PULSAR TIMING PROBES OF SMALL SCALE STRUCTURE

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ARXIV: 1901.04490
WITH JEFF DROR, TANNER TRICKLE, KATHRYN ZUREK
TO APPEAR
WITH TANNER TRICKLE, KATHRYN ZUREK









This is Greenland





This is Greenland







This is Greenland

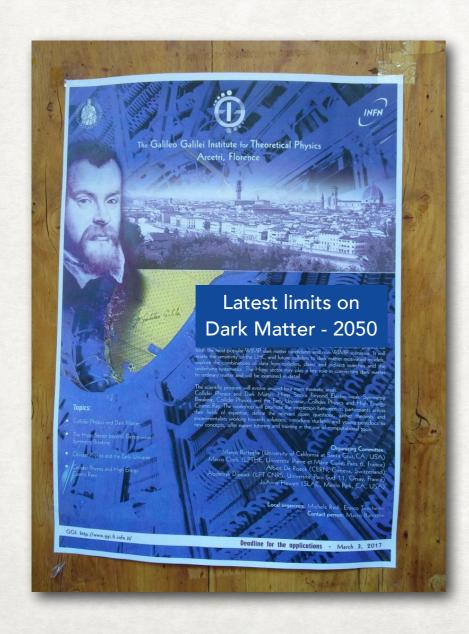


President Kanye West, Leader of the free world





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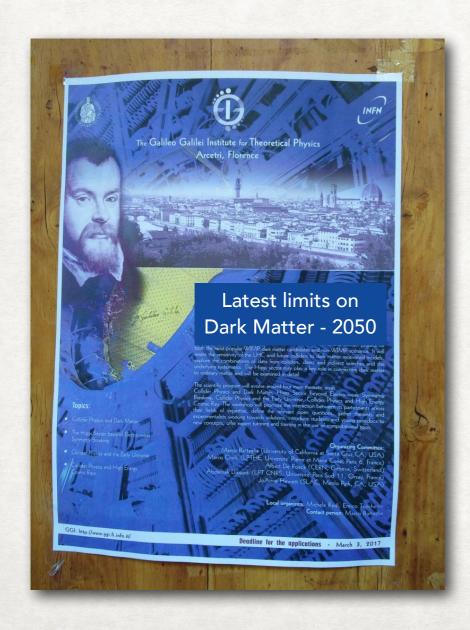


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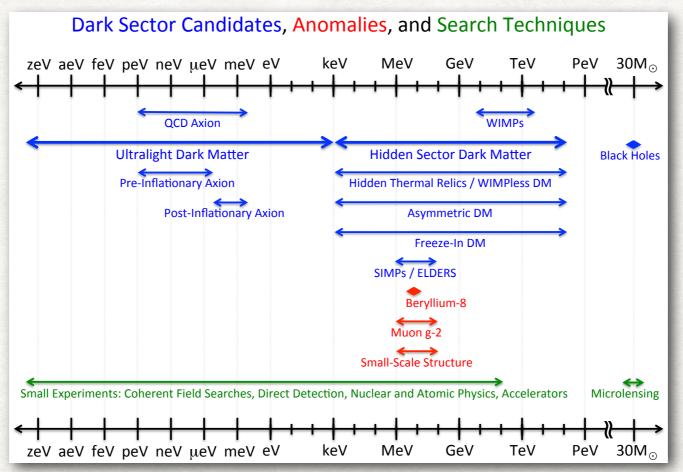
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This talk addresses one of these situations

WHAT DO WE KNOW ABOUT DARK MATTER?

- Ample Gravitational Evidence
- No confirmed positive signal in the lamp-post paradigm
- A bevy of promising experiments to probe interactions with SM and several more on the anvil



