

Primordial black holes



Christian Byrnes
University of Sussex, Brighton, UK



PBH collaborators include: J. Adamek, N. Bellomo, P. Cole, E. Copeland, A. Gow, M. Gosenca, A. Green, A. Hall, M. Hawkins, S. Hotchkiss, M. Hindmarsh, I. Musco, S. Patil, J. Peacock, D. Regan, M. Sasaki, D. Seery, S. Young

Cosmo19, Aachen, 2 September

Primordial Black holes (PBHs)

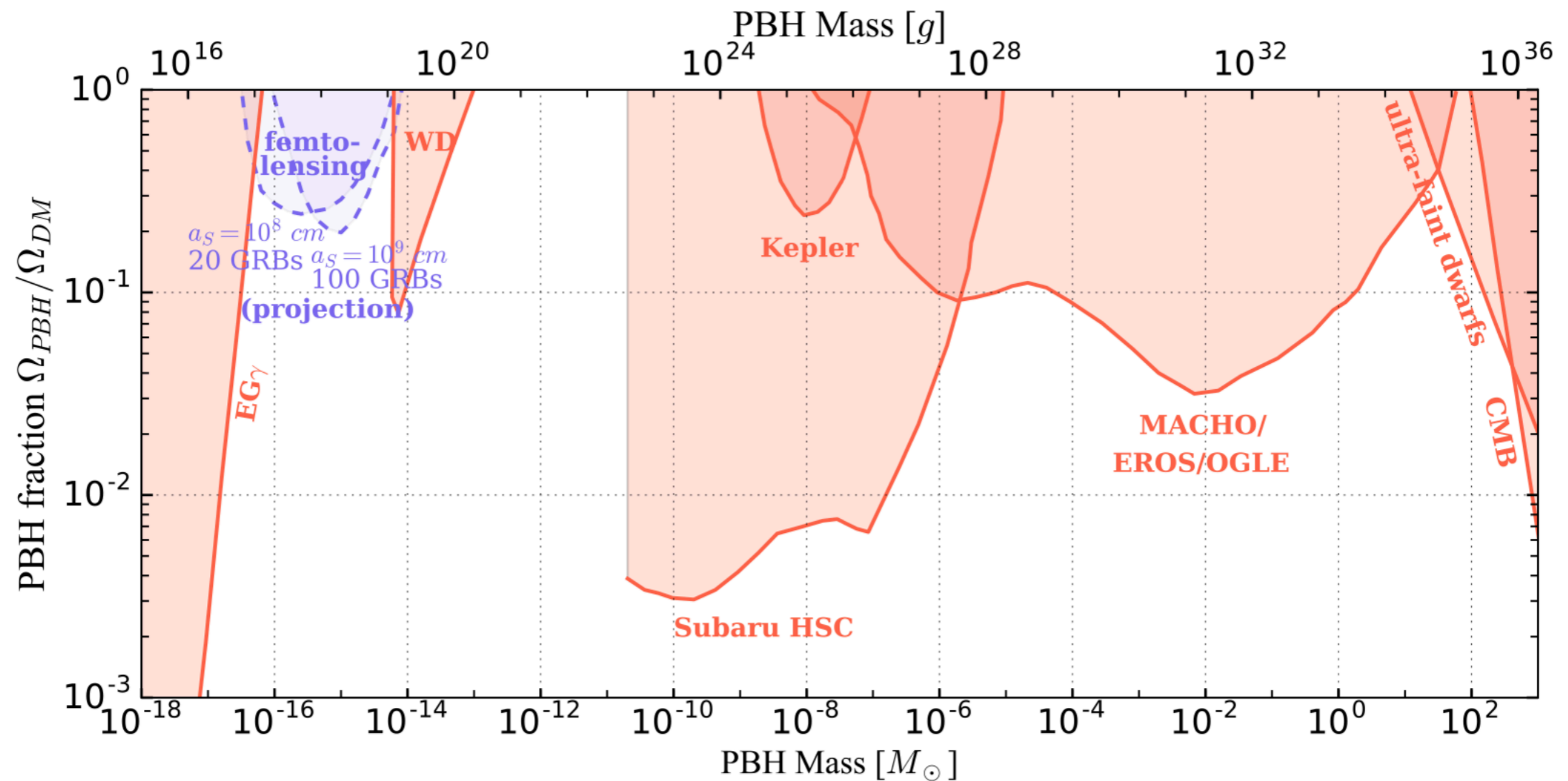
- Are not stellar remnants
- Naturally have all the properties required of dark matter
- Lots of interesting phenomenology: they could change early astrophysics, seed supermassive black holes, generate Hawking radiation...
- Require special initial conditions to form
- They might not exist, but are useful anyway

Zel'dovich and Novikov 1967; Hawking 1974; Carr and Hawking 1974

Contents

- Are PBHs the dark matter?
- Observational hints and constraints
- How can a PBH be distinguished from a stellar remnant?
- How can they form?
- The small scale perturbations
- Future directions

Constraints on PBHs as DM



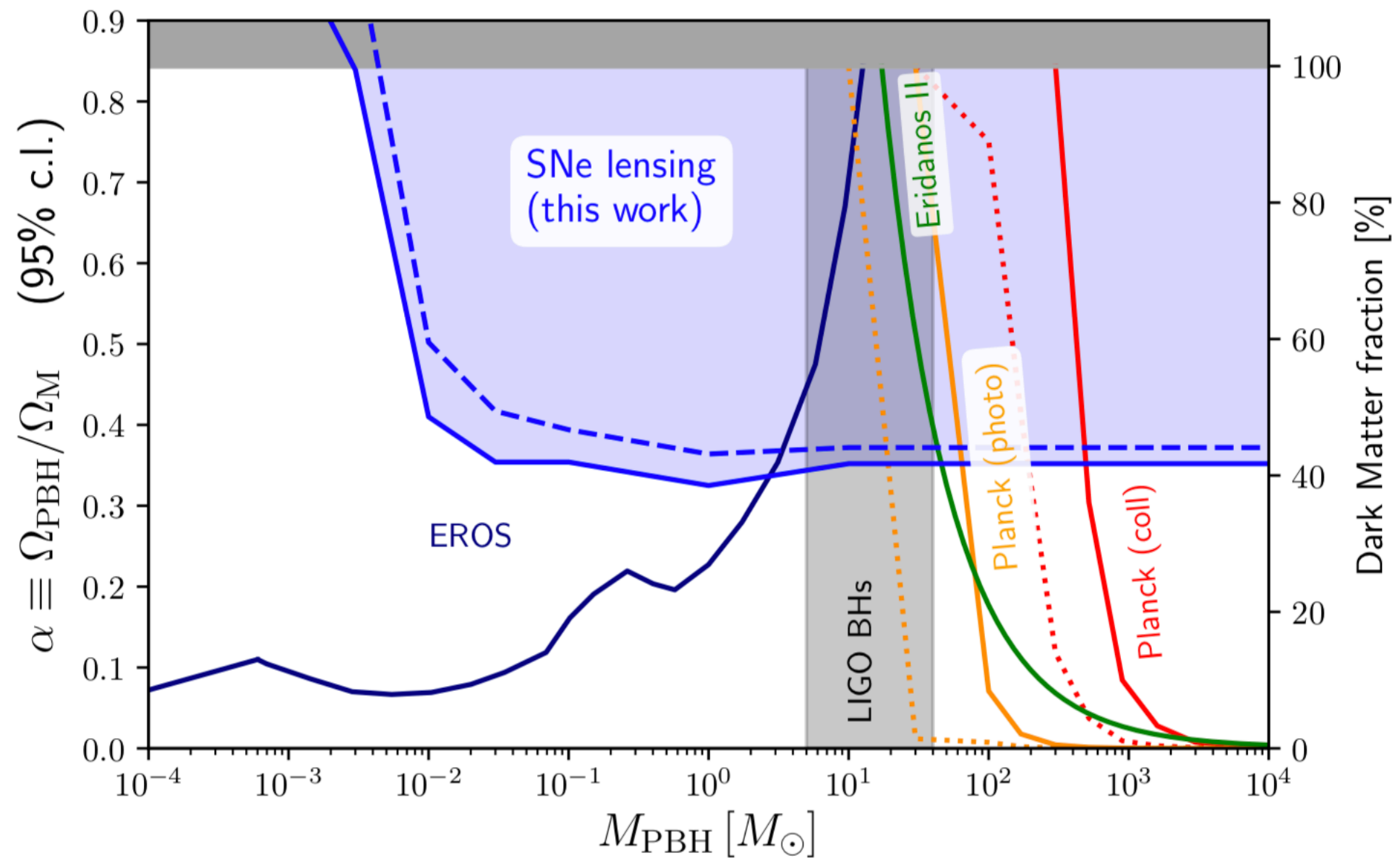
Katz et al 2018 “Femtolensing revisited”; see also Sasaki et al 2017 Review +++

The constraints have changed and in several cases weakened upon reanalyses
 Lots of people have worked on the constraints

Constraints on PBHs as DM

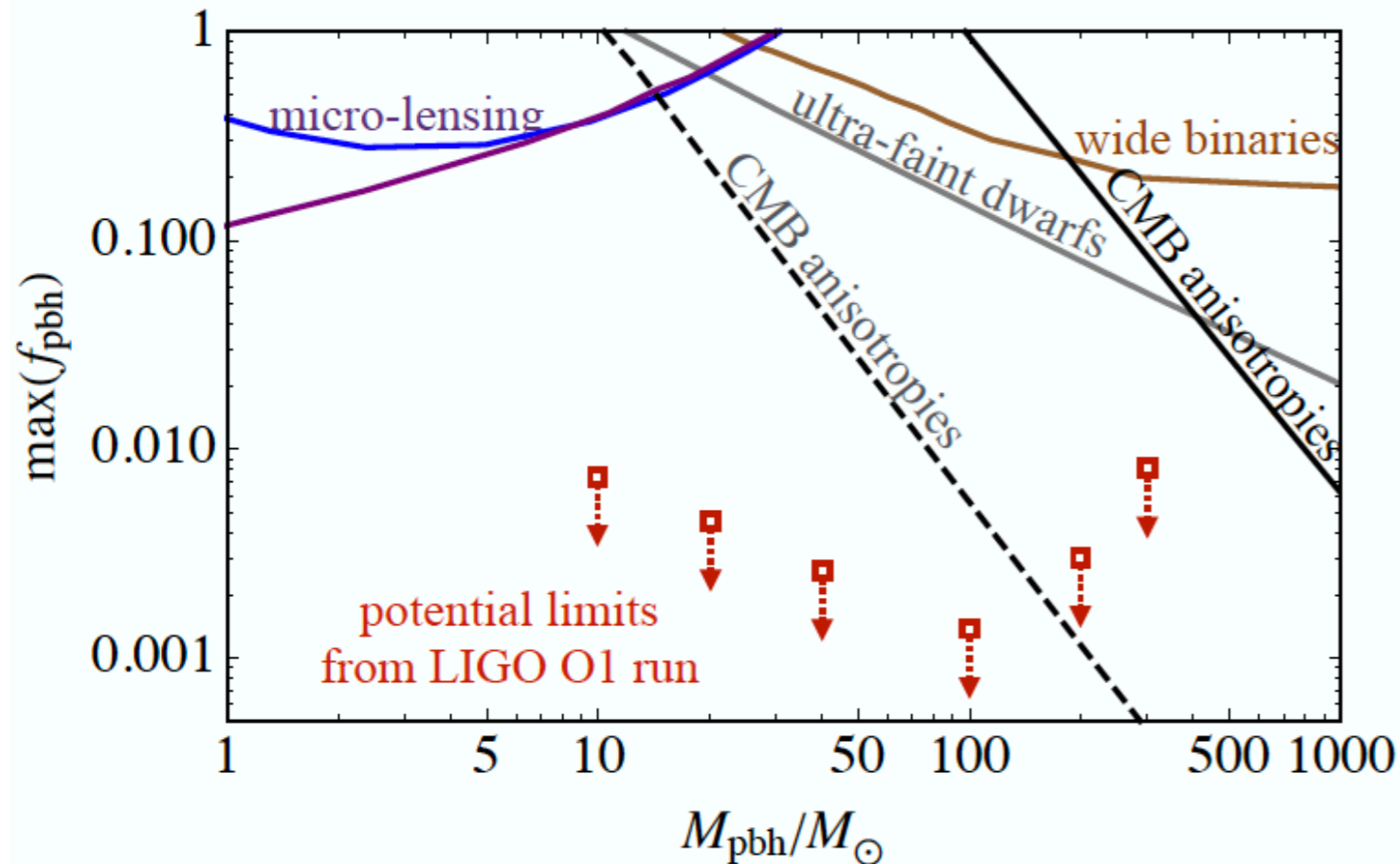
- There is evidence that PBHs cannot form all of the dark matter except in (most of) a mass range 10^{-15} - 10^{-11} M_{sun}
- Although “light”, this is the heaviest DM candidate by about 20 orders of magnitude
- Previous microlensing constraints had suggested this was already ruled out, but these neglected finite source and light wavelength effects. *Niikura et al 2018*
- Exception: BHs could leave Planck mass relics when decaying, these are a potentially untestable “nightmare” DM scenario. *MacGibbon 1987*
- Constraints normally assume an (unrealistic) monochromatic mass spectrum. Critical collapse creates a spread of masses. *Yokoyama '98, Nieyemeyer & Jedamzik '98*
- The constraints often tighten for broad mass spectra, except in some cases with multiple peaks tuned to “fit in the observational gaps”, e.g. *Bellomo et al 2018*

Constraints on the LIGO mass range



Zumalacárregui and Seljak 2018

The PBH merger rate places the tightest constraint



Haimoud et al 2017

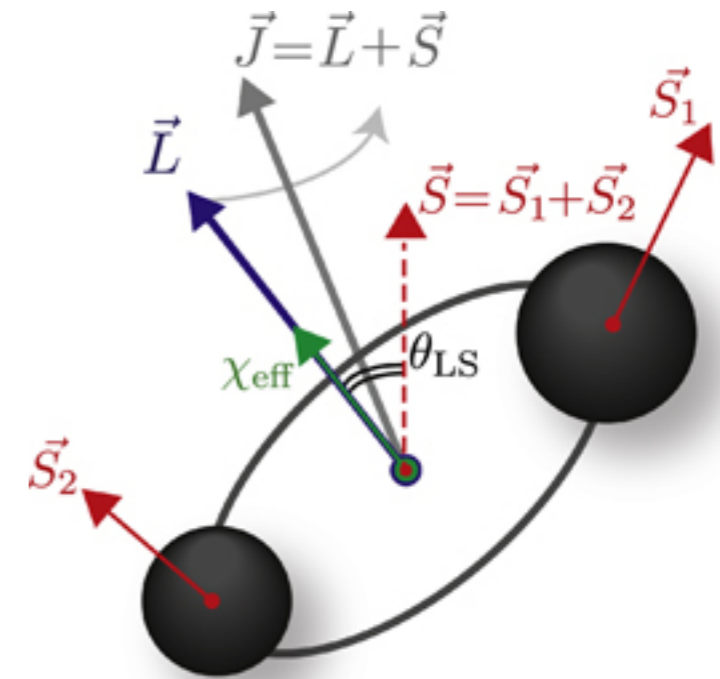
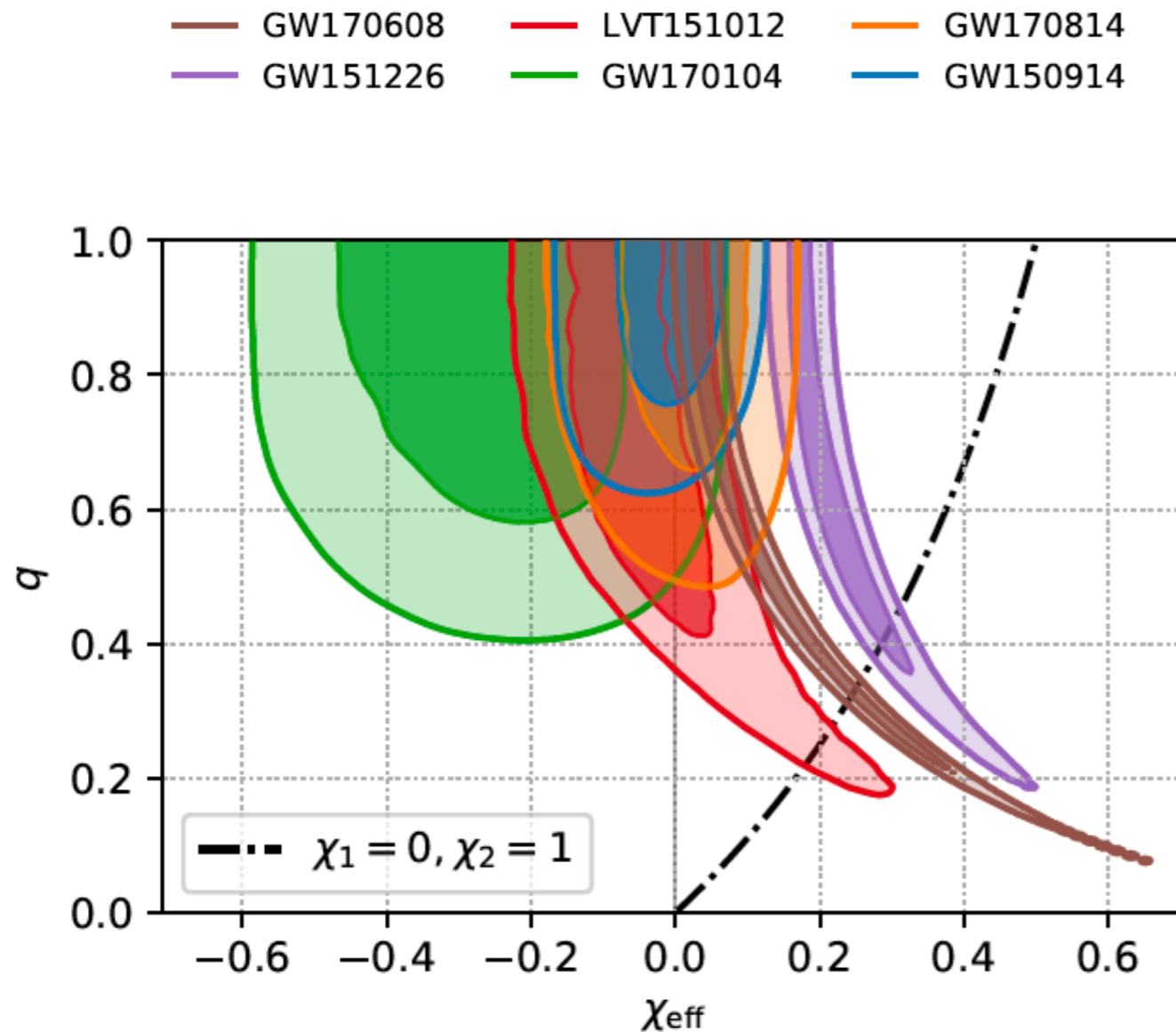
- Caveats:

