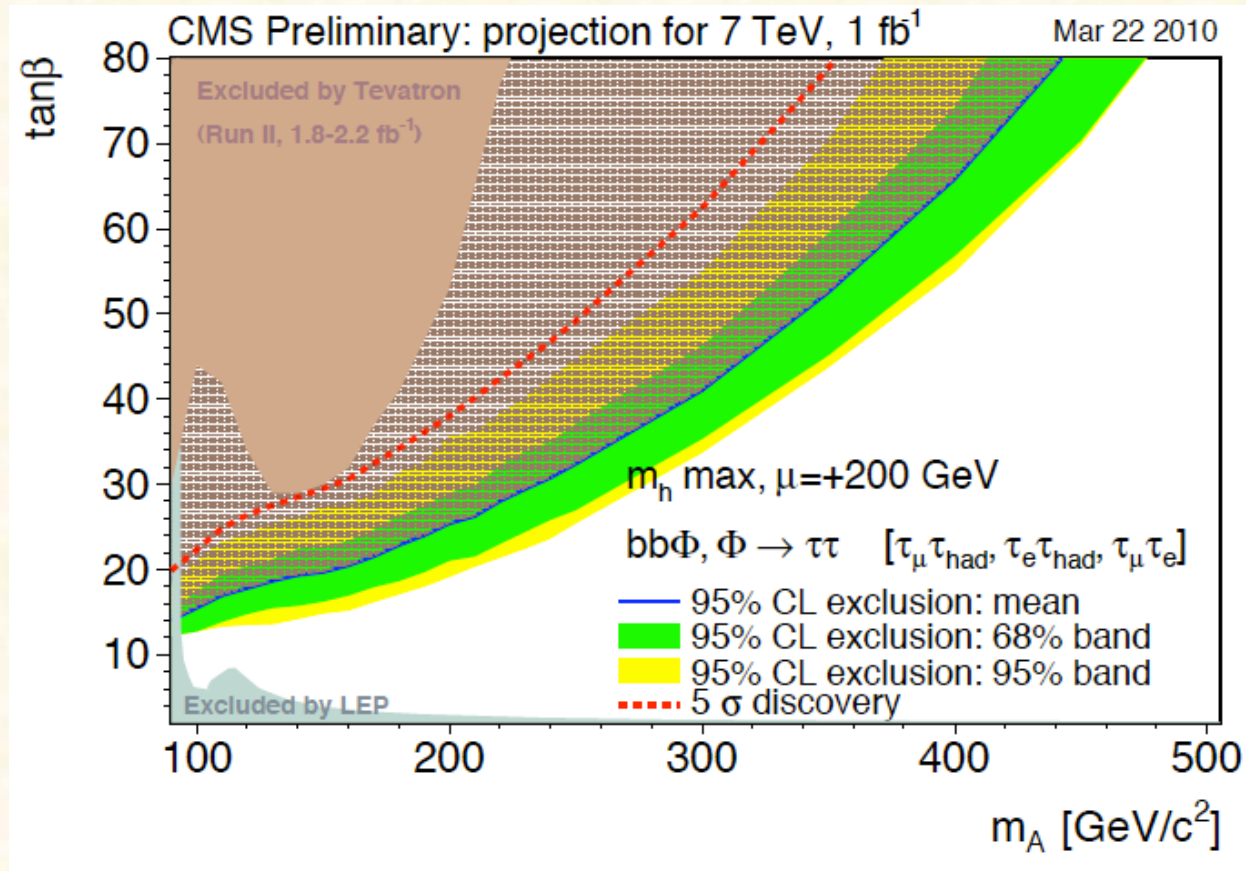


Excluded cross section (95% C.L. limits)

Combination of the CDF and D0 results on  $bb\Phi$ ,  $\Phi \rightarrow \tau\tau$



## Expectations at the LHC for $\sqrt{s} = 7 \text{ TeV}$ and $1 \text{ fb}^{-1}$



- Seems competitive at low mass
- Higher masses ....for later

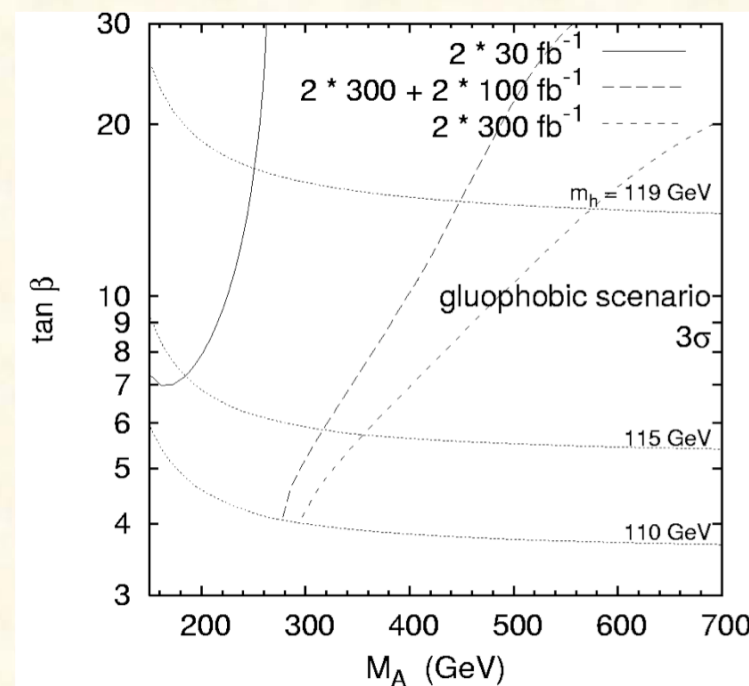
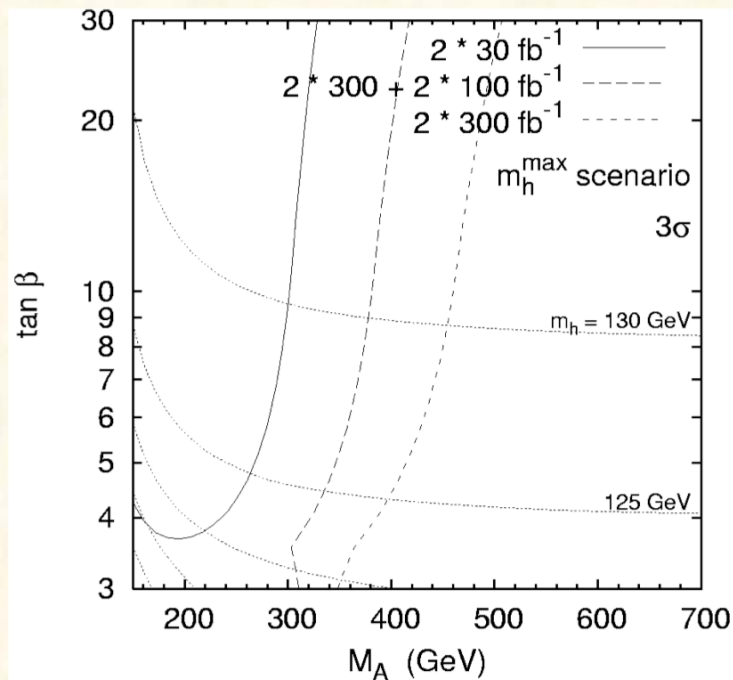
# A light SUSY Higgs or a light Standard Model Higgs?

