

How to relax the cosmological neutrino mass bound

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Neutrino oscillations → neutrinos massive

PMNS matrix:

Mass splittings:

Mass hierarchy:

Absolute mass scale:

Neutrino oscillations → neutrinos massive

PMNS matrix: ✓

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Troitsk (1994-2004):	$m_{\nu_e} < 2.05 \text{ eV}$	} Laboratory: Tritium β decay
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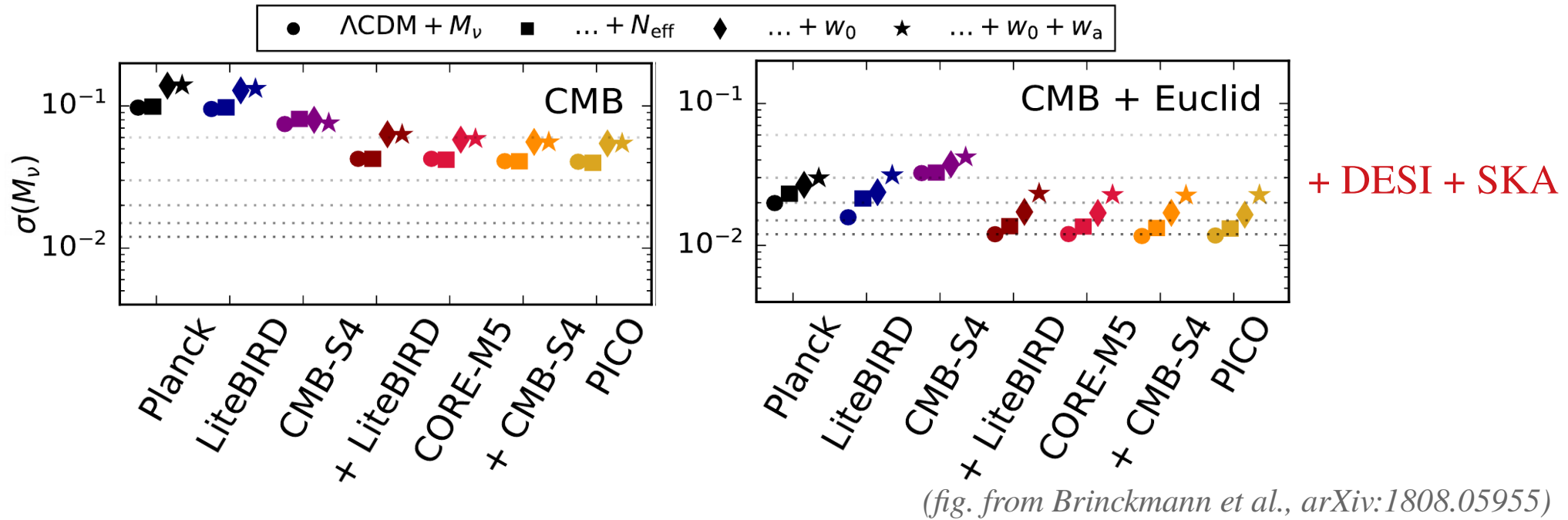
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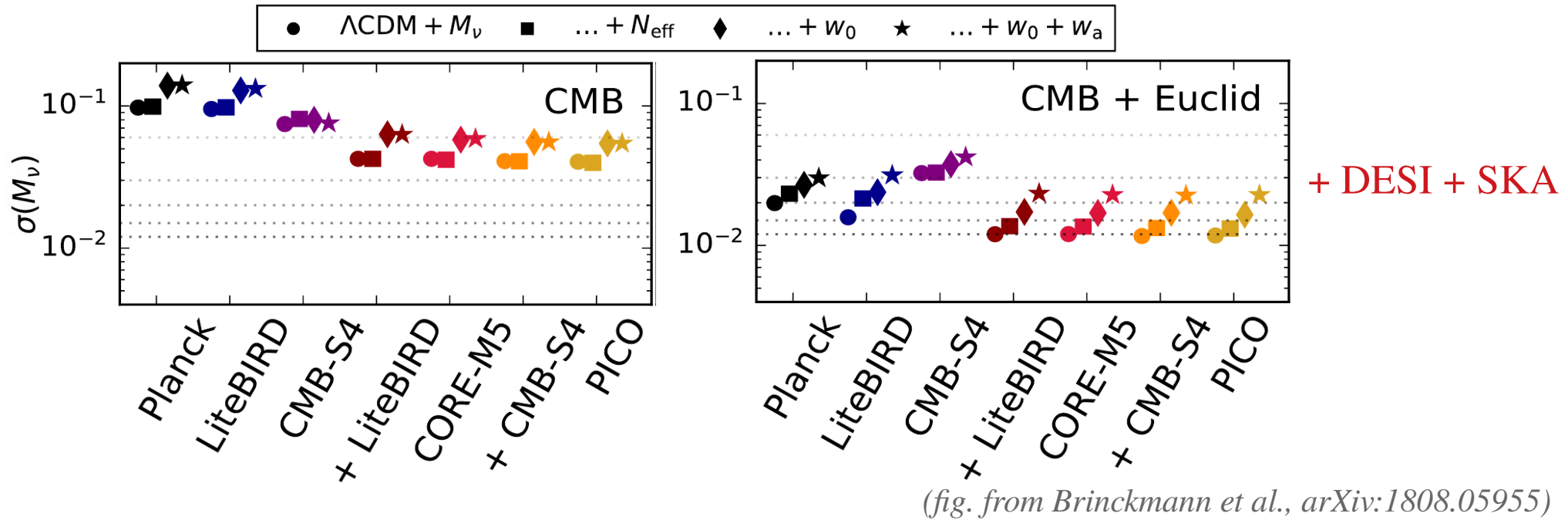
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	$\sum m_\nu < 0.54 \text{ eV}$	(TT+lowP)
<u>Cosmology (Planck 2018):</u>	$\sum m_\nu < 0.16 \text{ eV}$	(TT+lowP+BAO)
	$\sum m_\nu < 0.12 \text{ eV}$	(TT,TE,EE+lowP+lensing+BAO)

... and the future looks even brighter:

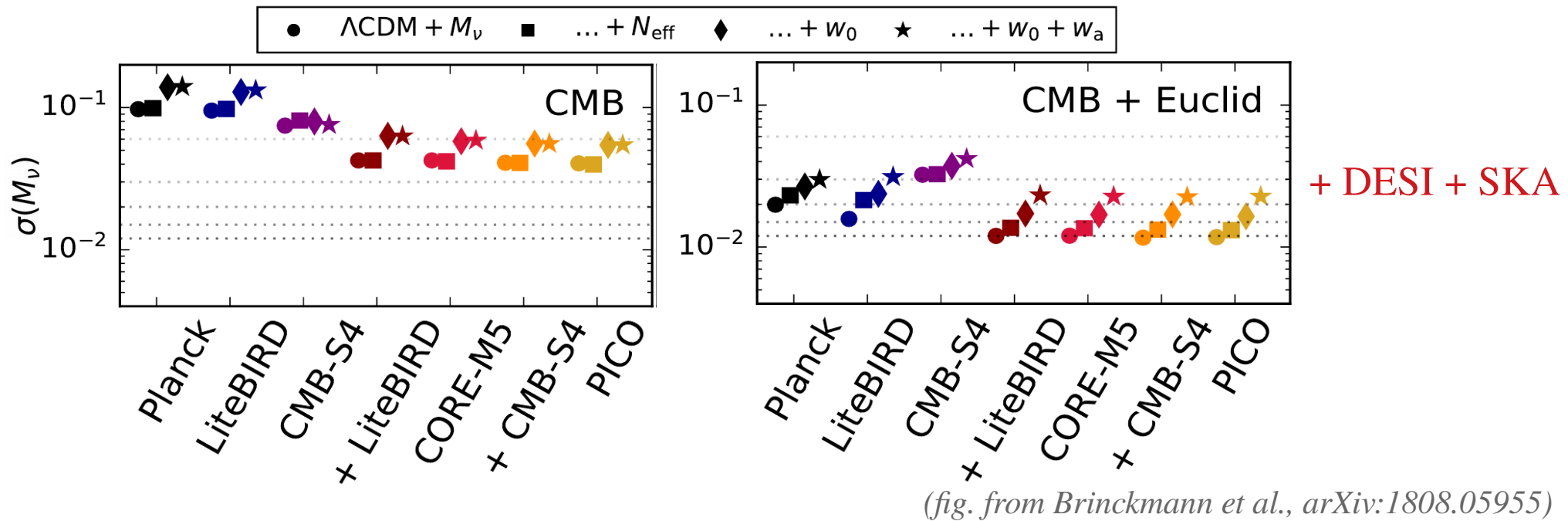


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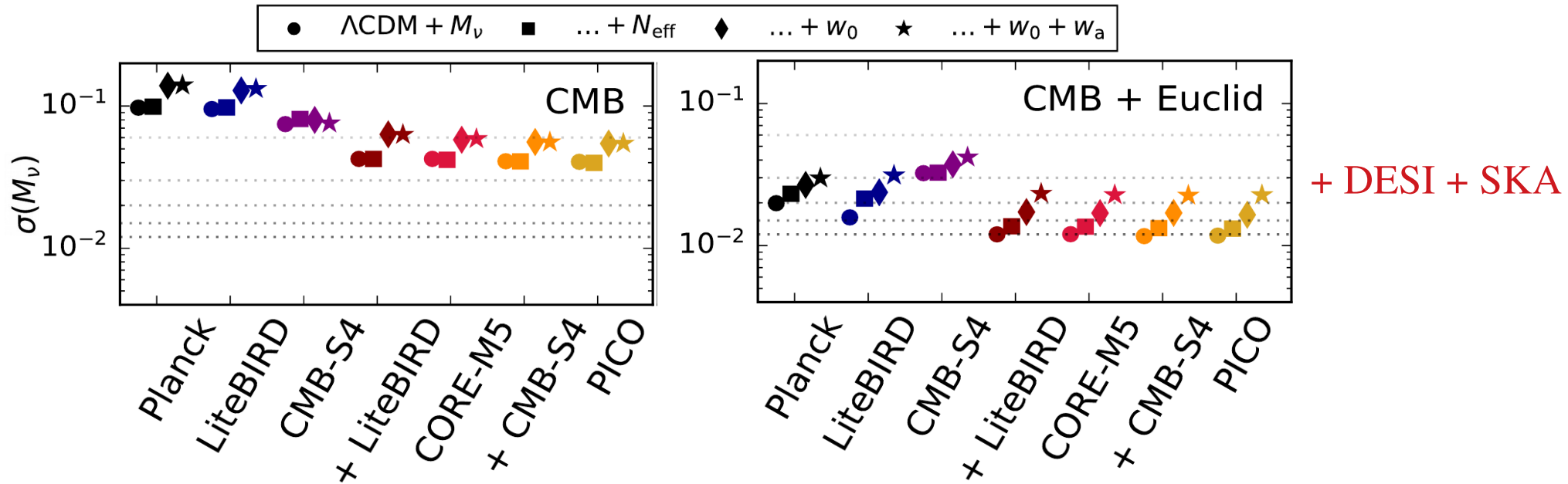


But how robust are those constraints?

Tightest constraints assume ΛCDM base model

→ parameter degeneracies (w , A_{lens} , τ_{reio} , N_{eff} ...)

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(fig. from Brinckmann et al., arXiv:1808.05955)

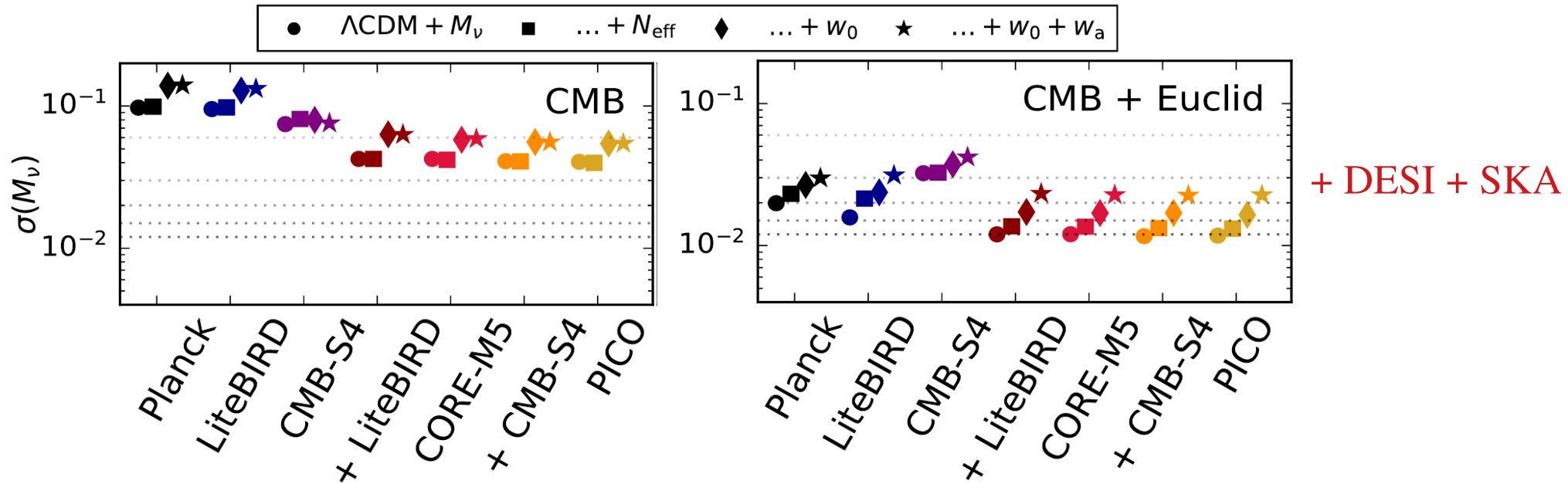
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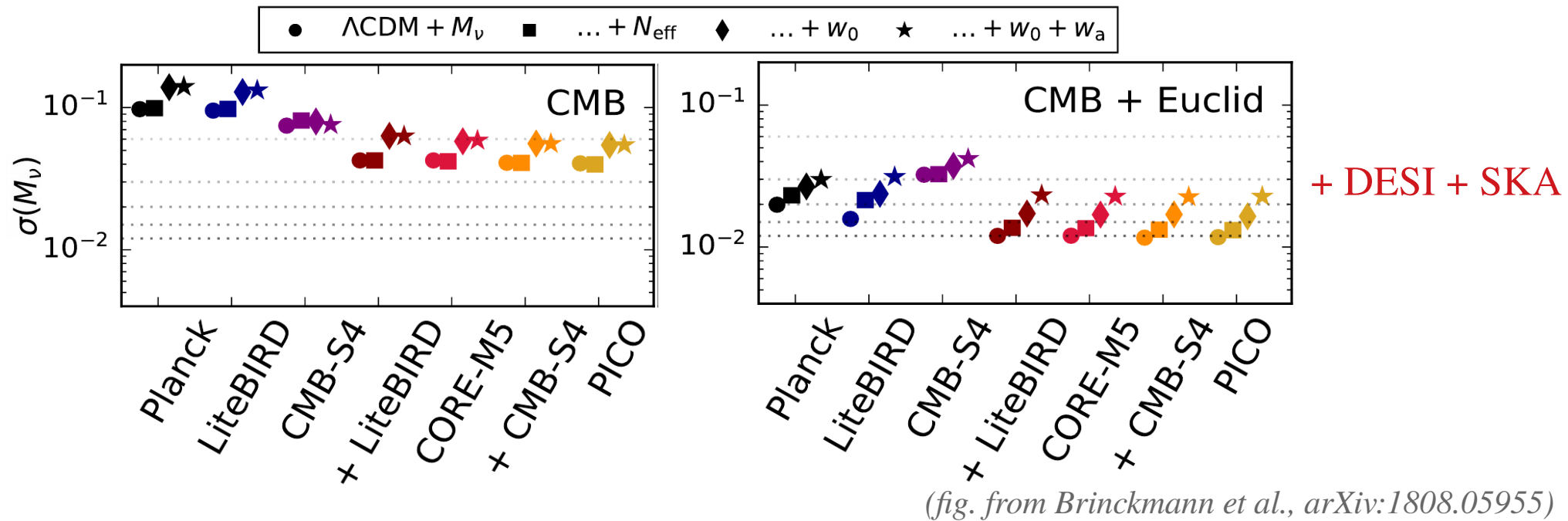
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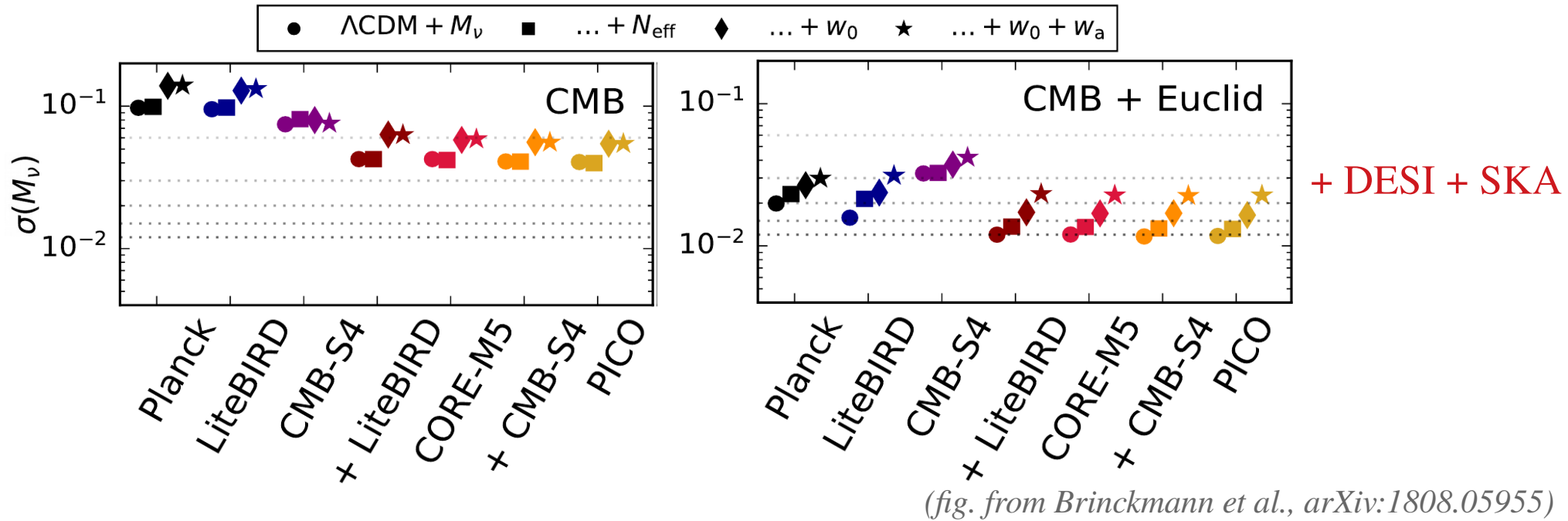
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} **This work:**
Less extreme
scenario₂

