SU(3) structures on Calabi-Yau manifolds

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Based on Larfors, Lukas, Ruehle 1805.08499, and work in progress

Why, what and how?

Motivation: 4D physics from string compactifications

- \bullet CY manifolds \longrightarrow large set of semi-realistic string vacua
- Still lack fully realistic compactifications: moduli, physical couplings, stability, cosmological constant,...

While CY geometry is useful it is not necessary.

This talk

- SU(3) structure ightarrow 4D $\mathcal{N}=1$ SUSY
- SUSY, BI, EOM constrain torsion
- Can we get a large class of example manifolds?

Idea:

- Construct explicit SU(3) structures on CY manifolds
- Bonus: get explicit metric
- How far can we get at satisfying all constraints?

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1 Motivation: 4D Heterotic $\mathcal{N} = 1$ Minkowski solutions

- 2 SU(3) structure
- 3 Construction of SU(3) structures on CY
- 4 SU(3) structures on CY: Torsion classes
- 5 SU(3) structures on CY: Strominger-Hull system
- 6 Conclusions and outlook

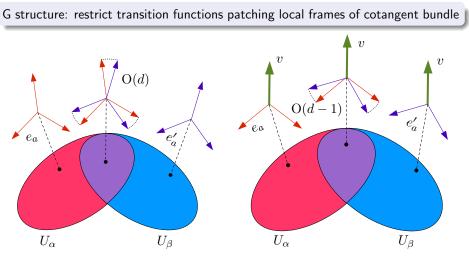
Motivation: 4D Heterotic $\mathcal{N} = 1$ Minkowski solutions

Candelas-Horowitz-Strominger-Witten:85,Strominger:86, Hull:86

Geometry			
SUSY equations, $H = 0$	\Rightarrow covariantly constant spinor η on X: $ abla \eta = 0$		
	$\iff X$ is Calabi–Yau		
SUSY equations, $H \neq 0$	\Rightarrow globally defined spinor η on X: $\nabla_T \eta = 0$		
	\iff SU(3) structure on X with torsion $T \sim H$		
	\iff SU(3) structure on X with torsion $T \sim H$		

Gauge field & vector bundle

SUSY equations \Rightarrow holomorphic vector bundle $V \rightarrow X$ with HYM connection Must also satisfy BI $dH = \frac{\alpha'}{4} (\operatorname{tr}(F \wedge F) - \operatorname{tr}(R^- \wedge R^-))$



picture from Koerber:10

$$\begin{split} \mathcal{M}_6 \text{ orientable with metric:} \quad & G = \mathsf{SO}(6) \subset \mathsf{GL}(6). \\ \mathcal{M}_6 \text{ spin:} & & \mathsf{SO}(6) \text{ lifts to } \mathsf{Spin}(6) \cong \mathsf{SU}(4). \end{split}$$

Let η Weyl, positive chirality: $\eta \in \mathbf{4}$ of SU(4). Choose basis:

$$\eta = \begin{pmatrix} 0 \\ 0 \\ 0 \\ \eta_0 \end{pmatrix} \text{ invariant under } \begin{pmatrix} U & \mathbf{0}_{3 \times 1} \\ \mathbf{0}_{1 \times 3} & 1 \end{pmatrix} \text{, } U \in \mathsf{SU}(3)$$

Globally defined $\eta \implies \mathsf{G} = \mathsf{SU}(3).$

All orientable, spin \mathcal{M}_6 admit a nowhere vanishing η ; torsion undetermined. cf. Bryant:05

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 $\eta \Leftrightarrow$ real two-form J and complex decomposable three-form Ω s.t.

$$\Omega \wedge J = 0,$$
 $\frac{3i}{4}\Omega \wedge \overline{\Omega} = J \wedge J \wedge J = 3! d$ vol

where
$$J_{mn} = -i\eta_+^{\dagger}\gamma_{mn}\eta_+$$
, $\Omega_{mnp} = -i\eta_-^{\dagger}\gamma_{mnp}\eta_+$

- Almost complex structure: $I_m{}^n \sim \epsilon^{nk_1..k_5} \text{Re}\Omega_{mk1k2} \text{Re}\Omega_{k3k4k5}$
- $J, \Omega \Rightarrow \text{metric } g_{mn} = I_m{}^p J_{pn}$ Hitchin:00

 J, Ω closed $\Leftrightarrow \mathcal{M}_6$ is Calabi–Yau.