# Discrimination between a SM or MSSM Higgs

In some regions of MSSM parameter space only one light Higgs boson is visible

⇒ Try to exclude MSSM using a  $\chi^2$  analysis of coupling fits

M. Dührssen et al., Phys. Rev. D70 (2004) 113009.



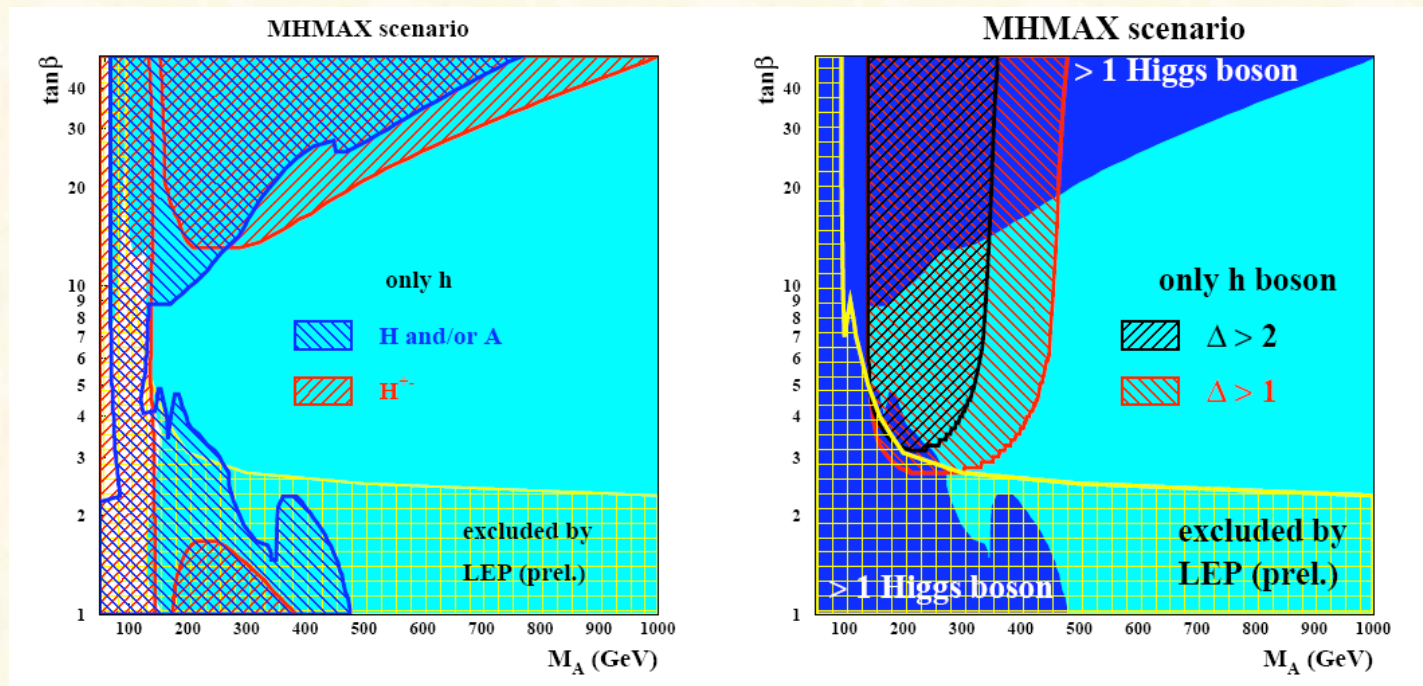
## Discrimination between a SM or MSSM Higgs (cont.)

Similar analysis based on direct comparison of ratios of rates in different final states, using VBF production

M. Schumacher et al., hep-ph/0410112

Consider variables  $R := \text{BR}(h \rightarrow \tau\tau) / \text{BR}(h \rightarrow WW)$

$$\Delta := (R_{\text{MSSM}} - R_{\text{SM}}) / \sigma_{\text{exp}}$$



30 fb<sup>-1</sup>

(only stat. errors considered so far,  $m_H$  assumed to be known with high precision)

Can **SUSY particles** be used to detect  
Higgs bosons ??

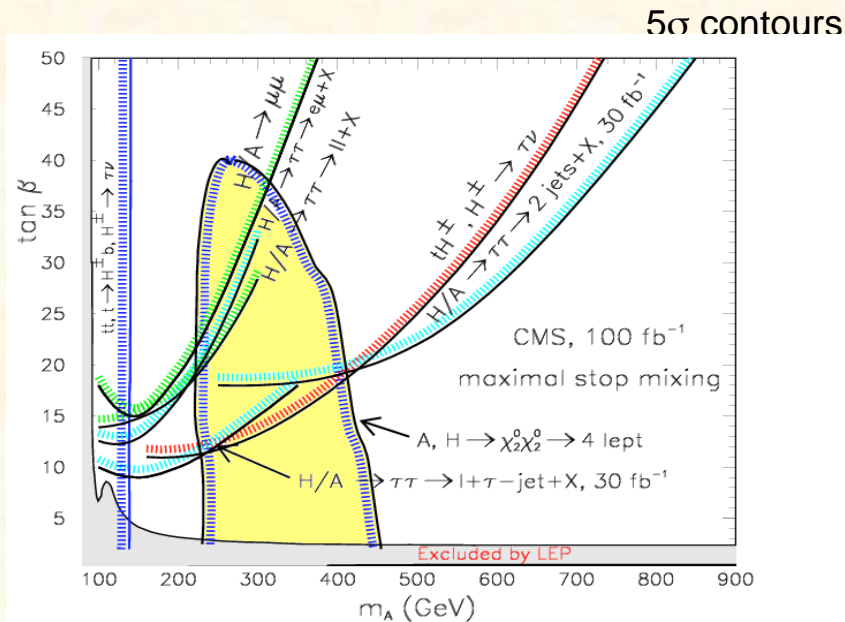
or

the interplay between the  
Higgs sector and SUSY particles

so far: SUSY particles have been assumed to be  
too heavy to play a role in Higgs boson decay phenomenology

# CMS study: MSSM scenario

$$H/A \rightarrow \chi_2^0 \chi_2^0 \rightarrow \ell\ell\chi_1^0 \ell\ell\chi_1^0$$



special choice in MSSM (no scan)

$$M_1 = 60 \text{ GeV}$$

$$M_2 = 110 \text{ GeV}$$

$$\mu = -500 \text{ GeV}$$

Exclusions depend on MSSM parameters  
(slepton masses,  $\mu$ )