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Simple take-home message:

Neutrinos can be much heavier if they have a larger average-momentum

Cosmic perturbation theory:

→ neutrino decouple at ~ 1 MeV (when still ultra-relativistic)

$$\text{Standard assumption: } f_\nu(x) = \frac{1}{e^x + 1}$$

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Model-independent parameterization for non-thermal distributions:

→ expansion in orthonormal polynomials $\int dx \frac{1}{e^x + 1} p_n(x) p_m(x) = \delta_{nm}$

$$\Rightarrow f_\nu(x) = \frac{1}{e^x + 1} \sum_{n=0}^{\infty} C_n p_n(x) \quad (\text{similar to } \text{Esposito, Miele, Pastor, Peloso, Pisanti 2000})$$

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Here: only $n < 3$

$$f_\nu(x) = N \cdot \frac{1}{e^x + 1} \left(p_0(x) + F_1 p_1(x) + F_2 p_2(x) \right)$$

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Eliminate degeneracy by fixing energy density at early times:

$$N \cdot \int dx \frac{x^3}{e^x + 1} \left(p_0(x) + F_1 p_1(x) + F_2 p_2(x) \right) \stackrel{!}{=} \int dx \frac{x^3}{e^x + 1}$$

$$\Rightarrow \text{fix the normalization: } N = \frac{5.682197}{6.825012 + 18.239415F_1 + 17.420046F_2}$$

