

Carving out the Swampland of 6D SCFTs

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[arXiv :1907.xxxxx](#)

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6D SCFT and Geometry

$6D$, $\mathcal{N} = (\geq 1, 0)$ SCFTs occupy a very special place in modern QFT [Nahm, '78]
From a Stringy point of view, SCFTs usually studied through their tensor branch.

- ▶ $B_{\mu\nu} \leftrightarrow$ String

Smoking gun for SCFTs : Tensionless strings.

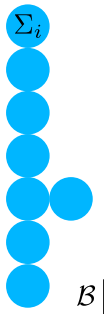
- ▶ F-theory : 6D strings from D3-branes wrapping $\Sigma_i \subset \mathcal{B}$.

- ▶ CFT point reached when contracting all curves to a point

$$T = \text{Vol}(\Sigma) \rightarrow 0 \quad \Sigma_i \cdot \Sigma_j \preceq 0$$

- ▶ All possible bases constructed gluing NHCs.
→ **Generalised quivers**

$$\mathcal{B} = \mathbb{C}^2/\Gamma, \quad \Gamma \subset U(2)$$



[Heckman, Morrison, Tomasiello, Rudelius, Vafa, '13-15]

What really happens at the SCFT point ?

Spectrum along the tensor branch reorganises into **superconformal multiplets**.

$$\begin{array}{ccc} SO(5,1) \times Sp(\mathcal{N}) & \longrightarrow & SO(6,2) \times Sp(\mathcal{N}) \\ \Phi, A_\mu, T_{\mu\nu} & \longrightarrow & ??? \end{array}$$

CFT operators are **gauge neutral**, e.g. $\mathcal{O}_k = \text{Tr}_G(\Phi^k)$.

- ▶ Goal : Use **unitarity** to obtain constraints **directly** at CFT point.

Possibly exclude theories not captured by the classification.

Dream : Exclude all spectra but countable set.

- ▶ Eliminate theories that do not even satisfy unitarity constraints (swampland of the swampland), e.g.

$$C_T \geq 0, \quad \Delta \geq \Delta_{\text{bound}}(\ell) = \begin{cases} 2, & \ell = 0 \\ \ell + 4, & \ell > 0 \end{cases}$$

- ▶ String Theory : unitarity very powerful (anomaly cancellation, $\Sigma_i \cdot \Sigma_j \preceq 0 \dots$)

The Conformal Bootstrap : New twist on an old Idea

- ▶ OPE : reduces 4-point functions as a sum of **conformal blocks**

$$\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle = \frac{1}{(x_{12} x_{34})^{\Delta_\phi}} \sum_{\mathcal{O} \in \phi \times \phi} \lambda_{\phi\phi\mathcal{O}}^2 g_{\Delta_{\mathcal{O}}, \ell_{\mathcal{O}}}(u, v)$$

Casimir Equation [Dolan, Osborn, '04+'11]

$$\mathcal{D} g_{\Delta, \ell} = c g_{\Delta, \ell}, \quad c = c_{\Delta, \ell} \in \mathbb{R}^+$$

In even dimension, $g_{\Delta, \ell}$ has known closed form.

The Conformal Bootstrap : New twist on an old Idea

- ▶ OPE should be **associative**

$$\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle = \begin{array}{c} \phi(x_1) \\ \diagdown \\ \text{---} \mathcal{O} \text{---} \\ \diagup \\ \phi(x_2) \end{array} \begin{array}{c} \phi(x_3) \\ \diagup \\ \text{---} \\ \diagdown \\ \phi(x_4) \end{array} = \begin{array}{c} \phi(x_1) \\ \diagdown \\ \text{---} \mathcal{O} \text{---} \\ \diagup \\ \phi(x_2) \end{array} \begin{array}{c} \phi(x_3) \\ \diagdown \\ \text{---} \\ \diagup \\ \phi(x_4) \end{array}$$

Sum Rule [Rattazzi,Rychkov,Tonni,Vichi,'08]

$$\sum_{\mathcal{O} \in \phi \times \phi} \lambda_{\mathcal{O}}^2 \underbrace{\left(v^{\Delta_{\phi}} g_{\Delta_{\mathcal{O}}, \ell_{\mathcal{O}}}(u, v) - u^{\Delta_{\phi}} g_{\Delta_{\mathcal{O}}, \ell_{\mathcal{O}}}(v, u) \right)}_{F_{\mathcal{O}}(u, v)} = 0$$

- ▶ By unitarity $\lambda_{\mathcal{O}} \in \mathbb{R}$.

