

Multi-Higgs solution to flavor anomalies and implications for neutrino physics



Martti Raidal

NICPB, Tallinn

[arXiv:1805.08189](https://arxiv.org/abs/1805.08189)

[arXiv:1901.08290](https://arxiv.org/abs/1901.08290)

Carlo Marzo, Luca Marzola

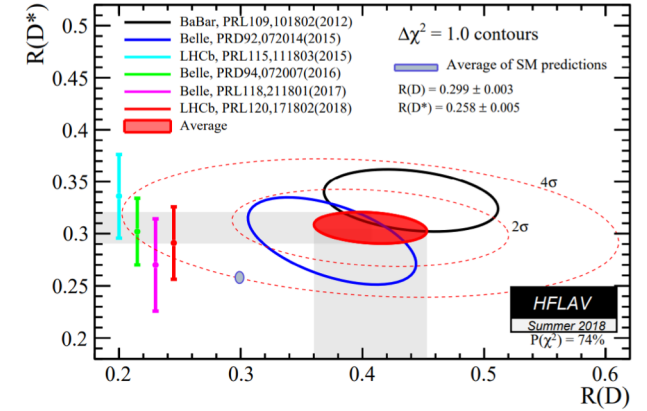
The R_D and R_K - before Moriond 2019

• R_D

$$R_{D^{(*)}} = \frac{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau \bar{\nu})}{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell \bar{\nu})},$$

$$R_D^{\text{SM}} = 0.300 \pm 0.008, \quad R_D^{\text{exp}} = 0.407 \pm 0.039 \pm 0.024,$$

$$R_{D^*}^{\text{SM}} = 0.252 \pm 0.003, \quad R_{D^*}^{\text{exp}} = 0.306 \pm 0.013 \pm 0.007,$$



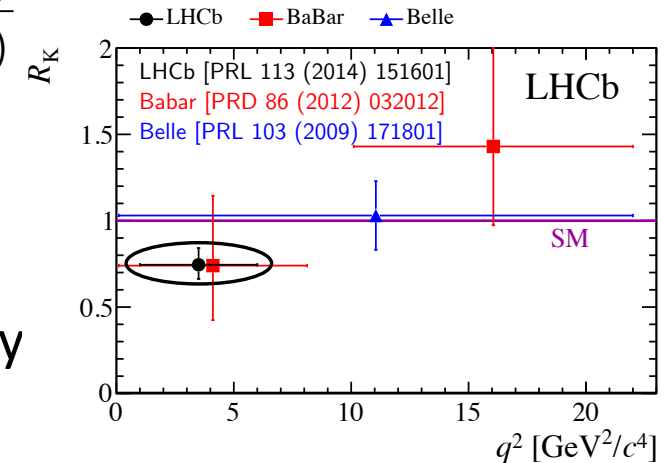
Several $\sim 2\sigma$ deviations by BaBar, Belle, LHCb combine to a 4.1σ anomaly

• R_K

$$R_K = \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)} \quad R_{K^*} = \frac{\mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{\mathcal{B}(B^0 \rightarrow K^{*0} e^+ e^-)}$$

$$R_K^{\text{SM}} = 1.0004 \pm 0.0002, \quad R_{K^*}^{\text{exp}} = 0.745_{-0.074}^{+0.090} \pm 0.036 \quad \text{for } 1 \text{ GeV}^2 \leq q^2 \leq 6 \text{ GeV}^2,$$

$$R_{K^*}^{\text{SM}} = \begin{cases} 0.926 \pm 0.003, \\ 0.9965 \pm 0.0005, \end{cases} \quad R_{K^*}^{\text{exp}} = \begin{cases} 0.66_{-0.07}^{+0.11} \pm 0.03 & \text{for } 0.045 \text{ GeV}^2 \leq q^2 \leq 1.1 \text{ GeV}^2, \\ 0.69_{-0.07}^{+0.11} \pm 0.05 & \text{for } 1.1 \text{ GeV}^2 \leq q^2 \leq 6 \text{ GeV}^2, \end{cases}$$



The LHCb measurements combine to almost $\sim 5\sigma$ anomaly

The R_D and R_K - after Moriond 2019

- R_D by Belle

$$\mathcal{R}_D^{\text{Belle}} = 0.307 \pm 0.037 \pm 0.016 \quad \text{and} \quad \mathcal{R}_{D^*}^{\text{Belle}} = 0.283 \pm 0.018 \pm 0.014, \quad 1.2\sigma \text{ dev.}$$

- R_K by LHCb and R_{K^*} by Belle

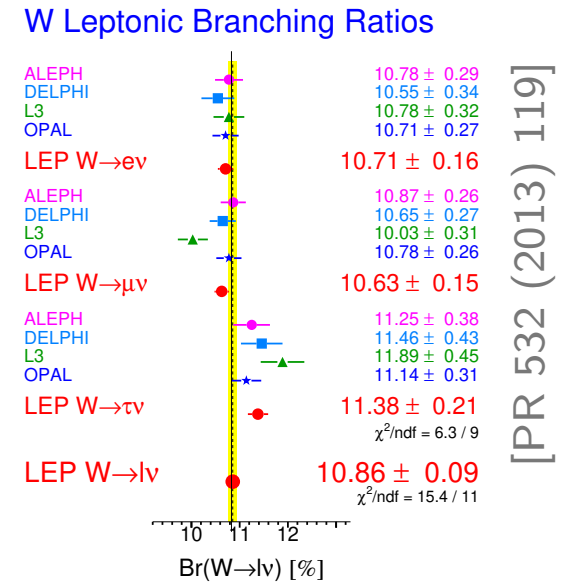
$$R_K = \frac{\text{BR}(B \rightarrow K\mu\mu)}{\text{BR}(B \rightarrow Kee)} = 0.846^{+0.060}_{-0.054} {}^{+0.016}_{-0.014}, \quad \text{for } 1.1 \text{ GeV}^2 < q^2 < 6 \text{ GeV}^2, \quad 2.5\sigma \text{ dev.}$$

$$R_{K^*} = \frac{\text{BR}(B \rightarrow K^*\mu\mu)}{\text{BR}(B \rightarrow K^*ee)} = \begin{cases} 0.90^{+0.27}_{-0.21} \pm 0.10, & \text{for } 0.1 \text{ GeV}^2 < q^2 < 8 \text{ GeV}^2, \\ 1.18^{+0.52}_{-0.32} \pm 0.10, & \text{for } 15 \text{ GeV}^2 < q^2 < 19 \text{ GeV}^2. \end{cases}$$

- The overall picture has not changed much. The significance has been mildly reduced, the same NP explanations are valid.