# Search for $H^\pm$ decays into SUSY particles

$$gb \rightarrow tH^\pm, H^\pm \rightarrow \chi_{2,3}^0 \chi_{1,2}^\pm \rightarrow 3\ell + E_T^{miss}$$

special choice in MSSM (no scan)

$$M_2 = 210 \text{ GeV}$$

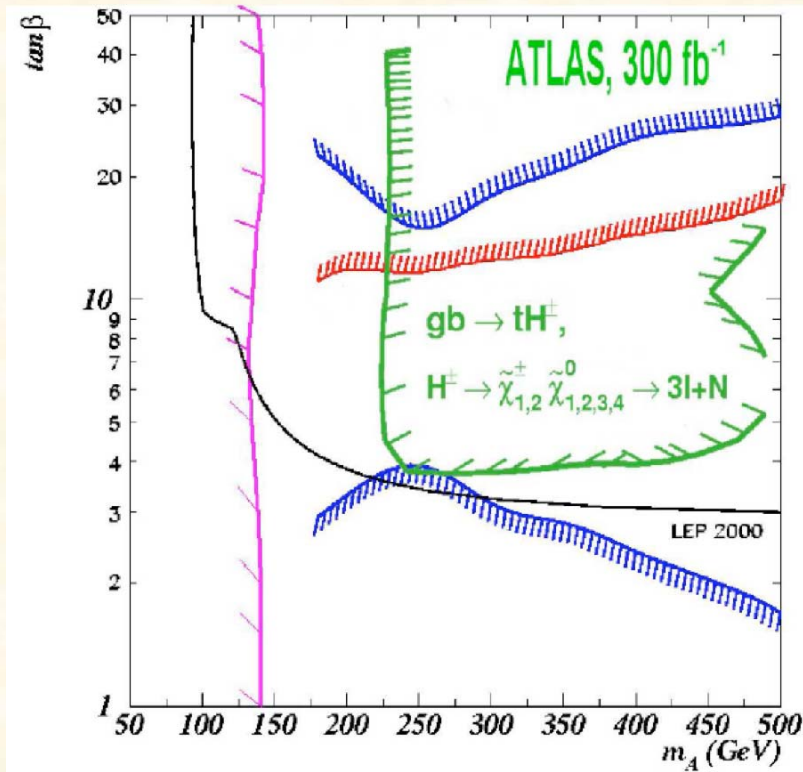
$$\mu = 135 \text{ GeV}$$

$$m(s-\ell_R) = 110 \text{ GeV}$$

$$m(s-\tau_R) = 210 \text{ GeV}$$

$$m_g = 800 \text{ GeV}$$

$$m_{\text{SUSY}} = 1000 \text{ GeV}$$



complementary discovery potential



## Excursion: Search for Supersymmetry

- If **SUSY** exists at the electroweak scale, a discovery at the LHC should be easy
- **Squarks** and **Gluginos** are strongly produced

They decay through cascades to the lightest SUSY particle (LSP)

⇒ **Heavy Higgs bosons might appear in cascade decays,**

**e.g. cascade decays of squarks and gluinos via heavy charginos and neutralinos**

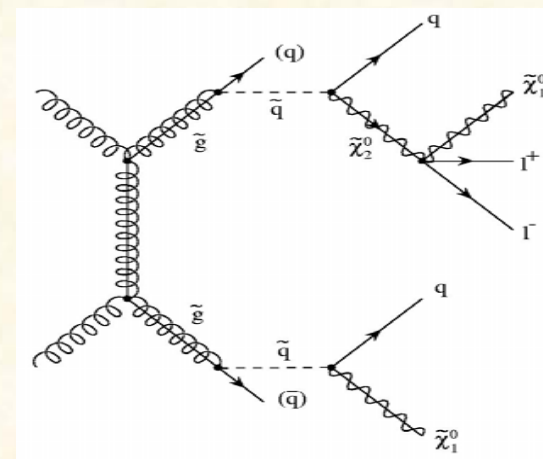
$$pp \rightarrow \tilde{g}\tilde{g}, \tilde{q}\tilde{q}, \tilde{q}\tilde{g} \rightarrow \chi_2^\pm, \chi_3^0, \chi_4^0 + X$$

$$\rightarrow \chi_1^\pm, \chi_2^0, \chi_1^0 + h, H, A, H^\pm + X$$

$$pp \rightarrow \tilde{g}\tilde{g}, \tilde{q}\tilde{q}, \tilde{q}\tilde{g} \rightarrow \chi_1^\pm, \chi_2^0 + X$$

$$\rightarrow \chi_1^0 + h, H, A, H^\pm + X$$

**Search for Higgs decays in standard channels:  $bb, \tau\tau, \tau\nu$**



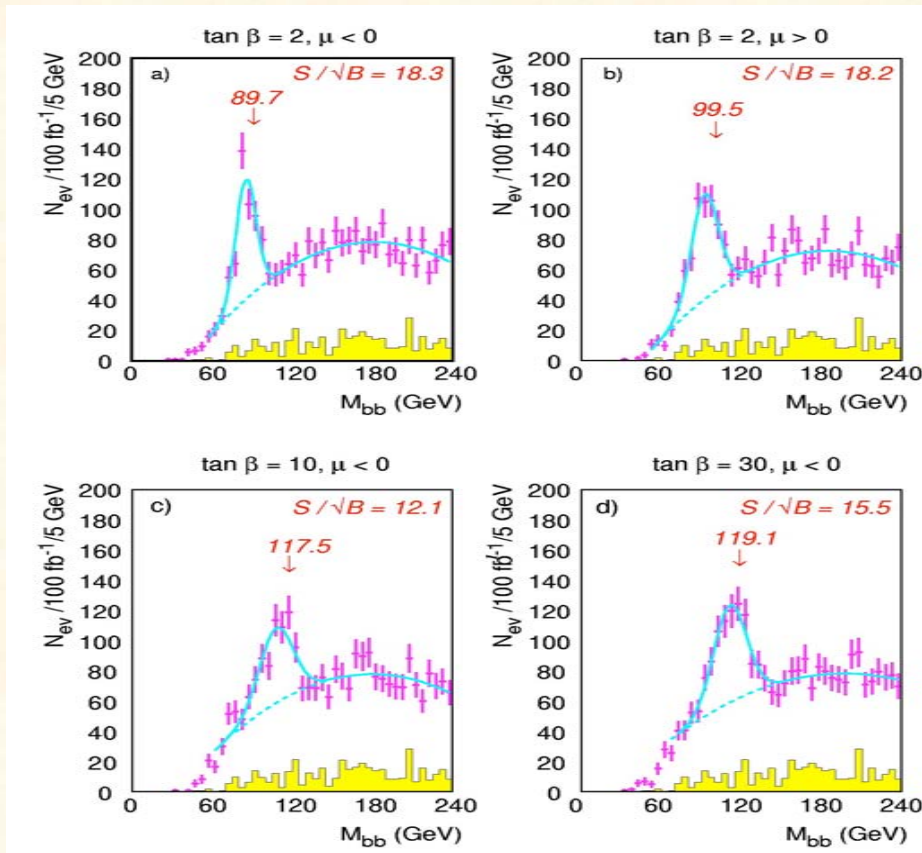
⇒ combination of  
**Jets, Leptons,  $E_T^{\text{miss}}$**