Otherwise non-zero torsion

Chiossi-Salamon:02

$$dJ = -\frac{3}{2} \mathsf{Im}(W_1 \overline{\Omega}) + W_4 \wedge J + W_3$$
$$d\Omega = W_1 J \wedge J + W_2 \wedge J + \overline{W}_5 \wedge \Omega$$

4D $\mathcal{N} = 1$ solutions from SU(3) structure manifolds

Remark: many Calabi–Yau \rightarrow many fluxless compactifications

Why so few explicit examples with flux?

SU(3) structure not enough: SUSY, BI and EOM selects W_i Complications:

- $W_1, W_2 \neq 0 \Rightarrow I_m{}^p$ not integrable (non-complex)
- $W_1, W_4, W_3 \neq 0$: not symplectic (non-Kähler)

Idea of this talk: Construct explicit SU(3) structures on CY manifolds.

Alternative: Construct non-explicit *SU*(3) structures as deformations of CY Witten–Witten:87, Li–Yau:05, Andreas–Garcia-Fernandez:12

4D Heterotic $\mathcal{N}=1$ Minkowski solutions: Equations

No flux: Calabi-YauCandelas-Horowitz-Strominger-Witten:85 $\mathcal{N} = 1$, Mkw, H = 0 $\iff X$ is Calabi-Yau, dilaton constant. $dJ = d\Omega = 0, H = 0$

With flux: Strominger-Hull systemStrominger:86, Hull:86 $\mathcal{N} = 1$, Mkw, $H \neq 0$ \iff SU(3) structure on X with torsion: $d(e^{-2\phi}J \wedge J) = d(e^{-2\phi}\Omega) = 0, H = i(\partial - \bar{\partial})J$ $W_0 = W_2 = 0, W_5 = 2W_4 = 2d\phi$.

Heterotic vector bundle

 $\mathcal{N}=1$ vector bundle V
ightarrow X with connection A and field strength F must satisfy

$$F \wedge \Omega = 0$$
, $F \wedge J \wedge J = 0$.

Must also satisfy BI

$$dH = rac{lpha'}{4} \left(\operatorname{tr}(F \wedge F) - \operatorname{tr}(R^- \wedge R^-) \right)$$

Motivational example: the quintic

• Hypersurface $X \subset \mathbb{P}^4$,

$$0 = P(x_0,..,x_4) = x_0^5 p(z_1,..,z_4) = 0 \;,\; z_a = rac{x_a}{x_0} \; {
m in} \; U_0: x_0
eq 0$$

Inherit Kahler form:

$$J_0 = \mathcal{J}|_X$$

FS Kahler form $\mathcal{J} = rac{i}{2\pi} \partial \bar{\partial} \ln \kappa, \kappa = 1 + \sum_{a=1}^{4} |z_a|^2$

• Inherit hol. top form:

$$\Omega_0 = rac{dz_1 \wedge dz_2 \wedge dz_3}{p_{,4}}$$

Check SU(3) structure conditions:

$$J_0 \wedge \Omega_0 = 0$$
 but $J_0 \wedge J_0 \wedge J_0 = rac{3i}{4} \mathcal{F} \ \Omega_0 \wedge ar{\Omega}_0$

Motivational example: the quintic

- Inherit Kahler form: $J_0 = \mathcal{J}|_X$ and holomorphic top form: Ω_0
- $J_0 \wedge \Omega_0 = 0$ but $J_0 \wedge J_0 \wedge J_0 = \frac{3i}{4} \mathcal{F} \ \Omega_0 \wedge \overline{\Omega}_0$
- Rescale forms to get SU(3) structure $J = \mathcal{F}^k J_0$, $\Omega = \mathcal{F}^{\frac{3k+1}{2}} \Omega_0$
- Complex, non-Kahler manifold
 - $W_1 = W_2 = W_3 = 0$, $W_4 = k d(\ln \mathcal{F})$, $W_5 = \frac{3k+1}{2} d(\ln \mathcal{F})$

• Strominger–Hull system if k = 1, with flux

$$H = i(\partial - \bar{\partial})J = i(\partial - \bar{\partial})\ln \mathcal{F} \wedge J$$

Method generalizes to any favourable CICY

$$X \sim \left[egin{array}{c|c} \mathbb{P}^{n_1} & q_1^1 & \cdots & q_K^1 \ dots & dots & dots \ \mathbb{P}^{n_m} & q_1^m & \cdots & q_K^m \end{array}
ight]_{\eta}^{h^{1,1},h^{2,1}}$$