The peculiar properties of the anomalies require an explanation

- Lepton flavor universality violation!
 - In the SM this refers to the **scalar sector interactions**.
- Very different scales of new physics
 - R_D - NP competes with tree level SM amplitude – **requires 1 TeV scale NP + large coupling**
 - R_K – NP competes with loop level SM amplitude – **requires 10 TeV scale NP**



Is a common NP explanation needed and possible?

Should we work on anomalies?

Should we work on anomalies?

We do not have a choice!

Every discovery is preceded by a deviation but the opposite is not true

Remain reasonable with interpretations, use your common sense

Is a common NP explanation possible?

- Need NP with many freely adjustable parameters
- Majority of flavor physics community has adopted **leptoquark** explanation to the anomalies

Extraordinary claim requires extraordinary evidence

What is my interest to the anomalies?

- To show that non-exotic, “SM-like NP” can simultaneously fit data
- To show that models with many scalars can fit the data
- To determine the minimal scalar model that can explain both observables

Notice: we do allow for GeV scale right-handed neutrinos – new interactions available

Model independent explanations to the anomalies

$$\mathcal{L}_{eff}^{b \rightarrow c \ell \bar{\nu}} = \frac{2G_F V_{cb}}{\sqrt{2}} \left(C_{VL}^{\ell} \mathcal{O}_{VL}^{\ell} + C_{AL}^{\ell} \mathcal{O}_{AL}^{\ell} + C_{SL}^{\ell} \mathcal{O}_{SL}^{\ell} + C_{PL}^{\ell} \mathcal{O}_{PL}^{\ell} + C_{VR}^{\ell} \mathcal{O}_{VR}^{\ell} \right. \\ \left. + C_{AR}^{\ell} \mathcal{O}_{AR}^{\ell} + C_{SR}^{\ell} \mathcal{O}_{SR}^{\ell} + C_{PR}^{\ell} \mathcal{O}_{PR}^{\ell} \right),$$

$$\mathcal{O}_{VX}^{\ell} = [\bar{c} \gamma^{\mu} b] [\bar{\ell} \gamma_{\mu} P_X \nu_{\ell}],$$

$$\mathcal{O}_{AX}^{\ell} = [\bar{c} \gamma^{\mu} \gamma_5 b] [\bar{\ell} \gamma_{\mu} P_X \nu_{\ell}]$$

$$\mathcal{O}_{SX}^{\ell} = [\bar{c} b] [\bar{\ell} P_X \nu_{\ell}],$$

$$\mathcal{O}_{PX}^{\ell} = [\bar{c} \gamma_5 b] [\bar{\ell} P_X \nu_{\ell}],$$

$$-\mathcal{L}_{bs} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (c_i \mathcal{O}_i + c'_i \mathcal{O}'_i),$$

$$\mathcal{O}_7 = (\bar{s} P_L b)(\bar{l} l),$$

$$\mathcal{O}'_7 = (\bar{s} P_R b)(\bar{l} l),$$

$$\mathcal{O}_8 = (\bar{s} P_L b)(\bar{l} \gamma_5 l),$$

$$\mathcal{O}'_8 = (\bar{s} P_R b)(\bar{l} \gamma_5 l),$$

$$\mathcal{O}_9 = (\bar{s} \gamma_{\mu} P_L b)(\bar{l} \gamma^{\mu} l),$$

$$\mathcal{O}'_9 = (\bar{s} \gamma_{\mu} P_R b)(\bar{l} \gamma^{\mu} l),$$

$$\mathcal{O}_{10} = (\bar{s} \gamma_{\mu} P_L b)(\bar{l} \gamma^{\mu} \gamma_5 l),$$

$$\mathcal{O}'_{10} = (\bar{s} \gamma_{\mu} P_R b)(\bar{l} \gamma^{\mu} \gamma_5 l).$$

$$C_{VL} = -C_{AL} = 1.$$

$$C_9 - C_9^{SM} \simeq -1.21 \text{ or } C_9 - C_9^{SM} = -\left(C_{10} - C_{10}^{SM}\right) \simeq -0.67$$

