

# A Tower WGC from IR consistency

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*based on SA, D. Junghans, T. Noumi, G. Shiu 1802.04287*



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# Outline

## Introduction and review

- Landscape and Swampland
- WGC and CHC

## Infrared Consistency and WGC

- Positivity Bounds (causality & analyticity)
- **A stronger CHC**

## Consistency under dimensional reduction

- (sub-)Lattice WGC
- **Tower WGC**

## Conclusions



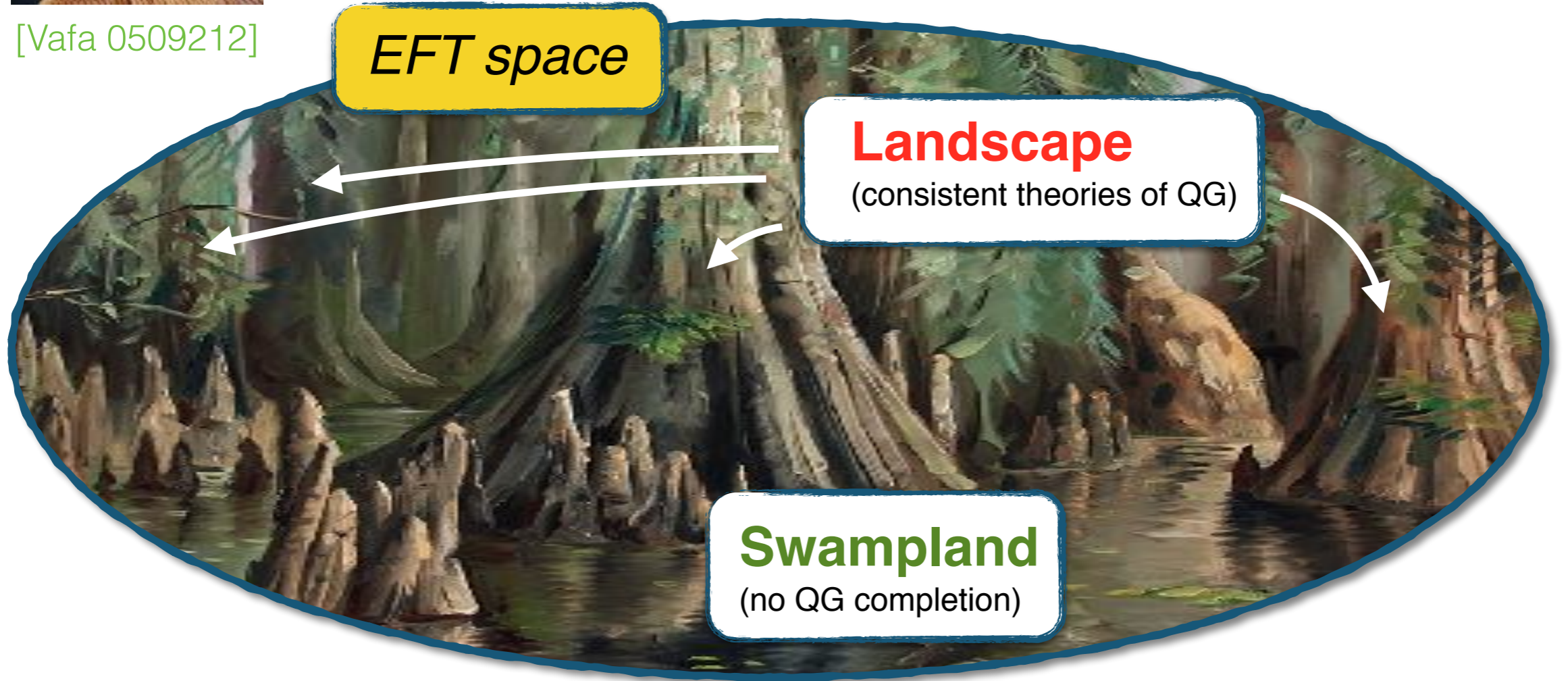
[Vafa 0509212]

SWAMPLAND!

*EFT space*

**Landscape**  
(consistent theories of QG)

**Swampland**  
(no QG completion)





[Vafa 0509212]

SWAMPLAND!