1. Assumes a monochromatic mass spectrum. Extended by *Chen & Huang '18, Raidal et al '18 ++*
2. Assumes PBHs are randomly placed initially, true if Gaussian initial conditions. Clustering does not help *Bringmann et al '18*
3. Assumes BH binaries are not disrupted. Recently tested to $z \sim 1000$ by simulations (*Raidal et al '18*) and even disrupted PBHs can merge *Vaskonen & Veermäe '19*
4. Neglects halo formation around the BHs. Not a huge effect overall *Kavanagh, Gaggero & Bertone '18*

Hints for PBHs

- PBHs might not exist, but there are some hints
Clesse & Garcia-Bellido 2017
- The “unexpected” masses and low spins of those LIGO/Virgo detected created an explosion of interest.
“Did LIGO detect dark matter” Bird et al;
see also Clesse & Garcia-Bellido; Sasaki et al; all 2016
- The existence of supermassive black holes in almost all galaxies, with unknown formation process. However, creating a heavy PBH seed requires non-Gaussian initial conditions to evade CMB μ -distortion constraints. *Carr, Nakama & Silk 2017*
- Correlation between the cosmic infrared background and unresolved cosmic X-ray background. *Kashlinsky et al Astro2020 white paper*

Black hole spin



$$\chi_{\text{eff}} = \frac{c}{G(m_1 + m_2)} \left(\frac{\vec{S}_1}{m_1} + \frac{\vec{S}_2}{m_2} \right) \cdot \vec{L}$$

$$a_* = \frac{c |\vec{S}|}{Gm^2} \leq 1$$

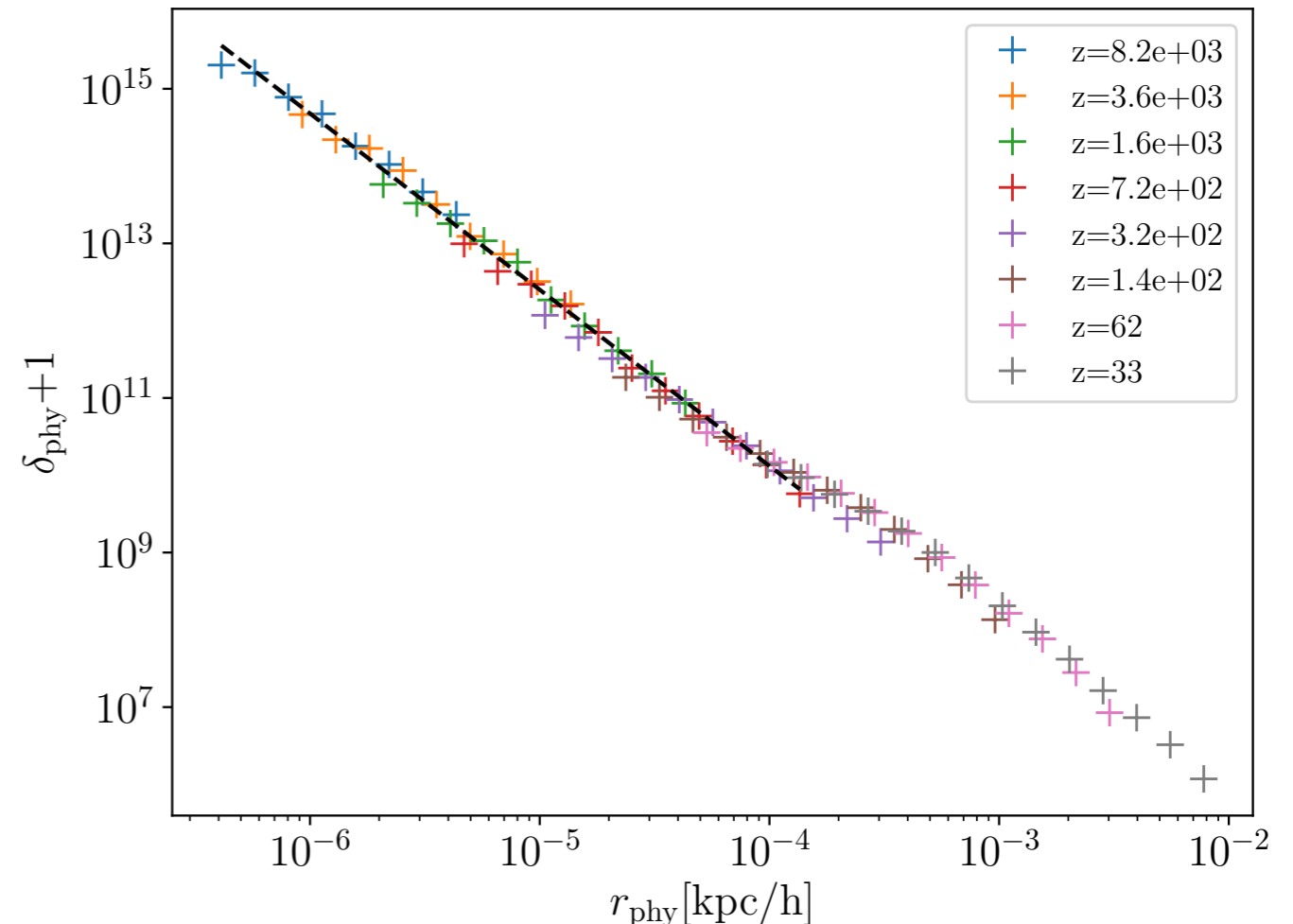
Negligible spin is expected for PBHs formed during radiation domination. Unlike astrophysical BHs, PBHs do not undergo much collapse before formation. See the talk by de Luca. *Chiba & Yokoyama 2017; Belczynski et al. 2017; Mirbabayi et al 2019; De Luca et al 2019; Fernandez & Profumo 2019*

What if one PBH was detected?

- We would know what some of the DM was (and what some of it is not)
- It would be the oldest relic detected, predating the primordial element abundance generated by BBN
- Other potential relics include: gravitational waves, topological defects such as cosmic strings, CMB spectral distortions, ultracompact minihalos
- Requires non-trivial inflationary dynamics, perhaps with an early-matter-dominated phase and/or topological defects.
PBH review article: Green 2015

WIMPs and PBHs are incompatible

- Assuming WIMPs have the standard, velocity independent cross section which gets the right abundance, and $M_{\text{PBH}} > 10^{-6} M_{\text{sun}}$. And annihilation channel into gamma rays.
- If $f_{\text{PBH}} < 1$, then another DM component is inevitable
- Steep and high density profiles form around PBHs (density $\sim r^{-9/4}$). WIMPs would rapidly annihilate in them.
- In contrast to ultracompact minihalos without a PBH seed. *Gosenca et al '17*, *Delos et al '17*
- See talk by Derek Inman on 3D simulations with PBHs + lighter DM
- A detection of WIMPs or PBHs may effectively rule out the existence of the other



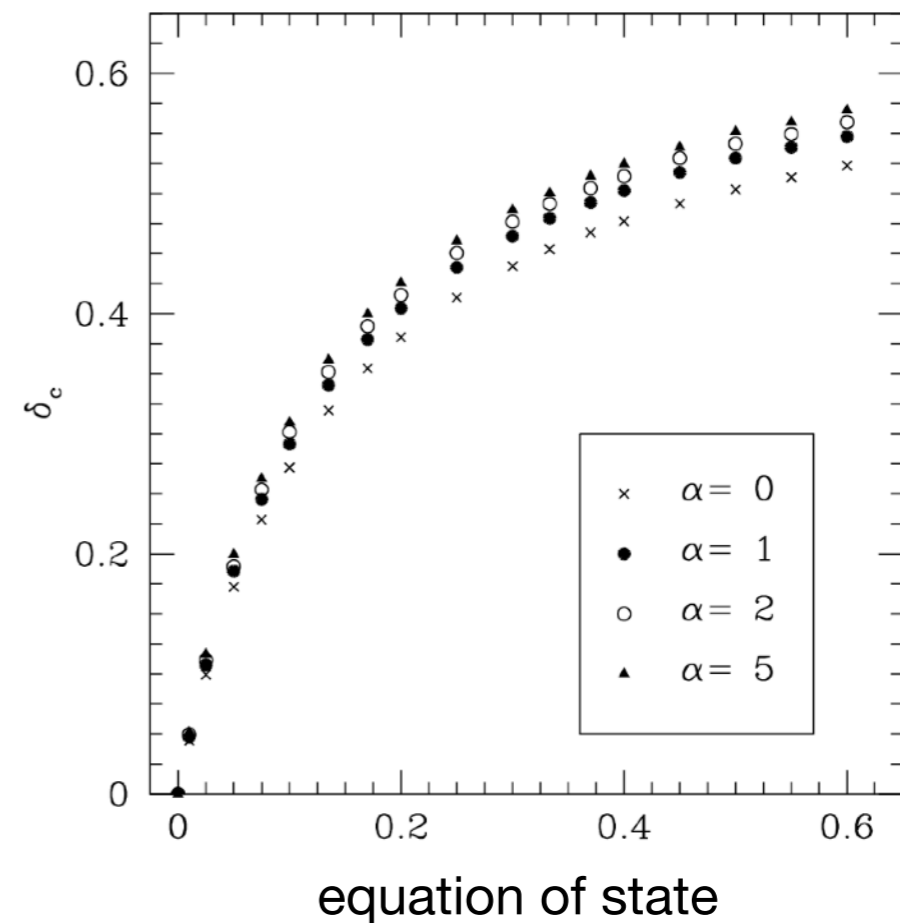
Adamek, CB, Gosenca & Hotchkiss 2019;

*Lacki & Beacom 2010; Eroshenko 2016;
Boucenna, Kühnel, Ohlsson & Visinelli 2017;
Bertone et al 2019*

PBH formation

1. They form from large amplitude density perturbations shortly after horizon entry
2. Causality prevents collapse while the perturbations are super-horizon
3. Approximate 1-to-1 relation exists between horizon entry time, horizon length and PBH mass

