

Portraying Double Higgs at the Large Hadron Collider



K.C. Kong
University of Kansas

**Symposium on
Artificial Intelligence for Science, Industry and Society
October 21-25, 2019
Mexico City, UNAM**

What is the universe made of?

- An age old question, but we live at a particularly interesting time:
 - We know how “big” the universe is.
 - We have no idea what most of it is made of.
- How does the universe work?

COSMOLOGY MARCHES ON



Physical Sciences

Mathematics

Physics and
Astronomy

Chemistry

Earth System Science

Particle
Physics

Atomic
Physics

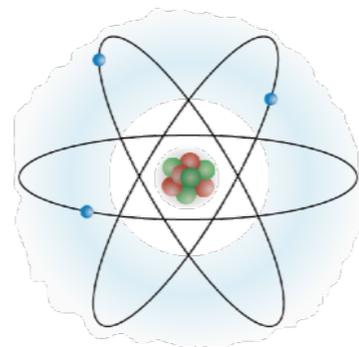
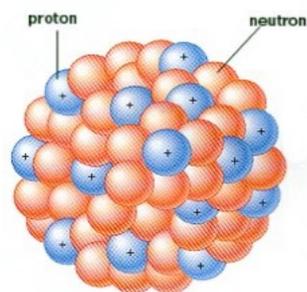
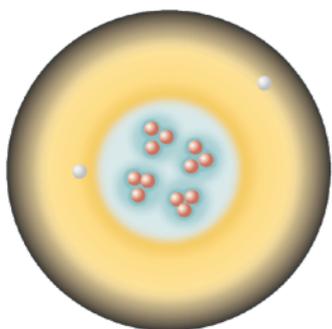
Biological
Physics

Cosmology

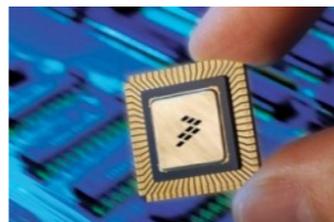
small

big

Nuclear
Physics



Condensed
Matter
Physics

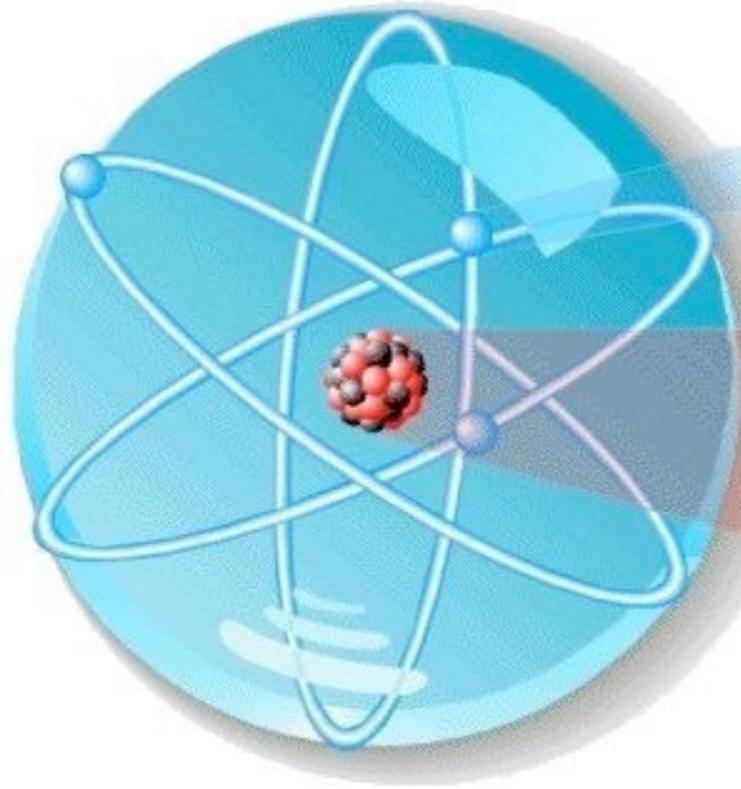


Astrophysics



SMALL: STATE OF THE ART

atom



10^{-10} meters
(thickness of human
hair $\sim 10^{-5}$ m)

electron
 $< 10^{-16}$ cm

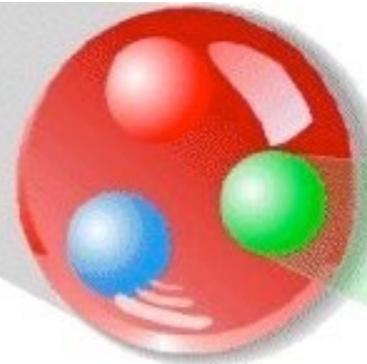


nucleus



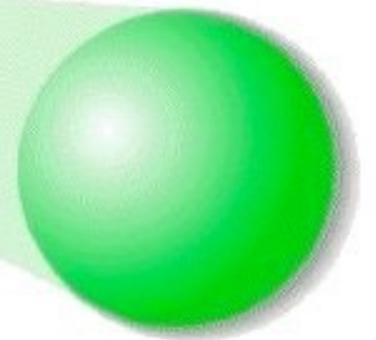
10^{-14}
meters

**proton
neutron**



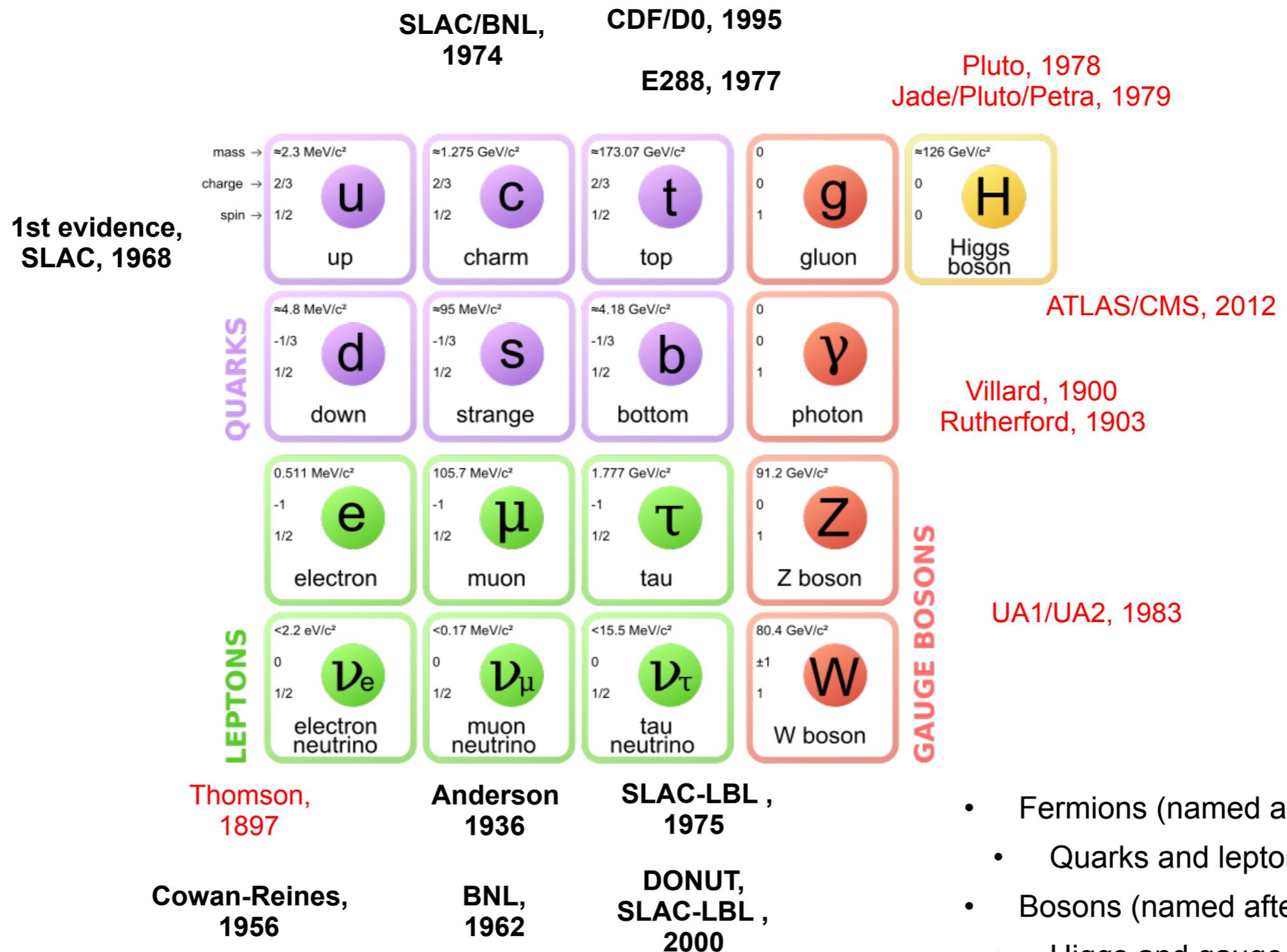
10^{-15}
meters

**up quark
down quark**



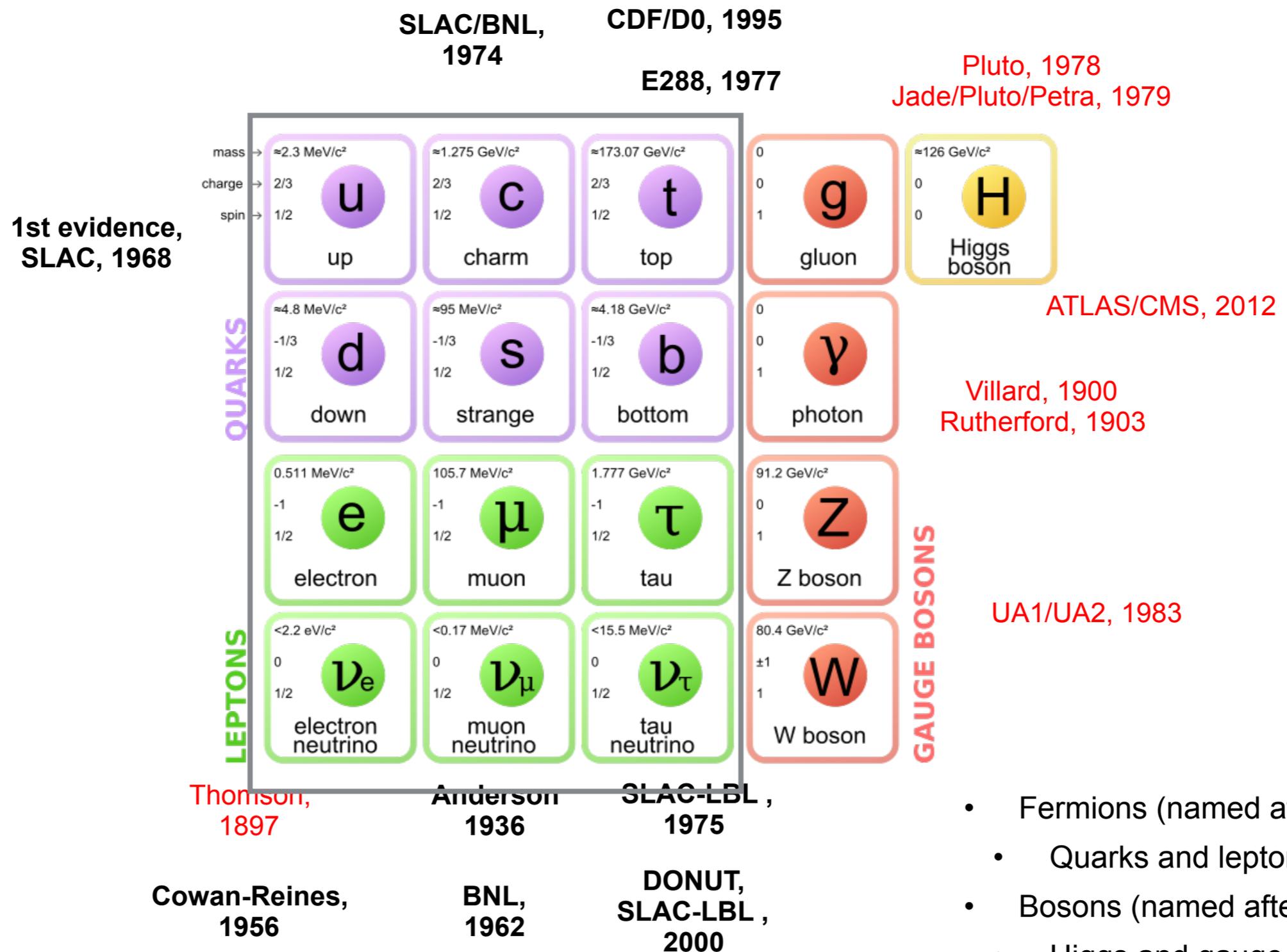
$< 10^{-18}$
meters

Standard Model (Periodic Table for Elementary Particles)



- Fermions (named after Fermi)
- Quarks and leptons
- Bosons (named after Bose)
- Higgs and gauge bosons

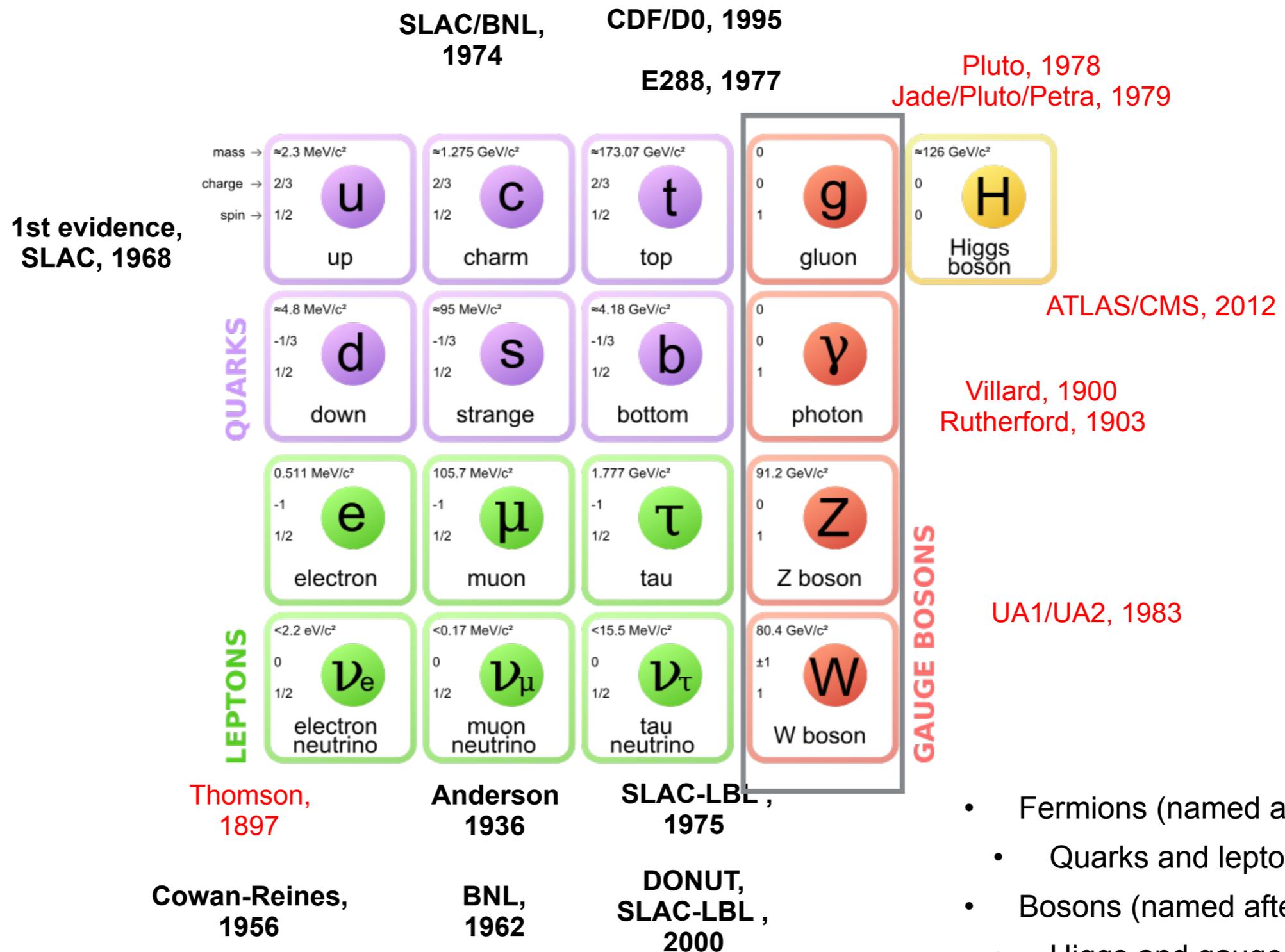
Standard Model (Periodic Table for Elementary Particles)



- Fermions (named after Fermi)
 - Quarks and leptons
- Bosons (named after Bose)
 - Higgs and gauge bosons

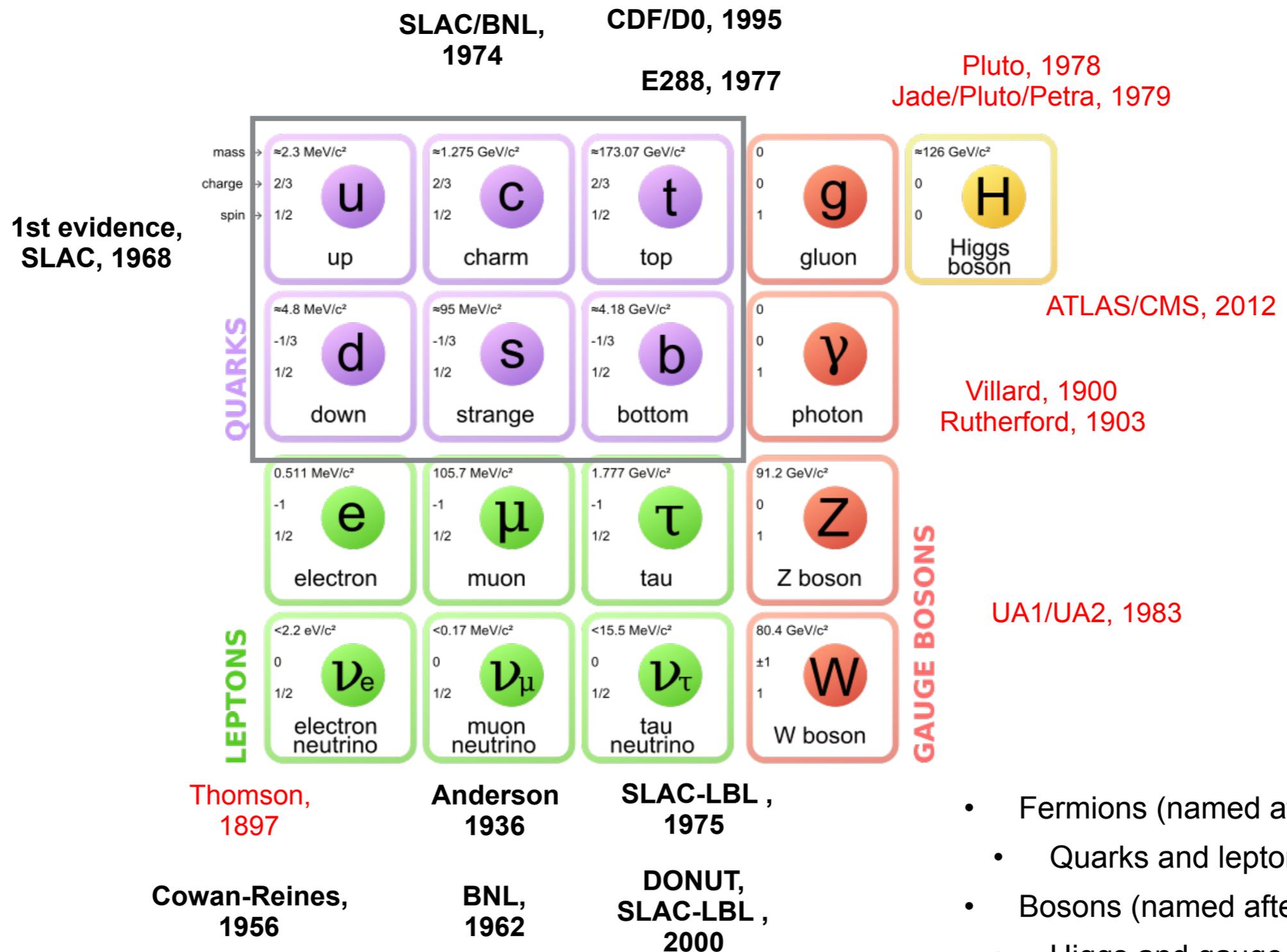
Standard Model

(Periodic Table for Elementary Particles)



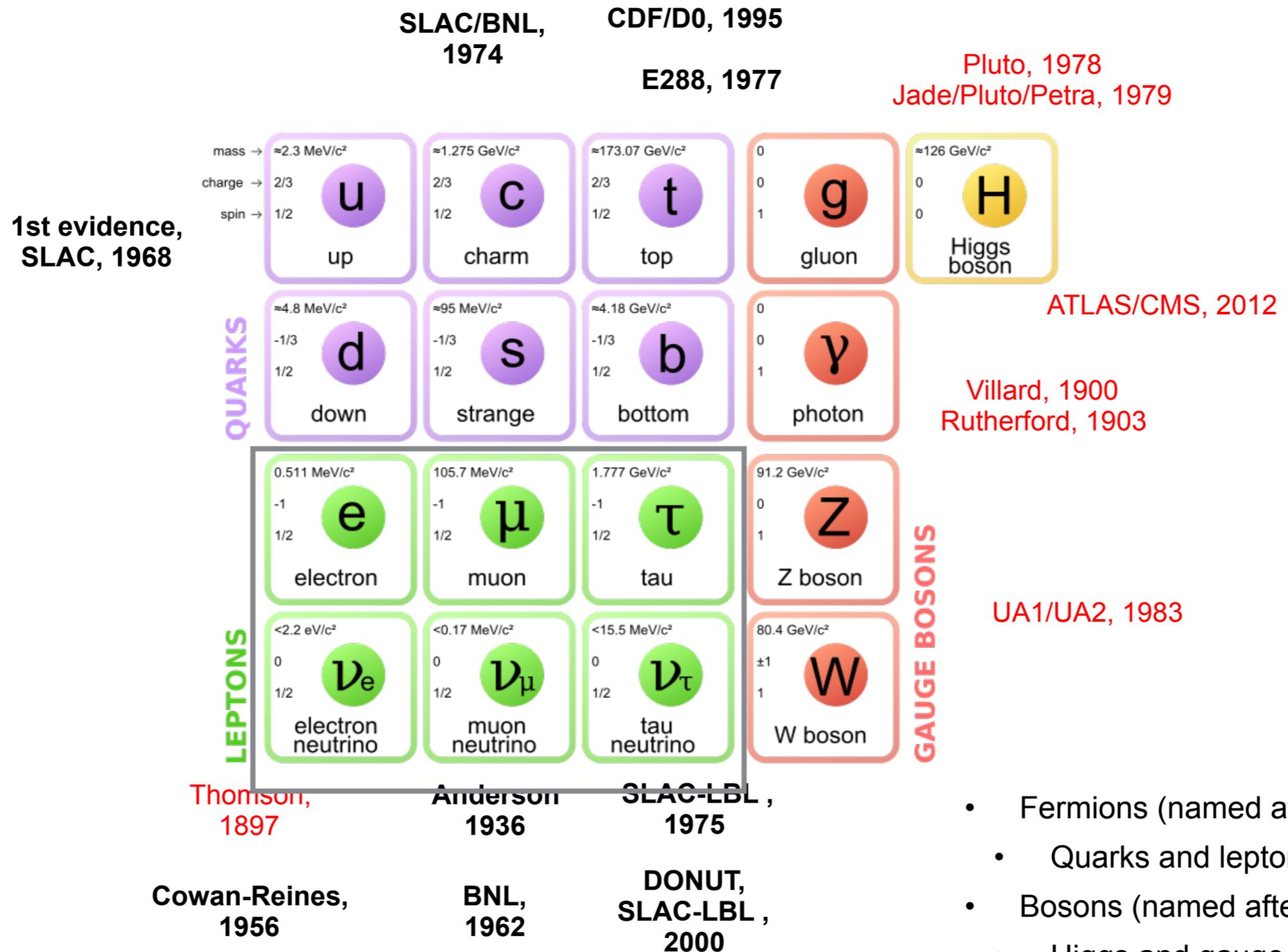
- Fermions (named after Fermi)
- Quarks and leptons
- Bosons (named after Bose)
- Higgs and gauge bosons

Standard Model (Periodic Table for Elementary Particles)