E.g. for co-dim 1 CICY

• 1 Kahler form $J_i = \mathcal{J}_i|_X$ from each $\mathbb{P}^{n_i} \subset \mathcal{A}$: $\left| J = \sum_{i=1}^m a_i J_i \right|$

• Holomorphic top form: $\begin{array}{c|c}
\Omega_0 = \hat{\Omega}|_X \\
\hat{\Omega} \wedge dP_1 \wedge \cdots \wedge dP_K = \mu_1 \wedge \cdots \wedge \mu_m, \quad \mu_i = \frac{1}{n_i!} \epsilon_{A_0 A_1 \cdots A_{n_i}} x_{iA_0} dx_{iA_1} \wedge \cdots \wedge dx_{iA_{n_i}}
\end{array}$ Check SU(3) structure:

•
$$J \wedge \Omega_0 = 0 \checkmark$$

• $J_i \wedge J_j \wedge J_k = \frac{3i}{4} \Lambda_{ijk} \Omega_0 \wedge \bar{\Omega}_0$
 $\Lambda_{ijk} = \frac{c_{ijk}}{6\pi^3} \left[\prod_{l=1}^m \frac{|\nabla_l P|^{2n_l}}{\sigma_l} \right] (|\nabla_i P|^2 |\nabla_j P|^2 |\nabla_k P|^2 \sigma_i \sigma_j \sigma_k)^{-1} , \sigma_i = \sum_{A=0}^{n_i} |x_{iA}|^2$

Example: SU(3) structure on tetraquadric

$$X \sim \begin{bmatrix} \mathbb{P}^{1} & 2 \\ \mathbb{P}^{1} & 2 \end{bmatrix}_{-128}^{4,68} \begin{array}{c} \mathbf{x}_{1} = (x_{10}, x_{11}) \\ \mathbf{x}_{2} = (x_{20}, x_{21}) \\ \mathbf{x}_{3} = (x_{30}, x_{31}) \\ \mathbf{x}_{4} = (x_{40}, x_{41}) \\ \mathbf{x}_{4} = \frac{x_{41}}{x_{40}} \end{array}$$

- Hypersurface in $(\mathbb{P}^1)^4$ set by $P(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4) = 0$
- 1 FS Kähler forms: $\mathcal{J}_i = \frac{i}{2\pi} \frac{dz_i \wedge d\overline{z}_i}{\kappa_i^2}$, $\kappa_i = 1 + |z_i|^2$

Restrict to tetra-quadric:

$$J_{\alpha} = \frac{i \, dz_{\alpha} \wedge d\bar{z}_{\alpha}}{2\pi\kappa_{\alpha}^2} , \qquad J_4 = \frac{i}{2\pi\kappa_4^2} \sum_{\alpha,\beta=1}^3 v_{\alpha} \bar{v}_{\beta} \, dz_{\alpha} \wedge d\bar{z}_{\beta} \quad \text{with} \quad v_{\alpha} := \frac{p_{,\alpha}}{p_{,4}} , \text{ on } U_0$$

• Holomorphic top form $\Omega_0 = rac{dz_1 \wedge dz_2 \wedge dz_3}{p_{,4}}$

Check SU(3) structure $J = \sum_{i=1}^{4} a_i J_i$, $\Omega = A \Omega_0$:

• $J \wedge \Omega_0 = 0 \checkmark$ • $\frac{3i}{4}\Omega \wedge \overline{\Omega} = J \wedge J \wedge J \iff |A|^2 = a_1 a_2 a_3 a_4 \sum_{i=1}^4 a_i^{-1} \Lambda_i$ $\frac{1}{6}\Lambda_I := \Lambda_{ijk} = \frac{1}{6\pi^3} \frac{|\rho_I|^2 \kappa_I^2}{\kappa_K^2 \kappa_K^2 \kappa_K^2}$ SU(3) structures on CY: Torsion classes

CICY SU(3) structure

$$J = \sum_{i=1}^{m} a_i J_i$$
 , $\Omega = A \Omega_0$

subject to
$$|A|^2 = \sum_{i,j,k=1}^m \Lambda_{ijk} a_i a_j a_k$$

Torsion classes easily computed:

$$dJ = \sum_{i=1}^m da_i \wedge J_i \;, \qquad d\Omega = d\ln(A) \wedge \Omega$$

 $\implies W_1 = W_2 = 0 ,$ $W_3 = \sum_i (da_i - W_4) \wedge J_i , \quad W_4 = \frac{1}{2} \sum_i J \lrcorner (da_i \wedge J_i) , \quad W_5 = d \ln(A).$

Integrable complex structure with exact W_5 ; rest set by a_i

Metric and torsion explicit, and slightly tuneable by choosing a_i .