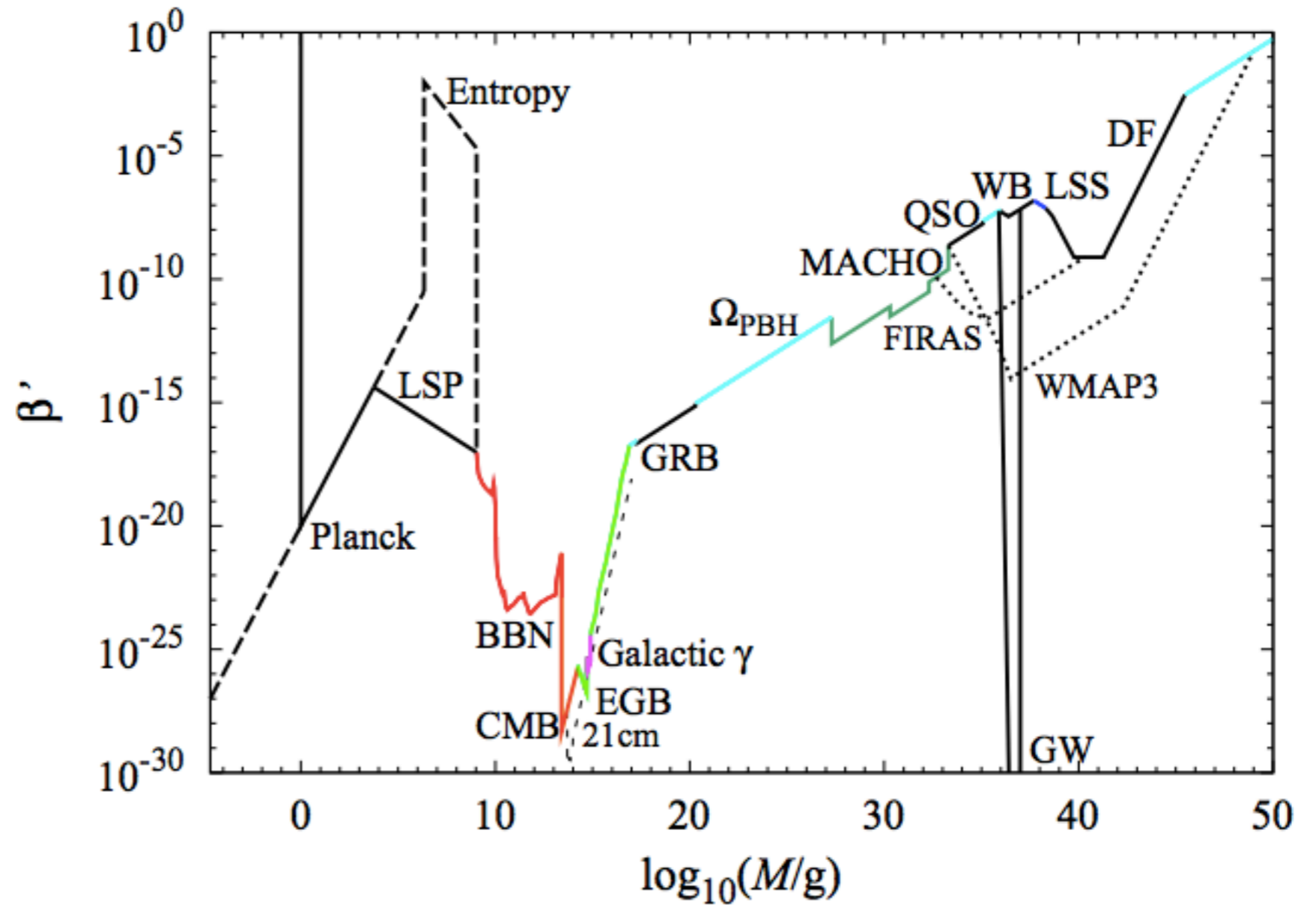
Collapse threshold



Musco and Miller 2013

The PBH fraction at formation is tiny

$$\beta' = e^{-\frac{\delta_c^2}{2\mathcal{P}_\delta}}$$



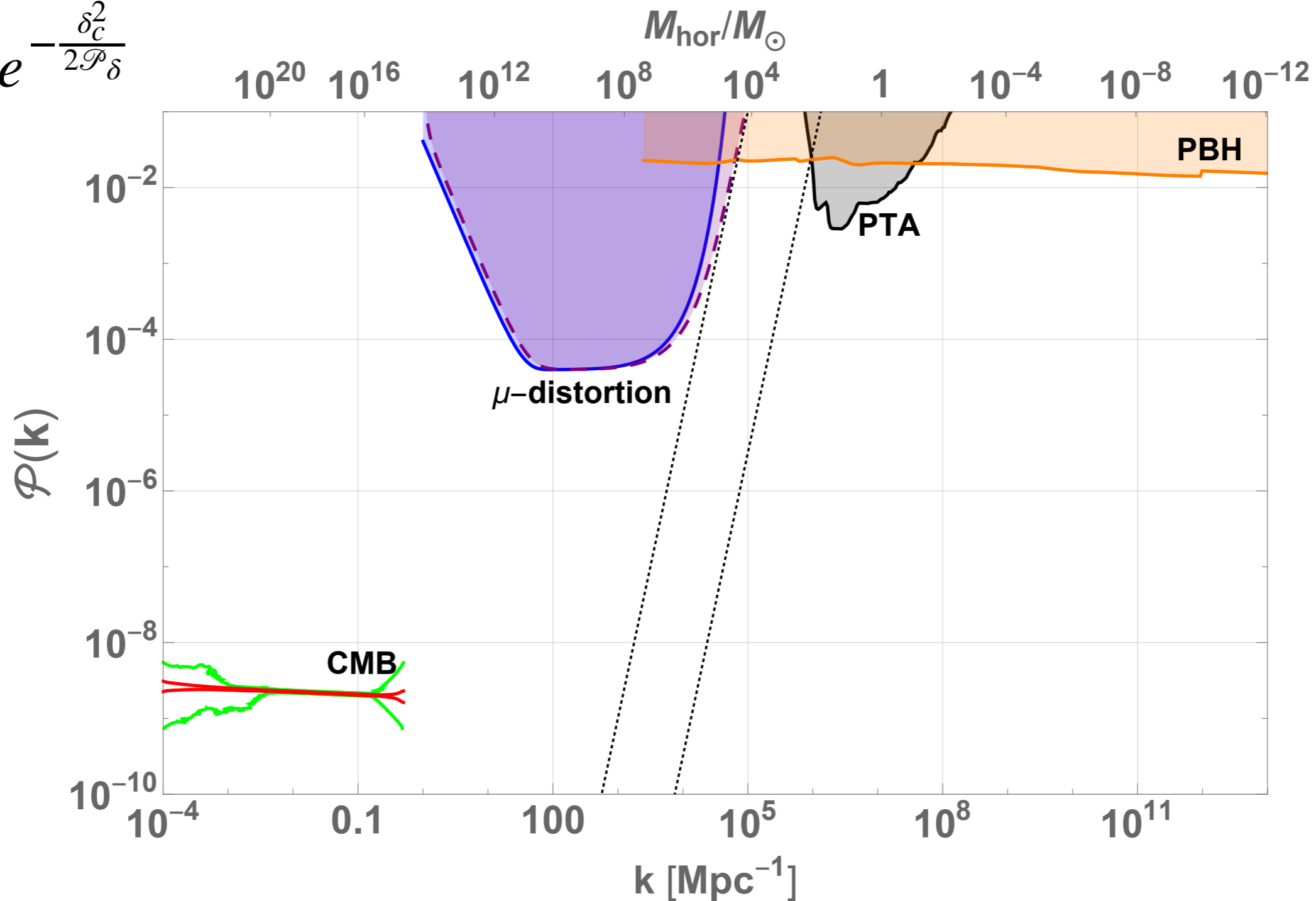
Carr et al 2009

$$\rho_{PBH} \propto a^{-3} \quad \rho_{rad} \propto a^{-4}$$

Ω_{PBH} grows like the scale factor from formation until radiation-matter equality

Power spectrum constraints

$$\beta' = e^{-\frac{\delta_c^2}{2\mathcal{P}_\delta}}$$



CB, Cole & Patil '18; See also Inomata and Nakama '18

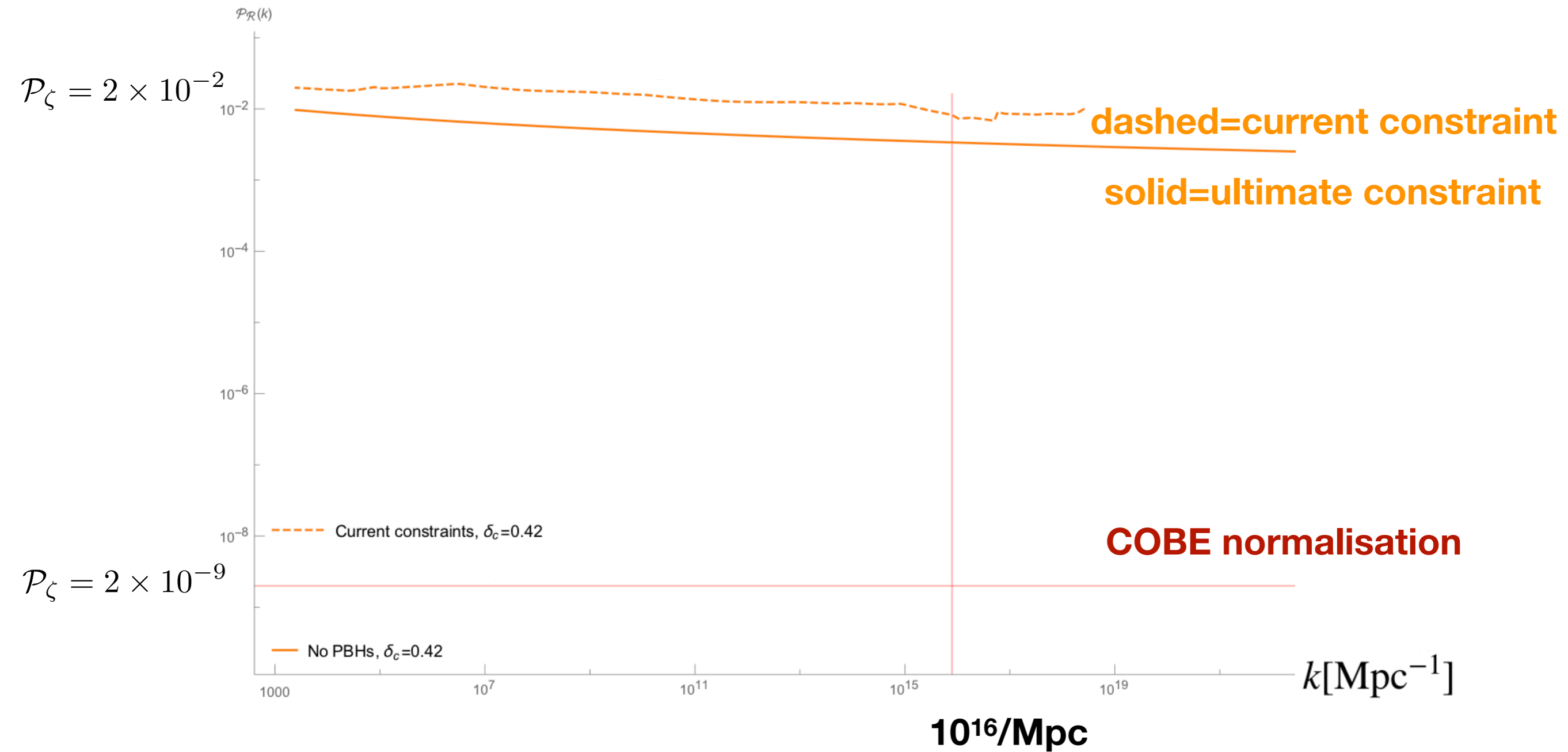
Small scale perturbations

- We have precision measurements from the CMB and large scale structure on scales > 1 Mpc (which are “linear” today)
- These probe about 2 decades in length scales, but inflation must have lasted much longer
- Radiation pressure/non-linearities erases the memory of smaller scale perturbations, unless there is a relic
- The large scales are well fit by a simple power-law spectrum
- The small scale perturbations could contain lots of new information

PBH formation comments

- The formation rate is exponentially sensitive to the amplitude of the power spectrum, and the collapse threshold
- Inflationary models posit an inflection point (ultra-slow-roll inflation) or other feature (see *talks by Pattison and Jain later today*).
- The power spectrum can't grow faster than about k^4 (in canonical single-field inflation), impacts the constraints. *Byrnes, Cole & Patil '18; Carrilho, Malik & Mulryne '19*
- PBHs are very rare - very sensitive to non-Gaussianity
- The formation criteria depends on the density profile. Many spherically symmetric simulations exist, e.g. *Niemeyer & Jedamjik, Musco & Miller, Harada ++, Nakama ++...*
- Extensive recent analytic work has been done to relate the power spectrum to PBH formation rate at, but (at least) an order unity uncertainty remains (= tens of orders of magnitude in terms of the formation rate). *Germani & Musco '17, Yoo et al '17, Kawasaki & Nakatsuka '19, de Luca et al '19, Young et al '19, Young '19, Kalaja et al '19*

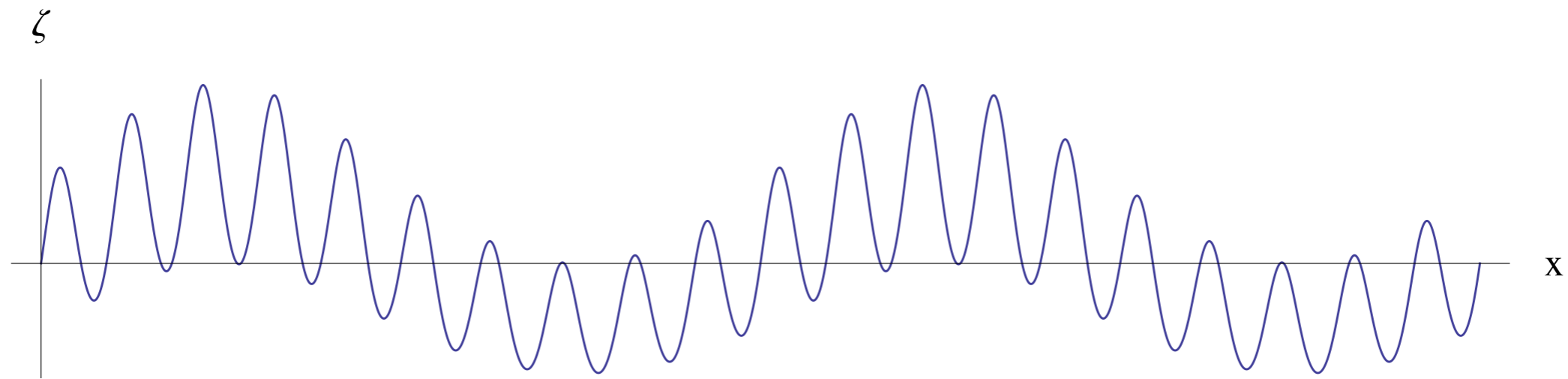
Ultimate constraints from PBHs



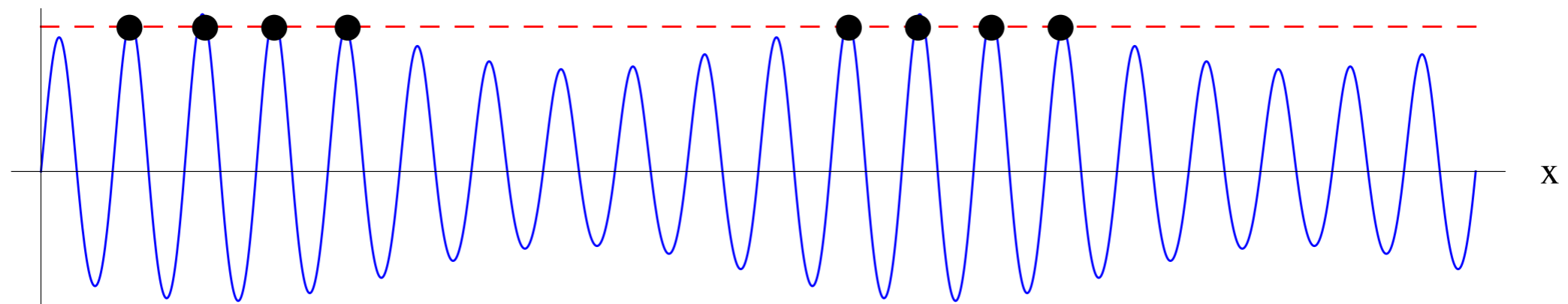
$$\beta' = e^{-\frac{\delta_c^2}{2\mathcal{P}_\delta}}$$

Cole & CB '17

Isocurvature constraints



ζ subtracting the long wavelength mode



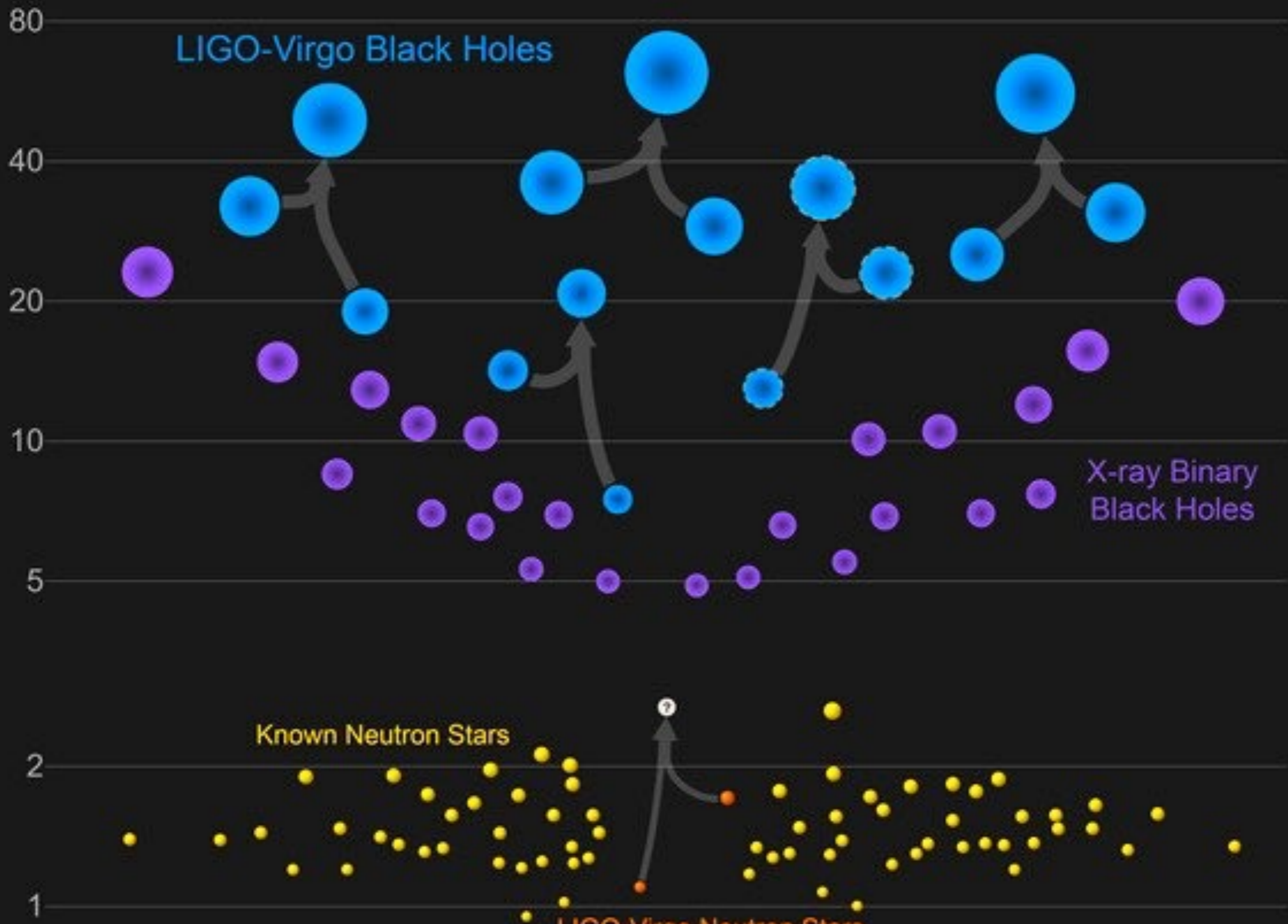
Mode coupling between CMB and PBH forming scales is ruled out, since that would generate a huge DM isocurvature perturbation - *Tada & Yokoyama 2015, Young & CB 2015*