# Search for $h \rightarrow bb$ in SUSY cascade decays

Applying a cut on  $E_T^{\text{miss}} \Rightarrow$  suppresses the Standard Model background (QCD-jets), dominant background from SUSY production

**h  $\rightarrow$  bb:**

**CMS study, mSUGRA**



important if  $\chi_2^0 \rightarrow \chi_1^0 h$  is open;

bb peak can be reconstructed in many cases

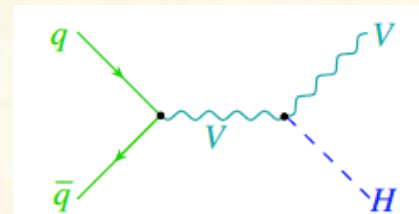
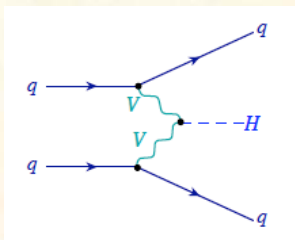
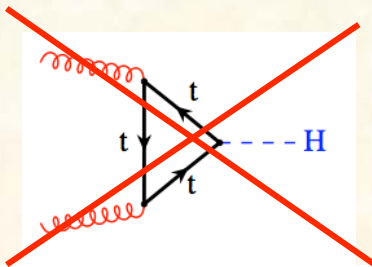
**Could be a Higgs discovery mode !**

# Can invisibly decaying

## Higgs bosons be detected ?

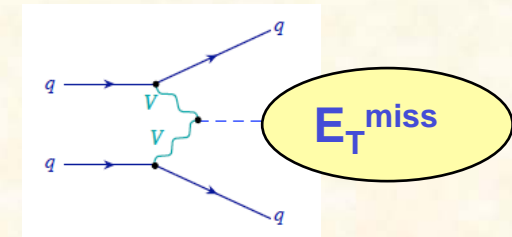
### Motivation:

- Invisibly decaying Higgs bosons appear in several extensions of the Standard Model
- e.g. Higgs bosons can decay into weakly interacting neutralinos, gravitinos, scalar particles, .....
- To detect invisibly decaying Higgs bosons, the associated production modes must be used; most promising are VBF and ZH associated production



## Invisible Higgs decays ?

Possible searches:  $tt \quad H \rightarrow \ell\nu b \bar{q}q + E_T^{\text{miss}}$   
 $W/Z \quad H \rightarrow \ell\nu (\ell\ell) + E_T^{\text{miss}}$   
 $qq \quad H \rightarrow qq + E_T^{\text{miss}}$



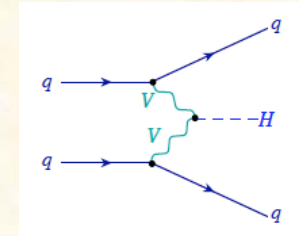
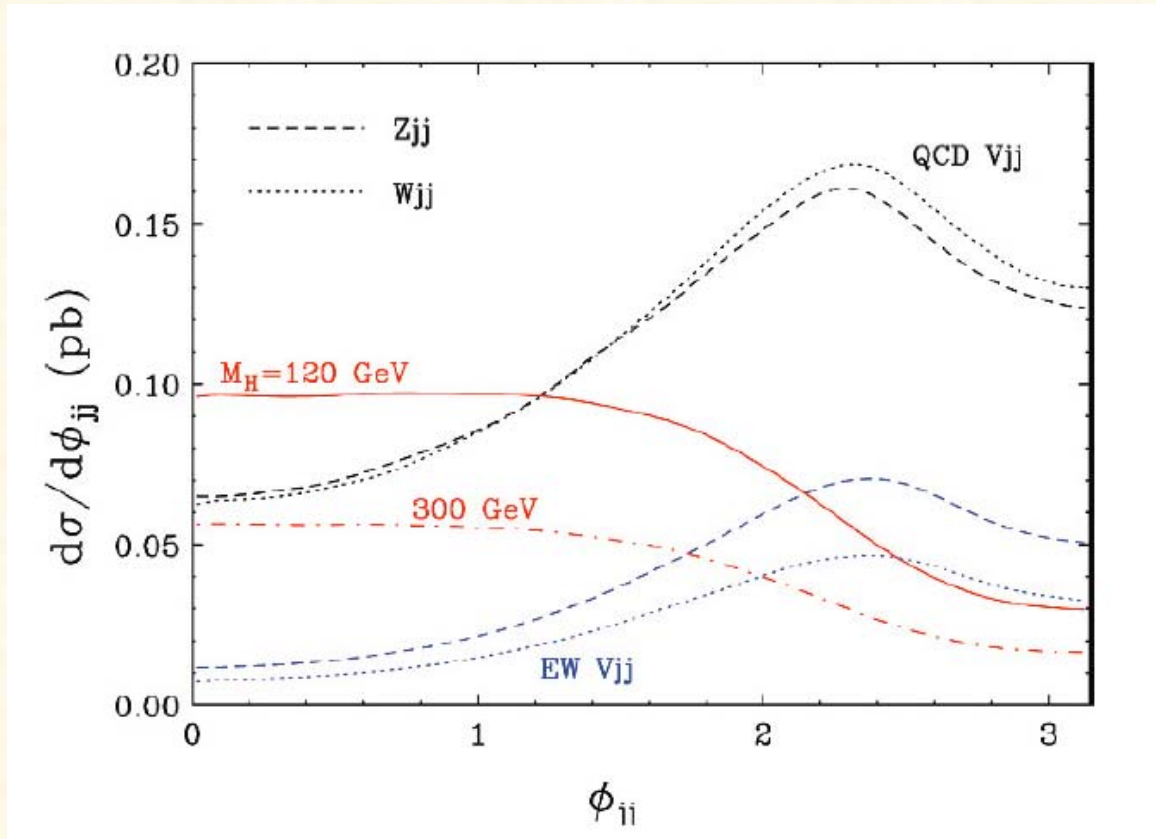
### ATLAS study:

Search for invisibly decaying Higgs boson in the VBF mode and in the associated ZH production