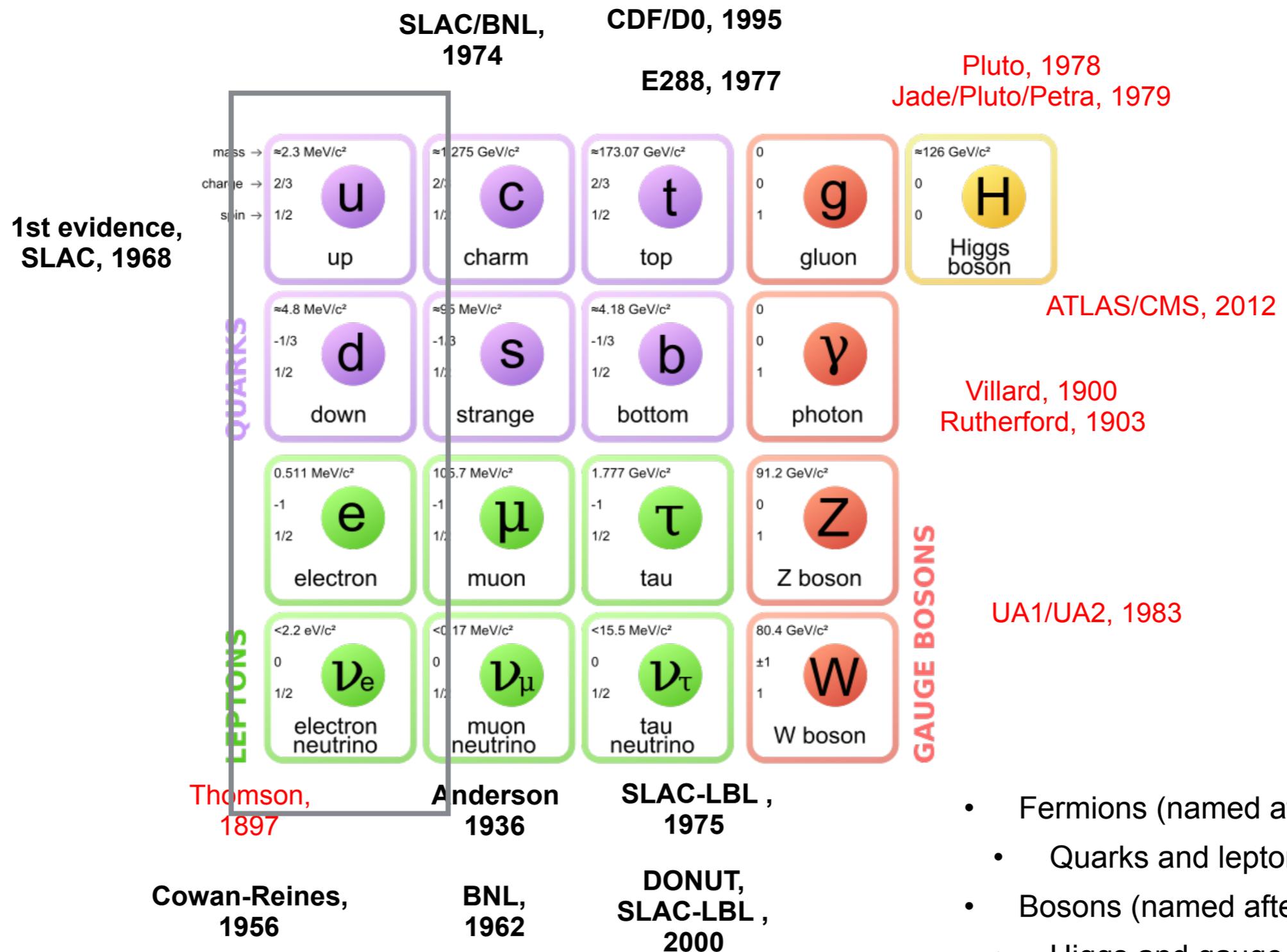
- Fermions (named after Fermi)
- Quarks and leptons
- Bosons (named after Bose)
- Higgs and gauge bosons

Standard Model (Periodic Table for Elementary Particles)



- Fermions (named after Fermi)
- Quarks and leptons
- Bosons (named after Bose)
- Higgs and gauge bosons

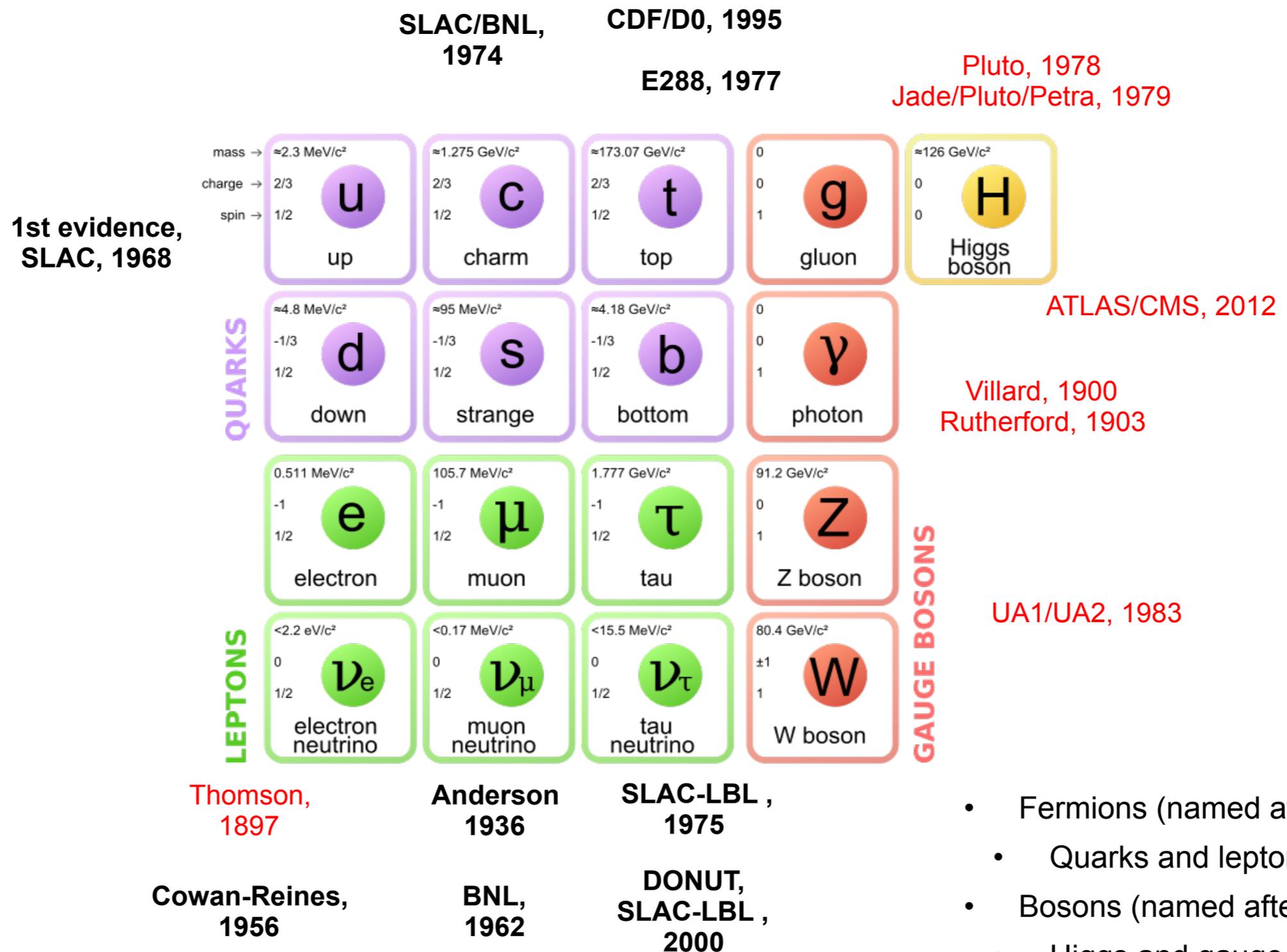
Standard Model (Periodic Table for Elementary Particles)



- Fermions (named after Fermi)
- Quarks and leptons
- Bosons (named after Bose)
- Higgs and gauge bosons

Standard Model

(Periodic Table for Elementary Particles)



- Fermions (named after Fermi)
- Quarks and leptons
- Bosons (named after Bose)
- Higgs and gauge bosons

Standard Model (Periodic Table for Elementary Particles)

	SLAC/BNL, 1974	CDF/D0, 1995 E288, 1977	Pluto, 1978 Jade/Pluto/Petra, 1979	
mass →	≈2.3 MeV/c ²	≈1.275 GeV/c ²	≈173.07 GeV/c ²	0
charge →	2/3	2/3	2/3	0
spin →	1/2	1/2	1/2	1
1st evidence, SLAC, 1968	u up	c charm	t top	g gluon
				H Higgs boson
QUARKS	≈4.8 MeV/c ²	≈95 MeV/c ²	≈4.18 GeV/c ²	0
	-1/3	-1/3	-1/3	0
	1/2	1/2	1/2	1
	d down	s strange	b bottom	γ photon
				Z Z boson
	0.511 MeV/c ²	105.7 MeV/c ²	1.777 GeV/c ²	91.2 GeV/c ²
	-1	-1	-1	0
	1/2	1/2	1/2	1
	e electron	μ muon	τ tau	W W boson
LEPTONS	<2.2 eV/c ²	<0.17 MeV/c ²	<15.5 MeV/c ²	80.4 GeV/c ²
	0	0	0	±1
	1/2	1/2	1/2	1
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson
				G Gauge bosons
Thomson, 1897	Anderson 1936	SLAC-LBL, 1975		
Cowan-Reines, 1956	BNL, 1962	DONUT, SLAC-LBL, 2000		
				Villard, 1900 Rutherford, 1903
				ATLAS/CMS, 2012
				UA1/UA2, 1983

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\
 & + i \bar{\Psi} \not{D} \Psi + h.c. \\
 & + \bar{\Psi} i \gamma_{ij} \Psi \phi + h.c. \\
 & + \frac{1}{2} \partial_{\mu} \phi^2 - V(\phi)
 \end{aligned}$$

Standard Model (Periodic Table for Elementary Particles)

	SLAC/BNL, 1974	CDF/D0, 1995 E288, 1977	Pluto, 1978 Jade/Pluto/Petra, 1979	ATLAS/CMS, 2012
mass →	≈2.3 MeV/c ²	≈1.275 GeV/c ²	≈173.07 GeV/c ²	≈126 GeV/c ²
charge →	2/3	2/3	2/3	0
spin →	1/2	1/2	1/2	0
QUARKS	u up	c charm	t top	H Higgs boson
	≈4.8 MeV/c ²	≈95 MeV/c ²	≈4.18 GeV/c ²	0
	-1/3	-1/3	-1/3	0
	1/2	1/2	1/2	1
	d down	s strange	b bottom	γ photon
	0.511 MeV/c ²	105.7 MeV/c ²	1.777 GeV/c ²	91.2 GeV/c ²
	-1	-1	-1	0
	1/2	1/2	1/2	1
	e electron	μ muon	τ tau	Z Z boson
LEPTONS	<2.2 eV/c ²	<0.17 MeV/c ²	<15.5 MeV/c ²	80.4 GeV/c ²
	0	0	0	±1
	1/2	1/2	1/2	1
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson
				GAUGE BOSONS

1st evidence, SLAC, 1968

Villard, 1900
Rutherford, 1903

UA1/UA2, 1983

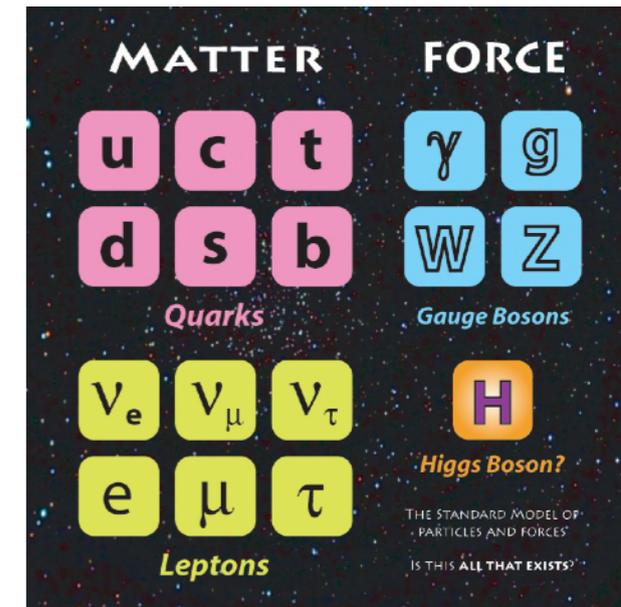
It is not just a table.

It contains profound mathematical formulation of our current understanding of the Universe.

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\
 & + i \bar{\Psi} \not{D} \Psi + h.c. \\
 & + \bar{\Psi}_i \gamma_{ij} \Psi_j \phi + h.c. \\
 & + \frac{1}{2} D_\mu \phi^2 - V(\phi)
 \end{aligned}$$

Current Particle Physics Model is INCOMPLETE

- Where are anti-particles in our universe?
- Where is Gravity in the table?
- Why there are three generations?
- Why are there three forces (besides gravity)?
- Why isn't only one kind of force (Grand Unified Theory)?
- Why do we need Higgs particle?
- Why some particles are heavy and some are light?
- Dark matter / Dark Energy
- How natural is this theory?
 - Naturalness; coincidence; fine tuning; multi-verse



del is



- Why isn't only one kind of force (Grand Unified Theory)?
- Why do we need Higgs particle?
- Why some particles are heavy and some are light?
- Dark matter / Dark Energy
- How natural is this theory?
 - Naturalness; coincidence; fine tuning; multi-verse

Erin

del is

MATTER FORCE
u c t γ g



- Why isn't only one kind of force?
- Why do we need Higgs particles?
- Why some particles are heavy?
- Dark matter / Dark Energy
- How natural is this theory?
 - Naturalness; coincidence;



Current Particle Physics Model is INCOMPLETE

- Where are anti-particles in our universe?
- Where is Gravity in the table?
- Why there are three generations?
- Why are there three forces (besides gravity)?
- Why isn't only one kind of force (Grand Unified Theory)?
- Why do we need Higgs particle?
- Why some particles are heavy and some are light?
- Dark matter / Dark Energy
- How natural is this theory?
 - Naturalness; coincidence; fine tuning; multi-verse



Higgs mass = bare mass - counter terms

Now we perform various experiments and among many other exciting experiments.....

Now we perform various experiments and among many other exciting experiments.....

LARGE HADRON COLLIDER

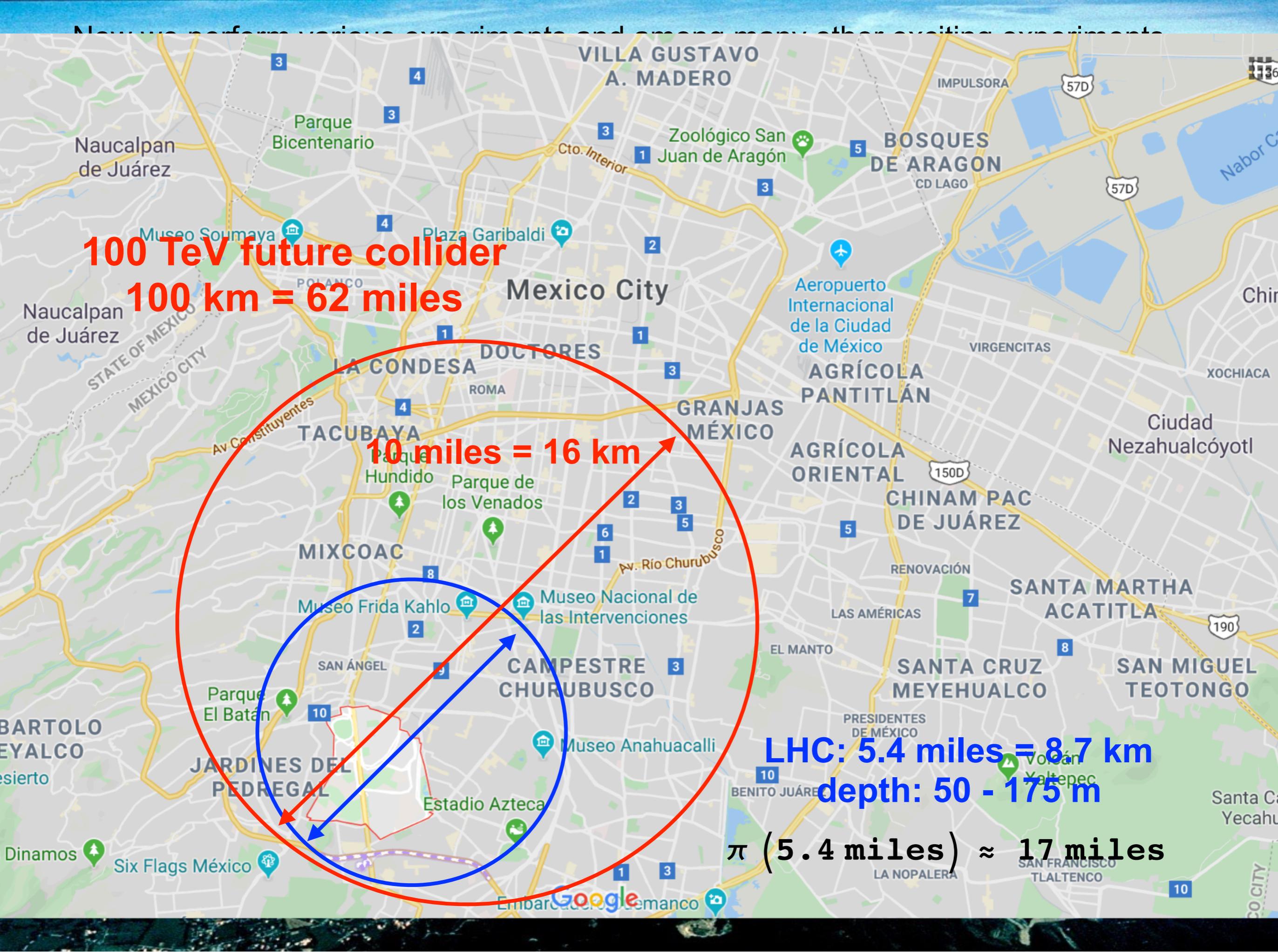
- Located at CERN, near Geneva
- Hadron: protons and nuclei
- Large: 17 miles of vacuum and superconducting magnets
- Accelerates protons to 99.9999999% the speed of light, 10,000 round trips per second
- Proton beams squeezed down to 64 microns
- 100 million proton-proton collisions per second
- More than 5000 Ph.D. physicists from > 100 countries
- ~\$10 billion project
- Conceived 1984, approved 1994
- Beams in Sept 2008, run I begun 2009, run II 2015@13TeV

100 TeV future collider
100 km = 62 miles

10 miles = 16 km

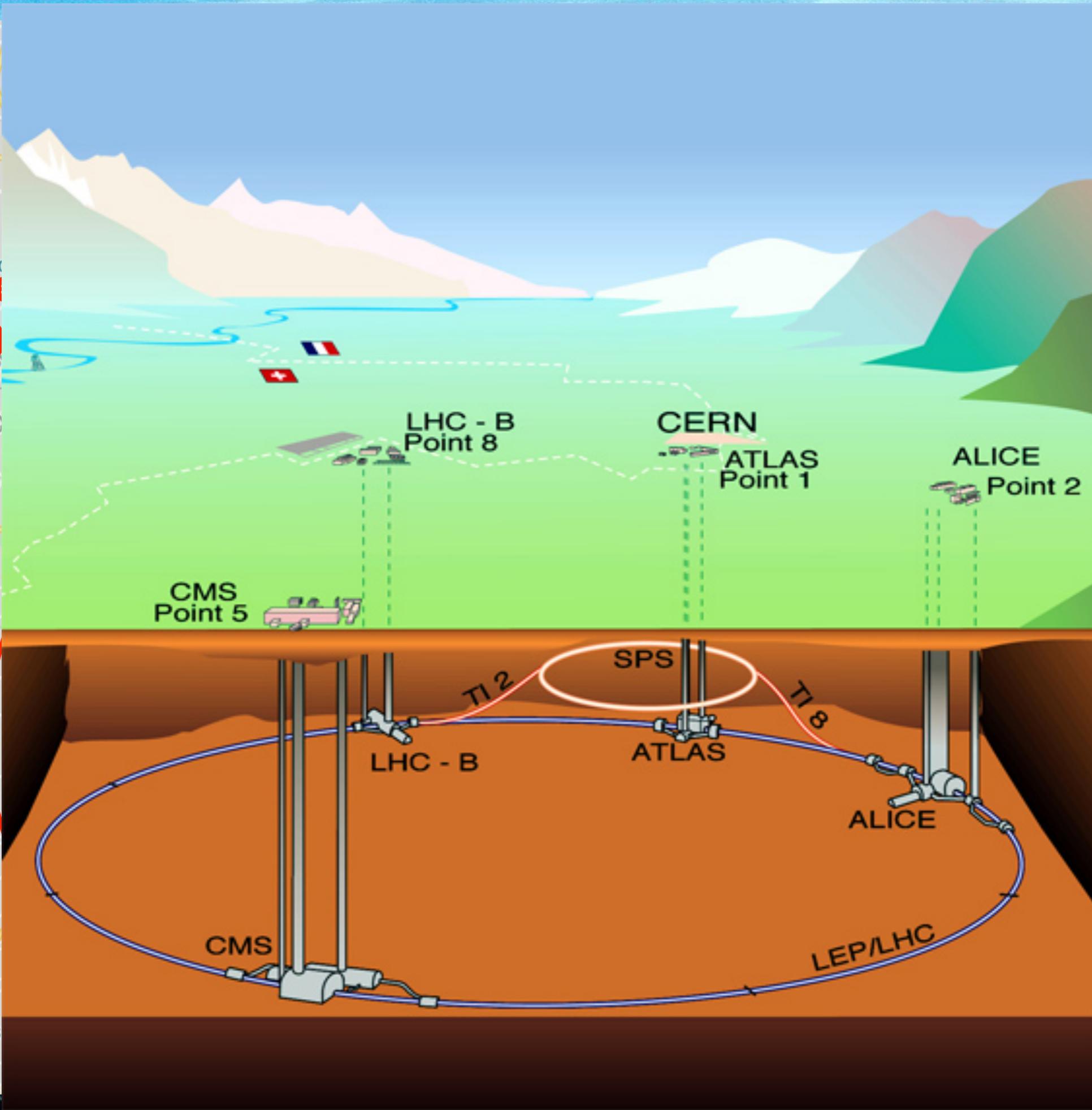
LHC: 5.4 miles = 8.7 km
depth: 50 - 175 m

$\pi (5.4 \text{ miles}) \approx 17 \text{ miles}$





100 Tera
100



km

iles

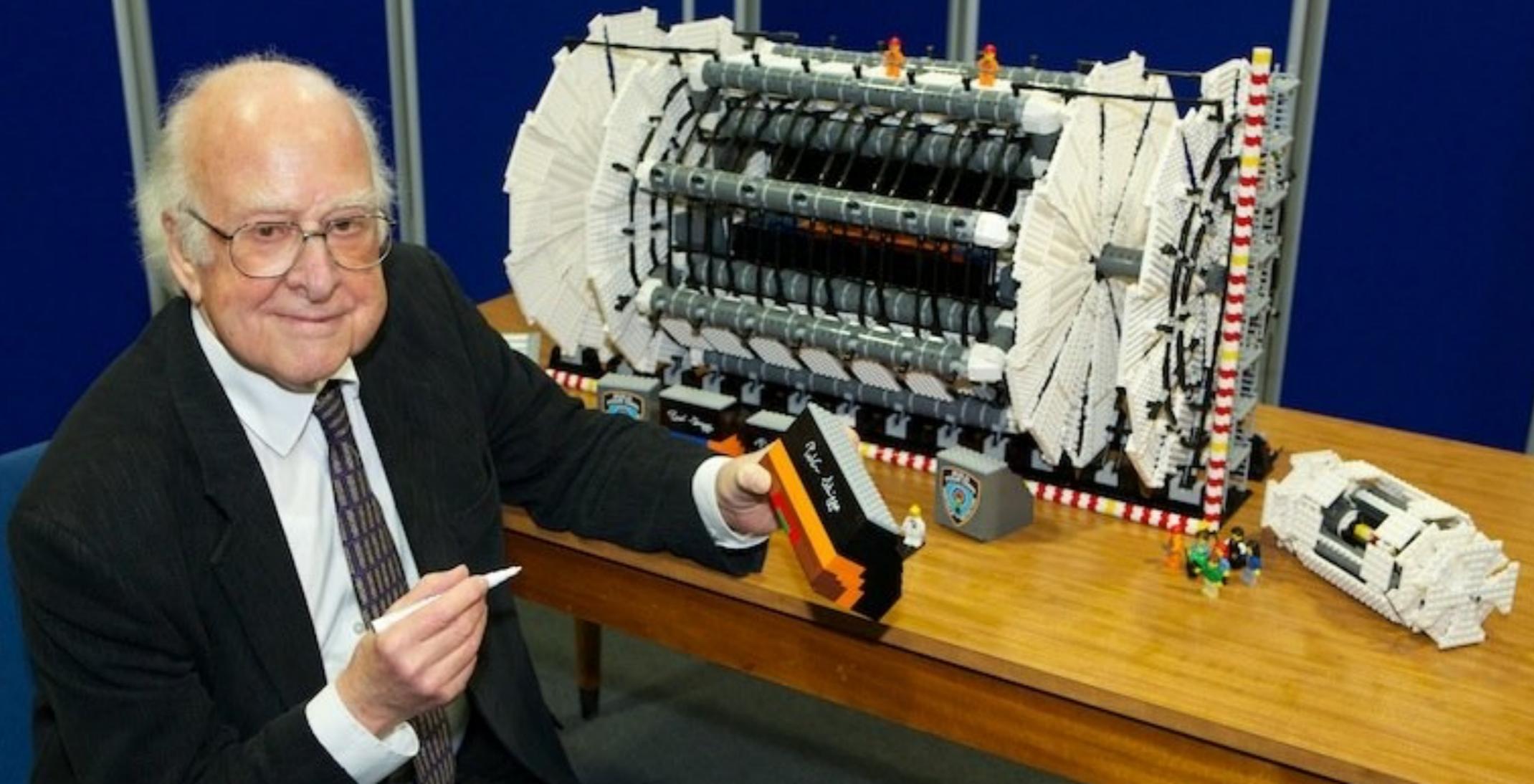
10

Object	Weight (tons)
Boeing 747 [fully loaded]	200
Endeavor space shuttle	368
ATLAS	7,000
Eiffel Tower	7,300
USS John McCain	8,300
CMS	12,500



- about 9500 Lego pieces
- roughly 50:1 in scale (close to scale with the LEGO man)
- material cost of about 2000 Euros
- about 1 m x 0.5 m x 0.5 m in size