SU(3) structures on CY: Torsion classes

Universal CICY SU(3) structure Choose $a_i = a t_i$ for $i = 1, \ldots, m$ $J = a J_0$, $J_0 := \sum_{i=1}^m t_i J_i$, $\Omega = A \Omega_0$ With $g_{0,\alpha\bar{\beta}} = -2iJ_{0,\alpha\bar{\beta}}$ get $|A|^2 = a^3 \mathcal{F} \ , \quad \text{where} \quad \mathcal{F} := \ \sum \ \Lambda_{ijk} t_i t_j t_k = |\text{det}B|^2 \text{det} \left(g_{0,\alpha\bar\beta}\right) > 0 \ .$ i.i.k=1

Torsion classes

$$W_1 = W_2 = W_3 = 0$$
, $W_4 = d \ln a$, $W_5 = d \ln A = \frac{3}{2} d \ln a + \frac{1}{2} d \ln \mathcal{F}$.

In summary: Any CICY allows a Universal SU(3) structure (J, Ω) with torsion

$$W_1 = W_2 = W_3 = 0$$
, $W_4 = d \ln a$, $W_5 = \frac{3}{2} d \ln a + \frac{1}{2} d \ln \mathcal{F}$,

where $\mathcal{F} := \sum_{i,j,k=1}^m \Lambda_{ijk} t_i t_j t_k = |\det B|^2 \det \left(g_{0,\alpha\bar{\beta}}\right) > 0$, and metric

 $g_{\alpha\bar{\beta}} = ag_{0,\alpha\bar{\beta}},$

Choose $a = \mathcal{F}$: reproduce torsion for Strominger–Hull system with

$$H = i(\partial - \bar{\partial})\mathcal{F} \wedge J_0$$

Any CICY allows Strominger–Hull type SU(3) structure (J, Ω) with torsion

$$W_1 = W_2 = W_3 = 0$$
, $W_4 = d \ln \mathcal{F}$, $W_5 = 2d \ln \mathcal{F}$,

Right torsion is not enough: must construct suitable vector bundle and solve BI.

SUSY and bundle stability — Work in progress

SUSY \iff holomorphic vector bundle $V \rightarrow X$ with HYM connection:

$$F \wedge \Omega = 0$$
, $F \wedge J \wedge J = 0$.

Li–Yau theorem: $V \to X$ allows HYM connection $\iff V \to X$ is stable.

Donaldson'85, Uhlenbeck-Yau'86, Li-Yau'87, Kobayashi'87, Hitchin

Here: know $J \rightsquigarrow$ solve for F.

Bianchi identity — Work in progress

$$2i\,\overline{\partial}\partial\mathcal{F}\wedge J_0=dH=rac{lpha'}{4}\left(\mathrm{tr}(F\wedge F)-\mathrm{tr}(R\wedge R)
ight)+...\,,$$

$\operatorname{tr}(R \wedge R)$

Invariant under conformal re-scaling: $tr(R \land R) = tr(R_0 \land R_0)$ Computable but lack manageable form for general CICY.

$\operatorname{tr}(F \wedge F)$

Construct stable vector bundle $V \rightarrow X$ with suitable connection A.

Example: $tr(R \land R)$ on tetraquadric

$$X \sim \begin{bmatrix} \mathbb{P}^1 & 2 \\ \mathbb{P}^1 & 2 \end{bmatrix}_{-128}^{4,68} \begin{array}{c} \mathbf{x}_1 = (x_{10}, x_{11}) & z_1 = \frac{x_{11}}{x_{10}} \\ \mathbf{x}_2 = (x_{20}, x_{21}) & z_2 = \frac{x_{21}}{x_{20}} \\ \mathbf{x}_3 = (x_{30}, x_{31}) & z_3 = \frac{x_{31}}{x_{30}} \\ \mathbf{x}_4 = (x_{40}, x_{41}) & z_4 = \frac{x_{41}}{x_{40}} \\ \end{bmatrix}$$