Consider would-be spectrum $\{\Delta_{\mathcal{O}}, \lambda_{\mathcal{O}}\}$, if we can find a functional such that

$$\alpha(F_{\mathcal{O}}) > 0, \quad \forall \ell, \Delta \geq \Delta_{\text{bound}} \quad \Rightarrow \quad \sum_{\mathcal{O}} \underbrace{\text{const}}_{>0} = 0 \quad \text{CONTRADICTION!}$$

→ **semi-definite program** (linear optimisation). \exists standard tools (SDPB).

[Poland, Simmons-Duffin, Vichi, '12; Simmons-Duffin, '15]

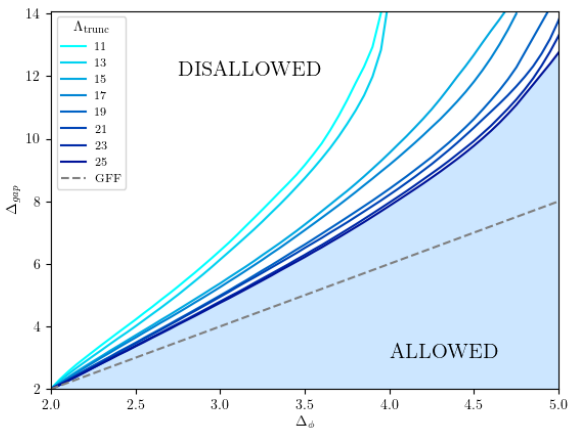
In practice : $\alpha_{\text{trunc}} = \sum_{m+n \leq \Lambda_{\text{trunc}}} \alpha_{m,n} \partial_u^m \partial_v^n F_{\mathcal{O}}|_{u=v}$

► As $\Lambda \rightarrow \infty$, more and more spectra can be excluded.

► Consider gapped spectrum

$$\Delta \geq \Delta_{\text{gap}} + 2, \quad \forall \text{ scalars}$$

► Only deviations from GFF for non-SUSY w/ $\Delta < 4$.



Generalised free theories :

$$\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle = \langle \phi(x_1)\phi(x_2) \rangle \langle \phi(x_3)\phi(x_4) \rangle + (\text{perms.}), \quad \langle \phi(x)\phi(y) \rangle = \frac{1}{|x-y|^\Delta}$$

Also used to bound OPE coefficients. Singling out energy-momentum tensor $T_{\mu\nu}$:

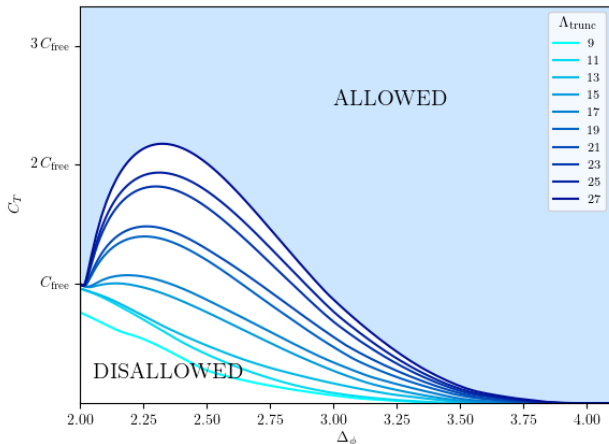
$$\lambda_{T_{\mu\nu}}^2 \underbrace{\alpha(F_{T_{\mu\nu}})}_{=1} = -\alpha(F_{\mathbb{1}}) - \sum_{\{\mathcal{O}\} \setminus \{\mathbb{1}, T_{\mu\nu}\}} \lambda_{\mathcal{O}}^2 \alpha(F_{\mathcal{O}}) \leq -\alpha(F_{\mathbb{1}})$$

► Maximise $\alpha(F_{\mathbb{1}})$

► Get lower bound :

$$\lambda_{T_{\mu\nu}}^2 = \frac{5}{6} \frac{\Delta_\phi^2}{C_T} \leq -\alpha(F_{\mathbb{1}})$$

► Weakest where supermultiplets appear.