Peter Higgs playing with ATLAS lego model, Nobel prize winner in 2013

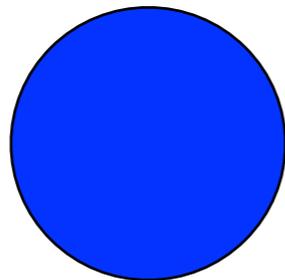




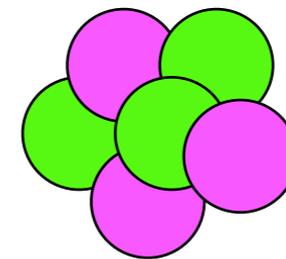
1:48 scale

MICROSCOPES

Higher energies \rightarrow shorter wavelengths



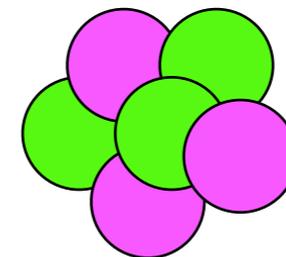
low resolution



high resolution

MICROSCOPES

Higher energies \rightarrow shorter wavelengths



high resolution

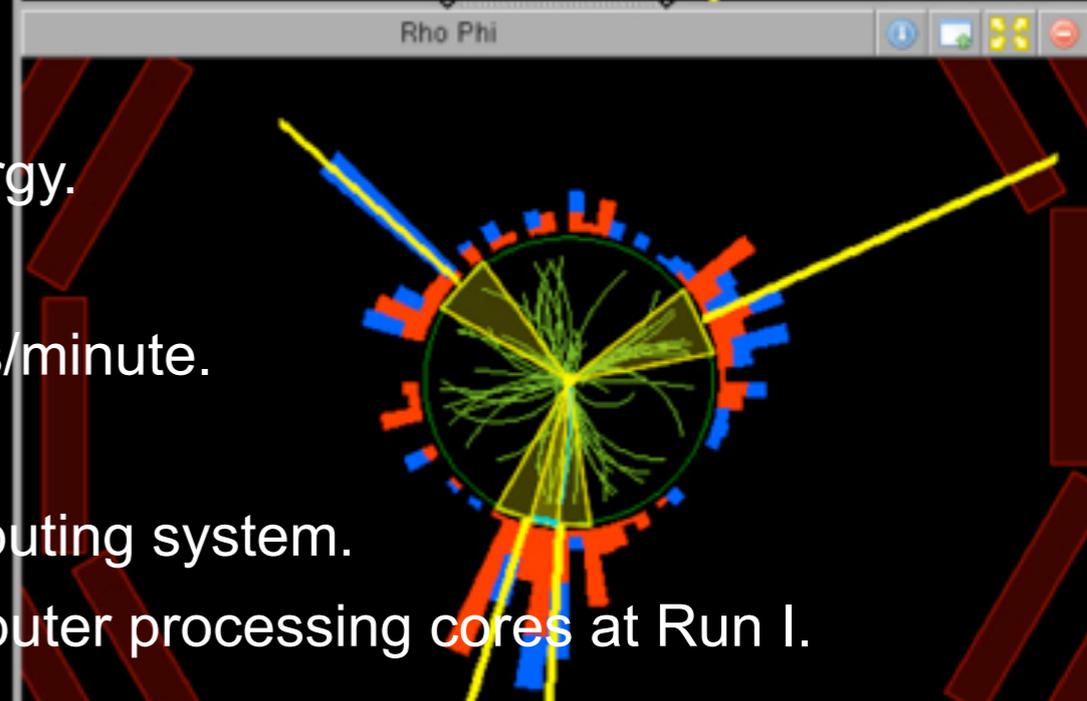
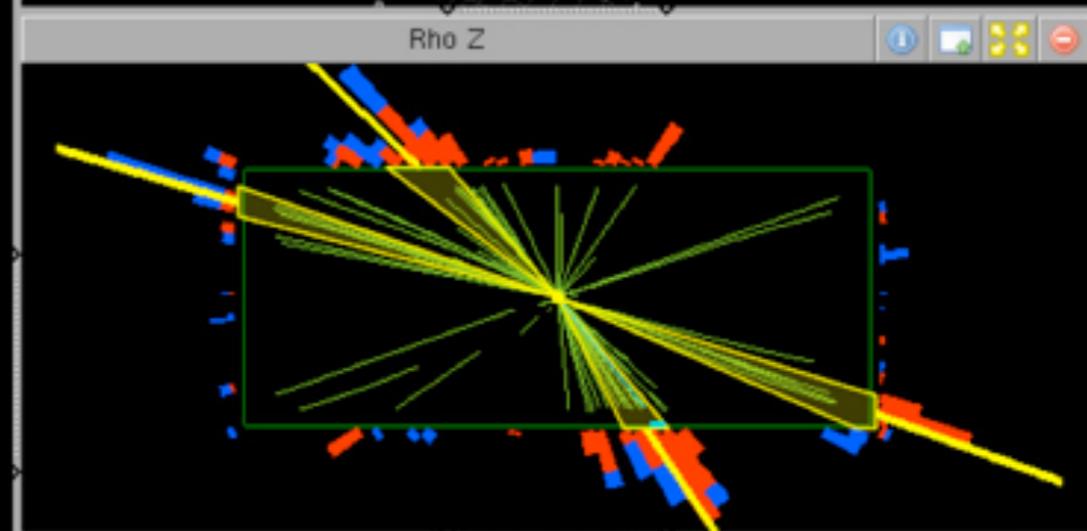
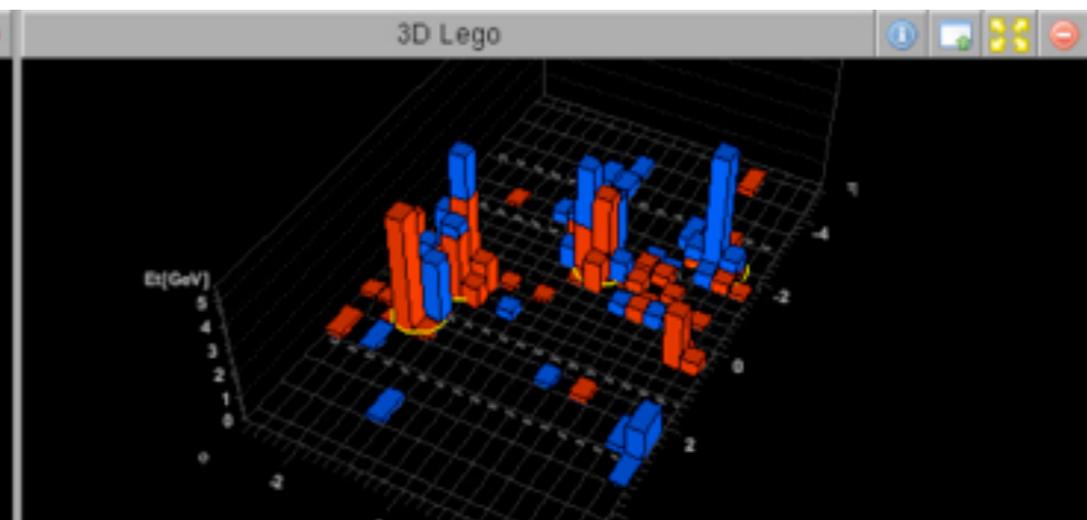
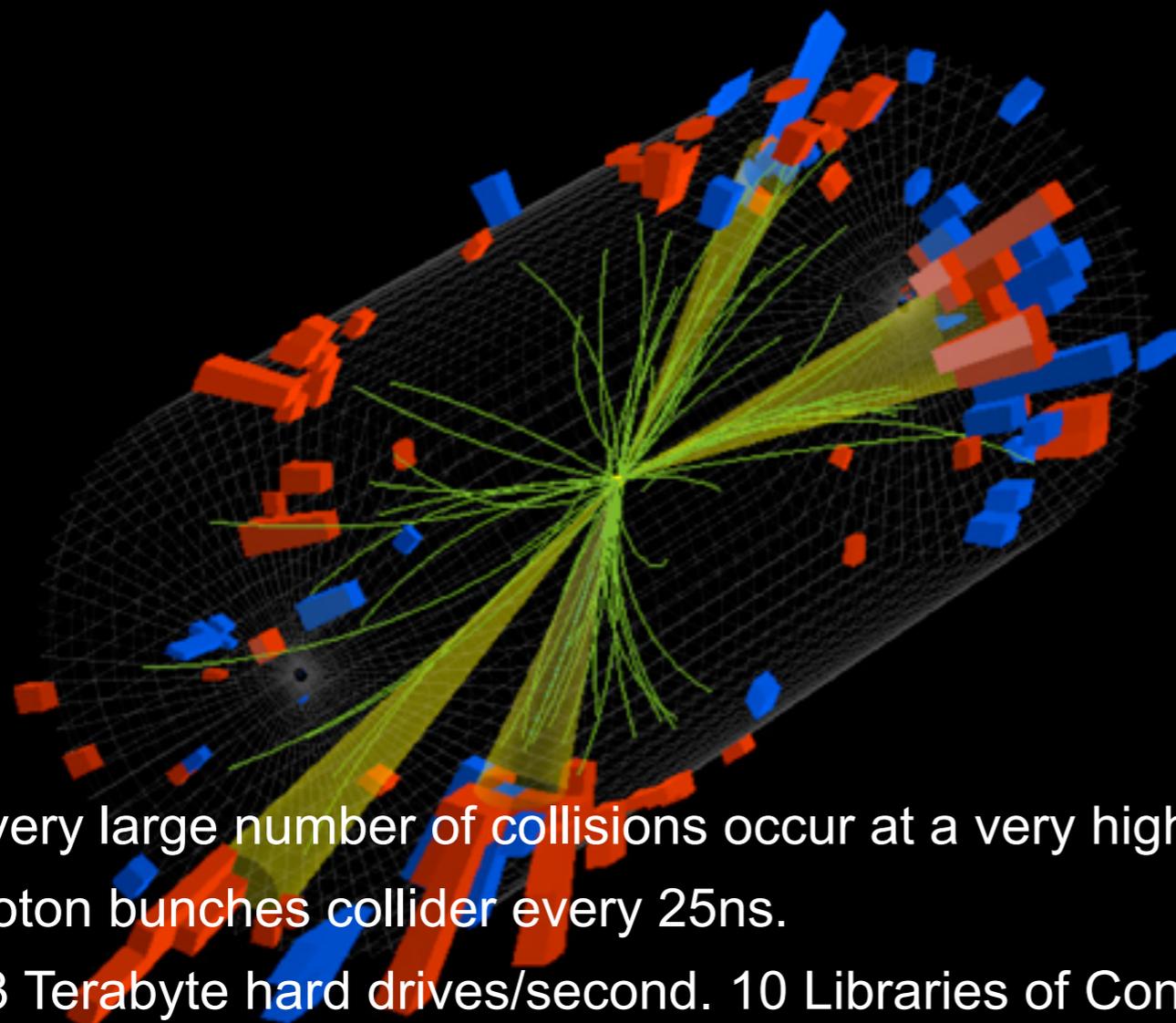


CMS Experiment at the LHC, CERN

Date Recorded: 2009-12-14

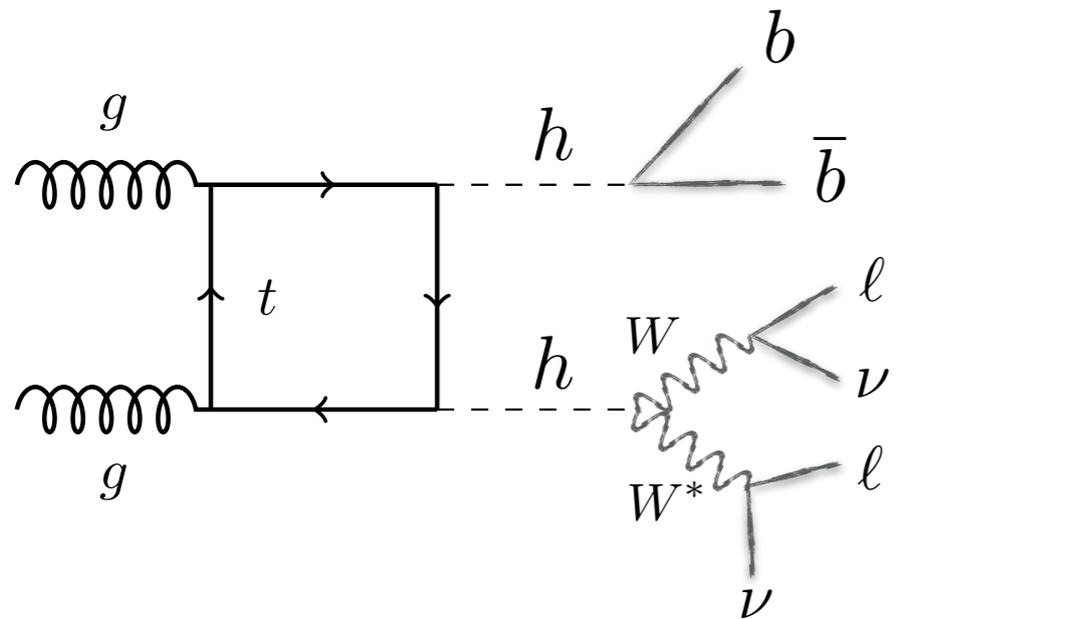
Run/Event: 124120/6613074

Candidate Multijet Event at 2.36 TeV

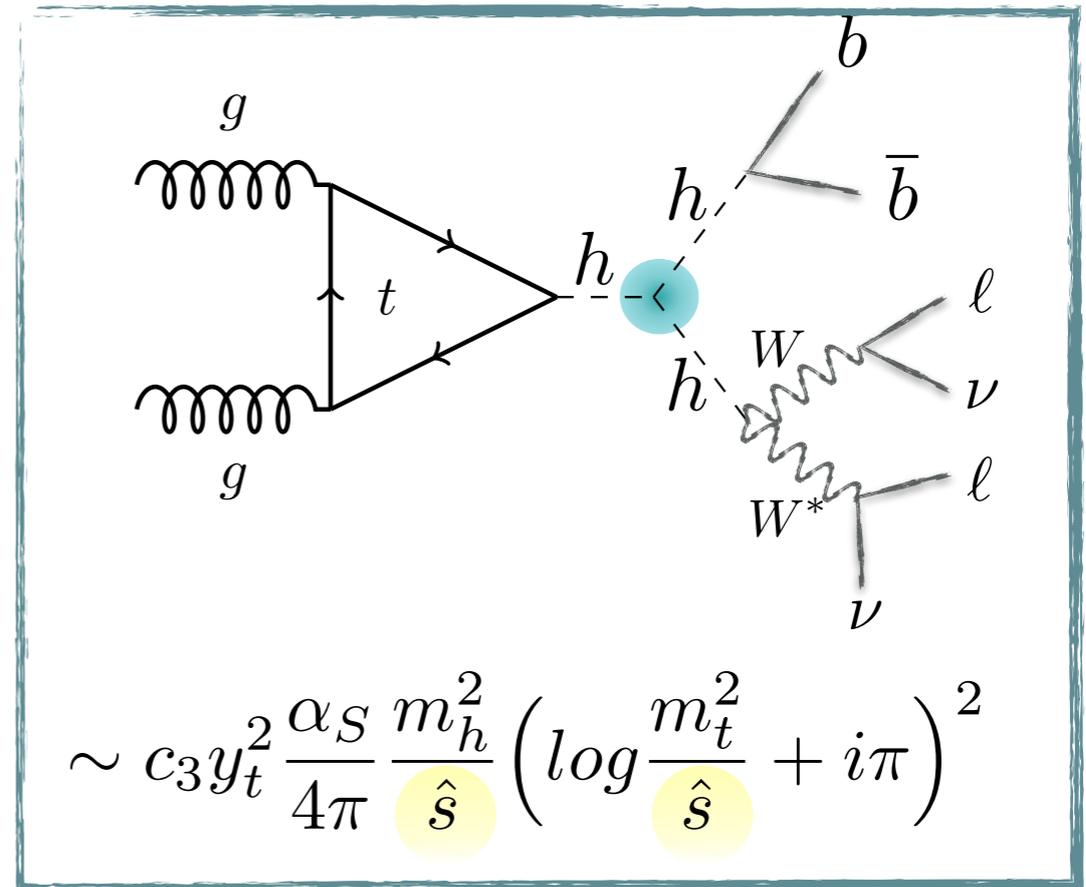


- A very large number of collisions occur at a very high energy.
- Proton bunches collide every 25ns.
- 3.3 Terabyte hard drives/second. 10 Libraries of Congress/minute.
- 100 full length DVD movies/second.
- Data analysis requires full use of the worldwide Grid computing system.
- Worldwide LHC Computing Grid used up to 485,000 computer processing cores at Run I.

Why double Higgs (hh) ?



+

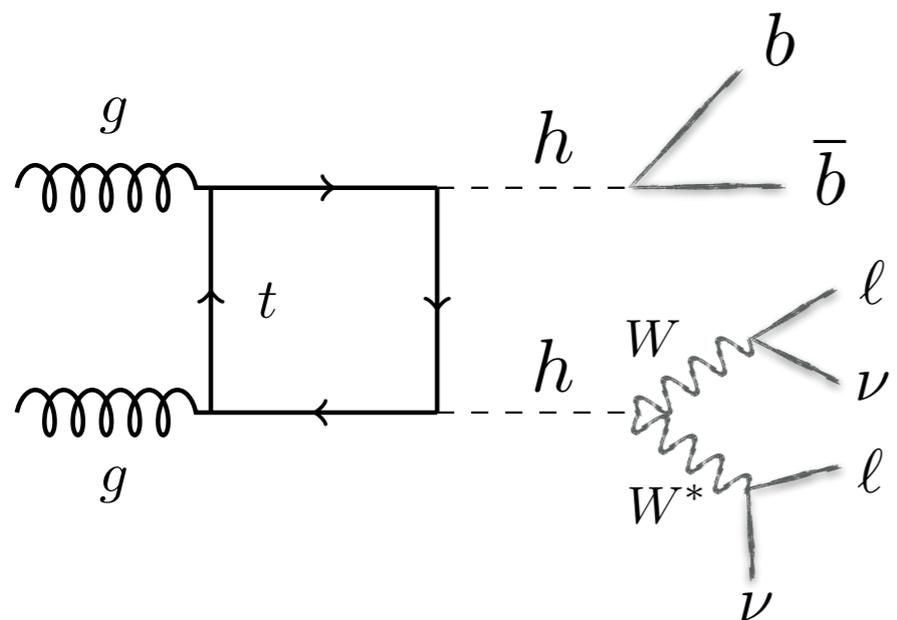


$$V_h = \frac{m_h^2}{2} h^2 + c_3 \frac{m_h^2}{2v} h^3 + c_4 \frac{m_h^2}{8v^2} h^4$$

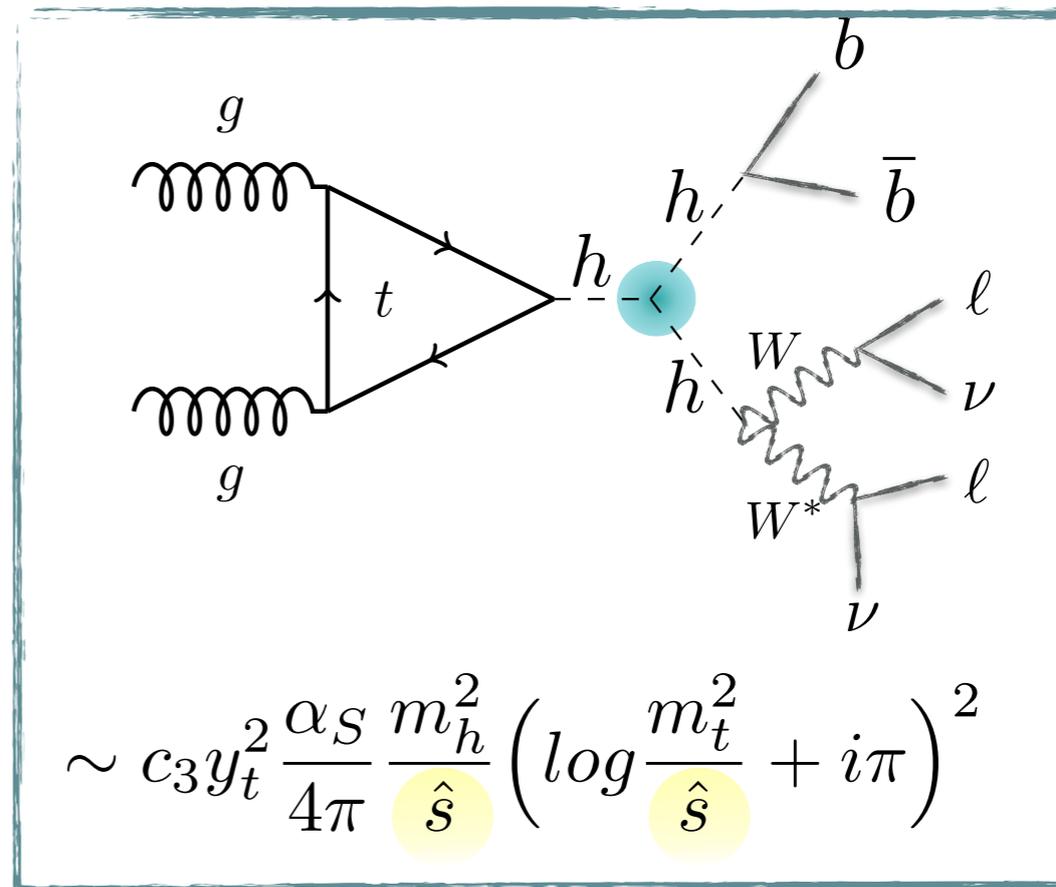
$$\sim c_3 y_t^2 \frac{\alpha_S}{4\pi} \frac{m_h^2}{\hat{s}} \left(\log \frac{m_t^2}{\hat{s}} + i\pi \right)^2$$

- We will focus on the triple Higgs self-coupling (c_3), with all other couplings set to their SM values. c_3 is accessible at the HL-LHC but c_4 needs 100 TeV collider.
- The knowledge of c_3 is crucial to reconstruct the Higgs potential for better understanding EWSB. Many BSM scenarios allow large deviations for c_3 .
- The c_3 is sensitive at lower-energy bins where the backgrounds are large.
- That's why probing c_3 is notoriously difficult.

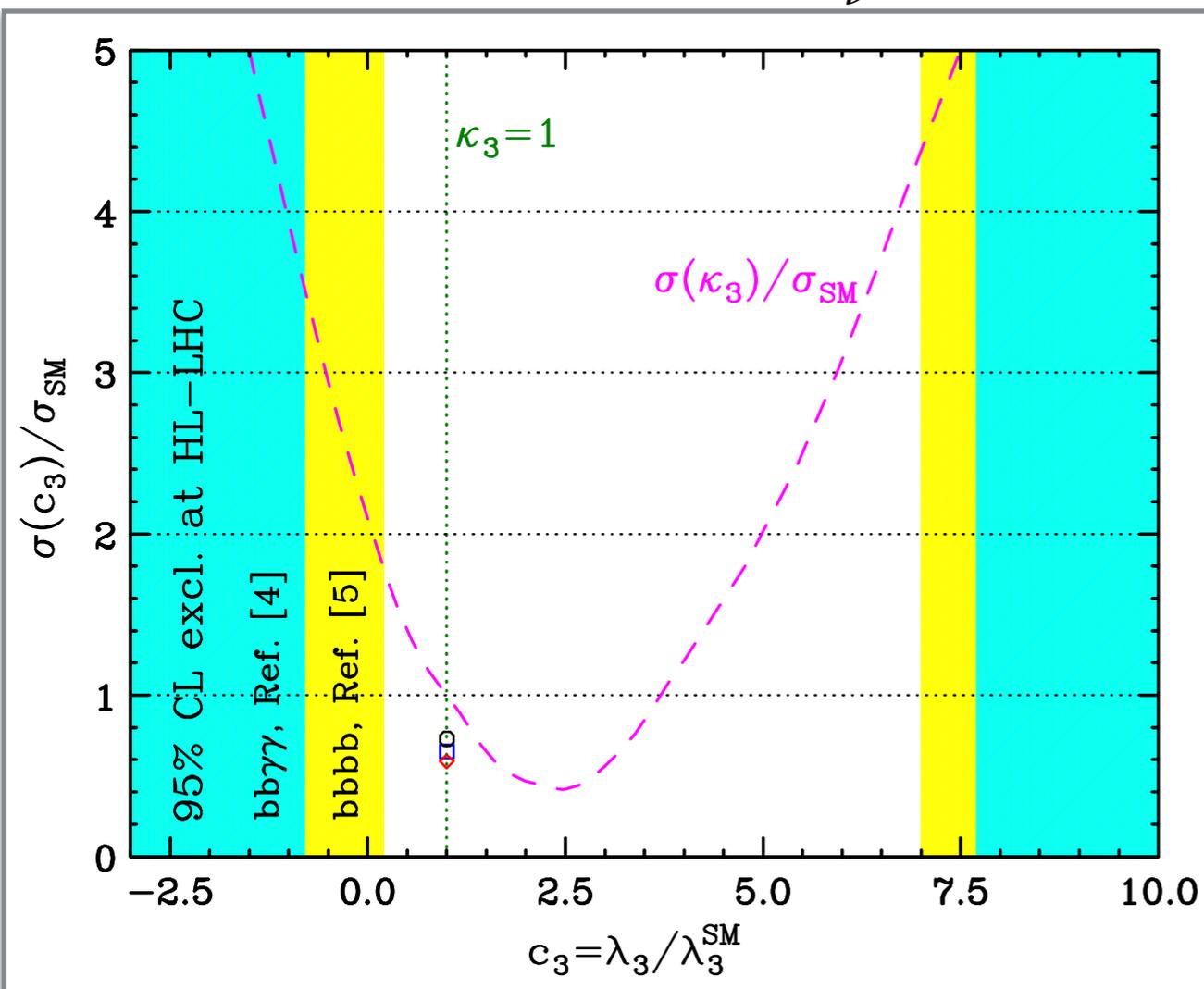
Why double Higgs (hh) ?



+



$$\sim c_3 y_t^2 \frac{\alpha_S}{4\pi} \frac{m_h^2}{\hat{s}} \left(\log \frac{m_t^2}{\hat{s}} + i\pi \right)^2$$



c_3 (c_3), with all other couplings set to their SM values, needs 100 TeV collider.

the Higgs potential for better understanding EWSB. or c_3 .

the backgrounds are large.

