

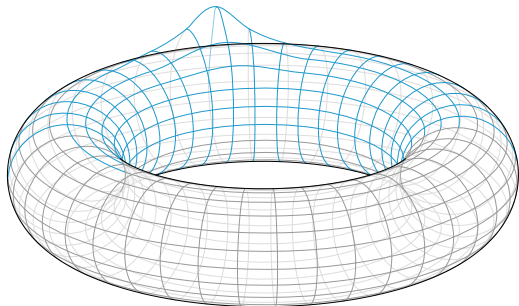
INSTANTONS AWAY FROM EQUILIBRIUM

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in collaboration with

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[1905.00219]



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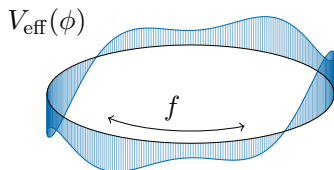
String Pheno 2019

Motivation

Non-perturbatively generated effective potentials

are an often crucial ingredient in string phenomenology:

Axion Inflation, Axionic Dark Matter, Moduli Stabilization, dS, Strong CP, ...



The **gauge symmetry** $\phi \sim \phi + f$ protects effective potential $V_{\text{eff}}(\phi)$, the vacuum energy with ϕ a fixed parameter. Useful QG lamppost!

However, $V_{\text{eff}}(\phi)$ often used off-shell or **out-of-equilibrium**.

Main Question | When does $V_{\text{eff}}(\phi)$ fully capture the dynamics?

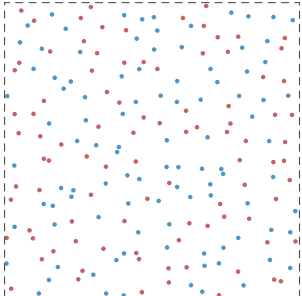
Mainly interested in **axion** or **natural inflation** [Freese, Frieman, Olinto '90]

$$S = \int d^4x \left(-\frac{f^2}{2} (\partial\varphi)^2 - \frac{1}{4g^2} F^2 + \frac{\varphi}{8\pi} F\tilde{F} + \mathcal{L}_{\text{cm}} \right)$$

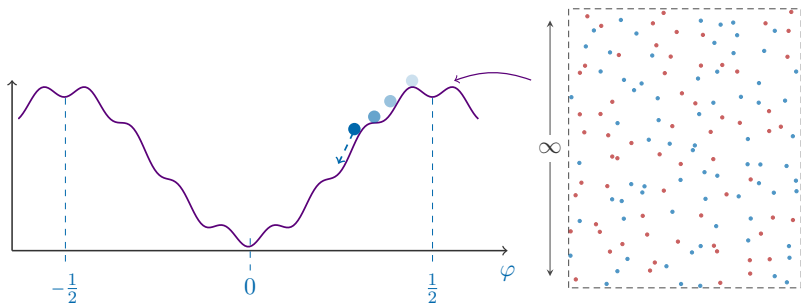
Key Ingredient | Both φ and A have topologically non-trivial field spaces:

$$\varphi \sim \varphi + 1 \quad A \sim A + \omega$$

Classically, $\varphi \rightarrow \varphi + \epsilon$ is a symmetry but broken by **instantons**.

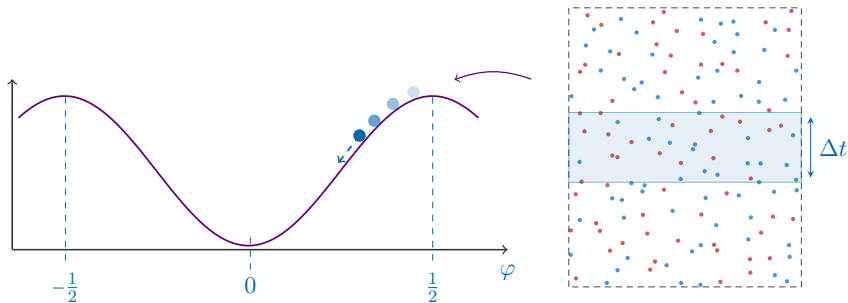
$$\int \mathcal{D}A e^{-S_A[\varphi]} \sim \beta \int \mathcal{D}A e^{-S_A[\varphi]} \sim \exp(-\beta V_{\text{eff}}(\varphi))$$


This potential drives inflation, determines φ 's semi-classical dynamics!