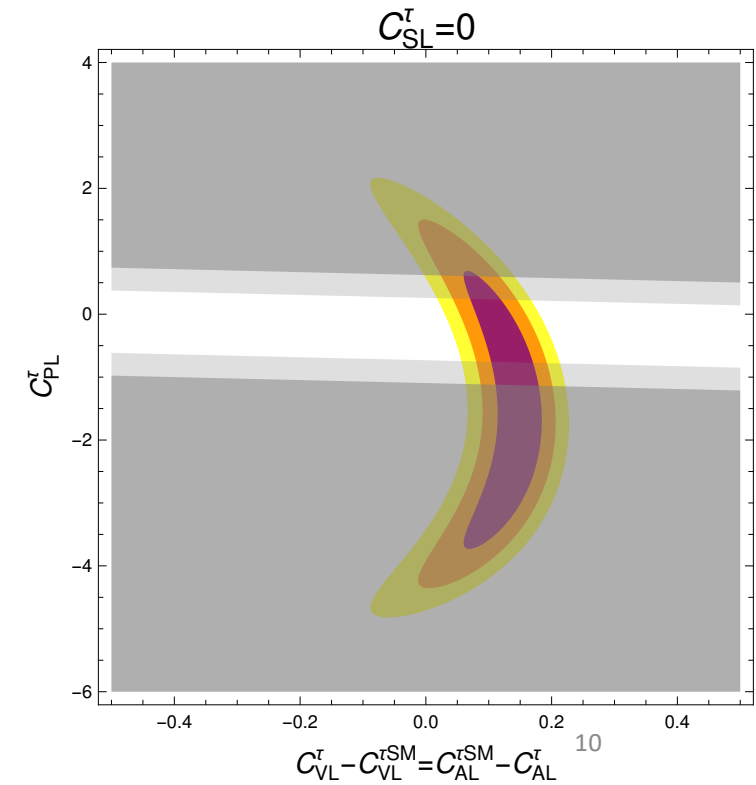
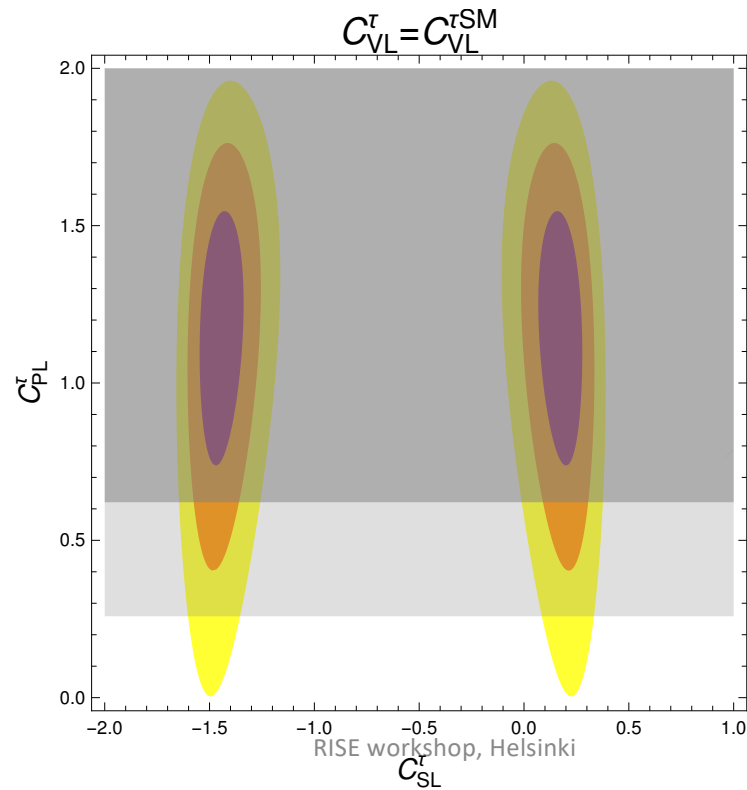
- 0. Vector (or axial vector) operators are strongly favoured by global fits to data in both cases
- 1. Small fraction of scalar contribution improves the global fit

R_D model independently

$$R_{D^{(*)}} = \frac{\mathcal{B}_\tau^{D^{(*)}}(m_\tau, \mathbf{C}^\tau)}{\mathcal{B}_{\mu/e}^{D^{(*)}}(m_{\mu/e}, \mathbf{C}^{\mu/e})}$$

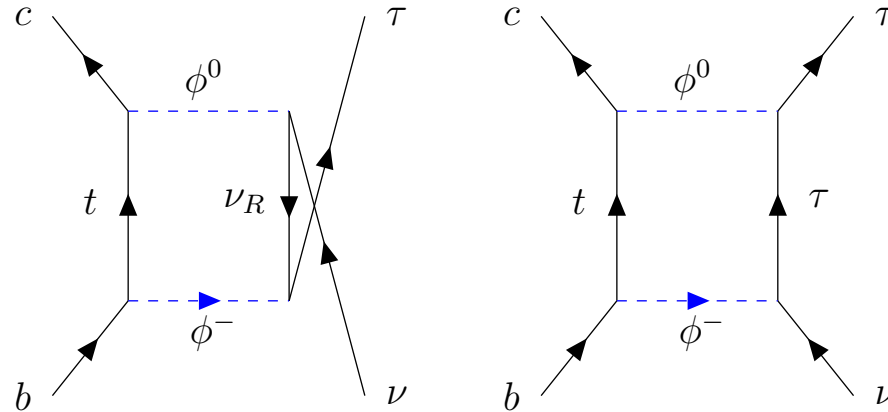
$$\mathcal{B}_{\tau\nu} = \frac{m_{B_c} m_\tau^2 f_{B_c}^2 G_F^2 |V_{cb}|^2}{8\pi \Gamma_{B_c^-}} \left(1 - \frac{m_\tau^2}{m_{B_c}^2}\right)^2 \left| \frac{m_{B_c}^2}{m_\tau(m_b + m_c)} C_{PL}^\tau - C_{AL}^\tau \right|^2$$

- Scalar operators strongly constrained by B_c lifetime allowing 10% (30%) deviations
- Vector operators are OK