Decays

$$1 \text{ fb} = 10^{-39} \text{ cm}^2$$

$$\sigma(hh)_{SM}^{NNLO} \simeq 40.7 \text{ fb} \quad (14 \text{ TeV})$$

higher branching ratios

cleaner final state

	bb	WW^*	$\tau\tau$	ZZ^*	$\gamma\gamma$
bb	33%				
WW^*	25%	4.6%			
$\tau\tau$	7.3%	2.7%	0.39%		
ZZ^*	3.1%	1.1%	0.33%	0.069%	
$\gamma\gamma$	0.26%	0.1%	0.028%	0.012%	0.0005%

1902.00134	Statistical-only		Statistical + Systematic	
	ATLAS	CMS	ATLAS	CMS
$HH \rightarrow bbbb$	1.4	1.2	0.61	0.95
$HH \rightarrow b\bar{b}\tau\tau$	2.5	1.6	2.1	1.4
$HH \rightarrow b\bar{b}\gamma\gamma$	2.1	1.8	2.0	1.8
$HH \rightarrow b\bar{b}VV(ll\nu\nu)$	-	0.59	-	0.56
$HH \rightarrow b\bar{b}ZZ(4l)$	-	0.37	-	0.37
combined	3.5	2.8	3.0	2.6
	Combined 4.5		Combined 4.0	

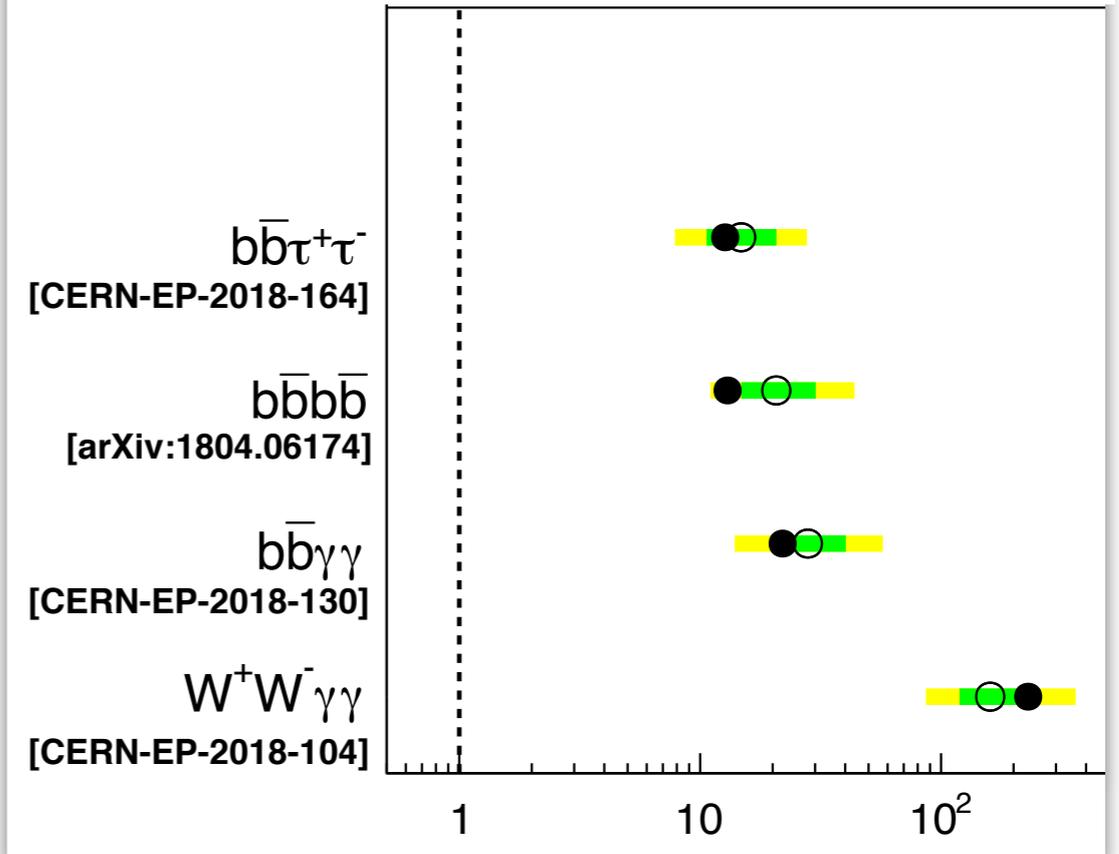
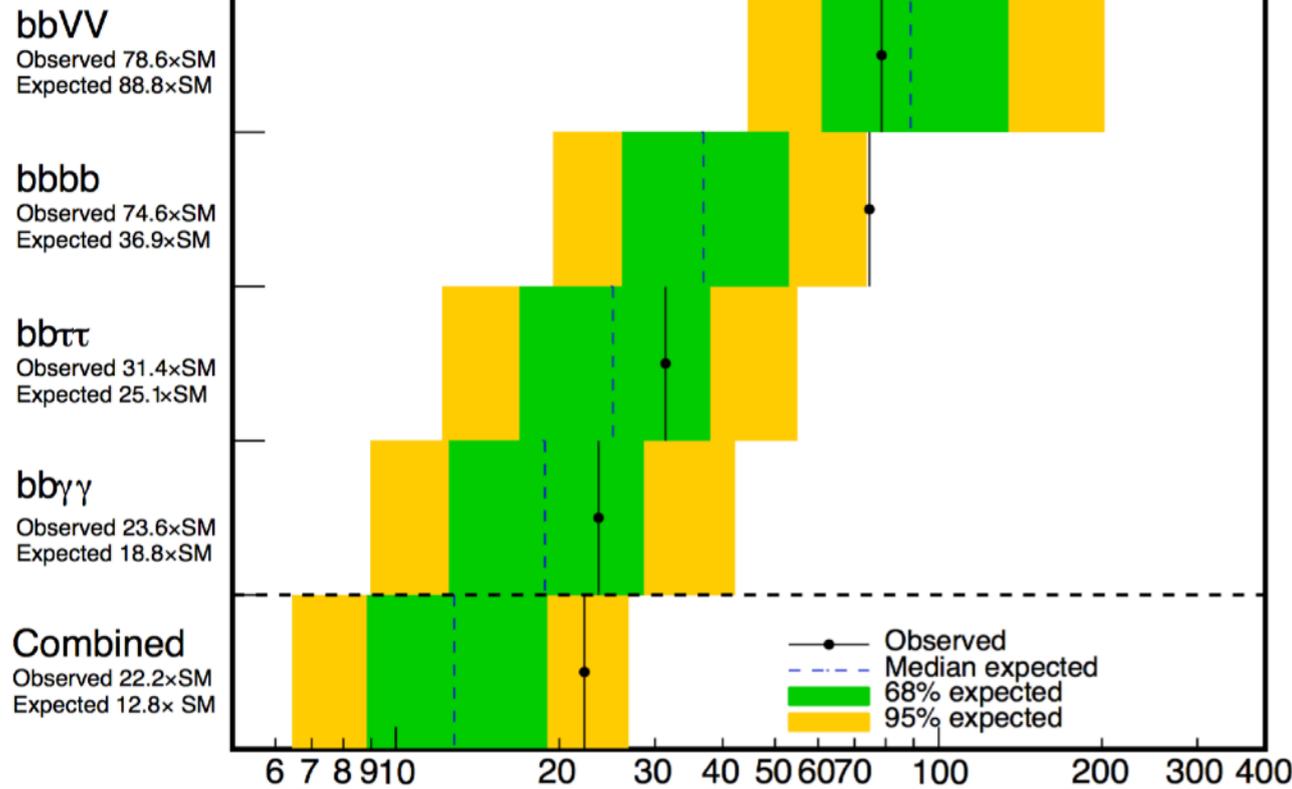
4 σ expected for ATLAS+CMS!

- These measurements are challenged by a low $\sigma(hh)$ and small branching ratios (BR).
- No single channel is expected to reach 5 sigma at HL-LHC.
- The combination of different channels is crucial. $bbWW$ has good potential for further improvement.

$$\sigma_{hh}^{observed} / \sigma_{hh}^{SM} = 22$$

$$\sigma_{hh}^{observed} / \sigma_{hh}^{SM} = 13$$

CMS preliminary $gg \rightarrow HH$ 35.9 fb⁻¹ (13 TeV)



95% CL on $\sigma_{hh} / \sigma_{hh}^{SM}$

1902.00134

	Statistical only		Statistical + Systematic	
	ATLAS	CMS	ATLAS	CMS
$HH \rightarrow bbbb$	1.4	1.2	0.61	0.95
$HH \rightarrow bb\tau\tau$	2.5	1.6	2.1	1.4
$HH \rightarrow bb\gamma\gamma$	2.1	1.8	2.0	1.8
$HH \rightarrow bbVV (ll\nu\nu)$	-	0.59	-	0.56
$HH \rightarrow bbZZ (4l)$	-	0.37	-	0.37
combined	3.5	2.8	3.0	2.6
	Combined		Combined	
	4.5		4.0	

4σ expected for ATLAS+CMS!

- These measurements are challenged by a low $\sigma (hh)$ and small branching ratios (BR).
- No single channel is expected to reach 5 sigma at HL-LHC.
- The combination of different channels is crucial. bbWW has good potential for further improvement.

Cross sections @ 14 TeV LHC

$\sigma_{hh} = 40.7 \text{ fb. (NNLO)}$ (only two studies in the dilepton channel)

$$\sigma_{hh} \cdot 2 \cdot \text{BR}(h \rightarrow b\bar{b}) \cdot \text{BR}(h \rightarrow WW^* \rightarrow \ell^+ \ell^- \nu \bar{\nu}) = 0.648 \text{ fb}$$

ℓ denotes an electron or a muon, including leptons from tau decays.

- tt: 953.6 pb (NNLO)
- tth: 611.3 fb (NLO)
- ttV (V=W,Z): 1.71 pb (NLO)
- DY: $k_{\text{QCD} \otimes \text{QED}}^{\text{NNLO, DY}} \approx 1$
- Irreducible jjllnunu: $k_{\text{NLO}} = 2$
- tWj: 0.51 pb (after cuts, including all relevant branching fractions)

$$\sigma_{\text{bkn d}} \sim 10^5 \sigma_{hh}$$

	Signal	$t\bar{t}$	$t\bar{t}h$	$t\bar{t}V$	$llbj$	$\tau\tau bb$	$tw + j$	$jjll\nu\nu$	σ	S/B
Baseline cuts: $p_T > 20 \text{ GeV}$, $p_{T,\ell} > 20 \text{ GeV}$, $\Delta R_{\ell\ell} < 1.0$, $p_{T,b} > 30 \text{ GeV}$, $\Delta R_{bb} < 1.3$, $m_{\ell\ell} < 65 \text{ GeV}$, $95 < m_{bb} < 140 \text{ GeV}$	0.01046	1.8855	0.0269	0.0179	0.0697	0.0250	0.2209	0.0113	0.38	0.0046

cross section in fb

tt: 84% tW: 9.8% DY+jets: 3.1% tth: 1.2% tautau + bb: 1.1% ttV: 0.8%

$hh \rightarrow bbWW^*$: dilepton channel

HL-LHC, 14 TeV, $L=3 \text{ ab}^{-1}$

Adhikary, Banerjee, Barman, Bhattacharjee, Niyogi JHEP 2017

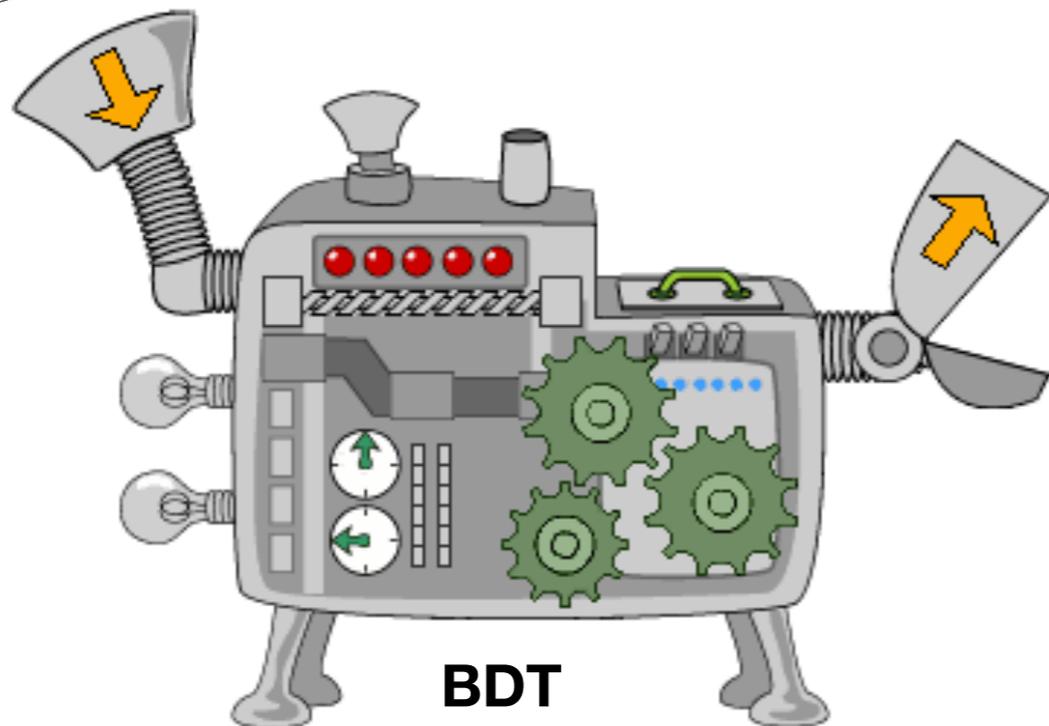
\cancel{E}_T

$p_{T,\ell_{1/2}} \Delta\phi_{bb\ell\ell}$

$\Delta R_{\ell\ell} \Delta R_{bb}$

$m_{bb} m_{\ell\ell}$

$p_{T,bb} p_{T,\ell\ell}$



BDT

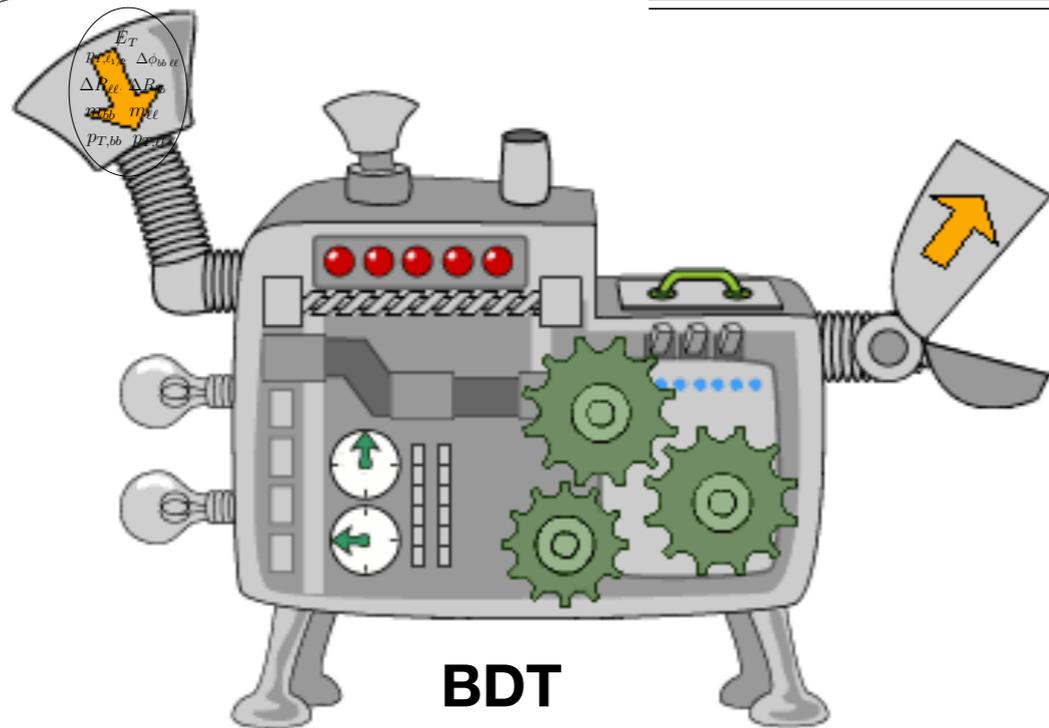
$hh \rightarrow bbWW^*$: dilepton channel

HL-LHC, 14 TeV, $L=3 \text{ ab}^{-1}$

Adhikary, Banerjee, Barman, Bhattacharjee, Niyogi JHEP 2017

- \cancel{E}_T
- $p_{T,\ell_{1/2}} \Delta\phi_{bb\ell\ell}$
- $\Delta R_{\ell\ell} \Delta R_{bb}$
- $m_{bb} m_{\ell\ell}$
- $p_{T,bb} p_{T,\ell\ell}$

Sl. No.	Process	Order	Events
Background	$t\bar{t} \text{ lep}$	NNLO [128]	2080.52
	$t\bar{t}h$	NLO [111]	131.66
	$t\bar{t}Z$	NLO [130]	106.31
	$t\bar{t}W$	NLO [129]	35.97
	$hb\bar{b}$	NNLO (5FS) + NLO (4FS) [111]	~ 0
	$\ell\ell b\bar{b}$	LO	842.72
	Total		3197.18
Signal ($hh \rightarrow b\bar{b}WW \rightarrow b\bar{b}\ell\ell + \cancel{E}_T$)		NNLO [70]	35.20
Significance (S/\sqrt{B})			0.62



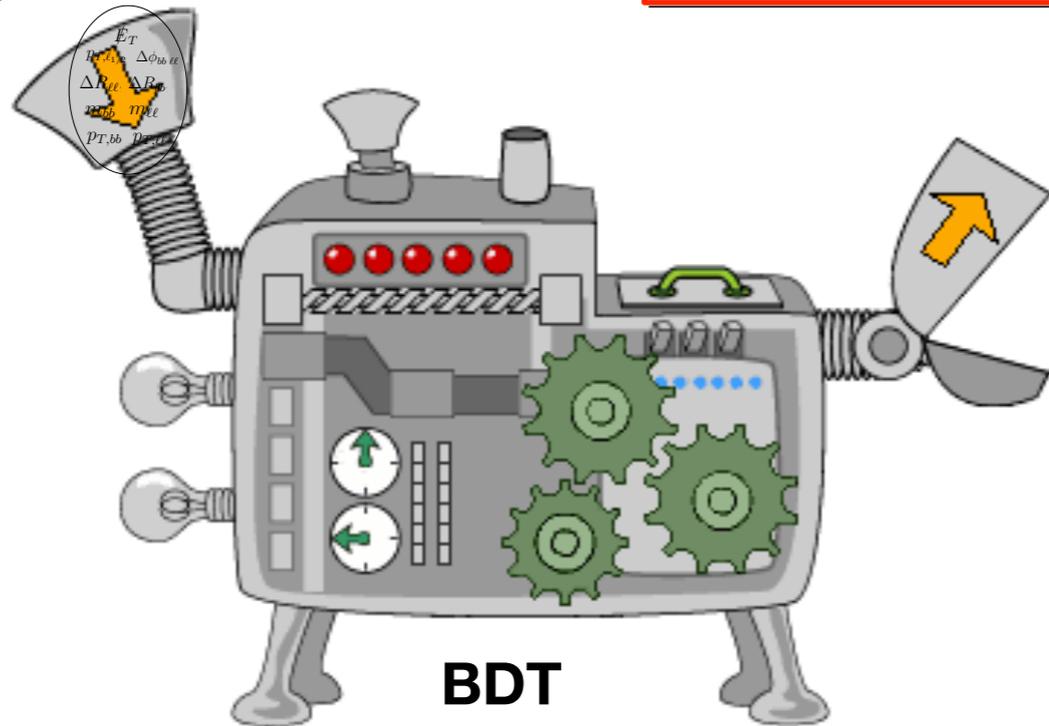
$hh \rightarrow bbWW^*$: dilepton channel

HL-LHC, 14 TeV, $L=3 \text{ ab}^{-1}$

Adhikary, Banerjee, Barman, Bhattacharjee, Niyogi JHEP 2017

\cancel{E}_T
 $p_{T,\ell_{1/2}} \Delta\phi_{bb\ell\ell}$
 $\Delta R_{\ell\ell} \Delta R_{bb}$
 $m_{bb} m_{\ell\ell}$
 $p_{T,bb} p_{T,\ell\ell}$

Sl. No.	Process	Order	Events
Background	$t\bar{t} \text{ lep}$	NNLO [128]	2080.52
	$t\bar{t}h$	NLO [111]	131.66
	$t\bar{t}Z$	NLO [130]	106.31
	$t\bar{t}W$	NLO [129]	35.97
	$hb\bar{b}$	NNLO (5FS) + NLO (4FS) [111]	~ 0
	$\ell\ell b\bar{b}$	LO	842.72
	Total		3197.18
	Signal ($hh \rightarrow b\bar{b}WW \rightarrow b\bar{b}\ell\ell + \cancel{E}_T$)	NNLO [70]	35.20
Significance (S/\sqrt{B})			0.62



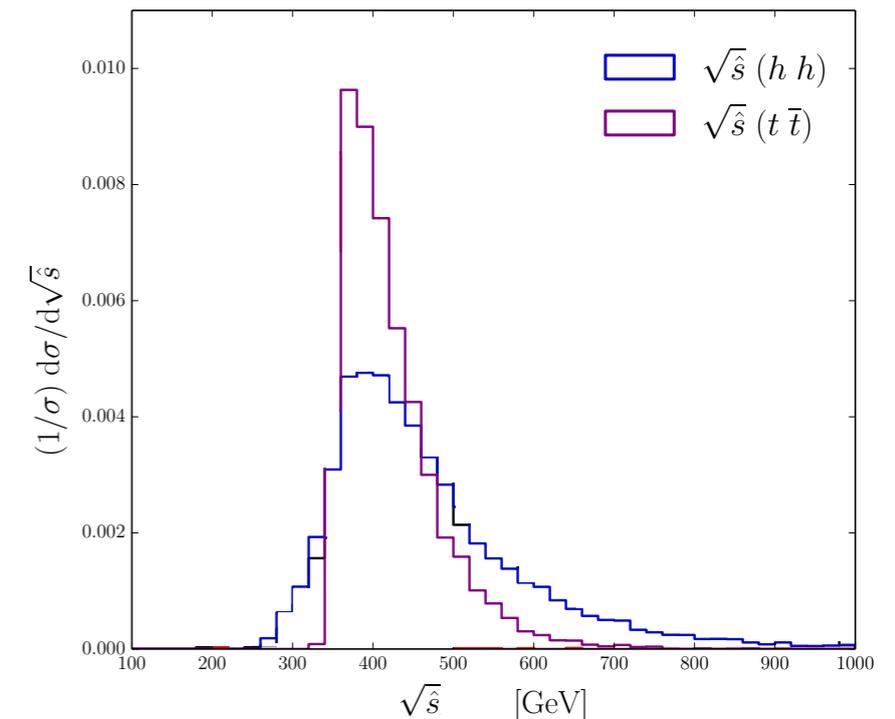
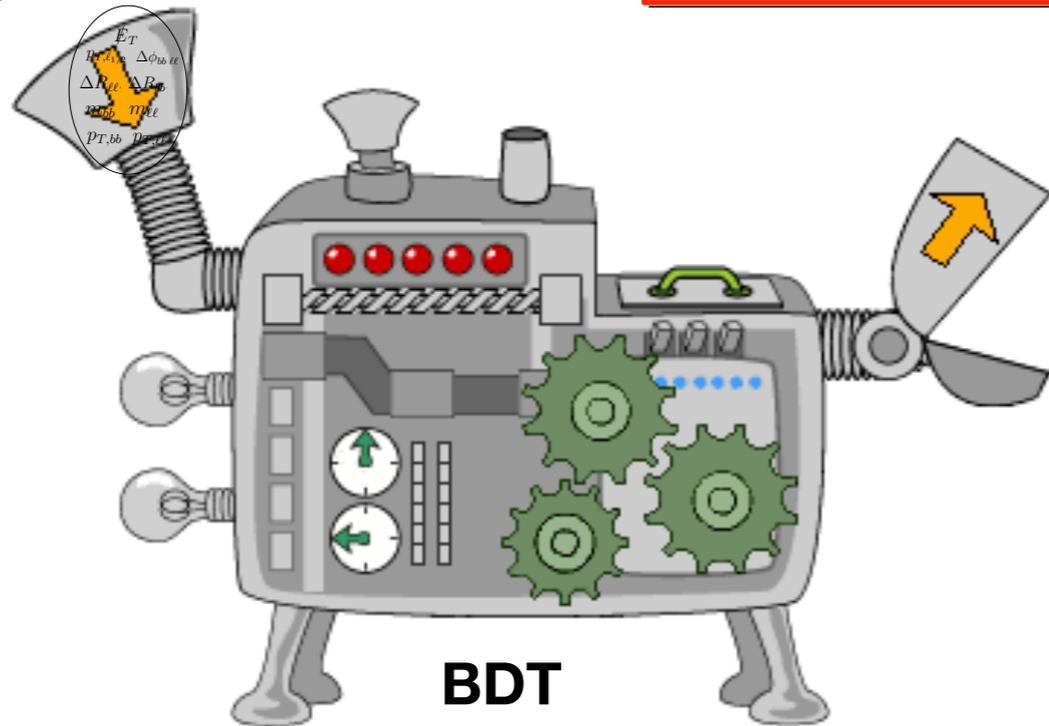
$hh \rightarrow bbWW^*$: dilepton channel