*EFT space*

**Boundaries ? Swampland criteria**

Recent work

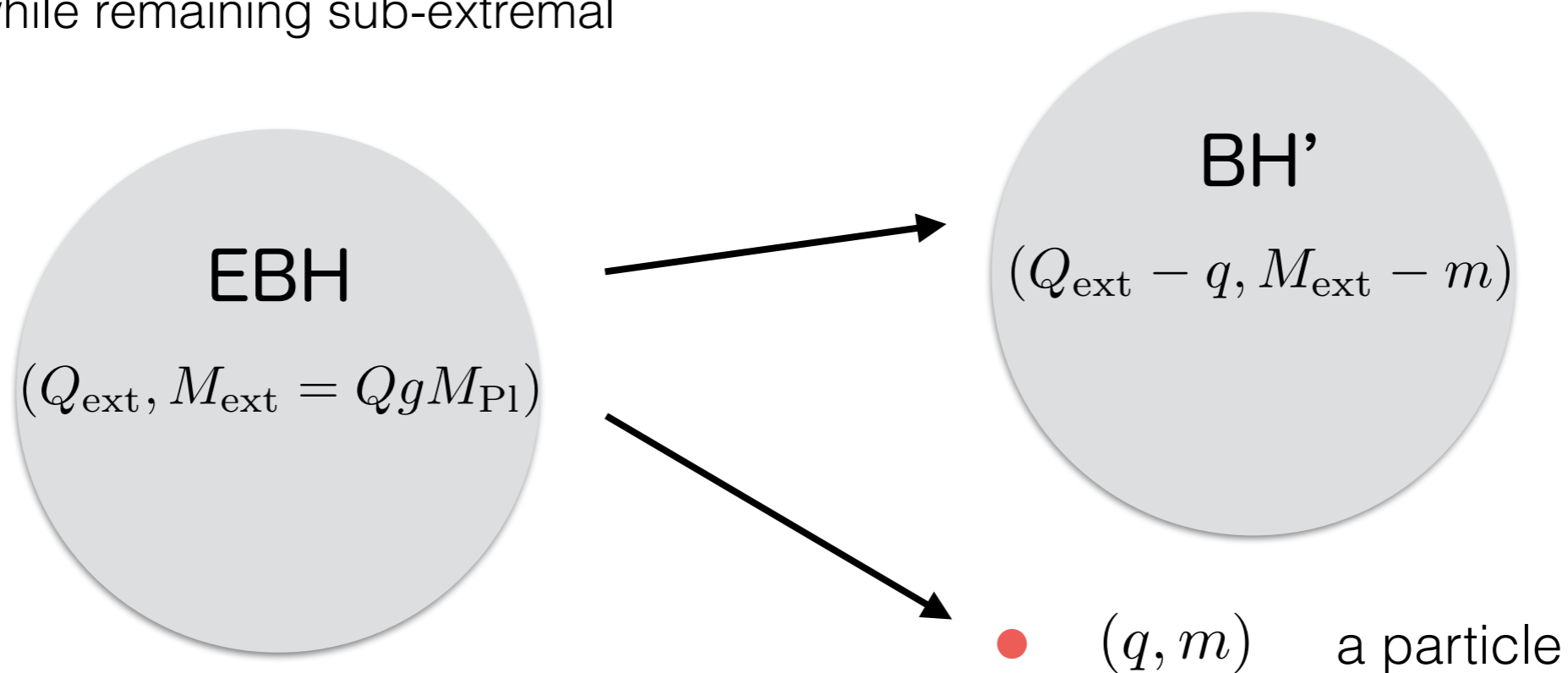
- ▶ **Distance conjecture** (see Palti's, Grimm's and Marchesano's talks)
- ▶ **dS** (see McAllister's and Grana's talks)
- ▶ **new criteria** (see Vafa's and Shiu's talks)
- ▶ **WGC** (see also Montero's and Heidenreich's talks)

# WGC & CHC

In any consistent EFT of gauge U(1) coupled to gravity  
There must exist a particle (m,q) with charge-mass ratio

$$z \equiv \frac{gqM_P}{m} \geq 1$$

To avoid troubles with remnants: EBH's must decay into charged object while remaining sub-extremal



$$g(Q_{\text{ext}} - q) \leq (M_{\text{ext}} - m)/M_{\text{Pl}} \iff gq \geq m/M_{\text{Pl}}$$

