



Bethe Center for  
Theoretical Physics



# Swampland Distance Conjecture for One-Parameter Calabi-Yau Manifolds

## A Test

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Based on arXiv:1903.00596 with Albrecht Klemm.

# Outline

- 1 Swampland Conjectures
  - Swampland Distance Conjecture
- 2 The Setup
  - Complex Structure Moduli Space  $M_{CS}$
  - BPS-States
  - One-Parameter Calabi-Yau Threefolds
- 3 Results
  - Period Calculation
  - Metric
  - Tower of States
  - Tower of States (Stability)
- 4 Conclusions

- 1 Swampland Conjectures
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# Swampland Conjectures

These are statements about the properties of a consistent quantum gravity effective field theory.

A consistent theory is a theory with a (string theoretic) UV completion.

**[Vafa, 2005]**



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- 1 Swampland Conjectures
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  - BPS-States
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  - Period Calculation
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## Swampland Distance Conjecture

Consider a point  $p_0$  on the moduli space defining an effective theory,  $\mathcal{L}_0$ . For another point  $p$ , the effective theory  $\mathcal{L}$  compared to  $\mathcal{L}_0$  develops a tower of exponentially light states with mass of the order

$$e^{-\alpha d(p_0, p)} ,$$

where  $d(p_0, p)$  is the geodesic distance. As this distance diverges, one finds infinitely many massless states leading to breakdown of the effective theory.

The parameter  $\alpha$  is arbitrary and conjectured to be  $\mathcal{O}(m_{Pl})$  in the refined version.

- 1 Swampland Conjectures
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**Special geometry** is a powerful toolkit to study  $M_{CS}$ .

The fundamental object of interest is the **period vector**  $\Pi$ .

$$\Pi = \begin{pmatrix} F_I(z) \\ X^I(z) \end{pmatrix} = \begin{pmatrix} \int_{B_I} \Omega(z) \\ \int_{A^I} \Omega(z) \end{pmatrix} .$$

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This allows us to find the **Kähler potential** and the **metric** on  $M_{CS}$  via

$$K(z, \bar{z}) = -\log(i\Pi^\dagger \Sigma \Pi) ,$$

$$G_{i\bar{j}} = \partial_i \bar{\partial}_{\bar{j}} K(z, \bar{z}) ,$$

where

$$\Sigma = \begin{pmatrix} 0 & \mathbb{1} \\ -\mathbb{1} & 0 \end{pmatrix} .$$

- 1 Swampland Conjectures
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- 2 The Setup
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- 3 Results
  - Period Calculation
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  - Tower of States
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The mass of the BPS 3-brane [Ceresole, D'Auria, Ferrara, Van Proeyen, 1995]

$$m_{\mathbf{q}} = |Z_{\mathbf{q}}| ,$$

where

$$\begin{aligned} Z_{\mathbf{q}} &= e^{K/2} \int_S \Omega \\ &= e^{K/2} \mathbf{q}^T \cdot \Sigma \cdot \Pi \end{aligned}$$

is called the central charge.



# BPS-States

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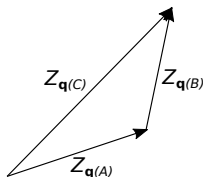
- 2 Are they **stable** against decay?

One uses following argument to **test** the stability.

Consider three 3-branes,  $A$ ,  $B$  and  $C$

$$Z_{\mathbf{q}(C)} = Z_{\mathbf{q}(A)} + Z_{\mathbf{q}(B)} ,$$

$$m_{\mathbf{q}(C)} \leq m_{\mathbf{q}(A)} + m_{\mathbf{q}(B)} .$$



Hence, if  $\arg(Z_{\mathbf{q}(A)}) = \arg(Z_{\mathbf{q}(B)})$ ,  $C$  can decay to  $A$  and  $B$  or vice-versa. Otherwise,  $C$  will be **stable**.

- 1 Swampland Conjectures
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- 2 The Setup
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We can get one-parameter Calabi-Yau manifolds by considering complete intersection of hyper-surfaces in weighted projective space  $\mathbb{P}^n$ .

The most interesting case is that of two cubics in  $\mathbb{P}^5$ , let's call it  $Y$ . It's mirror  $X$  is given by zero locus of [Libgober, Teitelbaum, 1993]

$$P_1 = x_1^3 + x_2^3 + x_3^3 - z^{-\frac{1}{6}} x_4 x_5 x_6 \text{ and}$$

$$P_2 = x_4^3 + x_5^3 + x_6^3 - z^{-\frac{1}{6}} x_1 x_2 x_3 ,$$

where  $z$  parameterises  $M_{CS}$ .

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# Period Calculation

The period vector can be calculated using the [Picard-Fuchs equation](#)

$$\left(\theta^4 - 3^6 z (\theta + 1/3)^2 (\theta + 2/3)^2\right) f_i(z) = 0 ,$$

where  $\theta = z \frac{d}{dz}$ .