Black holes have no hair

- Hence, a BH does not “remember” how it formed
- We can only measure the mass and spin
- Fortunately, stellar remnant black holes cannot form with any possible mass
- There is expected to be a lower bound around $5 M_{\text{sun}}$, and no compact objects (including neutron stars) below the Chandrasekhar mass, about $1 M_{\text{sun}}$
- A black hole merger at sufficiently high redshift would have to be primordial in origin. *Koushiappas & Loeb 2017*

Masses in the Stellar Graveyard

in Solar Masses



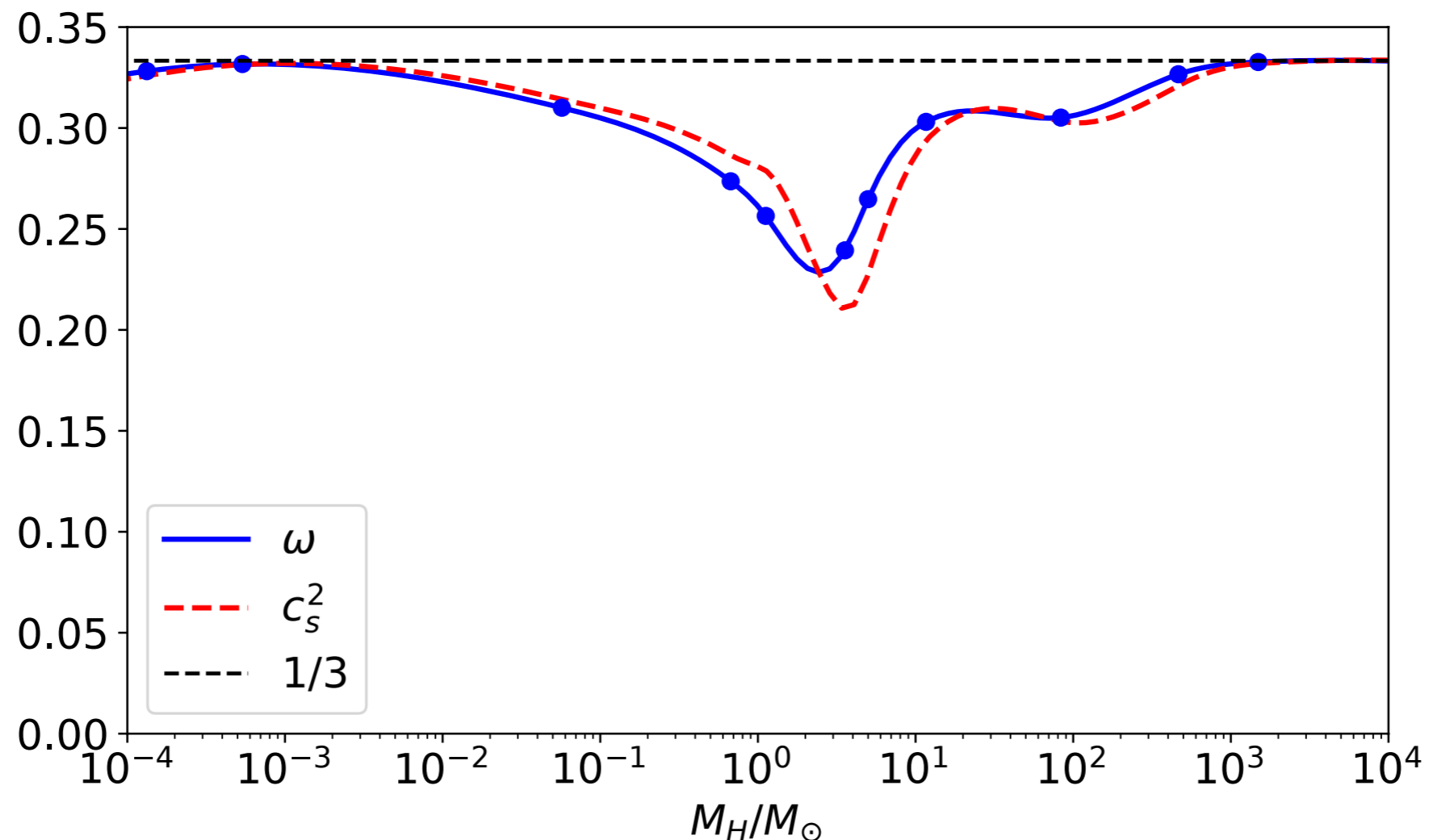
This plot is not up to date

LIGO-Virgo Neutron Stars

LIGO & Virgo collaboration

The QCD transition

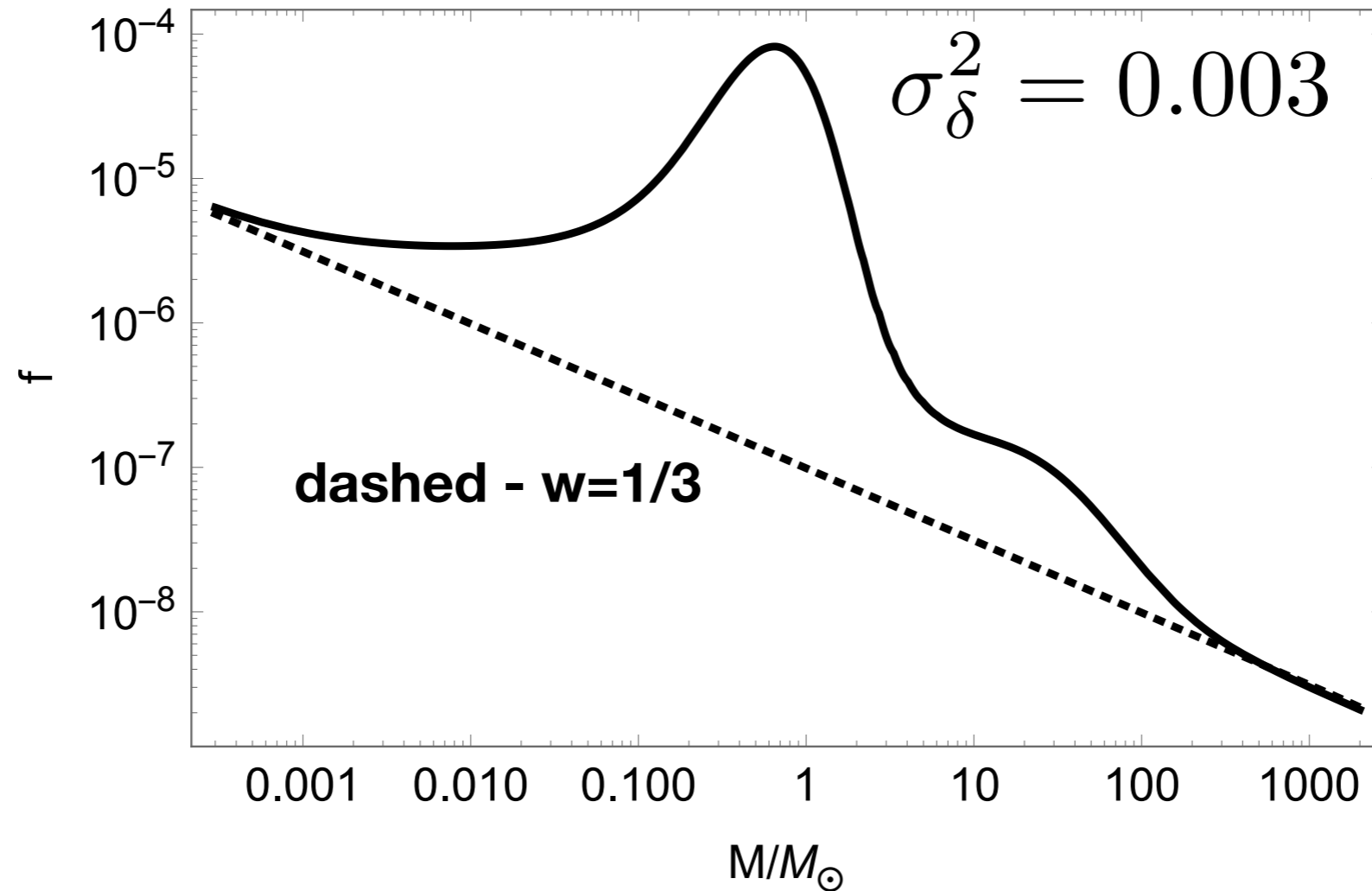
As the Universe cools below 1 GeV ($t \sim 10^{-6}$ s), strong interactions confine quarks into hadrons and the equation-of-state parameter w decreases. *Crawford & Schramm '82, Jedamzik '98*



CB, Hindmarsh, Young & Hawkins 2018 using Borsanyi et al 2016

The resultant PBH mass function

$$f(M) \propto M^{-1/2} e^{-\frac{\delta_c^2}{2\sigma_\delta^2}}$$

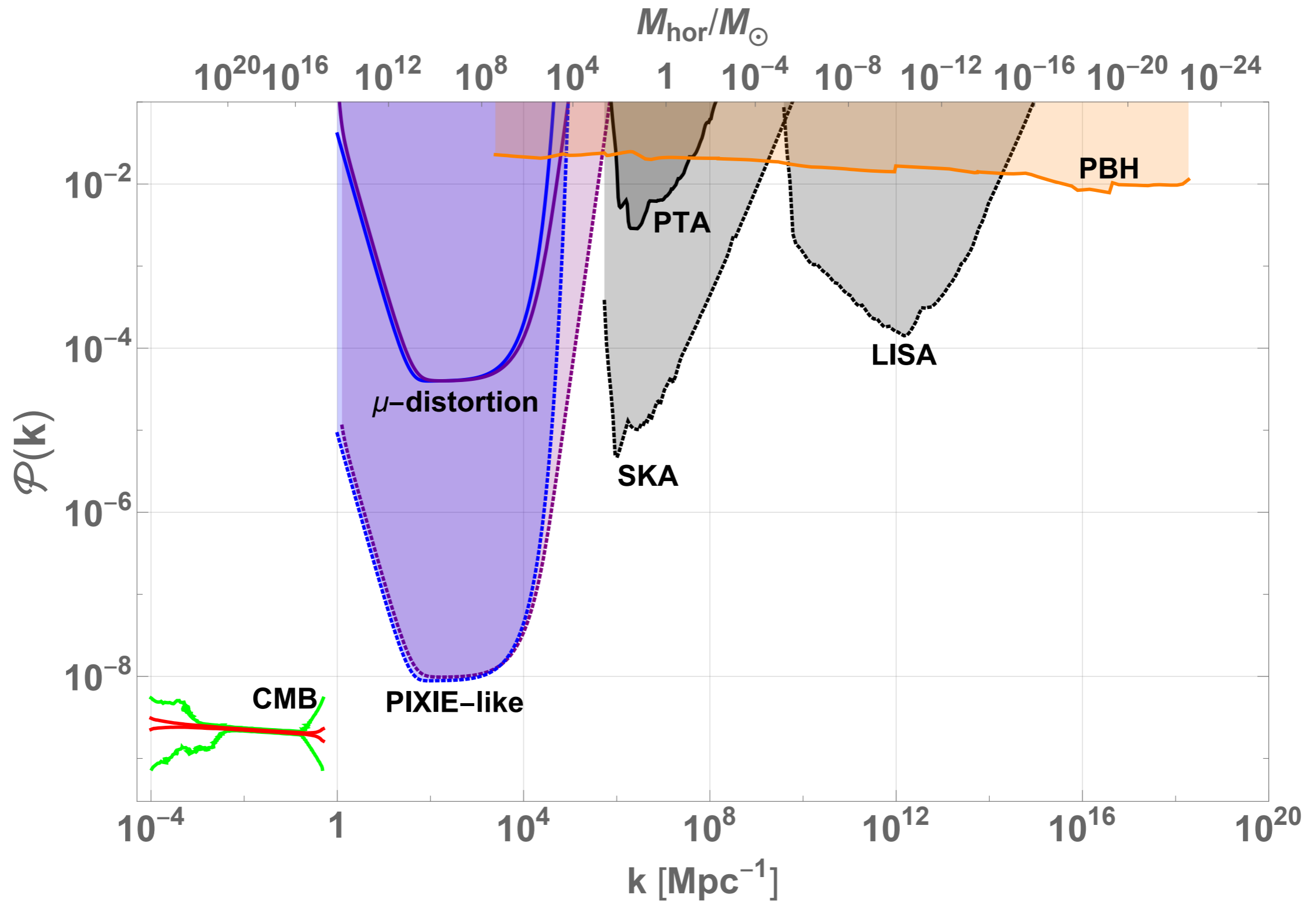


The QCD phase transition took place during the time when LIGO mass PBHs would have formed. **It boosts the formation rate of solar mass PBHs by 2 orders of magnitude**
No detection yet: *LIGO & Virgo collaboration 2019, Magee et al 2019*

Looking forwards

- LIGO/Virgo already have far more than the 10 analysed events, potentially including a BHNS merger and one in the upper mass gap (Tanja Hinterer's talk next)
- LISA and future ground based detectors will be sensitive to very high redshift mergers + better probe of the spins, mass ratio, etc
- Theoretical analysis of the formation criteria has greatly improved in recent years, but 3D simulations of PBH formation are still required (especially during early matter domination)
- See talks by Racco, Franciolini, Inomata, Terada and Dvorkin