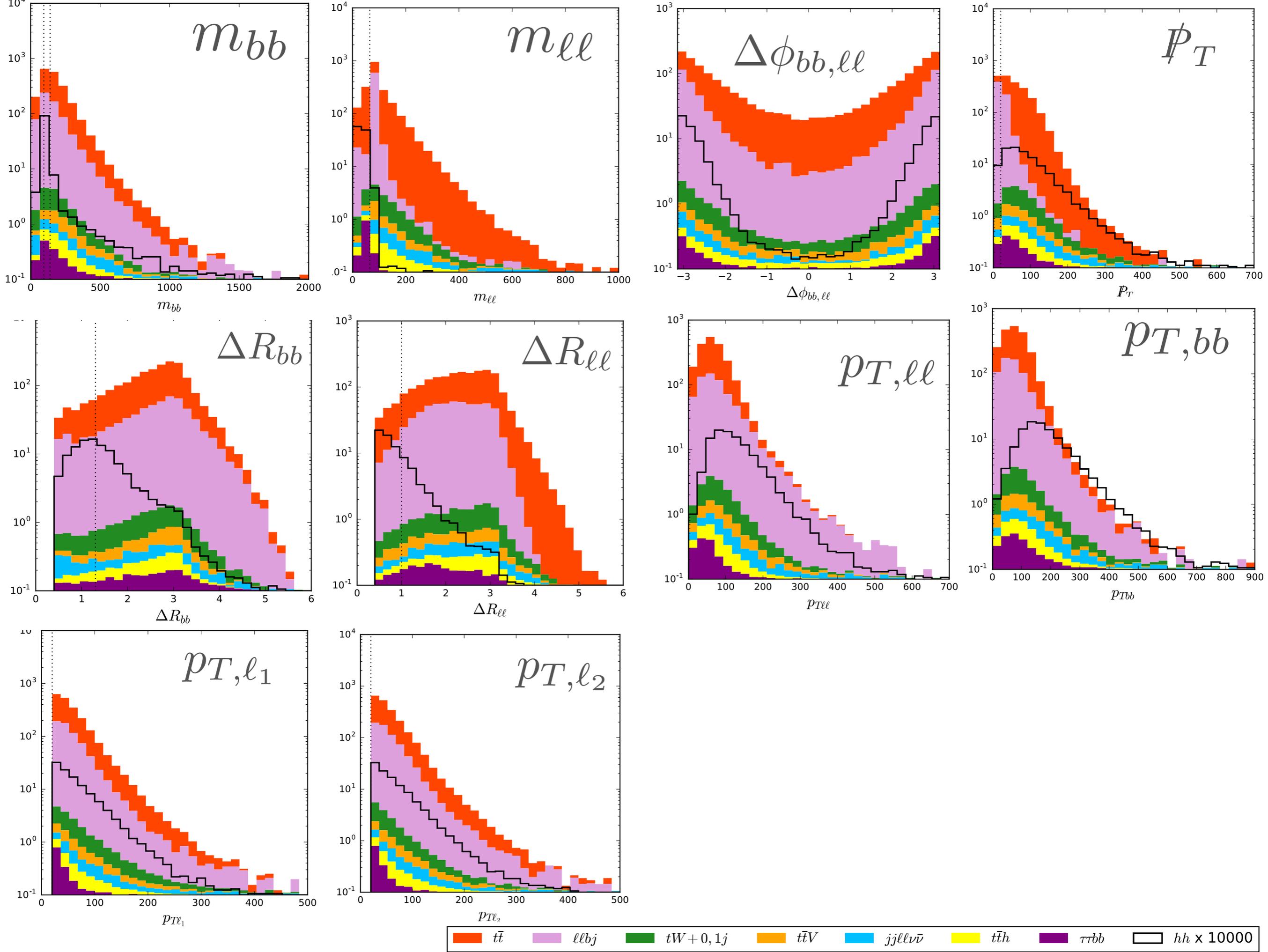
HL-LHC, 14 TeV, $L=3 \text{ ab}^{-1}$

Adhikary, Banerjee, Barman, Bhattacharjee, Niyogi JHEP 2017

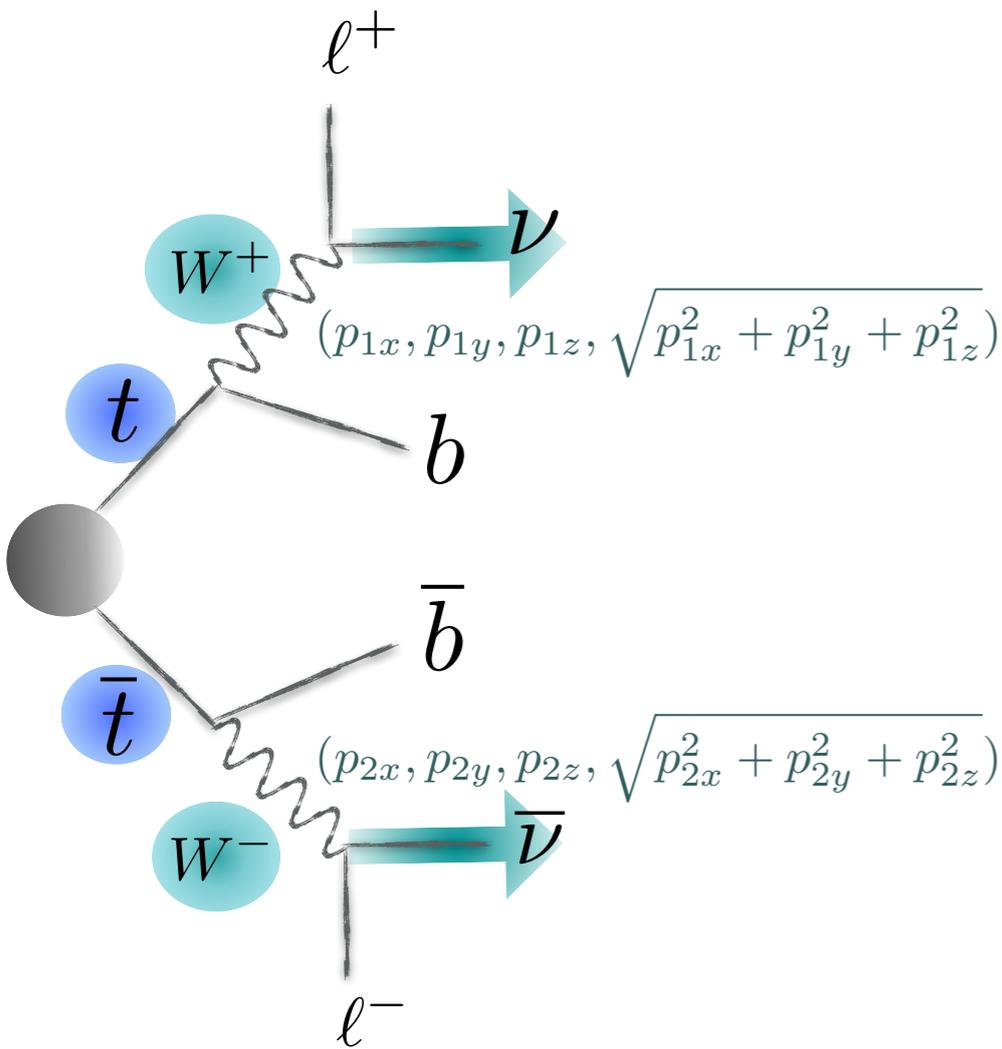
- \cancel{E}_T
- $p_{T,l_{1/2}} \Delta\phi_{bb \ell\ell}$
- $\Delta R_{\ell\ell} \Delta R_{bb}$
- $m_{bb} m_{\ell\ell}$
- $p_{T,bb} p_{T,\ell\ell}$

Sl. No.	Process	Order	Events
Background	$t\bar{t} \text{ lep}$	NNLO [128]	2080.52
	$t\bar{t}h$	NLO [111]	131.66
	$t\bar{t}Z$	NLO [130]	106.31
	$t\bar{t}W$	NLO [129]	35.97
	$hb\bar{b}$	NNLO (5FS) + NLO (4FS) [111]	~ 0
	$\ell\ell b\bar{b}$	LO	842.72
	Total		3197.18
	Signal ($hh \rightarrow b\bar{b}WW \rightarrow b\bar{b}\ell\ell + \cancel{E}_T$)	NNLO [70]	35.20
Significance (S/\sqrt{B})			0.62





Topness (T)



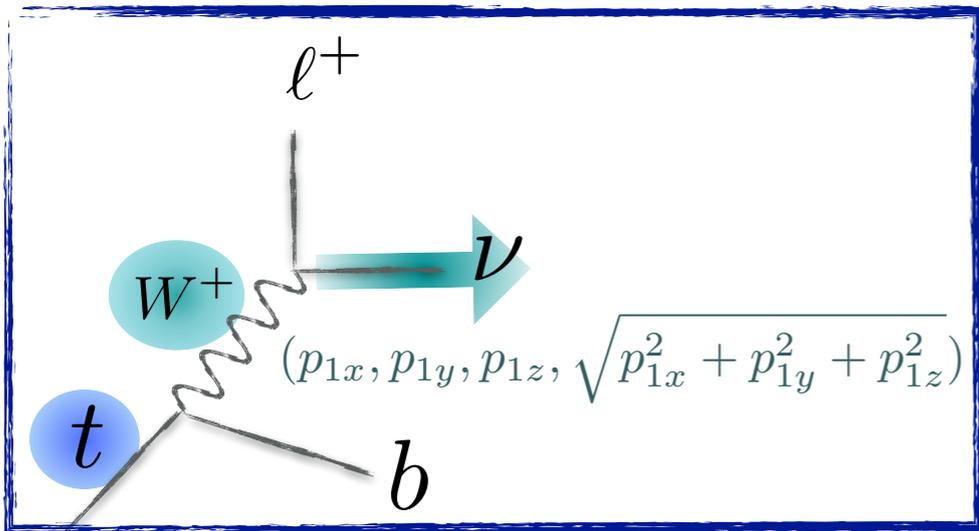
$$\chi_{ij}^2 \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} \left[\frac{(m_{b_i \ell^+ \nu}^2 - m_t^2)^2}{\sigma_t^4} + \frac{(m_{\ell^+ \nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{b_j \ell^- \bar{\nu}}^2 - m_t^2)^2}{\sigma_t^4} + \frac{(m_{\ell^- \bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} \right]$$

$$T \equiv \min(\chi_{12}^2, \chi_{21}^2)$$

two possible ways of pairing b and ℓ

- Topness provides a degree of consistency to dileptonic $t\bar{t}$ production.
- It scans over 6 unknowns of neutrino momenta with four on-shell masses and missing E_T constraints.
- And find a minimum of the likelihood function.

Topness (T)



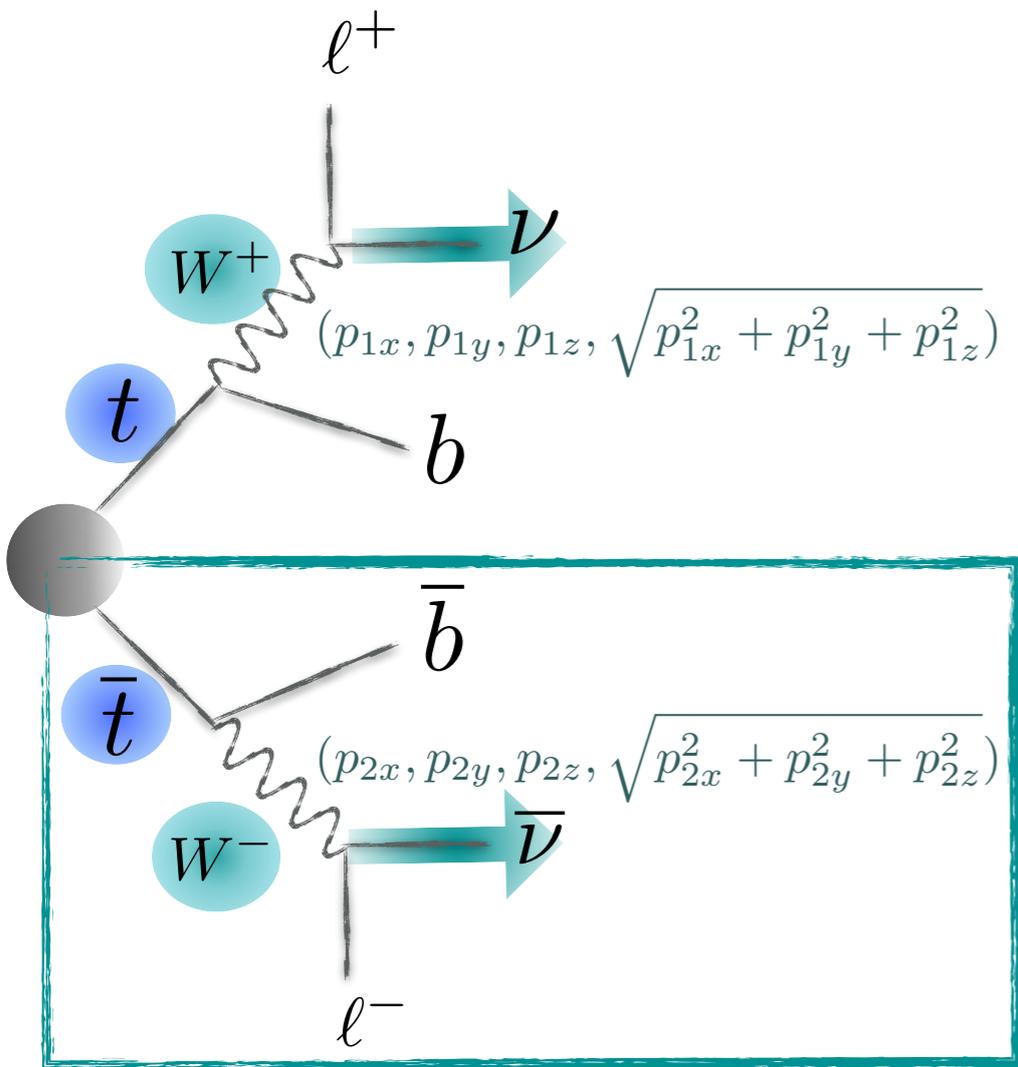
$$\chi_{ij}^2 \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} \left[\frac{(m_{b_i \ell^+ \nu}^2 - m_t^2)^2}{\sigma_t^4} + \frac{(m_{\ell^+ \nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{b_j \ell^- \bar{\nu}}^2 - m_t^2)^2}{\sigma_t^4} + \frac{(m_{\ell^- \bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} \right]$$

$$T \equiv \min(\chi_{12}^2, \chi_{21}^2)$$

two possible ways of pairing b and ℓ

- Topness provides a degree of consistency to dileptonic $t\bar{t}$ production.
- It scans over 6 unknowns of neutrino momenta with four on-shell masses and missing E_T constraints.
- And find a minimum of the likelihood function.

Topness (T)



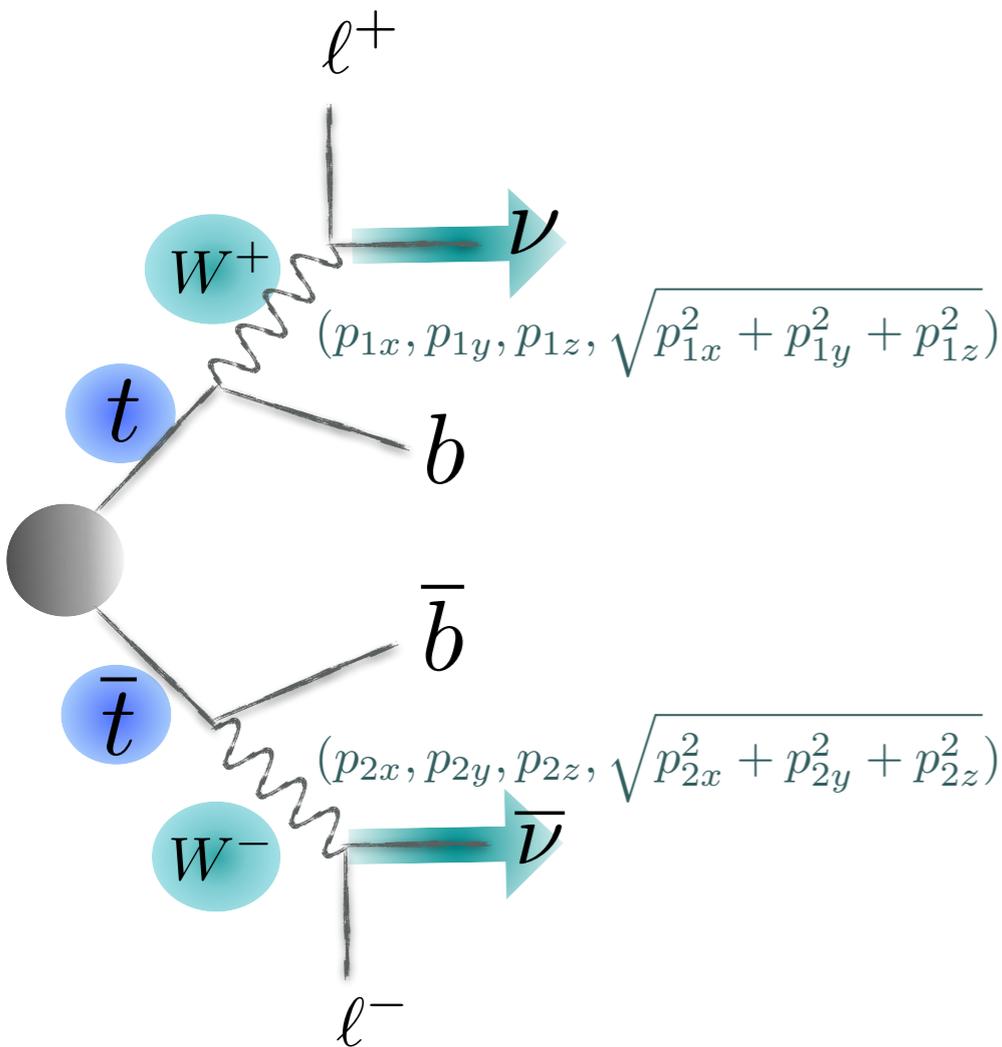
$$\chi_{ij}^2 \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} \left[\frac{\left(m_{b_i \ell^+ \nu}^2 - m_t^2 \right)^2}{\sigma_t^4} + \frac{\left(m_{\ell^+ \nu}^2 - m_W^2 \right)^2}{\sigma_W^4} \right. \\ \left. + \frac{\left(m_{b_j \ell^- \bar{\nu}}^2 - m_t^2 \right)^2}{\sigma_t^4} + \frac{\left(m_{\ell^- \bar{\nu}}^2 - m_W^2 \right)^2}{\sigma_W^4} \right]$$

$$T \equiv \min (\chi_{12}^2, \chi_{21}^2)$$

two possible ways of pairing b and ℓ

- Topness provides a degree of consistency to dileptonic $t\bar{t}$ production.
- It scans over 6 unknowns of neutrino momenta with four on-shell masses and missing E_T constraints.
- And find a minimum of the likelihood function.

Topness (T)



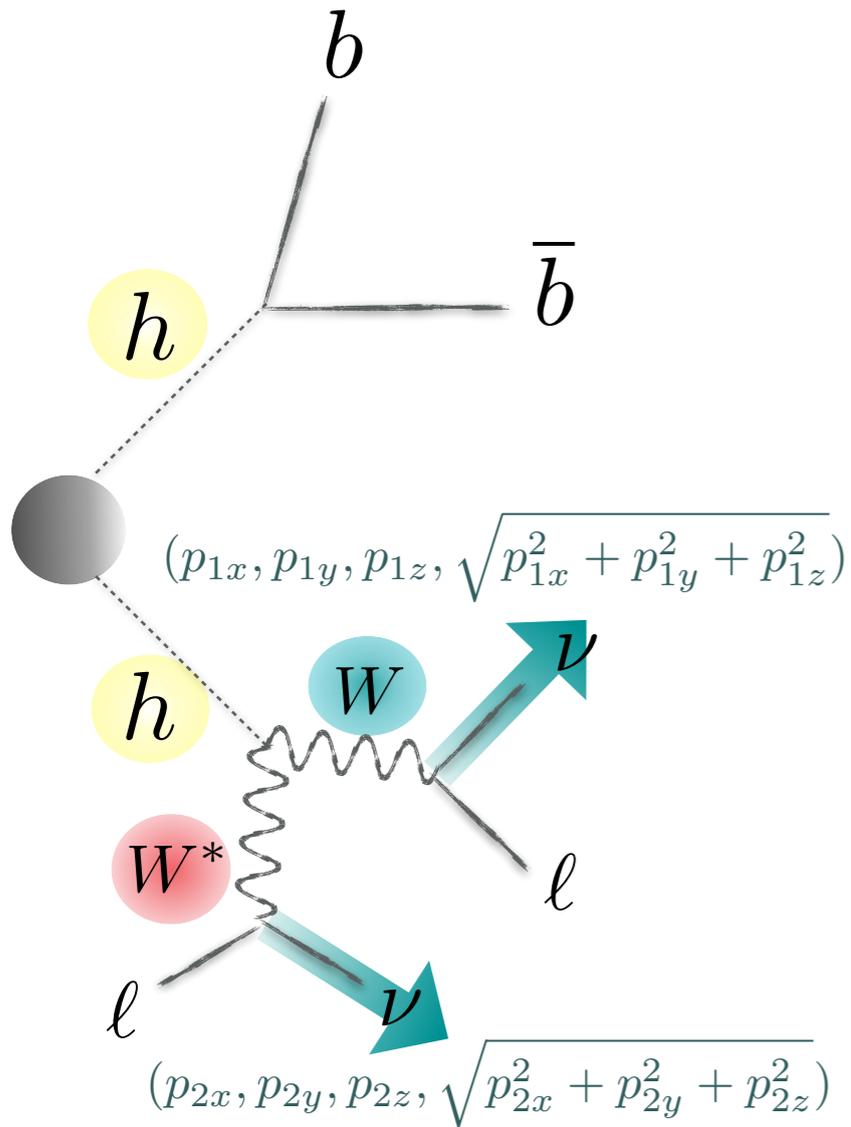
$$\chi_{ij}^2 \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} \left[\frac{\left(m_{b_i \ell^+ \nu}^2 - m_t^2 \right)^2}{\sigma_t^4} + \frac{\left(m_{\ell^+ \nu}^2 - m_W^2 \right)^2}{\sigma_W^4} + \frac{\left(m_{b_j \ell^- \bar{\nu}}^2 - m_t^2 \right)^2}{\sigma_t^4} + \frac{\left(m_{\ell^- \bar{\nu}}^2 - m_W^2 \right)^2}{\sigma_W^4} \right]$$

$$T \equiv \min (\chi_{12}^2, \chi_{21}^2)$$

two possible ways of pairing b and ℓ

- Topness provides a degree of consistency to dileptonic $t\bar{t}$ production.
- It scans over 6 unknowns of neutrino momenta with four on-shell masses and missing E_T constraints.
- And find a minimum of the likelihood function.

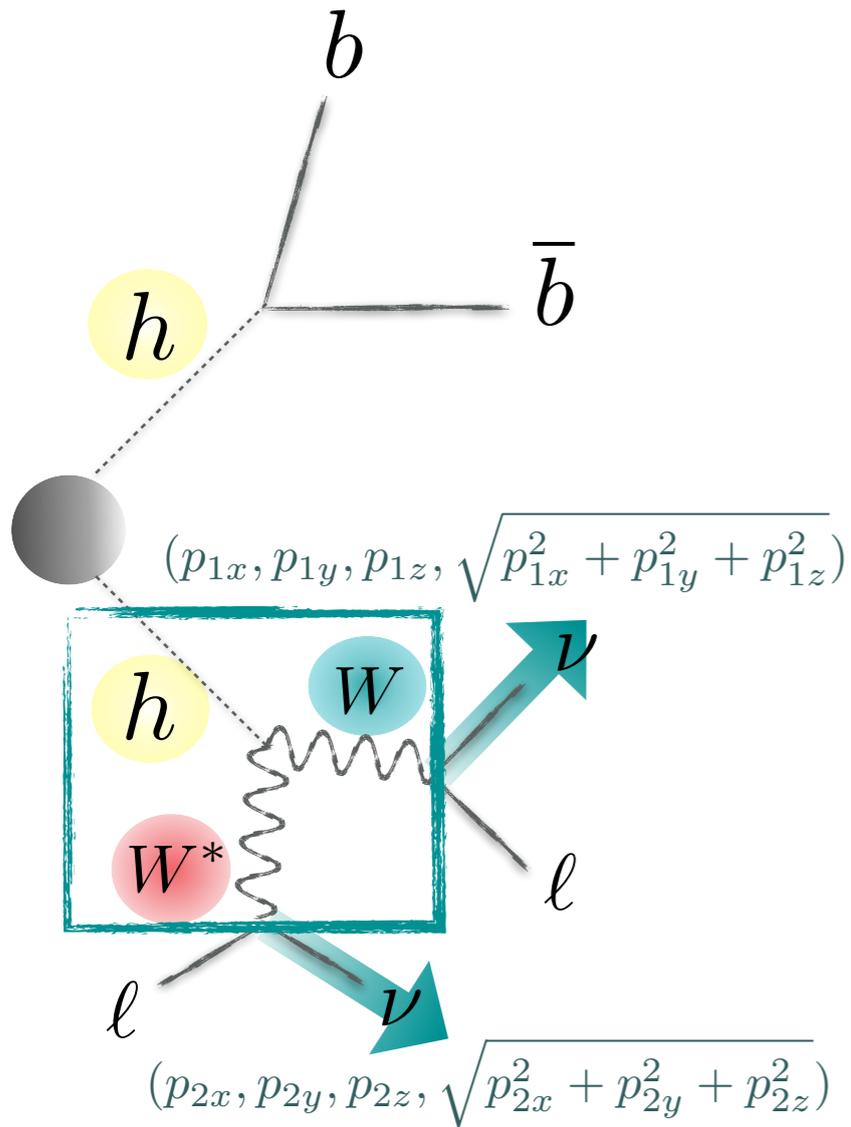
Higgsness (H)



$$\begin{aligned}
 H \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} & \left[\frac{(m_{\ell+\ell-\nu\bar{\nu}}^2 - m_h^2)^2}{\sigma_{h\ell}^4} + \frac{(m_{\nu\bar{\nu}}^2 - m_{\nu\bar{\nu},peak}^2)^2}{\sigma_{\nu}^4} \right. \\
 & + \min \left(\frac{(m_{\ell+\nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell-\bar{\nu}}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4}, \right. \\
 & \left. \left. \frac{(m_{\ell-\bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell+\nu}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4} \right) \right], \\
 & \sim m_h - m_W \text{ off-shell}
 \end{aligned}$$

- Higgsness provides a degree of consistency to dileptonic $h \rightarrow WW^*$ system.
- The off-shell W also has an end-point near $m_h - m_W$.
- Its distribution is wide, but there is a peak, which can constrain hh system further.