The idea

- Simplified model

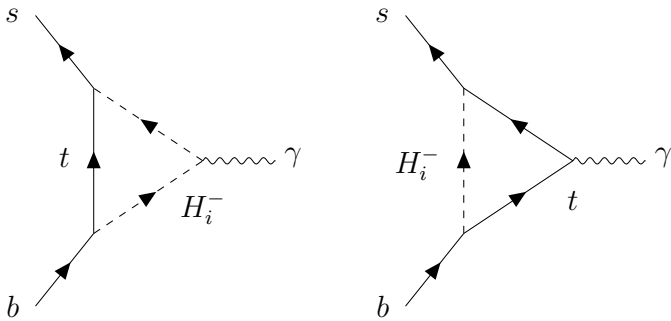


$$\begin{aligned}
 -\mathcal{L} \supset & \xi_{\bar{c}_L t_R} \phi^{0*} \bar{c}_L t_R + \xi_{\bar{t}_R b_L} \phi^+ \bar{t}_R b_L + \xi_{\bar{\tau}_L \tau_R} \phi^0 \bar{\tau}_L \tau_R - \xi_{\bar{\tau}_R \nu_L} \phi^- \bar{\tau}_R \nu_L + \\
 & + \xi_{\bar{\tau}_L \nu_R} \phi^- \bar{\tau}_L \nu_R + \xi_{\bar{\nu}_R \nu_L} \phi^0 \bar{\nu}_R \nu_L + h.c., \\
 C_{VL}^1 = -C_{AL}^1 = & - \left(\frac{m_W^2}{8\pi^2 V_{cb} g_W^2} \right) (\xi_{c_L t_R} \xi_{t_R b_L} \xi_{\nu_R \nu_L} \xi_{\tau_L \nu_R}) D_{dd00}[m_{\nu_R}^2, m_t^2, m_{\phi^0}^2, m_{\phi^-}^2] \\
 C_{VL}^2 = -C_{AL}^2 = & - \left(\frac{m_W^2}{8\pi^2 V_{cb} g_W^2} \right) (\xi_{c_L t_R} \xi_{t_R b_L} \xi_{\tau_L \tau_R} \xi_{\tau_R \nu_L}) D_{dd00}[m_t^2, m_{\phi^0}^2, m_{\phi^-}^2, m_\tau^2].
 \end{aligned}$$

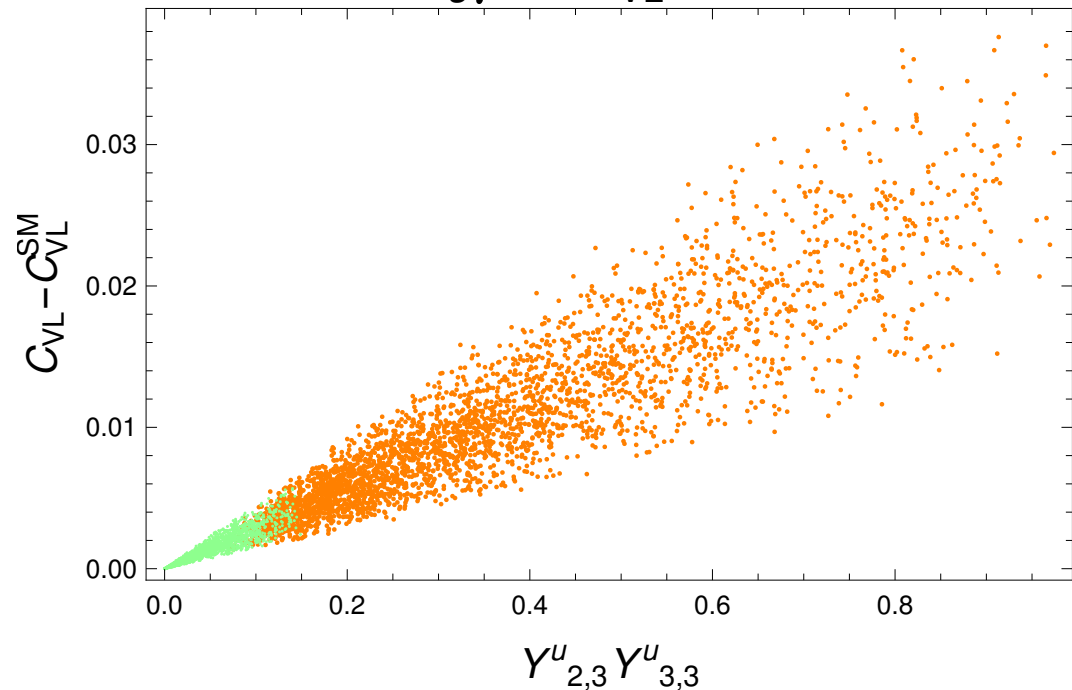
**New RH interactions
Enhancement by light, GeV mass RHNs**

Does 2HDM+ v_R work?

$$B \rightarrow X_s \gamma$$



$B \rightarrow X_s \gamma$ vs C_{VL} in 2HDM



NO! Existing constraints are too strong. Not enough free parameters (Yukawas) to fit all data!

3HDM+v_R

$$H_0 = \begin{pmatrix} H_0^+ \\ H_0^0 \end{pmatrix}, \quad H_1 = \begin{pmatrix} H_1^+ \\ H_1^0 \end{pmatrix}, \quad H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix}$$

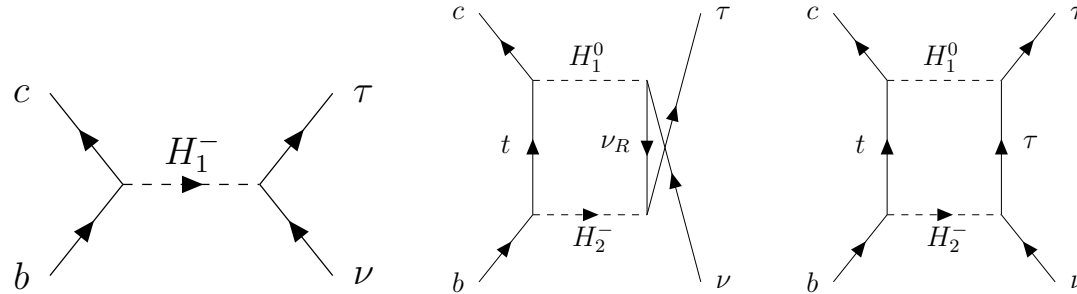
$$-\mathcal{L} \supset \bar{Q}_L \tilde{H}_1 \mathcal{Y}_1^u u_R + \bar{Q}_L \tilde{H}_2 \mathcal{Y}_2^u u_R + \bar{L}_L H_1 \mathcal{Y}_1^\ell \ell_R + \bar{L}_L H_2 \mathcal{Y}_2^\ell \ell_R + \bar{L}_L \tilde{H}_1 \mathcal{Y}_1^\nu \nu_R + \bar{L}_L \tilde{H}_2 \mathcal{Y}_2^\nu \nu_R + \bar{Q}_L H_1 \mathcal{Y}_1^d d_R + h.c.$$