Future constraints

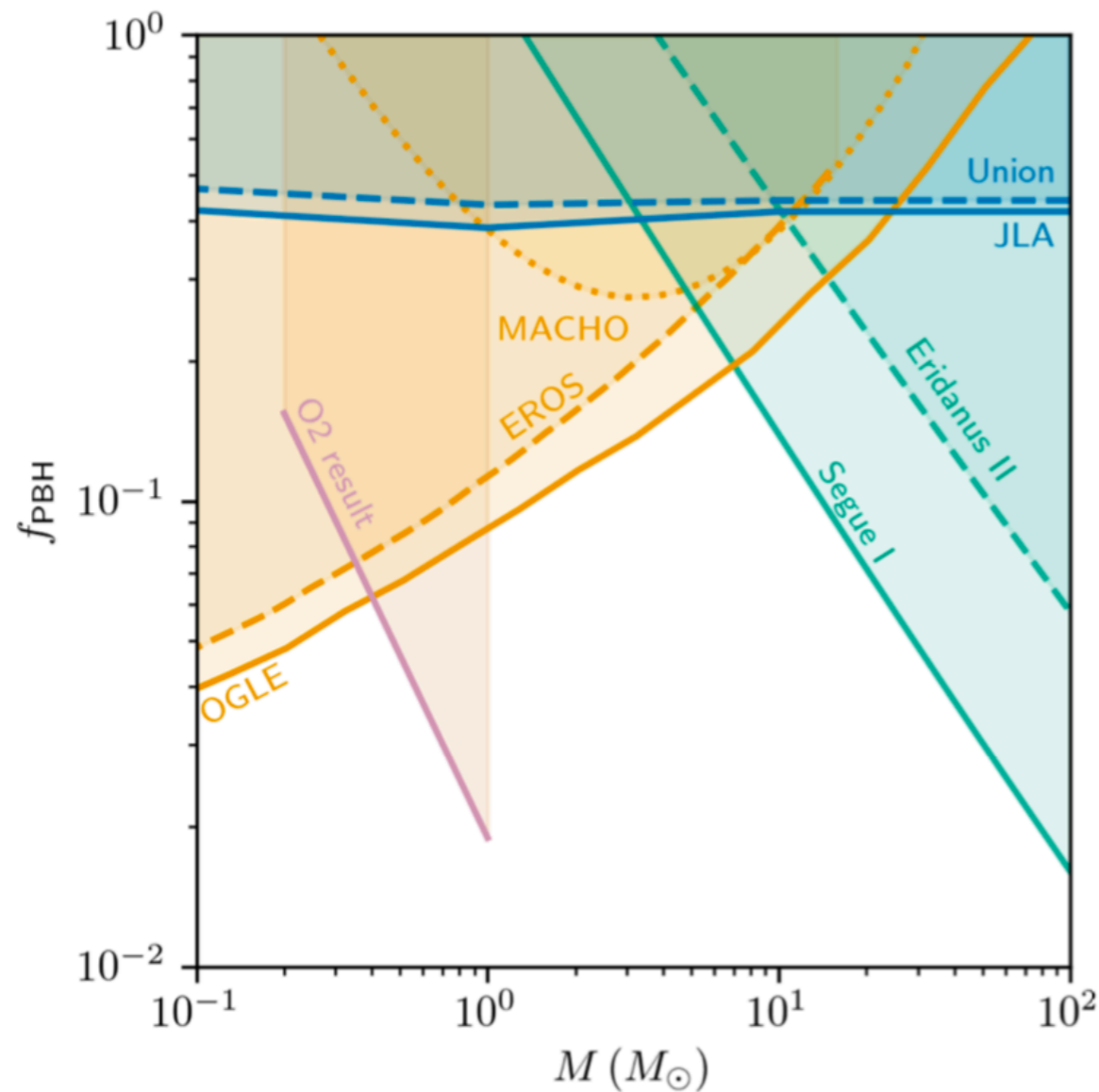


Summary

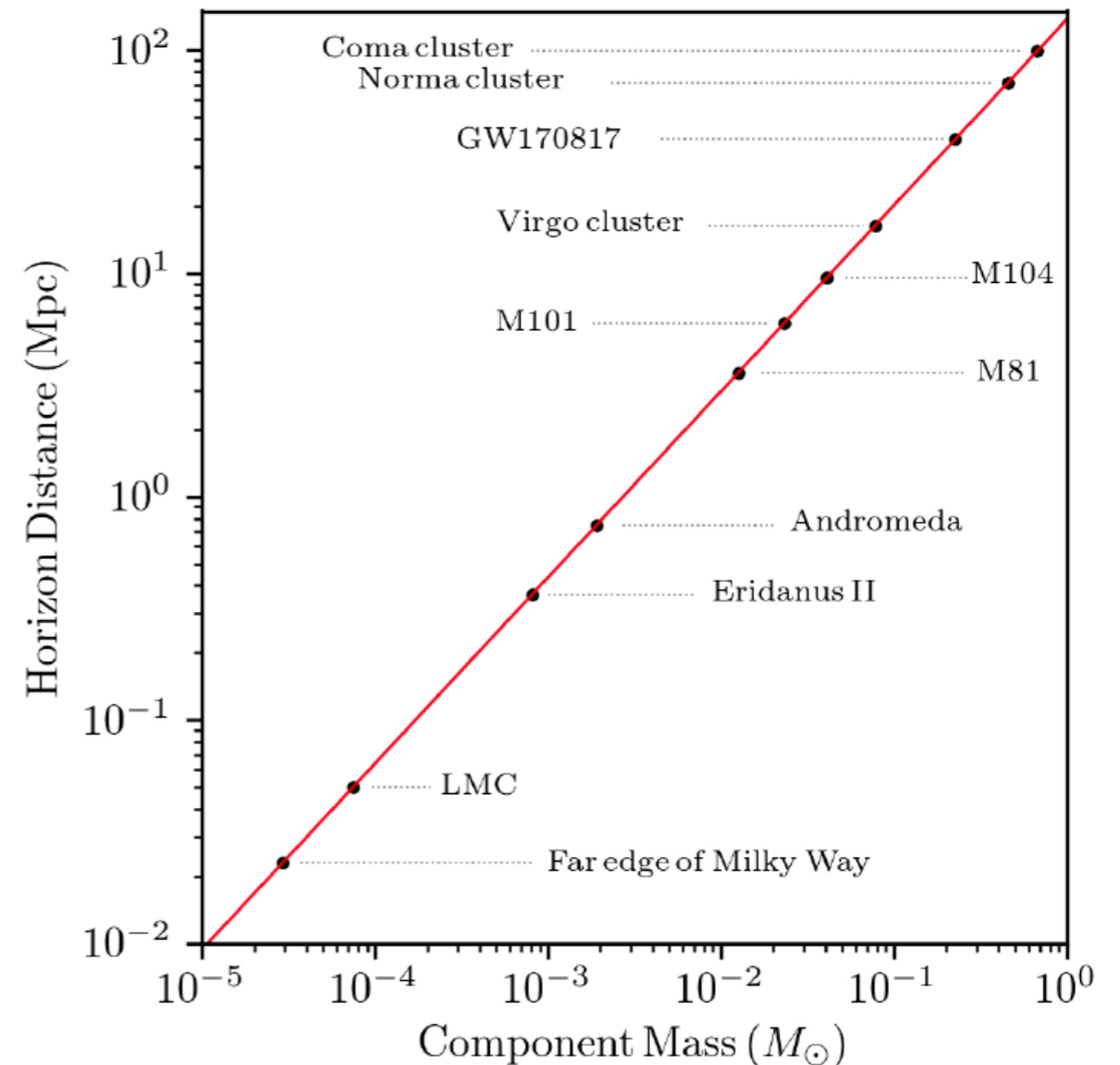
- Interest in PBHs has exploded since the first GW detection
- PBHs are used to probe the initial conditions and constituents of the Universe
- PBHs in the LIGO/Virgo mass range cannot be more than $\sim 1-10\%$ of the DM
- Asteroid mass PBHs could be all the DM, and LISA can test this scenario
- The detection of one PBH would be huge, and several “smoking gun” scenarios exist
- Even without a detection, PBHs constrain a greater range of scales than any other probe

Sub-solar mass GW searches

GW searches have been made, with no detections so far
 These are below the Chandrasekhar mass, hence potential proof of a primordial origin



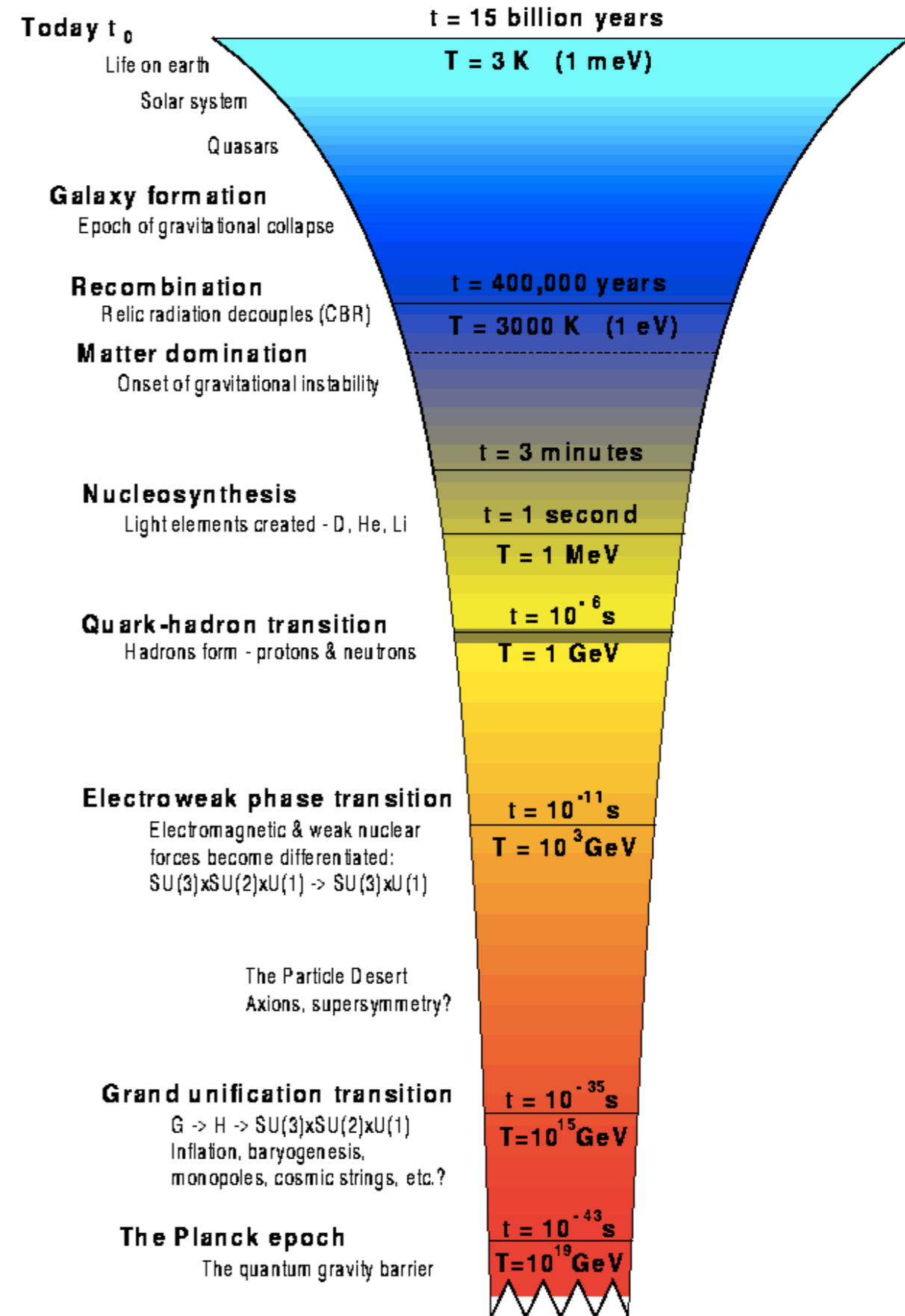
LIGO & Virgo collaboration 2019



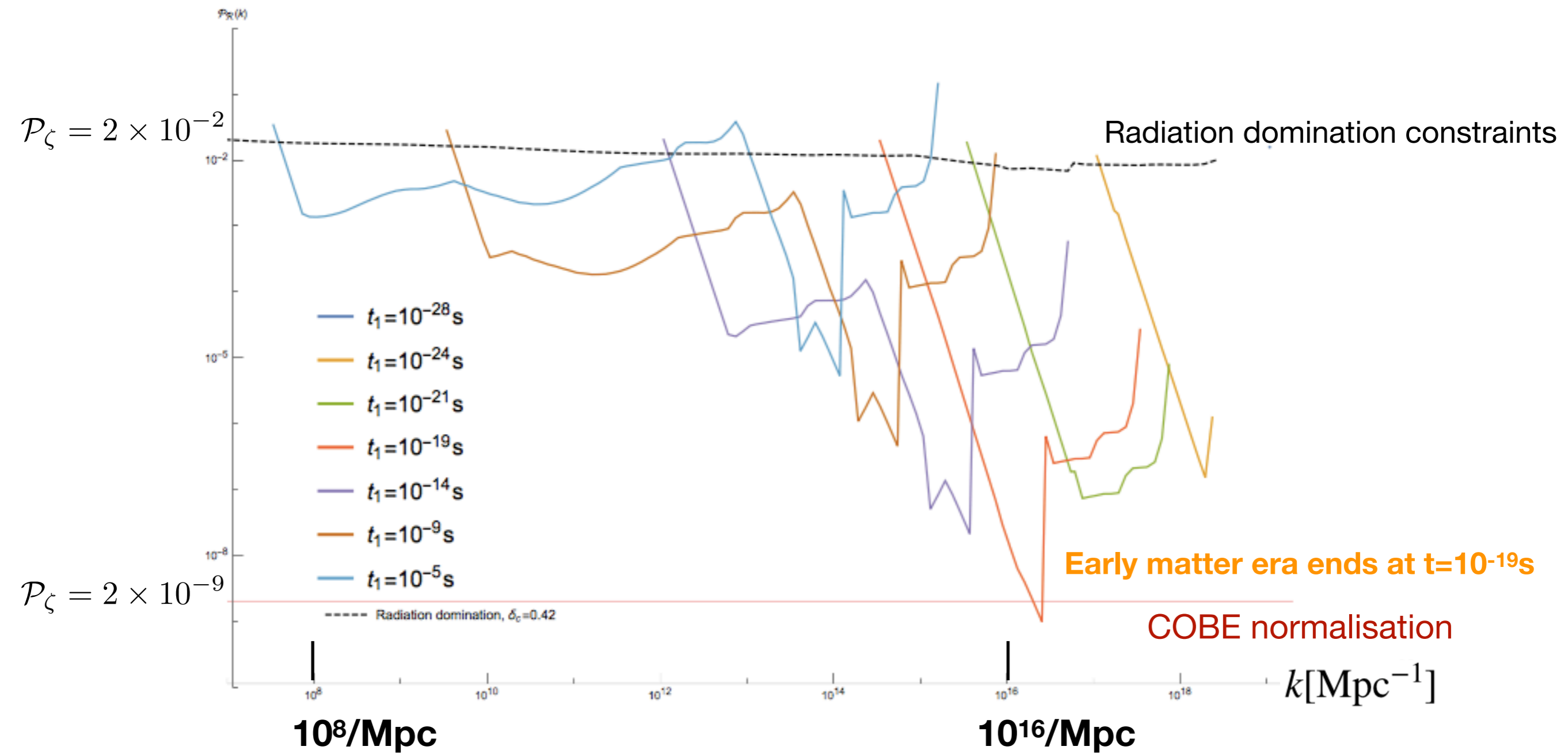
Magee et al 2018

A cosmological coincidence

- The QCD transition occurs during the time when LIGO mass PBHs formed.
- The horizon mass has grown by about 50 orders of magnitude since the end of inflation.
- QCD transition: $t \sim 10^{-6}$ s, $T \sim 200$ MeV, $M \sim 1 M_{\odot}$, $k \sim 10^7$ Mpc $^{-1}$

