



**Faculty  
of Physics**

WARSAW UNIVERSITY OF TECHNOLOGY



**ALICE**



# Probing space-time evolution at the femtometer scale in pp and Pb-Pb collisions with ALICE

**Łukasz Graczykowski  
for the ALICE Collaboration**

CERN-LHC Seminar  
CERN

5 March 2019



# Femtosecond technique

PRL **96**, 166101 (2006)

PHYSICAL REVIEW LETTERS

week ending  
28 APRIL 2006

## Laser-Induced Microexplosion Confined in the Bulk of a Sapphire Crystal: Evidence of Multimegabar Pressures

S. Juodkazis,<sup>1</sup> K. Nishimura,<sup>1</sup> S. Tanaka,<sup>1</sup> H. Misawa,<sup>1</sup> E. G. Gamaly,<sup>2</sup> B. Luther-Davies,<sup>2</sup>  
L. Hallo,<sup>3</sup> P. Nicolai,<sup>3</sup> and V. T. Tikhonchuk<sup>3</sup>

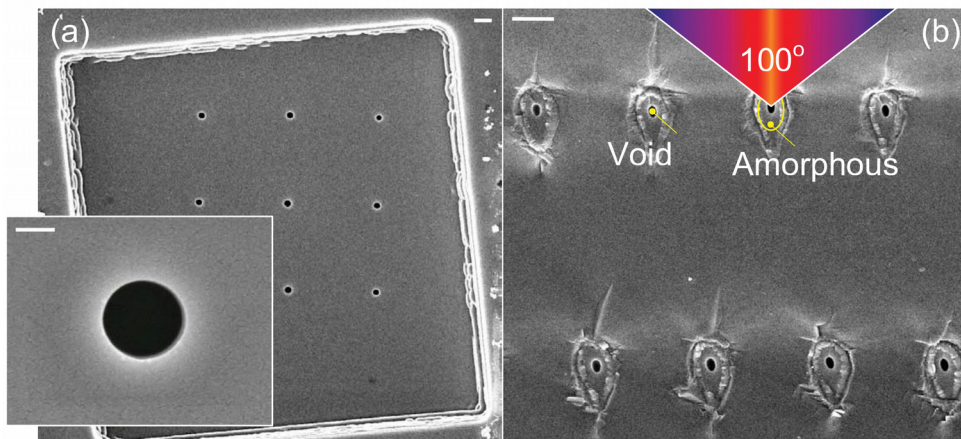
<sup>1</sup>CREST-JST and Research Institute for Electronic Science, Hokkaido University, N21-W10,  
CRIS Building, Kita-ku, Sapporo 001-0021, Japan

<sup>2</sup>Centre for Ultrahigh Bandwidth Devices for Optical Systems, Laser Physics Centre,  
Research School of Physical Sciences and Engineering, The Australian National University, Canberra ACT 0200, Australia

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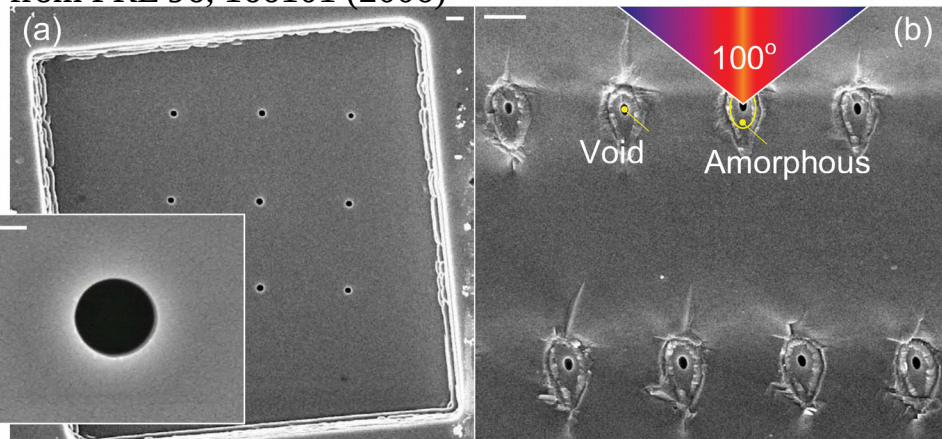
(Received 24 November 2005; published 25 April 2006)

Extremely high pressures ( $\sim 10$  TPa) and temperatures ( $5 \times 10^5$  K) have been produced using a single laser pulse (100 nJ, 800 nm, 200 fs) focused inside a sapphire crystal. The laser pulse creates an intensity over  $10^{14}$  W/cm<sup>2</sup> converting material within the absorbing volume of  $\sim 0.2$   $\mu\text{m}^3$  into plasma in a few fs. A pressure of  $\sim 10$  TPa, far exceeding the strength of any material, is created generating strong shock and rarefaction waves. This results in the formation of a nanovoid surrounded by a shell of shock-affected material inside undamaged crystal. Analysis of the size of the void and the shock-affected zone versus the deposited energy shows that the experimental results can be understood on the basis of conservation laws and be modeled by plasma hydrodynamics. Matter subjected to record heating and cooling rates of  $10^{18}$  K/s can, thus, be studied in a well-controlled laboratory environment.

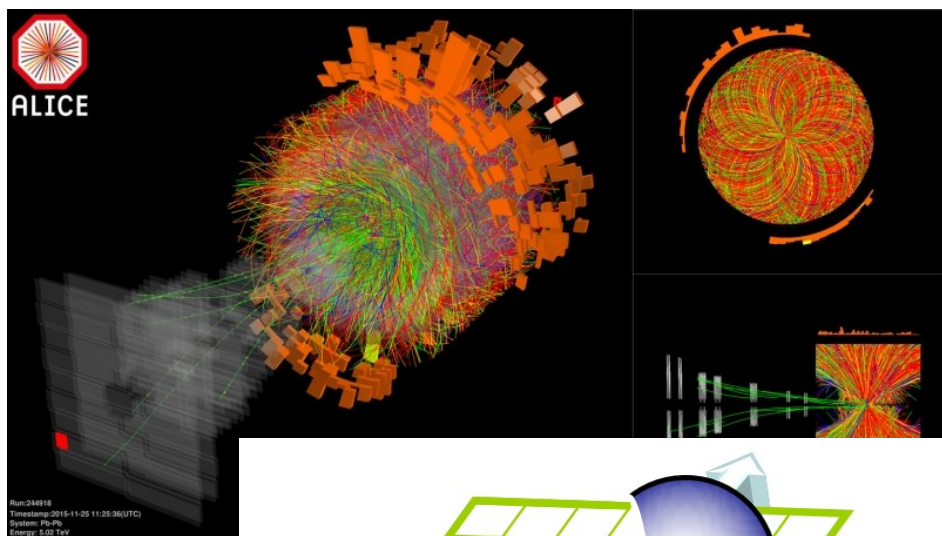


# Femtosecond technique

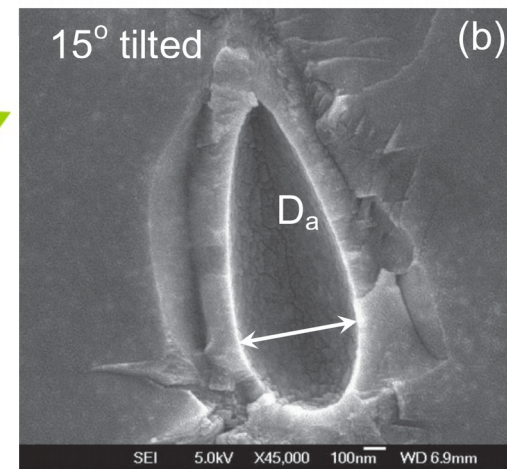
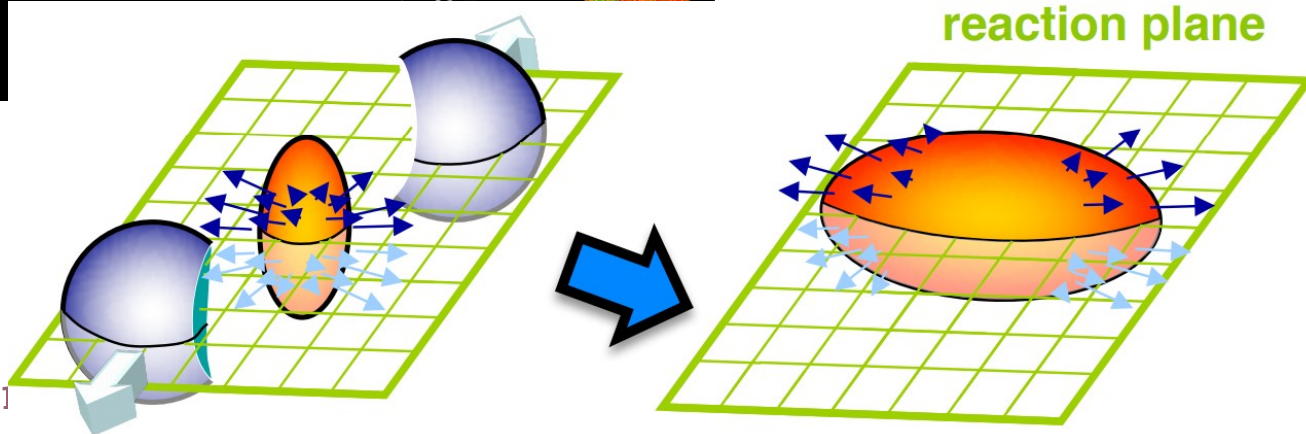
from PRL 96, 166101 (2006)



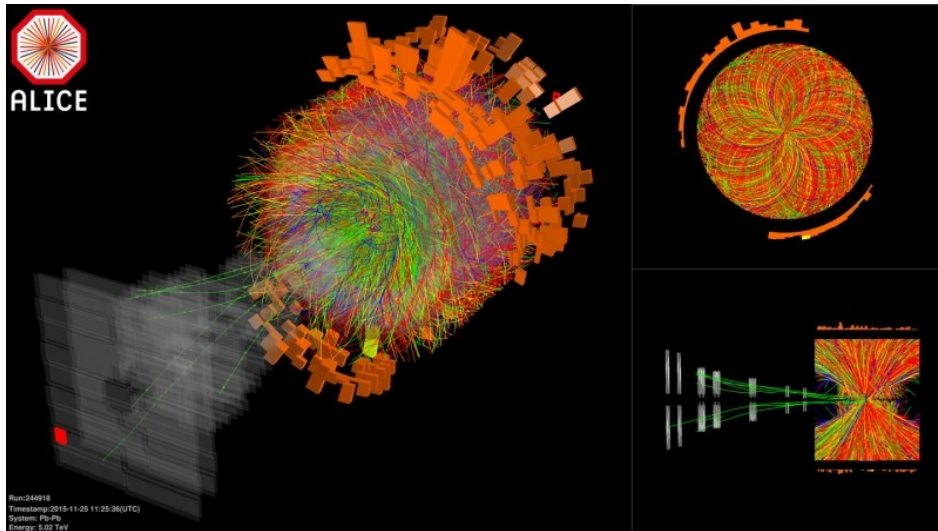
- Does it look similar?
- Let's see:
  - energy quickly deposited
  - enter plasma phase
  - expand hydrodynamically



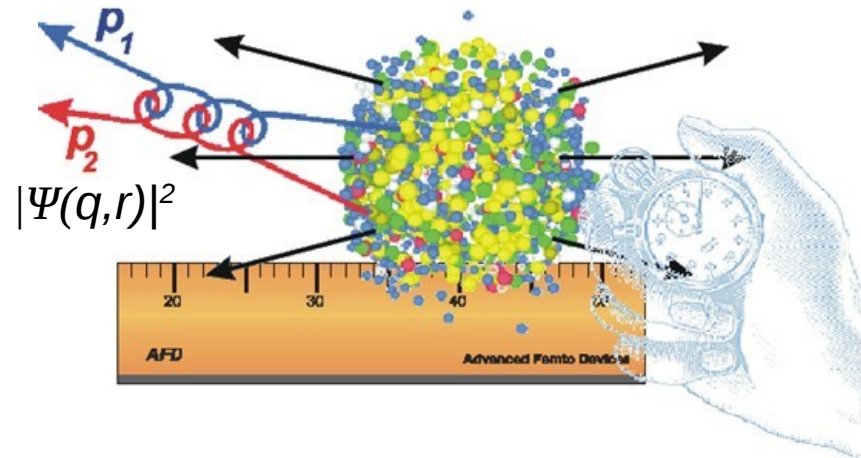
- We can do a “post mortem” analysis to investigate i.e. the source geometry



# Femtoscscopy technique



from M. Lisa and S. Pratt



- Femtoscscopy – measures space-time characteristics of the source using particle correlations in momentum space

$$C(\vec{p}_1, \vec{p}_2) = \frac{P_{12}(\vec{p}_1, \vec{p}_2)}{P_1(\vec{p}_1)P_2(\vec{p}_2)}$$

$$\vec{q} = \vec{p}_1 - \vec{p}_2$$

$$\vec{r} = \vec{x}_1 - \vec{x}_2$$

experiment

$$C(\vec{q}) = \frac{A(\vec{q})}{B(\vec{q})}$$

$A(\vec{q})$  - correlated pairs (“same events”)

$B(\vec{q})$  - uncorrelated pairs (“mixed events”)

theory (models)

$$C(\vec{q}) = \frac{\int d^3 r S_{12}(\vec{q}, \vec{r}) |\Psi(\vec{q}, \vec{r})|}{\int d^3 x_1 S_1(\vec{x}_1, \vec{p}_1) \int d^3 x_2 S_2(\vec{x}_2, \vec{p}_2)}$$

$$C(\vec{q}) = \int d^3 r S(\vec{q}, \vec{r}) |\Psi(\vec{q}, \vec{r})|$$



# Why are particles correlated?

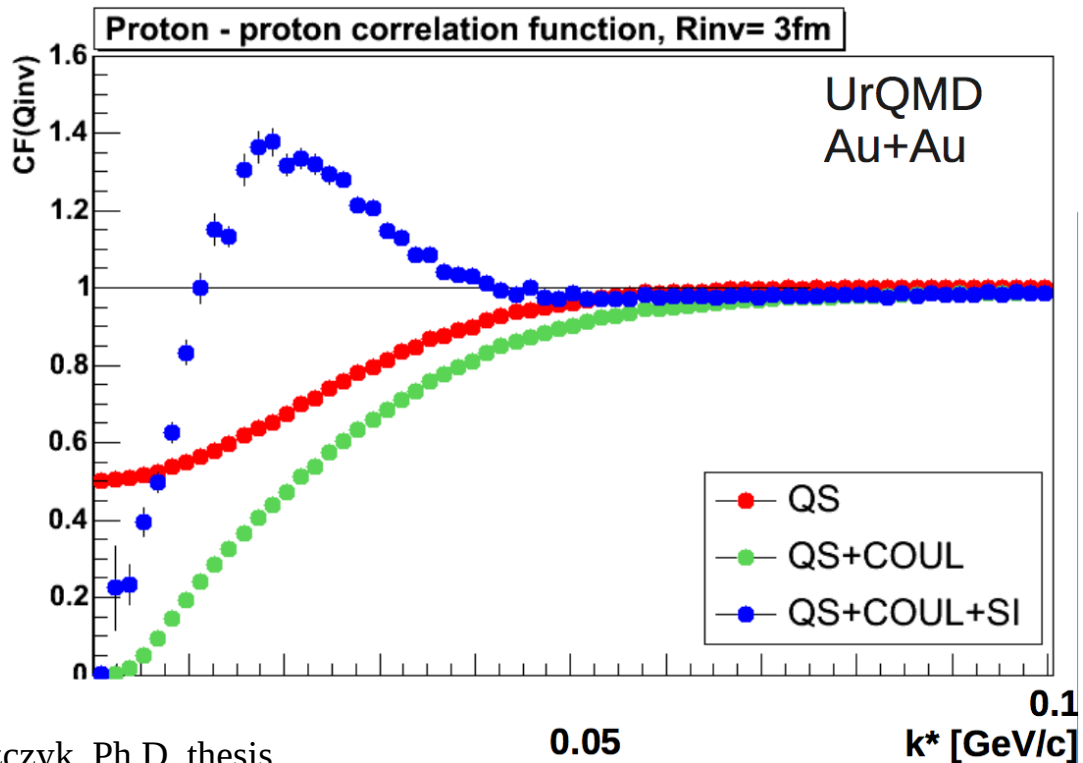
- **Main sources of correlations:**

- **Quantum statistics (QS)**

- pairs of identical bosons (i.e. pions) – Bose-Einstein QS
- pairs of identical fermions (i.e. protons) – Fermi-Dirac QS

- **Final-state interactions (FSI)**

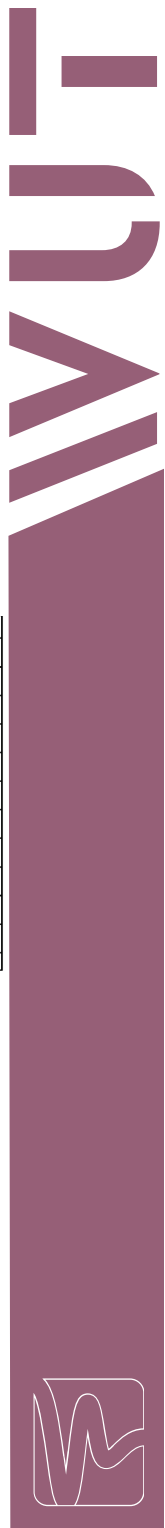
- strong interaction
- Coulomb interaction



H. Zbroszczyk, Ph.D. thesis

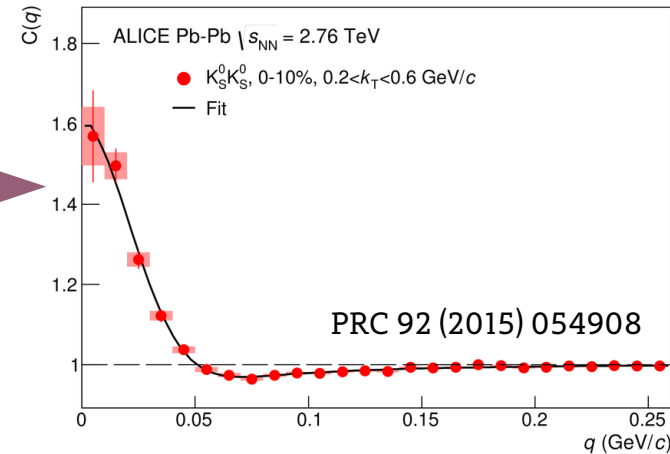


# How does it look in data?



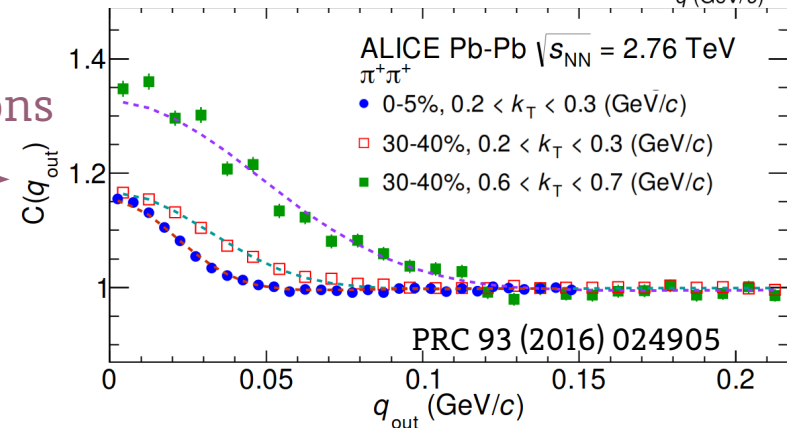
Quantum Statistics

identical neutral kaons



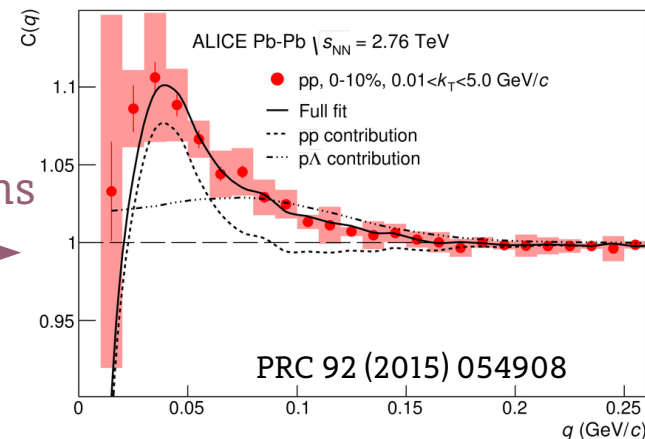
Quantum Statistics + Coulomb FSI

identical charged pions

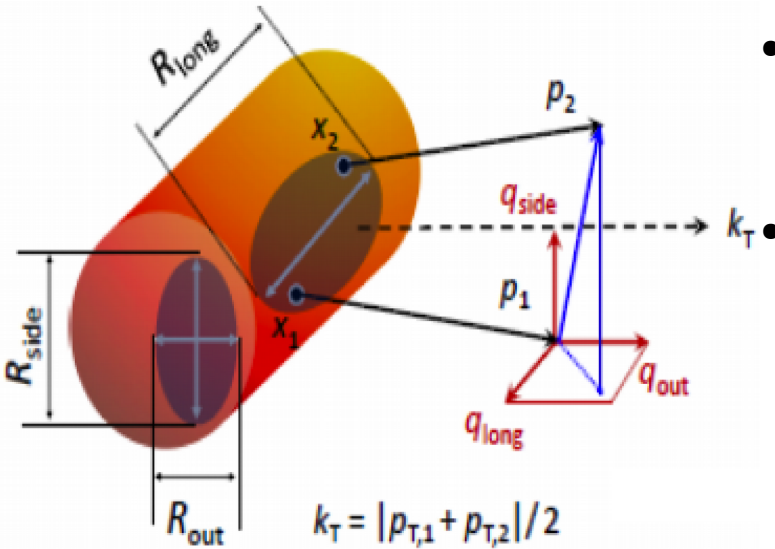
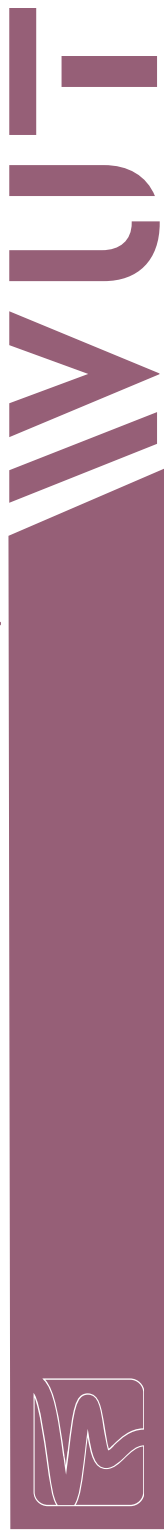


Quantum Statistics + Coulomb FSI + strong FSI

identical protons



# Know your reference frame



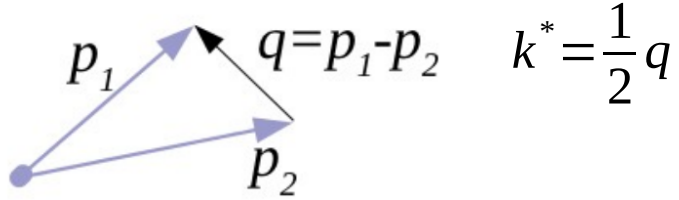
- 3-dimensional **Longitudinally Co-Moving System (LCMS)**
- Decomposition of  $q$ :
  - Long along the beam: sensitive to longitudinal dynamics and evolution time
  - Out along  $k_T$ : sensitive to geometrical size, emission time and space-time correlation
  - Side perpendicular to Long and Out: sensitive to geometrical size

$$m_T = \sqrt{k_T^2 + m_\pi^2}$$

LCMS frame:

$$p_{1,long} = -p_{2,long}$$

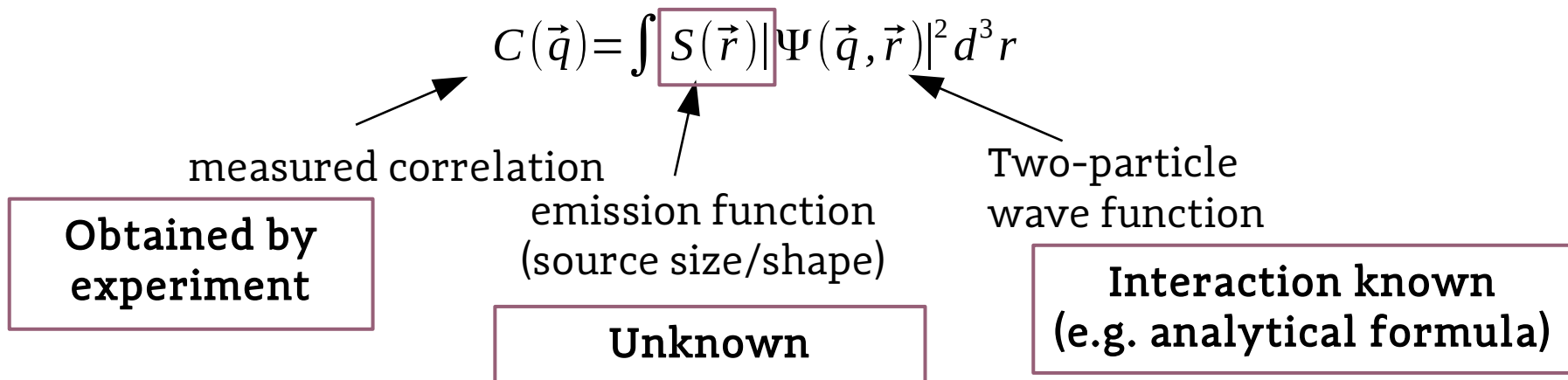
PRF frame:



- For statistically challenged analyses, measurement in one dimension (giving only one size) in **Pair Rest Frame (PRF)**

# “Traditional” femtoscopy

## Measuring the system size



- **Two main topics in this part:**
  - collectivity and spatial geometry of the system
  - temporal particle emission from QGP



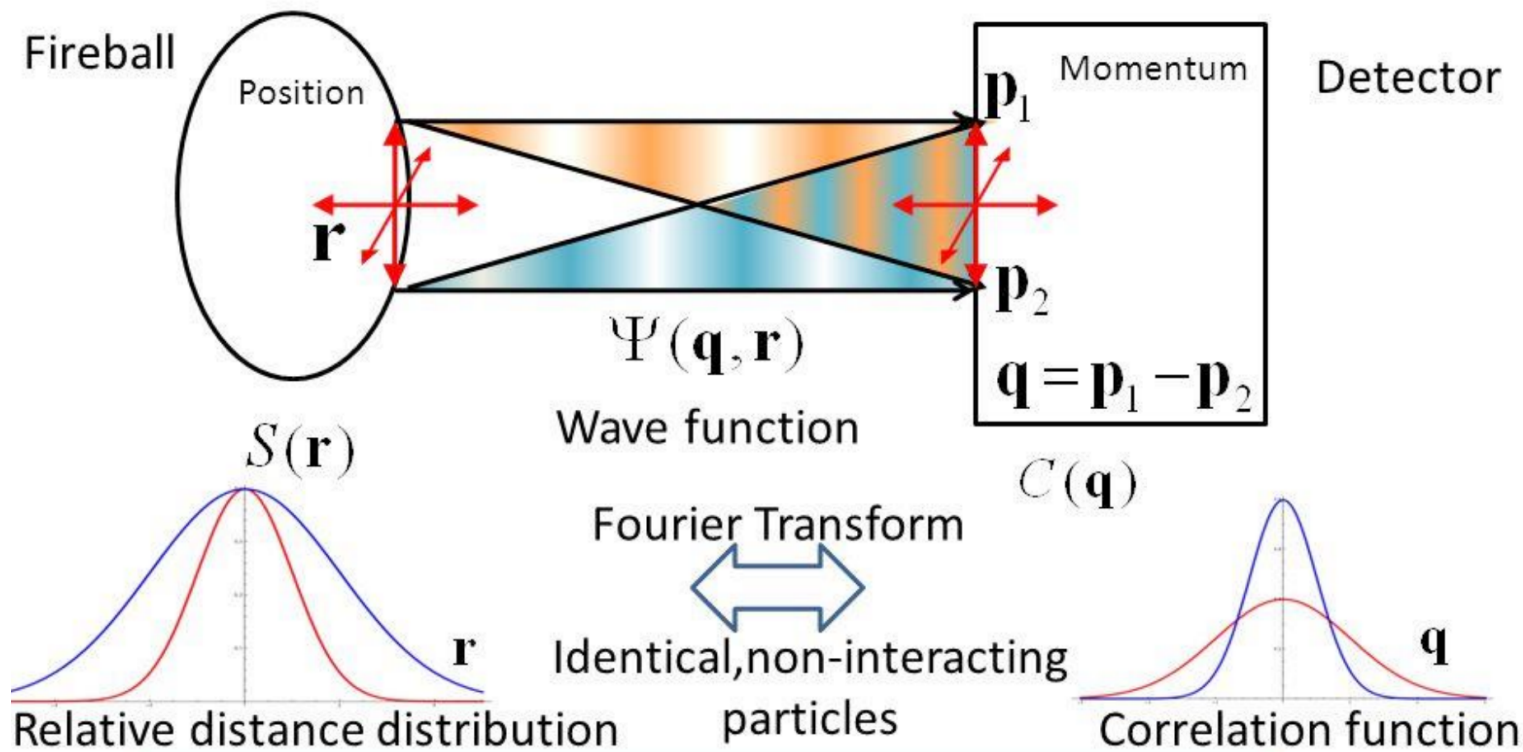


# HBT radius

In case of identical bosons, CF is a Fourier transformation of WF

$$S(\vec{r}) \sim \exp\left(-\frac{r_{out}^2}{4R_o^2} - \frac{r_{side}^2}{4R_s^2} - \frac{r_{long}^2}{4R_l^2}\right) \quad \left. \begin{array}{l} \text{Gaussian source} \\ \\ \text{Bose-Einstein QS} \end{array} \right\} C = 1 + \lambda \exp(-R_o^2 q_o^2 - R_s^2 q_s^2 - R_l^2 q_l^2)$$

$$|\Psi(\vec{q}, \vec{r})|^2 = 1 + \cos(\vec{q} \cdot \vec{r})$$



from Michiel de Kock

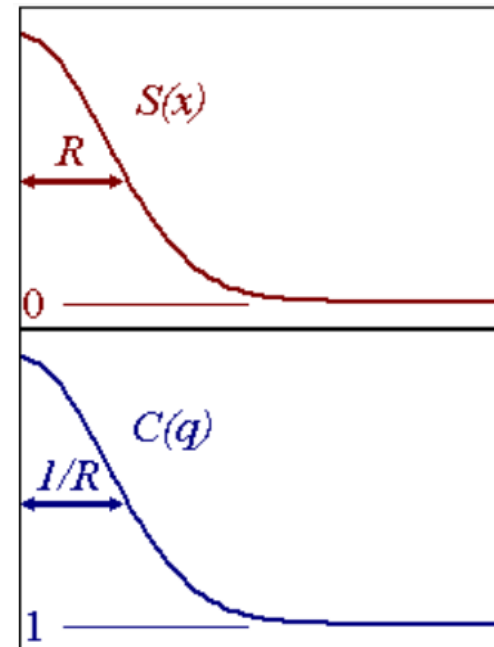
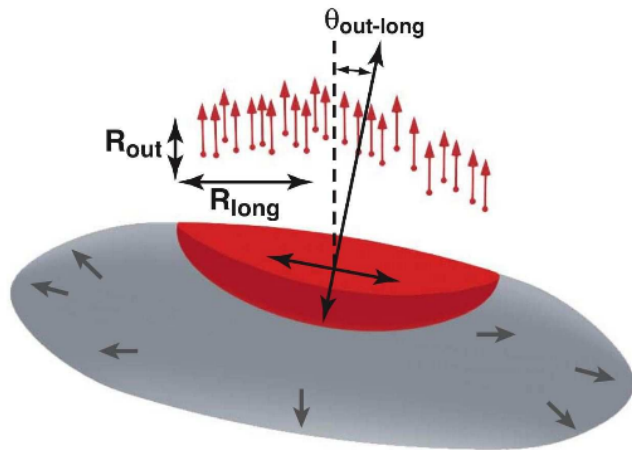