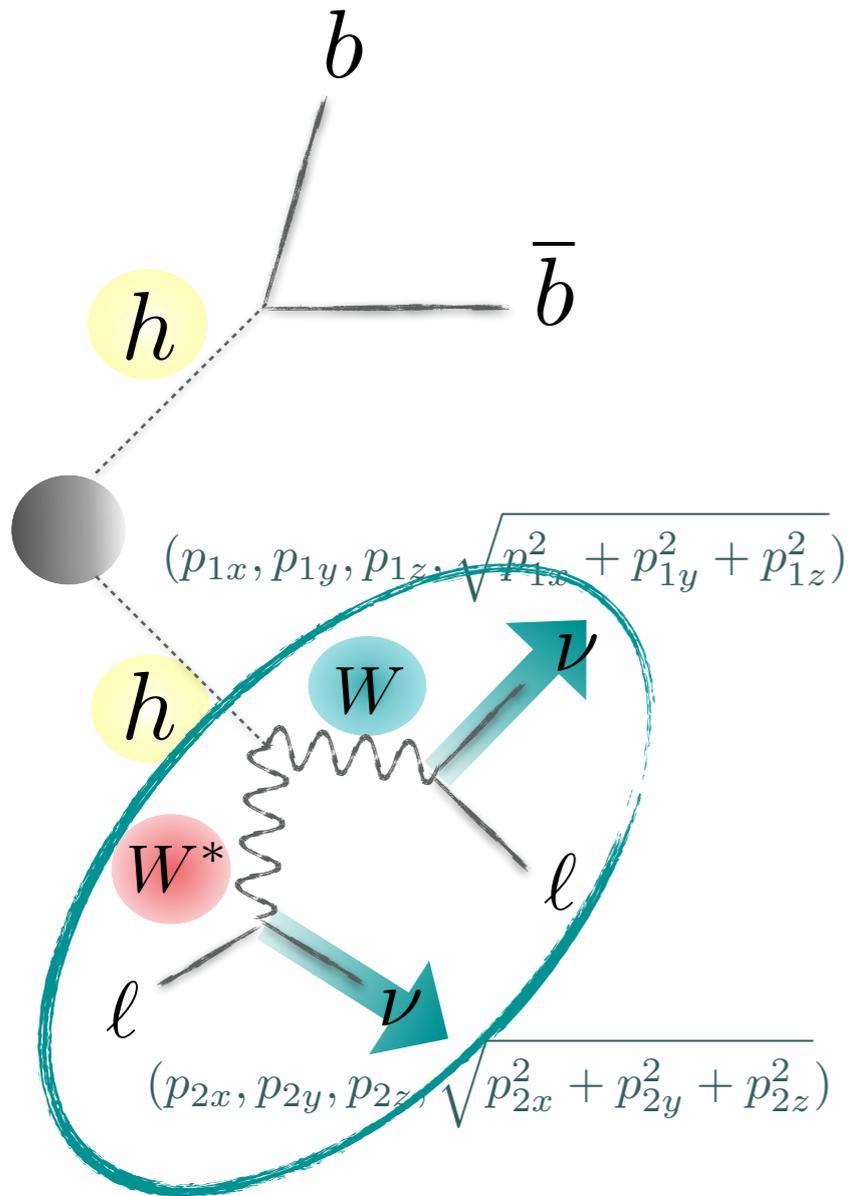
Higgsness (H)



$$\begin{aligned}
 H \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} & \left[\frac{(m_{\ell+\ell-\nu\bar{\nu}}^2 - m_h^2)^2}{\sigma_{h\ell}^4} + \frac{(m_{\nu\bar{\nu}}^2 - m_{\nu\bar{\nu},peak}^2)^2}{\sigma_{\nu}^4} \right. \\
 & + \min \left(\frac{(m_{\ell+\nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell-\bar{\nu}}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4}, \right. \\
 & \left. \left. \frac{(m_{\ell-\bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell+\nu}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4} \right) \right], \\
 & \sim m_h - m_W \text{ off-shell}
 \end{aligned}$$

- Higgsness provides a degree of consistency to dileptonic $h \rightarrow WW^*$ system.
- The off-shell W also has an end-point near $m_h - m_W$.
- Its distribution is wide, but there is a peak, which can constrain hh system further.

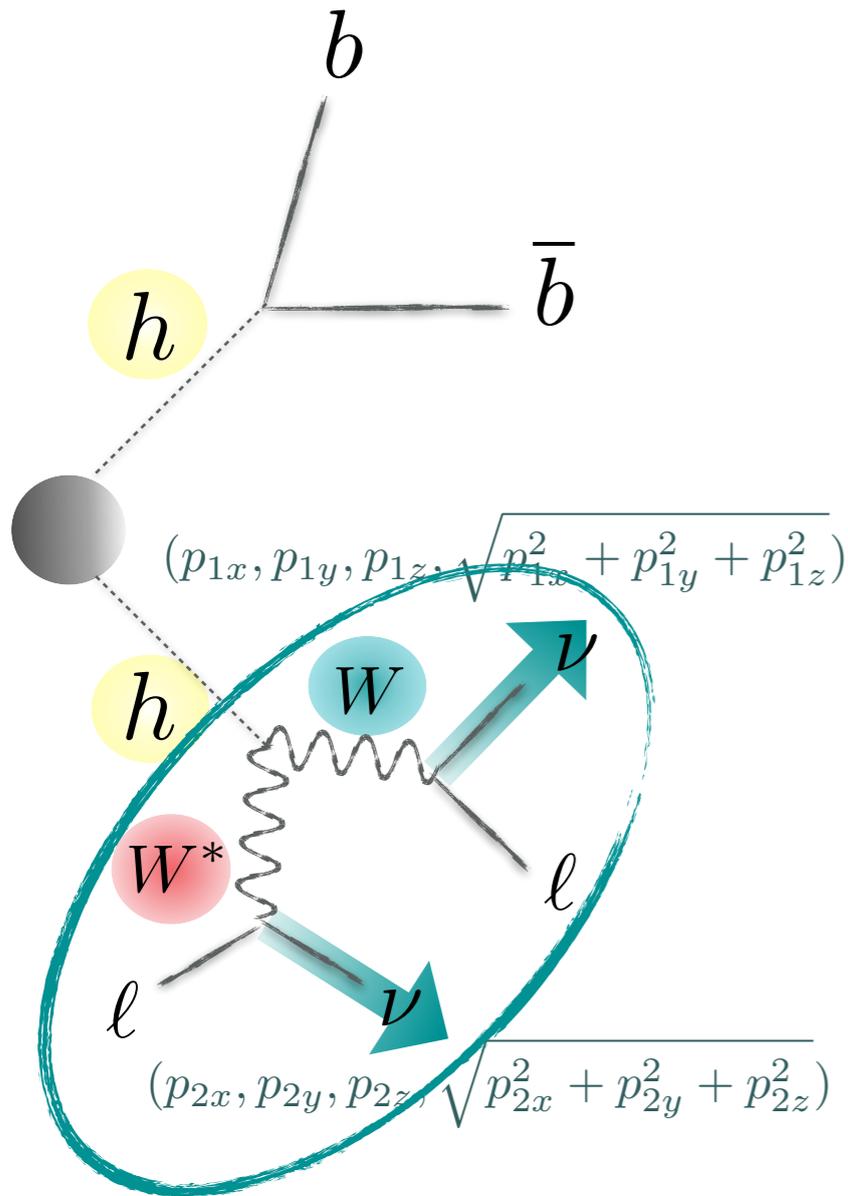
Higgsness (H)



$$\begin{aligned}
 H \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} & \left[\frac{(m_{\ell+\ell-\nu\bar{\nu}}^2 - m_h^2)^2}{\sigma_{h\ell}^4} + \frac{(m_{\nu\bar{\nu}}^2 - m_{\nu\bar{\nu},peak}^2)^2}{\sigma_{\nu}^4} \right. \\
 & + \min \left(\frac{(m_{\ell+\nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell-\bar{\nu}}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4} \right. \\
 & \left. \left. \frac{(m_{\ell-\bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell+\nu}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4} \right) \right], \\
 & \sim m_h - m_W \text{ off-shell}
 \end{aligned}$$

- Higgsness provides a degree of consistency to dileptonic $h \rightarrow WW^*$ system.
- The off-shell W also has an end-point near $m_h - m_W$.
- Its distribution is wide, but there is a peak, which can constrain hh system further.

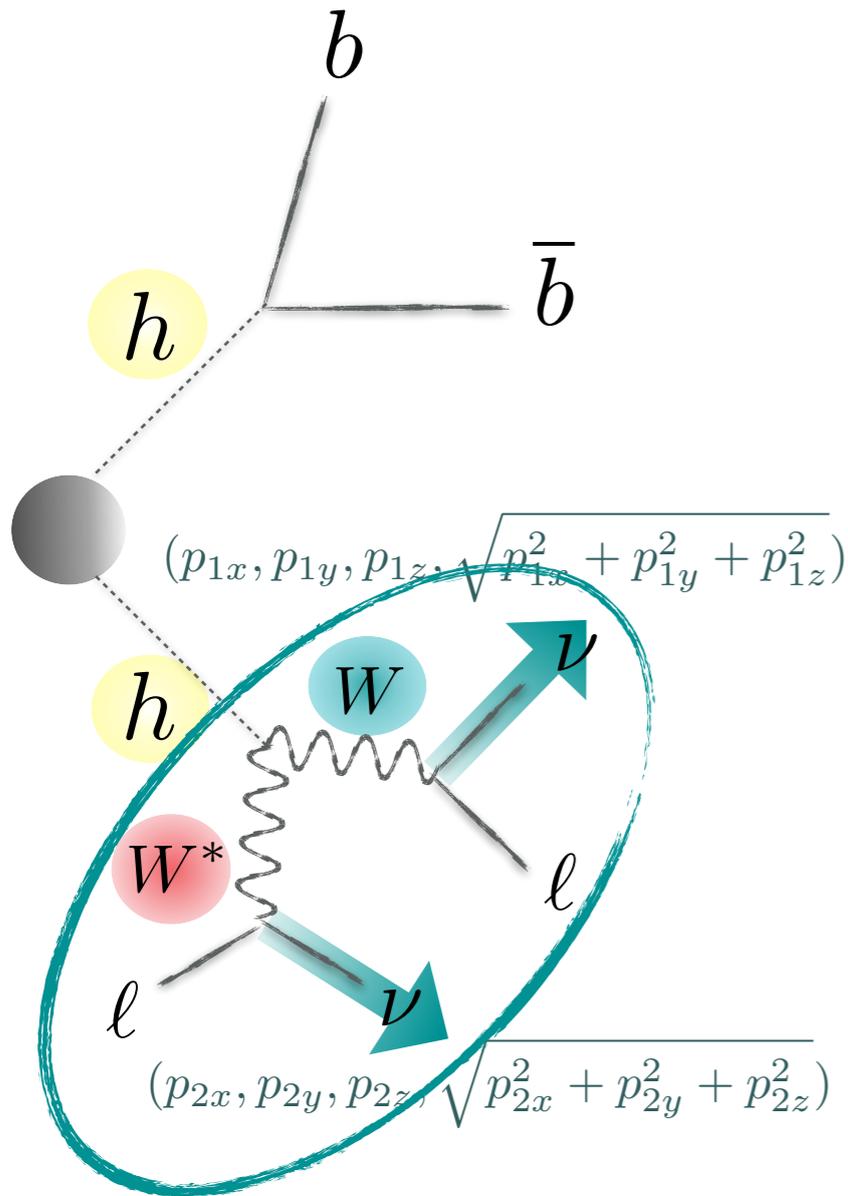
Higgsness (H)



$$\begin{aligned}
 H \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} & \left[\frac{(m_{\ell^+ \ell^- \nu \bar{\nu}}^2 - m_h^2)^2}{\sigma_{h\ell}^4} + \frac{(m_{\nu \bar{\nu}}^2 - m_{\nu \bar{\nu}, peak}^2)^2}{\sigma_{\nu}^4} \right. \\
 & + \min \left(\frac{(m_{\ell^+ \nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell^- \bar{\nu}}^2 - m_{W^*, peak}^2)^2}{\sigma_{W^*}^4}, \right. \\
 & \left. \left. \frac{(m_{\ell^- \bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell^+ \nu}^2 - m_{W^*, peak}^2)^2}{\sigma_{W^*}^4} \right) \right], \\
 & \sim m_h - m_W \text{ off-shell}
 \end{aligned}$$

- Higgsness provides a degree of consistency to dileptonic $h \rightarrow WW^*$ system.
- The off-shell W also has an end-point near $m_h - m_W$.
- Its distribution is wide, but there is a peak, which can constrain hh system further.

Higgsness (H)



$$H \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} \left[\frac{(m_{\ell^+ \ell^- \nu \bar{\nu}}^2 - m_h^2)^2}{\sigma_{h\ell}^4} + \frac{(m_{\nu \bar{\nu}}^2 - m_{\nu \bar{\nu}, peak}^2)^2}{\sigma_{\nu}^4} \right]$$

$$+ \min \left(\frac{(m_{\ell^+ \nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell^- \bar{\nu}}^2 - m_{W^*, peak}^2)^2}{\sigma_{W^*}^4} \right),$$

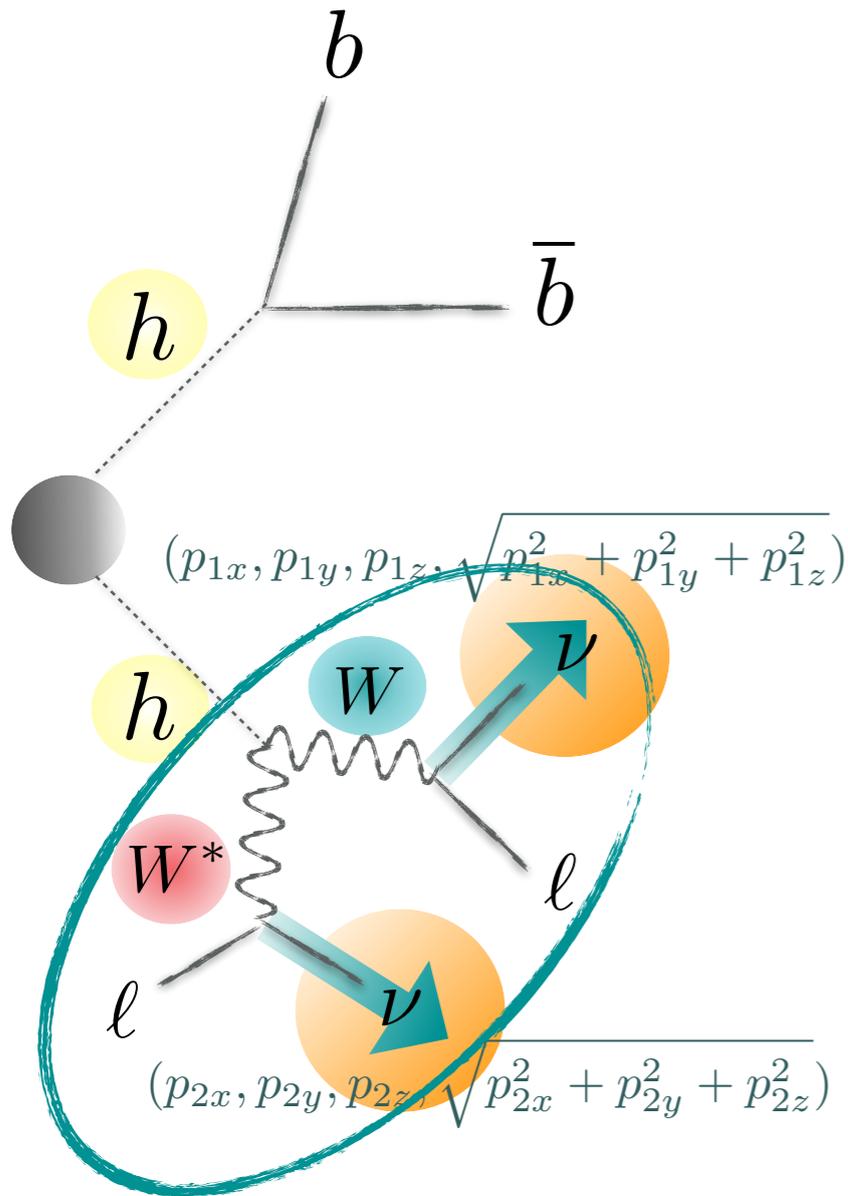
$$\left[\frac{(m_{\ell^- \bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell^+ \nu}^2 - m_{W^*, peak}^2)^2}{\sigma_{W^*}^4} \right],$$

two possible ways of paring ν and ℓ

$\sim m_h - m_W$
off-shell

- Higgsness provides a degree of consistency to dileptonic $h \rightarrow WW^*$ system.
- The off-shell W also has an end-point near $m_h - m_W$.
- Its distribution is wide, but there is a peak, which can constrain hh system further.

Higgsness (H)



$$H \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} \left[\frac{(m_{\ell+\ell-\nu\bar{\nu}}^2 - m_h^2)^2}{\sigma_{h\ell}^4} + \frac{(m_{\nu\bar{\nu}}^2 - m_{\nu\bar{\nu},peak}^2)^2}{\sigma_{\nu}^4} \right]$$

$$+ \min \left(\frac{(m_{\ell+\nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell-\bar{\nu}}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4} \right),$$

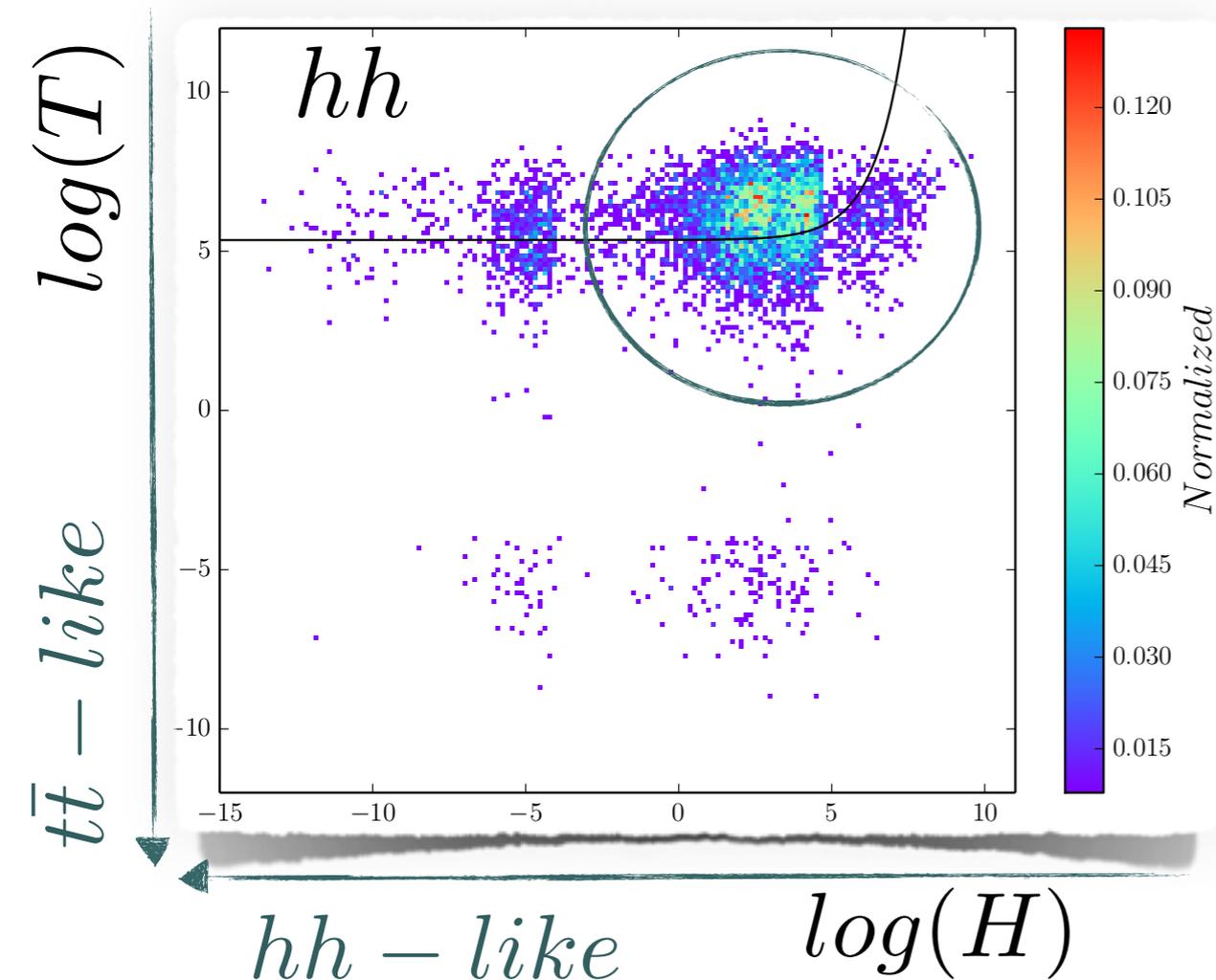
$$\left[\frac{(m_{\ell-\bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell+\nu}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4} \right],$$

two possible ways of paring ν and ℓ

$\sim m_h - m_W$
off-shell

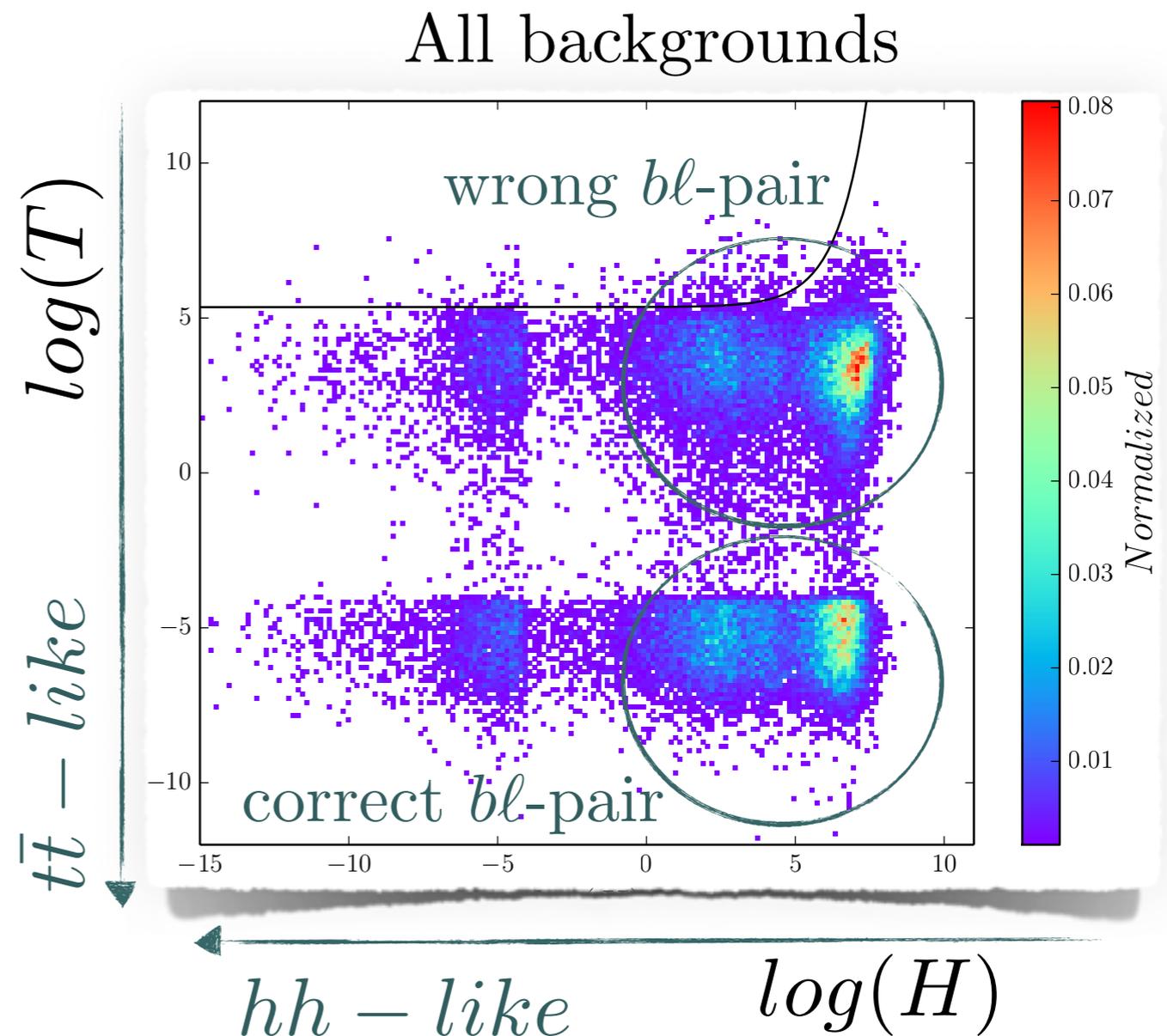
- Higgsness provides a degree of consistency to dileptonic $h \rightarrow WW^*$ system.
- The off-shell W also has an end-point near $m_h - m_W$.
- Its distribution is wide, but there is a peak, which can constrain hh system further.

Distributions of $(\log H, \log T)$ after baseline selection cuts

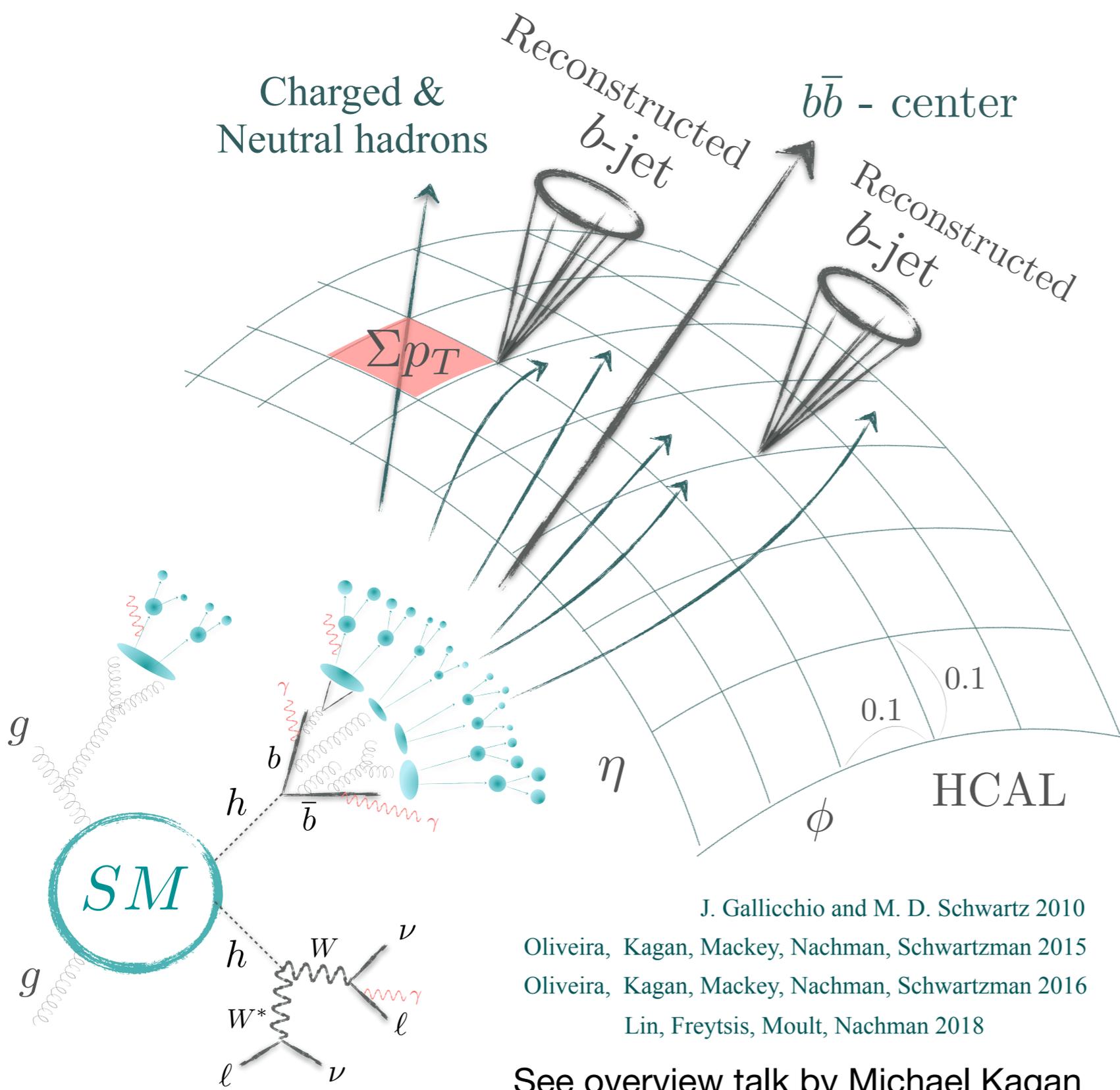


- A clear separation between hh and backgrounds ($t\bar{t}$ is dominant)

- The Topness (Higgsness) variable measures the degree of consistency of a given event with the kinematics of dilepton $t\bar{t}$ (hh) production.