$$\mathcal{Y}_1^u = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & f_{\bar{c}_L t_R} \\ 0 & f_{\bar{b}_L c_R} & 0 \end{pmatrix}, \quad \mathcal{Y}_1^d = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & f_{\bar{c}_L b_R} \\ 0 & 0 & 0 \end{pmatrix}, \quad \mathcal{Y}_2^u = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & g_{\bar{b}_L t_R} \end{pmatrix},$$

$$\mathcal{Y}_1^\nu = \begin{pmatrix} 0 & 0 & 0 \\ 0 & f_{\bar{\nu}_L \nu_R} & 0 \\ 0 & 0 & f'_{\bar{\nu}_L \nu_R} \end{pmatrix}, \quad \mathcal{Y}_1^\ell = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & f_{\bar{\tau}_L \tau_R} \end{pmatrix}, \quad \mathcal{Y}_2^\nu = \begin{pmatrix} 0 & 0 & 0 \\ 0 & g_{\bar{\nu}_L \nu_R} & 0 \\ 0 & 0 & g'_{\bar{\nu}_L \nu_R} \end{pmatrix}, \quad \mathcal{Y}_2^\ell = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & g_{\bar{\nu}_L \tau_R} \end{pmatrix}$$

Textures constructed to fit all observables simultaneously consistently with constraints

R_D in our model



$$C_{VL}^{\tau(1)} = -C_{AL}^{\tau(1)} = \left(-\frac{m_W^2}{8\pi^2 V_{cb} g_W^2} \right) (f_{\bar{c}_L t_R} g_{\bar{b}_L t_R} f'_{\bar{\nu}_L \nu_R} g'_{\bar{\nu}_L \nu_R}) D_{dd00}[m_{\nu_R}^2, m_t^2, m_{H_1^0}^2, m_{H_2^-}^2],$$

$$C_{VL}^{\tau(2)} = -C_{AL}^{\tau(2)} = \left(-\frac{m_W^2}{8\pi^2 V_{cb} g_W^2} \right) (f_{\bar{c}_L t_R} g_{\bar{b}_L t_R} f_{\bar{\tau}_L \tau_R} g_{\bar{\nu}_L \tau_R}) D_{dd00}[m_t^2, m_{H_1^0}^2, m_{H_2^-}^2, m_{\tau}^2],$$

Must be maximized

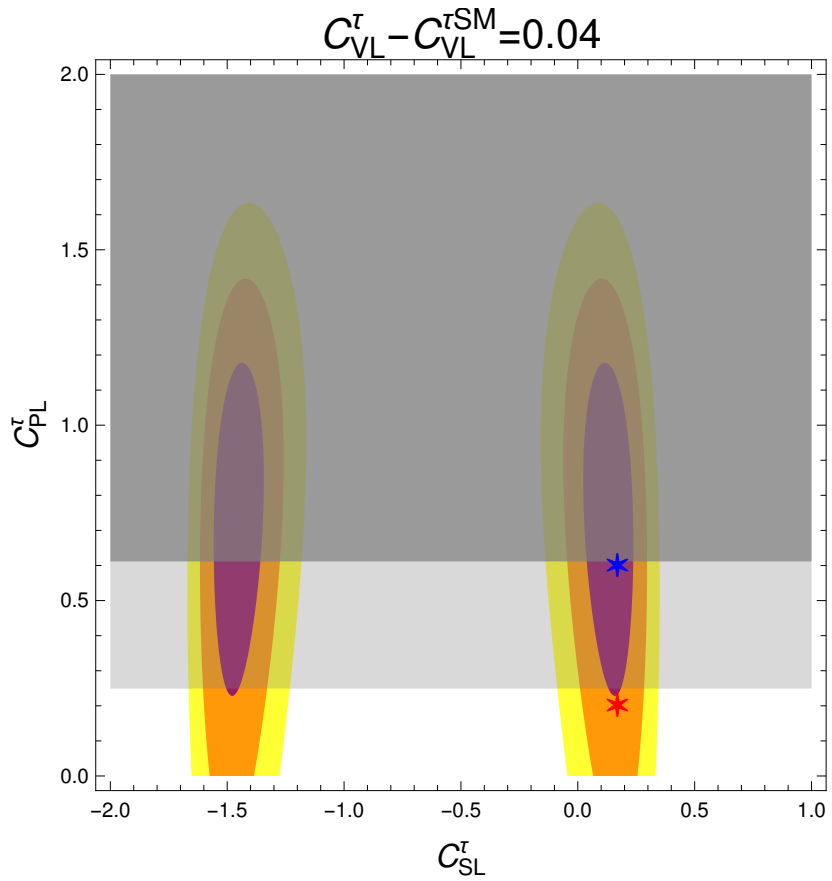
$$C_{SL}^{\tau} = -\frac{2 m_W^2}{V_{cb} g_W^2 m_{H_1^-}^2} f_{\bar{\nu}_L \tau_R} (f_{\bar{b}_L c_R} - f_{\bar{c}_L b_R}),$$

$$C_{PL}^{\tau} = -\frac{2 m_W^2}{V_{cb} g_W^2 m_{H_1^-}^2} f_{\bar{\nu}_L \tau_R} (f_{\bar{b}_L c_R} + f_{\bar{c}_L b_R}),$$