# HBT radius

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$$|\Psi(\vec{q}, \vec{r})|^2 = 1 + \cos(\vec{q} \cdot \vec{r})$$

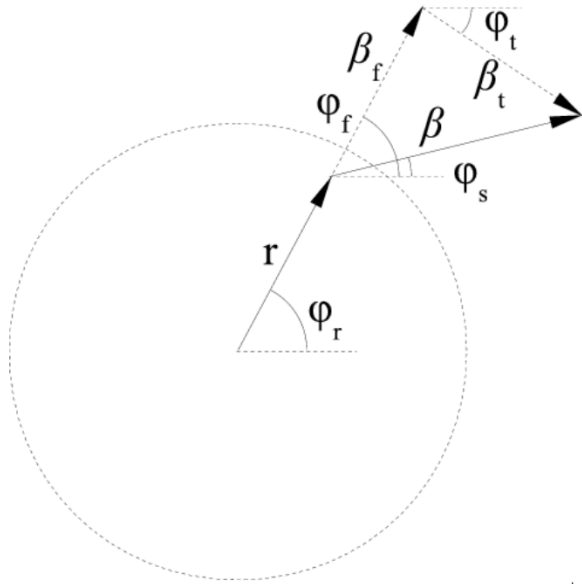


from Ann.Rev.Nucl.Part.Sci.55:357-402,2005

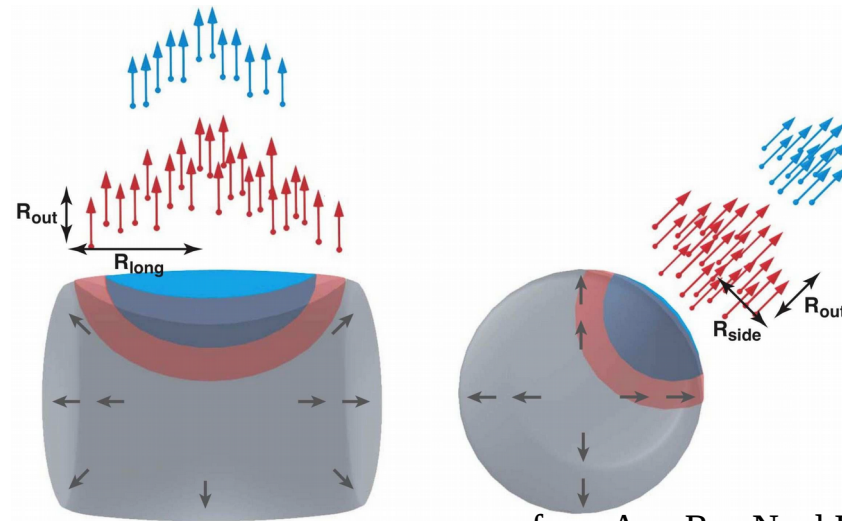
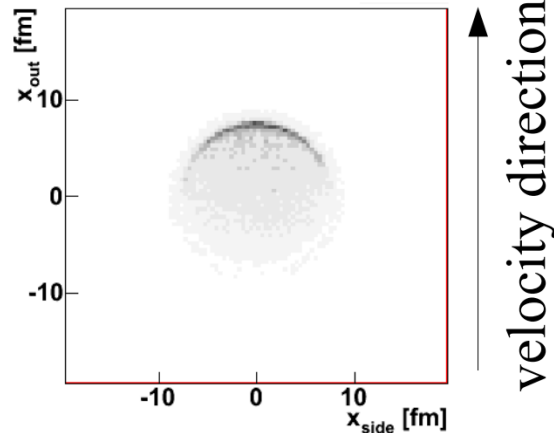
- The size (or sizes in 3D)  $R$  is referred to as the “HBT radius”  
 → size of the “region of homogeneity”  
 region from which particles are emitted with similar velocity



# Probing collectivity



- Particle emitted from the medium will have a collective (flow) velocity  $\beta_f$  and thermal (random) velocity  $\beta_t$
- While  $m_T$  (transverse momentum or particle mass) grows,  $\beta_t \sim 1/\sqrt{m_T}$  decreases, the region from where particles are emitted gets smaller and shifted towards the outside of the source



**blue** – heavier  
(or higher  $p_T$ )

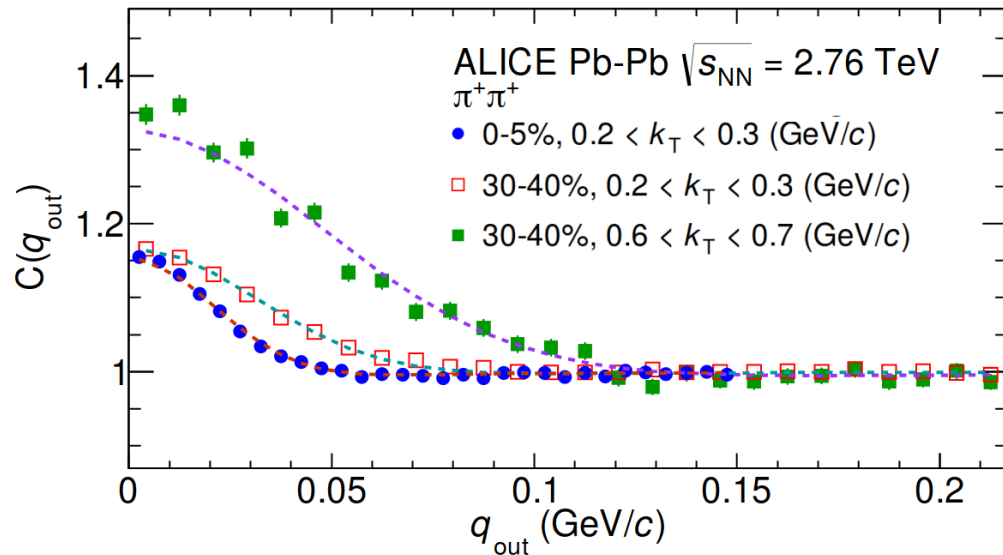
**red** – lighter  
(or lower  $p_T$ )

from Ann.Rev.Nucl.Part.Sci.55:357-402,2005

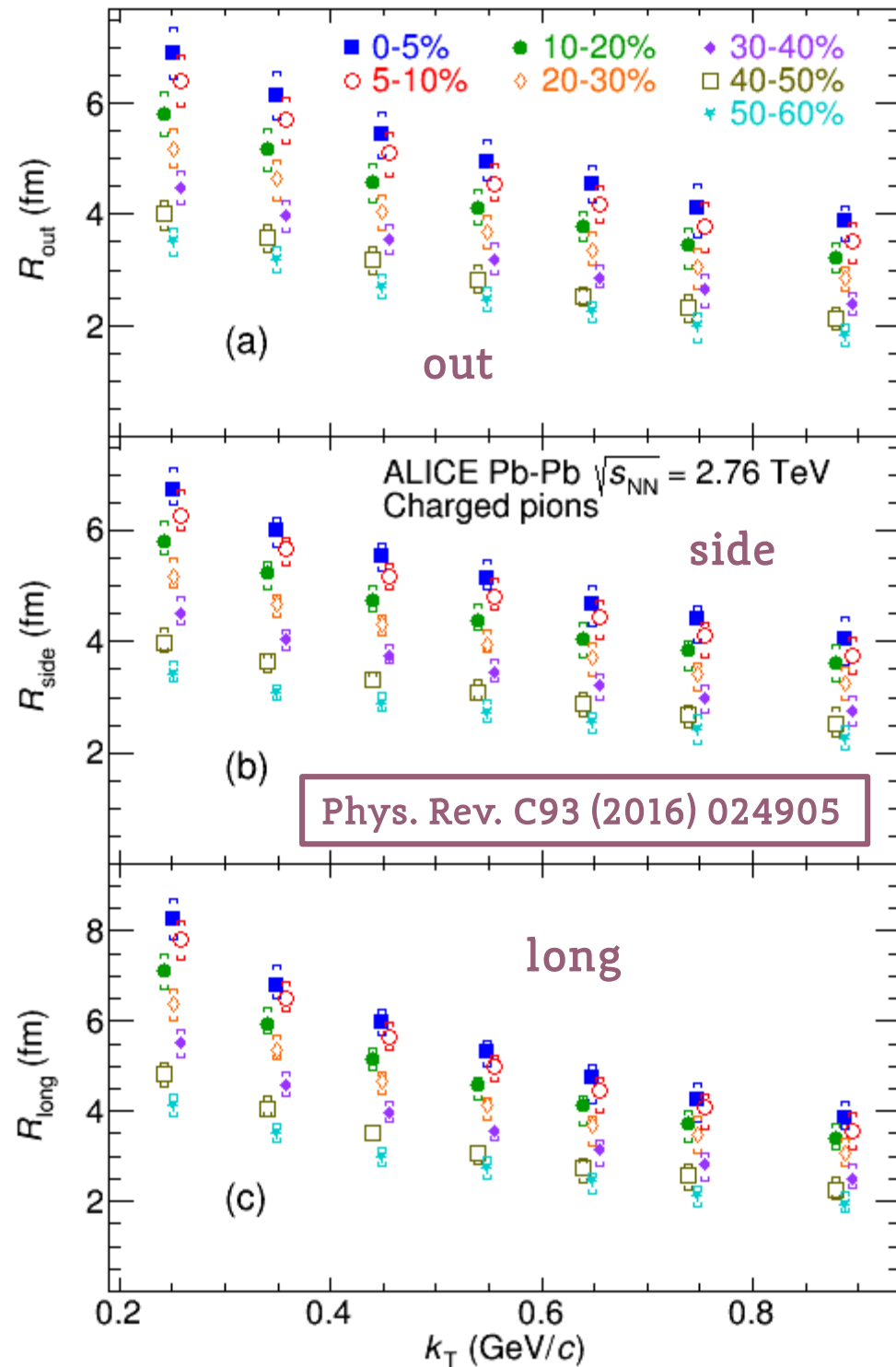
- Conclusion: HBT radii should get smaller with increasing  $k_T$  or  $m_T$



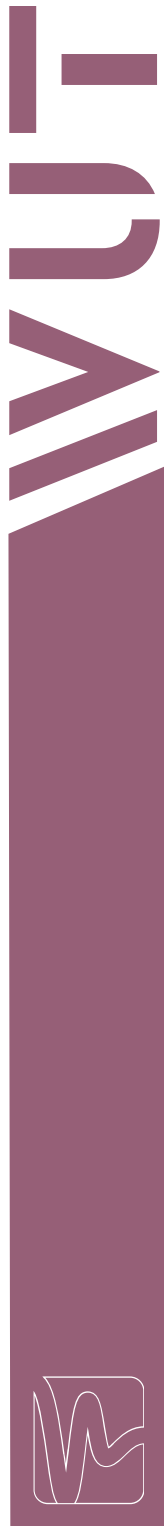
# Pion femtoscopy



- Very good description of the measured correlation function with the **Gaussian source**
- Radii universally grow with event multiplicity and fall with pair transverse momentum  $k_T$



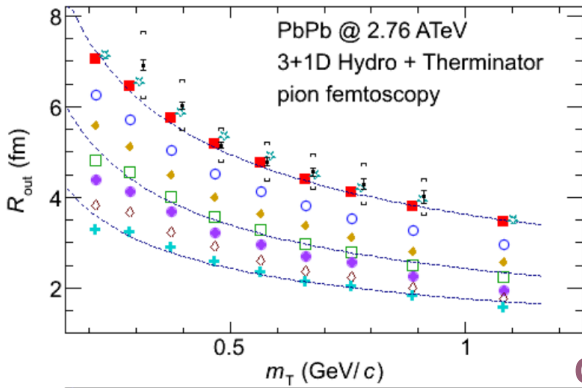
# Pion femtoscopy



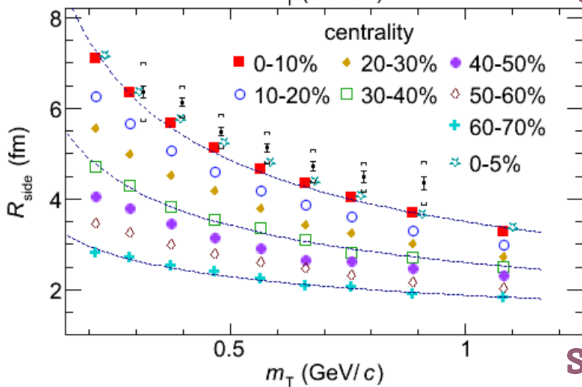
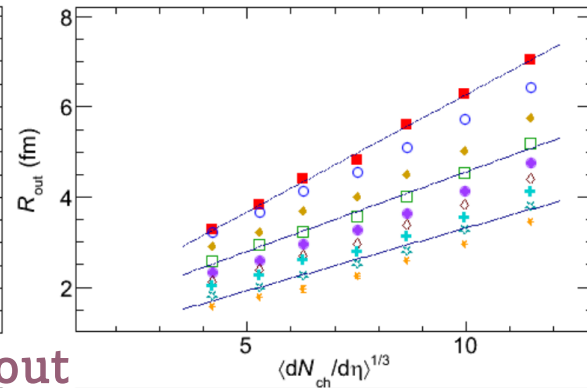
$m_T$  scaling

multiplicity scaling

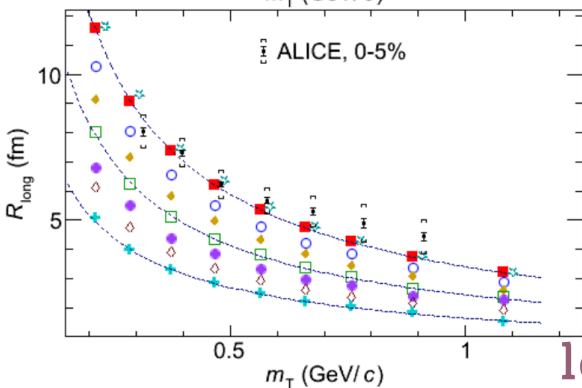
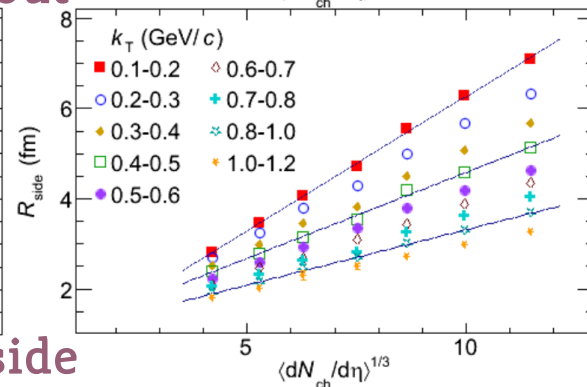
- Both scalings in agreement with calculations from collective (hydrodynamic) models – qualitatively and quantitatively



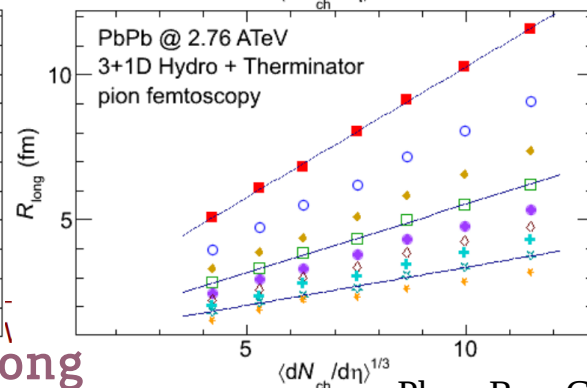
out



side



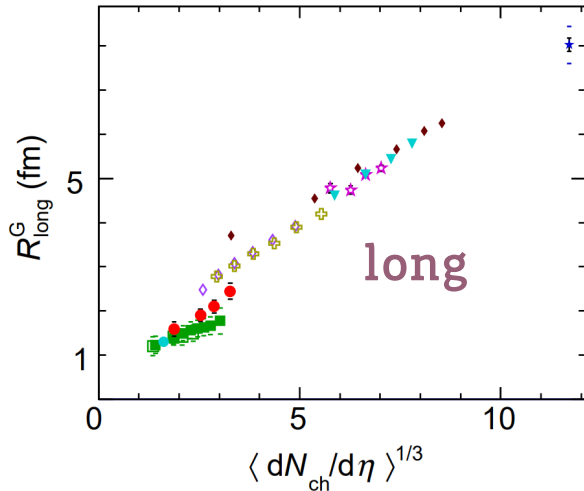
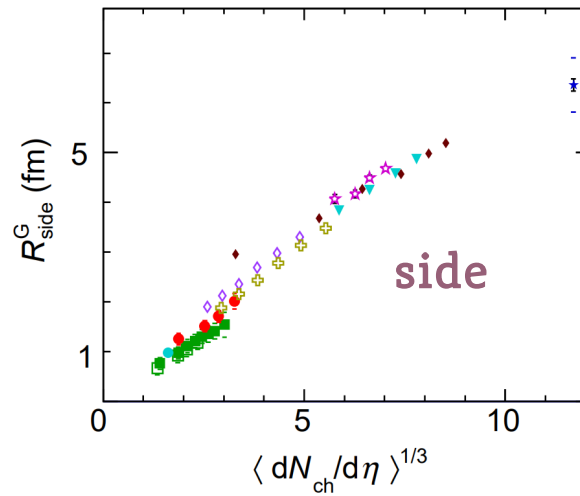
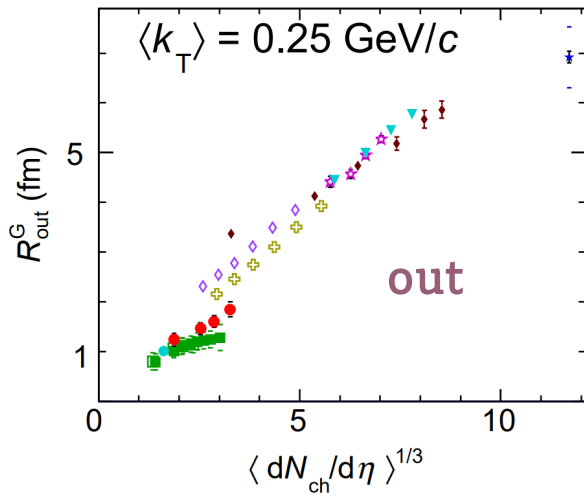
long



Phys. Rev. C 90, 064914 (2014)

# Pion femtoscopy – small systems

- Remarkable agreement between many different experiments and energies
- Linearity extends to pp and p-Pb collisions – although slopes are different
- Initial collision geometry does influence the final measured radii



- ◆ STAR Au-Au  $\sqrt{s_{NN}} = 200 \text{ GeV}$
- ⊕ STAR Cu-Cu  $\sqrt{s_{NN}} = 200 \text{ GeV}$
- ▼ STAR Au-Au  $\sqrt{s_{NN}} = 62 \text{ GeV}$
- ◇ STAR Cu-Cu  $\sqrt{s_{NN}} = 62 \text{ GeV}$
- ★ CERES Pb-Au  $\sqrt{s_{NN}} = 17.2 \text{ GeV}$
- ★ ALICE Pb-Pb  $\sqrt{s_{NN}} = 2760 \text{ GeV}$
- ALICE pp  $\sqrt{s} = 7000 \text{ GeV}$
- ALICE pp  $\sqrt{s} = 900 \text{ GeV}$
- STAR pp  $\sqrt{s} = 200 \text{ GeV}$
- ALICE p-Pb  $\sqrt{s_{NN}} = 5020 \text{ GeV}$

$$R \propto N^{1/3}$$

$$R^3 \propto V \propto N$$

$$\frac{N}{V} = \rho = \text{constant}$$

Density is a function of the measured slope

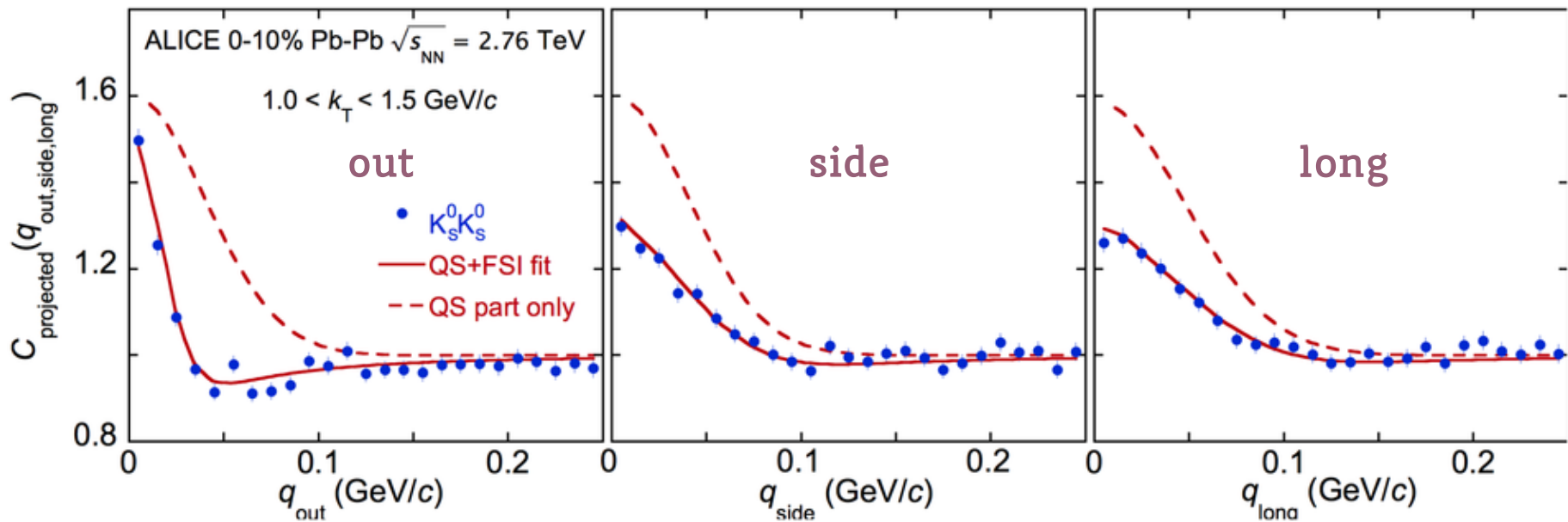
PRC 91 (2015) 044906



# Kaon femtoscopy

- Cleaner signal than pions (less affected by resonance decays)
- Consistency check
  - charged and neutral kaons
  - different detection techniques

Phys. Rev. C96 (2017) 064613

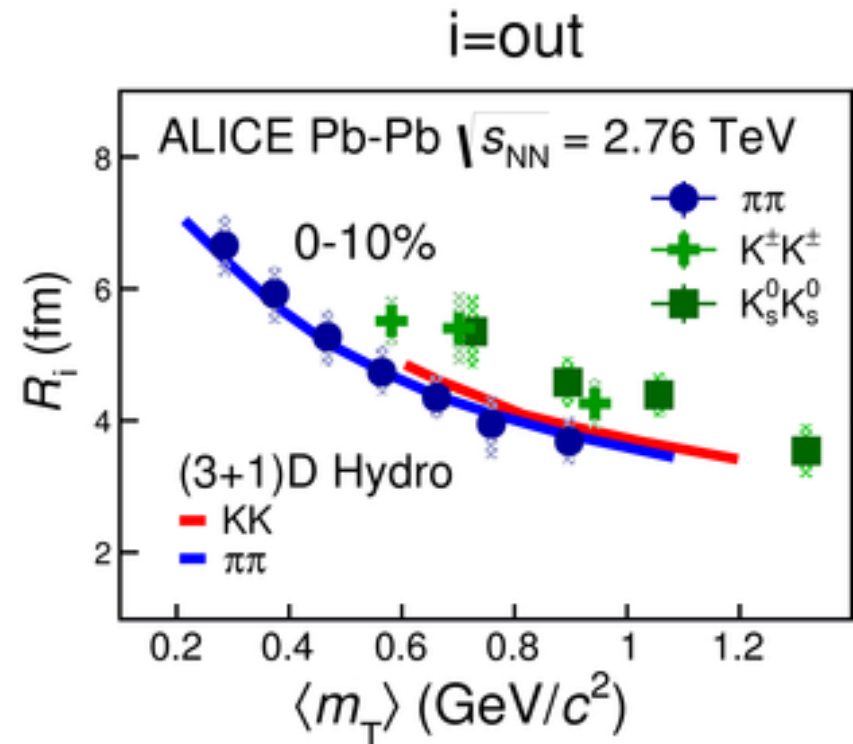


# Kaon femtoscopy

- **Results:**

- **broken  $m_T$ -scaling** → radii larger than for pions (possibly due to rescattering through  **$K^*$  resonance**)
- comparison to hydro **without rescattering phase** → **scaling present**

Phys. Rev. C96 (2017) 064613





# Kaon femtoscopy



- **Results:**

Phys. Rev. C96 (2017) 064613

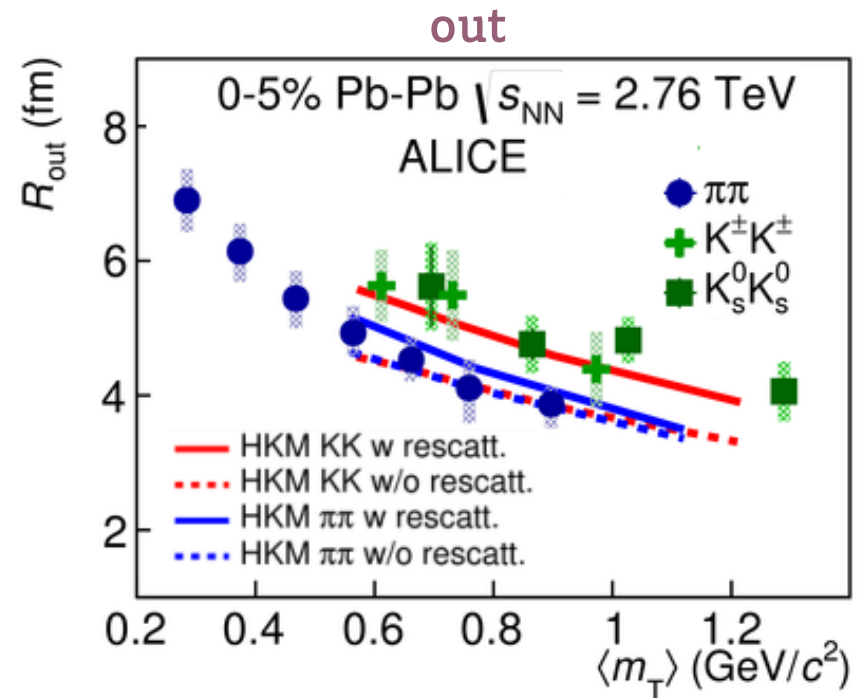
- **broken  $m_T$ -scaling**  $\rightarrow$  radii larger than for pions (possibly due to rescattering **through  $K^*$  resonance**)
- comparison to hydro **with rescattering phase**  $\rightarrow$  **scaling broken**
- model does describe the data