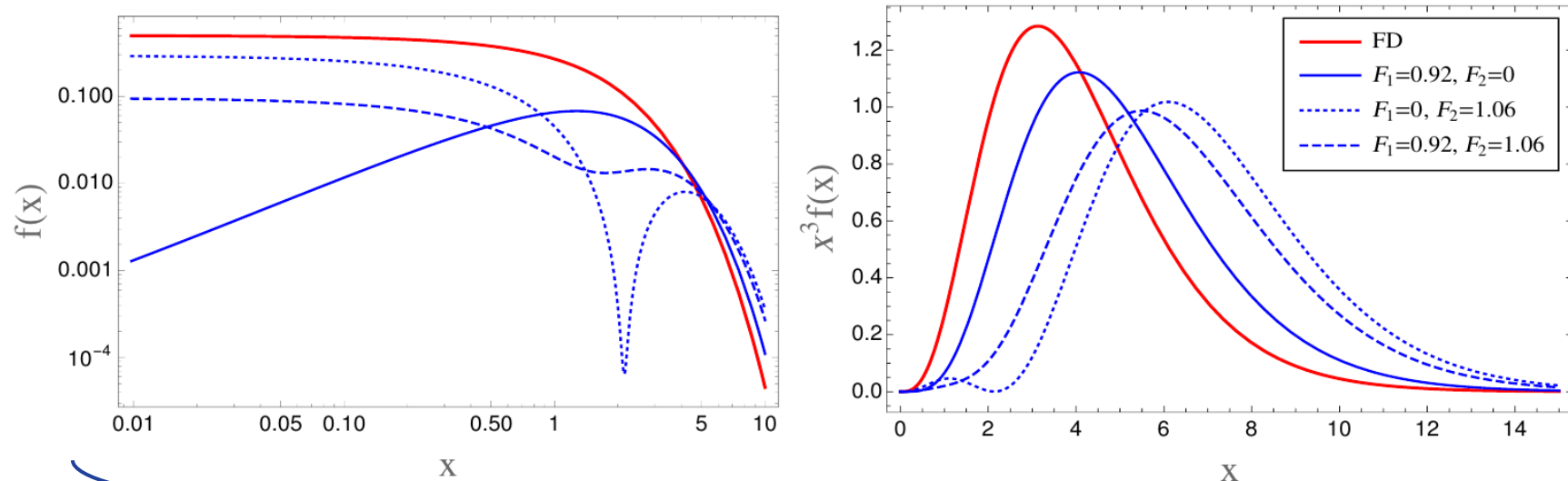
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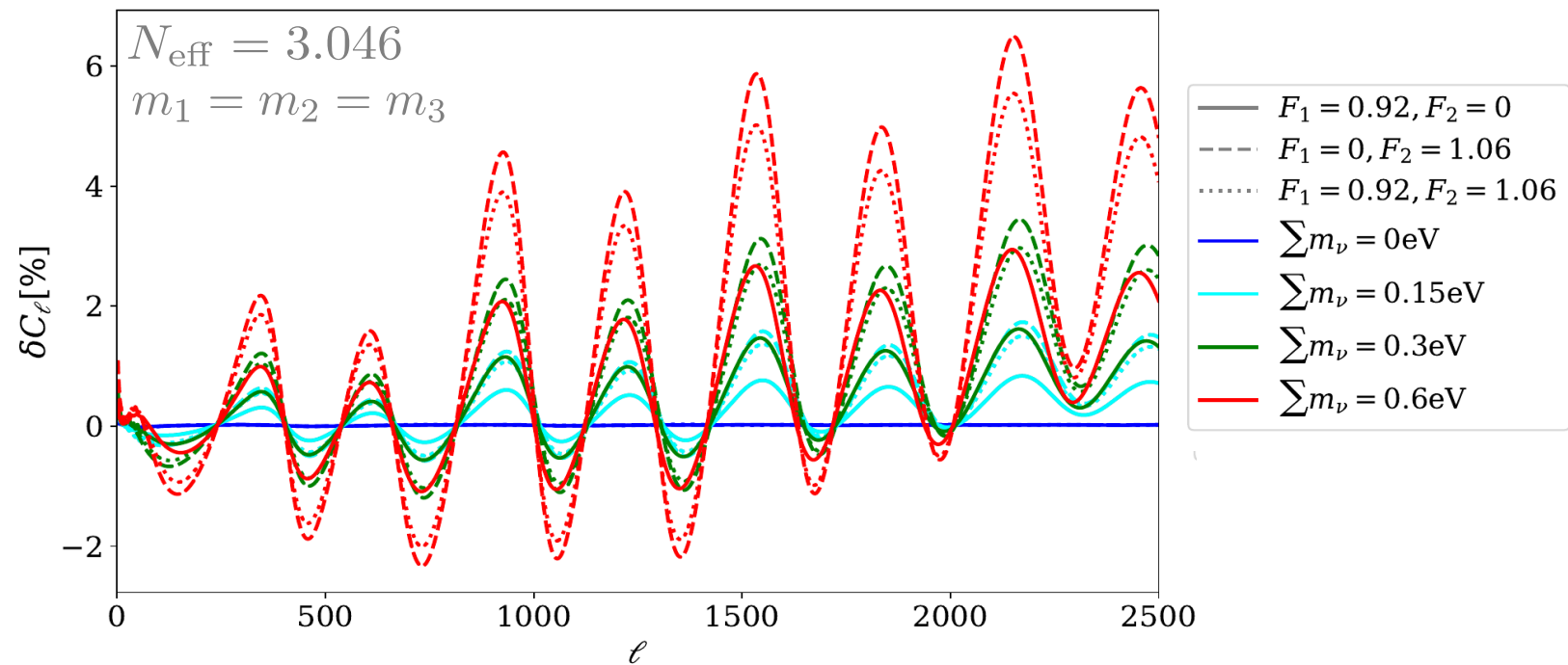
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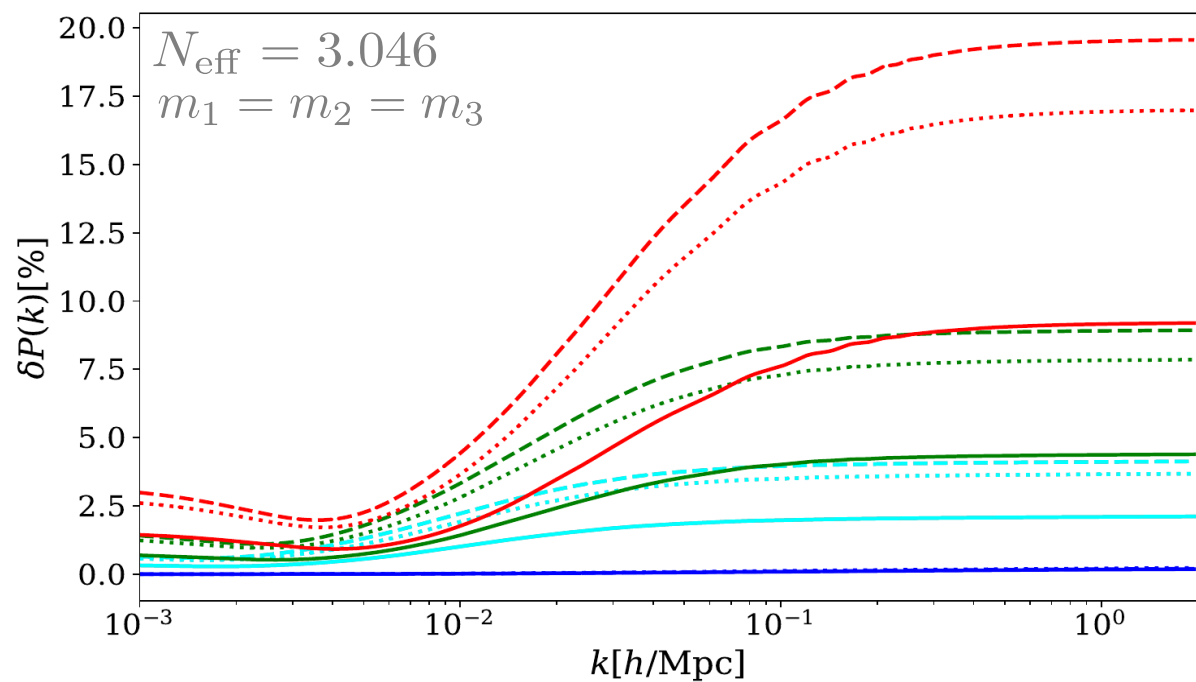
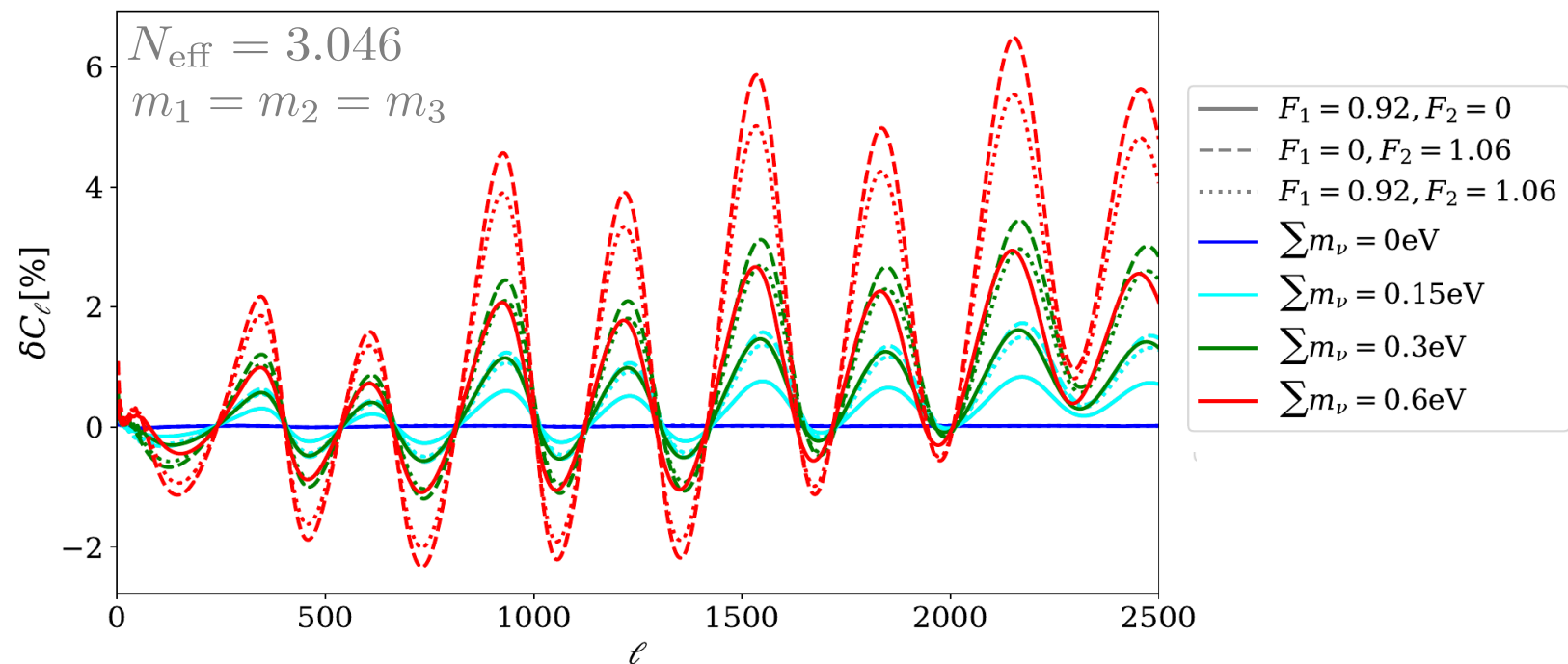
$$N \cdot \int dx \frac{x^3}{e^x + 1} \left(p_0(x) + F_1 p_1(x) + F_2 p_2(x) \right) \stackrel{!}{=} \int dx \frac{x^3}{e^x + 1}$$