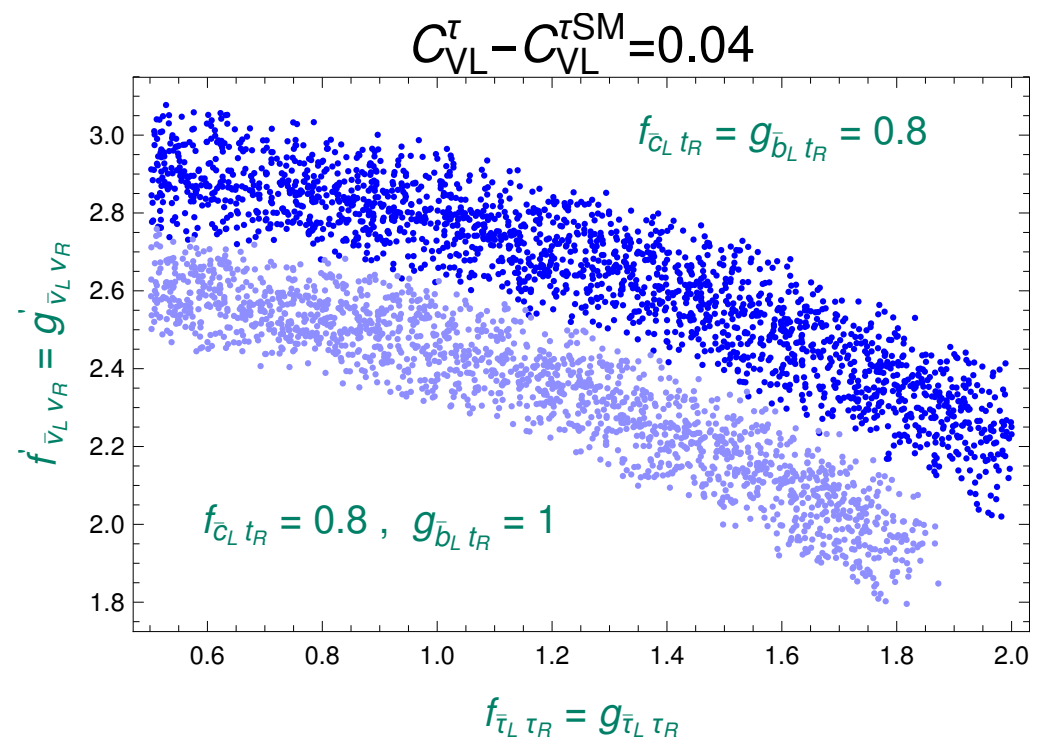
Must be minimized

We require that all existing exp. constraints are satisfied

Benchmark points for R_D and vector operators



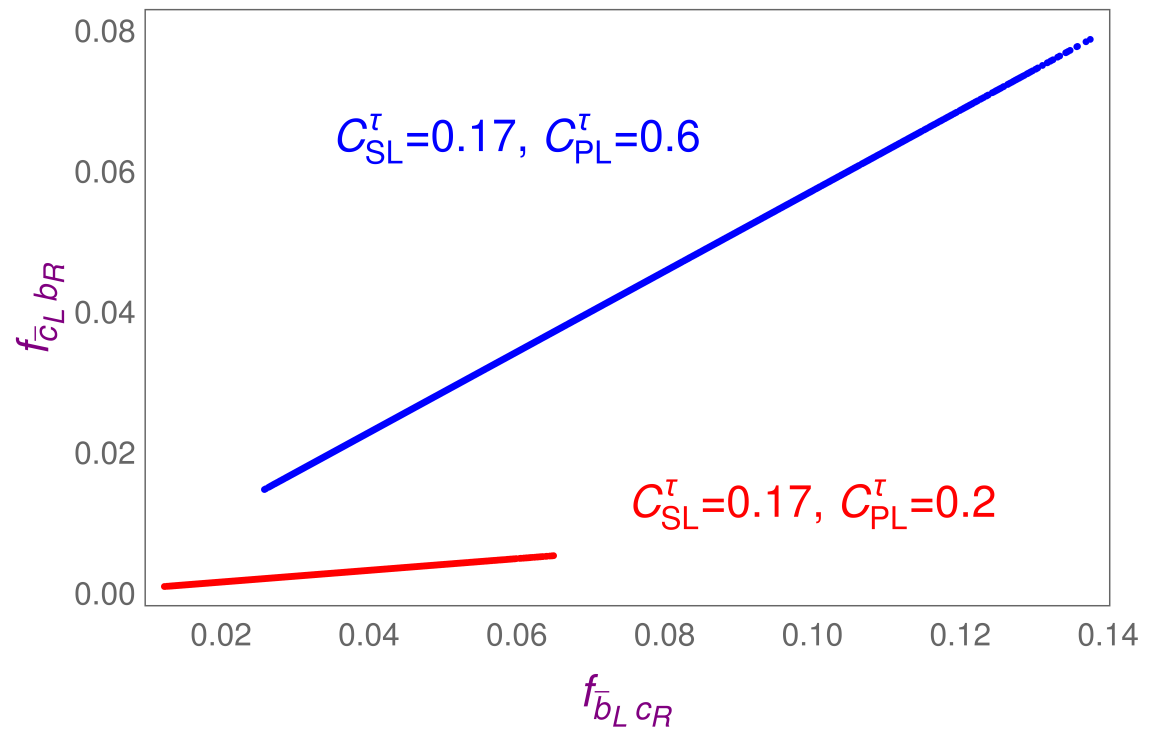
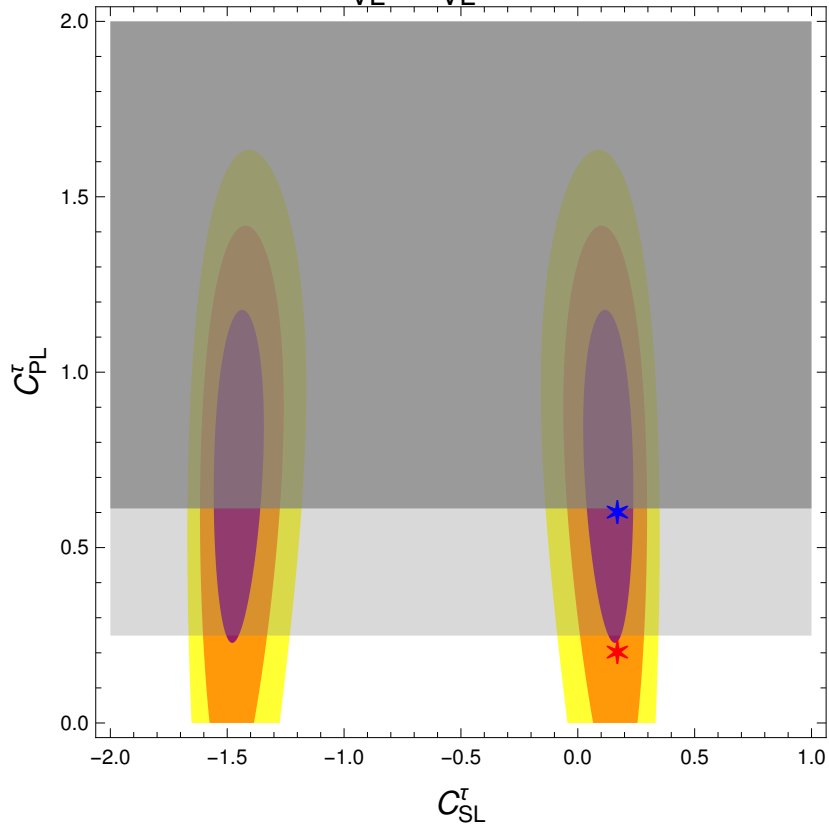
$350 \text{ GeV} < M < 400 \text{ GeV}$