- **From  $m_T$  scaling of  $R_{long}$  we can extract times for pions and kaons:**

$$R_{long}^2 = \tau_{max}^2 \frac{T_{max}}{m_T \cosh y_T} \left( 1 + \frac{3T_{max}}{2m_T \cosh y_T} \right)$$

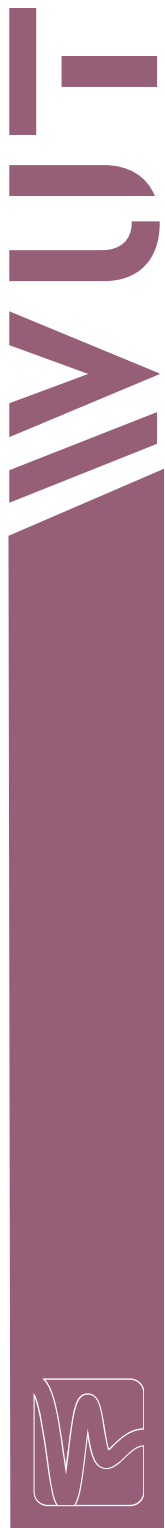
Y. Sinyukov, *et al* Nucl. Phys. A946 (2016) 227–239

- **emission time of kaons is larger than that of pions by 2.1 fm**



method	$T$ (GeV)	$\tau_\pi$ (fm/c)	$\tau_K$ (fm/c)
fit with Eq. (9)	0.144	$9.5 \pm 0.2$	$11.6 \pm 0.1$

# Pion-kaon femtoscopy



$$C(\vec{q}) = \int |\Psi(\vec{q}, \vec{r})|^2 S(\vec{r}) d^3 r$$

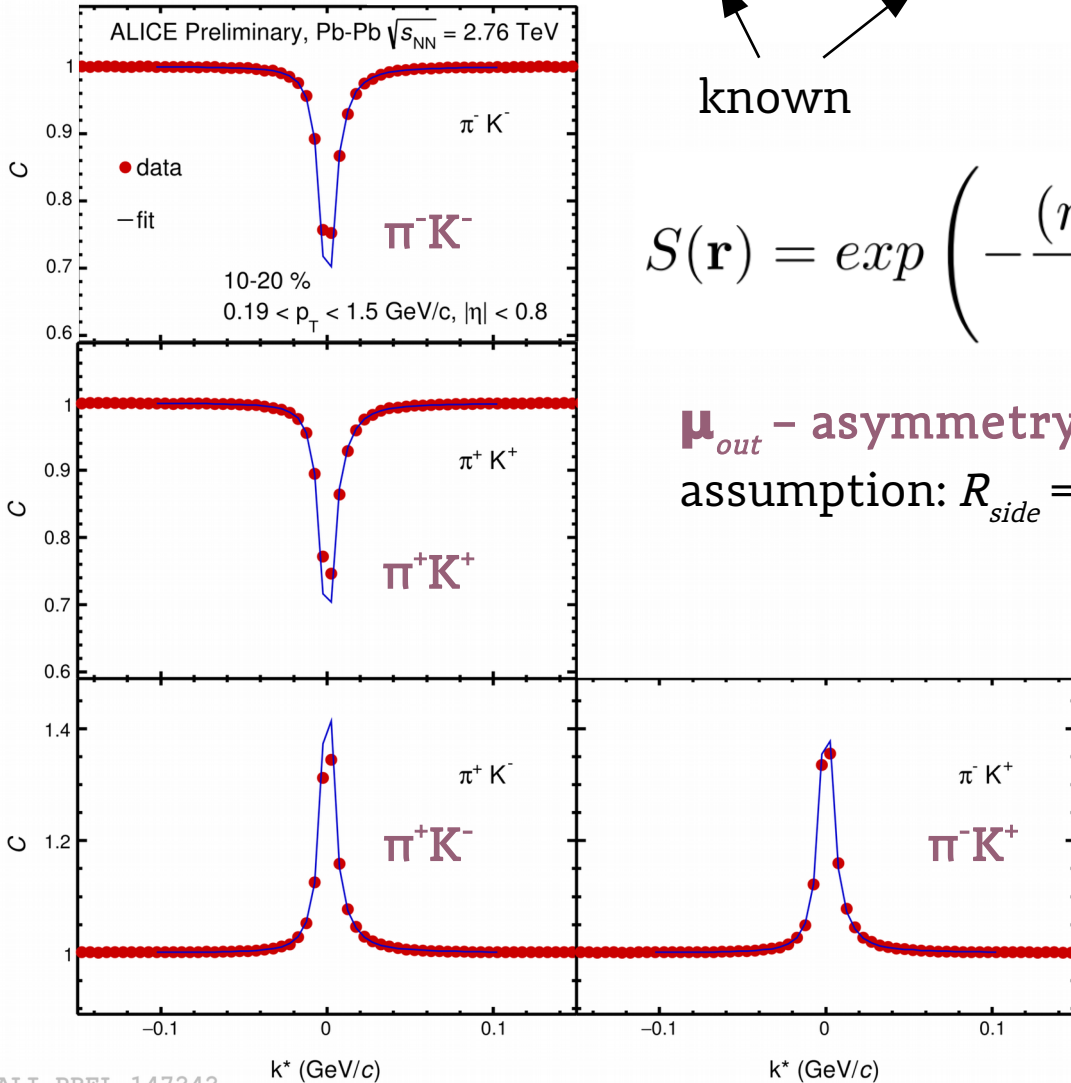
known

unknown

$$S(\mathbf{r}) = \exp\left(-\frac{(r_{out} - \mu_{out})^2}{R_{out}^2} - \frac{r_{side}^2}{R_{side}^2} - \frac{r_{long}^2}{R_{long}^2}\right)$$

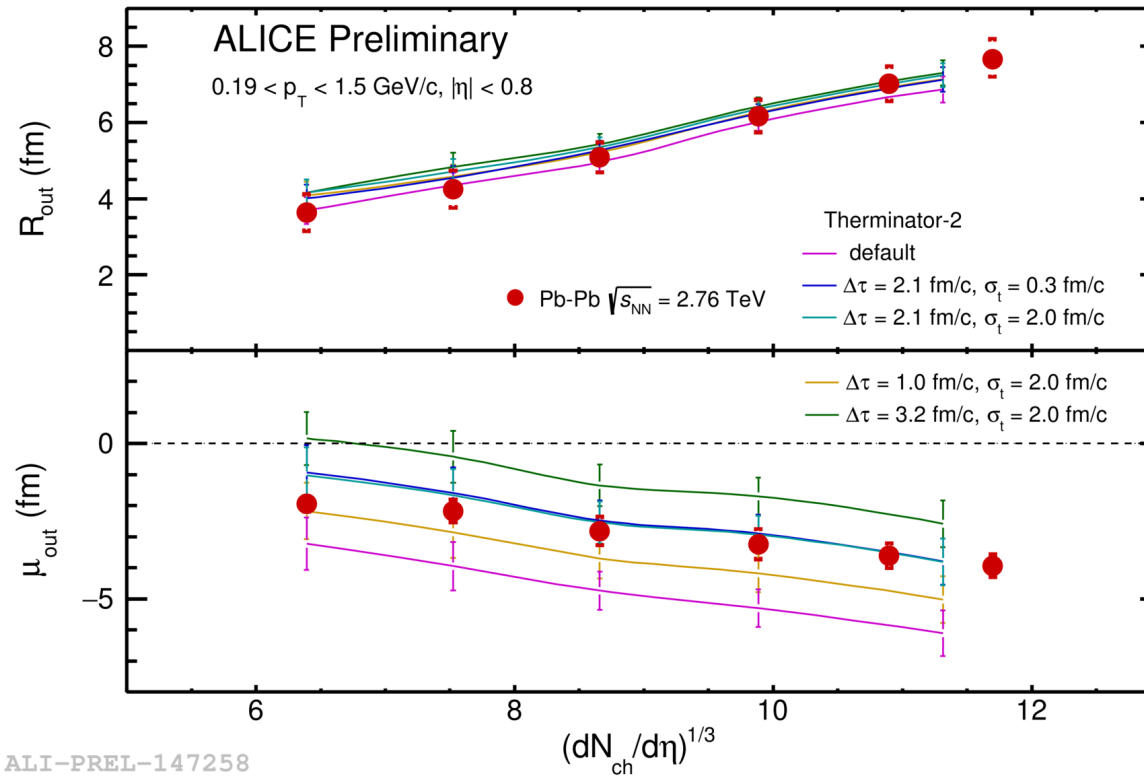
$\mu_{out}$  - asymmetry in out direction

assumption:  $R_{side} = R_{out}$ ,  $R_{long} = 1.3R_{out}$



- 4 different pair combinations
- Independent check of rescattering phase (emission asymmetry)

# Pion-kaon femtoscopy

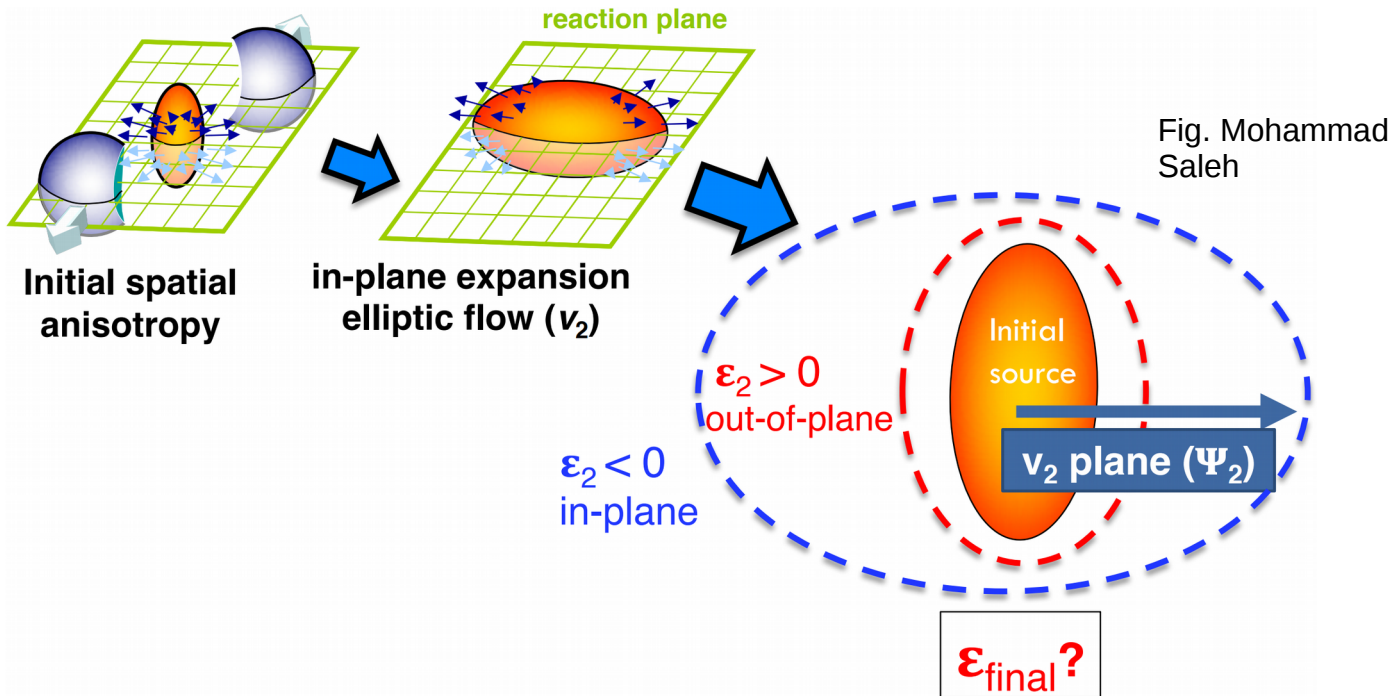


- **ALICE data (points):** significant negative pion-kaon emission asymmetry is observed which increases with centrality
- **Model (lines):** kaon emission delay of 2.1 fm/c introduced in Therminator 2 model predicts the measured asymmetry

→ different particle species freeze out at different times



# Azimuthally differential pion femtoscopy



- Can we measure the final source eccentricity and how does it relate to anisotropy in initial stage?
- We expect the source to be more spherical at freeze-out wrt initial state due to stronger in-plane expansion



# Azimuthally differential pion femtoscopy

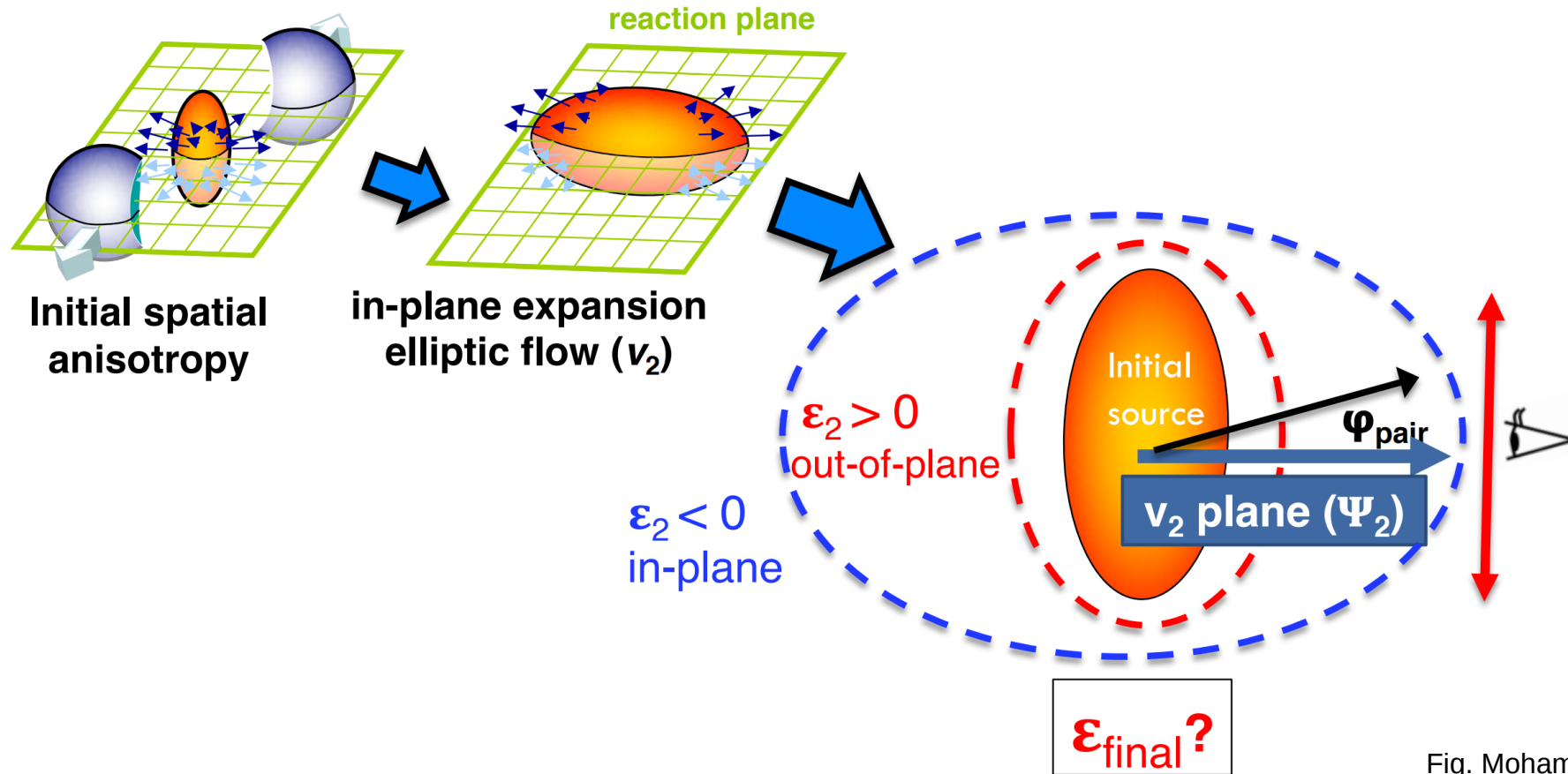


Fig. Mohammad Saleh

- We can use femtoscopy and look at the source size from different angles!



# Azimuthally differential pion femtoscopy

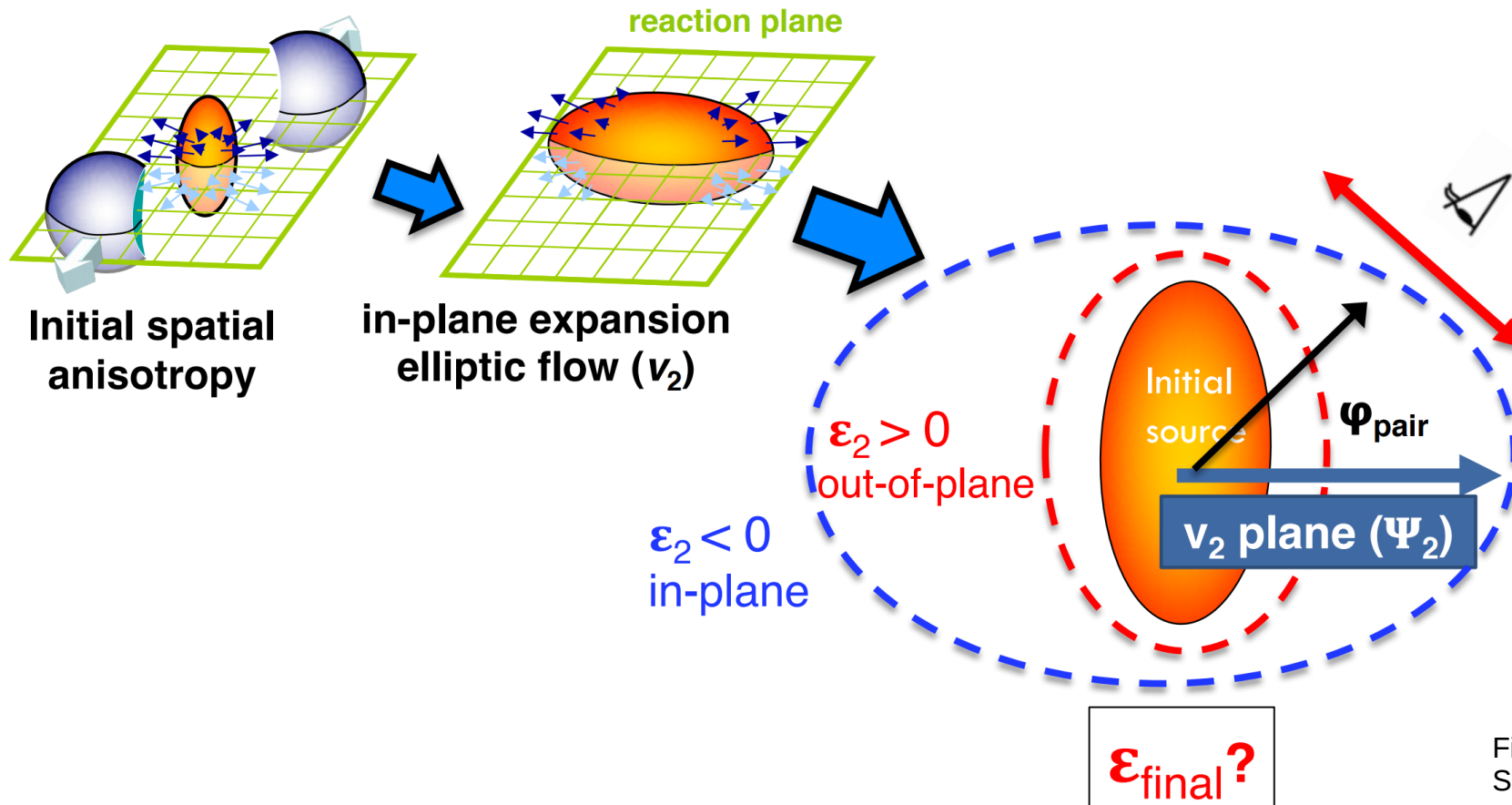


Fig. Mohammad Saleh

- We can use femtoscopy and look at the source size from different angles!



# Azimuthally differential pion femtoscopy



Fig. Mohammad Saleh

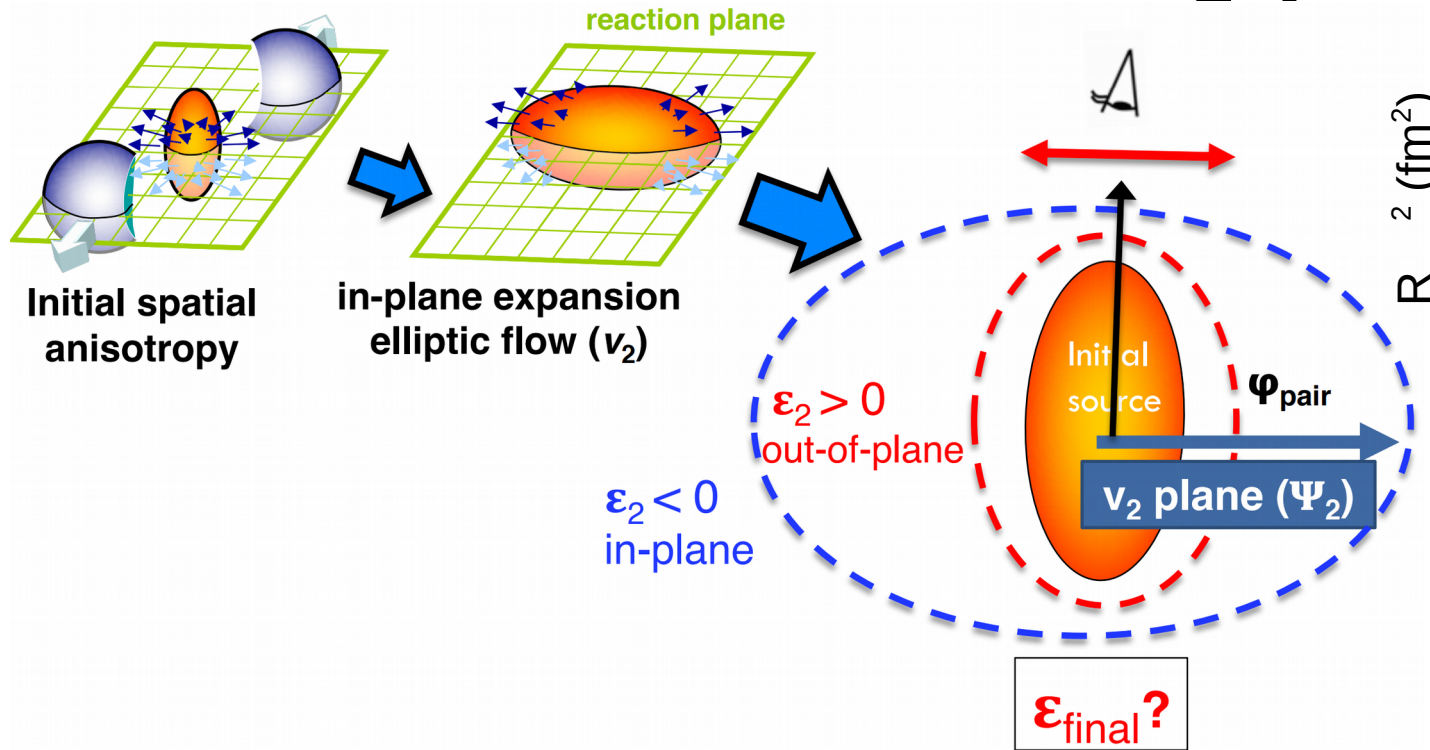
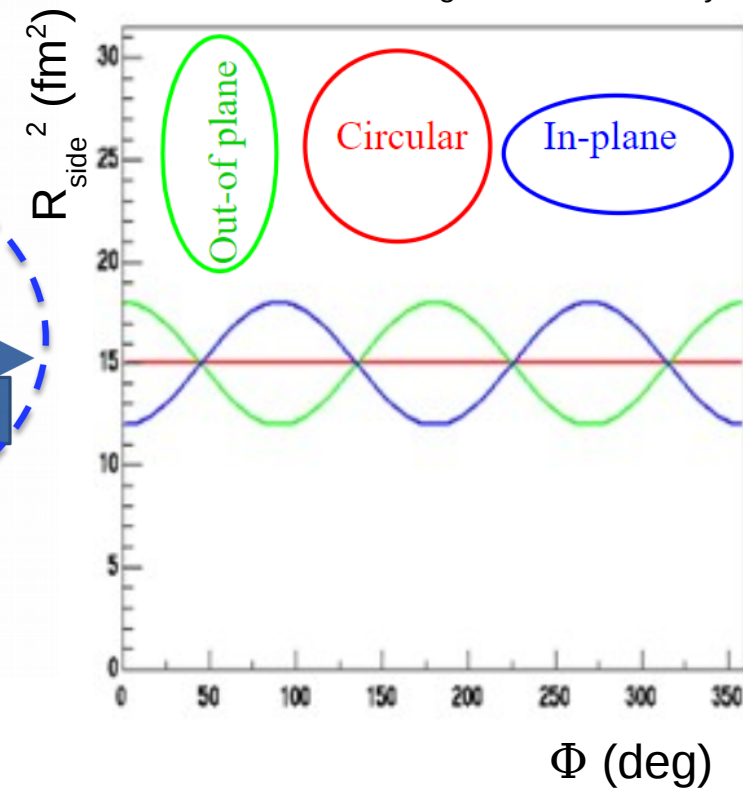


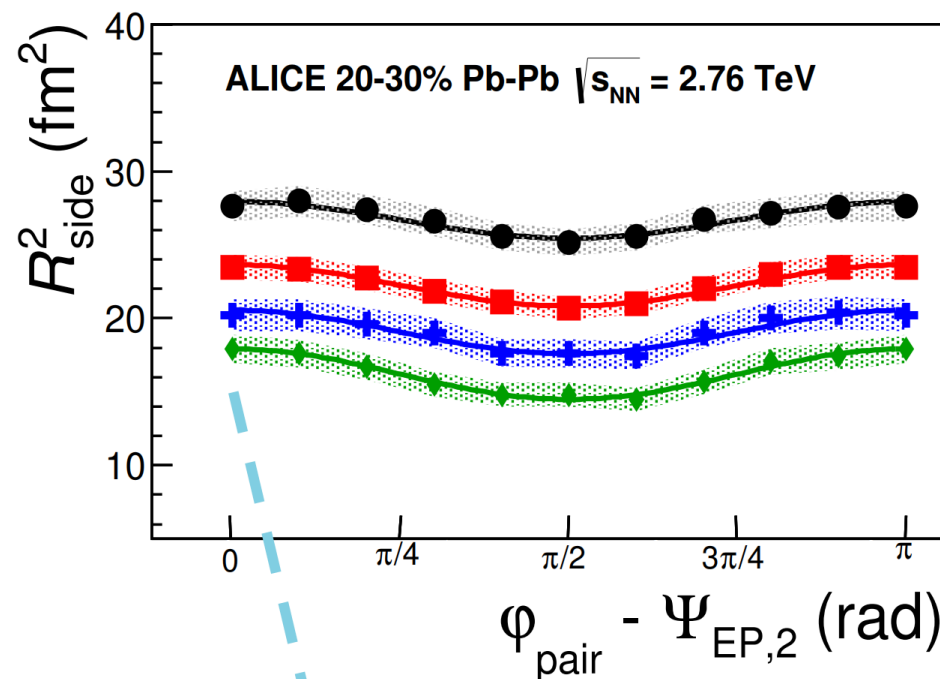
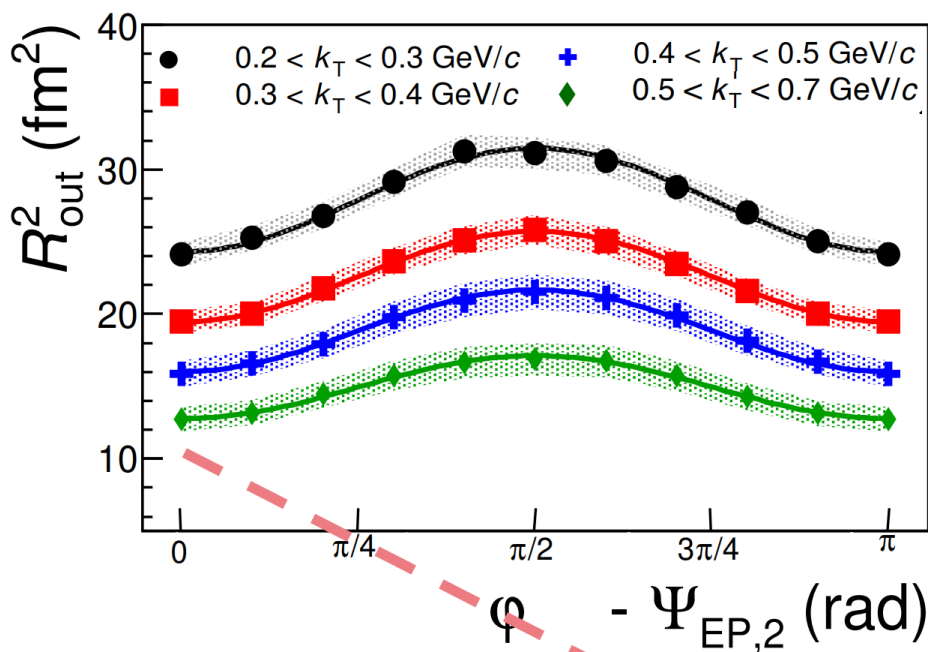
Fig. Richard Lednicky