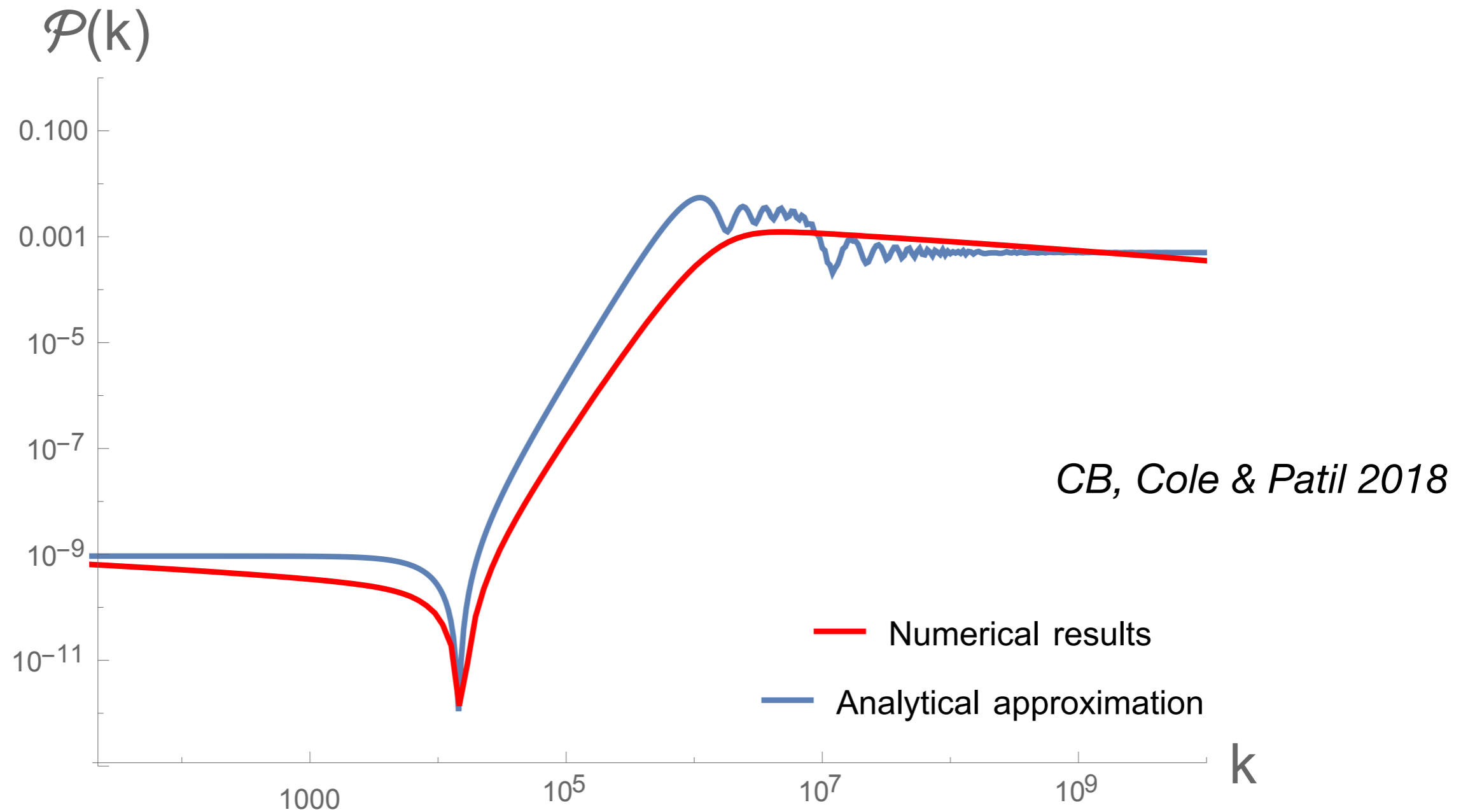


Early matter domination dramatically tightens constraints



Cole & CB '17

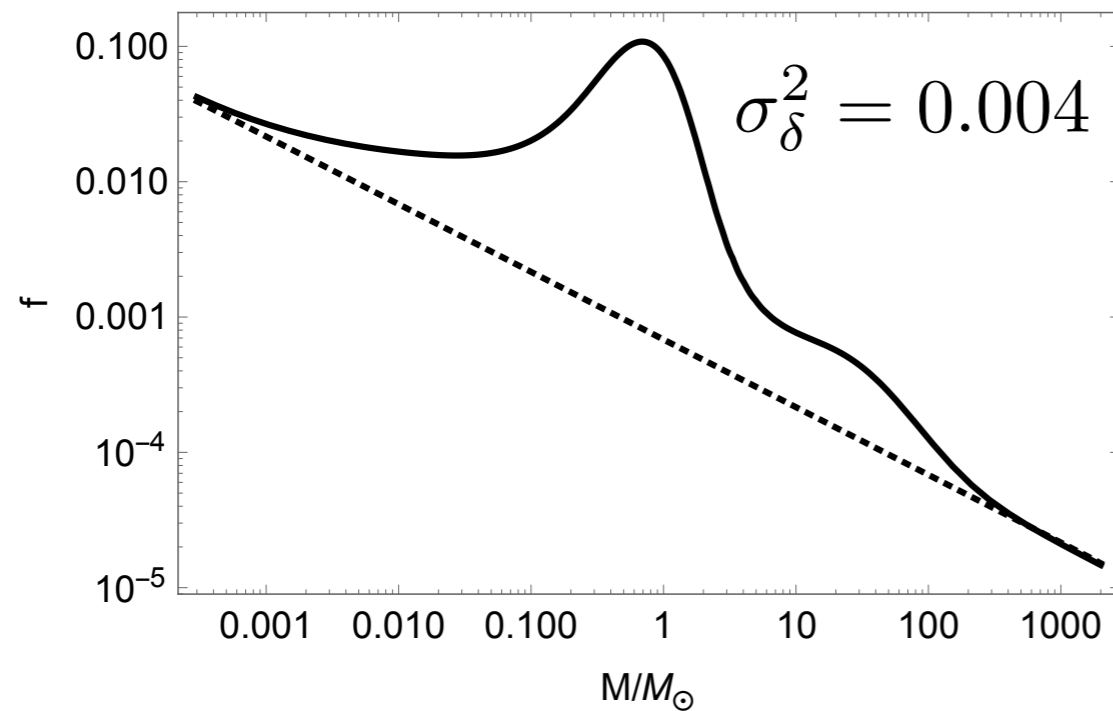
“Realistic” model with a smooth potential



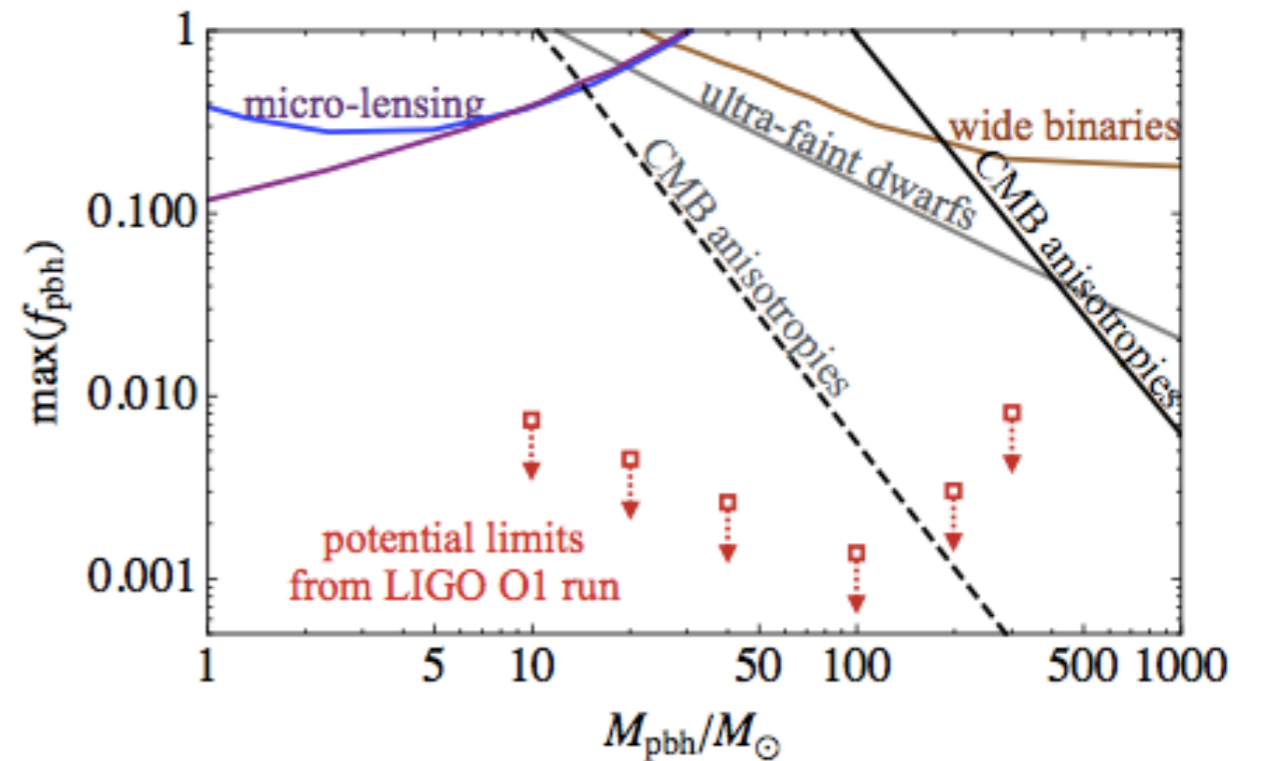
The red line is a full result from a smooth potential (*Germani & Prokopec '17*), while the blue line is based on a piecewise analytic calculation. Calculated using CPPTransport created by *David Seery 2016*

The resultant PBH mass function

$$f(M) = \frac{1}{\Omega_{\text{CDM}}} \frac{d\Omega_{\text{PBH}}}{d \ln M_{\text{H}}}$$



CB, Hindmarsh, Young & Hawkins 2018



Haimoud et al 2017

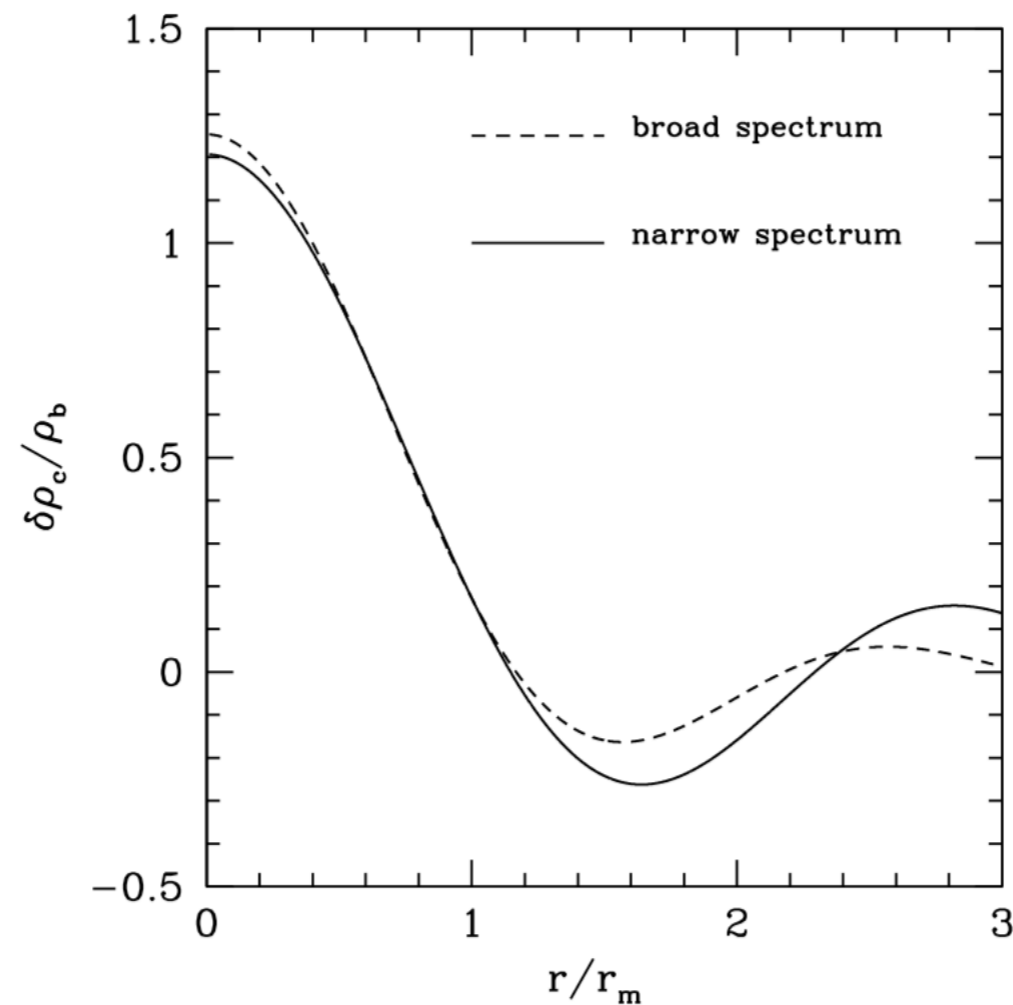
For the left plot, approx 10% of DM is made up of \sim solar mass PBHs and 0.1% lies in the LIGO mass range - enough to get the merger rate LIGO detects
Sasaki et al + Haimoud et al + Chen & Huang + Raidal et al + many more

The density profiles are related to the power spectrum shape

Due to the 2-point correlation function which tells you the density near peaks and shows that spherical symmetry is a good approximation

(*BBKS 1986* classic paper, non-spherical effects *Kühnel & Sandstad 2016*)

The density profile does not change strongly assuming a smooth peak in the primordial power spectrum, independently of the width



Germani & Musco 2018

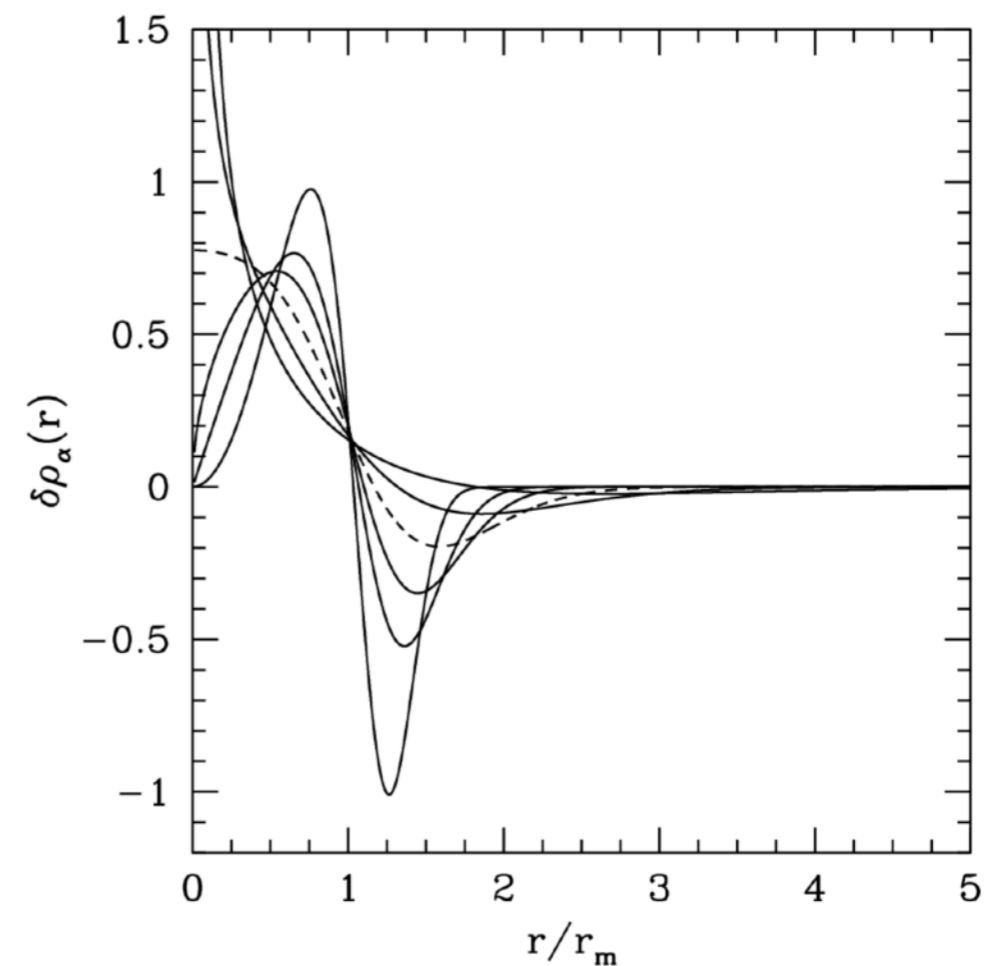
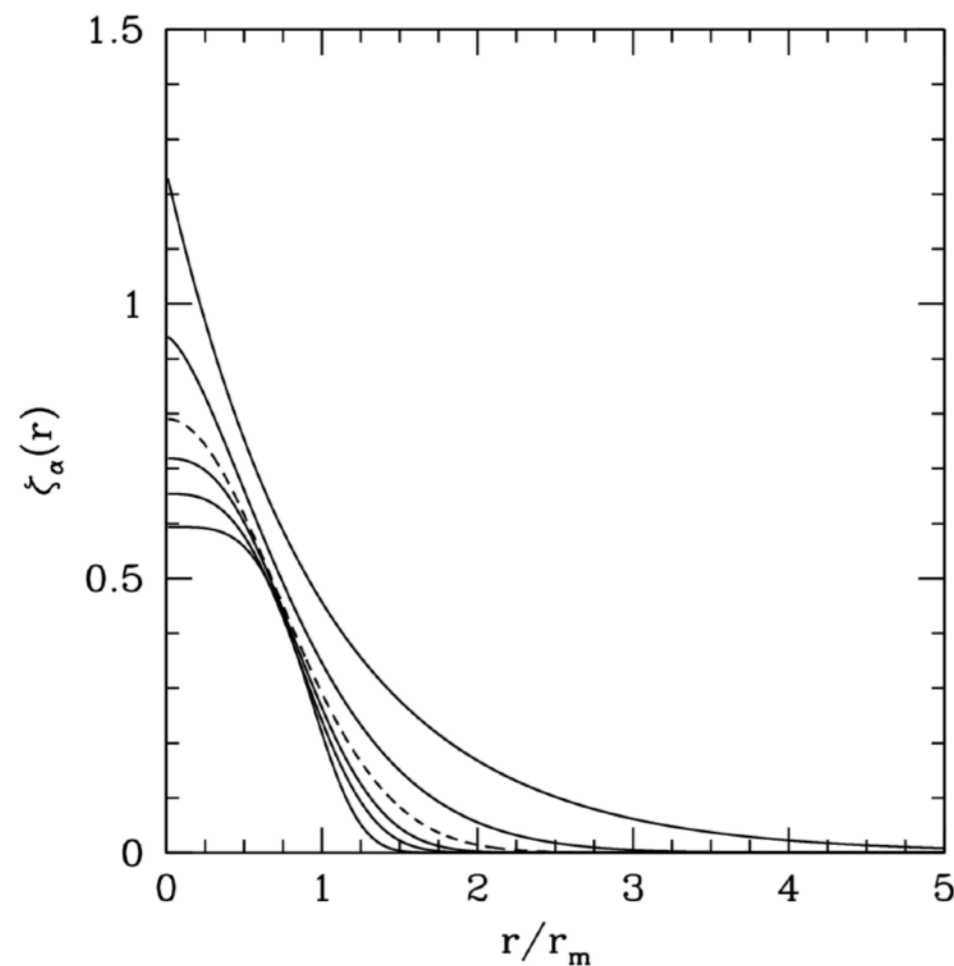
The density profiles are related to the power spectrum shape