# WGC & CHC

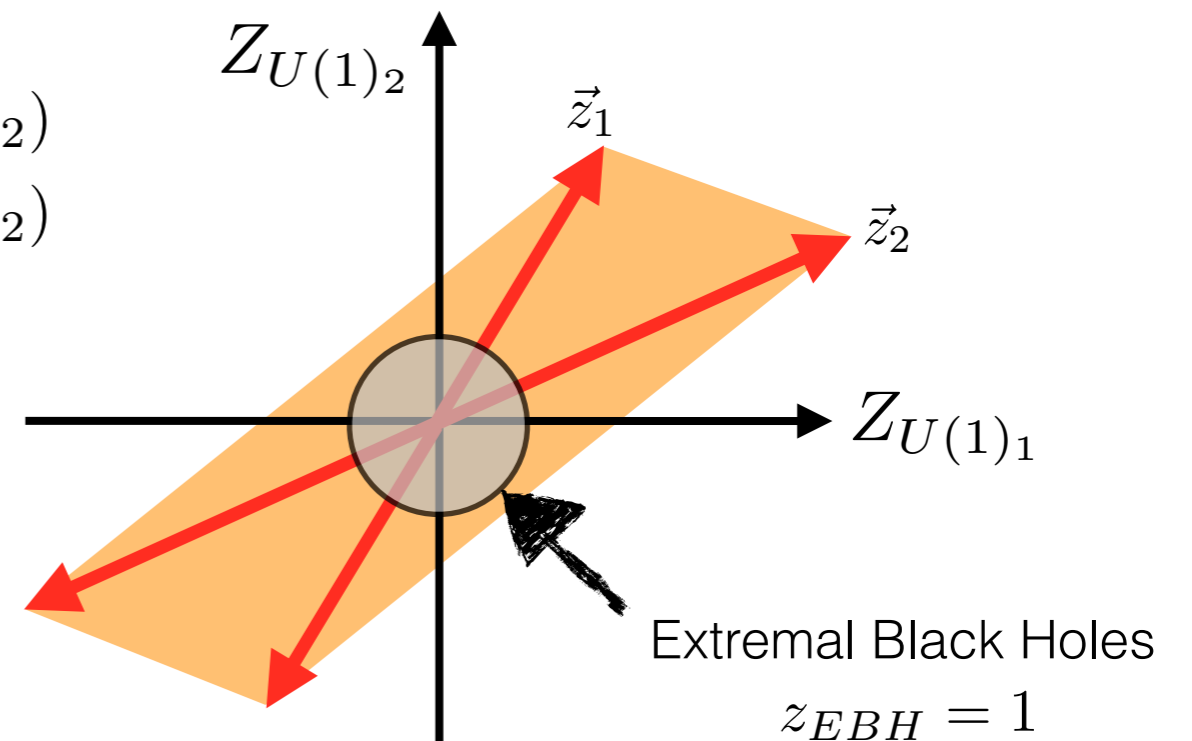
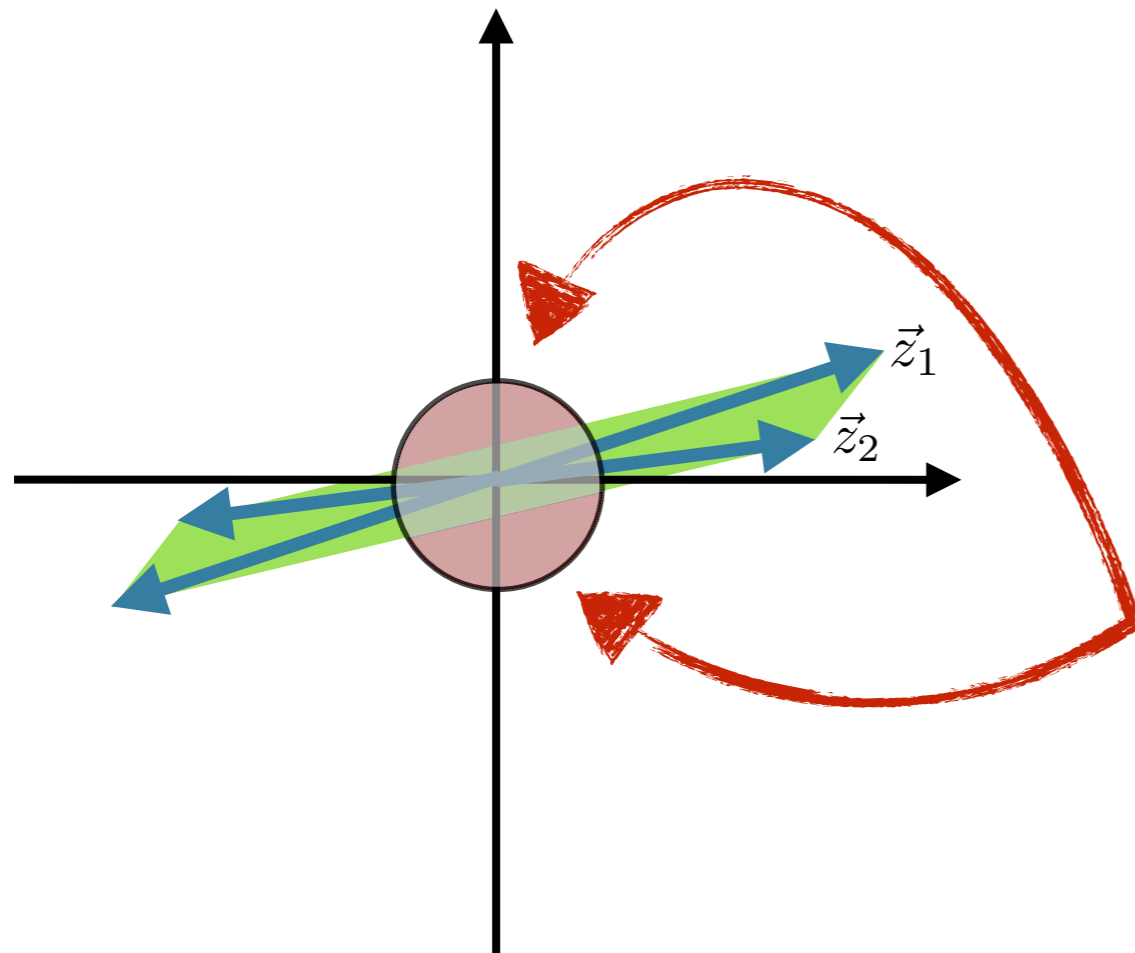
## Multiple U(1)'s:

(sub-)Extremal Black holes can decay if they lie inside the Convex Hull!

Ex.:  $U(1)^2$ , 2 particles

$$\vec{z}_1 = (z_{11}, z_{12})$$

$$\vec{z}_2 = (z_{21}, z_{22})$$



There are no states which can discharge Black Holes in the red regions!

# Positivity bounds (IR)

An EFT with HO operators = higher derivatives corrections

Eg. 
$$\mathcal{L}_{1-loop} = \frac{M_P}{2} R - \frac{F^2}{4e^2} + CF^4$$

is “IR consistent” i.e. respects:

**causality** (subluminal fluctuations)

**analyticity of S-matrix**  $\mathcal{M}(s)$

(non-anal in  $s$  where particles are produced)

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(non-anal in  $s$  where particles are produced)

if

$$C > 0$$



# Causality (eg 3D)

Dualisation:  $\mathcal{L}_1 = \frac{M_3}{2} R - \frac{(\partial\phi)^2}{2} + 4C(\partial\phi)^4 \quad F \sim *d\phi$

Require the speed of propagation of fluctuations  $\phi = \bar{\phi} + \varphi$  is sub-luminal (in any locally flat frame  $\eta_{ab}$ )

EOMs  $\xrightarrow{\text{Fourier space}}$  dispersion relation

$$(\eta_{ab} + 16C \overline{\partial_a\phi\partial_b\phi}) k^a k^b = 0 \quad \text{True for any bg choice}$$

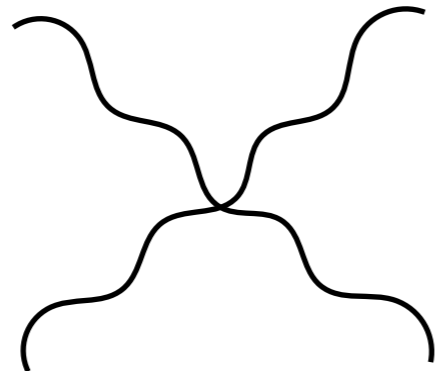
Simplest bg: constant EM field  $\overline{\partial_a\phi} = w_a = (w_0, \vec{w}) = \text{const}$

$$v = \frac{k_0}{|\vec{k}|} = 1 - 8C(w_0 - \vec{w} \cdot \hat{k})^2 \quad \longrightarrow \quad \boxed{C > 0}$$