- We can use femtoscopy and look at the source size from different angles!
- We expect radii oscillations, depending on the final anisotropy of the source

# Azimuthally differential pion femtoscopy

PRL 118 (2017), 222301



## Results:

- out and side radii oscillate out-of-phase, as it should be!
- Pion source at the freeze-out is elongated in the out-of-plane direction

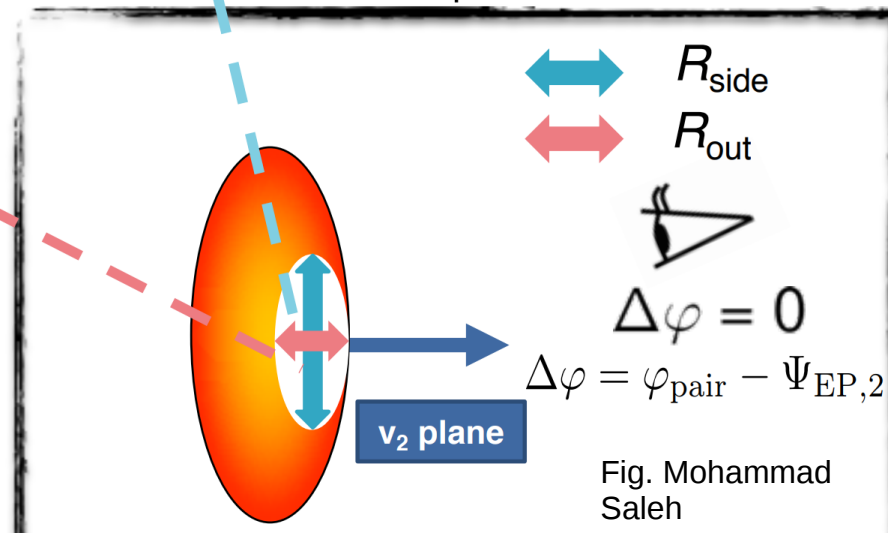


Fig. Mohammad Saleh





# Azimuthally differential pion femtoscopy

PRL 118 (2017), 222301

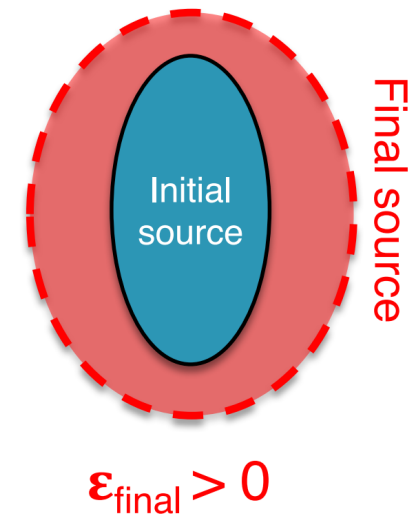
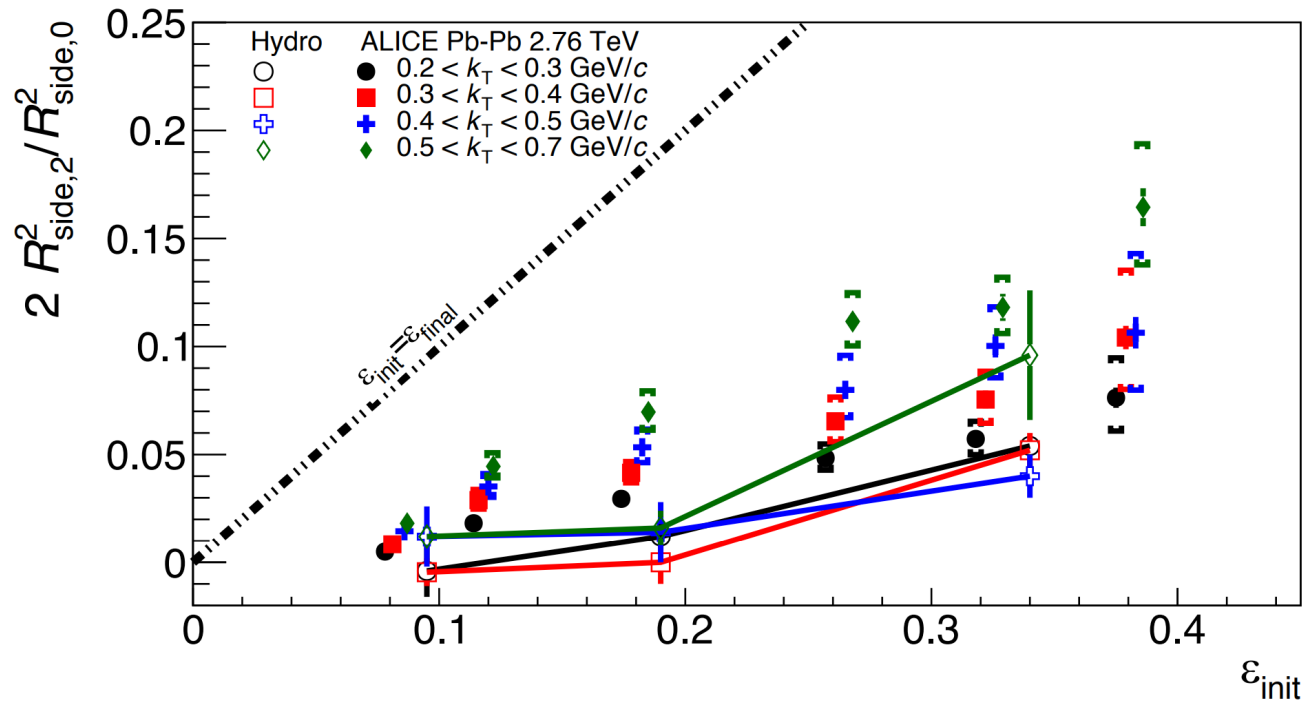


Fig. Mohammad Saleh

## Results:

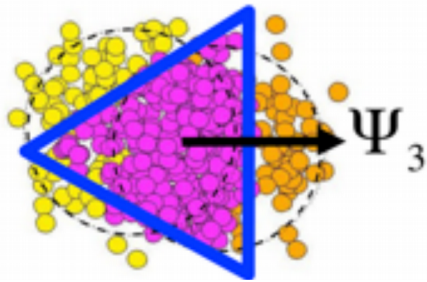
- The final source eccentricity is **smaller than initial anisotropy**, but still exhibits an **out-of-plane elongated source**, even after stronger in-plane expansion
- Hydrodynamic models agree qualitatively but predict **more isotropic source**
- **BUT, can we do more?**



# Azimuthally differential pion femtoscopy

PLB 785 (2018) 320

Initial geometry



Final source shape

??

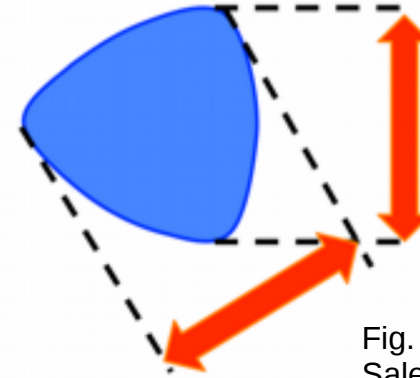
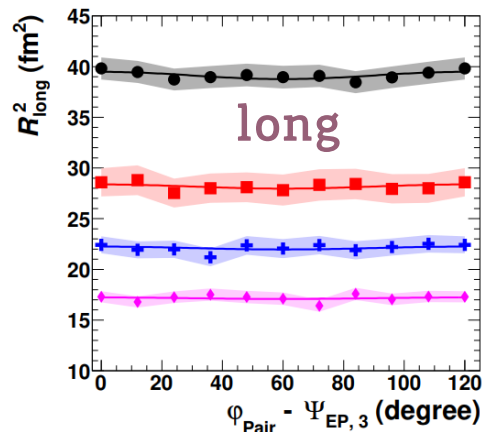
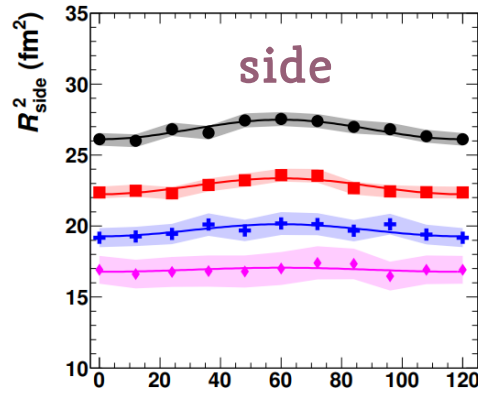
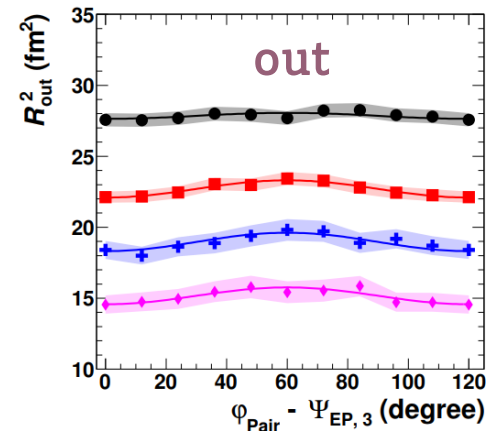


Fig. Mohammad Saleh



ALICE Pb-Pb  $\sqrt{s_{NN}} = 2.76$  TeV  
 Centrality 20-30%

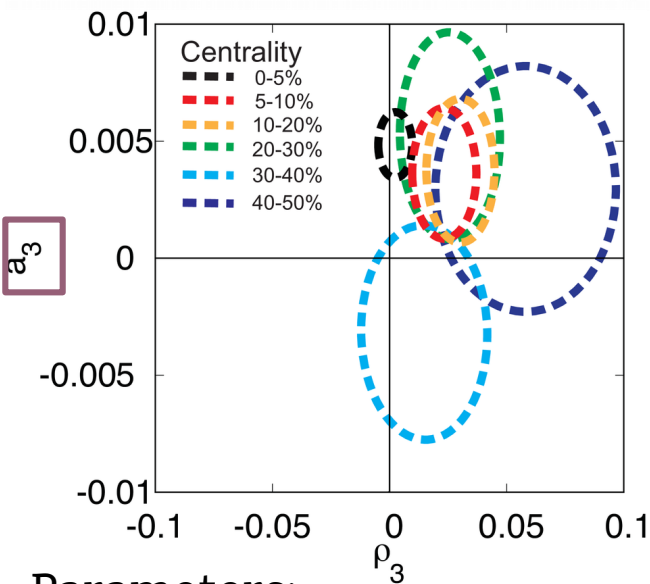
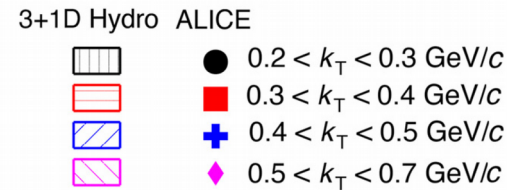
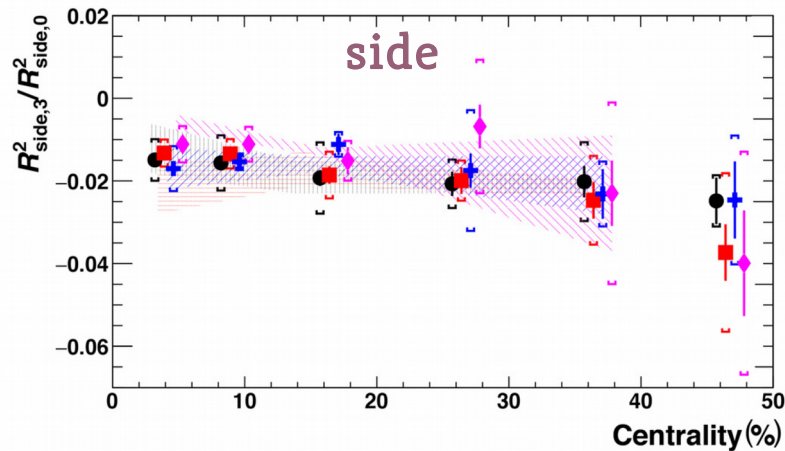
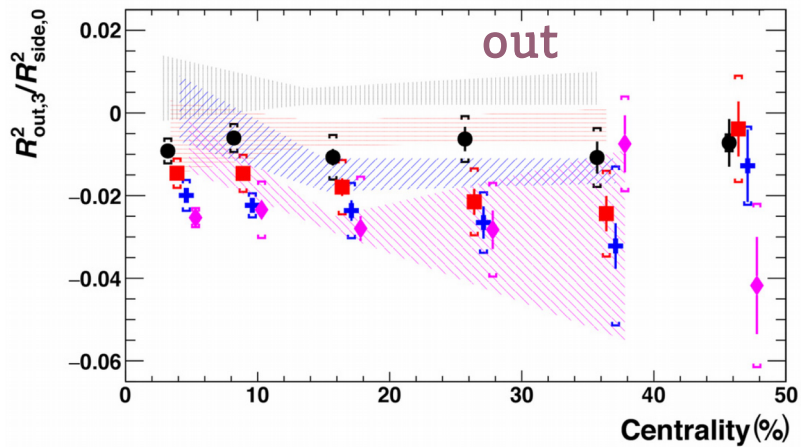
- $0.2 < k_T < 0.3$  GeV/c
- $0.3 < k_T < 0.4$  GeV/c
- ⊕  $0.4 < k_T < 0.5$  GeV/c
- ◆  $0.5 < k_T < 0.7$  GeV/c

- Pion source size relative to the 3<sup>rd</sup> harmonic event plane in Pb-Pb at  $\sqrt{s_{NN}}=2.76$  TeV
- Out and side radii oscillate in-phase
- No oscillations for long radius
- Such oscillations are consistent if the radial expansion includes third harmonic modulation
- For triangular but static source no oscillations are expected



# Azimuthally differential pion femtoscopy

PLB 785 (2018) 320



Parameters:

$a_3$  – final source anisotropy

$\rho_3$  – transverse flow

- Side radii oscillations agree well with hydrodynamic models, others only qualitatively
- We can also calculate the final anisotropy:
 
$$R(\phi) = R_0 \left( 1 - \sum_{n=2}^{\infty} a_n \cos(n(\phi - \Psi_n)) \right)$$
- $a_3$  is close to zero, which is significantly smaller than the initial triangular eccentricities which are on the order of 0.2-0.3
- **Conclusion: initial state triangularity is washed out at freeze-out due to triangular flow**

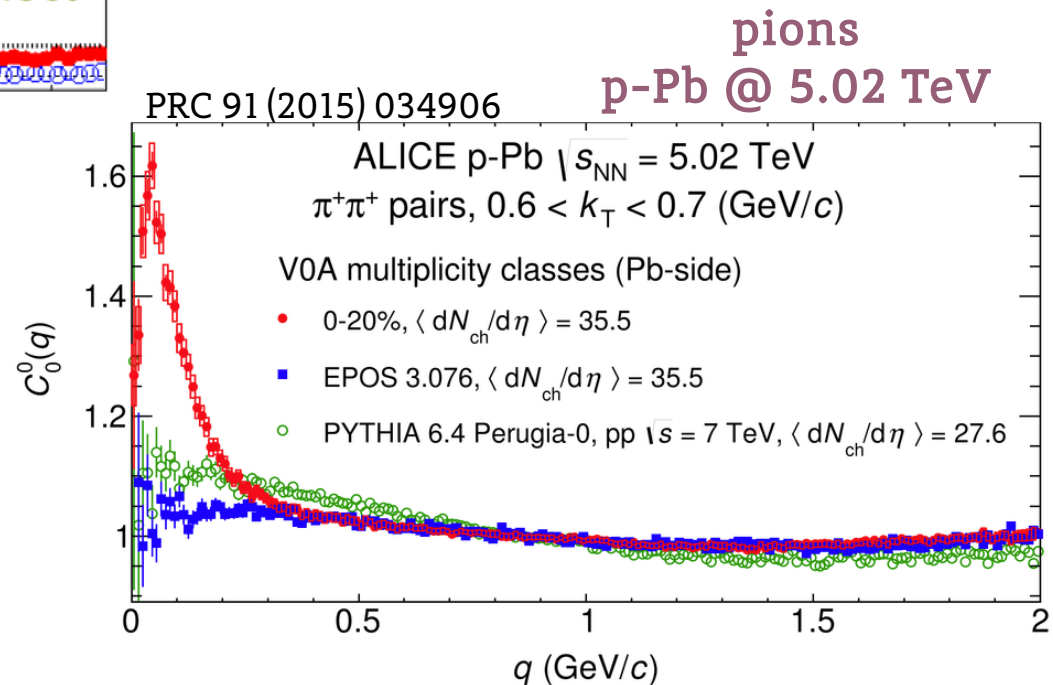
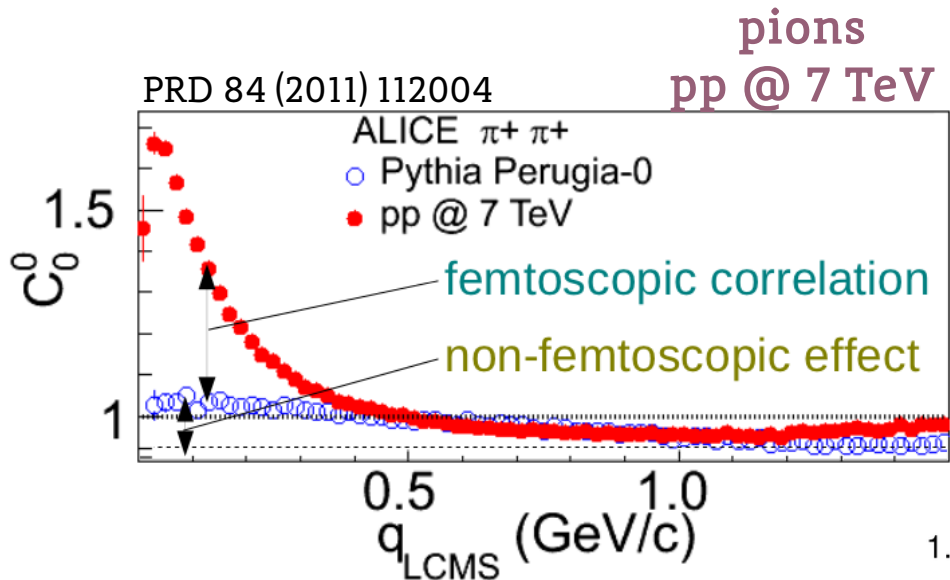
**What about pp and p-Pb frontier?**

**New developments in small systems**



# Non-femtoscopic correlations

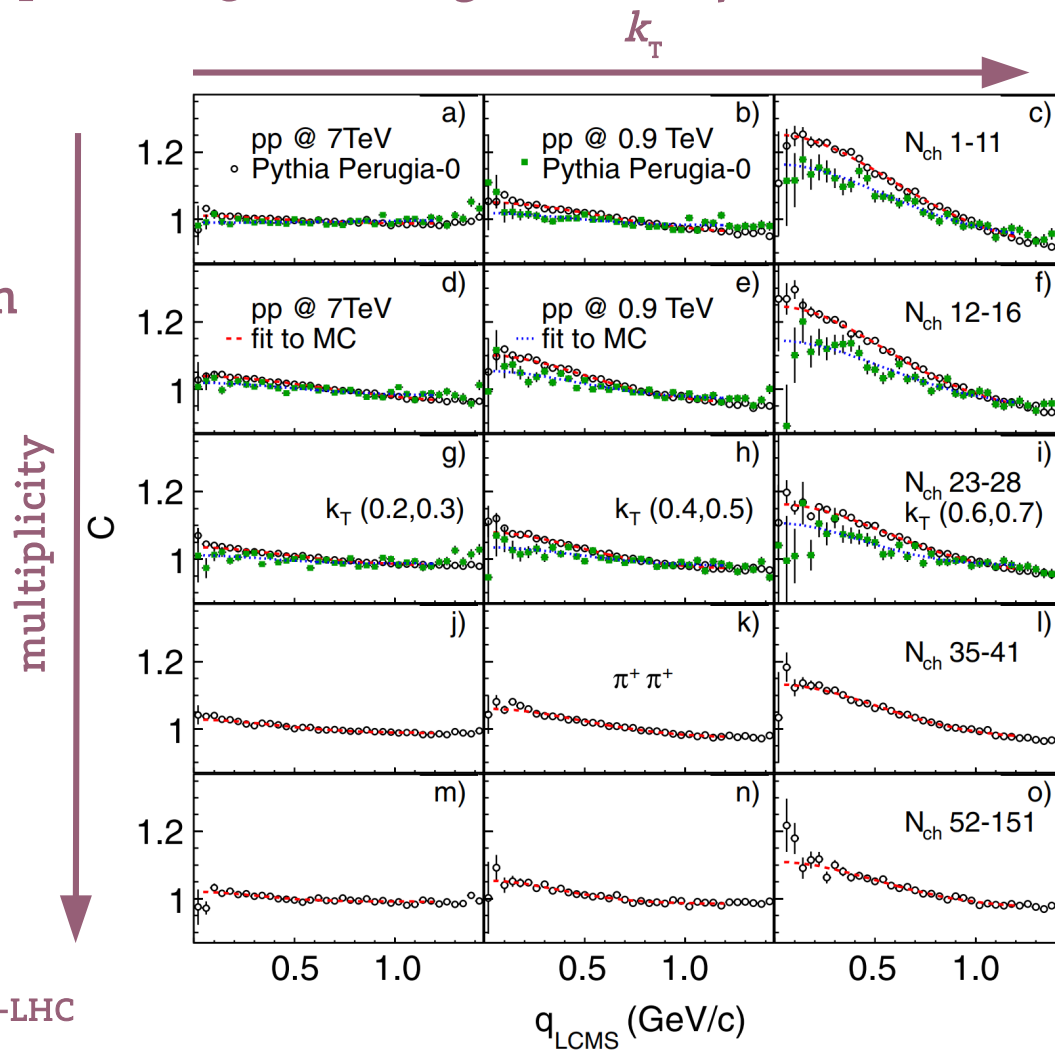
- In small systems we are looking for **signs of collectivity**
- In pp and p-Pb significant non-femtoscopic correlations are present – **hypothesis of (mini-)jet origin**



# Non-femtoscopic correlations

- In small systems we're looking for signs of collectivity
- In pp and p-Pb significant non-femtoscopic correlations are present – hypothesis of (mini-)jet origin
- Non-femtoscopic background **significantly limits** the accessible range in  $k_T$

pions  
Pythia simulation



# How do (mini-)jets look like?

## Let's look at angular space

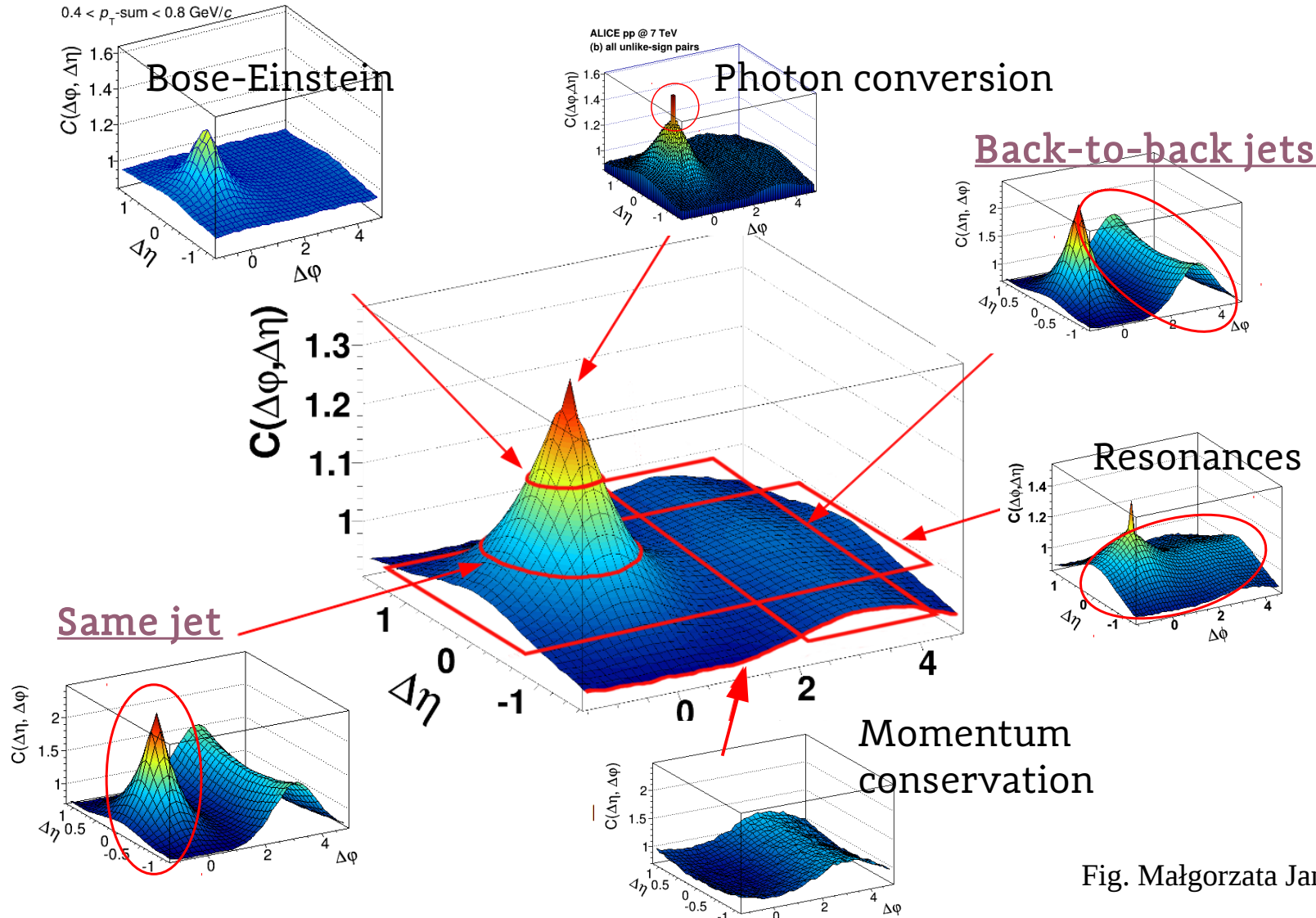


Fig. Małgorzata Janik

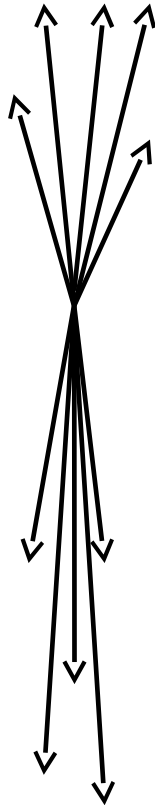
Jets produce streams of particles close to each other – like femto correlations!



