

Severin Lüst

CEA Saclay and École Polytechnique

# Uplifting Runaways

with I. Bena, E. Dudas, and M. Graña [arXiv:1809.06861]

June 27, 2019

String Phenomenology 2019, CERN

# de Sitter vacua in string theory

Three-step procedure [Kachru, Kallosh, Linde, Trivedi '03]:

1. warped IIB with CS-moduli stabilized by three-form fluxes including a region with strong warping [Giddings, Kachru, Polchinski '01]  
described by the Klebanov Strassler throat [Klebanov, Strassler '00]  
→ large hierarchy of scales
2. Stabilize Kähler moduli by non-perturbative effects  
→ supersymmetric AdS-vacuum
3. Supersymmetry breaking by an  $\overline{D3}$ -brane at the bottom of the throat  
→ exponentially suppressed uplift to dS due to strong warping

(See also many other talks during this conference.)

# de Sitter vacua in string theory

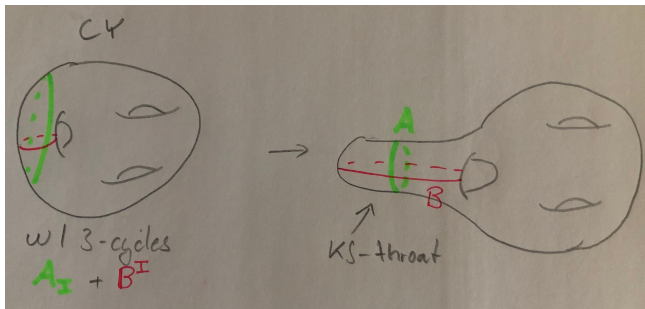
Three-step procedure [Kachru, Kallosh, Linde, Trivedi '03]:

1. warped IIB with CS-moduli stabilized by three-form fluxes including a region with strong warping [Giddings, Kachru, Polchinski '01]  
described by the Klebanov Strassler throat [Klebanov, Strassler '00]  
→ large hierarchy of scales
2. Stabilize Kähler moduli by non-perturbative effects  
→ supersymmetric AdS-vacuum
3. Supersymmetry breaking by an  $\overline{D3}$ -brane at the bottom of the throat  
→ exponentially suppressed uplift to dS due to strong warping

(See also many other talks during this conference.)

# Warped CY

- Metric:  $ds_{10}^2 = e^{2A} ds_4^2 + e^{-2A} ds_{CY_3}^2$
- Fluxes fix the sizes of the 3-cycles:  $\int_{A^I} F_3 = M^I, \int_{B^I} H_3 = K_I$



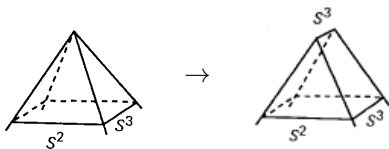
- Choose a configuration such that one cycle is exponentially large.
- Klebanov-Strassler throat.

# Deformed conifold

- In the region of high warping, the six-dimensional geometry is given by the deformed conifold.
- embedding of the deformed conifold into  $\mathbb{C}^4$ :

$$\sum_{a=1}^4 z_a^2 = S.$$

- Replace the singularity of the conifold ( $S = 0$ ) by a  $S^3$  of size  $|S|$

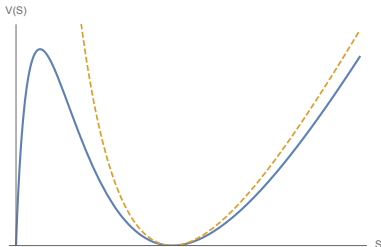


[Candelas, Ossa '89]

- $S$  is a complex structure modulus of the deformed conifold.

# Potential for $S$

- Fluxes  $M$  and  $K$  along the two three-cycles of the conifold generate a potential  $V_{KS}(S)$  [Douglas, Shelton, Torroba '07, '08]:



(dotted without warping effects)

- (Supersymmetric) minimum at  $s_{KS} = \Lambda_0^3 \exp\left(-\frac{2\pi K}{g_s M}\right)$ .
  - Relative warp factor:  $\Lambda_0/\Lambda_{IR} \sim |s_{KS}|^{\frac{1}{3}}$ .
- Large hierarchy for suitable values of  $K$ ,  $M$ , and  $g_s$  [Giddings et al. '01].

# Mass of $S$

- The mass of  $S$  at the minimum  $s_{KS}$  can be computed by

$$m_S^2 \equiv \frac{1}{M_{pl}^2} G^{S\bar{S}} \partial_S \partial_{\bar{S}} V \Big|_{S=s_{KS}}$$

- Including the effects of the warping we find:

$$m_S^2 \sim \frac{s_{KS}^{2/3}}{\alpha'^2}$$

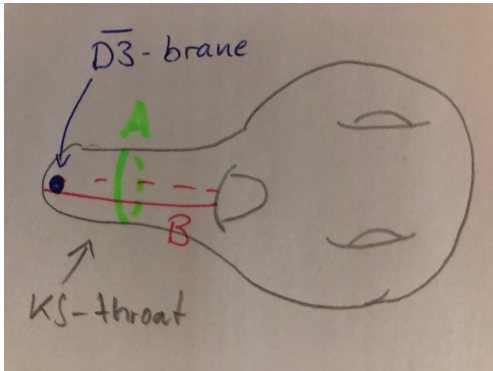
(c.f. [Blumenhagen, Herschmann, Wolf '16] for  $m_S^2$  without warping)

- If  $s_{KS}$  is exponentially small,  $S$  becomes exponentially light.
- $S$  cannot be integrated out before uplifting with an anti-brane.