Event selection: 2 tag jets,  $(P_T, \Delta\eta, M_{jj} > 1200 \text{ GeV})$   
 $P_T^{\text{miss}} > 100 \text{ GeV}$   
Lepton and Jet veto (no jets with  $P_T > 20 \text{ GeV}$ )

Main backgrounds: W jj production ( $W \rightarrow \ell\nu$ )  
Z jj production ( $Z \rightarrow \nu\nu$ )  
QCD jet production, fake  $E_T^{\text{miss}}$

Discriminating variable:  $\Delta \phi_{jj}$  (separation between tag jets)  
 expect differences due to Higgs coupling structure:

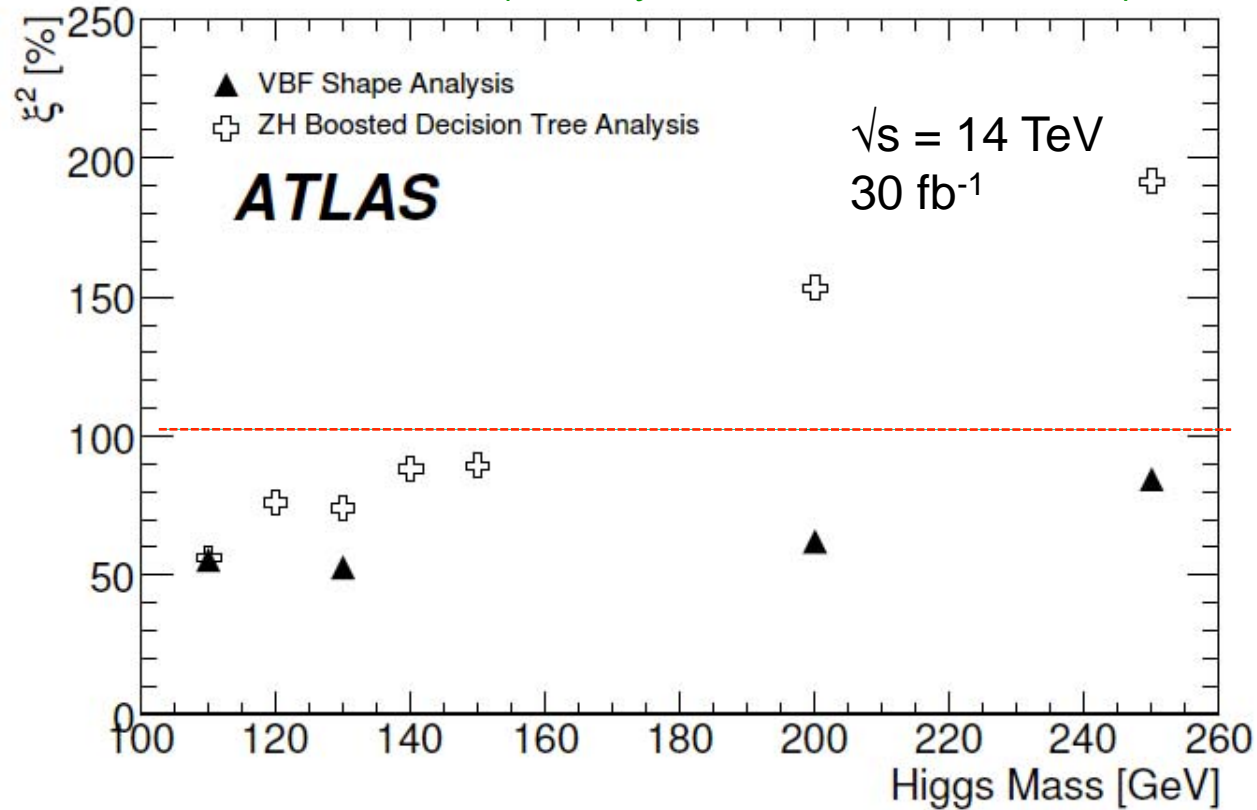


D. Zeppenfeld et al.

## Sensitivity:

$$\xi^2 = Br(H \rightarrow Inv.) \frac{\sigma_{qq \rightarrow qqH}}{\sigma_{qq \rightarrow qqH}|_{SM}}$$

95% C.L. limits (incl. systematic uncertainties)

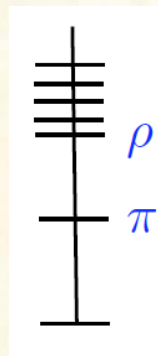


- Sensitivity to large invisible branching ratios (SM production) in VBF channel
- Proof that the nature of the invisibly decaying object is a Higgs boson is difficult

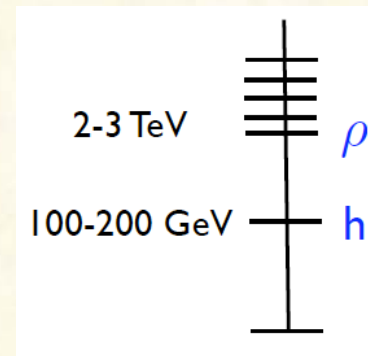
Composite Higgs Bosons ?

# Composite Higgs models

- Composite Higgs models are inspired by QCD where one observes that the (pseudo)scalar mesons are the lightest states
  - pseudo-Nambu-Goldstone bosons, mass protected by global symmetries in QCD



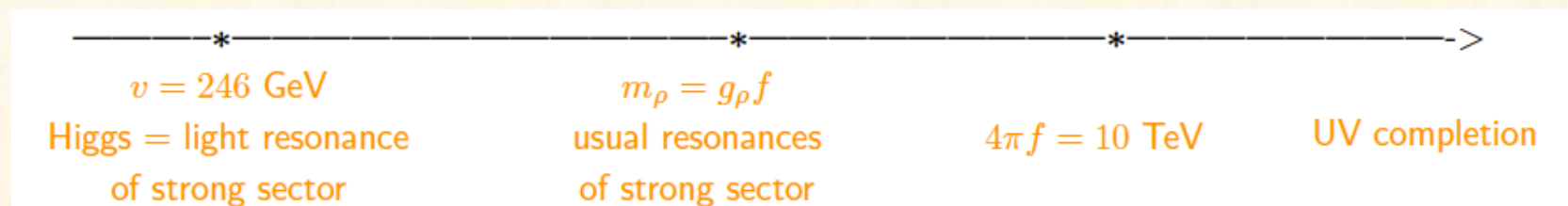
QCD



QCD like sector

- in analogy, a light Higgs can be obtained as a Goldstone boson of a globally broken symmetry at a scale  $f$

- Energy / mass scales:





# Composite Higgs models (cont.)\*

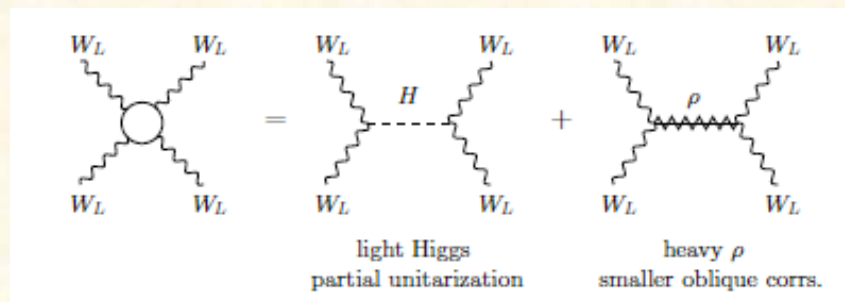
- Couplings of composite Higgs are altered, deviate from Standard Model couplings

$$\text{HWW: } a \frac{2M_W^2}{v}$$

$$\text{HHWW: } b \frac{2M_W^2}{v^2}$$

**Standard Model:  $a = b = 1$**

- Composite models:
- deviations from 1,
  - non perfect cancellation of unitarity problems by light Higgs
  - heavier  $\rho$ -like states needed



\*)

- Giudice, Grojean, Pomarol, Ratazzi
- Contino et al.
- Espinosa, Grojean, Mühlleitner
- ....

# Composite Higgs models (cont.)

- Continuous interpolation between the Standard Model and Technicolour:

$$\xi = \left(\frac{v}{f}\right)^2 = \left(\frac{\text{weak scale}}{\text{strong coupling scale}}\right)^2$$

$\xi = 0$  : Standard Model

$\xi = 1$ : Technicolour

Higgs decouples, vector resonances like in TC

- Modified couplings in composite models  
e.g. Minimal Composite Higgs Models (MCHM):

MCHM4:  $g_{HVV} = g_{HVV}^{SM} \sqrt{1 - \xi}$

$$g_{Hff} = g_{Hff}^{SM} \sqrt{1 - \xi}$$

universal shift of couplings,  
no change in branching ratios

MCHM5:  $g_{HVV} = g_{HVV}^{SM} \sqrt{1 - \xi}$

$$g_{Hff} = g_{Hff}^{SM} \frac{1 - 2\xi}{\sqrt{1 - \xi}}$$

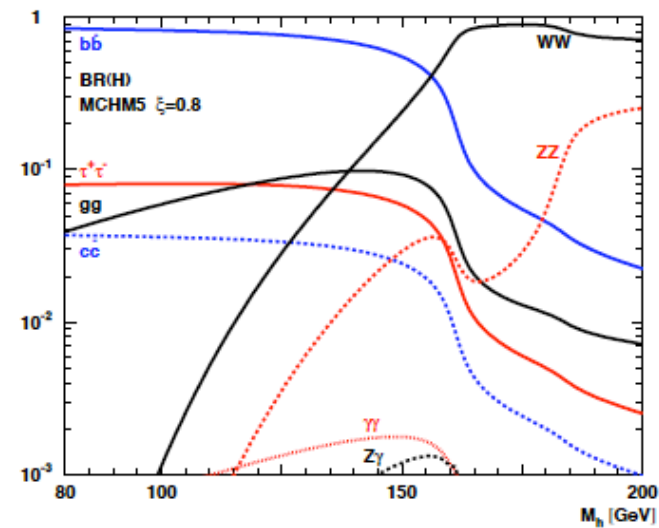
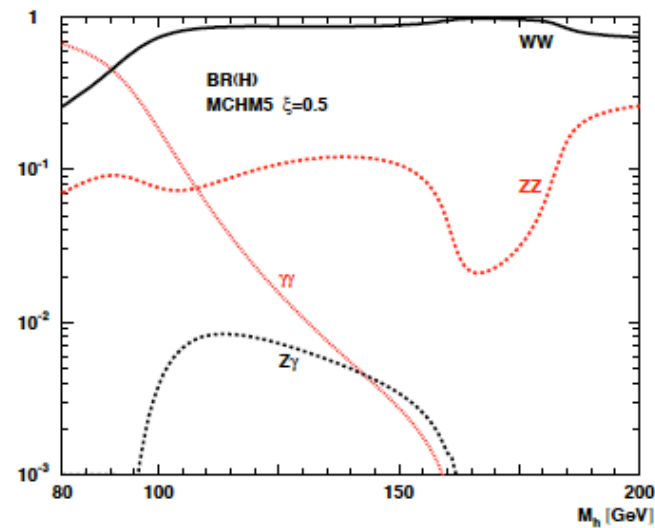
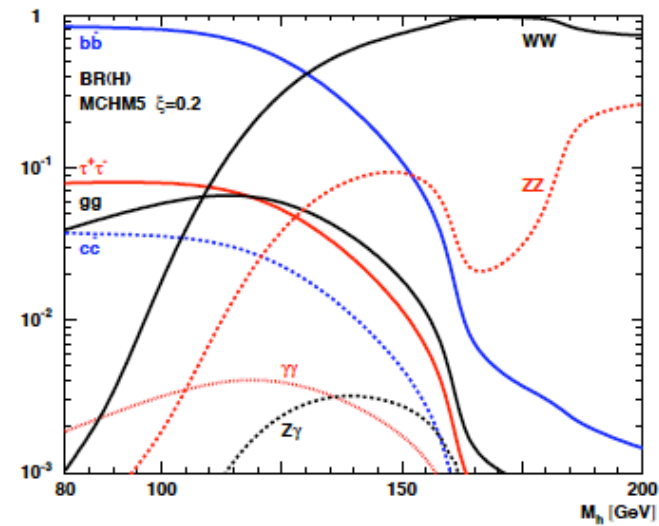
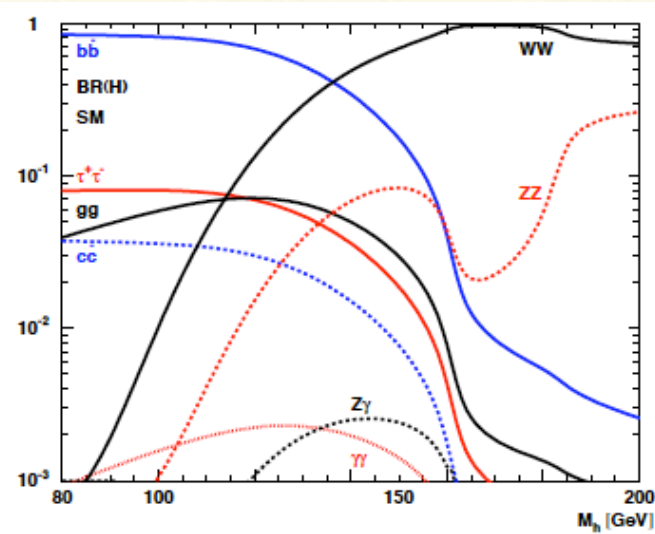
change in BR, depend on  $\xi$ ,  
vanishes for  $\xi=0.5$

...and: lower production rates !!