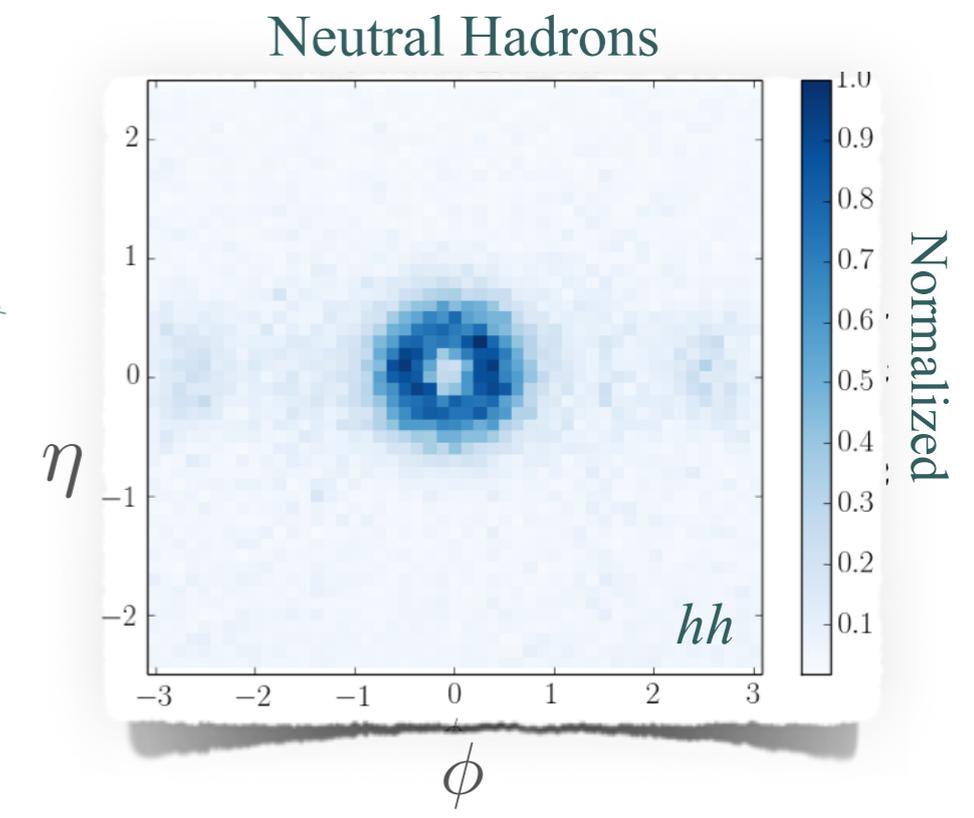
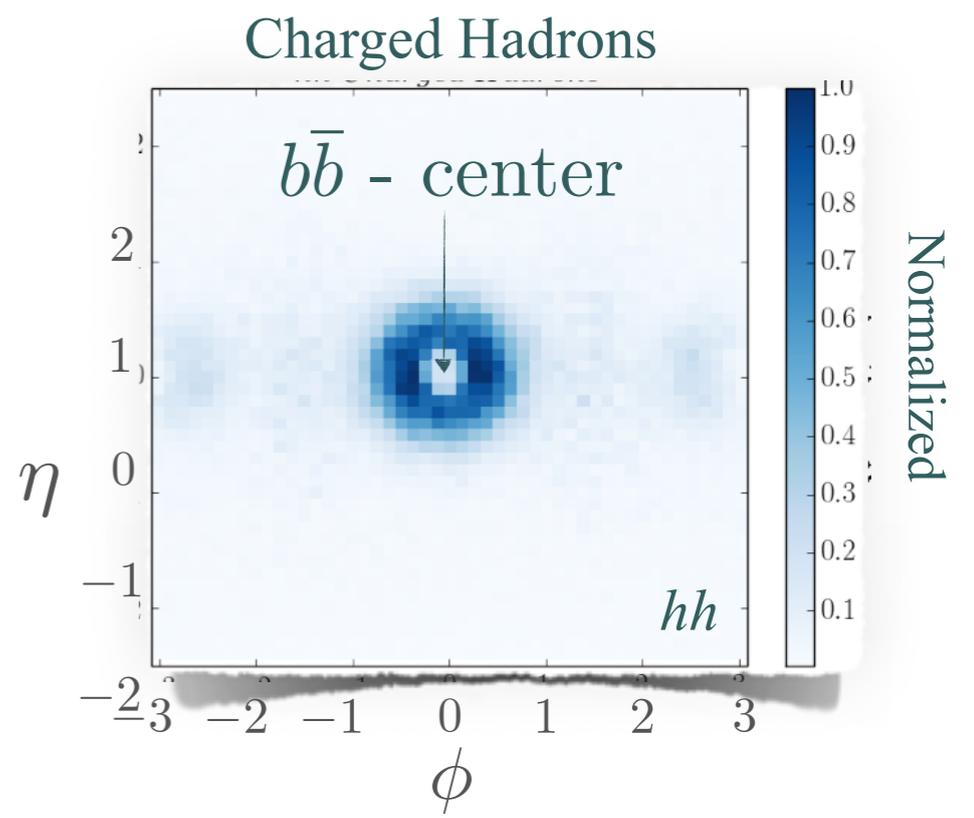


Processing Hadron Images (hh)



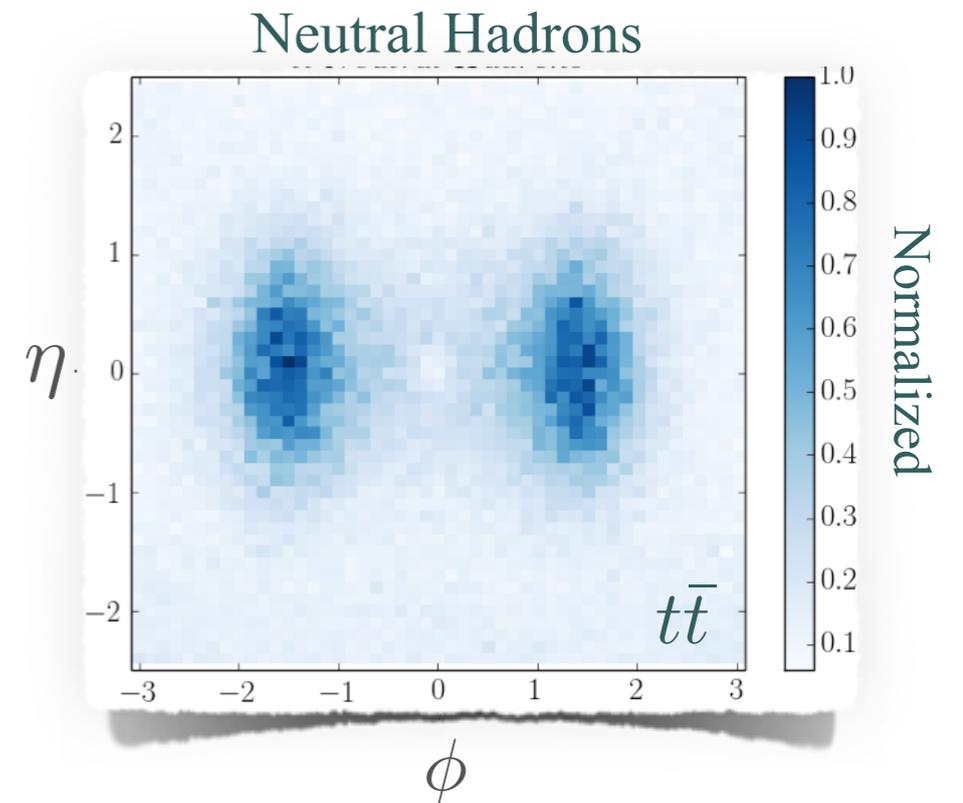
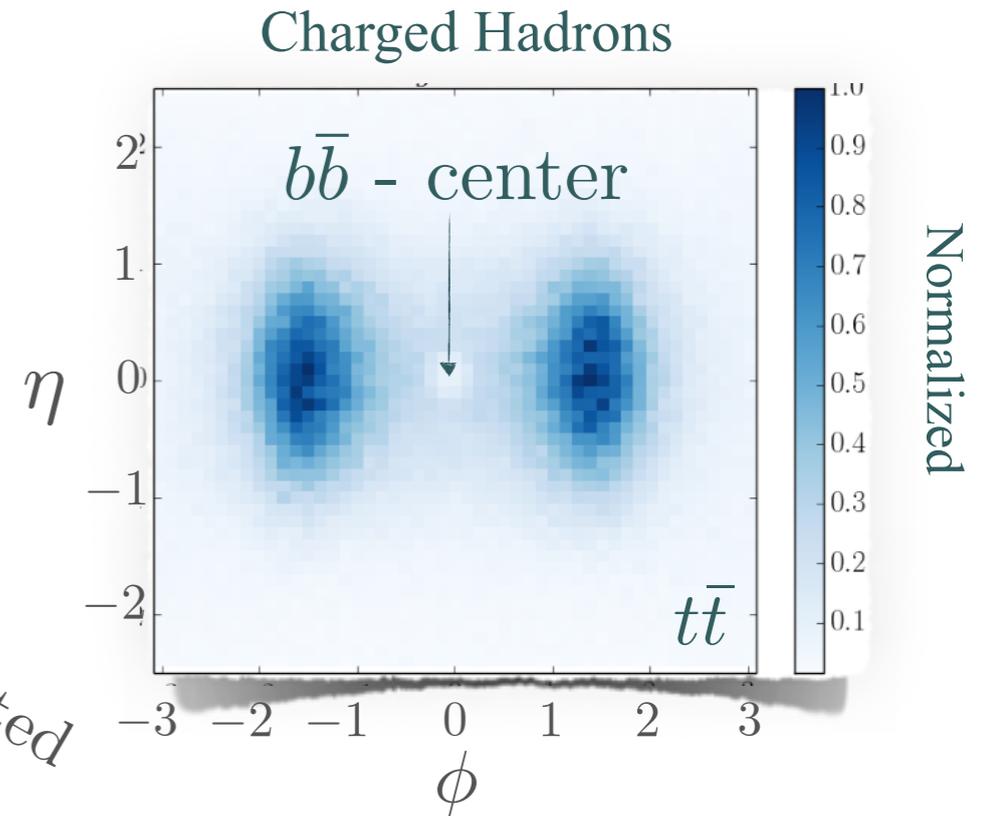
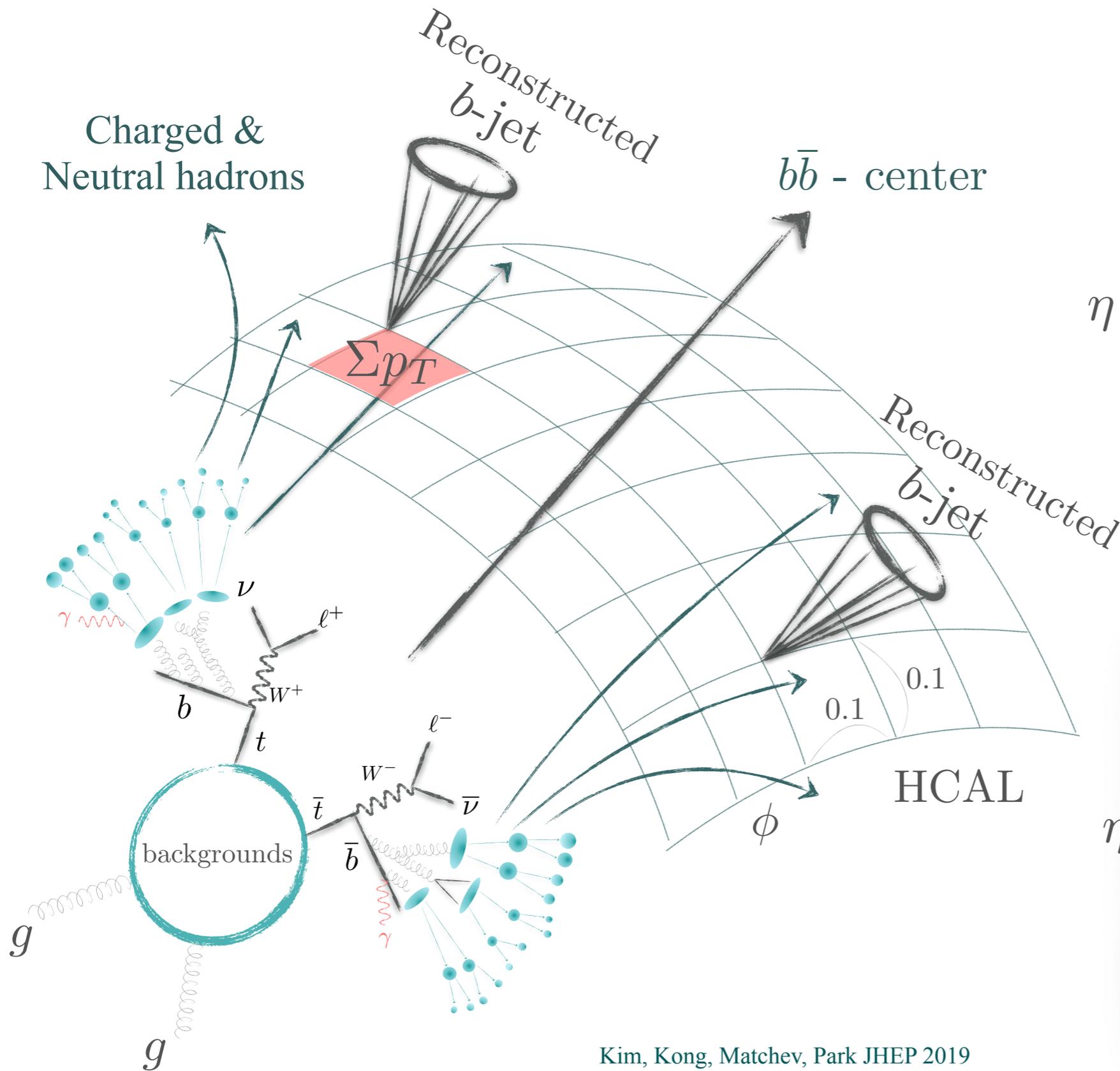
J. Gallicchio and M. D. Schwartz 2010
 Oliveira, Kagan, Mackey, Nachman, Schwartzman 2015
 Oliveira, Kagan, Mackey, Nachman, Schwartzman 2016
 Lin, Freytsis, Moutl, Nachman 2018

See overview talk by Michael Kagan



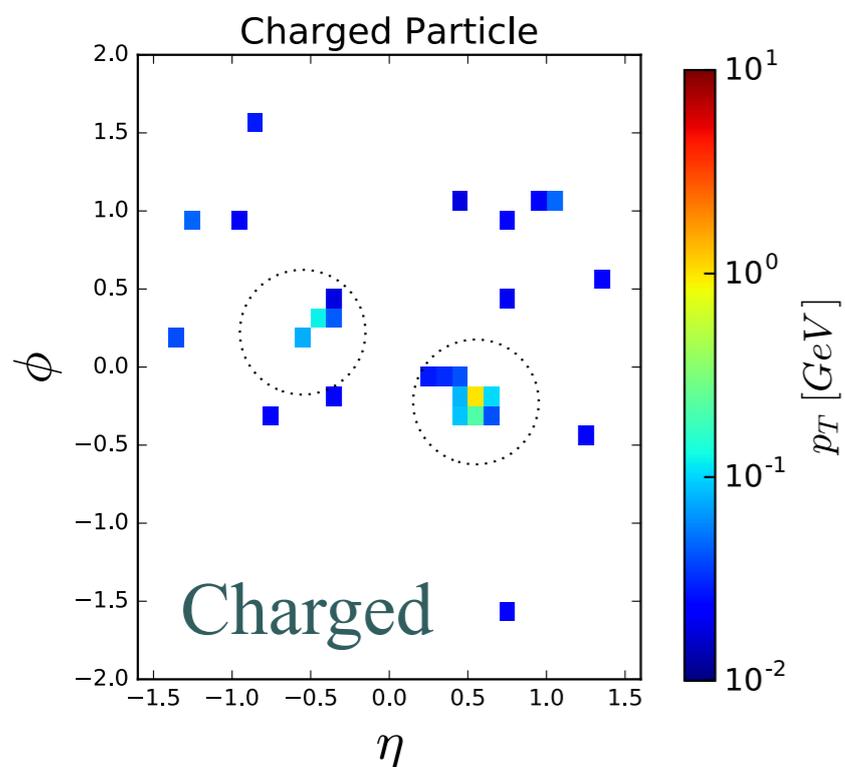
Kim, Kong, Matchev, Park JHEP 2019

Processing Hadron Images (tt)



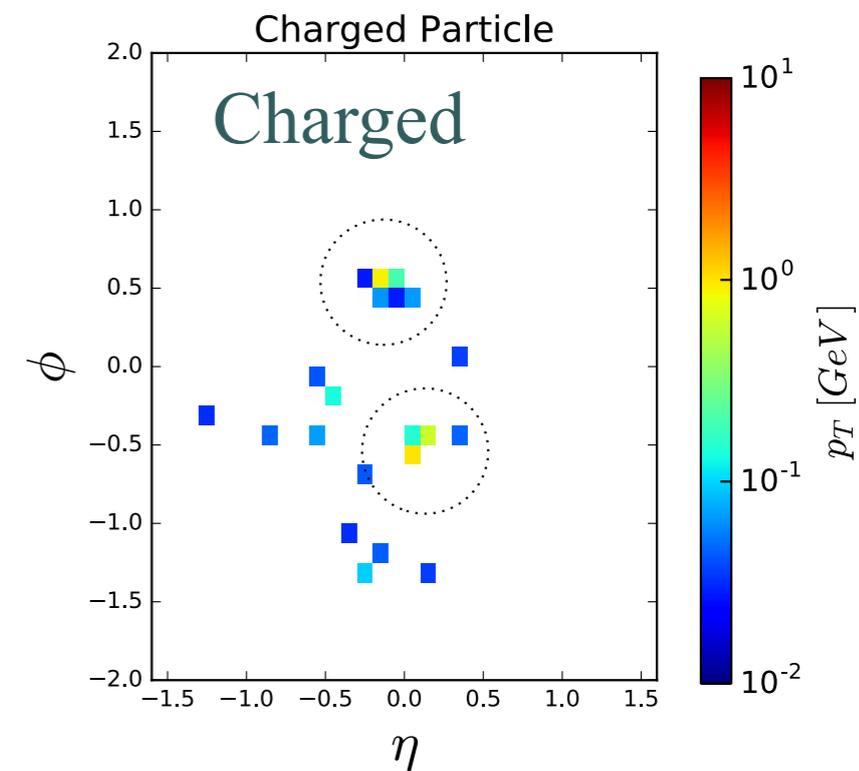
Jet images: event-by-event

Each hh event

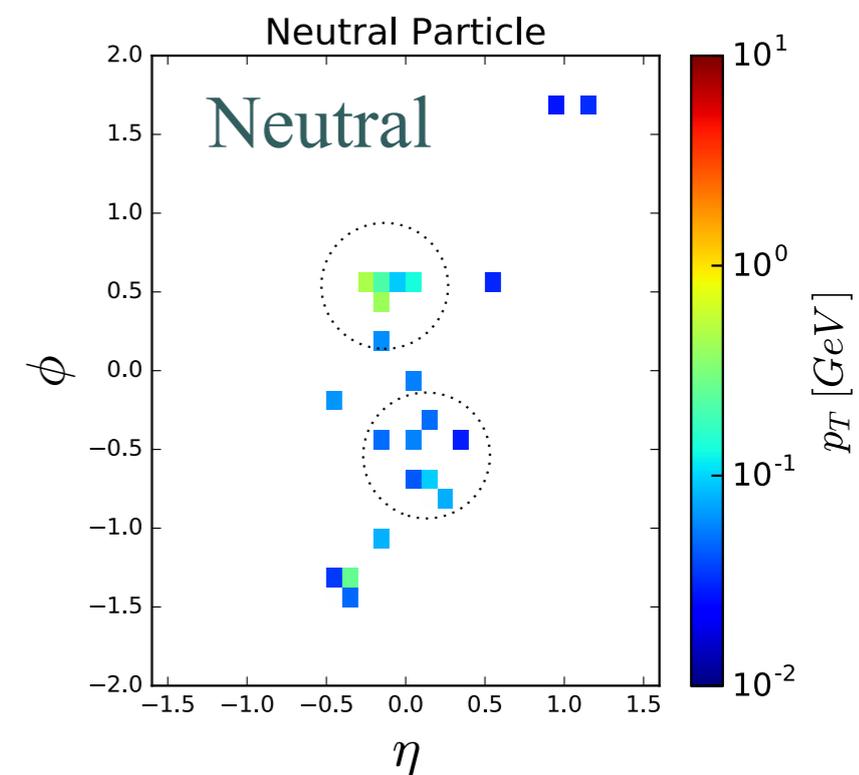
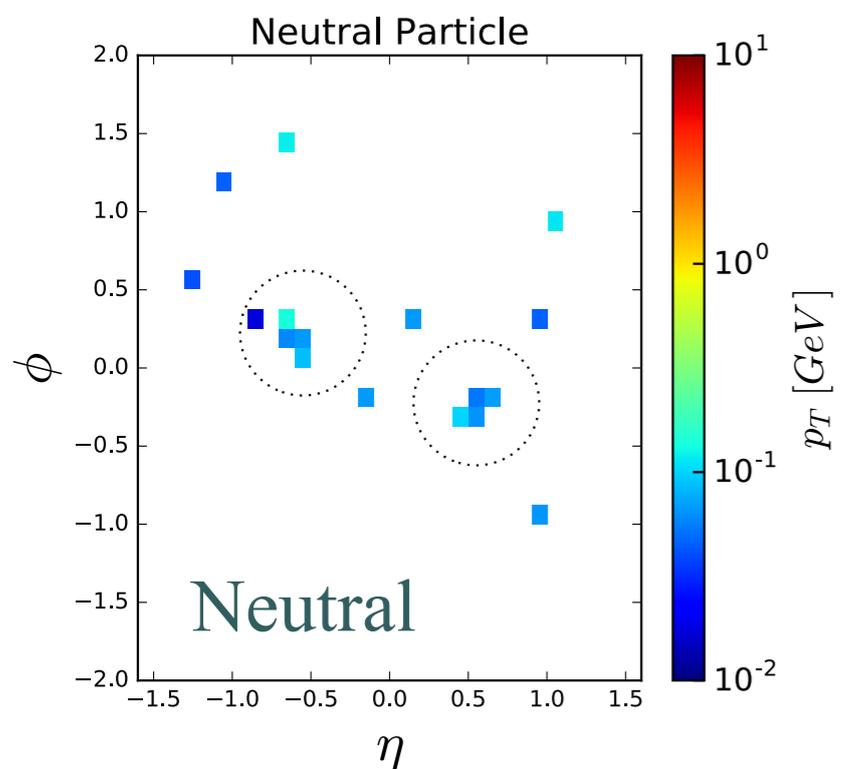


- It is difficult to distinguish them using the cut-and-count method...

Each tt event

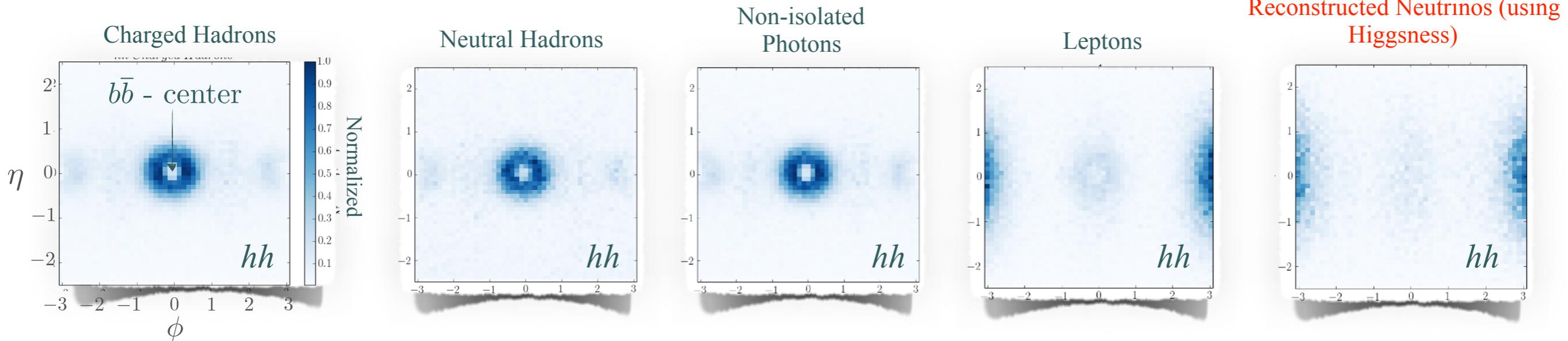


- We train the machine using Deep Neural Network (DNN) architecture.

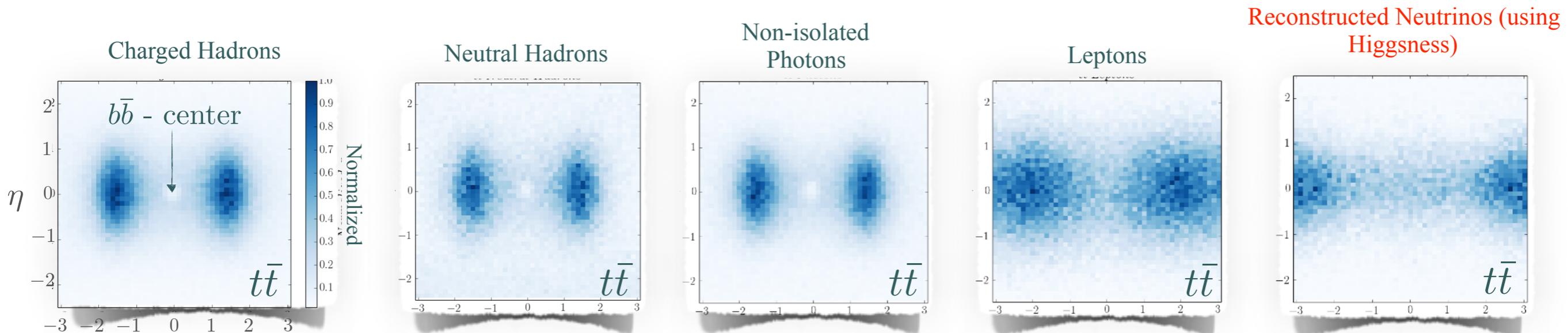


The Di-Higgs Photography

- Totally hadrons, lepton, photon, and neutrino images are shown for hh .