Curvature perturbation profiles:

$$\zeta_\alpha(r) = \mathcal{A} \exp \left[- \left(\frac{r}{r_m} \right)^{2\alpha} \right]$$

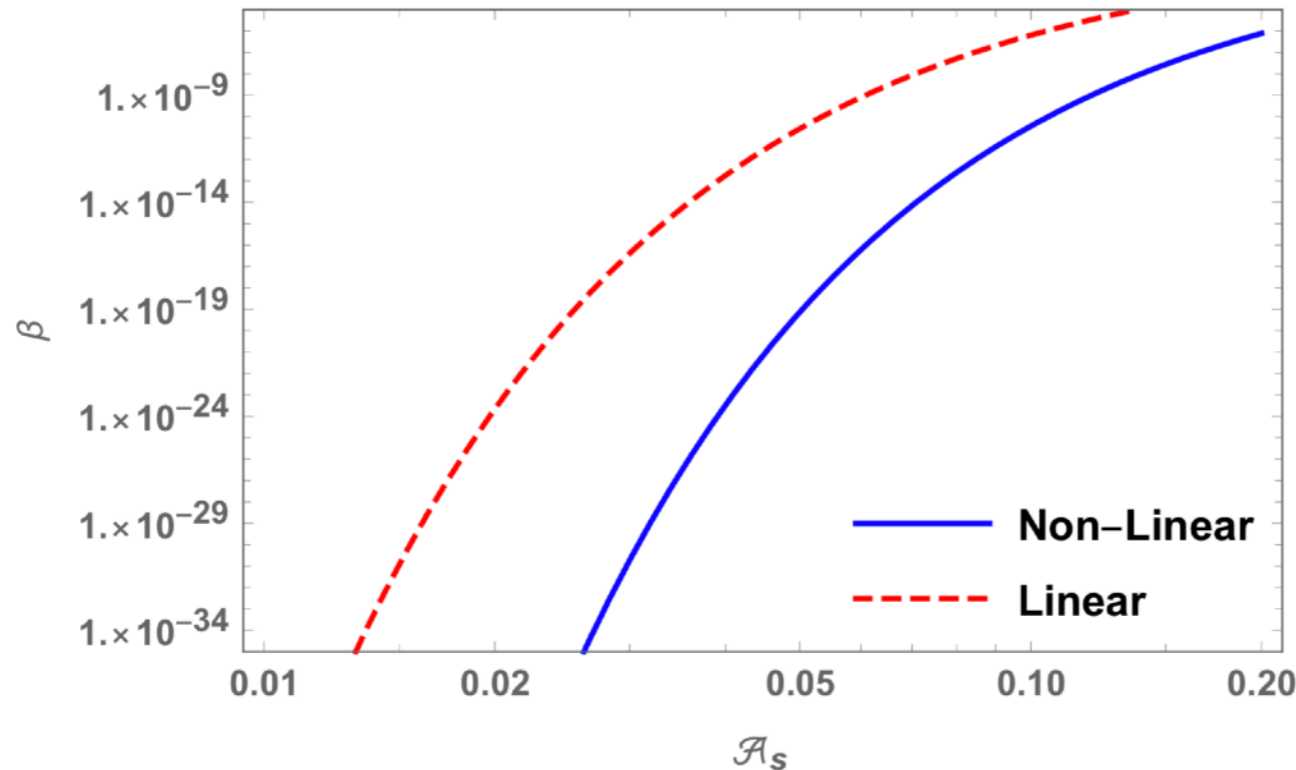
r_m is the scale at which the compaction function is maximised $\mathcal{C}(r, t) \equiv 2 \frac{M(r, t) - M_b(r, t)}{R(r, t)}$

Profile dependence: *Musco 2018* (see also *Nakama, Harada, Polnarev & Yokoyama 2014*)



Young, Musco, CB 2019

Power spectrum constraints are weakened by a factor ~ 2



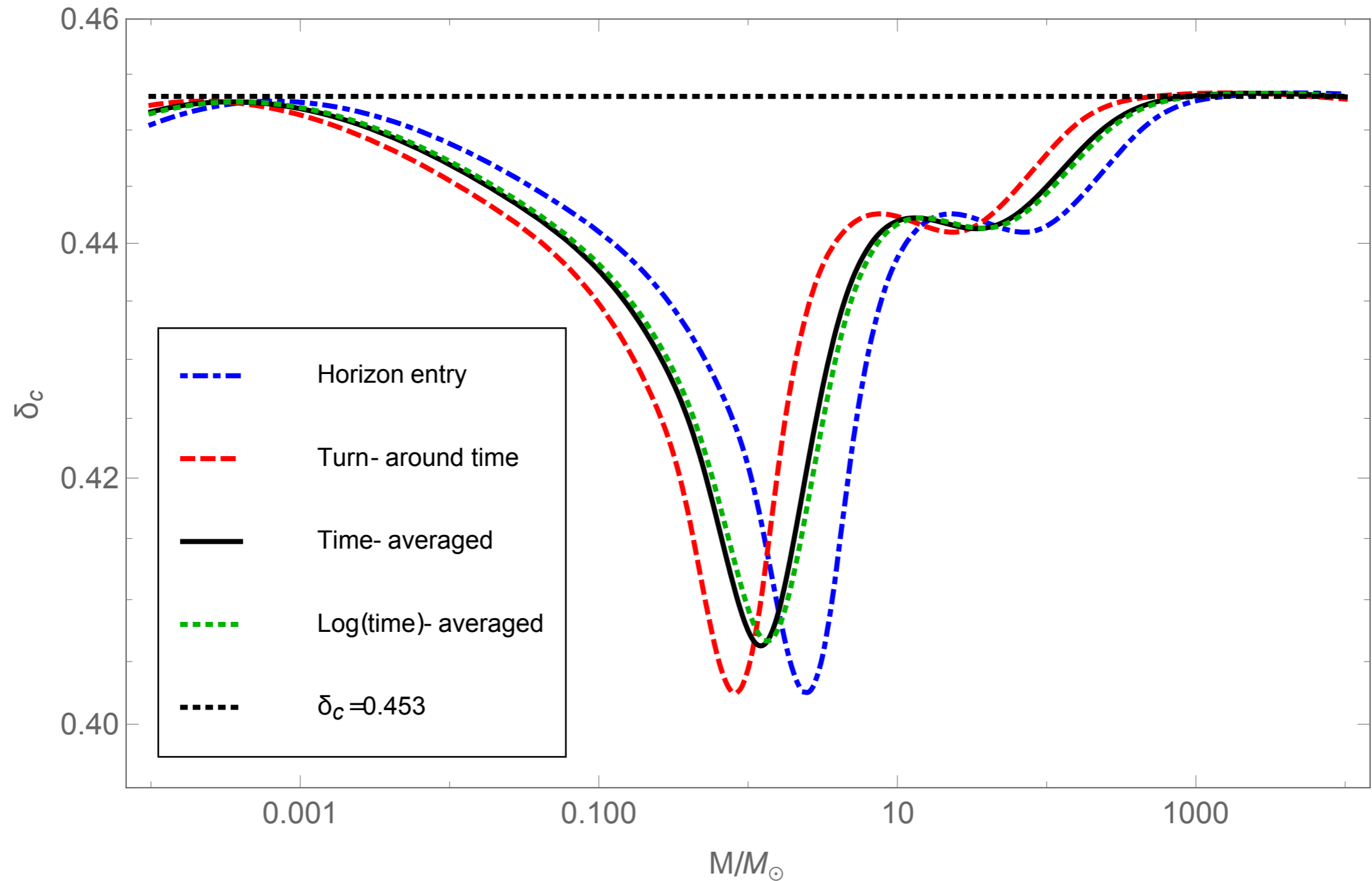
Delta function power spectrum
Young, Musco, CB '19

In order to generate the same number of PBHs when taking the non-linear (NL) relation into account, compared to the normal/wrong case that you use the linear relation, the power spectrum amplitude needs to increase by the ratio

$$1.5 \lesssim \frac{\mathcal{A}_{NL}}{\mathcal{A}_L} = \frac{16 \left(1 - \sqrt{\frac{2 - 3\delta_c}{2}} \right)^2}{9\delta_c^2} \lesssim 4$$

For the typical value of $\delta_c \sim 0.55$, power spectrum constraints are weakened by a factor of 2

Collapse threshold vs horizon mass



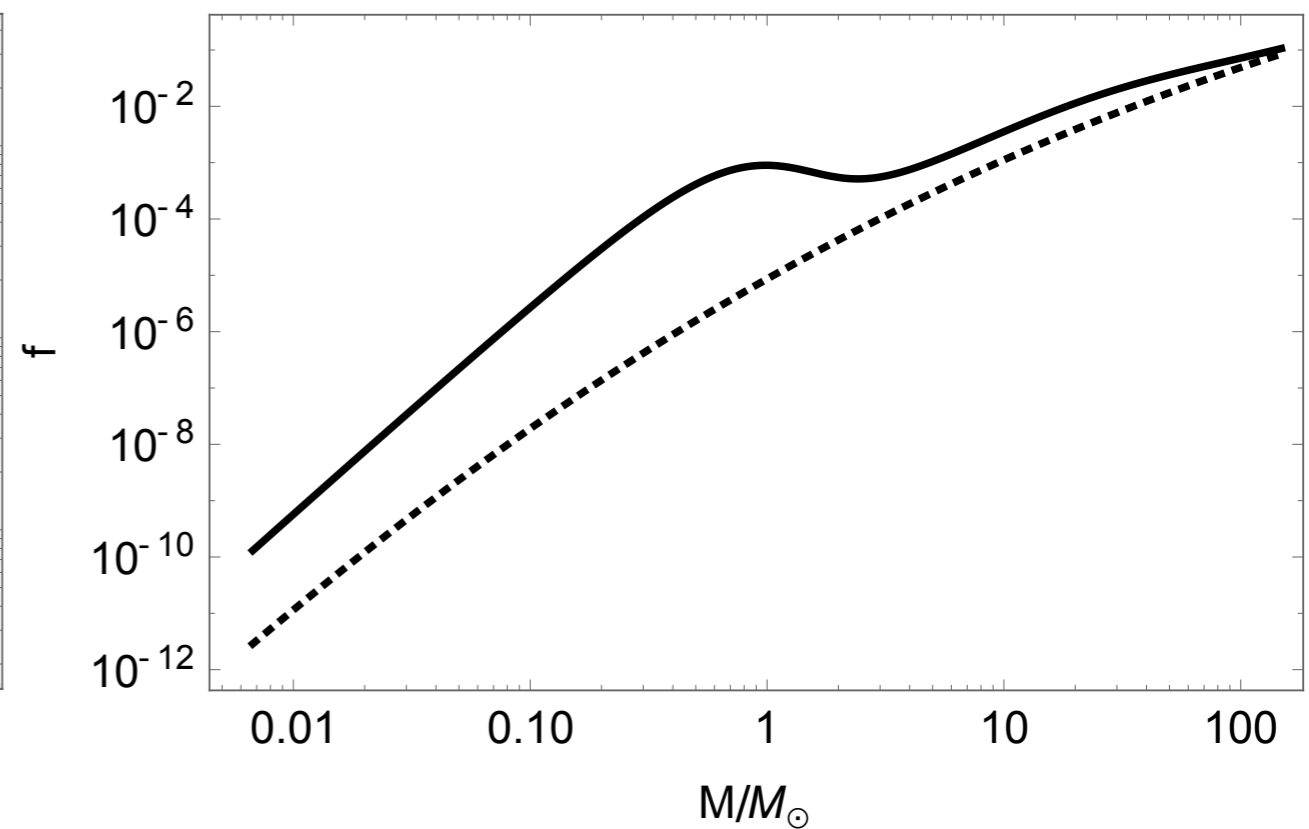
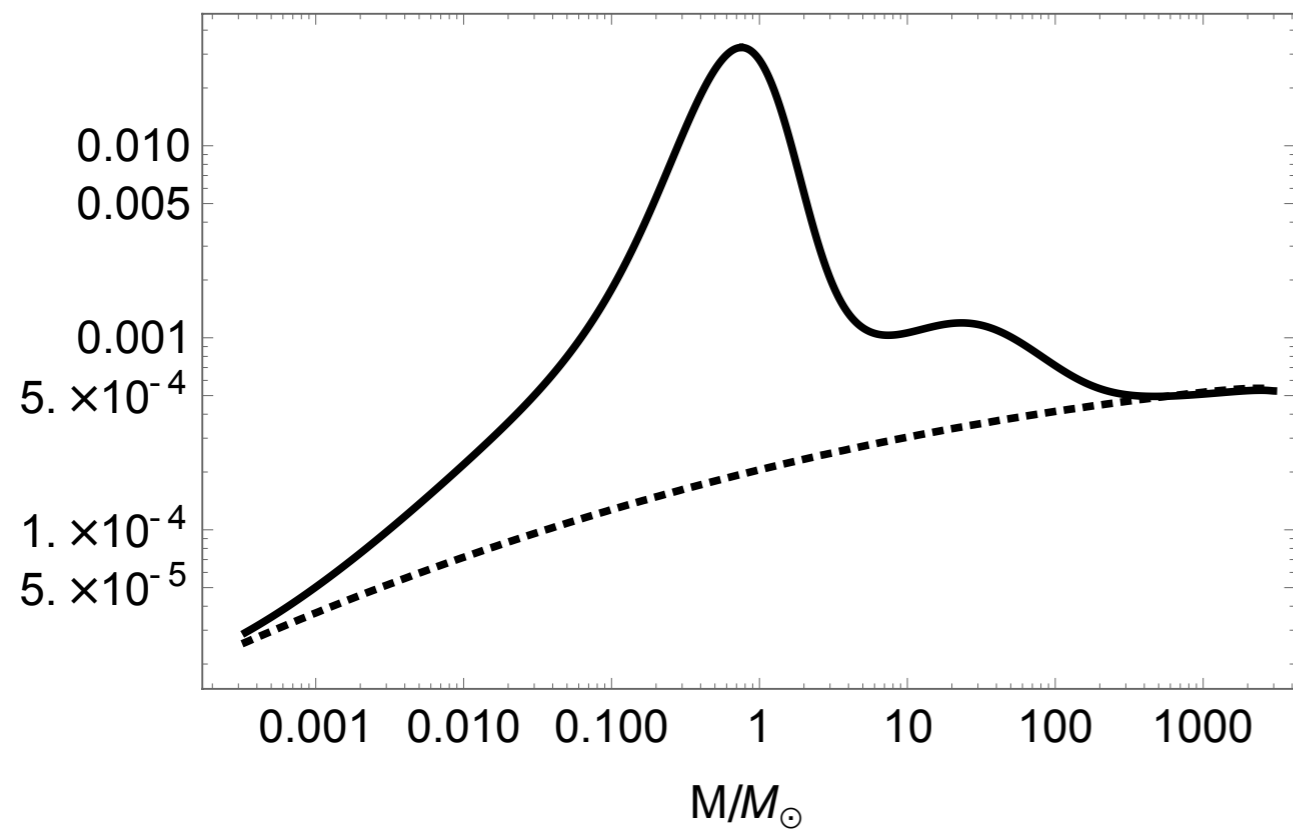
CB, Hindmarsh, Young & Hawkins 2018