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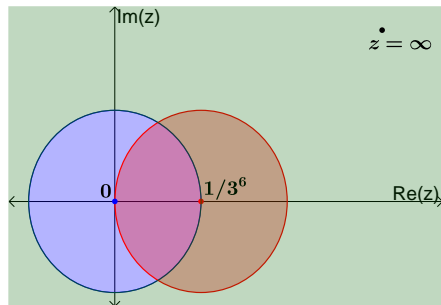
where  $\theta = z \frac{d}{dz}$ .

$$\tilde{\Pi} = \begin{pmatrix} f_1(z) \\ f_2(z) \\ f_3(z) \\ f_4(z) \end{pmatrix} \xrightarrow{\text{base change}} \Pi = T \tilde{\Pi} = \begin{pmatrix} F_0(z) \\ F_1(z) \\ X^0(z) \\ X^1(z) \end{pmatrix} .$$

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Three singular points

$$z = \left\{ 0, \frac{1}{3^6}, \infty \right\}.$$

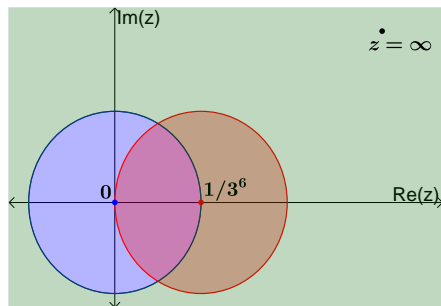


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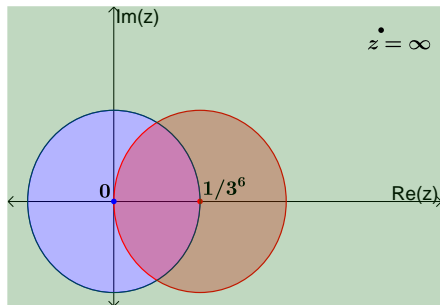
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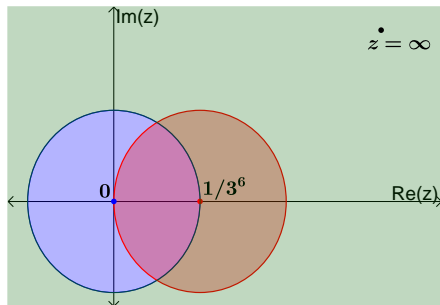
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We solve the differential equation around other singular points to find  $\tilde{\Pi}(z)$ .

Using analytic continuation, we match them to get symplectic basis

$$\tilde{\Pi}(z) \rightarrow \Pi(z).$$

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  - Swampland Distance Conjecture
  
- 2 The Setup
  - Complex Structure Moduli Space  $M_{CS}$
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- 3 Results
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# Results: Metric

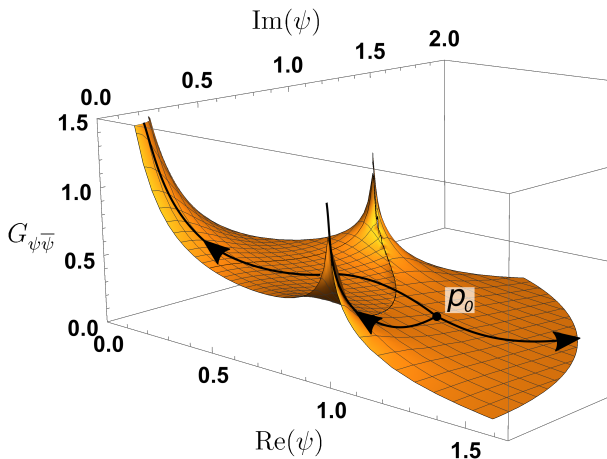
Define:  $\psi = \frac{1}{3z^{1/6}}$  .      Singular points:  $\overbrace{\left\{0, \frac{1}{3^6}, \infty\right\}}^z \rightarrow \overbrace{\{\infty, 1, 0\}}^\psi$  .

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Define:  $\psi = \frac{1}{3z^{1/6}}$  .

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The resulting **metric**  
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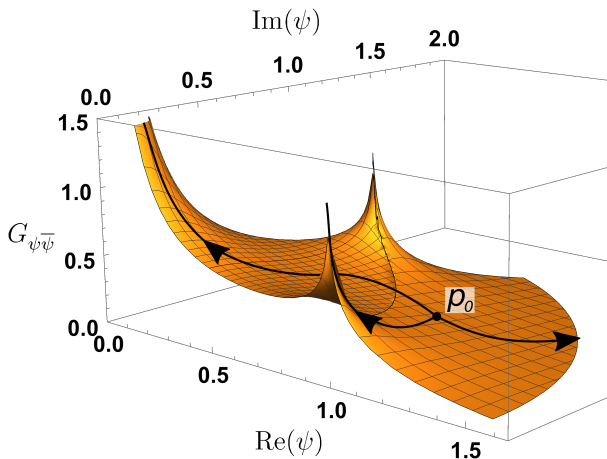
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$\psi$	$G_{\psi\bar{\psi}}$	Distance
0	$\infty$	$\infty$
1	$\infty$	Finite
$\infty$	0	$\infty$