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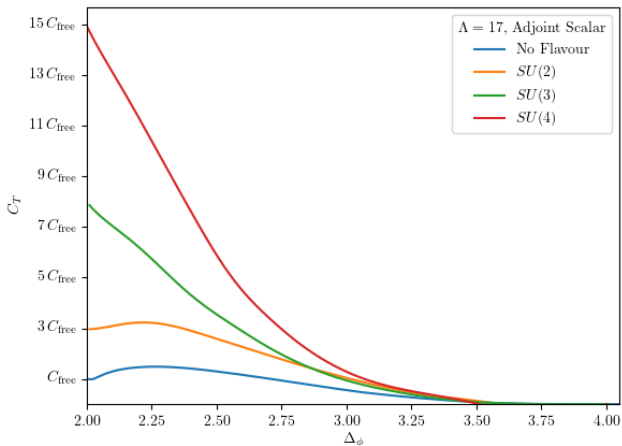
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Bootstrapping \mathcal{D} -multiplets

While the bootstrap program has grown into a field of physics of its own, the literature on 6D SCFT is sparse, but has yielded interesting results

- ▶ $\mathcal{N} = (2, 0)$ with $\mathcal{D}[2, 0] \ni T_{\mu\nu}$ [Beem,Lemos,Rastelli,van Rees,'16]

$$C_T^{\min} \rightarrow 25, \Lambda_{\text{trunc}} \rightarrow \infty \quad (A_1 \text{ theory!})$$

- ▶ $\mathcal{N} = (1, 0)$ with $\mathcal{D}[1] \ni J_\mu$ [Chang,Lin,'17]

$$C_J^{\min} \rightarrow 150, \Lambda_{\text{trunc}} \rightarrow \infty \quad (\text{E-String!})$$

Superblocks satisfy a **super Casimir equation**

$$(\mathcal{D}_{\text{bos}} + \mathcal{D}_{\text{SUSY}}) G_\chi = (c_{\text{bos}} + c_{\text{SUSY}}) G_\chi, \quad G_\chi = \sum_{m,n} f_{m,n} g_{\Delta+m, \ell+n}$$

Can use bosonic Casimir equation to find prefactors.

Bootstrapping \mathcal{D} -multiplets

\mathcal{D} -multiplets correlators most constrained (shortness conditions)

\Rightarrow Casimir equation simpler ($Q^n \mathcal{O} = 0$ for some operators inside multiplet)

► Generalises known results for $\mathcal{D}[1]$ to arbitrary $\mathcal{D}[J_r]$ ($\forall 2 < D \leq 6$ w/ $8Q$)
[Chang, Lin, '17 ; Bobev, Lauria, Mazáč, '17]

► Also for 4-pt functions involving different $\mathcal{D}[J_R]$

\rightarrow Enables bootstrap for **mixed** correlators.

► Will put strong constraints on theories without known F-theory realisation

$$G_F = SU(3) \quad \text{with} \quad C_T = \frac{19488}{5}, \quad C_J = \frac{195}{5}$$

(stay tuned!)

[Shimizu, Tachikawa, Zafir'17]

Conclusion

- ▶ The bootstrap is a very powerful tool to study SCFTs
- ▶ Bootstrap “knows” about supersymmetry and String theory!
- ▶ We found the blocks for 4-pt function of any \mathcal{D} -multiplets

A lot of exciting direction to explore :

- ▶ Relate the bounds to **geometric quantities**?
- ▶ How much freedom is there for **non F-theoretic SCFTs**?
- ▶ Are SCFTs embeddable into compact geometries special (**swampland**)?
- ▶ ...

Thank you for your attention !