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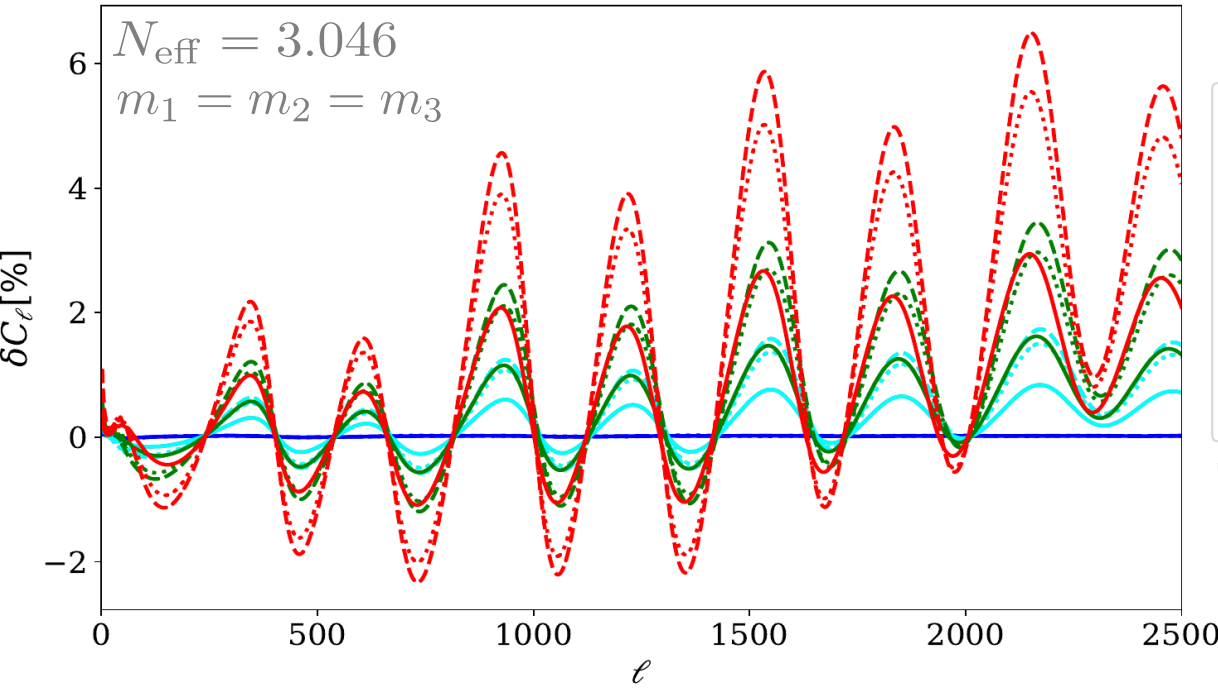


Model independent parameterization allows some variation,
but all distributions still relatively close to Fermi-Dirac...

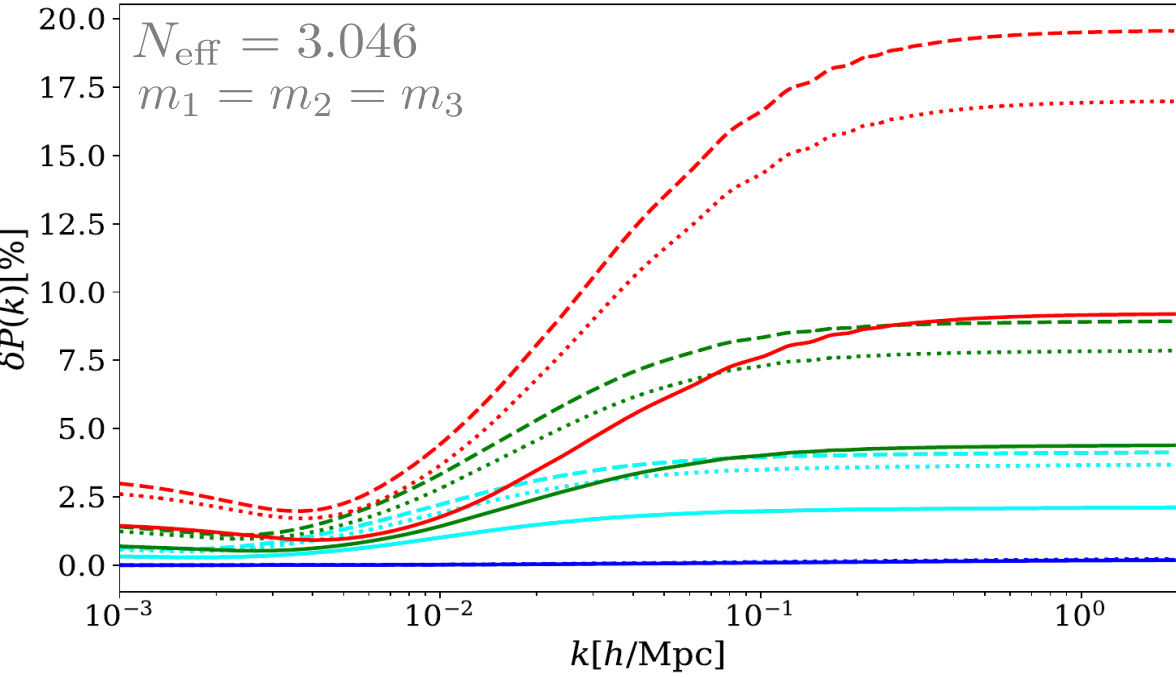




Using CLASS
(Blas, Lesgourgues, Tram 2011)



- $F_1 = 0.92, F_2 = 0$
- - - $F_1 = 0, F_2 = 1.06$
- ⋯ $F_1 = 0.92, F_2 = 1.06$
- $\sum m_\nu = 0 \text{ eV}$
- $\sum m_\nu = 0.15 \text{ eV}$
- $\sum m_\nu = 0.3 \text{ eV}$
- $\sum m_\nu = 0.6 \text{ eV}$



Is this a unique feature of the distribution function???

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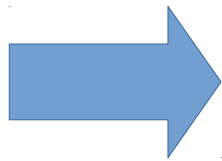
Energy density at *late times* can be fixed by adopting the neutrino mass:

$$\int dx x^2 \sqrt{x^2 + \frac{m_\nu^2}{T_{\nu 0}^2}} \left(\frac{1}{e^x + 1} \right) \stackrel{!}{=} N \int dx x^2 \sqrt{x^2 + \frac{m_\nu^{*2}}{T_{\nu 0}^2}} \frac{1}{e^x + 1} \left(p_0(x) + F_1 p_1(x) + F_2 p_2(x) \right)$$