## Results: Metric

Two infinite distance points, two infinite tower of states according to Swampland Distance Conjecture.

$$\boxed{\psi \rightarrow 0} \quad \text{and} \quad \boxed{\psi \rightarrow \infty}$$

- 1 Swampland Conjectures
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- 2 The Setup
  - Complex Structure Moduli Space  $M_{CS}$
  - BPS-States
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- 3 Results
  - Period Calculation
  - Metric
  - **Tower of States**
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- 4 Conclusions

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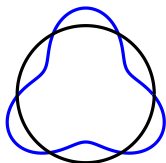
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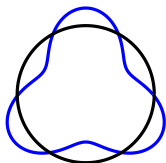
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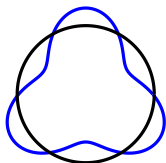
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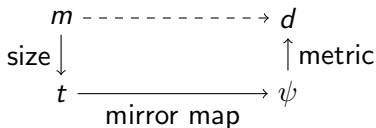
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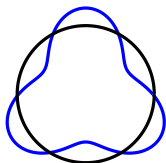
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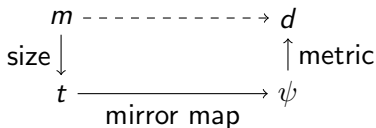
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$$m \sim m_0 e^{-\alpha d}, \quad \alpha = \frac{1}{\sqrt{3}}$$

[Blumenhagen, Klaewer, Schlechter, Wolf, 2018]

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There exists  $\mathbf{q}$ :

$$\mathbf{q} = \begin{pmatrix} 3n \\ 3m - 4n \\ m \\ n \end{pmatrix}, \quad n, m \in \mathbb{Z}$$

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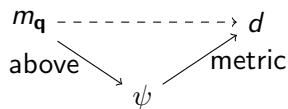
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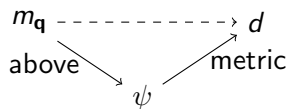
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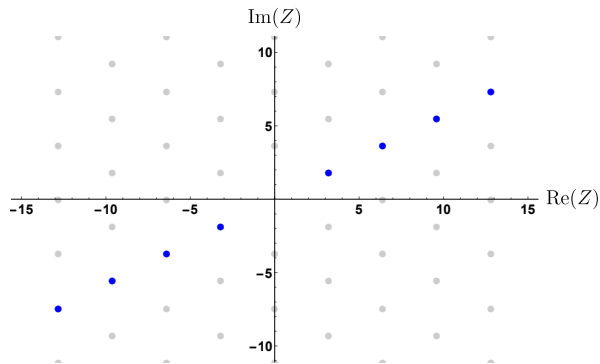
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For stability, consider the set of all possible  $\mathbf{q}$  which leads to all possible  $Z_{\mathbf{q}}$ .

- Consider this subset.

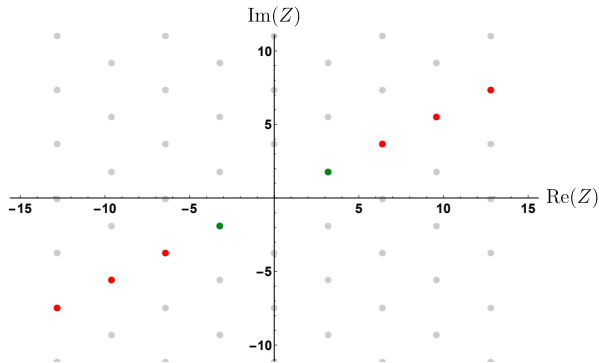


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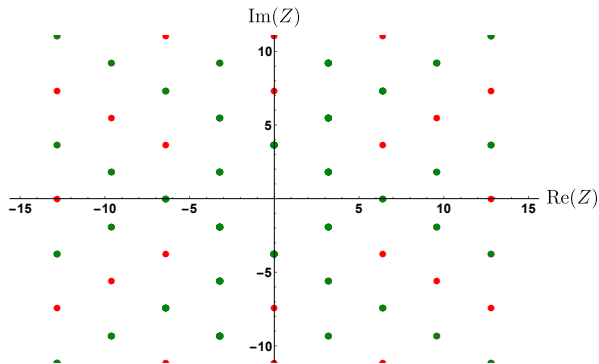


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- Green: infinite set of stable states.



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  - Swampland Distance Conjecture
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In principle, [four types of singularities](#) can occur specified by their limiting mixed Hodge structure. Two of them, called  $K(a,a,b,b)$  and  $M(a,a,a,a)$ , are at infinite distance.

We studied all 14 cases and the one presented here is one of the interesting ones since it has [both type of infinite distance singularities](#).  $K$  at  $\psi = 0$  and  $M$  at  $\psi = \infty$ .

# Conclusions

We find a **charge sub-lattice** of infinite light BPS-states at these infinite distance points (In addition to KK states at large volume point.) showing the **validity of Swampland Distance Conjecture**.

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**Thank you for listening. Questions?**