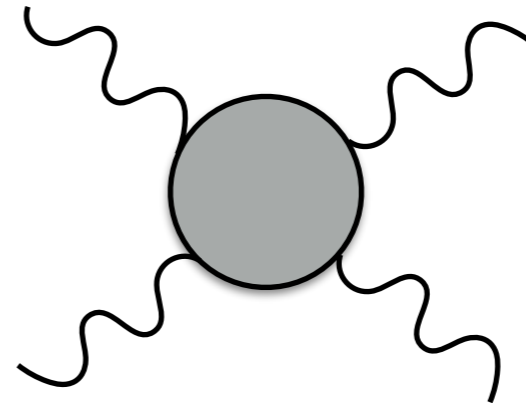
# Analyticity - 3D/4D

Consider 4-pts photon scatterings

IR



from



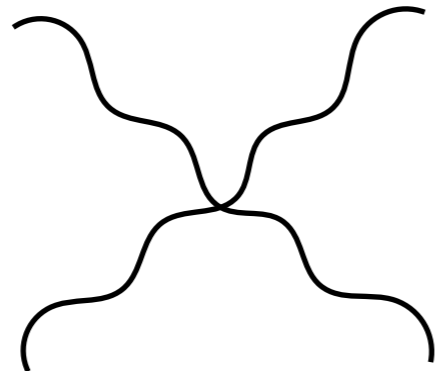
In the forward limit  $\mathcal{M}(s, t) = 8C(s^2 + t^2 + u^2) \xrightarrow[t \rightarrow 0]{} \mathcal{M}(s) = 16Cs^2$

We can extract C from

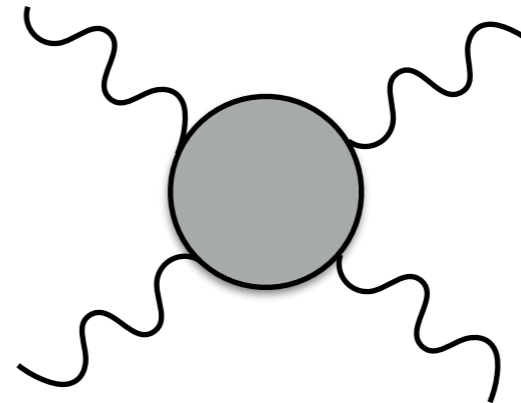
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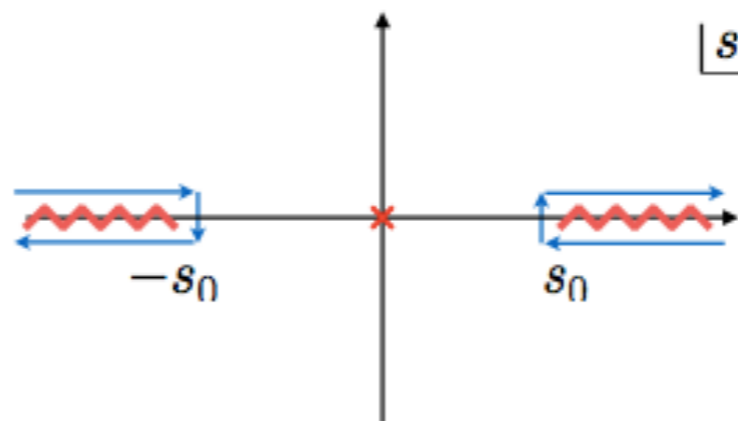
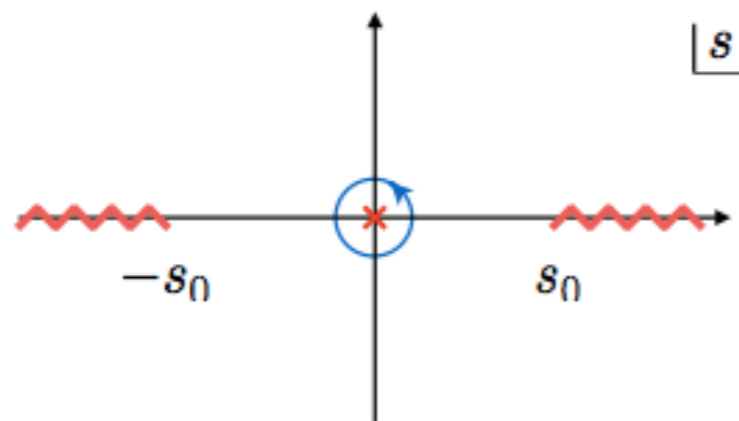


$$16C = \oint_{\gamma} \frac{ds}{2\pi i} \frac{\mathcal{M}(s)}{s^3} = \left( \int_{-\infty}^{-s_0} + \int_{s_0}^{\infty} \right) \frac{ds}{2\pi i} \frac{\text{Disc}[\mathcal{M}(s)]}{s^3} = \frac{2}{\pi} \int_{s_0}^{\infty} ds \frac{\text{Im}\mathcal{M}(s)}{s^3} > 0$$

contour def  
+ **analyticity**  
+ Froissart bound

analyticity  
+ Schwarz reflection  
+ crossing symm

Opt Th  
 $\text{Im}\mathcal{M}(s) = s\sigma(s)$



$$C > 0$$

The S-matrix is **analytical** along the real axis  $|s| < s_0$ , up to the lowest energy where on-shell intermediate states are created (= red discontinuities)

Positivity bounds



WGC ?