Varying the primordial perturbations

If the primordial power spectrum is not scale invariant on the relevant scales then the mass function changes, but a peak remains

$$n_s - 1 = -0.05$$

$$n_s - 1 = -0.2$$



The PBH merger rate

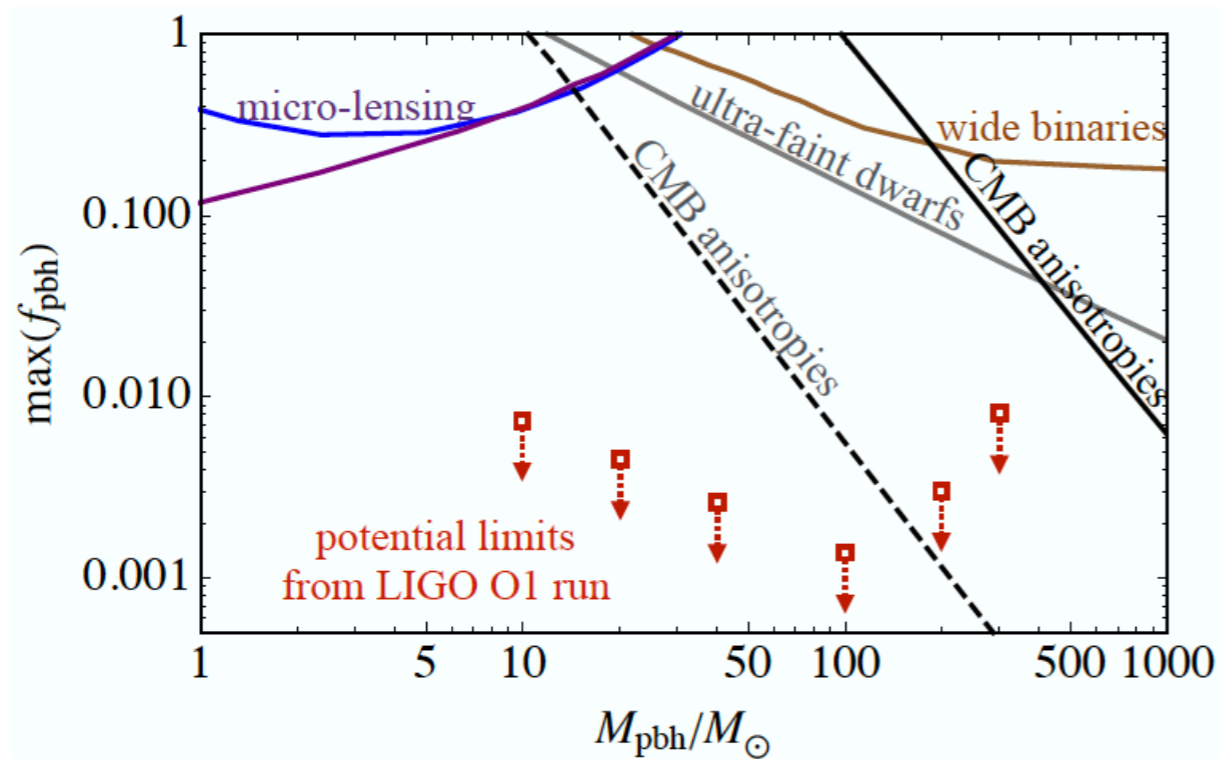


FIG. 7. Potential upper bounds on the fraction of dark matter in PBHs as a function of their mass, derived in this paper (red arrows), and assuming a narrow PBH mass function. These bounds need to be confirmed by numerical simulations. For comparison we also show the microlensing limits from the EROS [21] (purple) and MACHO [20] (blue) collaborations (see Ref. [74] for caveats and Ref. [32] for a discussion of uncertainties), limits from wide Galactic binaries [22], ultra-faint dwarf galaxies [25], and CMB anisotropies [24].

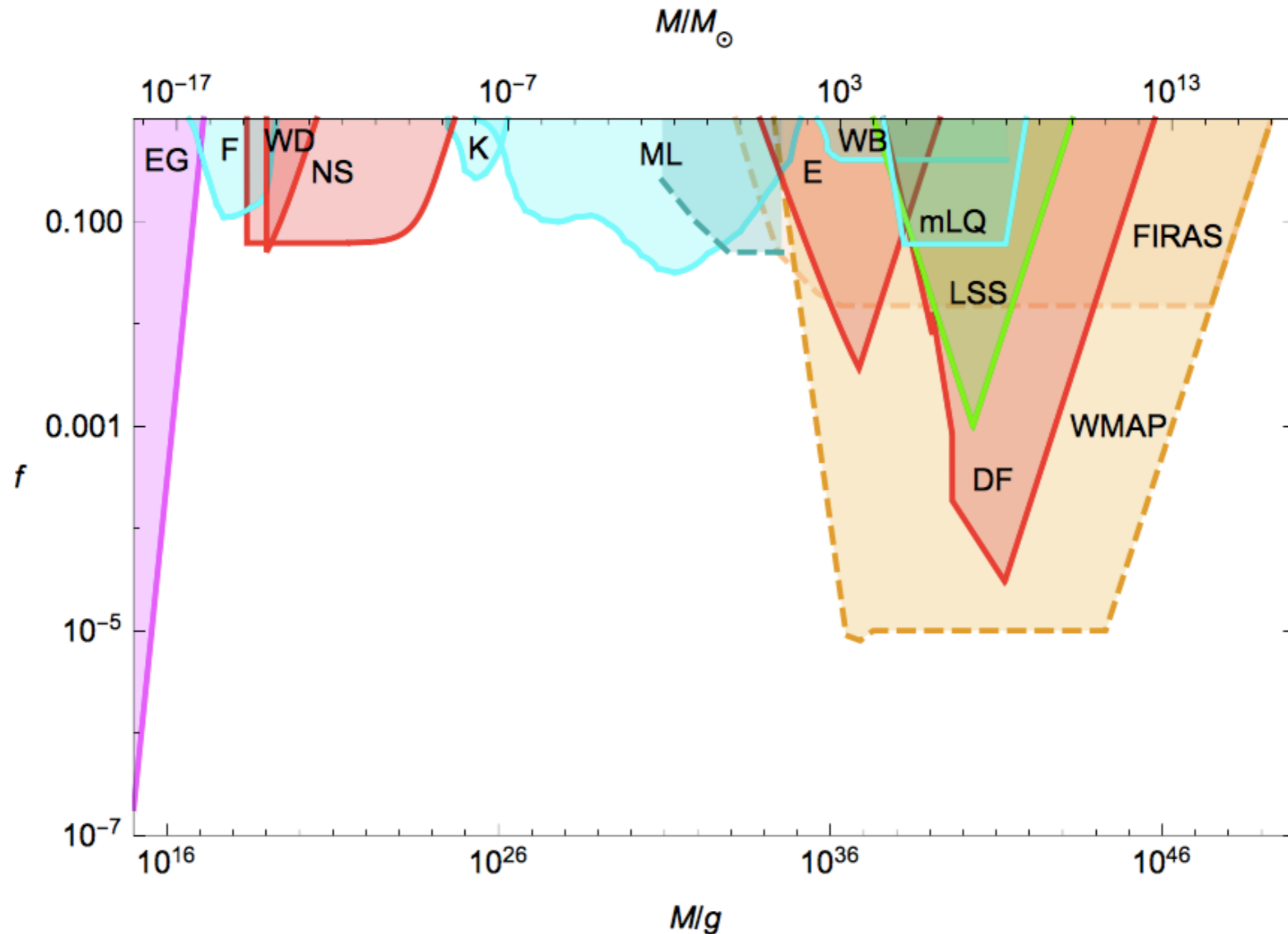
- Caveats:

1. Assumes a monochromatic mass spectrum
2. Assumes PBHs are randomly placed initially
3. Assumes BH binaries are not disrupted
4. Neglects accretion around the BHs

Haimoud et al 2017

First cosmological evolution simulations
Raidal, Spethmann, Vaskonen & Veermäe '18

Allowed fraction of PBHs in DM



Carr, Kuhnel & Sandstad 2017

See also Barack et al 2018; Black holes, gravitational waves and fundamental physics: a roadmap

(Disputed) observational constraints

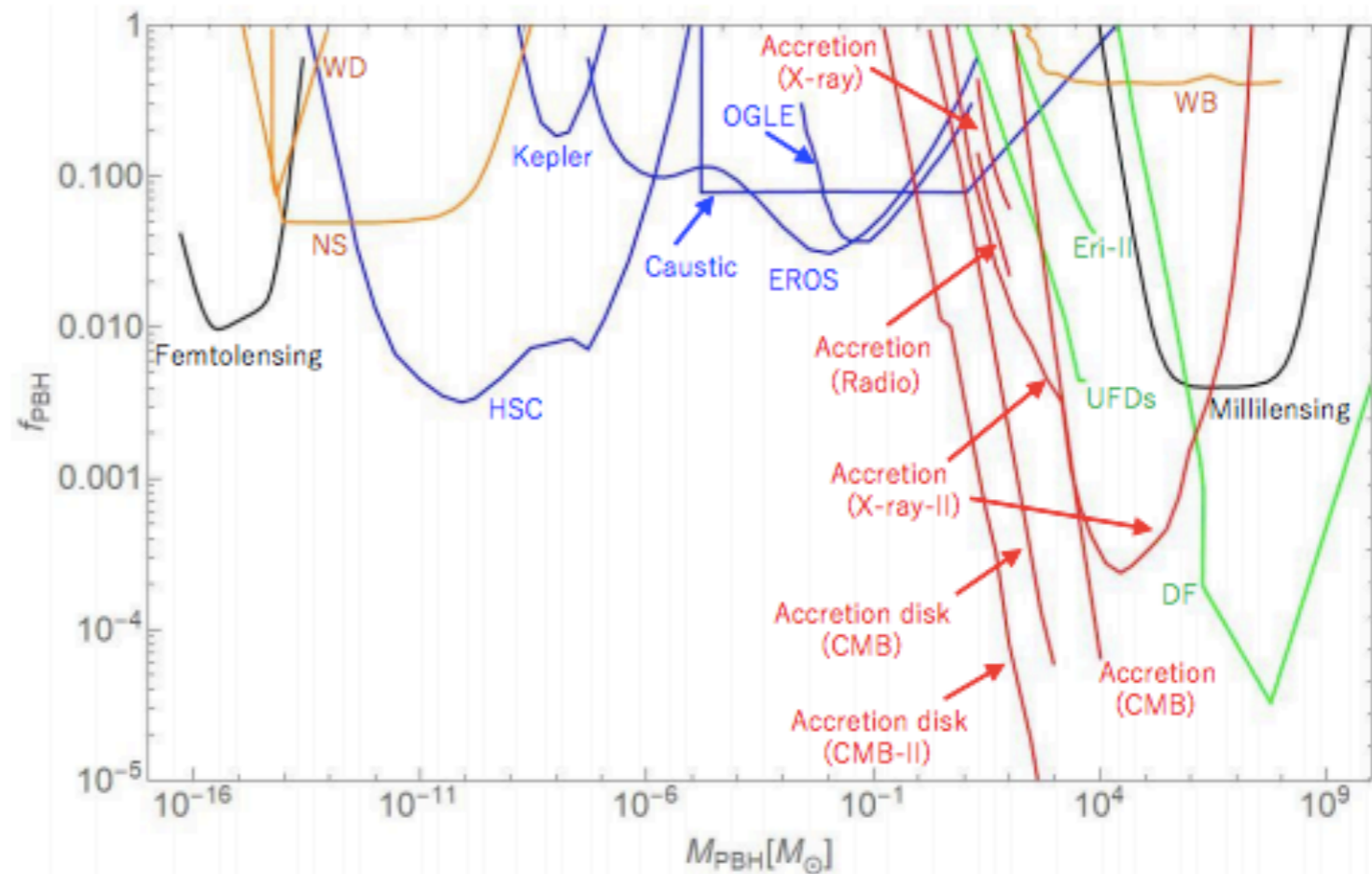
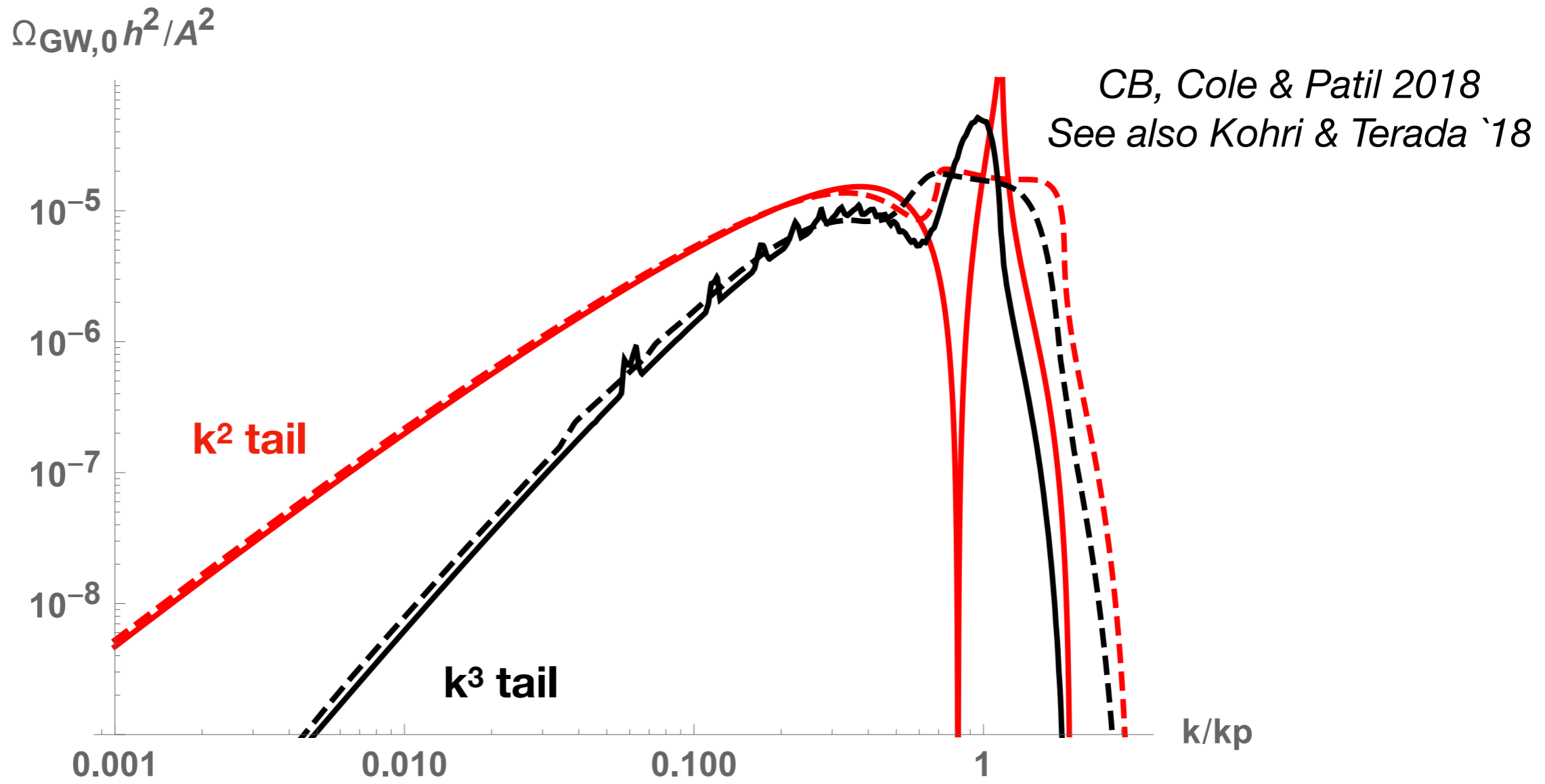


Figure 11: Upper limit on $f_{\text{PBH}} = \Omega_{\text{PBH}}/\Omega_{\text{DM}}$ for various PBH mass (assuming monochromatic mass function). Blue curves represent lensing constraints by EROS [116], OGLE [119], Kepler [122], HSC [123] and Caustic [125] (see 3.1.1). Black curves represent constraints by the millilensing [132] (3.1.2) and the femtolensing [138] (3.1.3). Orange curves represent dynamical constraints obtained by requiring that existent compact objects such as white dwarfs (WDs) [141] (3.2.1) and neutron stars (NSs) [142] (3.2.2) as well as the wide binaries (WBs) [151] (3.2.3) are not disrupted by PBHs. Green curves represent constraints by the dynamical friction (DF) on PBHs [152] (3.2.6), the ultra-faint dwarfs (UFDs) [153], and Eridanus II [153] (3.2.5). Red curves represent constraints by the accretion onto the PBHs such as CMB for the case of the spherical accretion [166] and the case of the accretion disk [171] with two opposite situations where the sound speed of the baryonic matter is greater (labeled by CMB) or smaller (labeled by CMB-II) than the relative baryon-dark matter velocity (3.3.1), radio, and X-rays [173, 180] (3.3.2).

GW spectrum



Red line - delta function scalar power spectrum

Black line - k^4 scalar power spectrum with cut off

Surprisingly, the delta function scalar power spectrum has a broader GW spectrum
This is unphysical and a warning against using delta function power spectra