Yukawa couplings of order 1 are needed

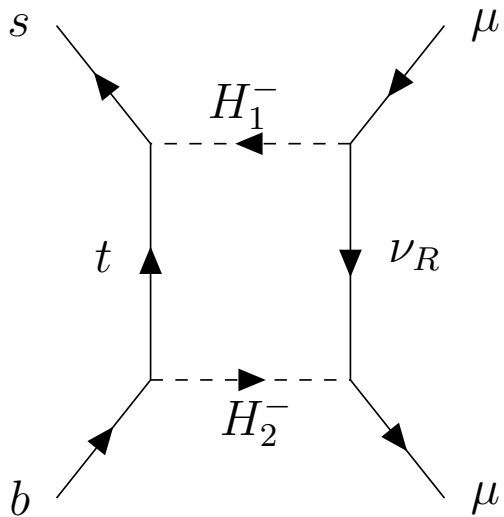
Benchmark points for R_D and scalar operators

$$C_{VL}^T - C_{VL}^{TSM} = 0.04$$



The related Yukawa couplings must be small

R_K in our model



$$-\mathcal{L} \supset f_{\bar{\mu}_L \nu_R} H_1^- \bar{\mu}_L \nu_R + g_{\bar{\mu}_L \nu_R} H_2^- \bar{\mu}_L \nu_R + h.c.,$$

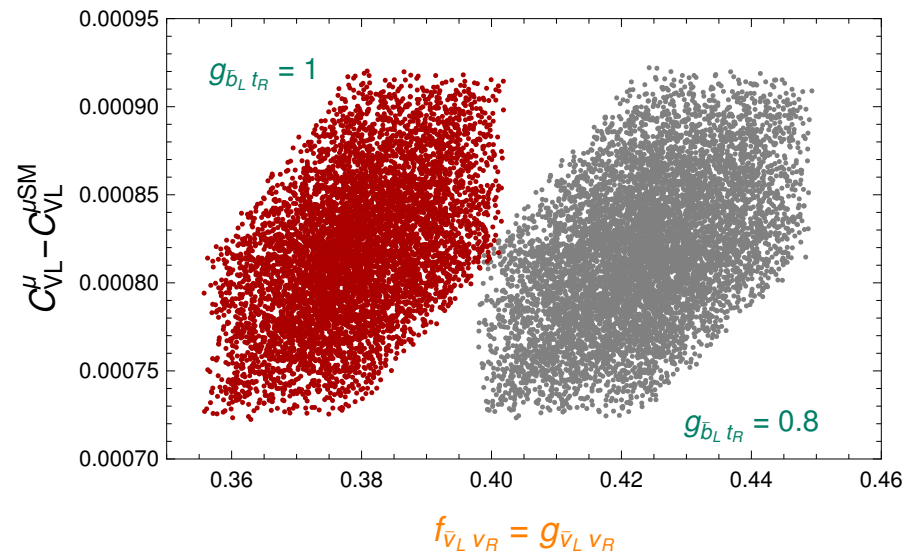
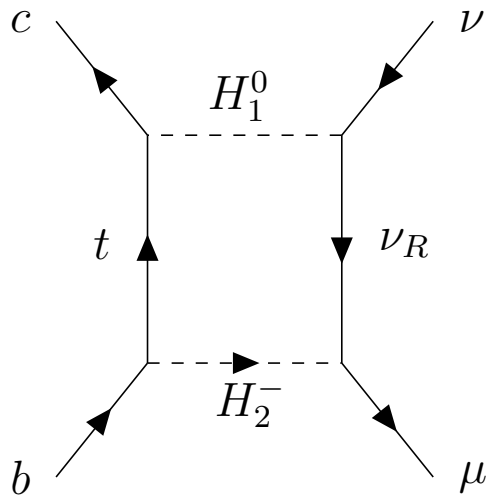
$$-C_9 = C_{10} = -\frac{m_W^2}{4\pi\alpha g_W^2} \frac{V_{cs}^*}{V_{ts}^*} (f_{\bar{c}_L t_R} g_{\bar{b}_L t_R} f_{\bar{\mu}_L \nu_R} g_{\bar{\mu}_L \nu_R}) D_{dd00}[m_t^2, m_{H_1^-}^2, m_{H_2^-}^2],$$

$$O_9 = -O_{10} = -0.67,$$

Is needed to explain the R_K anomaly

Can R_K be explained simultaneously? **Yes!**

- Couplings inducing R_K enter to denominator of $R_{D^{(*)}} = \frac{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau \bar{\nu})}{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell \bar{\nu})}$



