

Dark energy without fine tuning

arXiv:1807.04359

"Dark energy as a remnant of inflation and electroweak symmetry breaking"

arXiv:1905.00045

"Dark energy without fine tuning"

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Aachen



Outline

- 1 Introduction
- 2 The conspiracy of scales
- 3 Assisted relaxation
- 4 Conclusions

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The Einstein equation

$$G_{\mu\nu}M_{\text{P}}^2 = T_{\mu\nu} - g_{\mu\nu}\rho_\Lambda$$

$$\rho_\Lambda \equiv \Lambda_0 M_{\text{P}}^2 + \text{'slow' fields} + \text{vacuum energy}$$

- Observations consistent with:

$$\rho_\Lambda^{\text{obs}} \approx 2.6 \cdot 10^{-47} \text{GeV}^4$$

Problems with ρ_Λ

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0-point energy with a cut-off

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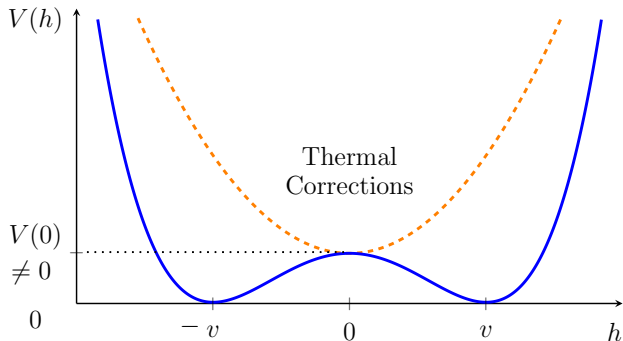
$$\int^{M_{\text{P}}} d^3k \omega \sim M_{\text{P}}^4 \ll \rho_\Lambda^{\text{obs}}$$

**The quintessentially
incorrect calculation!**

The Higgs vacuum energy

The Higgs potential:

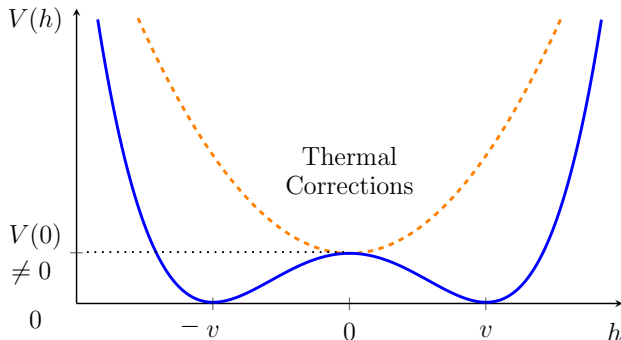
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The Higgs vacuum energy

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$$V(h_{\min}) = \begin{cases} V(0) = \frac{\lambda}{4} v^4, & T > T_{\text{EW}} \\ V(v) = 0, & T < T_{\text{EW}} \end{cases} \Rightarrow \Delta V = \frac{\lambda}{4} v^4 \sim 10^{55} \rho_{\Lambda}^{\text{obs}}$$

Severe fine tuning! (but not of $\mathcal{O}(10^{120})$)

Dark energy sourced by fields

- Postulate a field with $V(\phi)$ and a cancellation:

$$\Lambda_0 M_{\text{P}}^2 + \text{vacuum energy} = 0 \quad (\textit{strictly})$$

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Quintessence: [Wetterich (87); Ratra (87); Peebles (87)]

Pros

- Avoids fine tuning of Λ_0

Cons

- $m_{\phi} \sim 10^{-33} \text{eV}$
- $\lambda_{\phi} \sim 0$ (no 5th forces)
- Sensitivity to initial cond.

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A very active field of research

The conspiracy of scales

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The conspiracy of scales

- Starobinsky inflation and the Higgs

$$\sqrt{-g}\mathcal{L} = \frac{M_{\text{P}}^2}{2}R + \frac{1}{16\alpha^2}R^2 - V(h); \quad V(h) = \begin{cases} \frac{\lambda}{4}v^4, & T > T_{\text{EW}} \\ 0, & T < T_{\text{EW}} \end{cases}$$

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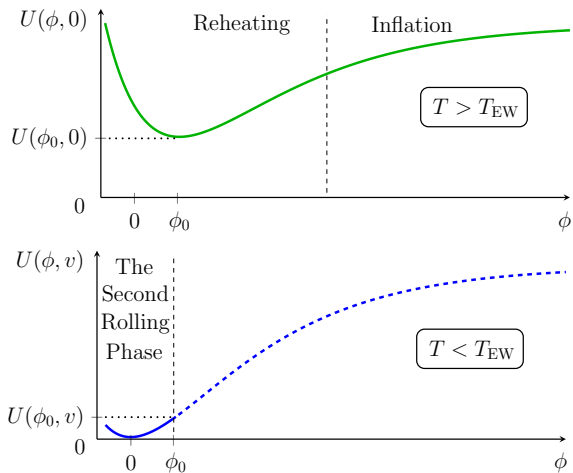
- The *scalaron* ϕ couples to $V(h)$!

$$\sqrt{-g}\mathcal{L} = \frac{M_{\text{P}}^2}{2}R - \frac{1}{2}(\partial_{\mu}\phi)^2 - \alpha^2 M_{\text{P}}^4 \left(1 - e^{-\sqrt{\frac{2}{3}}\phi/M_{\text{P}}}\right)^2 - e^{-\sqrt{\frac{8}{3}}\phi/M_{\text{P}}}V(h)$$