Comparison with Kähler moduli masses: [Blumenhagen, Kläwer, Schlechter '19]

# $\overline{D3}$ -brane in the KS throat

- Place an anti-D3 brane at the bottom of the throat



- Positive contribution to the energy  $\rightarrow$  uplift to de Sitter



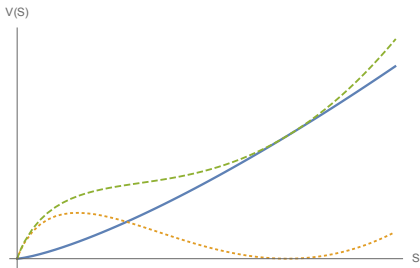
# $\overline{D3}$ -brane in the KS throat

- The  $\overline{D3}$ -brane gives a contribution to the potential:

$$V_{\overline{D3}}(S) \propto e^{4A} \propto \frac{|S|^{4/3}}{(\alpha' g_s M)^2}$$

with  $e^{4A}$  the warp factor of the Klebanov-Strassler solution.

- Plot of the potential:



(dotted lines represent the KS potential and their superposition)

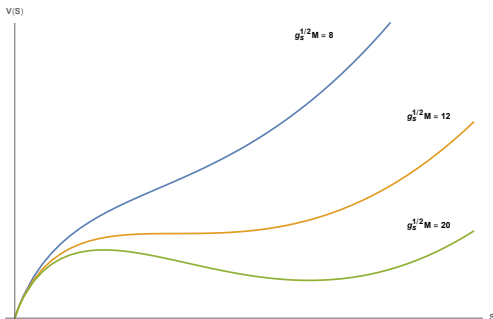
# Stability with one $\overline{D3}$ -brane

- A stable minimum of  $V_{KS} + V_{\overline{D3}}$  with  $S > 0$  exists iff

$$g_s M^2 > M_{min}^2 \quad \text{with} \quad M_{min} \approx 12.$$

(see also [Blumenhagen, Kläwer, Schlechter '19])

- Superposition of the potentials:



# Implications on the maximal hierarchy

- Warping creates a hierarchy of scales

$$h = 3 \ln \frac{\Lambda_0}{\Lambda_{IR}} = \frac{2\pi K}{g_s M}$$

- Tadpole cancellation:

$$M^l K_l + Q_3^{loc} = 0,$$

where  $Q_3^{loc}$  is the D3-charge of localized sources.

- Stability of the KS throat + tadpole cancellation:

$$h = 2\pi \frac{MK}{g_s M^2} < 2\pi \frac{|Q_3^{loc}|}{M_{min}^2} \approx \frac{|Q_3^{loc}|}{23}$$

# Tadpole cancellation in IIB

- For CY orientifolds with O3-planes and D3-branes:

$$Q_3^{loc} = N_{D3} - \frac{1}{4} N_{O3}$$

- Largest number of O3-planes:  $T^6/\mathbb{Z}_2$ :  $Q_3^{loc} \leq 32$

→ No large hierarchy possible.

- O7-planes and D7-branes:

$$Q_3^{loc} = \frac{1}{24} \chi(D7) + \frac{1}{6} \chi(O7) - (\text{gauge})$$

- $\chi$ : Euler number of the 4-cycles wrapped by the D7s/O7s.

→ Large tadpole possible, but D7-moduli need to be stabilized.

# Tadpole cancellation in F-theory

- Tadpole cancellation for F-theory on a Calabi-Yau four-fold  $CY_4$  with four-form flux  $G$ :

$$N_{D3} + \frac{1}{2} \int G \wedge G = \frac{\chi(CY_4)}{24}$$

- $\chi(CY_4)$ : Euler number of the CY  $\rightarrow$  can be very large  
(largest known example [Klemm et al. '97]:  $\chi = 1\,820\,448 = 24 \cdot 75\,852$ )
- But: Large  $\chi$  implies a lot of moduli:

$$\chi(CY_4) = 6(8 + h^{1,1} + h^{3,1} - h^{2,1})$$

- $h^{3,1}$ : complex structure of  $CY_4 \rightarrow$  must be stabilized by flux:

$$\int G \wedge G = \mathcal{O}(h^{3,1}) ?$$

# Conclusions

- With a large hierarchy the KS-modulus becomes exponentially light.
- One  $\overline{D3}$  makes a Klebanov-Strassler throat unstable unless  $g_s^{1/2} M > 12$ .
- Due to tadpole-cancellation: Constraints on the hierarchy.

# Conclusions

- With a large hierarchy the KS-modulus becomes exponentially light.
- One  $\overline{D3}$  makes a Klebanov-Strassler throat unstable unless  $g_s^{1/2} M > 12$ .
- Due to tadpole-cancellation: Constraints on the hierarchy.

Thank You!