Main Question | When does $V_{\text{eff}}(\varphi)$ fully capture the dynamics?

But, at short times the system may not see all instantons!



Main Question | When does $V_{\text{eff}}(\varphi)$ fully capture the dynamics?

How slowly must φ move to see the effective potential?

Does it ever?

Toy Model

We want to mimic

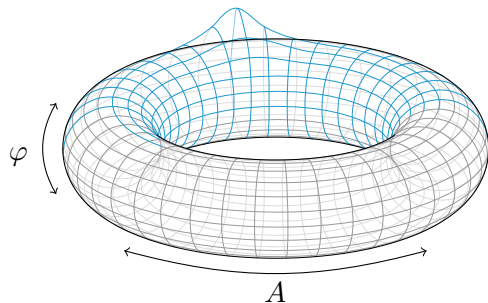
$$S = \int d^4x \left[-\frac{f^2}{2} (\partial\varphi)^2 - \frac{1}{4g^2} F^2 + \frac{\varphi}{8\pi} F\tilde{F} + \mathcal{L}_{\text{cm}} \right]$$

$$\begin{aligned} \varphi &\sim \varphi + 1 \\ &\text{and} \\ A &\sim A + \omega \end{aligned}$$

Consider a particle on a torus in an electromagnetic field,

$$S = \int dt \left[\frac{f^2}{2} \dot{\varphi}^2 + \frac{1}{2g^2} \dot{A}^2 + 2\pi\hbar\varphi\dot{A} - V_c(A) \right]_{\text{periodic}}$$

$$\begin{aligned} \varphi &\sim \varphi + 1 \\ &\text{and} \\ A &\sim A + 1 \end{aligned}$$



Interested in zero modes:

Like $\varphi(t, \mathbf{x}) = \varphi(t)$ and

$C_3 = A(t) dx \wedge dy \wedge dz$,

where $F \wedge F = dC_3$.

When \mathcal{A} remains "near" its ground state, how does φ behave?

Described by an **effective** Schrödinger equation,

$$i \partial_t \psi(\varphi, t) = \left(\frac{1}{2f^2} (p_\varphi - \mathcal{A}_B)^2 + V_{\text{eff}}(\varphi) + V_B(\varphi) \right) \psi(\varphi, t) \\ + \sum_{\substack{\text{excited} \\ \text{states, } n}} \left(\mathcal{F}_{0,n}(\varphi) (p_\varphi - \mathcal{A}_B) + V_{0,n}(\varphi) \right) \psi_n(\varphi, t)$$

The gauge field's vacuum energy acts as an
effective potential, $V_{\text{eff}}(\varphi)$

When A remains "near" its ground state, how does φ behave?

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Non-trivial **Berry potential** \mathcal{A}_B can act like a magnetic field,
and can (sometimes) be gauged away

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Berry scalar potential $V_B(\varphi)$ corrects $V_{\text{eff}}(\varphi)$

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Can not be seen from standard vacuum Euclidean path integral!

[Moody, Shapere, Wilczek '89]

$$V_B(\varphi) = \frac{1}{f^2} \sum_{\substack{\text{excited} \\ \text{states, } n}} \left| \frac{\langle 0 | \partial_\varphi \mathcal{H}_A | n \rangle}{V_{\text{eff}}(\varphi) - E_n(\varphi)} \right|^2$$

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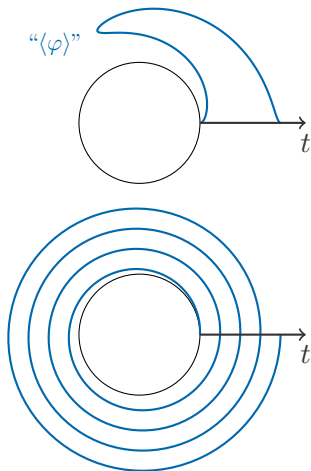
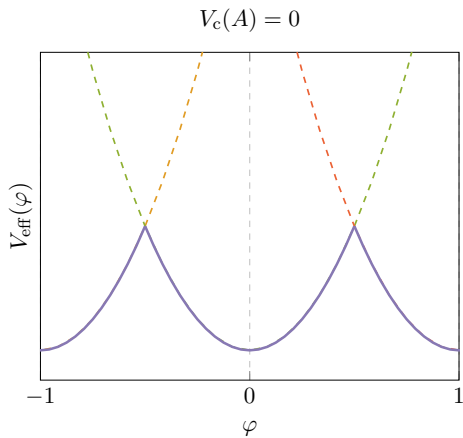
Coupling to other modes imposes an **axion speed limit**