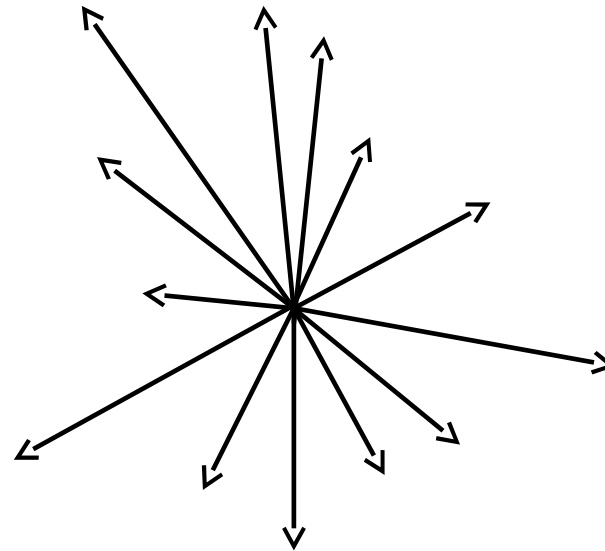
# Sphericity

- How about selecting events which do have (mini-)jets or not?

jetty



spherical



$$S_{xy}^L = \frac{1}{\sum_i p_{Ti}} \sum_i \frac{1}{p_{Ti}} \begin{pmatrix} p_{xi}^2 & p_{xi}p_{yi} \\ p_{yi}p_{xi} & p_{yi}^2 \end{pmatrix}$$

$$S_T = \frac{2\lambda_2}{\lambda_1 + \lambda_2} \Rightarrow S_T = \begin{cases} \approx 0 & \text{Jet-like} \\ \approx 1 & \text{Spherical} \end{cases}$$

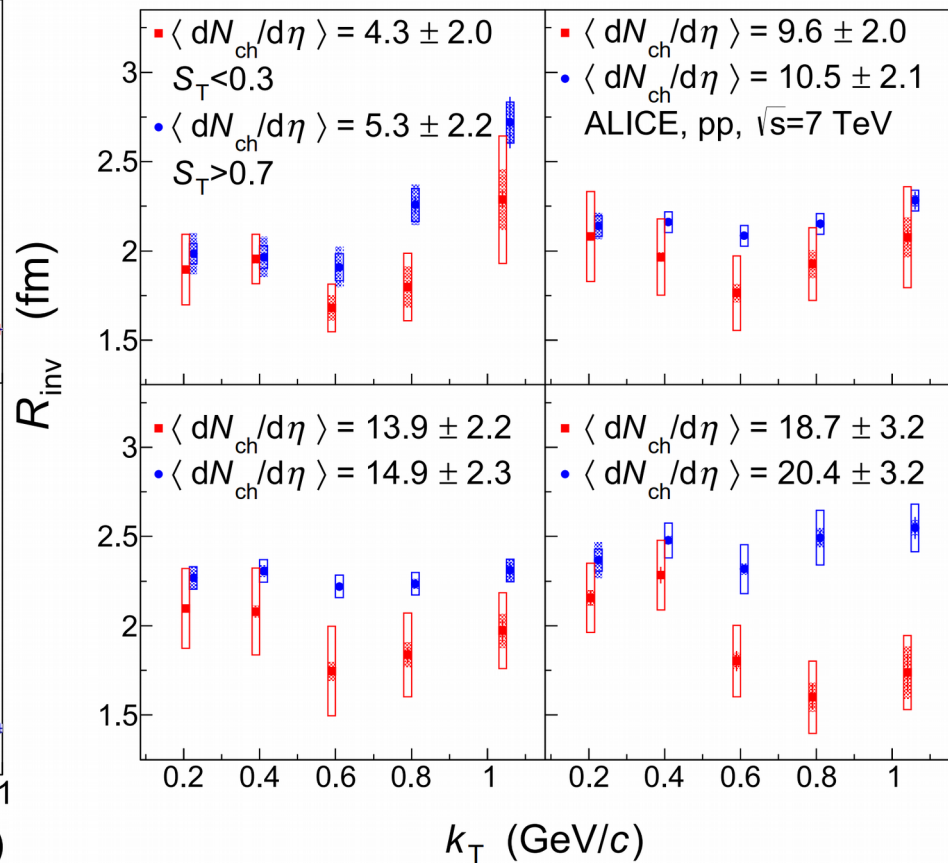
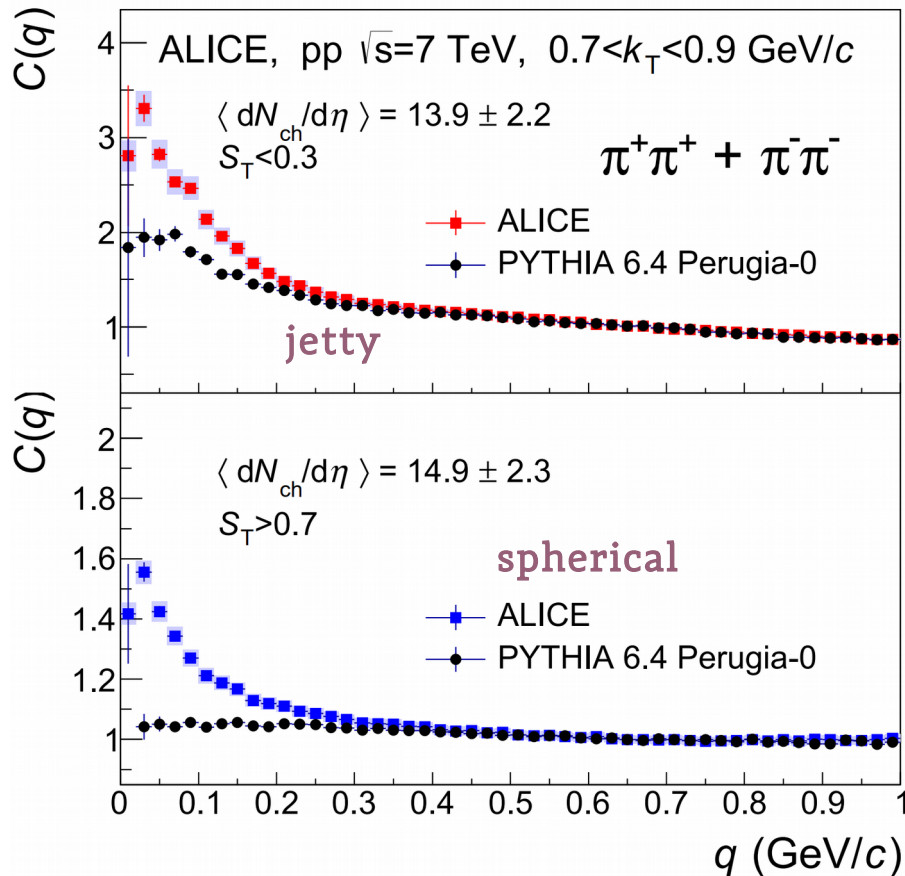




# Sphericity

- (Mini-)jet contribution is the main source of underlying correlation
- Sphericity analysis allows to extend the studied  $k_T$  range
- $k_T$  dependence of 1D radii in PRF not well explored
- **Next step: 3D analysis in LCMS to check for the  $k_T$  scaling**

arXiv:1901.05518



# Going beyond the system size



# Beyond the system size

$$C(q) = \int S(r) |\Psi(q, r)|^2 d^4 r$$

$$q = 2 \cdot k^* = p_1 - p_2$$

measured correlation

emission function  
(source size/shape)

pair wave function  
(includes cross section)

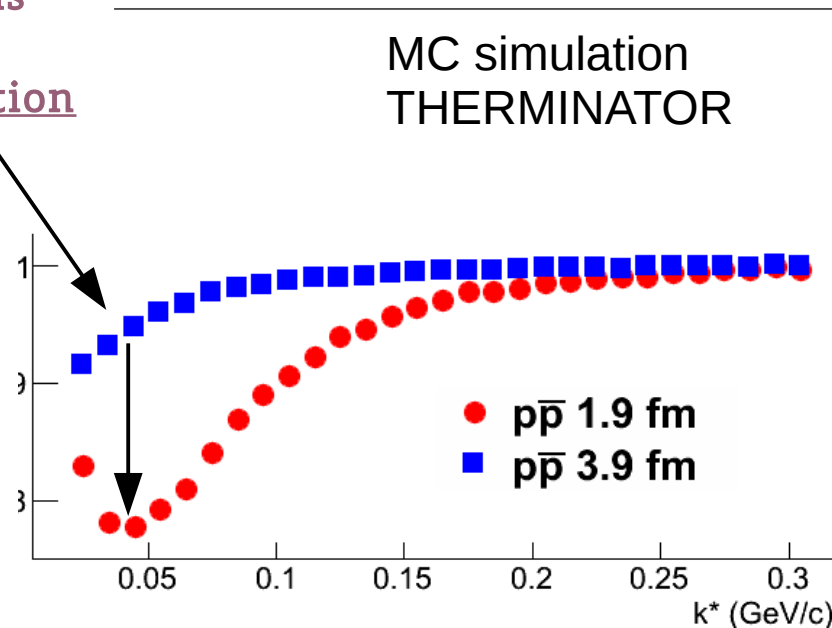
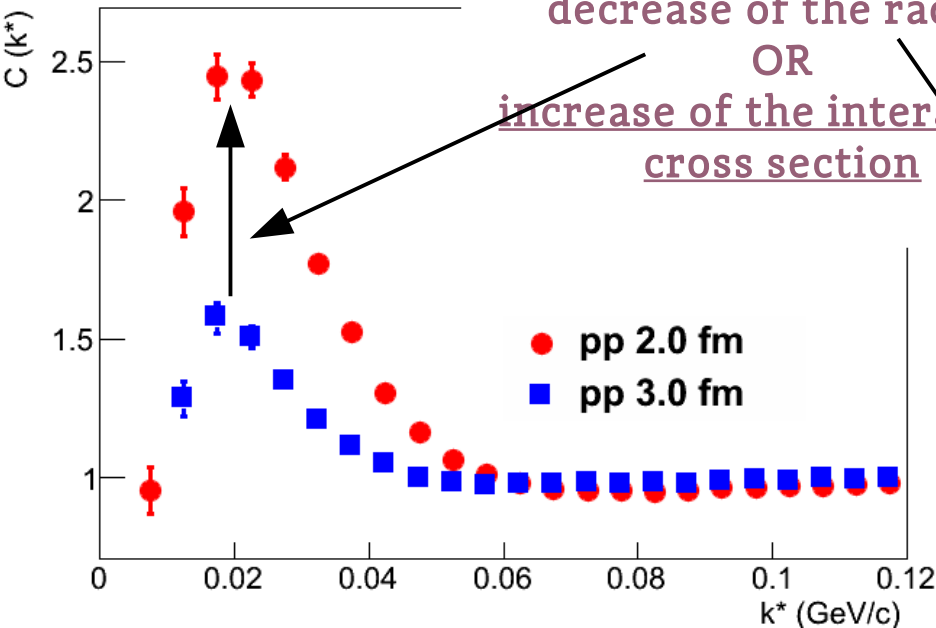
increase of (anti)correlation

=

decrease of the radius

OR

increase of the interaction  
cross section



- One can use femtoscopy to measure interactions for those pairs where it is poorly known or not known at all!
- Powerful tool when scattering experiments are not possible



# Beyond the system size

$$C(q) = \int S(r) |\Psi(q, r)|^2 d^4 r \quad q = 2 \cdot k^* = p_1 - p_2$$

measured correlation      emission function (source size/shape)      pair wave function (includes cross section)

pair wave function  $\longrightarrow \Psi = \exp(-ik^* r) + f \frac{\exp(ik^* r)}{r}$  s-wave scattering approximation

scattering amplitude  $\longrightarrow f^{-1}(k^*) = \frac{1}{f_0} + \frac{1}{2} d_0 k^{*2} - ik^*$  effective range approximation

- If only Strong FSI is present:

Lednický equation

$$C(k^*) = 1 + \sum_s \rho_s \left[ \frac{1}{2} \left| \frac{f^s(k^*)}{R} \right|^2 \left( 1 - \frac{d_0^s}{2\sqrt{\pi}R} \right) + \frac{2\Re f^s(k^*)}{\sqrt{\pi}R} F_1(2k^*R) - \frac{\Im f^s(k^*)}{R} F_2(2k^*R) \right]$$

Sov. J. Nucl. Phys., 35, 770 (1982)

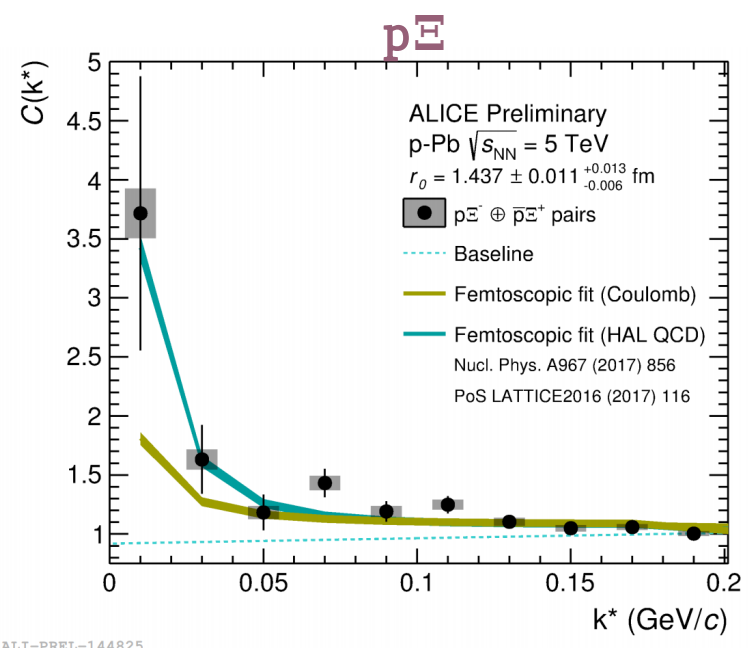
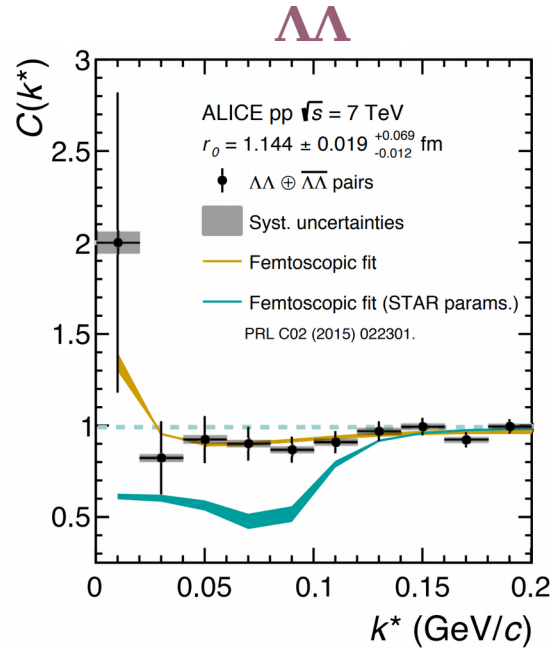
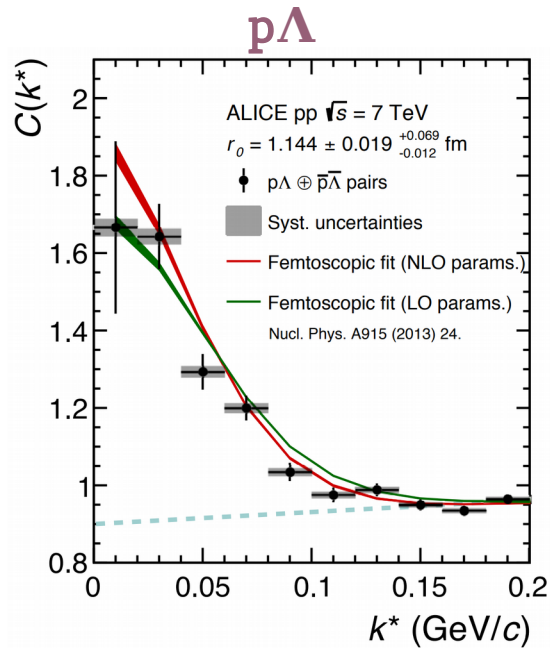
where  $\rho_s$  are the spin fractions

- The correlation function is characterized by **three parameters**:

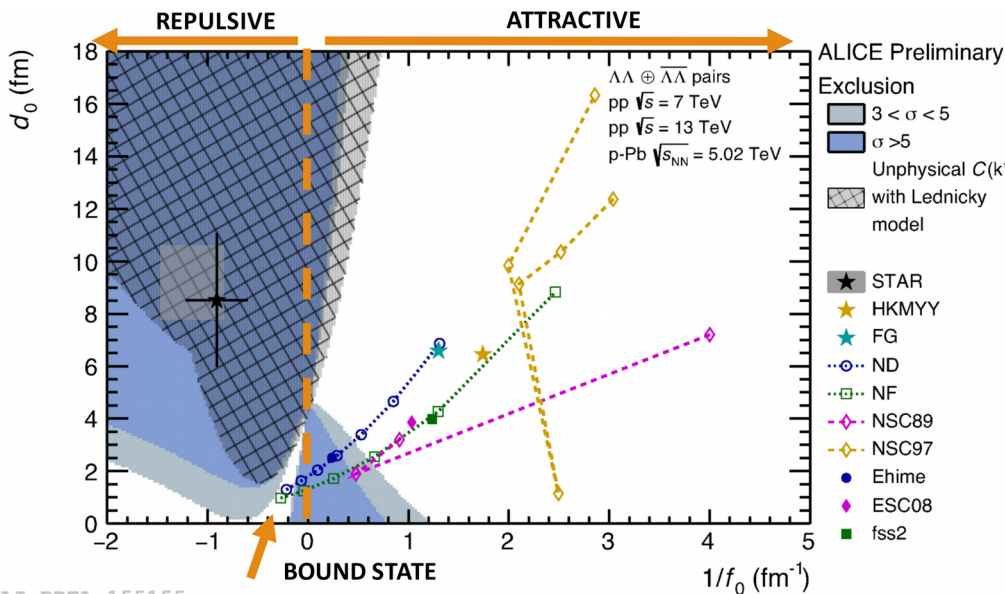
- radius  $R$ , scattering length  $f_0$ , and effective radius  $d_0$
- cross section  $\sigma$  (at low  $k^*$ ) is simply:  $\sigma = 4\pi |f|^2$



# Baryon-baryon correlations



ALI-PREL-144825



ALI-PREL-155155

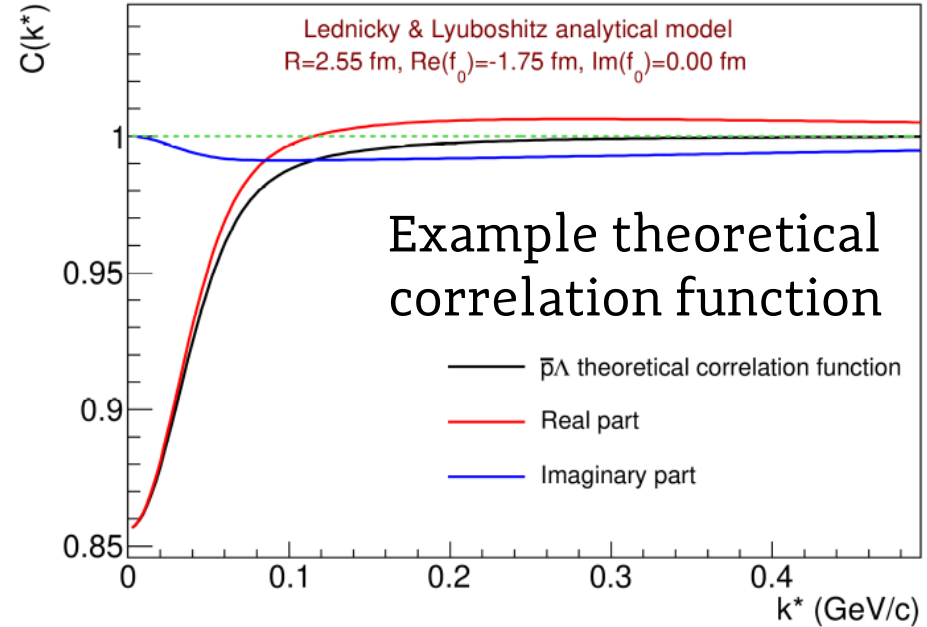
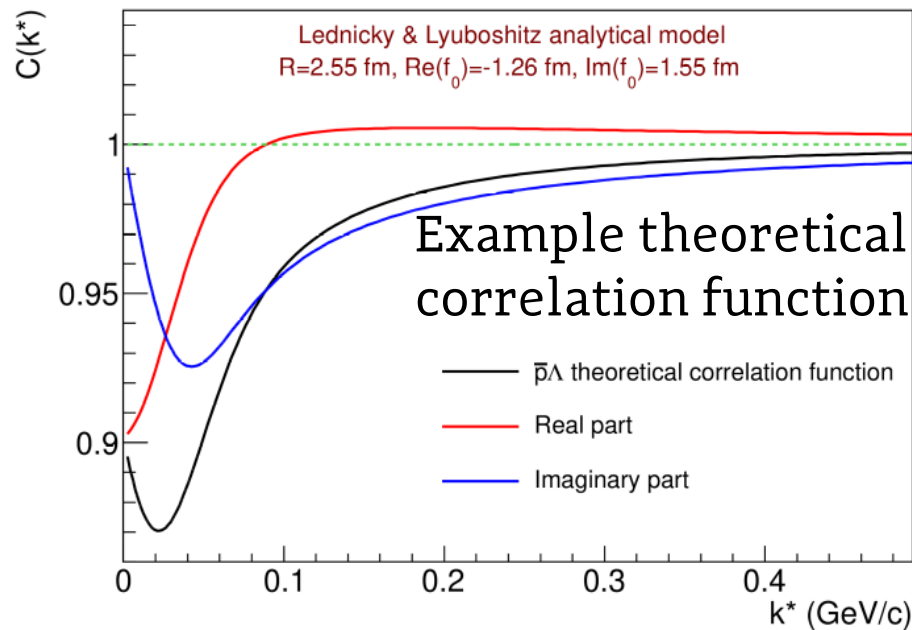
- Constraining lambda-lambda scattering parameters and bound states (H-dibaryon)
- First measurement of attractive proton- $\Xi$  potential

For details see CERN-LHC Seminar  
 25 Sep 2018  
<https://indico.cern.ch/event/749074/>

PRC 99 (2019) 024001



# Baryon-antibaryon correlations

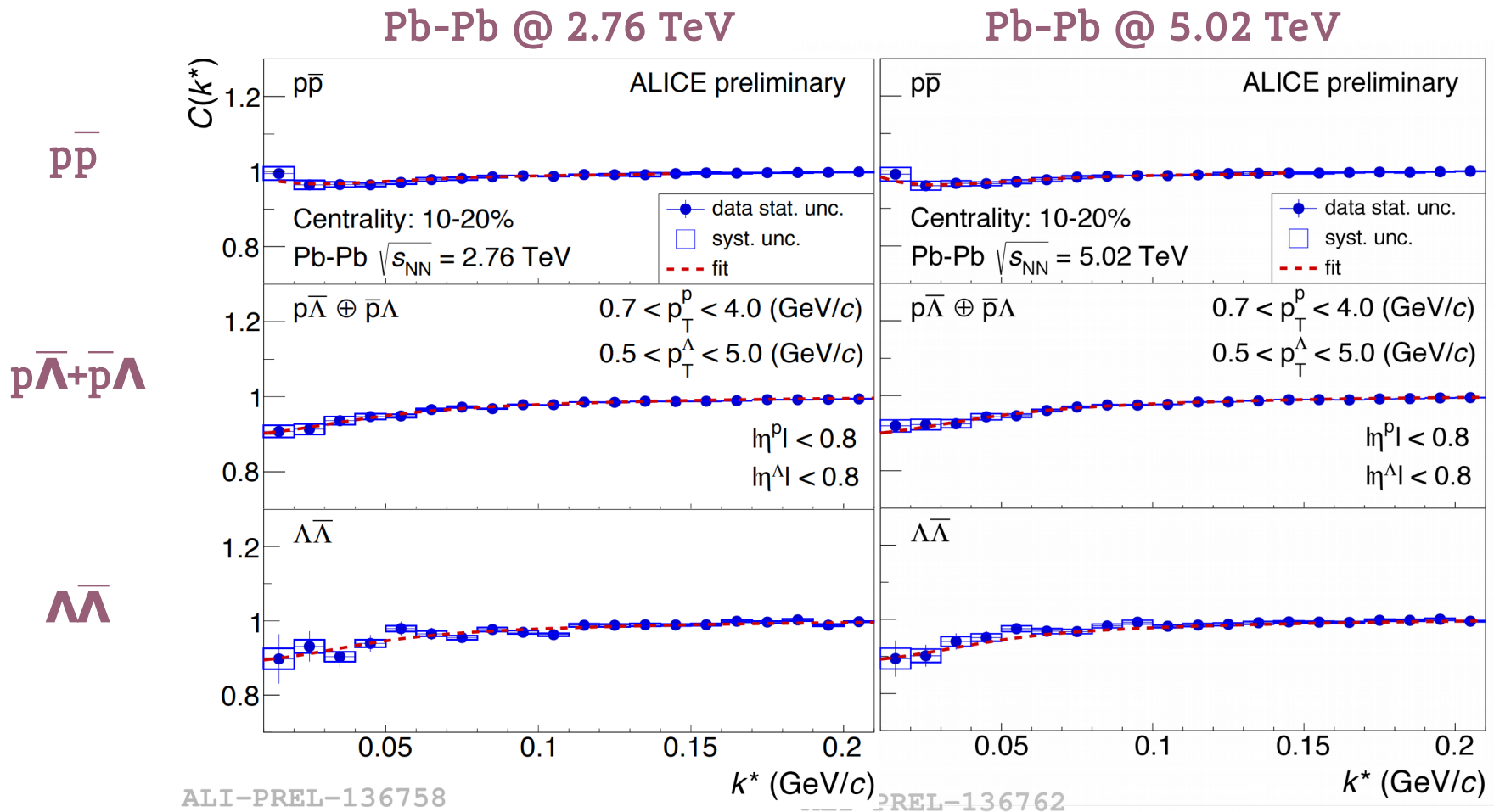


- Real and imaginary part of scattering length have **distinctively different contributions**
- Contribution from  $\text{Re}(f_0)$  is either positive or negative but **very narrow** (up to 100 MeV/c) in  $k^*$
- The  $\text{Im}(f_0)$  accounts for baryon-antibaryon annihilation and produces a **wide** (hundreds of MeV) **negative correlation**

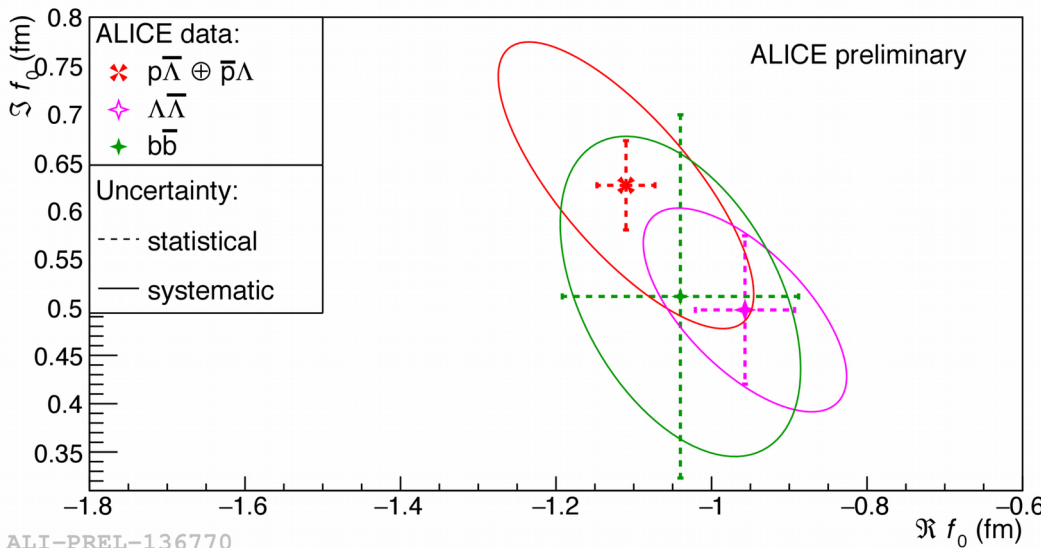
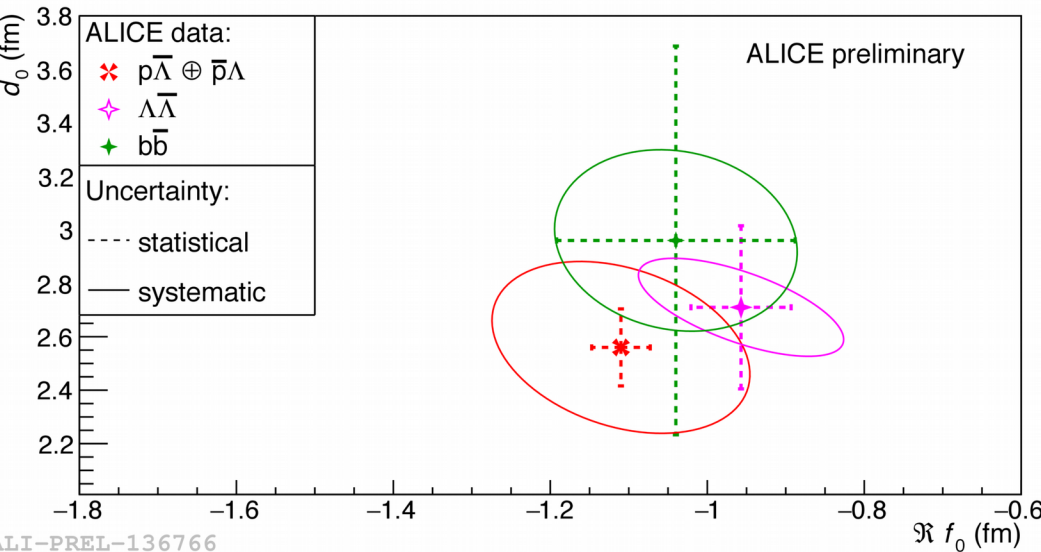


# Baryon-antibaryon correlations

- Correlation functions measured for two collision energies
- $\Lambda\bar{\Lambda}$  correlations measured for the first time
- Wide negative correlation seen for all pairs