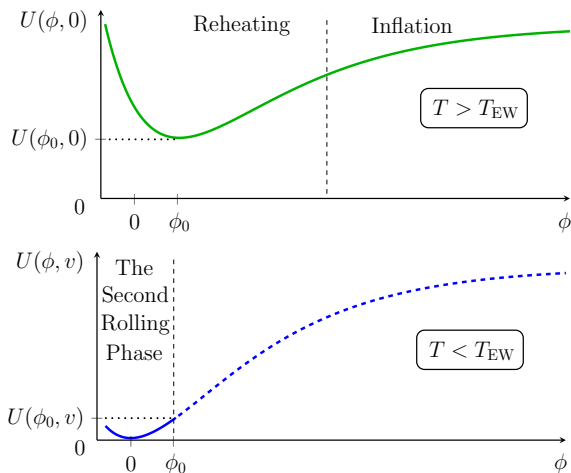
- **Minimum at $\phi \neq 0$ when $T > T_{\text{EW}}$:**

$$e^{\sqrt{\frac{2}{3}}\phi_0/M_{\text{P}}} = 1 + \frac{\lambda v^4}{4\alpha^2 M_{\text{P}}^4}$$

The conspiracy of scales



The conspiracy of scales



- Initial potential energy at $T = T_{EW}$

$$U(\phi_0, v) \approx \frac{V^2(0)}{\alpha^2 M_{\text{P}}^4} = \frac{\lambda^2 v^8}{16\alpha^2 M_{\text{P}}^4} \approx 4.0 \cdot 10^{-48} \text{GeV}^4$$

The conspiracy of scales

- A (very) surprising relation:

$$\rho_{\Lambda} \approx \frac{v^8}{\mathcal{P}_{\zeta} M_{\text{P}}^4}$$

$$(v \approx 246\text{GeV}, \mathcal{P}_{\zeta} \approx 2.2 \cdot 10^{-9}, M_{\text{P}} \approx 2.44 \cdot 10^{18}\text{GeV})$$

- Similar to

$$v \approx \sqrt{\rho_{\Lambda}^{1/4} M_{\text{P}}}$$

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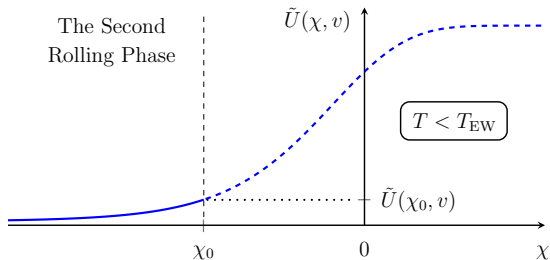
- Unfortunately, the scalaron is heavy

$$m_\phi \gg H$$

Need to prevent further rolling of ϕ !

The bait-and-switch

- Make kinetic term non-canonical



$$(\partial_\mu \phi)^2 \quad \longrightarrow \quad \left(b \frac{M_{\text{P}}}{\phi} \right)^2 (\partial_\mu \phi)^2 \equiv (\partial_\mu \chi)^2$$

- Potential gets stretched

$$\tilde{U}(\chi_0, v) \sim \rho_\Lambda, \quad \tilde{U}''(\chi_0, v) \sim H^2$$

The bait-and-switch

- Minimum at $\chi_0 \sim 10^2 M_{\text{P}}$
- ⇒ Potential cannot be stretched before the minimum is reached

Bait-and-switch mechanism: [arXiv:1807.04359]

Trigger stretching via the Higgs at $T = T_{\text{EW}}$

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Pros

- New parameters $\mathcal{O}(1)$
- All couplings get suppressed (no 5th forces)
- No sensitivity to initial conditions

Cons

- No top-down motivation
- Higgs coupling sensitive to colliders bounds

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Let's try a multi-field model!

Assisted relaxation

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Assisted relaxation

- A model with *two* fields, $\mathcal{L} = \mathcal{L}_\phi + \mathcal{L}_\sigma$; [arXiv:1807.04359]
 - Inflaton, $V(\phi) = \frac{1}{2}m^2\phi^2$, and 'darkon' with:

$$-\sqrt{-g}\mathcal{L}_\sigma = b^2 \frac{1}{2} \frac{M_{\text{P}}^2}{\phi^2 + \sigma^2} \partial_\mu \sigma \partial^\mu \sigma + \frac{1}{2} m^2 \sigma^2 + c \frac{\sigma}{M_{\text{P}}} V(\phi)$$

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Same attractive features as previously, but avoids collider bounds

Benefits from multiple fields

- Many realizations when

$$\mathcal{L}_\phi = -G_{ab}(\phi)\partial_\mu\phi_a\partial^\mu\phi_b - V(\phi)$$

⇒ Rich model-building possibilities

- **Talks:** Achúcarro, Sfakianakis, Turzynski, Welling
- **Posters:** Fumagalli, Christodoulis, Chiovoloni, Pinol, Ronayne, Wang

- Generalized Einstein theories

$$\mathcal{L}_{\text{GE}} = \frac{M_{\text{P}}^2}{2}R - \frac{1}{2}\partial_\mu\phi\partial^\mu\phi - e^{-2F(\phi)}\frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma - U(\phi, \sigma)$$

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Thank You!