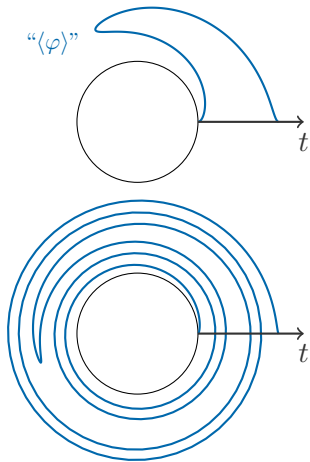
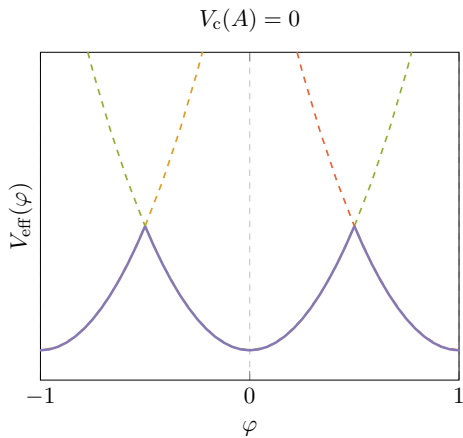
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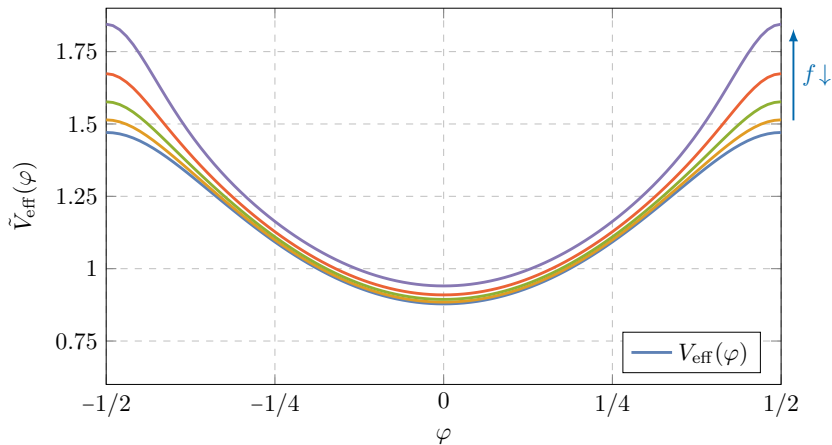
$$i \partial_t \psi(\varphi, t) = \left(\frac{1}{2f^2} (p_\varphi - \mathcal{A}_B)^2 + V_{\text{eff}}(\varphi) + V_B(\varphi) \right) \psi(\varphi, t) \\ + \sum_{\substack{\text{excited} \\ \text{states, } n}} \left(\mathcal{F}_{0,n}(\varphi) (p_\varphi - \mathcal{A}_B) + V_{0,n}(\varphi) \right) \psi_n(\varphi, t)$$

Coupling to other modes, **backreacts** on the gauge field





$$V_c(A) \propto 1 - \cos(2\pi A)$$



Conclusions

Main Question | When does $V_{\text{eff}}(\varphi)$ fully capture the dynamics?

- In this model, almost never
- Quantum adiabatic limit can be richer than classical, in ways that cannot be seen from the standard vacuum-to-vacuum Euclidean path integral
- These corrections can compete with the quantum mechanically generated effective potential

Future | How do we consistently incorporate non-perturbative effects in an EFT off-shell and away from equilibrium?

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Future | How do we consistently incorporate non-perturbative effects in an EFT off-shell and away from equilibrium?

Thanks!