# Baryon-antibaryon correlations

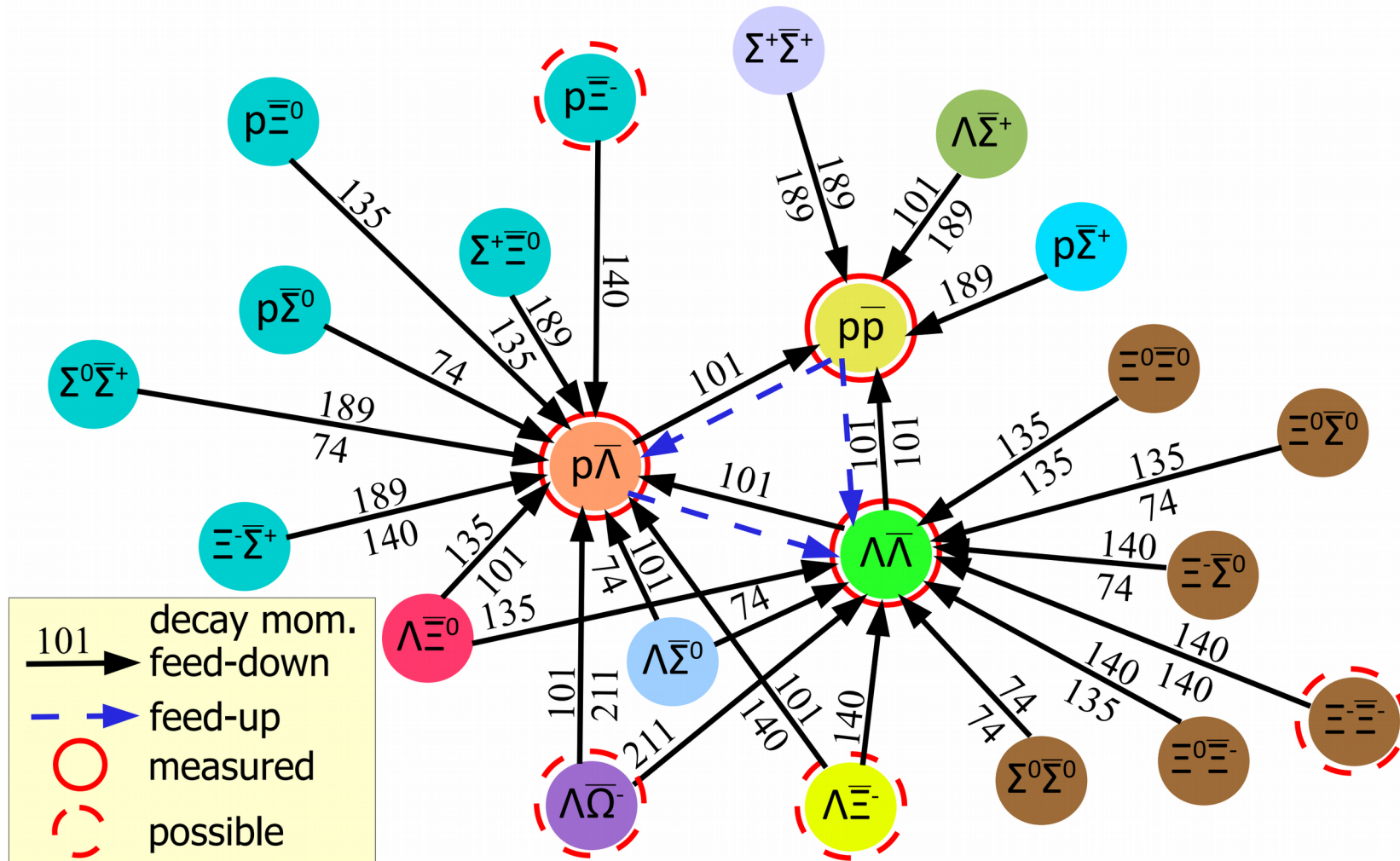


## Conclusions from fitting:

- Interaction parameters are measurable
- Scattering parameters for **all baryon-antibaryon pairs are similar to each other**
- We observe a **negative real part of scattering length** → repulsive strong interaction or creation of a bound state (existence of baryon-antibaryon bound states?)
- Significant **positive imaginary part of scattering length** – presence of a non-elastic channel – annihilation



# Residual correlations and other possibilities



# Motivation for $K^0_s K^\pm$ analysis

- Which sources of correlations are present in kaon systems?
  - Quantum Statistics (QS) – both  $K^0_s K^0_s$  and  $K^\pm K^\pm$
  - Coulomb FSI –  $K^\pm K^\pm$
  - **Strong FSI –  $K^0_s K^0_s$  (via  $f_0(980)/a_0(980)$  resonances)**
- Why are  $K^0_s K^\pm$  pairs interesting?
  - only Strong FSI:
    - $f_0(980)$  resonance is isospin = 0  $\rightarrow$  no  $f_0(980)$  strong interaction
    - $a_0(980)$  resonance is isospin = 1 as is the kaon pair  $\rightarrow$  only  $a_0(980)$  strong interaction present
- $a_0(980)$  resonance is a proposed tetraquark state (PRC 75 (2007) 045206)
- $a_0(980)$  mass and coupling par. (in GeV) from fits to  $\phi$  decay experiments

	$m_{a_0}$	$\Upsilon_{a_0 \rightarrow K\bar{K}}$	$\Upsilon_{a_0 \rightarrow \pi\eta}$	Reference
“Martin”	0.974	0.3330	0.2220	Nucl. Phys. B 121, 514 (1977)
“Antonelli”	0.985	0.4038	0.3711	arXiv: hep/ex-0209069 (2002)
“Achasov1”	0.992	0.5555	0.4401	Phys. Rev. D 68, 014006 (2003)
“Achasov2”	1.003	0.8365	0.4580	Phys. Rev. D 68, 014006 (2003)

# Motivation for $K^0_s K^\pm$ analysis

PHYSICAL REVIEW D 79, 074014 (2009)

## Global aspects of the scalar meson puzzle

Amir H. Fariborz,<sup>1,\*</sup> Renata Jora,<sup>2,†</sup> and Joseph Schechter<sup>3,‡</sup>

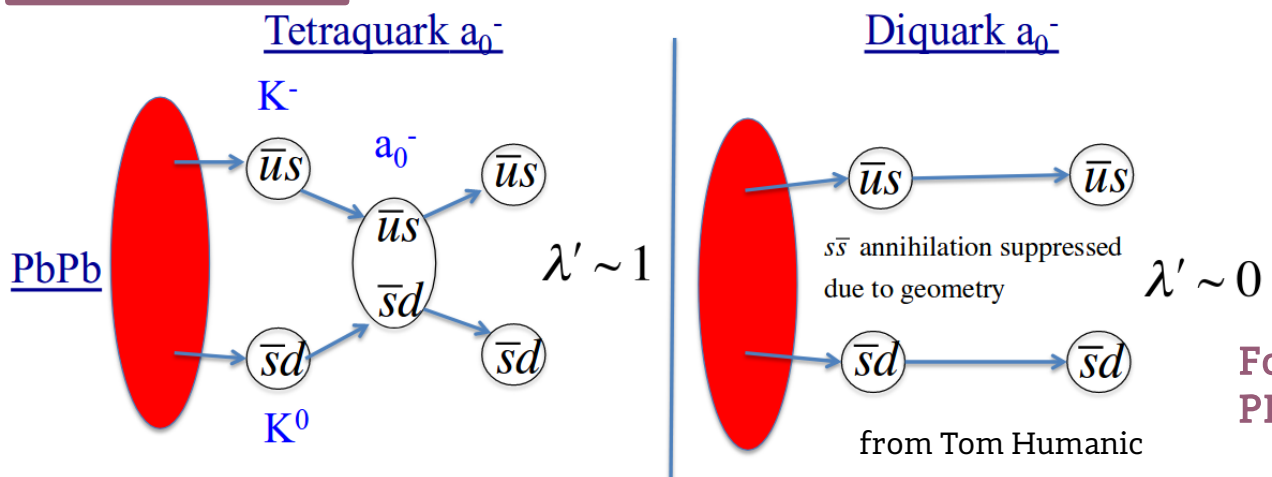
TABLE II.  $m_a$  and  $m_{a'}$  are inputs. Typical predicted properties of scalar states:  $\bar{q}q$  percentage (2nd column),  $\bar{q}\bar{q}qq$  (3rd column) and masses (last column).

State	$\bar{q}q\%$	$\bar{q}\bar{q}qq\%$	$m$ (GeV)
$a$	24	76	0.984
$a'$	76	24	1.474
$\kappa$	8	92	1.067

$a_0$  predicted to be 76% tetraquark  
PDG still lists it as a light meson

• How femtoscopy can help to determine which state it is?

$\lambda' \equiv \lambda_{K^0 K^-} / \lambda_{KK}$  for  $\bar{u}s\bar{s}d$  vs.  $\bar{u}d$   $a_0^-$  expected from geometry



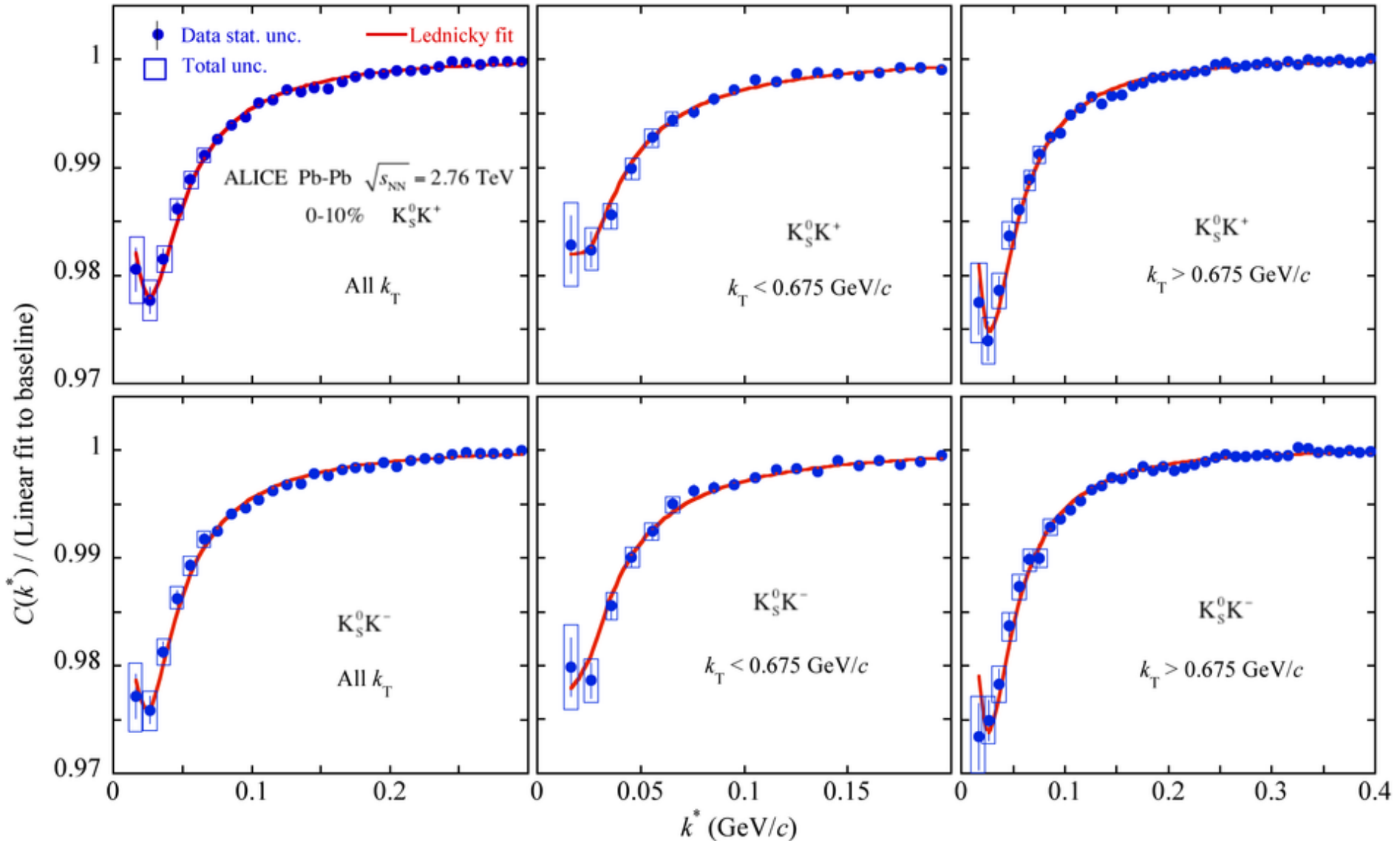
For details:  
PLB 774 (2017) 64-77



# Measured correlation functions

PLB 774 (2017) 64-77

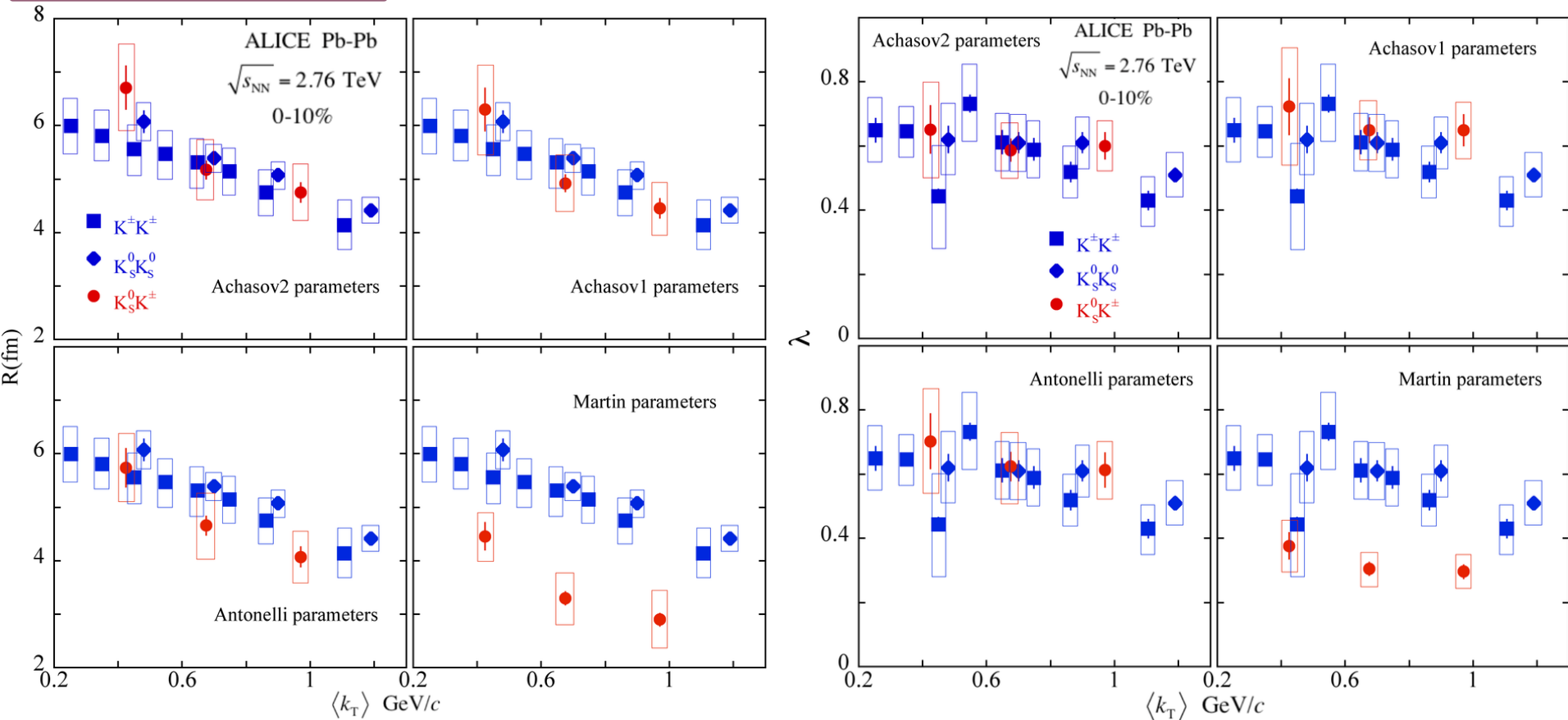
$C_{\text{raw}}(k^*) / (\text{linear fit})$



The  $a_0(980)$  final interaction gives an excellent fit to the data!



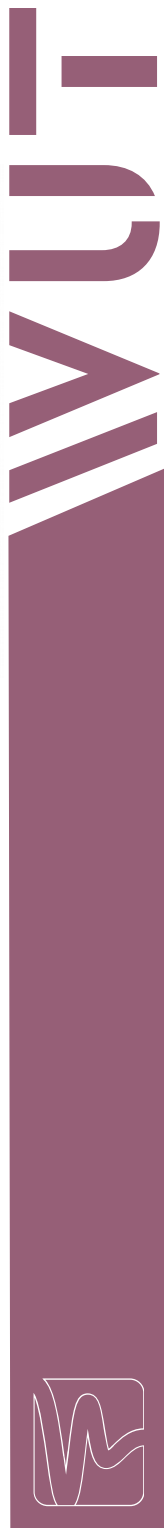
# Results of the fit



- “Achasov” parameter fits give best agreements with  $K_s^0K_s^0$  and  $K^{\pm}K^{\pm}$
- “Antonelli” parameter fits are somewhat lower
- “Martin” parameter fits much lower
- Present results favor higher  $a_0(980)$  parameters  $\rightarrow a_0(1000)$

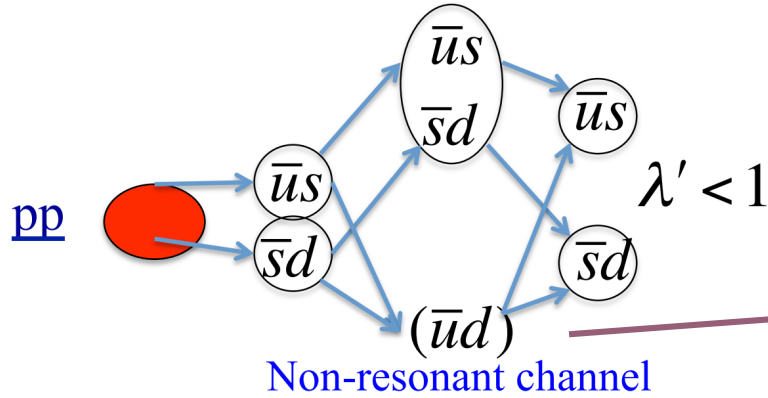


# Can pp collisions tell us more?



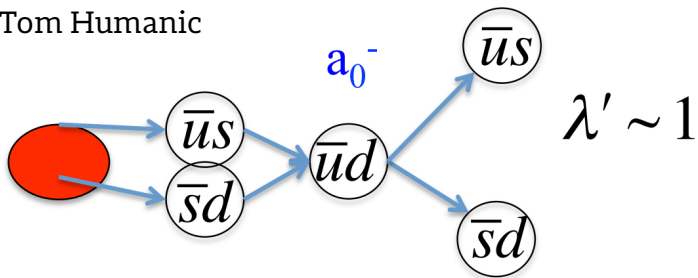
PLB 790 (2019) 22

Tetraquark  $a_0^-$

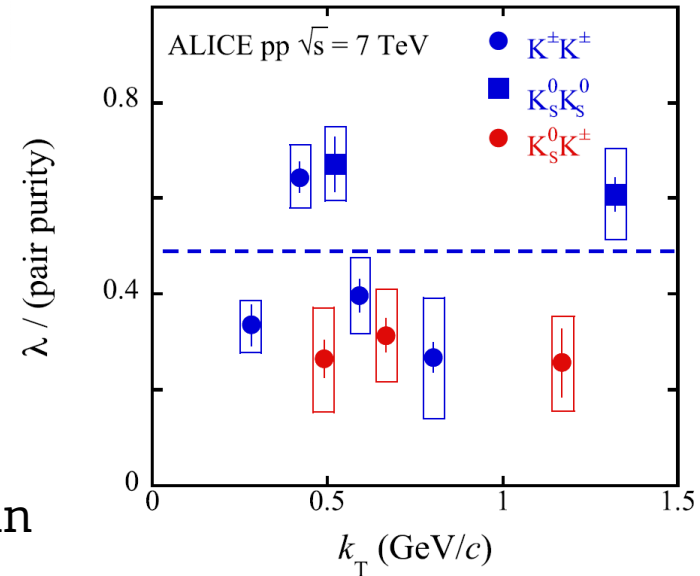
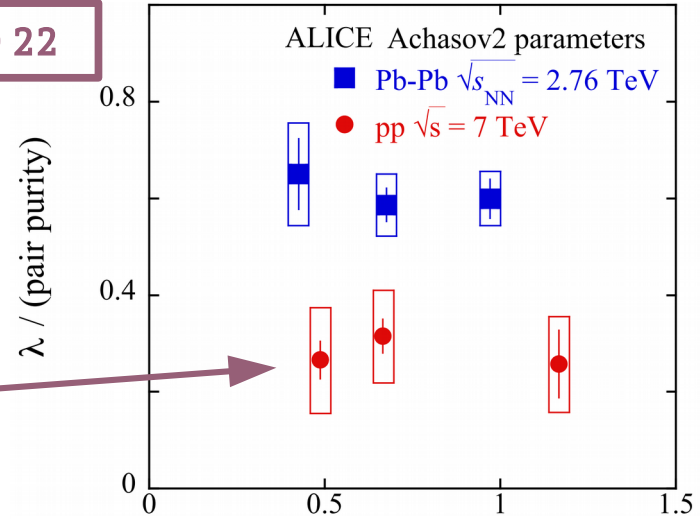


Diquark  $a_0^-$

from Tom Humanic



- $\lambda$  parameters from pp are significantly lower than Pb-Pb (non-resonant channel also present)
- They also tend to be lower than for identical-kaon analysis
- **Results from both pp and Pb-Pb favor  $a_0$  to be a tetraquark**



# Status of femtoscopy in ALICE

- Two-pion Bose-Einstein correlations in pp collisions at  $\sqrt{s}=900$  GeV, Phys. Rev. D 82 (2010) 052001
- Two-pion Bose-Einstein correlations in central Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Lett. B 696 (2011) 328-337
- Femtoscopy in pp a 0.9 and 7 TeV, Phys. Rev. D 84 (2011) 112004
- K0sK0s correlations in pp collisions at  $\sqrt{s}=7$  TeV from the LHC ALICE experiment, Phys. Lett. B 717 (2012) 151-161
- Charged kaon femtosopic correlations in pp collisions at  $\sqrt{s}=7$  TeV, Phys. Rev. D 87 (2013) 052016
- Two and Three-Pion Quantum Statistics Correlations in Pb-Pb Collisions at  $\sqrt{s_{NN}}=2.76$  TeV at the LHC Phys. Rev. C 89 (2014) 024911
- Freeze-out radii extracted from three-pion cumulants in pp, p-Pb and Pb-Pb collisions at the LHC, Phys. Lett. B 739 (2014) 139-151
- Pion femtoscopy in p-Pb Phys. Rev. C 91 (2015) 034906
- Multipion Bose-Einstein correlations in pp, p-Pb, and Pb-Pb collisions at the LHC, Phys. Rev. C 93 (2016) 054908
- Centrality dependence of pion freeze-out radii in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV Phys. Rev. C 93 (2016) 024905
- 1D pion, kaon , proton femtoscopy in Pb-Pb Phys. Rev. C 92 (2015) 054908
- Azimuthally differential pion femtoscopy in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Rev. Lett. 118 (2017) 222301
- Kaon femtoscopy in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Rev. C96 (2017) 064613
- Measuring K0sK± interactions using Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Lett. B 774 (2017) 64
- Azimuthally-differential pion femtoscopy relative to the third harmonic event plane in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Lett. B 785 (2018) 320-331
- Measuring K0sK± interactions using pp collisions at  $\sqrt{s}=7$  TeV, Phys. Lett. B 790 (2019) 22
- p-p, p- $\Lambda$  and  $\Lambda$ - $\Lambda$  correlations studied via femtoscopy in pp reactions at  $\sqrt{s}=7$  TeV, Phys. Rev. C 99 (2019) 024001
- Event-shape and multiplicity dependence of freeze-out radii in pp collisions at  $\sqrt{s}=7$  TeV, arXiv:1901.05518



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- Measuring K0sK± interactions using pp collisions at  $\sqrt{s}=7$  TeV, Phys. Lett. B 790 (2019) 28
- p-p, p- $\Lambda$  and  $\Lambda$ - $\Lambda$  correlations studied via femtoscopy in pp reactions at  $\sqrt{s}=7$  TeV, Phys. Rev. C 99 (2019) 024908
- Event-shape and multiplicity dependence of freeze-out radii in pp collisions at  $\sqrt{s}=7$  TeV, arXiv:1901.05518

**18 papers in total**  
**~7.5% of all ALICE papers**  
**(241, 3 March 2019)**  
**3 more coming soon**





# Summary

- **Femtoscscopy in ALICE has produced a rich set of results**
  - Remarkable scaling of radii from various experiments is kept also at LHC energies
- **More differential and complex analysis are being carried out right now**
  - Kaon-kaon and pion-kaon femtoscopy suggest that **kaons are emitted later than pions**. Data show **2.1 fm/c delay** between pion and kaon average emission time
  - Pion source at the freeze-out is **elongated in the out-of-plane direction**. Initial state **triangularity is washed-out at freeze out**
  - In small systems – sphericity allows for better handle of underlying non-femtoscopic correlations
- **Femtoscscopy has also become a powerful tool for hadron physics**
  - Unique experimental environment at LHC → **“matter-antimatter pair factory”**
  - Possibility to study exotic states such as **tetraquarks**
- **More new and exciting analyses are on the way – stay tuned!**





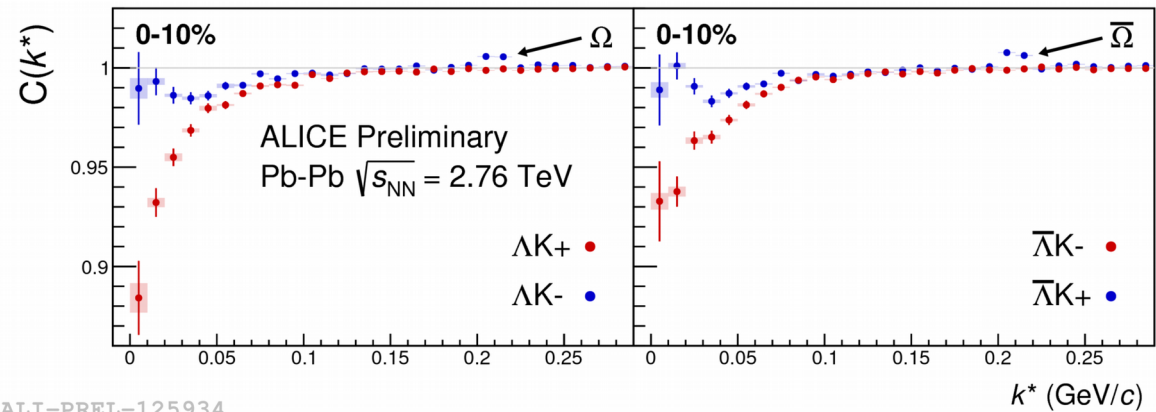
**Thank you!**

# Backup



# Other interesting pairs

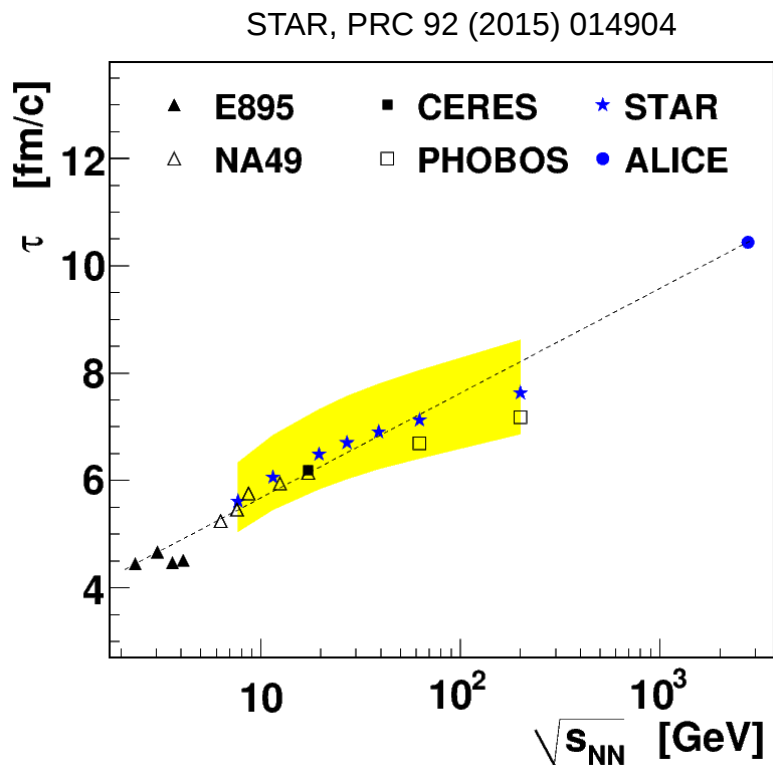
- Many other interesting correlations are currently studied by ALICE
- Lambda-kaon (both charged and neutral) pairs
  - scattering parameters measured for the first time
- $\Lambda K^+$  shows greater suppression at low  $k^*$  compared to:  $\Lambda K^-$ :
  - effect arising from  $s\bar{s}$  annihilation compared to  $u\bar{u}$ ?
  - or  $S=0$   $\Lambda K^+$  system has more interaction channels than  $S=-2$   $\Lambda K^-$ ?