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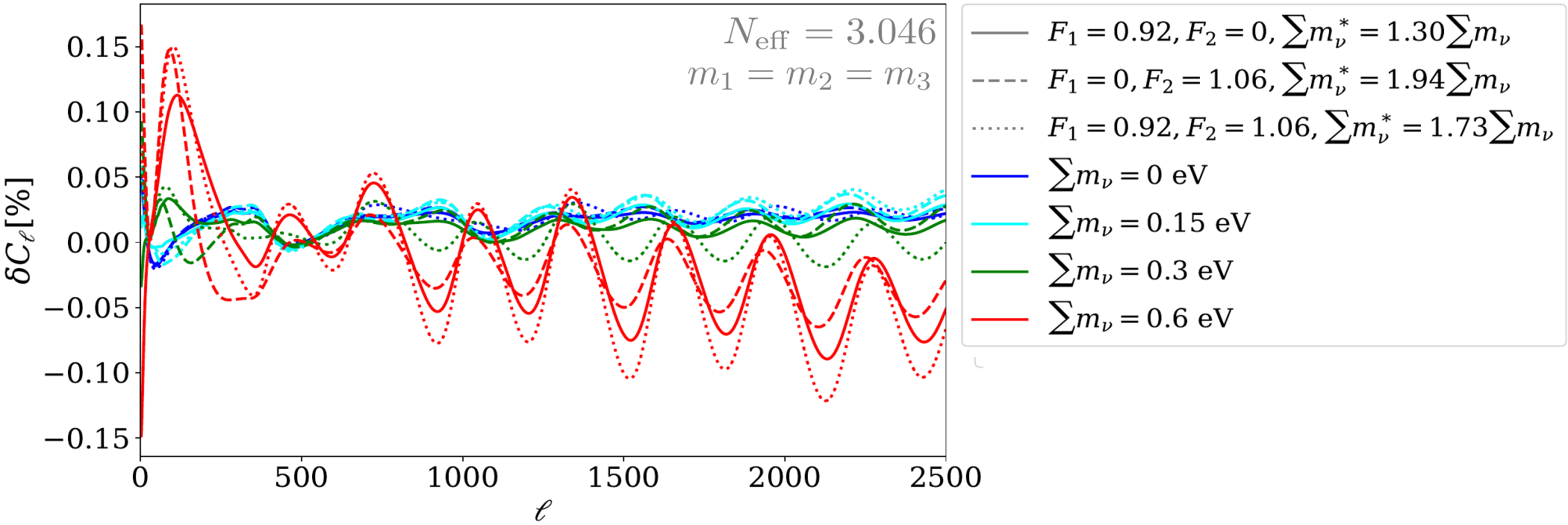
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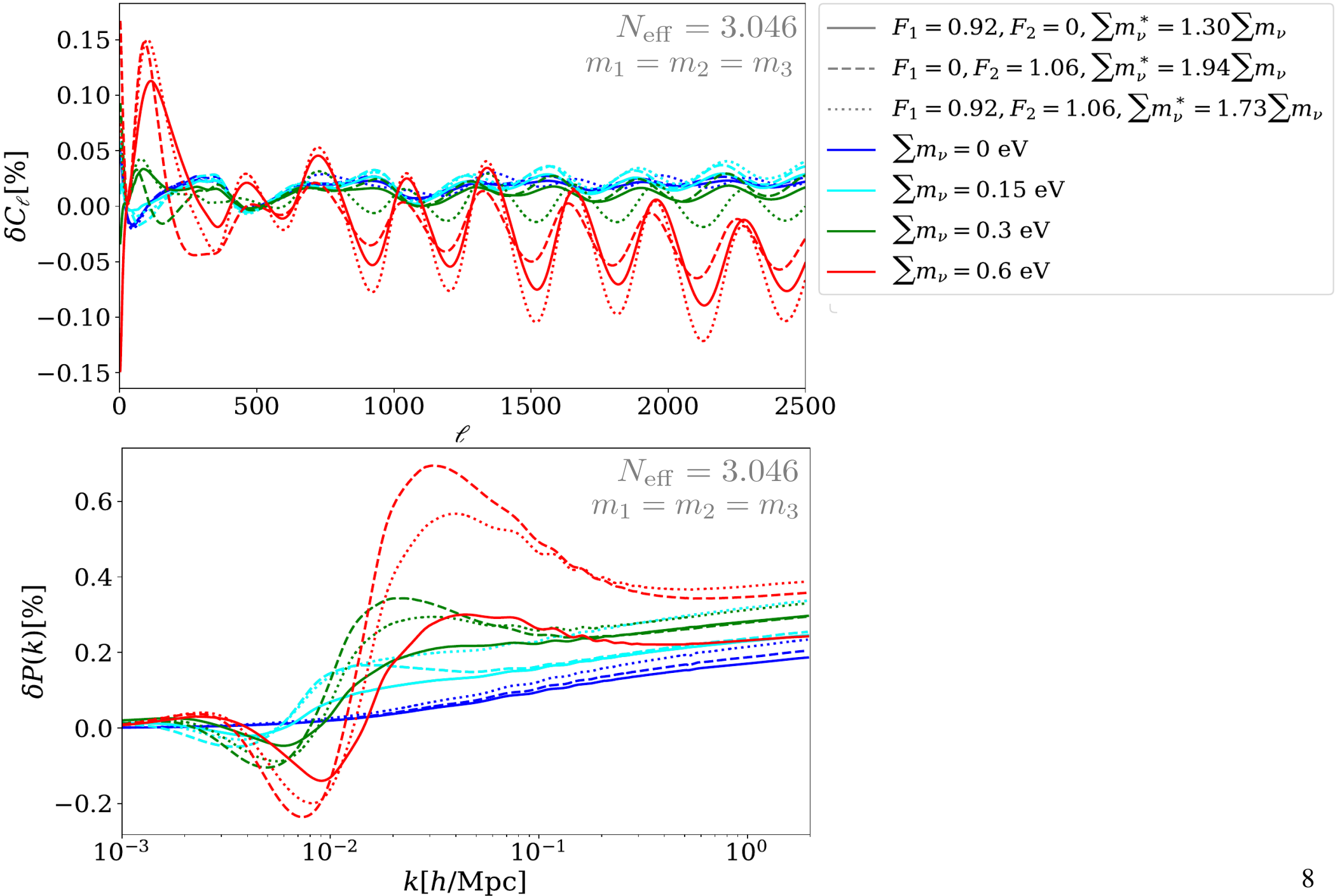


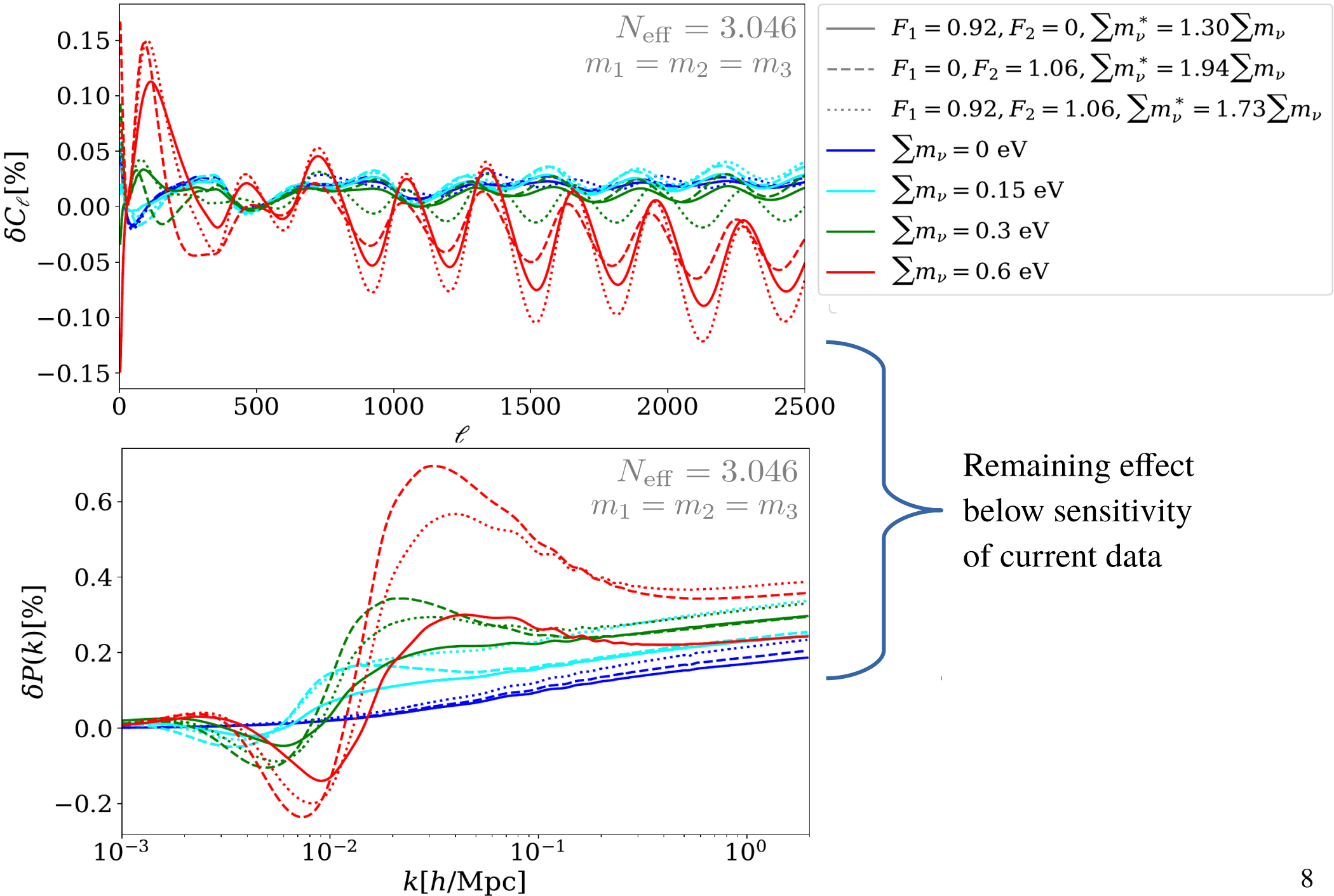
$$F_1 = 0.92, \quad F_2 = 0 \quad : \quad m_\nu^* = 1.30 \cdot m_\nu$$