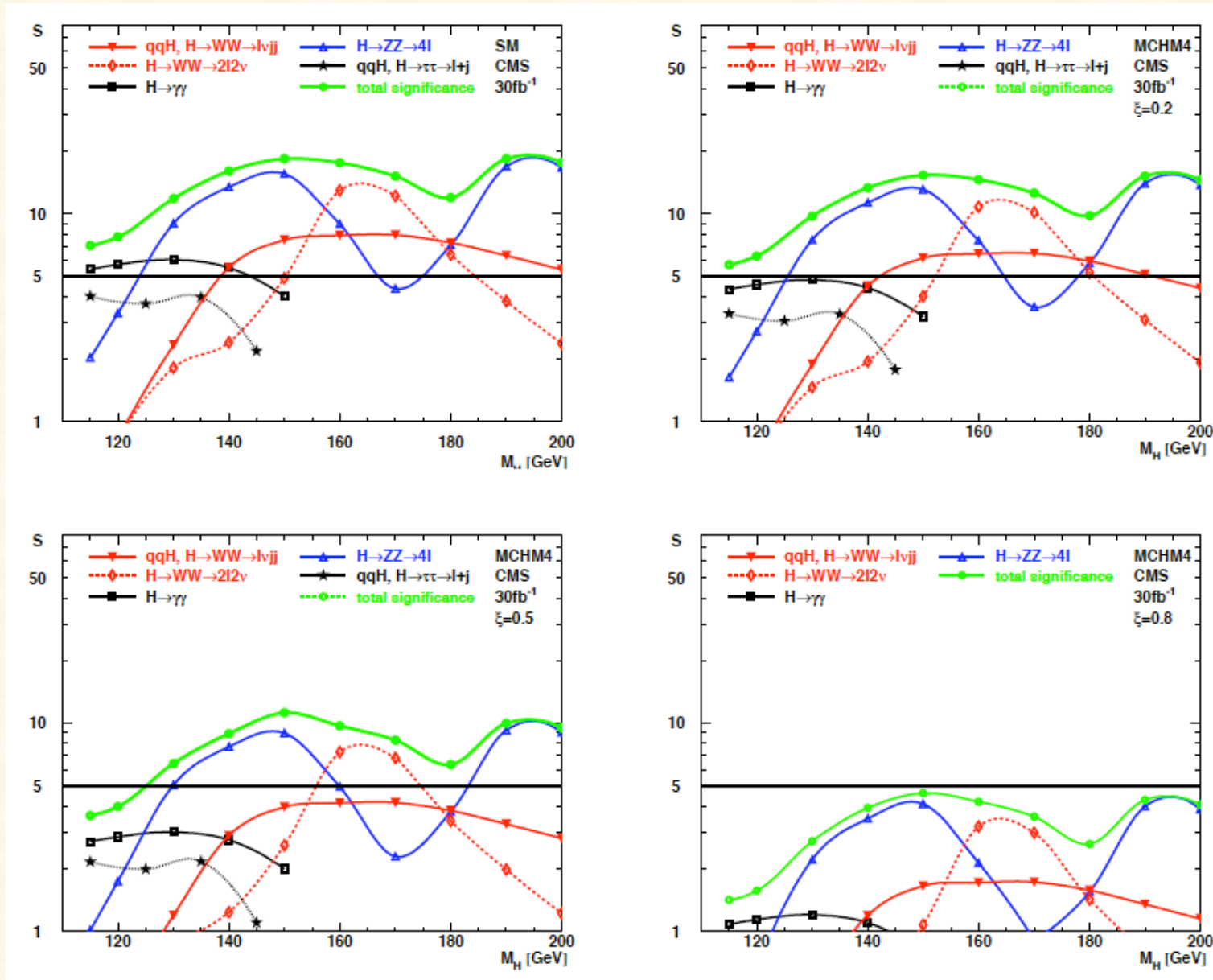
# Branching ratios in MCHM5 model

M. Mühlleitner, Physics at the LHC, Hamburg June 2010

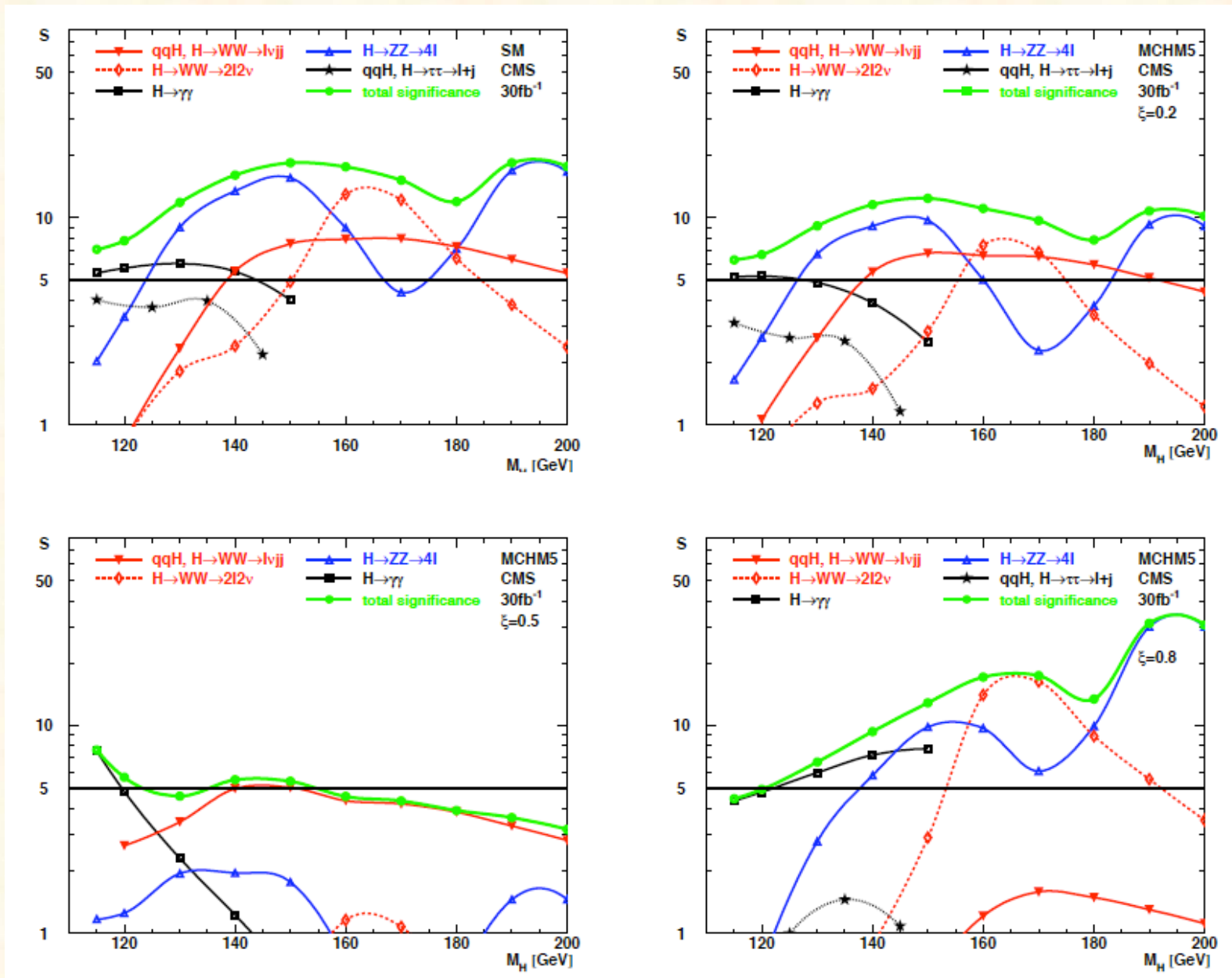
Espinosa, Grojean, Mühlleitner



# LHC sensitivity to such models: MCHM4



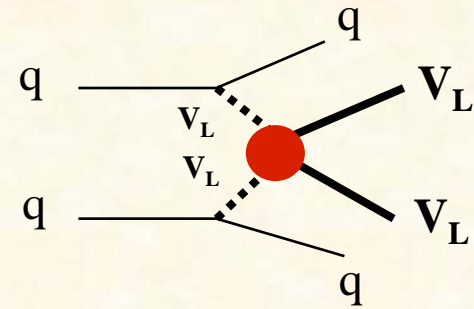
# LHC sensitivity to such models: MCHM5



# What if the LHC does

**not** find the

## Higgs Boson ?



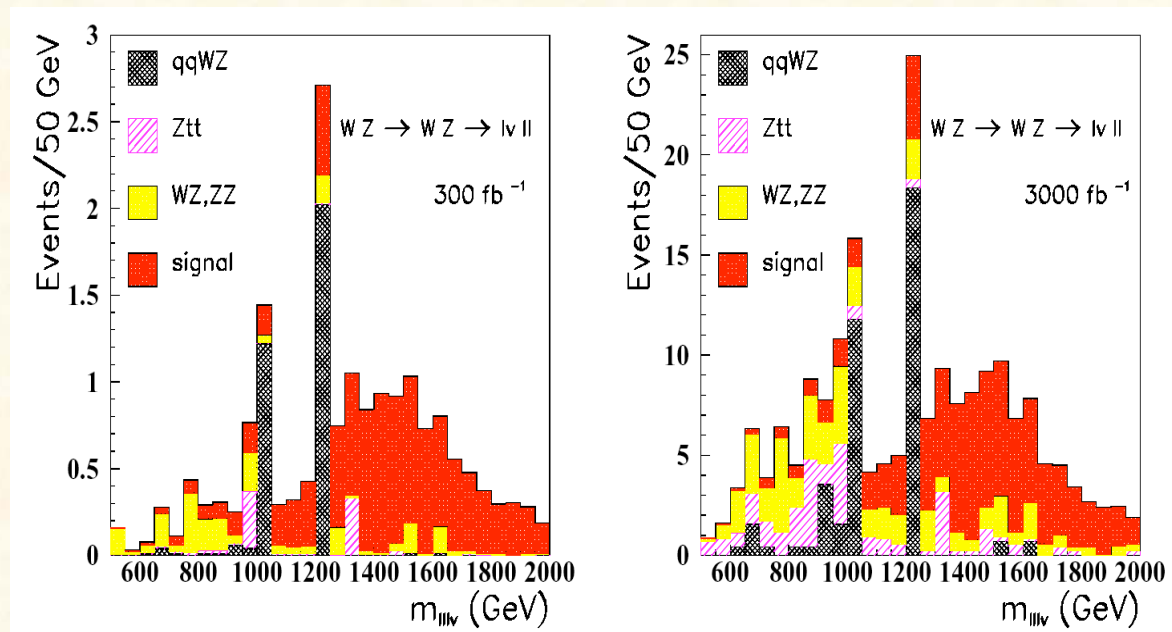
- Study of longitudinal gauge boson scattering is the key  
High luminosities, i.e. sLHC, required to make quantitative measurements  
(strong physics case)
- **If no Higgs**, expect strong  $V_L V_L$  scattering (resonant or non-resonant) at  $\sim 1\text{TeV}$
- Also the question of a composite Higgs boson will profit from these measurements

# WZ resonances in Vector Boson Scattering

Example: Vector resonance ( $\rho$ -like) in  $W_L Z_L$  scattering  
from Chiral Lagrangian model

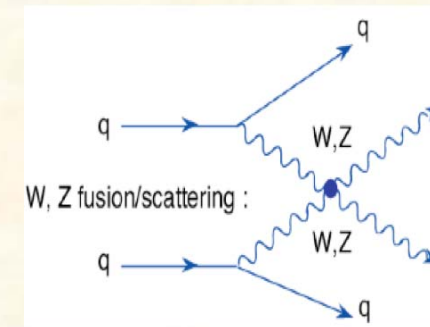
$m = 1.5 \text{ TeV} \Rightarrow 300 \text{ fb}^{-1} \text{ (LHC) vs. } 3000 \text{ fb}^{-1} \text{ (sLHC)}$

Lepton cuts:  $p_T^1 > 150 \text{ GeV}$ ,  $p_T^2 > 100 \text{ GeV}$ ,  $p_T^3 > 50 \text{ GeV}$ ;  $E_T^{\text{miss}} > 75 \text{ GeV}$