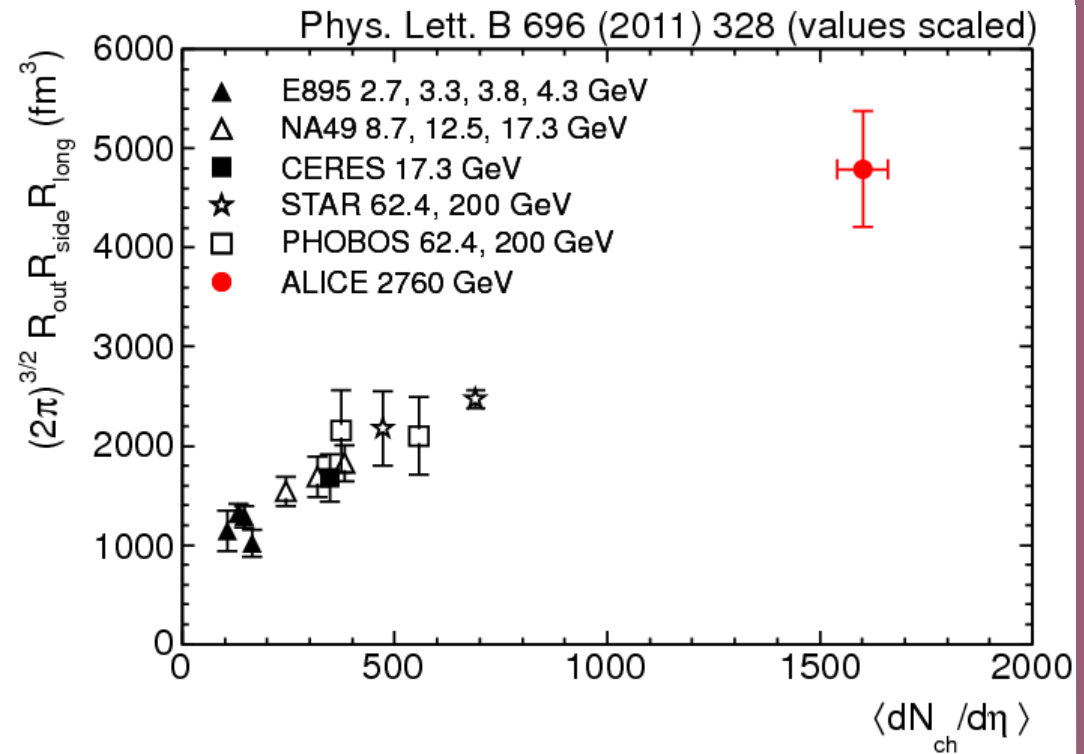
# Pion femtoscopy

- Lifetime can be estimated from the longitudinal radius
- Clear increase of the system volume and lifetime with collision energy, at LHC system **twice as large** and **living 30% longer** than at top RHIC energy (good conditions for QGP studies)

Lifetime



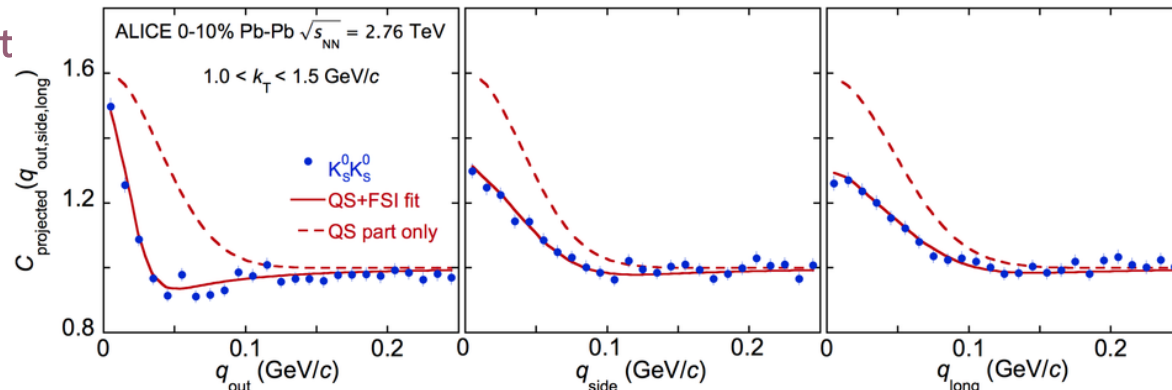
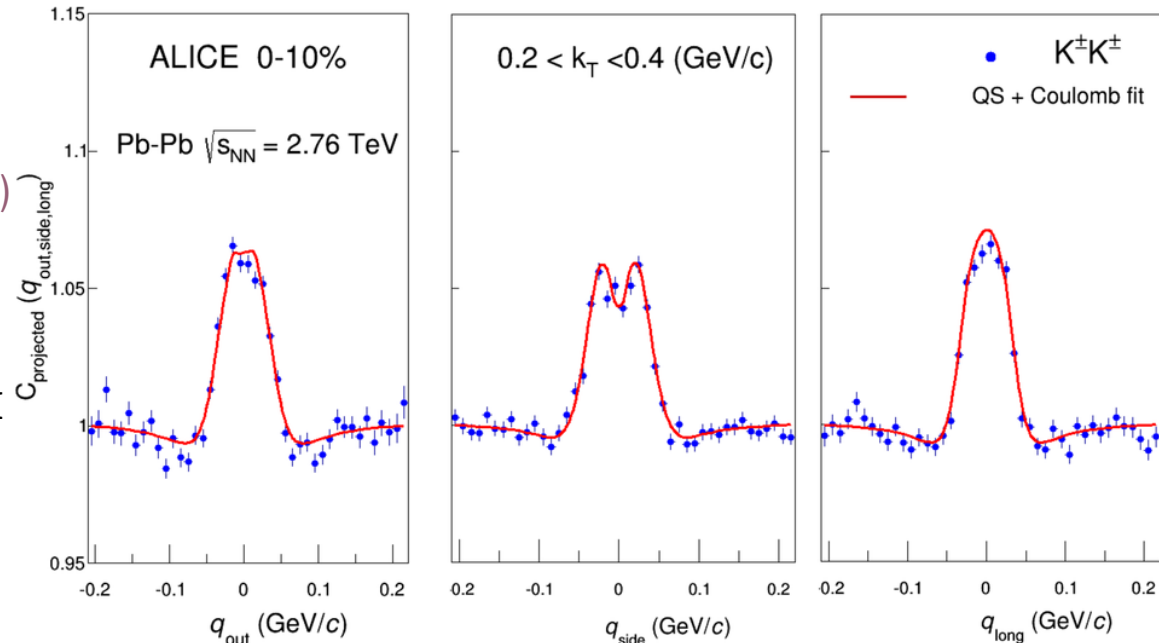
Volume



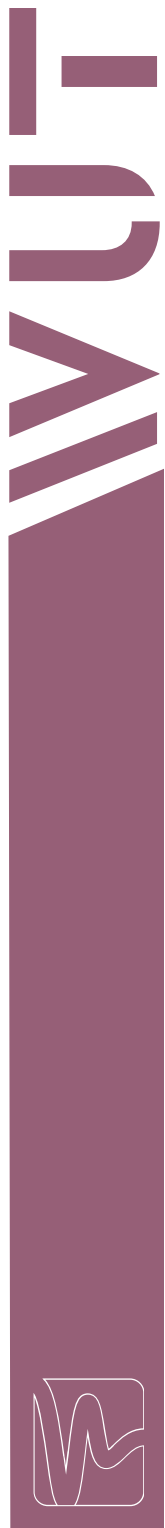
# Kaon femtoscopy

Phys. Rev. C96 (2017) 064613

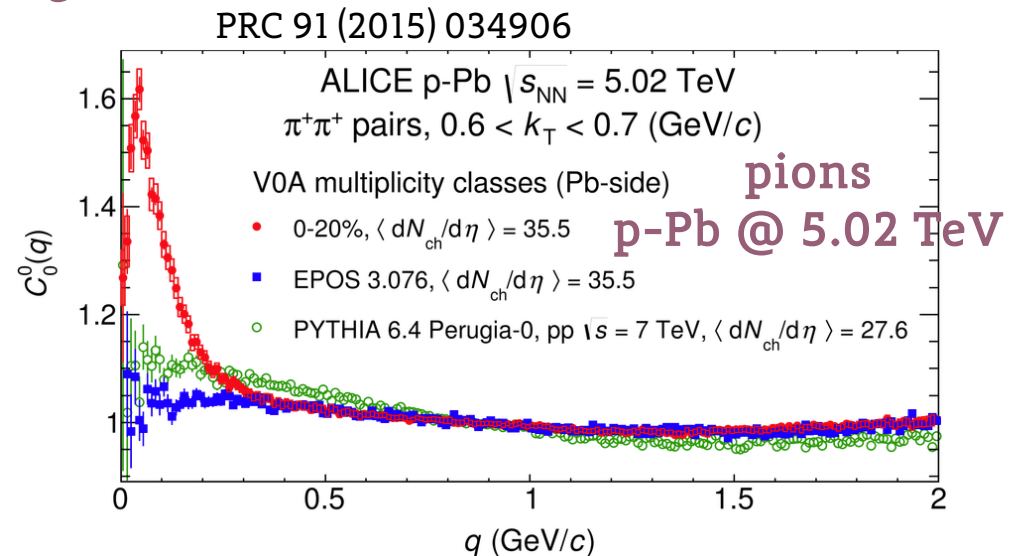
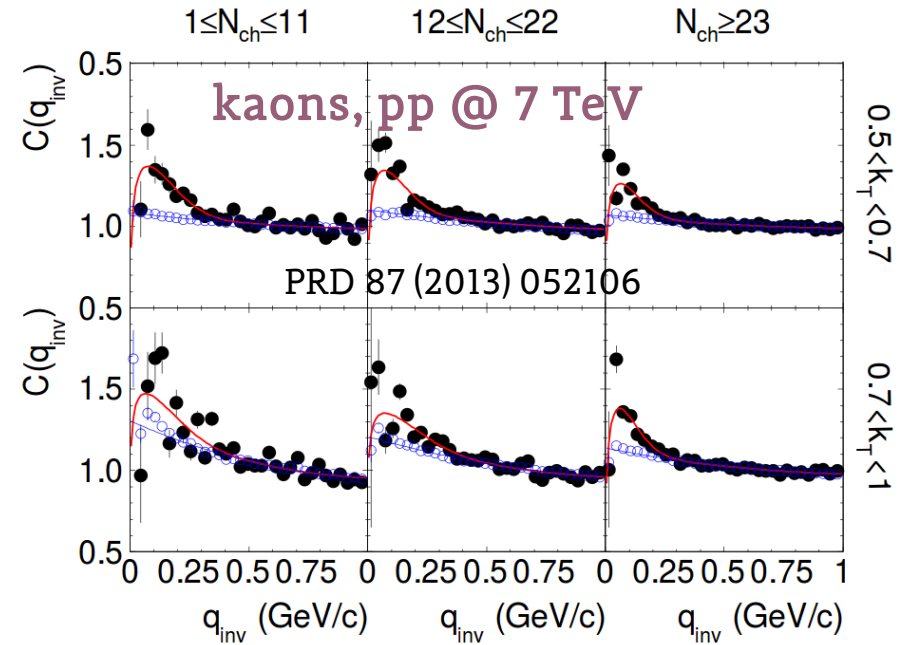
- **Cleaner signal** compared to pions (less affected by resonances)
- Studying **neutral and charged kaons** together provides a convenient **consistency check** (different experimental techniques)
  - Charged Kaons: QS + Strong and Coulomb FSI
  - Neutral Kaons: QS + Strong FSI (including resonances)
- Models which describe pions well, should describe kaons with **equal precision**
- **Rescattering phase has different influence on pions and kaons:**
  - can lead to broken  $m_T$ -scaling
  - good probe of the rescattering phase effects



# Non-femtoscopic correlations



- Non-femtoscopic correlations visible in small systems for pions and kaons:
  - Grow with increasing  $k_T$
  - Grow with decreasing multiplicity
  - Significant source of systematics in the fitting procedure
- So far hypothesis of (mini-)jet origin
- How can we test it?



# Beyond the system size

$$C(q) = \int S(r) |\Psi(q, r)|^2 d^4 r$$

$q = 2 \cdot k^* = p_1 - p_2$

measured correlation      emission function  
(source size/shape)      pair wave function  
(includes cross section)

- Pair wave function  $\Psi$  can be parametrized with **scattering length**  $f_0$ , and **effective radius**  $d_0$  parameters.

The correlation function is characterized by **three parameters**:

- **radius**  $R$ , **scattering length**  $f_0$ , and **effective radius**  $d_0$
- **cross section**  $\sigma$  (at low  $k^*$ ) is simply:  $\sigma = 4 \pi |f|^2$



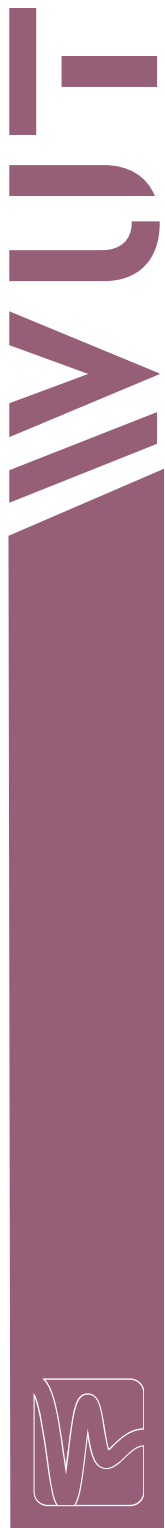


# Potential applications

- **Input to models with re-scattering phase (eg. UrQMD):**
  - annihilation cross sections only measured for  $p\bar{p}$ ,  $\bar{p}n$ , and  $\bar{p}d$  pairs – UrQMD currently **guesses it for other systems** from  $pp$  pairs
- **Structure of baryons/search for CPT violation**  
STAR, Nature 527, 345-348 (2015)
- **Search for H-dibaryon**  
ALICE, PLB 752 (2016) 267-277
- **Hypernuclear structure theory**  
Nucl.Phys. A914 (2013) 377-386
- **Neutron star equation of state**  
Nucl.Phys. A804 (2008) 309-321
- **Relativistic heavy-ion collisions at LHC or RHIC produce very similar number of baryons and antibaryons, “matter-antimatter pair factories”**



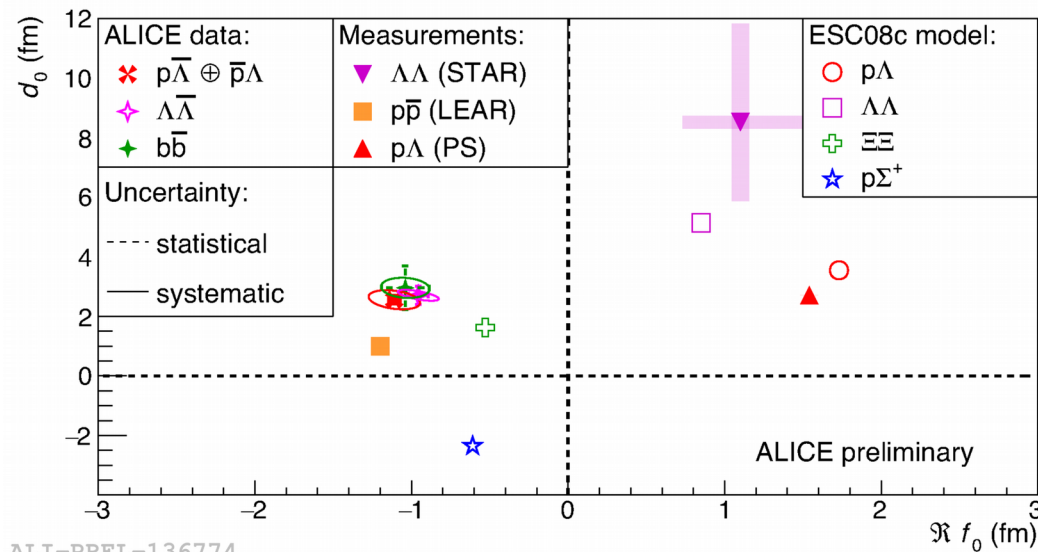
# Baryon-antibaryon correlations



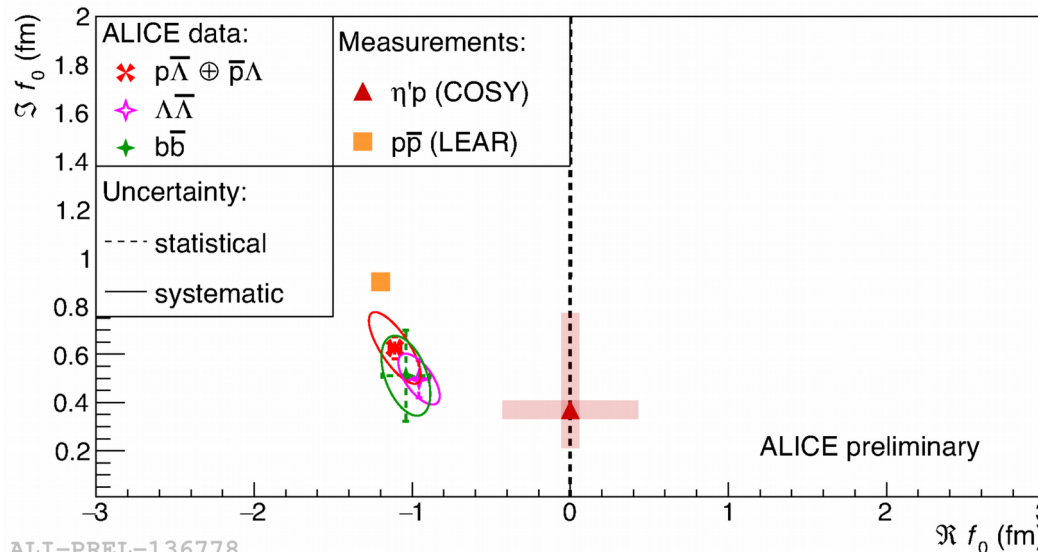
## Conclusions from fitting:

- Interaction parameters are measurable
- Scattering parameters for **all baryon-antibaryon pairs are similar to each other** (UrQMD assumption is valid)
- We observe a **negative real part of scattering length** → repulsive strong interaction or creation of a bound state (existence of baryon-antibaryon bound states?)
- Significant **positive imaginary part of scattering length** – presence of a non-elastic channel – annihilation

**Next steps:** try to look for baryon-antibaryon bound states



ALI-PREL-136774



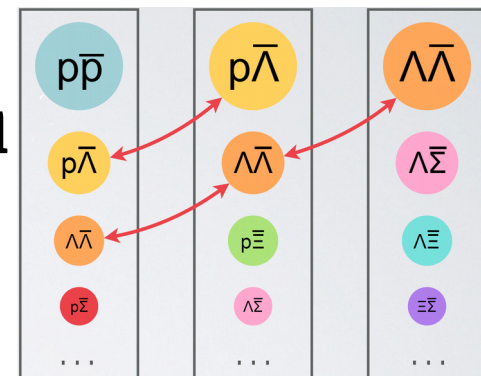
ALI-PREL-136778

Jeremi Niedziela, PhD thesis

# Baryon-antibaryon correlations

## Explanation of the fitting procedure:

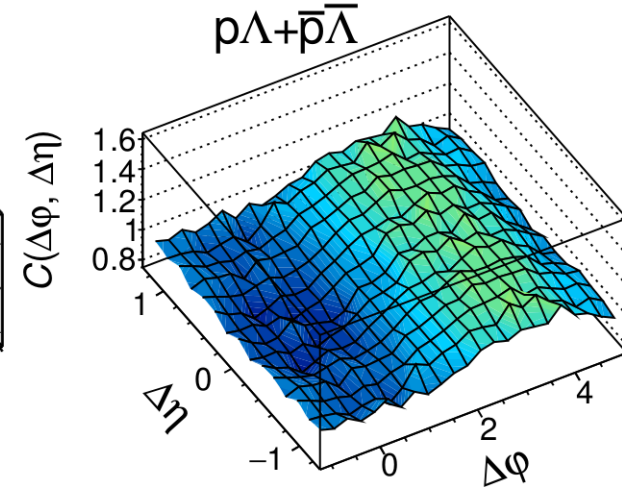
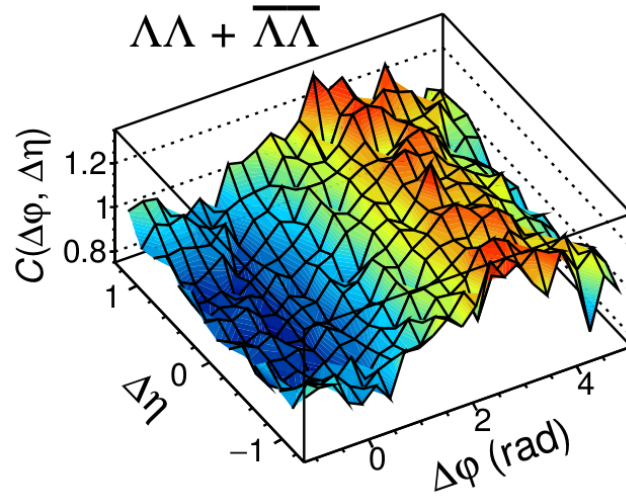
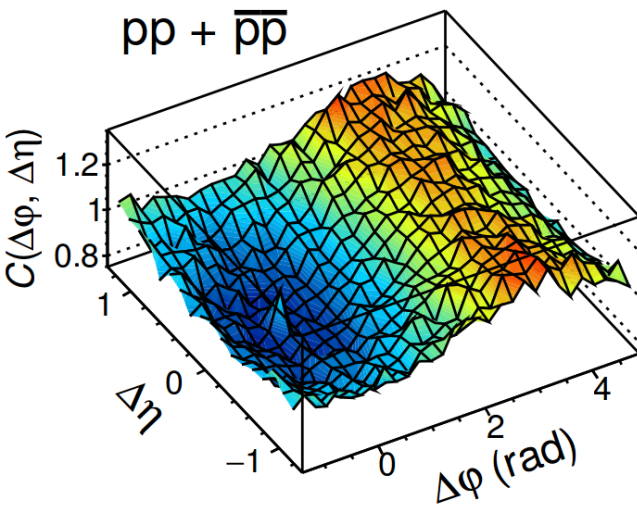
- $\chi^2$  is calculated from a “global” fit to all functions:  
2 data sets, 3 pair combinations, 6 centrality bins (**total 36 functions**)
- simultaneous fit accounts for parameters **shared** between different systems (such as  $\Lambda\bar{\Lambda}$  scattering length)
- **radii scale with multiplicity** for a given system  $R_{inv} = a \cdot \sqrt[3]{N_{ch}} + b$
- for different systems we assume **radii scaling with  $m_T$**
- Fractions of **residual pairs** taken from AMPT



# $\Delta\eta\Delta\phi$ of baryons

Eur.Phys.J. C77 (2017) 8, 569

Małgorzata Janik, ŁG



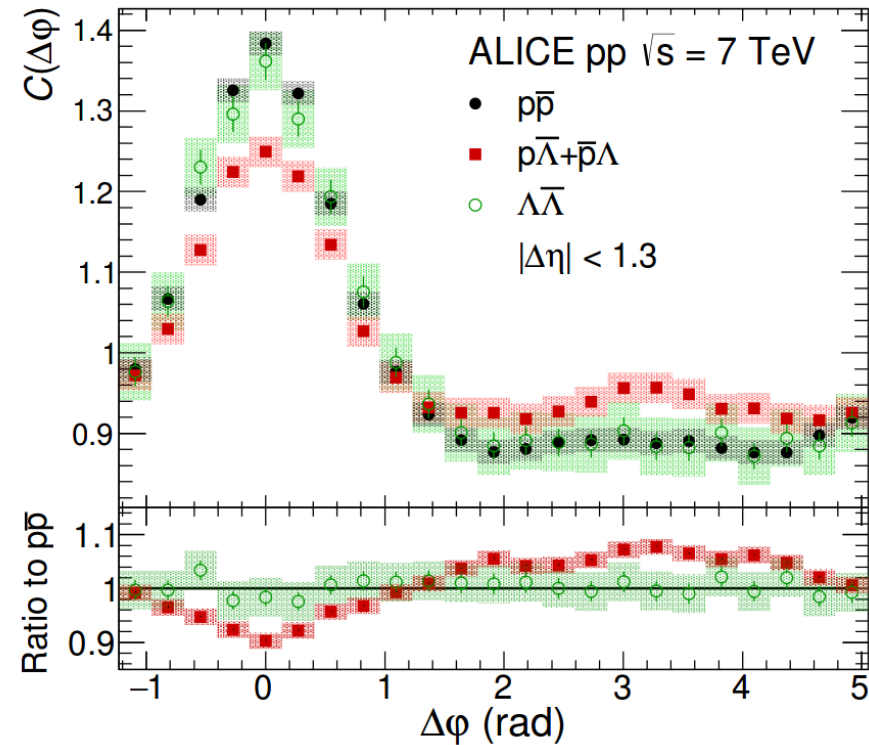
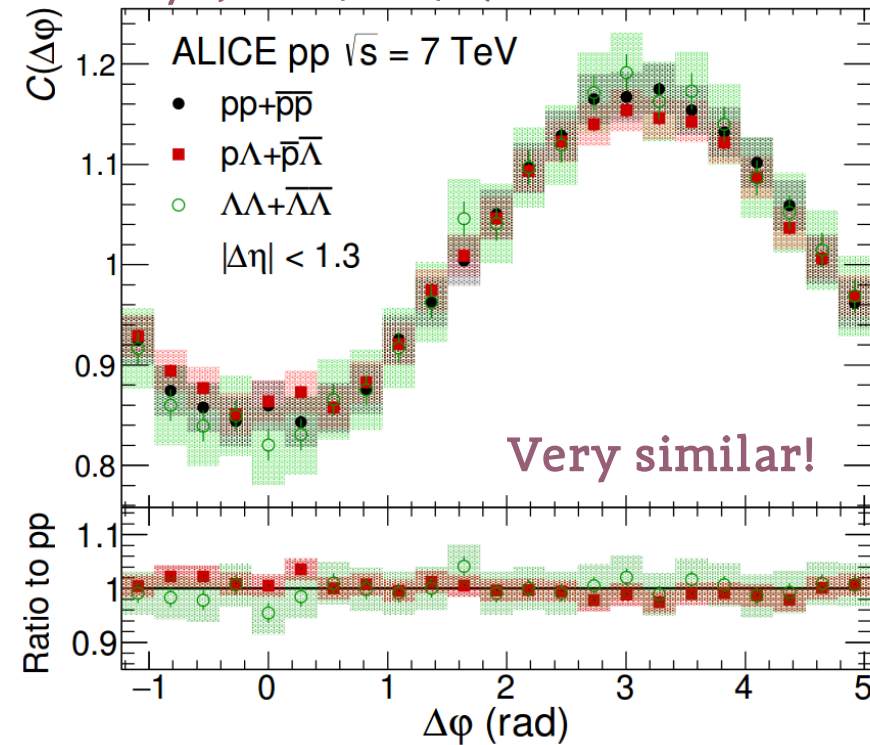
- We found that all baryon-baryon pairs show a **depression** instead of a typical near-side peak



# $\Delta\eta\Delta\phi$ of baryons

Małgorzata Janik, ŁG

Eur.Phys.J. C77 (2017) 8, 569



- Projections show how similar are baryon-baryons pairs to each other
- Similarity between pairs, to a lesser extent, is also observed in the baryon-antibaryon case

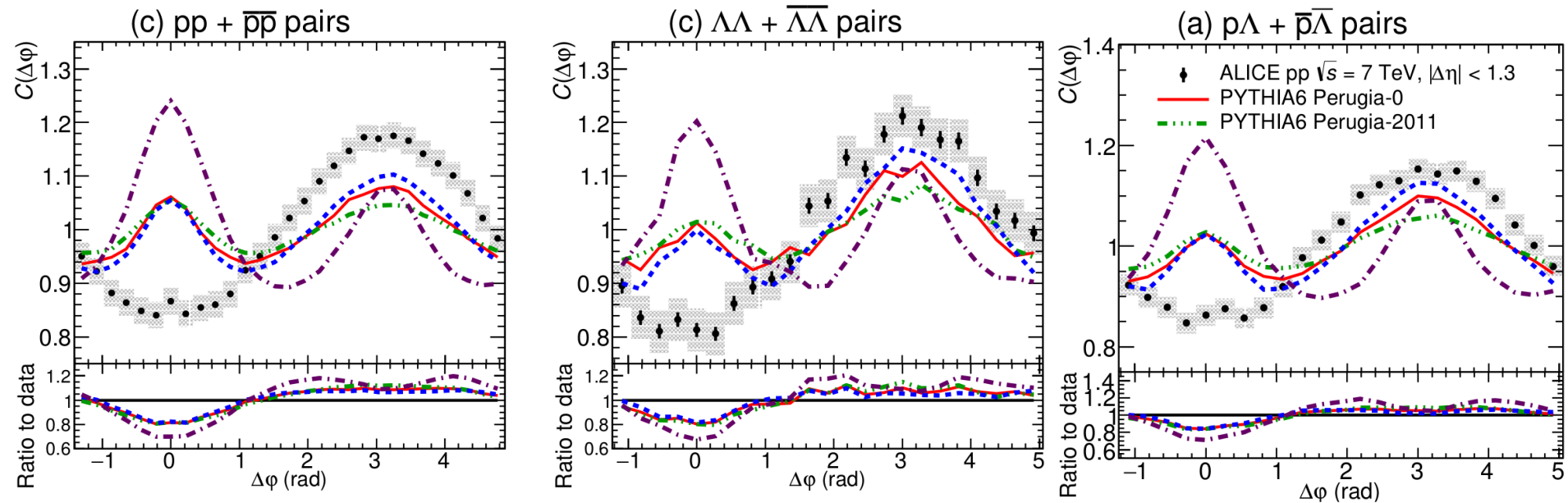
## Possible explanations:

- Fermi-Dirac Quantum Statistics? **NO (non-identical particles)**
- Coulomb repulsion? **NO (uncharged particles)**
- **Strong Final-State Interactions?** (see next slides)



# Are baryons interesting?

Eur.Phys.J. C77 (2017) 8, 569



- **Azimuthal correlations: depletion** instead of a near-side peak visible for **baryon-baryon pairs** in pp collisions @ 7 TV
- **New baryon production mechanism needed** (T. Sjöstrand, QM 2018)
- It seems we do not understand some fundamental properties of baryons

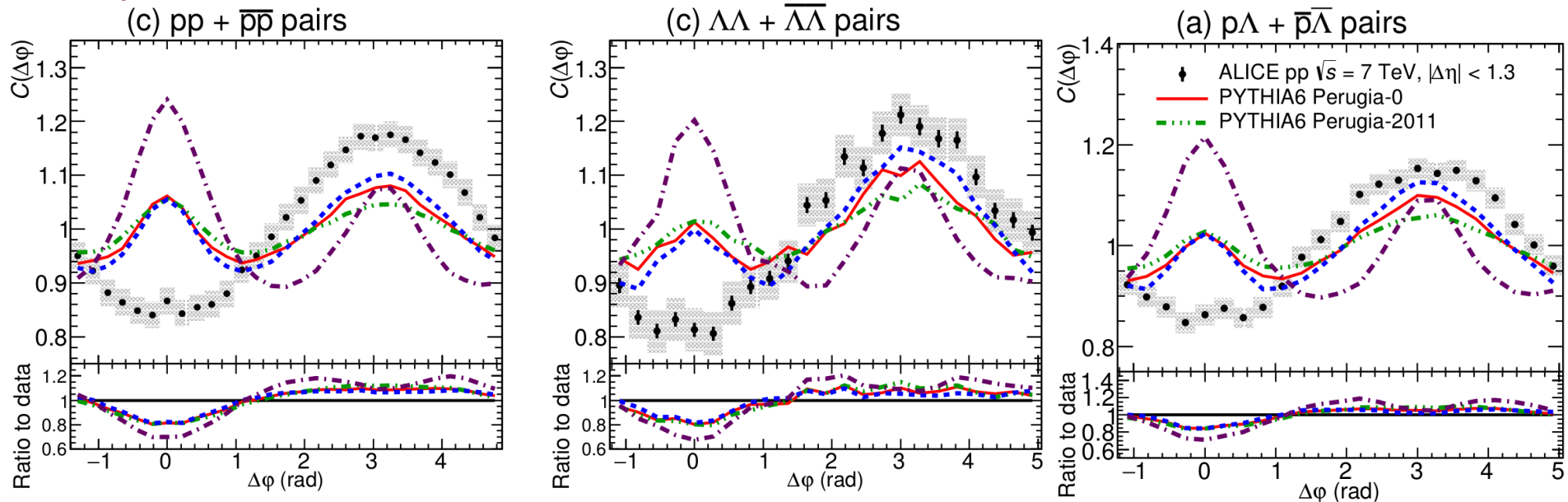
For details see CERN-LHC seminar by  
Małgorzata Janik  
30 May 2017  
<https://indico.cern.ch/event/632396/>



# $\Delta\eta\Delta\phi$ of baryons

Małgorzata Janik, ŁG

Eur.Phys.J. C77 (2017) 8, 569



- None of studied MC models (PYTHIA, PHOJET, EPOS, HERWIG) agrees with the data even qualitatively
- What can be the explanation of this effect?