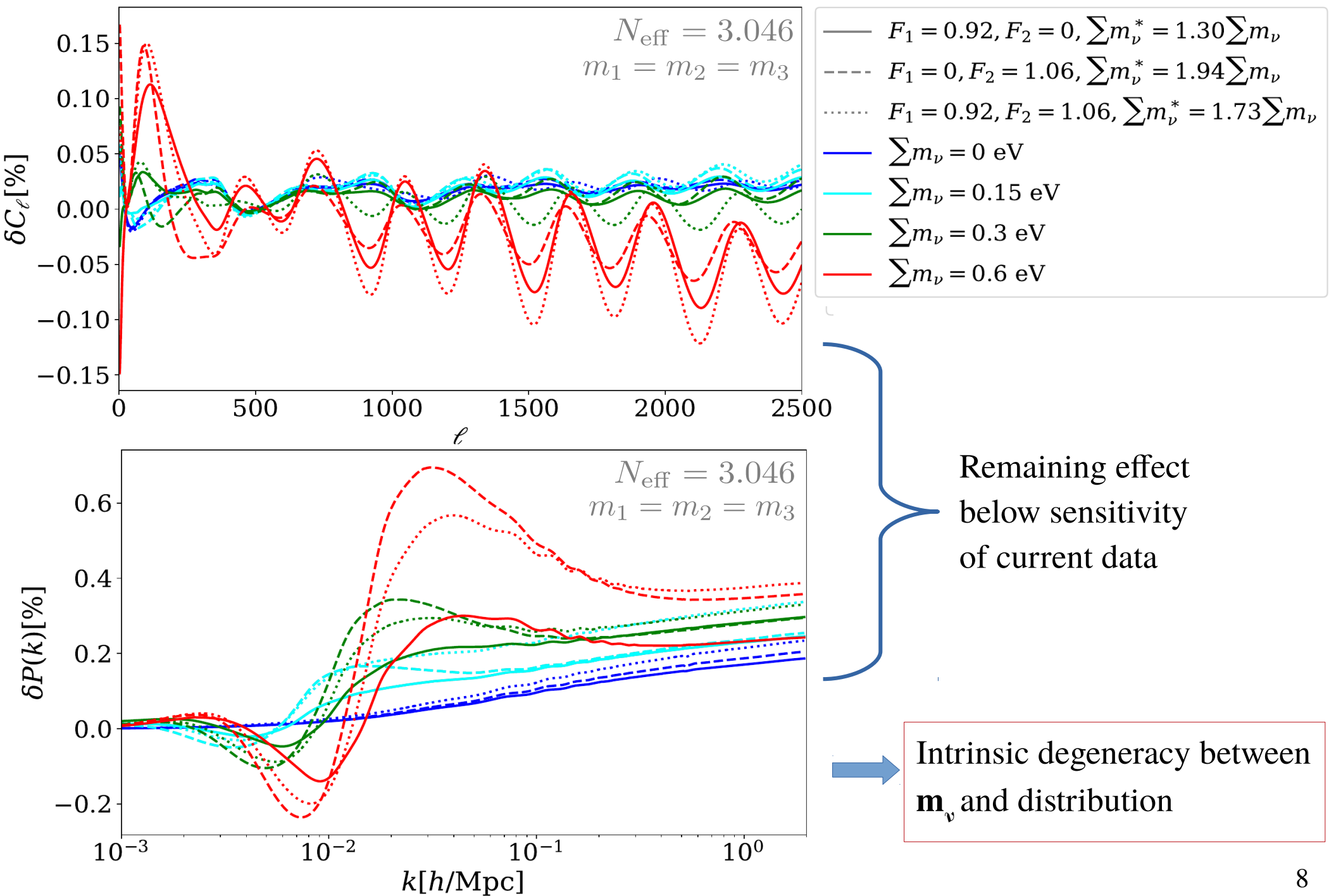
$$F_1 = 0, \quad F_2 = 1.06 \quad : \quad m_\nu^* = 1.94 \cdot m_\nu$$

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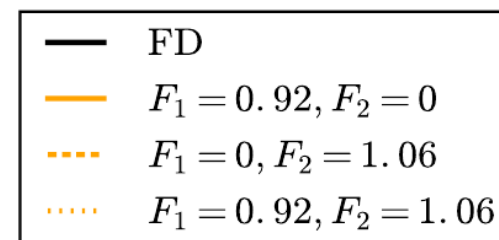
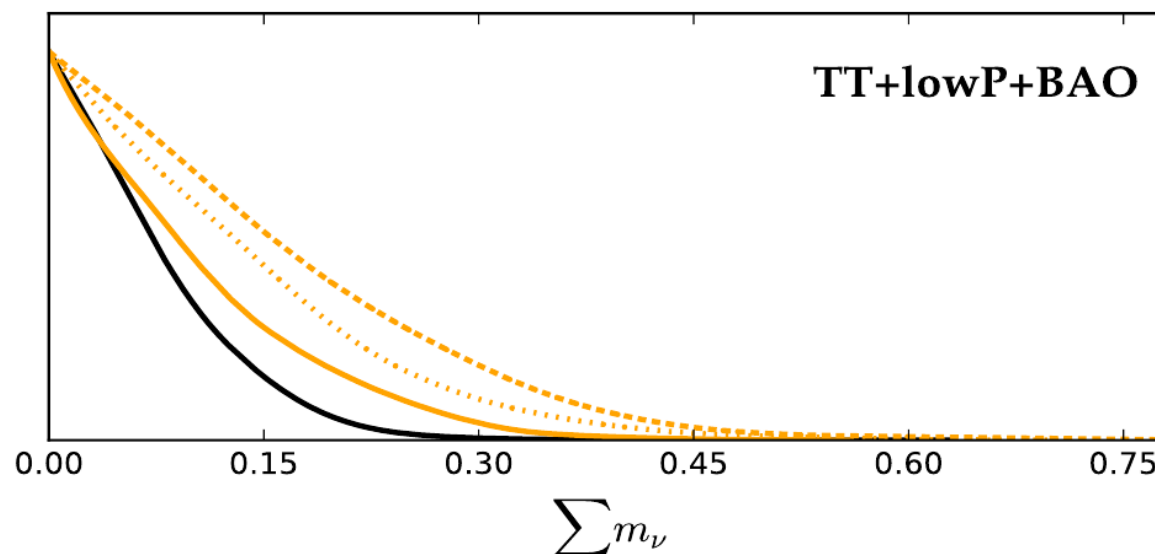
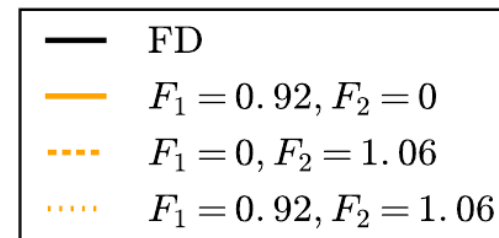
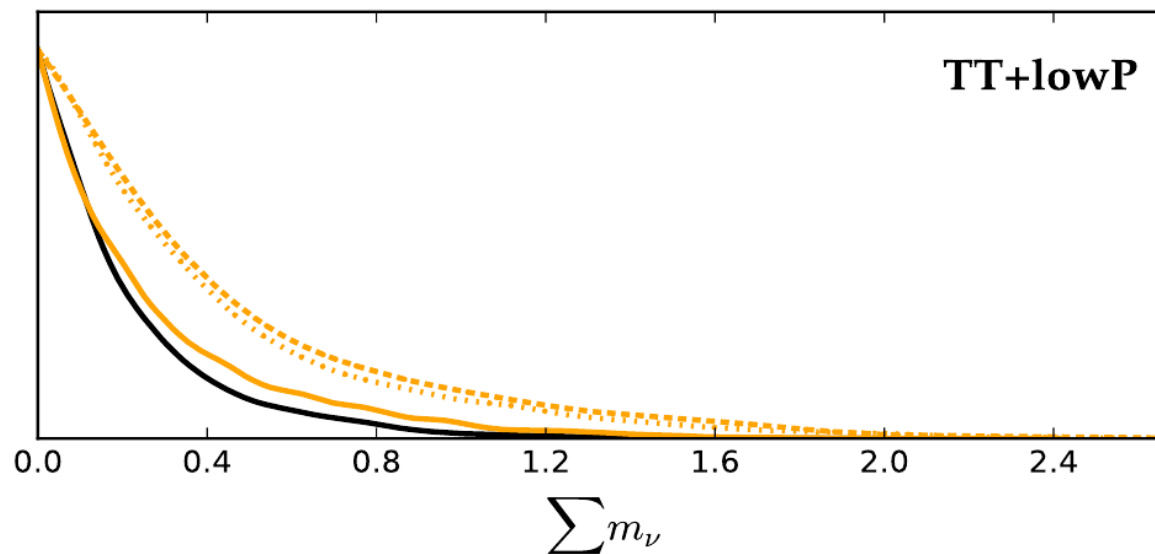






Monte-Carlo-Markov-Chain analysis

*Using MontePython
(Audren, Brinckmann,
Lesgourgues, Benabed, Prunet)*



➔ Around 100% relaxed neutrino mass bounds

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 - **Cosmological neutrino mass bound strongly depends on our assumptions about the relic distribution function...**
- Here: Relaxed neutrino mass bounds by around **100%**.
 - **More drastic scenarios may relax the mass bound in a more drastic way**