$$C > 0$$



**bound for z ?**

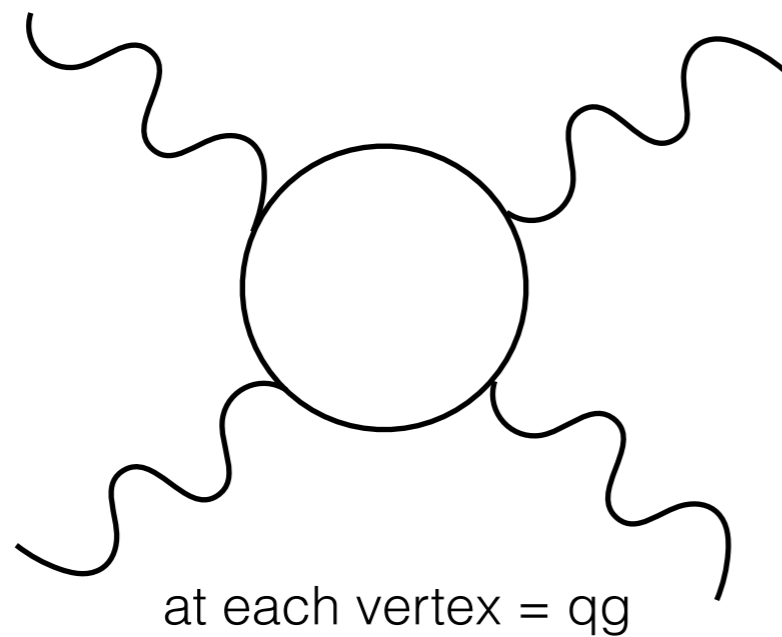
$$f(z) > 0$$

[Cheung-Remmen '14]

**YES! \***

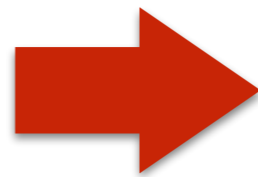
\* caveats...

Suppose the 1-loop EFT is obtained by **integrating out charged matter**



$$C(z)$$

$$C > 0$$



Conditions on  
 $z$

## Setup (3d)

EFT multiple scalar/fermions charged under multiple U(1)'s

$$\Gamma = \int d^3x \sqrt{-g} \left[ \frac{M_3}{2} R - \frac{1}{4} \sum_i F_i^2 \right] + \Gamma_{s/f} + H.O. \quad \boxed{\text{UV}}$$

$$\Gamma_s = \int d^3x \sqrt{-g} \sum_a (-|D_\mu \phi_a|^2 - m_a^2 |\phi_a|^2)$$

$$\Gamma_f = \int d^3x \sqrt{-g} \sum_a \bar{\psi}_a (-\Gamma^\mu D_\mu - m_a) \psi_a$$

$$H.O. = \sum_{ijkl} c_{ijkl} (F_i \cdot F_j)(F_k \cdot F_l)$$

**UV-Physics dof**  
(kept **generic/unknown**)

$\Lambda$



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**UV-Physics dof**  
(kept **generic/unknown**)

“Elephant in the room”...


**Nonetheless, there is a regime where we can extract interesting results!**

$\Lambda$



## Setup (3d)

Integrating out matter, we obtain

$\min(m_a)$  

$$\Gamma_1 = \int d^3x \sqrt{-g} \left[ \frac{M_3}{2} R - \frac{1}{4} \sum_{i,j} \delta_{ij} F_i \cdot F_j + \sum_{i,j,k,l} C_{ijkl} (F_i \cdot F_j) (F_k \cdot F_l) \right]$$



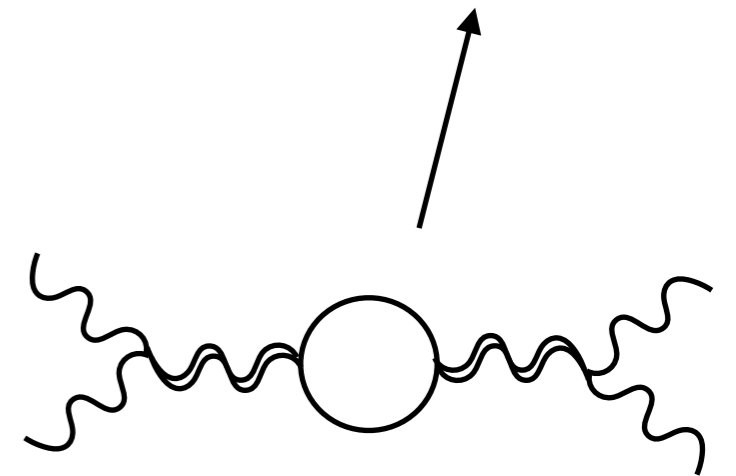
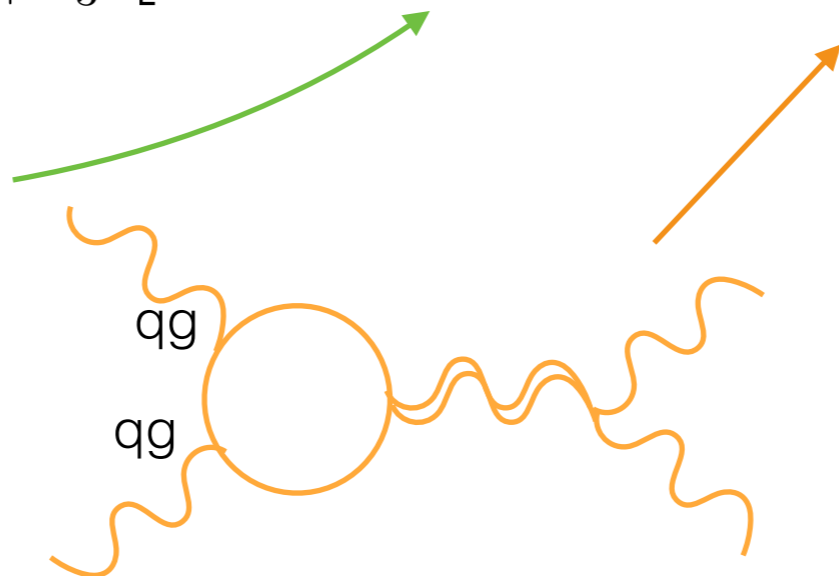
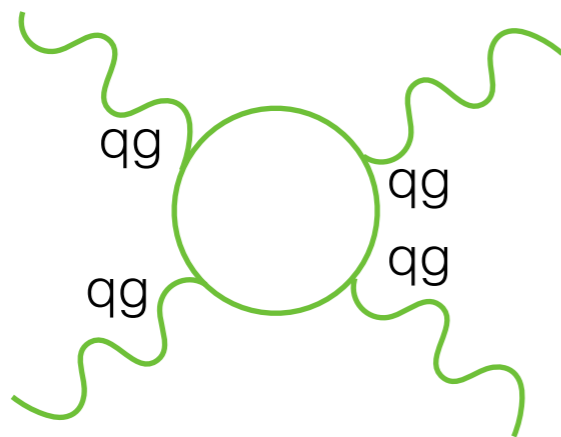
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$$C_{ijkl}^s = C_{ijkl} + \sum_a \frac{1}{1920\pi |m_a| M_3^2} \left[ \frac{7}{8} z_{ai} z_{aj} z_{ak} z_{al} + \frac{3}{2} z_{ai} z_{aj} \delta_{kl} - z_{ai} z_{ak} \delta_{jl} + \frac{1}{2} \delta_{ij} \delta_{kl} + \delta_{ik} \delta_{jl} \right]$$