At LHC:  $S = 6.6 \text{ events}$ ,  
 $B = 2.2 \text{ events}$

At sLHC:  $S/\sqrt{B} \sim 10$



These studies require  
both forward jet  
tagging and central jet  
vetoing!  
Expect degradation of  
sLHC performance

# Conclusions

- Should a SM Higgs boson exist, it cannot escape detection at the LHC
- MSSM parameter space can be covered for several benchmark scenarios (incl. LHC-phobic scenarios)
- Tevatron might have a  $3\text{-}\sigma$  evidence window;  
Information on  $b\bar{b}$  decays important complementary information to LHC
- On the longer term, LHC can perform first, important measurements of Higgs boson parameters, which are needed to probe the underlying theory

“Exiting times ahead of us”





## Acknowledgements:

I am grateful to:

K. Assamagan, R. Bernhard, S. Dittmaier, A. Djouadi, C. Grojean,  
R. Harlander, C. Mariotti, M. Mühlleitner, H. Murayama, A. Pomarol,  
M. Schumacher, M. Spira, R. Thorne, G. Watts, G. Weiglein, C. Weiser,  
D. Zeppenfeld

for useful discussions on the various Higgs topic....and for their “indirect”  
contributions to these lectures via their excellent talks at conferences and  
workshops, from which I profited a lot.