Any correlation is numerically small

ϵ'/ϵ anomaly?

$$\text{Re}(\epsilon'/\epsilon) = 1.38(5.15)(4.59) \times 10^{-4}, \quad \text{Re}(\epsilon'/\epsilon)_{\text{exp}} = (16.6 \pm 2.3) \times 10^{-4}.$$

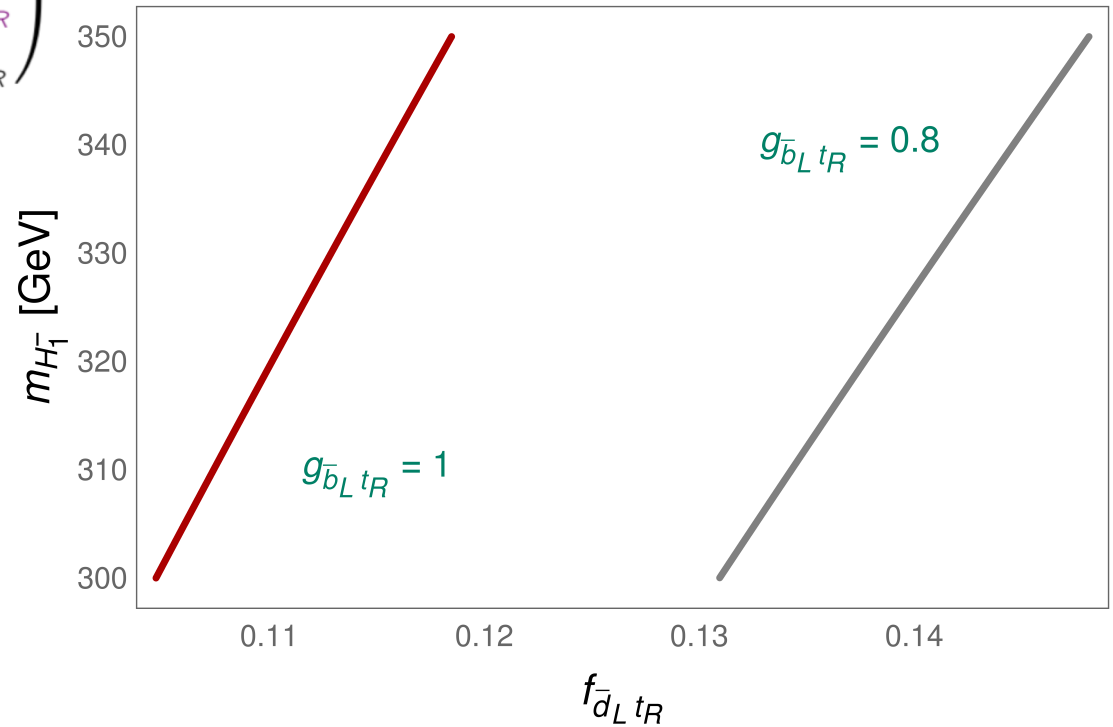
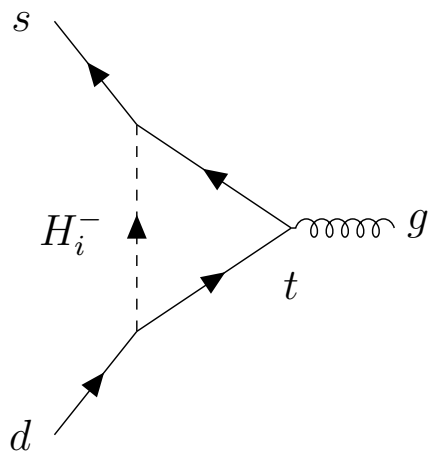
If there is an anomaly, no problem to explain it simultaneously with others. There is enough free parameters.

$$\left| \text{Re}(\epsilon'/\epsilon)_{8g} \sim - (1.85 \times 10^5 \text{ GeV}) \times \text{Im}(C_{8g}^-) \right.,$$

$$C_{8g}^- = -\frac{G_F}{\sqrt{2}} V_{ts}^* V_{td} (m_d C'_{8g} - m_s C_{8g}) .$$

ϵ'/ϵ anomaly in our model

$$\mathcal{Y}_1^d = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & f_{\bar{c}_L b_R} \\ 0 & 0 & 0 \end{pmatrix} \rightarrow \mathcal{Y}_1^d = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & f_{\bar{c}_L b_R} \\ 0 & 0 & f_{\bar{s}_L t_R} \end{pmatrix}$$



$$Re(\epsilon'/\epsilon)_{8g} = 10^{-3}$$

Can easily be explained in our model

Distinctive prediction of the scenario

- RH neutrino masses must be below 10 GeV

$$m_{\nu_R} < 10 \text{ GeV}$$

- Flavour anomalies are related to the leptogenesis and neutrino masses
- ν_R can be produced and discovered at SHIP

This scenario is predictive and testable

Conclusions

- All the considered flavor anomalies can be simultaneously explained in $3\text{HDM}+\nu_R$
- RH neutrinos must be light, induce baryogenesis and neutrino masses, and may show up at SHIP experiment