$$C_{ijkl}^f = C_{ijkl} + \sum_a \frac{1}{1920\pi |m_a| M_3^2} \left[ z_{ai} z_{aj} z_{ak} z_{al} + z_{ai} z_{aj} \delta_{kl} - \frac{3}{2} z_{ai} z_{ak} \delta_{jl} - \frac{1}{2} \delta_{ij} \delta_{kl} + \frac{3}{2} \delta_{ik} \delta_{jl} \right]$$



Implications:

- **single particle**, charged under **single U(1)**

$$C > 0 \left\{ \begin{array}{l} \text{scalar: } z^2 \left( z^2 + \frac{4}{7} \right) + \mathcal{O}_s(z^0) > 0 \\ \text{fermion: } z^2 \left( z^2 - \frac{1}{2} \right) + \mathcal{O}_f(z^0) > 0 \end{array} \right.$$

Implications:

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$$C > 0 \left\{ \begin{array}{l} \text{scalar: } z^2 \left( z^2 + \frac{4}{7} \right) + \mathcal{O}_s(z^0) > 0 \quad \longrightarrow \quad \text{trivial} \\ \text{fermion: } z^2 \left( z^2 - \frac{1}{2} \right) + \mathcal{O}_f(z^0) > 0 \quad \longrightarrow \quad z > \frac{1}{\sqrt{2}} \end{array} \right.$$

**WGC!**

In the regime where  $\mathcal{O}(z^0)$  is negligible

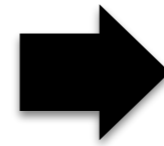
$$\mathcal{O}(z^0) \sim$$



$$+ \text{[Feynman diagram: a circle with wavy lines on both sides]} \longrightarrow 0$$

Implications:

- **more particles, multiple U(1)'s:**

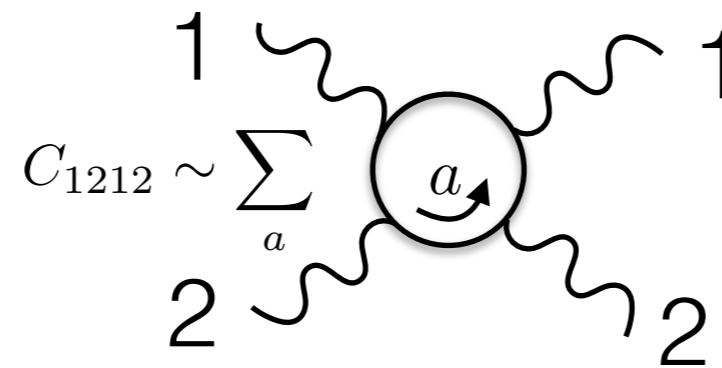


Several positivity conditions on

$$C_{ijkl}$$

Strongest positivity conditions are given by **mixed** scatterings

e.g.: **U(1)xU(1)**



a=matter field  $\phi_a, \psi_a$   
 $m_a, \vec{q}_a = (q_{1a}, q_{2a})$

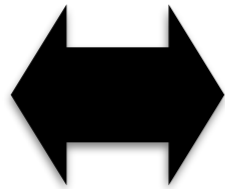
$$C_{(12)(12)} > 0 \left\{ \begin{array}{l} \text{scalar: } \sum_a \frac{1}{3840\pi|m_a|} \left[ z_{1a}^2 z_{2a}^2 - \frac{2}{7} z_{1a}^2 - \frac{2}{7} z_{2a}^2 \right] > 0 \\ \text{fermion: } \sum_a \frac{1}{480\pi|m_a|} \left[ z_{1a}^2 z_{2a}^2 - \frac{3}{8} z_{1a}^2 - \frac{3}{8} z_{2a}^2 \right] > 0 \end{array} \right.$$