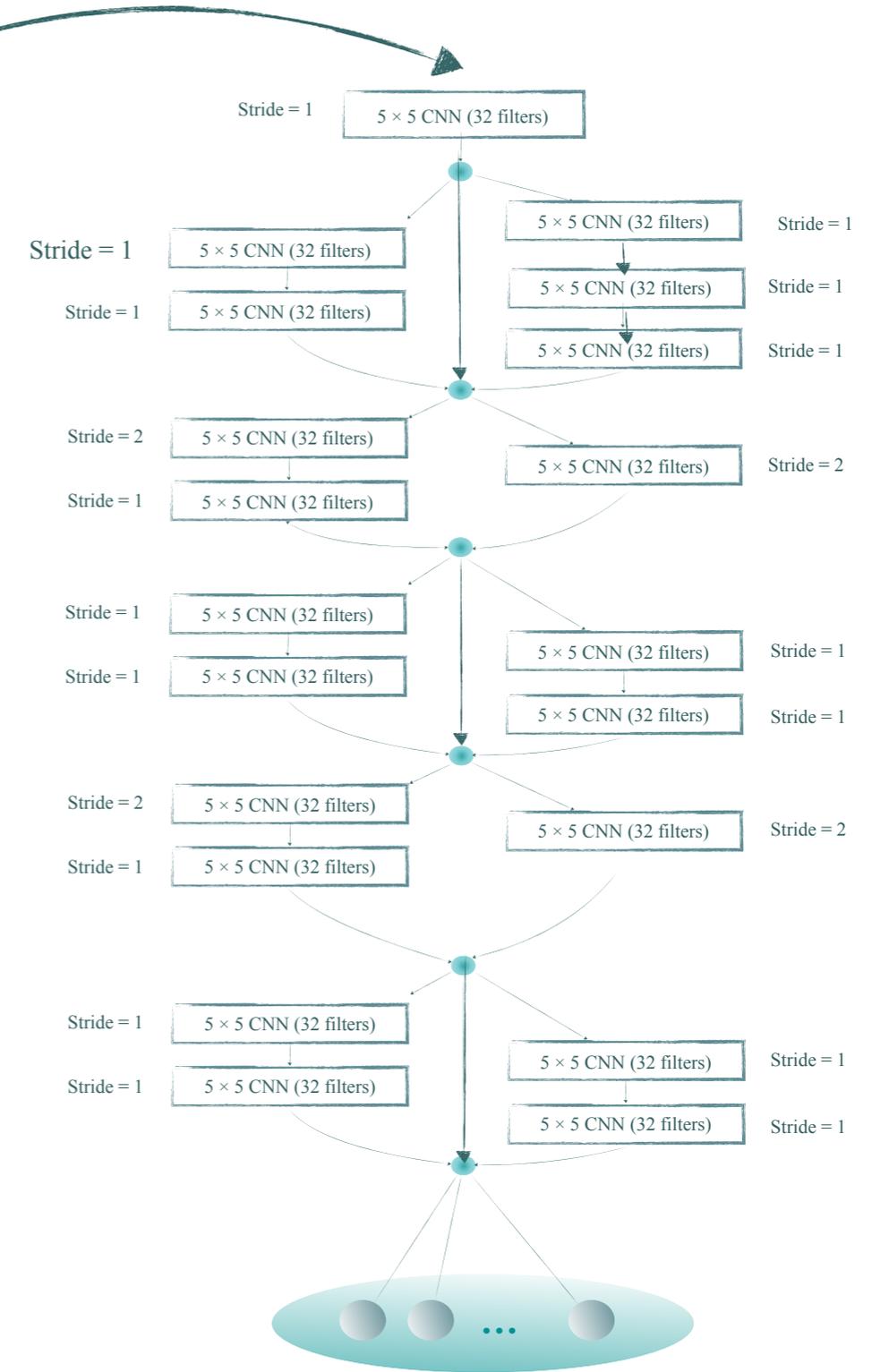
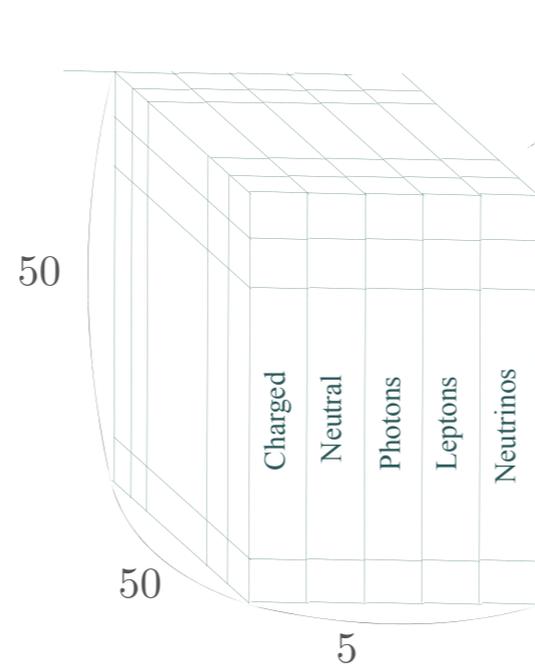
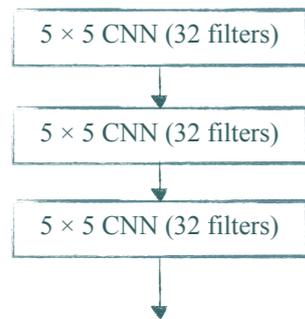
- Analogous five images are shown for $t\bar{t}$ background.



- A sharp difference between hh and $t\bar{t}$.

Our ResNet

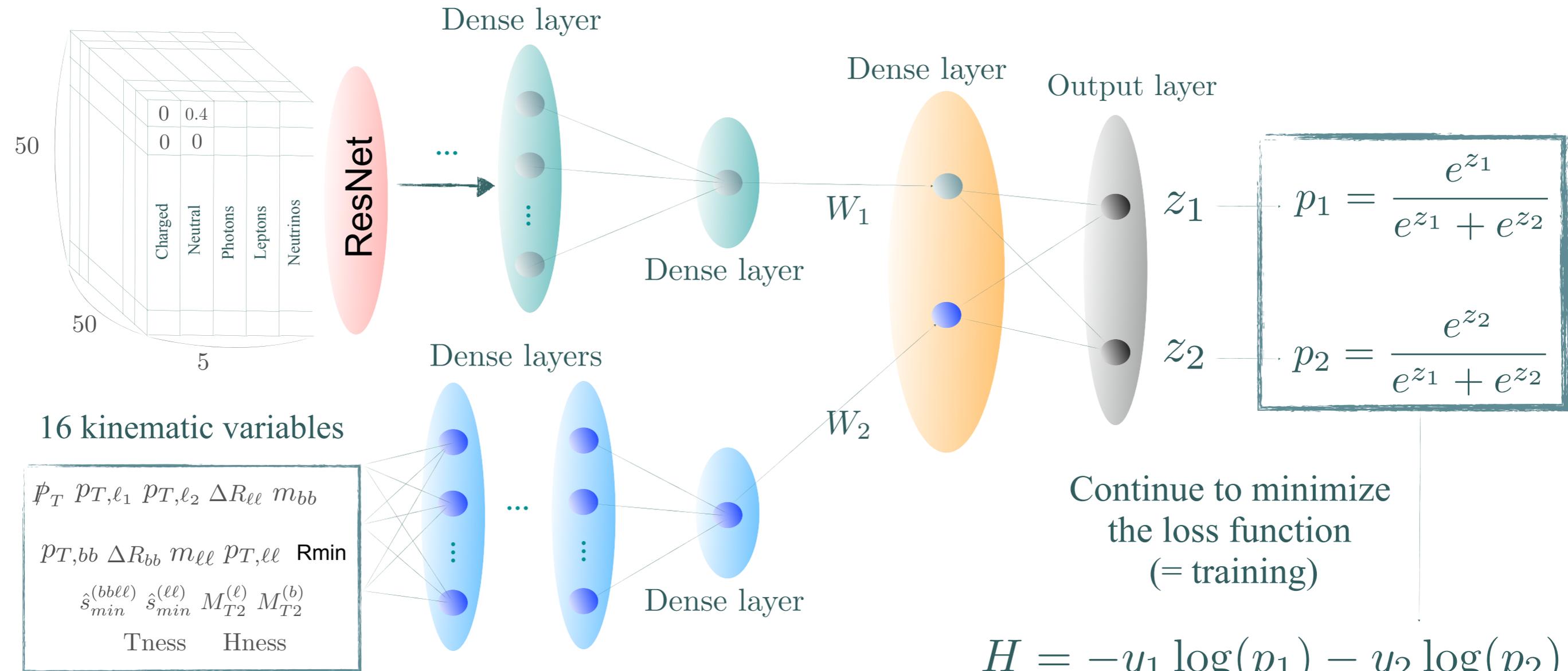
Classic CNN



- As the network becomes deeper, the performance of CNN gets saturated or even starts degrading rapidly.
- The classic CNN starts to degrade after 3 layers, with our 5-image data.
- Absolutely not suitable.

- Motivated by ResNet, we could design a much deeper network.
- The machine will be able to resolve much detailed features of 5-image data.
- This ensemble-like topology is much resilient against the change in network hyper-parameter.

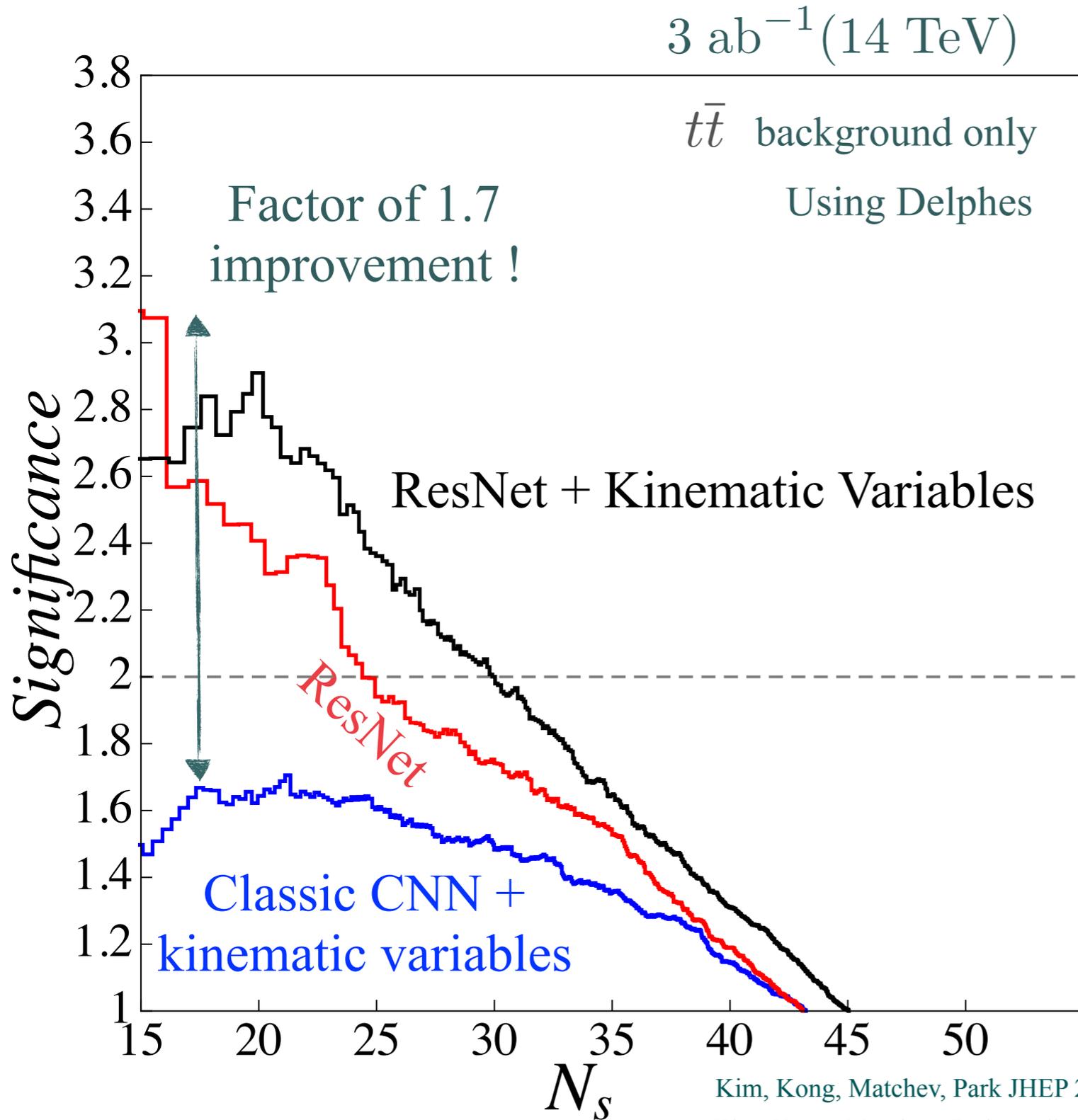
Combining image data and high level kinematic variables in dense neural networks



Signal : $(y_1, y_2) = (1, 0)$

Backgrounds : $(y_1, y_2) = (0, 1)$

Preliminary Results



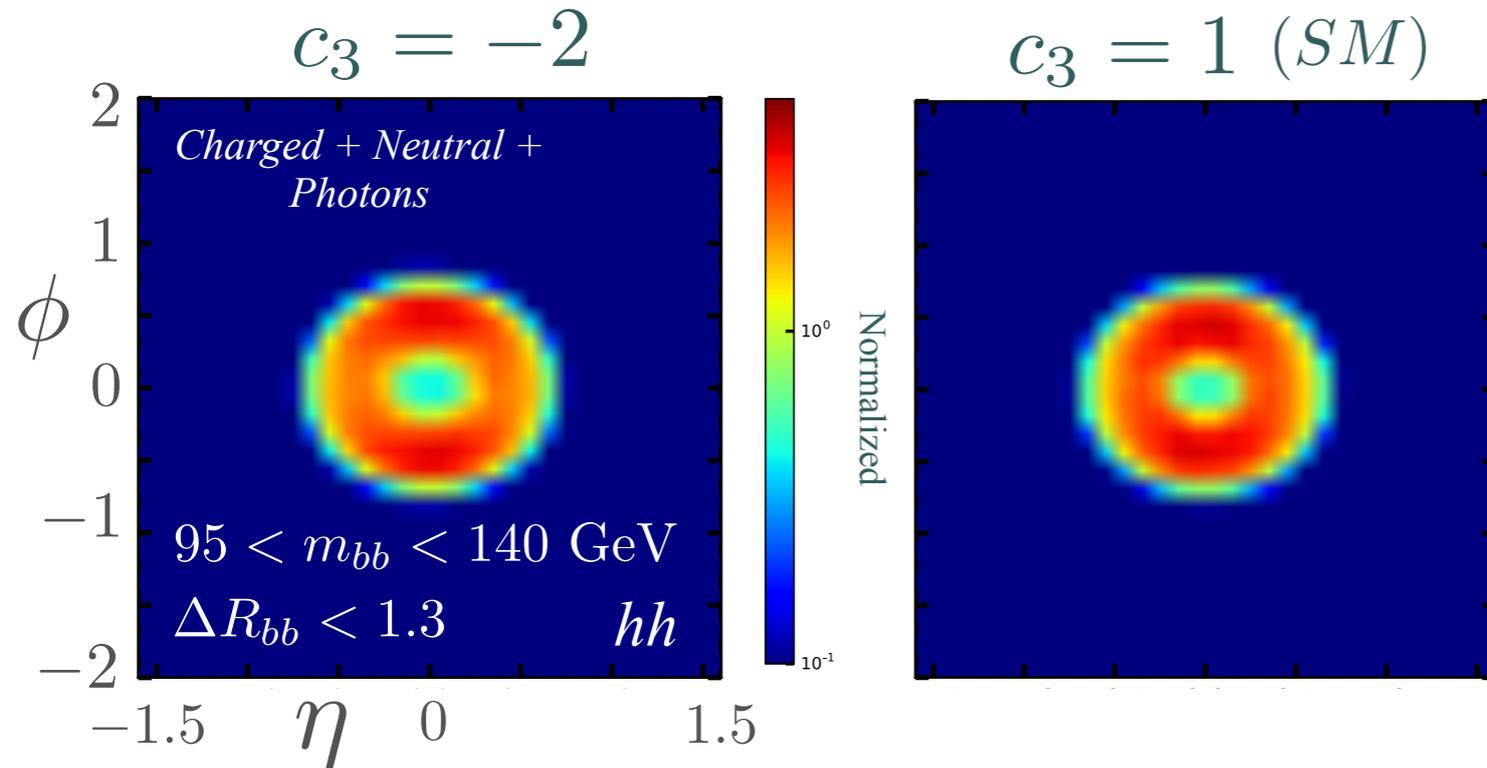
Kim, Kong, Matchev, Park JHEP 2019

Kim, Kong, Matchev, Park, preliminary

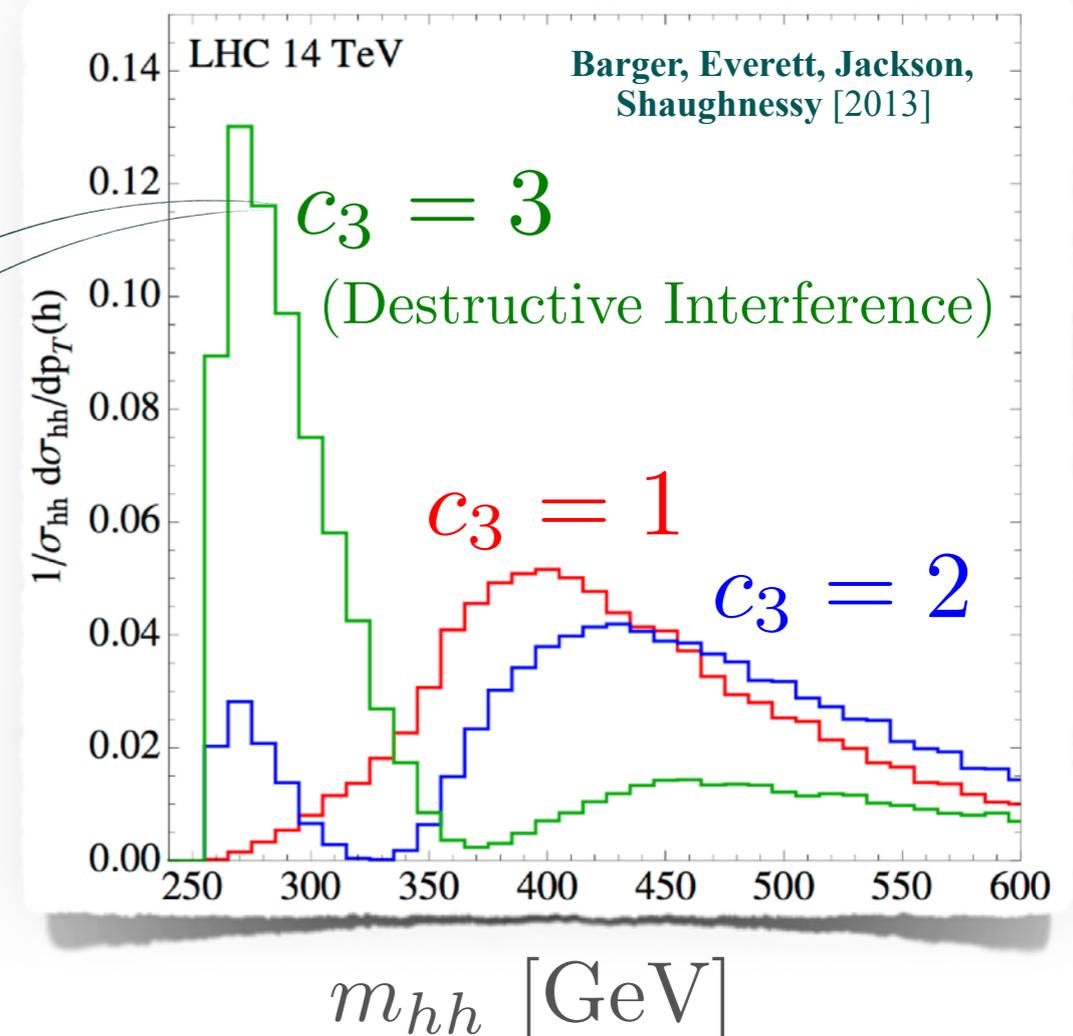
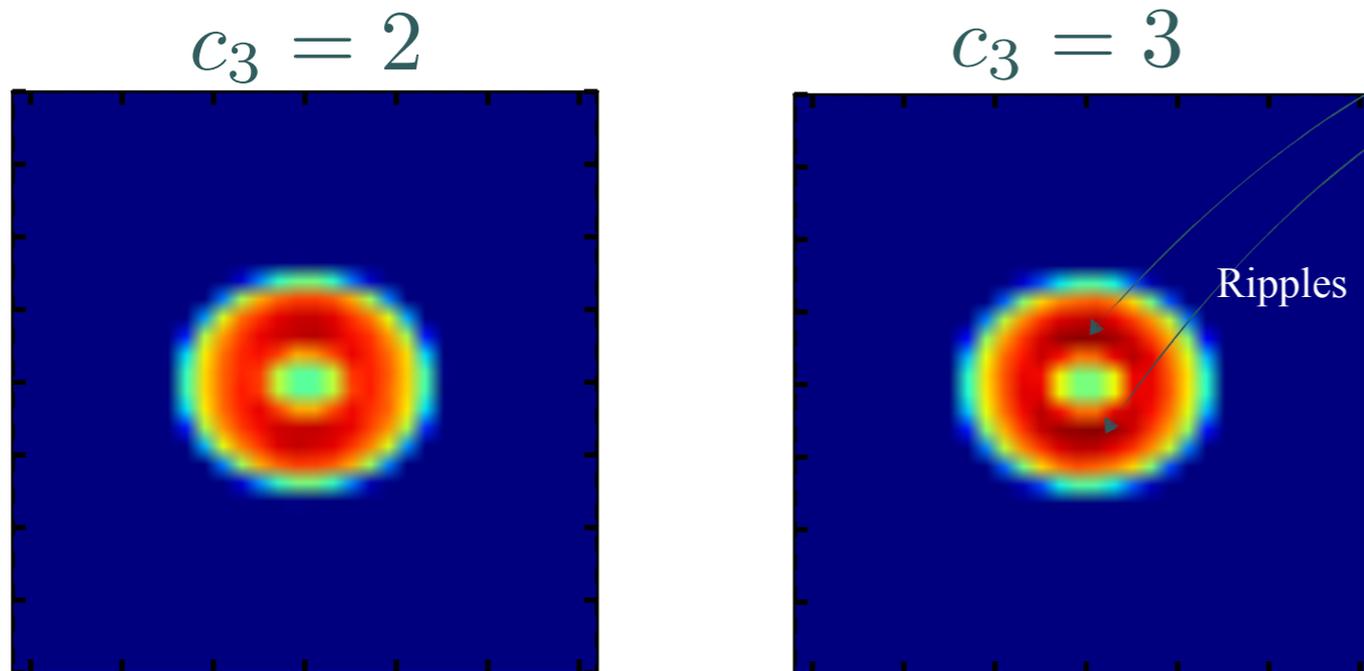
- As a preliminary study, we included only $t\bar{t}$ background.
- Significance with kinematic variables only or jet images only give signal significance below 1.
- We find that the classic CNN + kinematic variables can reach at most 1.7 sigma.
- Our ResNet can reach up to 2.9 sigma (factor of ~ 1.7 improvement)
- Our ResNet can be further reinforced by combining reconstructed kinematic variables.
- More recent architecture, Capsule Nets also provides a similar performance.

Shifting the Higgs triple coupling c_3

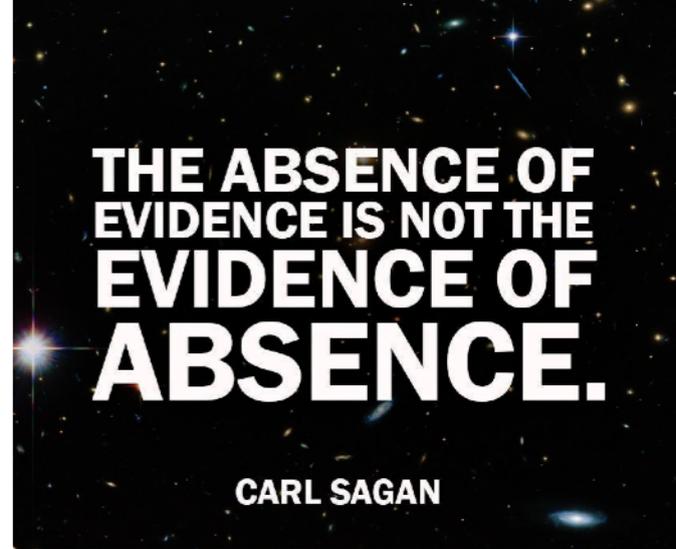
- How does the hh image vary by shifting the Higgs triple coupling c_3 ? $\mathcal{L} \supset -c_3 \frac{m_h^2}{2v} h^3$



- Interestingly, the gradient of images change as moving into the region of a destructive interference.
- We are working on how much neural network can be sensitive to this change (we don't know yet, sorry).



Summary



- LHC has great opportunity to study new physics beyond the SM.
- Measurement of triple Higgs self coupling reveals the nature of electroweak symmetry breaking.
- Double Higgs production is challenging due to small signal cross section / large SM backgrounds, and strong correlation among many kinematic variables.
- Multivariate analysis could benefit from deep neural networks using jet images and Topness / Higgsness with mass information.
- $bbWW$ dilepton channel would make a significant contribution in triple Higgs coupling measurement. There is room for further improvement.
- Semi-leptonic channel would be similar.