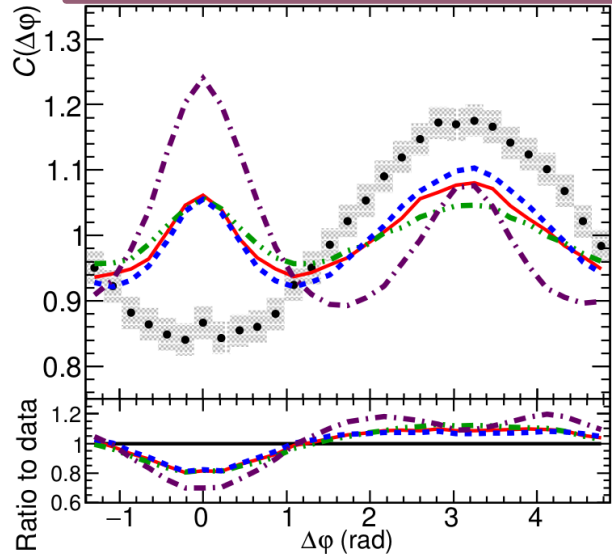
Let's look at similar studies in  $e^+e^-$  collisions at  $\sqrt{s} = 29 \text{ GeV}$  (SLAC-PEP) from late 80's



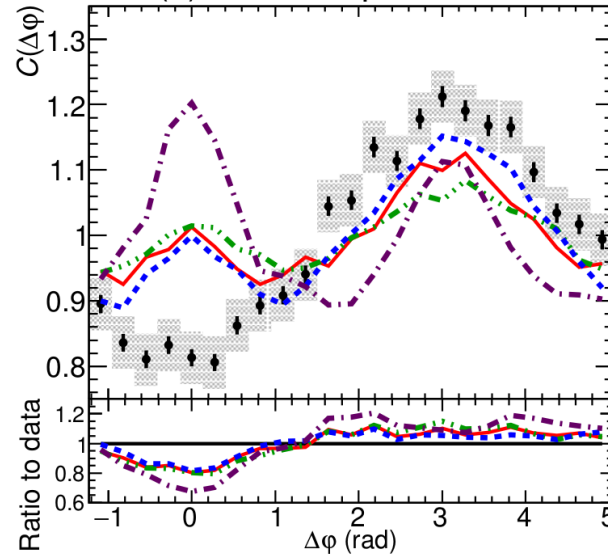
# $\Delta\eta\Delta\phi$ of baryons

Eur.Phys.J. C77 (2017) 8,

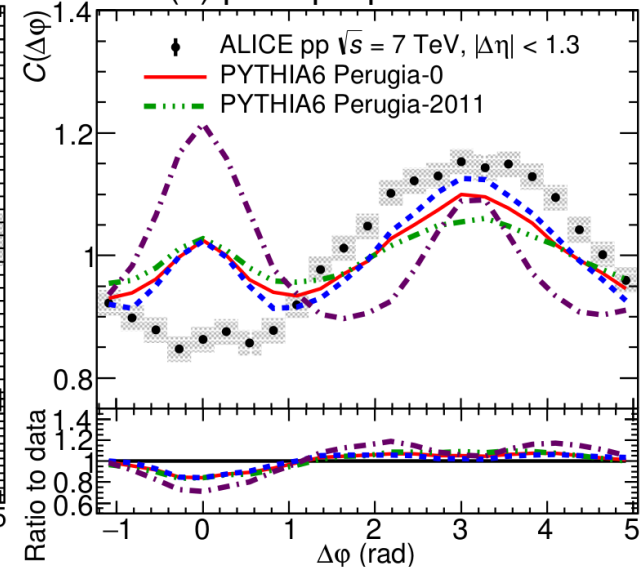
(c)  $pp + \bar{p}\bar{p}$  pairs <sup>569</sup>



(c)  $\Lambda\Lambda + \bar{\Lambda}\bar{\Lambda}$  pairs



(a)  $p\Lambda + \bar{p}\bar{\Lambda}$  pairs



- **Depletion** instead of a near-side peak visible for **baryon-baryon pairs** in azimuthal correlations
- None of studied Monte Carlo models (PYTHIA, PHOJET, EPOS, HERWIG) agrees with the data even qualitatively
- Results challenge current hadronization models
- Possible explanation → AMPT with string melting and new coalescence model describes qualitatively our data





# $\Delta\eta\Delta\phi$ of baryons

arXiv:  
1808.10641  
31 Aug 2018

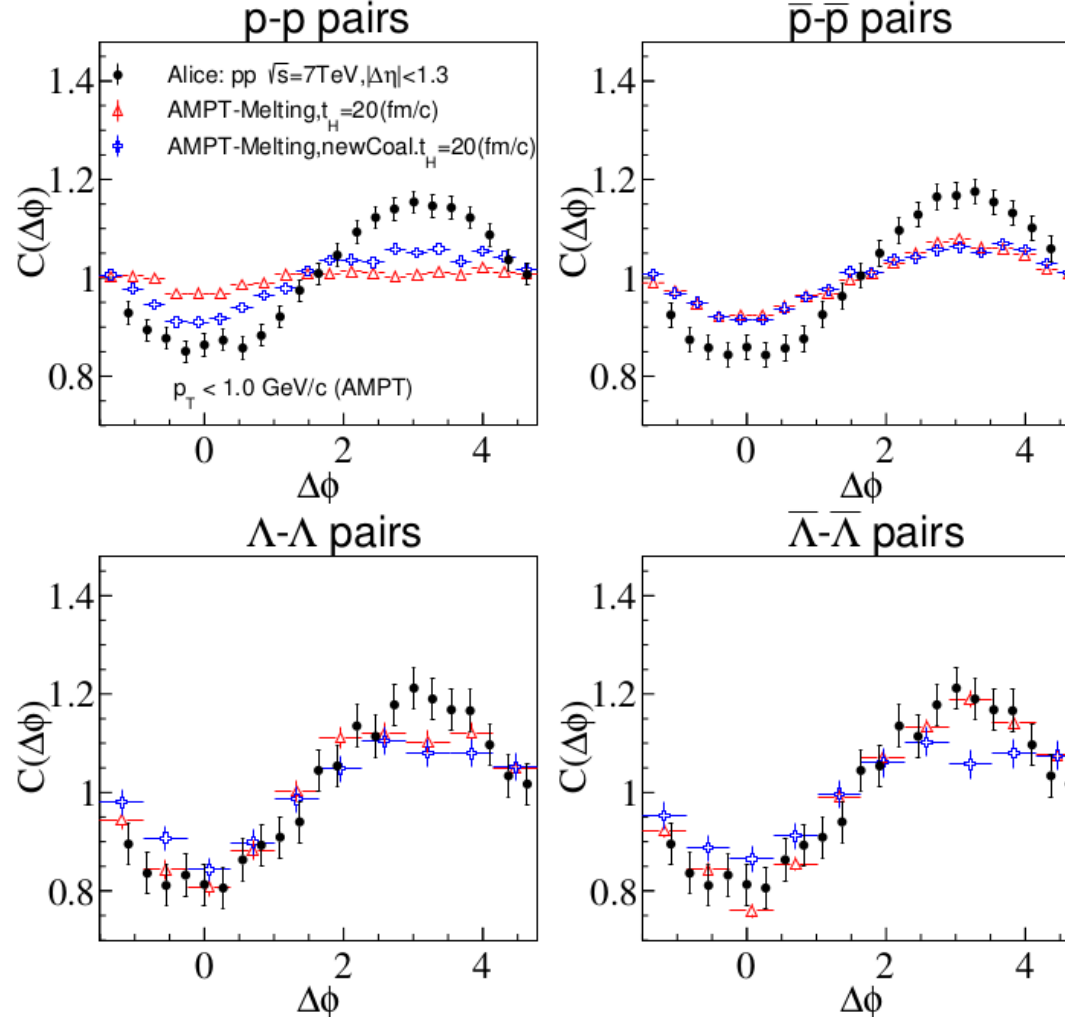


FIG. 6: One-dimensional  $\Delta\phi$  correlation functions of p-p,  $\bar{p}-\bar{p}$ ,  $\Lambda-\Lambda$  and  $\bar{\Lambda}-\bar{\Lambda}$  for  $p_T < 1.0 \text{ GeV}/c$  from the AMPT model. Open symbols represent AMPT calculations with different configurations as illustrated in the figure. Solid points are experiment data [14].

- Possible explanation → AMPT with string melting and new coalescence model describes qualitatively ALICE data



- Conventional pp generators successful, with MPI + CR generating some collectivity, but now cracks.
- Need new framework for baryon production.
- String close-packing likely to influence hadronization, before (shoving), during (ropes) and after (rescattering).
- Currently no known unique solution, so free to explore.
- Several recent & ongoing studies look promising, but much work and few active with pp generator outlook.
- Further experimental input crucial!

**Whole new field of study opening up!**

# Non-femtoscopic correlations



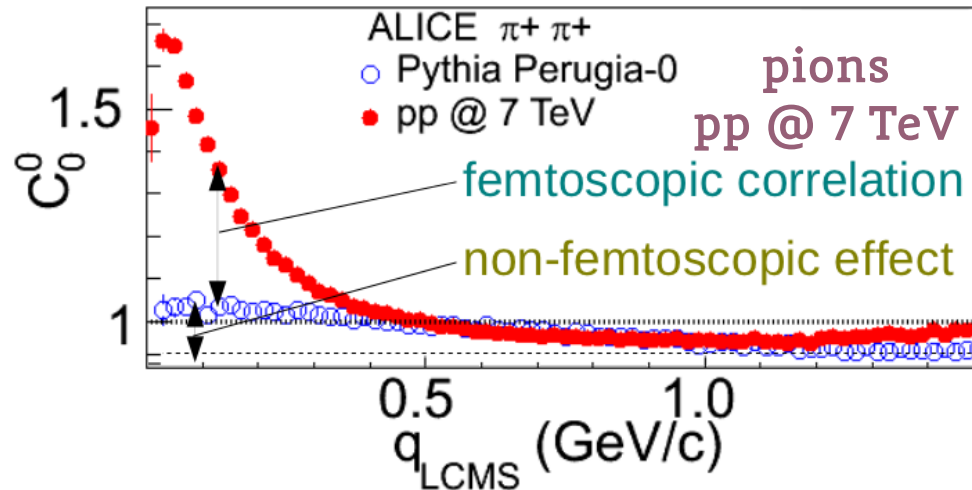
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- Significant source of systematics in the fitting procedure

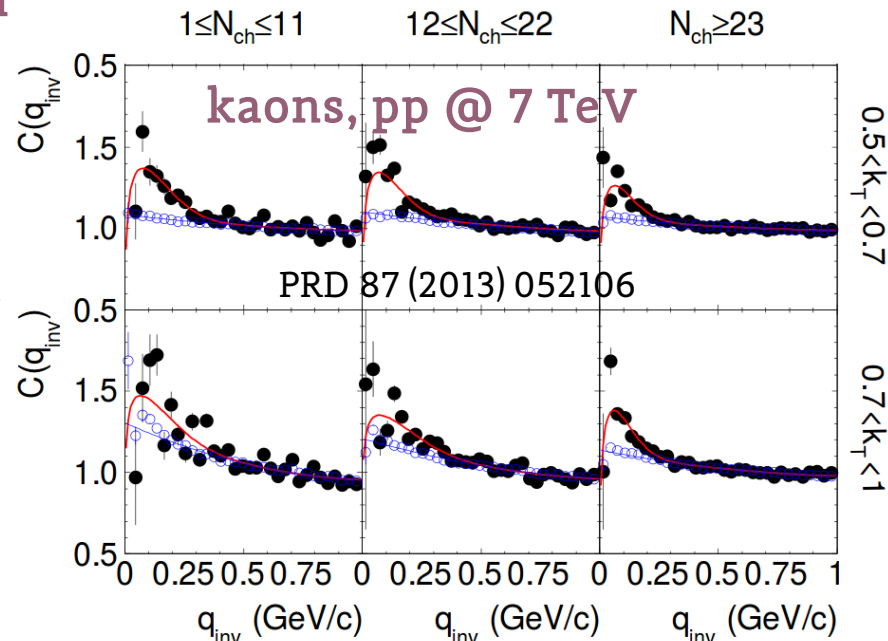
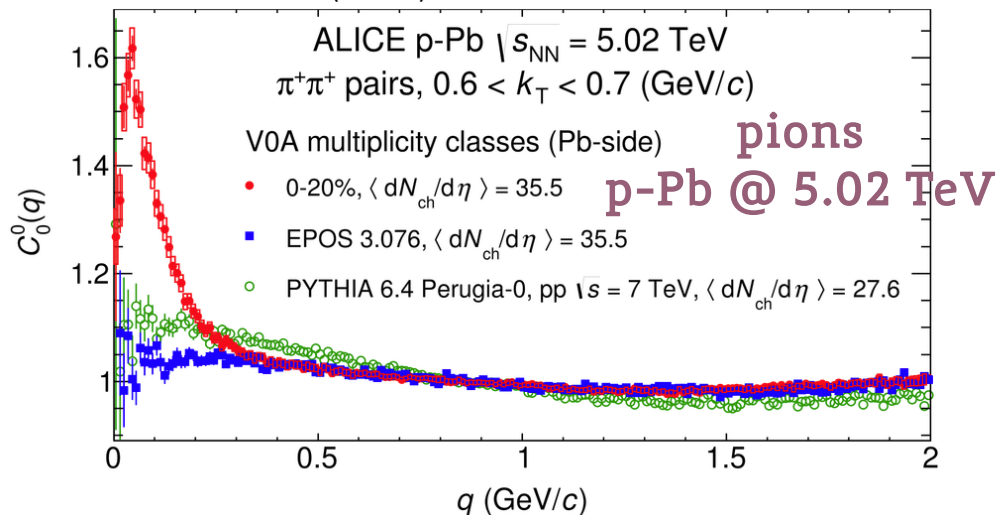
- So far only hypothesis of (mini-)jet origin

- How do baryon correlations look like in pp?

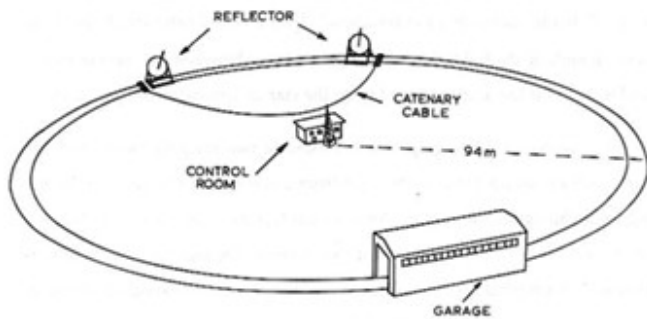
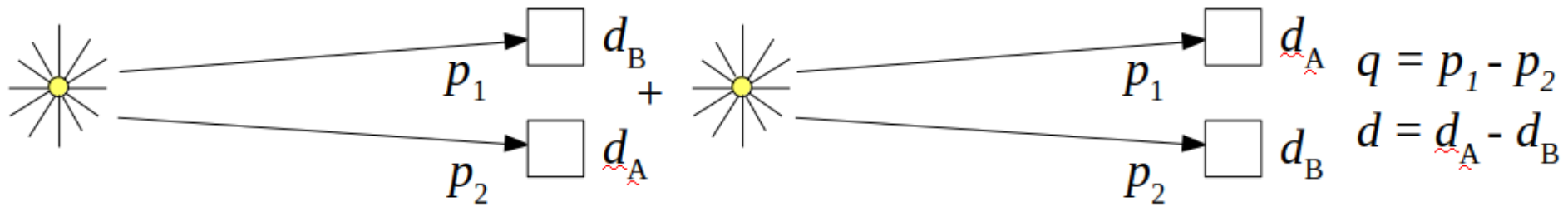
PRD 84 (2011) 112004



PRC 91 (2015) 034906



# Long history of HBT



- In astronomy angular size of the star is measured via photon correlation vs. spatial separation of detectors
- The momentum spread can be inferred, which is transformed into angular size of the star
- The mathematical formalism is similar
- The first measurement was done by Hanbury-Brown and Twiss – HBT !

Figure 1. Aerial photo and illustration of the original HBT apparatus. They have been extracted from Ref.[1].



# Are baryons interesting?

LETTER

BASE experiment

OPEN

doi:10.1038/nature14861

nature

Search for potential CPT symmetry breaking



## High-precision comparison of the antiproton-to-proton charge-to-mass ratio

S. Ulmer<sup>1</sup>,  
Y. Matsud

LETTER

STAR

doi:10.1038/nature15724

Invariance formation model of fundamental are identical invariance to be invariant although it and Lorentz pendulum only a few fundamental we report a few antiproton pair correlations among data collected by the STAR experimenter to that for at the Relativistic Heavy Ion Collider (RHIC)<sup>3</sup>, where gold ions are collided with a centre-of-mass energy of 200 GeV per nucleon pair. Antiprotons are abundantly produced in such collisions, thus making it feasible to study details of the antiproton-antiproton interaction. By applying a technique similar to Hanbury Brown and Twiss intensity interferometry<sup>4</sup>, we show that the force between two antiprotons is attractive. In addition, we report two key parameters that characterize the corresponding strong interaction: the scattering length and the effective range of the interaction. Our measured parameters are consistent within error with the corresponding values for proton-proton interactions. Our results provide direct information on the interaction between two antiprotons, one of the simplest systems of antinucleons, and are fundamental to understanding the structure of more-compl

## Measurement of interaction between antiprotons

The STAR Collaboration\*

One of the primary goals of nuclear physics is to understand the force between nucleons, which is a necessary step for understanding the structure of nuclei and how nuclei interact with each other. Rutherford discovered the atomic nucleus in 1911, and the large body of knowledge about the nuclear force that has since been acquired was derived from studies made on nucleons or nuclei. Although antinuclei up to antihelium-4 have been discovered and their masses measured, little is known directly about the nuclear force between antinucleons. Here, we study antiproton-antiproton pair correlations among data collected by the STAR experimenter to that for at the Relativistic Heavy Ion Collider (RHIC)<sup>3</sup>, where gold ions are collided with a centre-of-mass energy of 200 GeV per nucleon pair. Antiprotons are abundantly produced in such collisions, thus making it feasible to study details of the antiproton-antiproton interaction. By applying a technique similar to Hanbury Brown and Twiss intensity interferometry<sup>4</sup>, we show that the force between two antiprotons is attractive. In addition, we report two key parameters that characterize the corresponding strong interaction: the scattering length and the effective range of the interaction. Our measured parameters are consistent within error with the corresponding values for proton-proton interactions. Our results provide direct information on the interaction between two antiprotons, one of the simplest systems of antinucleons, and are fundamental to understanding the structure of more-compl

nature physics

ALICE

LETTERS

PUBLISHED ONLINE: 17 AUGUST 2015 | DOI: 10.1038/NPHYS3432

OPEN

## Precision measurement of the mass difference between light nuclei and anti-nuclei

ALICE Collaboration<sup>†</sup>

The measurement of the mass differences for systems bound by the strong force has reached a very high precision with protons and anti-protons<sup>1,2</sup>. The extension of such measurements from (anti-)baryons to (anti-)nuclei allows one to probe any difference in the interactions between nucleons and antinucleons encoded in the (anti-)nuclei masses. This force is a remnant of the underlying strong interaction among quarks and gluons and can be described by effective theories<sup>3</sup>, but cannot yet be directly derived from quantum chromodynamics. Here we report a measurement of the difference between the ratios of the mass and charge of deuterons (d) and anti-deuterons ( $\bar{d}$ ), and <sup>3</sup>He and  $\bar{3}\text{He}$  nuclei carried out with the ALICE (A Large Ion Collider Experiment)<sup>4</sup> detector in Pb-Pb collisions at a centre-of-mass energy per nucleon pair of 2.76 TeV. Our direct measurement of the mass-over-charge differences confirms

and specific energy loss (dE/dx) measurements, and the TOF (time of flight)<sup>23</sup> detector to measure the time  $t_{\text{TOF}}$  needed by each track to traverse the detector. The combined ITS and TPC information is used to determine the track length (L) and the rigidity (p/z, where p is the momentum and z the electric charge in units of the elementary charge e) of the charged particles in the solenoidal 0.5 T magnetic field of the ALICE central barrel (pseudorapidity  $|\eta| < 0.8$ ). On the basis of these measurements, we can extract the squared mass-over-charge ratio  $\mu_{\text{TOF}}^2 \equiv (m/z)_{\text{TOF}}^2 = (p/z)^2 [(t_{\text{TOF}}/L)^2 - 1/c^2]$ . The choice of this variable is motivated by the fact that  $\mu^2$  is directly proportional to the square of the time of flight, allowing to better preserve its Gaussian behaviour.

The high precision of the TOF detector, which determines the arrival time of the particle with a resolution of 80 ps (ref. 20), allows us to measure a clear signal for (anti-)protons, (anti-)deuterons and

# ALICE studies possible light tetraquark



**ALICE**

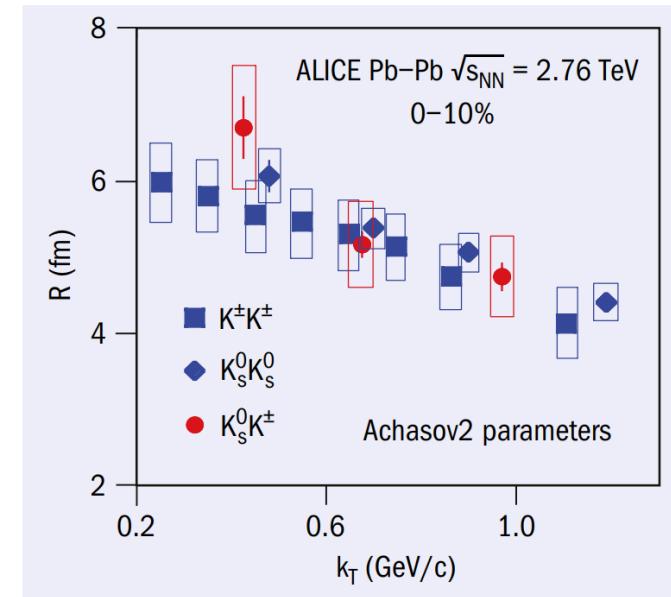
The  $a_0(980)$  resonance is formally classified by the Particle Data Group as a light diquark (quark + antiquark) meson similar to the pion.

However, it has long been considered as a candidate tetraquark state made up of two quarks and two antiquarks. Existing experimental evidence based on the radiative decay of the  $\phi$  meson has not been convincing, so the ALICE collaboration took a different approach to study the  $a_0$  by measuring  $K_S^0-K^\pm$  correlations in lead–lead collisions at the LHC. Since the kaons are not identical there is no Hanbury–Brown–Twiss interferometry enhancement, and since the  $K_S^0$  is uncharged there is no Coulomb effect. Nevertheless, because the rest masses of the two kaons reach the threshold to produce the  $a_0$  it is expected that there is a strong final-state interaction between the two kaons through the  $a_0$  resonant channel.

Using the data from central lead–lead collisions with a nucleon–nucleon energy

of 2.76 TeV, ALICE fitted the experimental two-kaon yield to extract the radius and emission strength of the kaon source assuming only a final-state interaction through the  $a_0$  (see figure).

Both the radii and the emission strength from the  $K_S^0-K^\pm$  analysis agree with the identical kaon results, suggesting that the final-state interaction between the  $K_S^0$  and  $K^\pm$  goes solely through the  $a_0$  resonance without any competing non-resonant channels. A tetraquark  $a_0$  is expected to couple more strongly to the two kaons, since it has the same quark content, while the formation of a diquark state requires the annihilation of the strange quarks, which is suppressed due to geometric effects and a selection rule. Although there are no quantitative predictions for the magnitude of this suppression that would result for a diquark form of  $a_0$ , the qualitative expectation is that this would open up non-resonant channels that would compete with the  $a_0$  final-state interaction, making it smaller than the



*Radius parameters versus average transverse kaon-pair momentum determined from  $K_S^0-K^\pm$  correlations and identical-kaon correlations in central ALICE lead–lead collisions.*

identical-kaon values. The ALICE result of the final-state interaction going solely via the  $a_0$  thus favours the interpretation of the  $a_0$  as a tetraquark state.

## ● Further reading

ALICE Collaboration 2017 *Phys. Lett. B* **774** 64.



# Status of femtoscopy in ALICE

## • Previous results

- Multipion Bose-Einstein correlations in pp, p-Pb, and Pb-Pb collisions at the LHC, Phys. Rev. C 93 (2016) 054908
- Centrality dependence of pion freeze-out radii in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV Phys. Rev. C 93 (2016) 024905
- 1D pion, kaon, proton femtoscopy in Pb-Pb Phys. Rev. C 92 (2015) 054908
- Pion femtoscopy in p-Pb Phys. Rev. C 91 (2015) 034906
- Freeze-out radii extracted from three-pion cumulants in pp, p-Pb and Pb-Pb collisions at the LHC, Phys. Lett. B 739 (2014) 139-151
- Two and Three-Pion Quantum Statistics Correlations in Pb-Pb Collisions at  $\sqrt{s_{NN}}=2.76$  TeV at the LHC Phys. Rev. C 89 (2014) 024911
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- K0sK0s correlations in pp collisions at  $\sqrt{s}=7$  TeV from the LHC ALICE experiment Phys. Lett. B 717 (2012) 151-161
- Femtoscopy in pp at 0.9 and 7 TeV: Phys. Rev. D 84 (2011) 112004,
- Two-pion Bose-Einstein correlations in central Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV Phys. Lett. B 696 (2011) 328-337
- Two-pion Bose-Einstein correlations in pp collisions at  $\sqrt{s}=900$  GeV, Phys. Rev. D 82 (2010) 052001

## • Newly published papers:

- Azimuthally-differential pion femtoscopy relative to the third harmonic event plane in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, arxiv: 1803.10594
- Azimuthally differential pion femtoscopy in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Rev. Lett. 118 (2017) 222301
- Kaon femtoscopy in Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Rev. C 96 (2017) 064613
- Measuring K0SK± interactions using Pb-Pb collisions at  $\sqrt{s_{NN}}=2.76$  TeV, Phys. Lett. B 774 (2017) 64

## • Preliminary results:

- Baryon results:
  - $p\bar{p}, \bar{p}p$ , pp from Run2
  - Baryon-baryon correlations ( $p\Lambda, \bar{p}\Lambda, \Lambda\Lambda$ , and  $\bar{\Lambda}\bar{\Lambda}$ ) from Run1 and Run2
  - Baryon-antibaryon correlations ( $p\bar{p}$ , and  $\bar{p}\Lambda, \bar{p}\Lambda$ , and  $\Lambda\bar{\Lambda}$ ) from Run1 and Run2
  - Analysis of heavier baryons (eg.  $p\Xi, \bar{p}\Xi$ )
- Lambda-K+, Lambda-K-, and Lambda-K0s
- Kaon-proton



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  - Analysis of heavier baryons (eg.  $p\Xi, \bar{p}\Xi$ )
- Lambda-K+, Lambda-K-, and Lambda-K0s
- Kaon-proton

New since last  
WPCF