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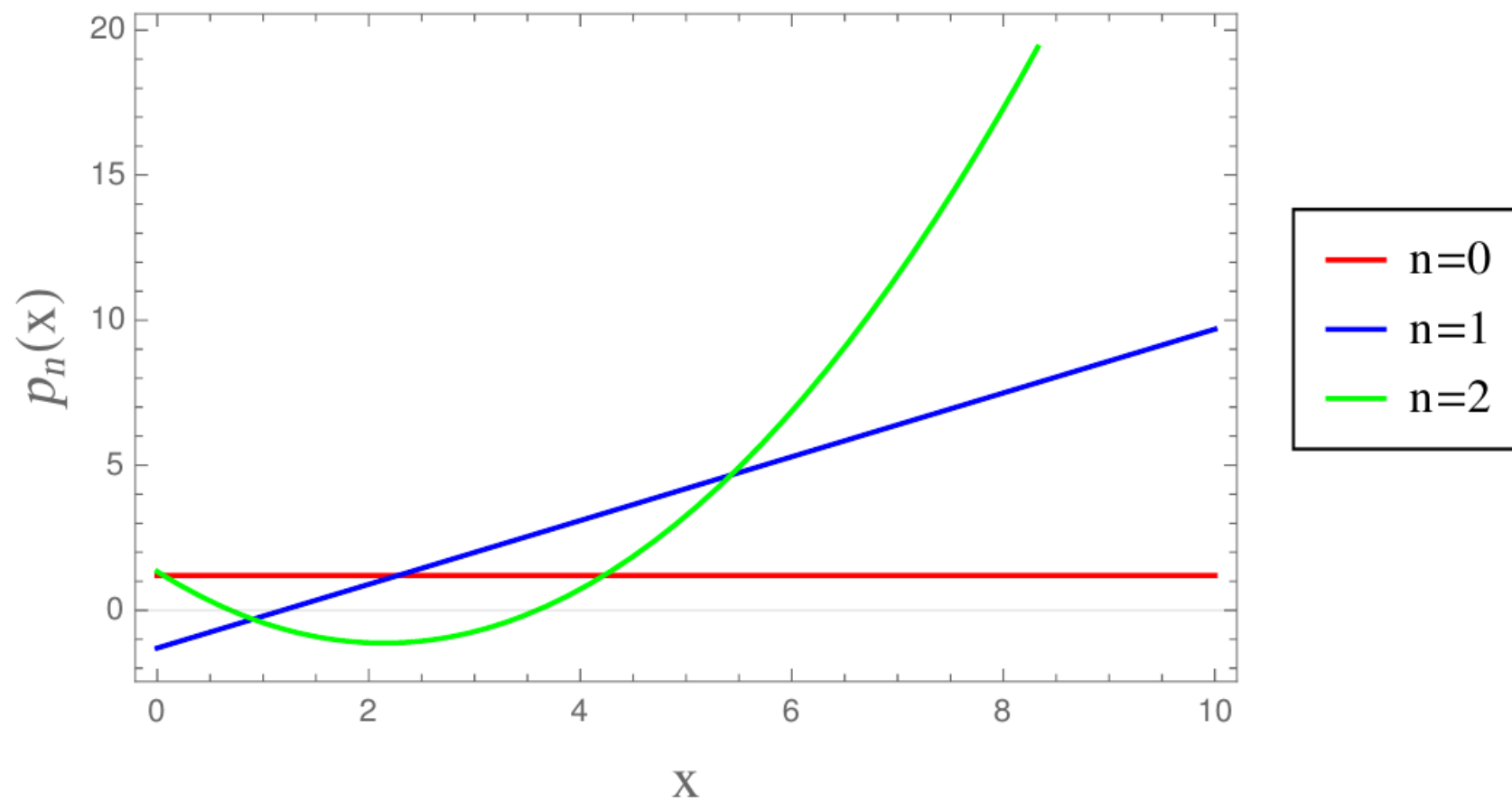
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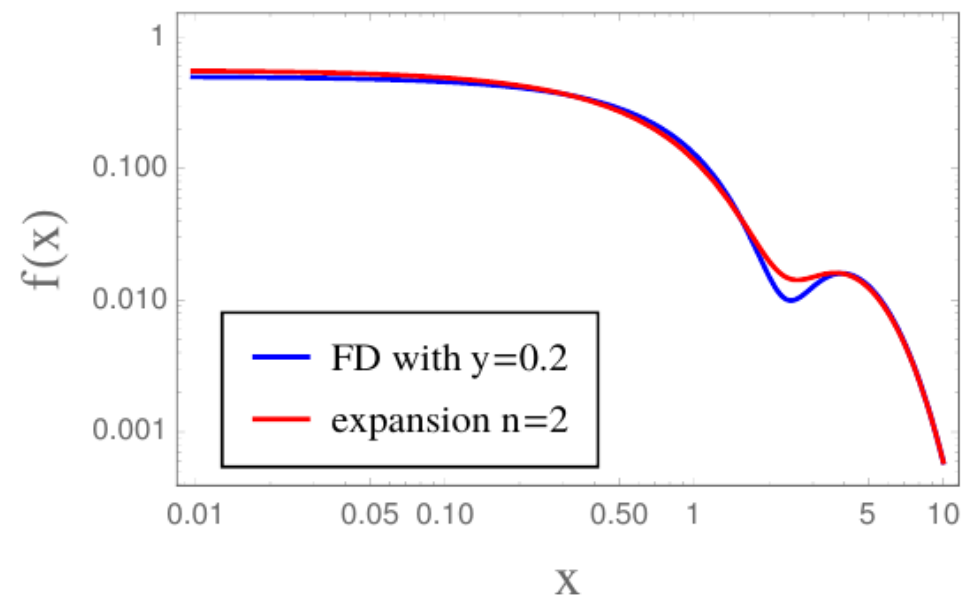
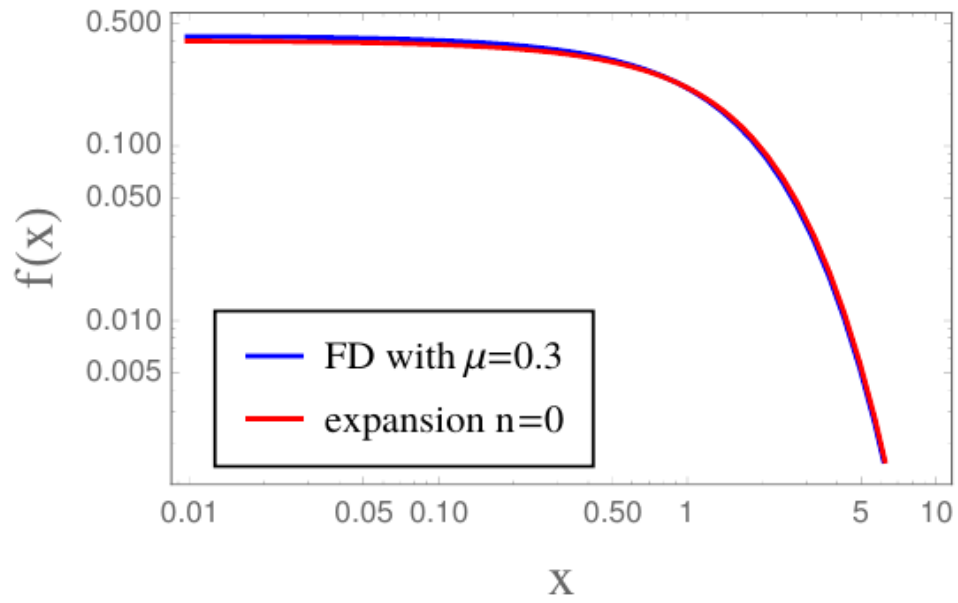
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Thank you
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	TT+lowP (95 % CL)	TT+lowP+BAO (95 % CL)
FD	$\sum m_\nu < 0.73 \text{ eV}$	$\sum m_\nu < 0.18 \text{ eV}$
$F_1 = 0.92, F_2 = 0$	$\sum m_\nu < 0.95 \text{ eV}$	$\sum m_\nu < 0.26 \text{ eV}$
$F_1 = 0, F_2 = 1.06$	$\sum m_\nu < 1.45 \text{ eV}$	$\sum m_\nu < 0.37 \text{ eV}$
$F_1 = 0.92, F_2 = 1.06$	$\sum m_\nu < 1.34 \text{ eV}$	$\sum m_\nu < 0.32 \text{ eV}$