$\mathcal{O}(z^0) \rightarrow 0$

Presence of **bifundamentals** is crucial for IR consistency

# A stronger CHC (3d/4d)

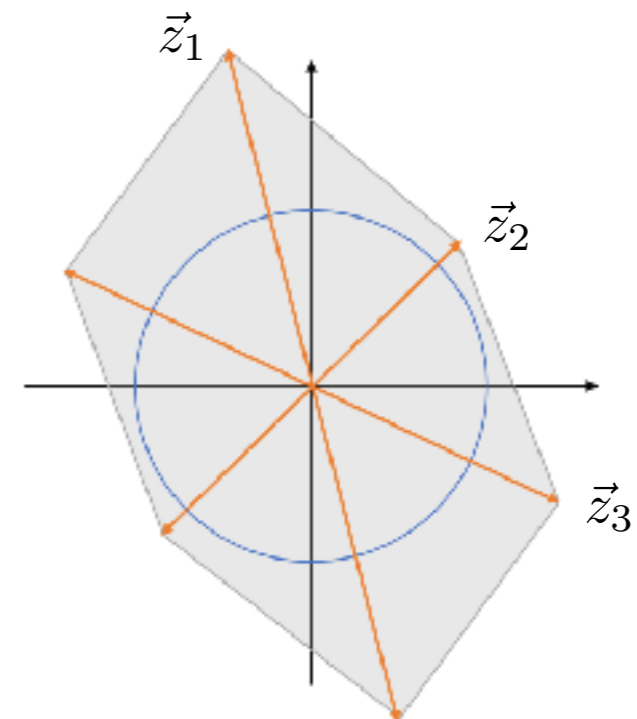
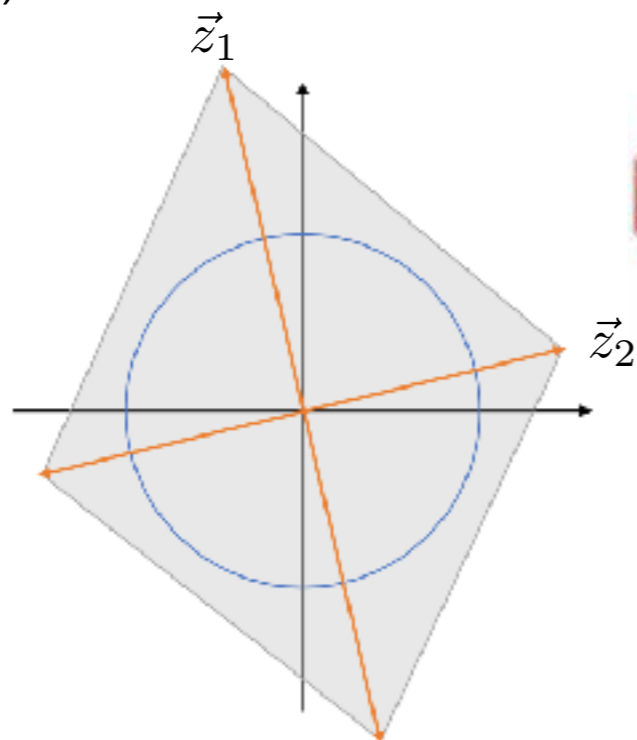
Positivity conditions



1) CHC is satisfied

2) There must exist (at least) a bifundamental particle for any (orthogonal) basis choice for the U(1)'s gauge fields

E.g., U(1)xU(1)



# Dimensional reduction

Is the WGC **consistent** under dimensional reduction ?

[B. Heidenreich, M. Reece, T. Rudelius 1509.06374]

**D** dim: 1 particle, single U(1)  $z_0 \geq 1$  WGC ✓

$S^1_{(r)}$  ↓

**D-1** dim: KK tower,  $U(1) \times U(1)_{\text{KK}}$  WGC (CHC) ? **NOT** always!  
(problem  $r \rightarrow 0$  limit)

Unless...



**A super-extremal particle  $z > 1$  should exist for every charge in the charge lattice**



Using **positivity bounds** ?

Same problem in D-1 for  $r \rightarrow 0$

$$\vec{Z}_{(n)} = (z_F, z_{KK}) = \left( \frac{q}{\sqrt{m^2 + n^2/r^2}}, \frac{n}{\sqrt{m^2 r^2 + n^2}} \right)$$

- lowest mode (n=0):  $(z_F, z_{KK}) = (q/m, 0)$

- KK modes (n≠0):  $(z_F, z_{KK}) \sim (0, 1)$

Absence of bifundamentals



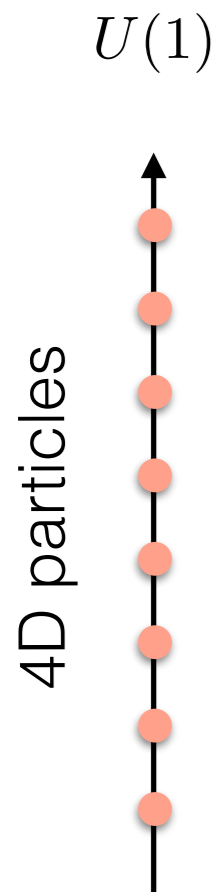
**IR inconsistent**

*Solution...*

## Tower-WGC\*

(\*in absence  
of fermions)

Replace the 4D field with a **tower** of 4D fields  $\Phi_l$  charged under  $U(1)$  with masses and charges  $(m_l, q_l)$  s.t. **bifundamental** contributions (at any  $r$ ) saves IR cons.