Set
$$z_4 = f(z_1, z_2, z_3)$$
, $v_{\alpha} = p_{,\alpha}/p_{,4}$, $\Omega_{\alpha}^{\beta} = \frac{t_1 t_2 t_3 t_4}{t_{\beta}} \bar{\partial} \left(\frac{\Lambda_4 \Lambda_{\beta}}{\mathcal{F} \Lambda_{\alpha}} \partial \left(\frac{\Lambda_{\alpha}}{\Lambda_4} \right) \right)$.
Compute curvature 2-form:

$$\mathcal{R}^{eta}_{lpha} = -4\pi i J_{lpha} \delta^{eta}_{lpha} - rac{m{v}_{lpha}}{m{v}_{eta}} \Omega^{eta}_{lpha}$$

 $\operatorname{tr} R = \partial \bar{\partial} \ln \det g_0 \checkmark$

Finally

$$\operatorname{tr}(R \wedge R) = \sum_{\alpha,\beta=1}^{3} (R_{\alpha}^{\beta} \wedge R_{\beta}^{\alpha}) + \operatorname{c.c.} = 8\pi i \sum_{\alpha=1}^{3} J_{\alpha} \wedge \Omega_{\alpha}^{\alpha} + \sum_{\alpha,\beta=1}^{3} \Omega_{\alpha}^{\beta} \wedge \Omega_{\beta}^{\alpha} + \operatorname{c.c.}.$$

Conclusions

All CY manifolds allow several SU(3) structures

CICY: ambient space provide building blocks for non-trivial SU(3) structures

- $J = \sum_{i=1}^{m} a_i J_i$, $\Omega = A \Omega_0$, subject to $|A|^2 = \sum_{i,j,k=1}^{m} \Lambda_{ijk} a_i a_j a_k$
- metric computable $g_{mn} = I_m^k J_{kn}$
- torsion computable:

$$W_3=\sum_i (da_i-a_iW_4)\wedge J_i\;,\quad W_4=rac{1}{2}\sum_i J\lrcorner(da_i\wedge J_i)\;,\quad W_5=d\ln(A)\;.$$

• With $a_i = a \quad \forall i \implies$ Strominger–Hull system, with $dH \neq 0$

Work on heterotic BI: compute $tr(R \land R)$, construct vector bundles, ...

Conclusions

Use ambient space forms to build complex SU(3) structures on CICY Necessary constraints for e.g. Strominger–Hull system fulfilled

- Construct vector bundles with HYM connections satisfying heterotic BI
- Explore "non-universal" SU(3) structures
- Type IIB vacua: "smeared" sources required?
- Extend construction: other types of CY manifolds, other dimensions, ...
- Generalise method: non-complex SU(3) structures, ...

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Conclusions

Use ambient space forms to build complex SU(3) structures on CICY Necessary constraints for e.g. Strominger–Hull system fulfilled

What we should do next:

- Construct vector bundles with HYM connections satisfying heterotic BI
- Explore "non-universal" SU(3) structures
- Type IIB vacua: "smeared" sources required?
- Extend construction: other types of CY manifolds, other dimensions, ...
- Generalise method: non-complex SU(3) structures, ...

Thank You

4D $\mathcal{N} = 1$ solutions from SU(3) structure manifolds

4D geometry	String vacuum	Non-vanishing torsion	SU(3) type
$\mathcal{N}=1$ Mkw	Heterotic, Type II (H_3)	$W_3, W_4 = d\phi, W_5 = 2W_4$	Complex
	Heterotic $(H_3 = 0)$	$W_i = 0$, $\forall i$	CY
$\mathcal{N}=1$ Mkw	Type IIB	$3W_4 = 2W_5$	Conf. CY
	$(H_3, F_3, F_5, O_3/O_7)$		
	Type IIB	$W_3, W_4 = d\phi, W_5 = 2W_4$	Complex
	$(F_3, O_5/O_9)$		
	Type IIB/F-theory	$W_4 = W_5 = d\phi$	Complex
	$(H_3, F_3, F_5, O_3/O_7)$		
$\mathcal{N}=1$ Mkw	Type IIA	$W_2, 3W_5 = d\phi$	Symplectic
	(F_2, F_4, O_6)		
$\mathcal{N}=1~AdS$	Type IIA	$W_1^+, W_2^+, dW_2^+ \propto \Omega^+$	Half-flat
	(H_3, F_{even})		