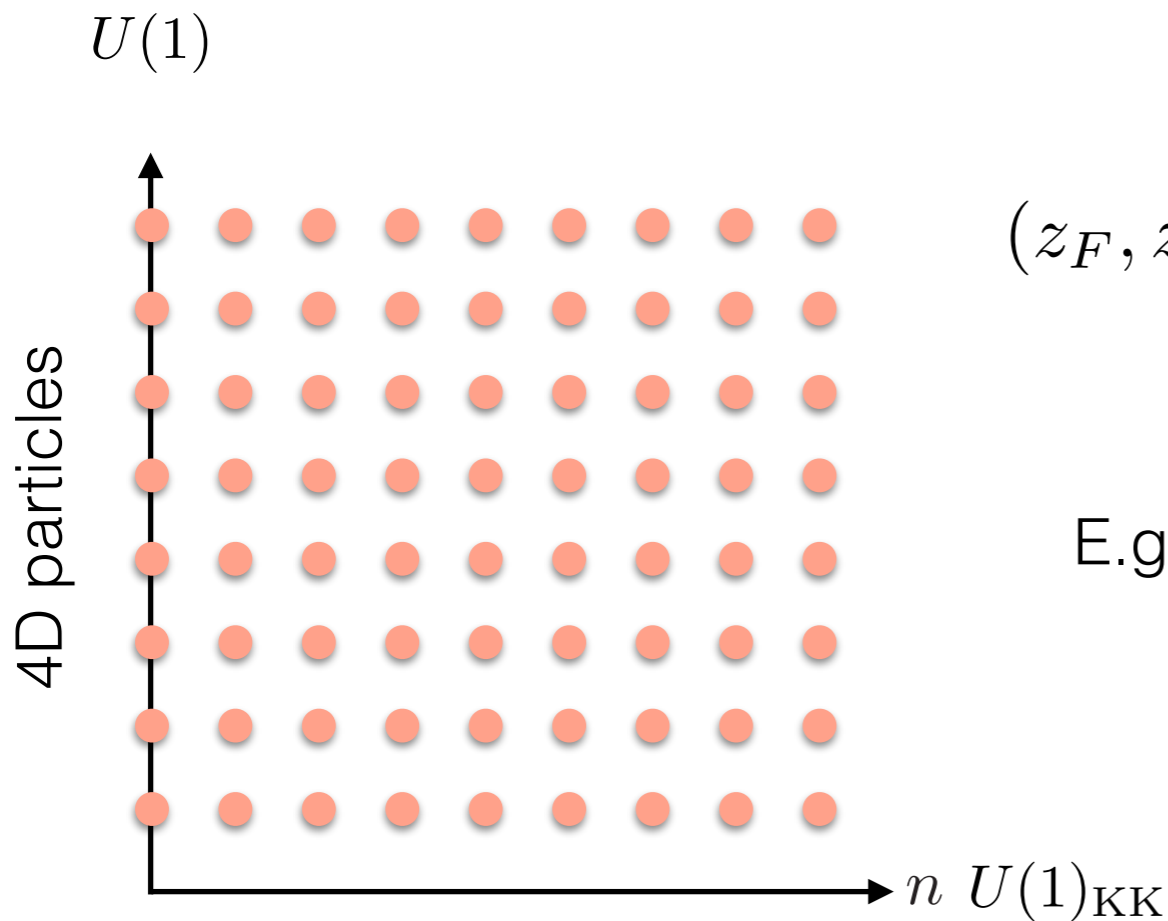




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$$(z_F, z_{KK}) = \left( \frac{q_l}{\sqrt{m_l^2 + n^2/r^2}}, \frac{n}{\sqrt{m_l^2 r^2 + n^2}} \right)$$

E.g.:  $m_l = \sqrt{m^2 + l^2 \mu^2}$        $q_l = (l + 1)q$



(depending on  $\mu$ )

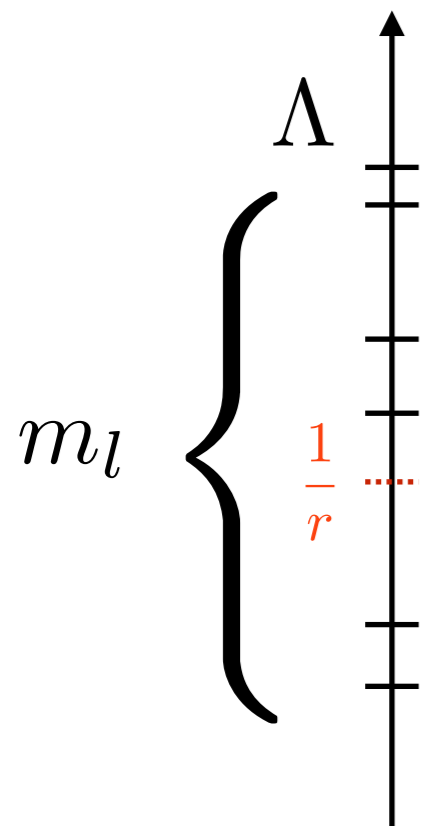
there may be bifundamentals at small radii: IR-OK

*Many other possibilities...*

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more precisely:



- There must exist particles with mass near the cut-off  $m_l \lesssim \Lambda$
- Such particles must have  $z_l \gtrsim \mathcal{O}(1)$
- In case the lightest particle has mass  $m \ll r^{-1}, \Lambda$  then the number of particles in the tower is *at least* of order  $(mr)^{-1}$

**TWGC** is *weaker* than the **LWGC**

no counterexample  
+ agreement with swampland distance conjecture

# Conclusions

## Relation IR consistency - WGC ?



Clear connection when **UV contribution** to the EFT in a certain regime. Possibly, compute UV contributions and check...



Existence of **bifundamentals** is crucial for multiple  $U(1)$ 's:

# stronger CHC

# under KK reduction: necessity of a tower of particles in the parent D-dimensional theory (**TWGC**)

# **TWGC** weaker than (sub-)LWGC and agrees with literature