Cosmic visions

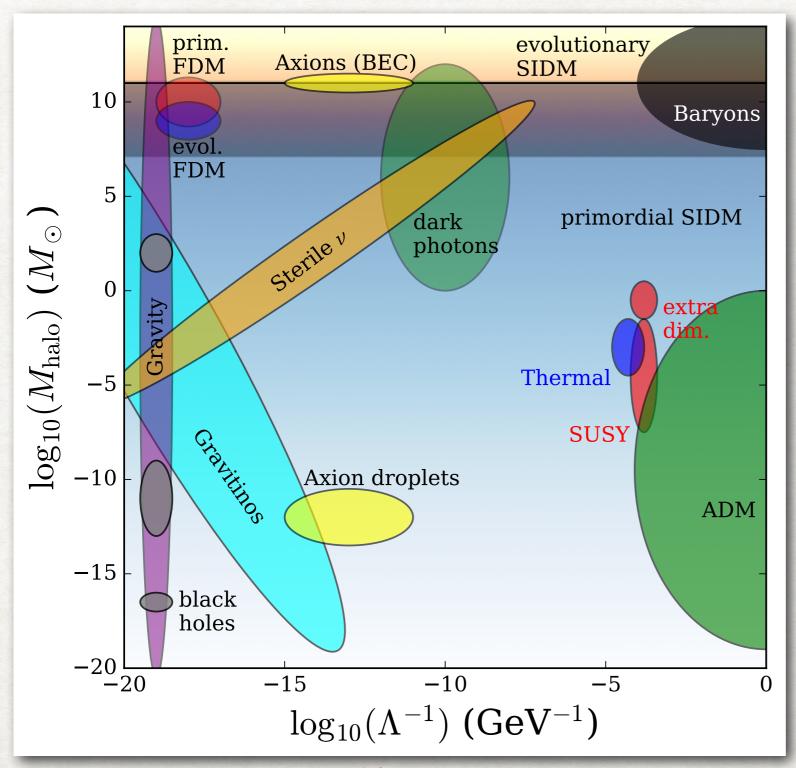
What about gravitational probes?

GRAVITATIONAL PROBES

Gravitational probes have already provided a wealth of information: sometimes even about particle nature

- Bullet Cluster self interactions
- Dwarf Galaxies elementary mass
- Super-radiance, other gravity probes of very light dark matter
- Clues from "small scale" challenges viz. core vs cusp, missing satellites etc.
- How about substructure at even small scales (intra-galactic)?

PROBING SUBSTRUCTURE - PROBING PARTICLE PHYSICS



Source: [Buckley, Peter:1712.06615]

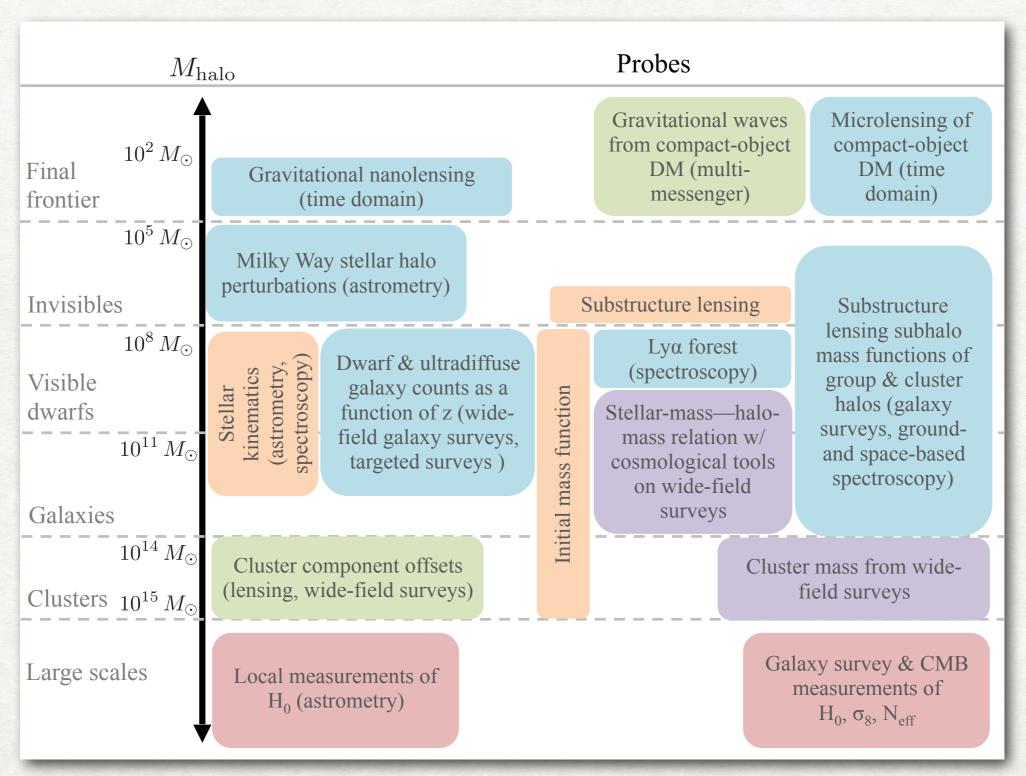
- Most models typically predict an extended halo mass function
- Minimum Halo mass and size set by
- Thermal History
- Collapse Redshift
- Early Matter Domination (See Nikita's talk)
- Inflationary vector production etc
- PQ breaking after inflation

predict enhanced structure in the small scales

SEVERAL UNKNOWNS

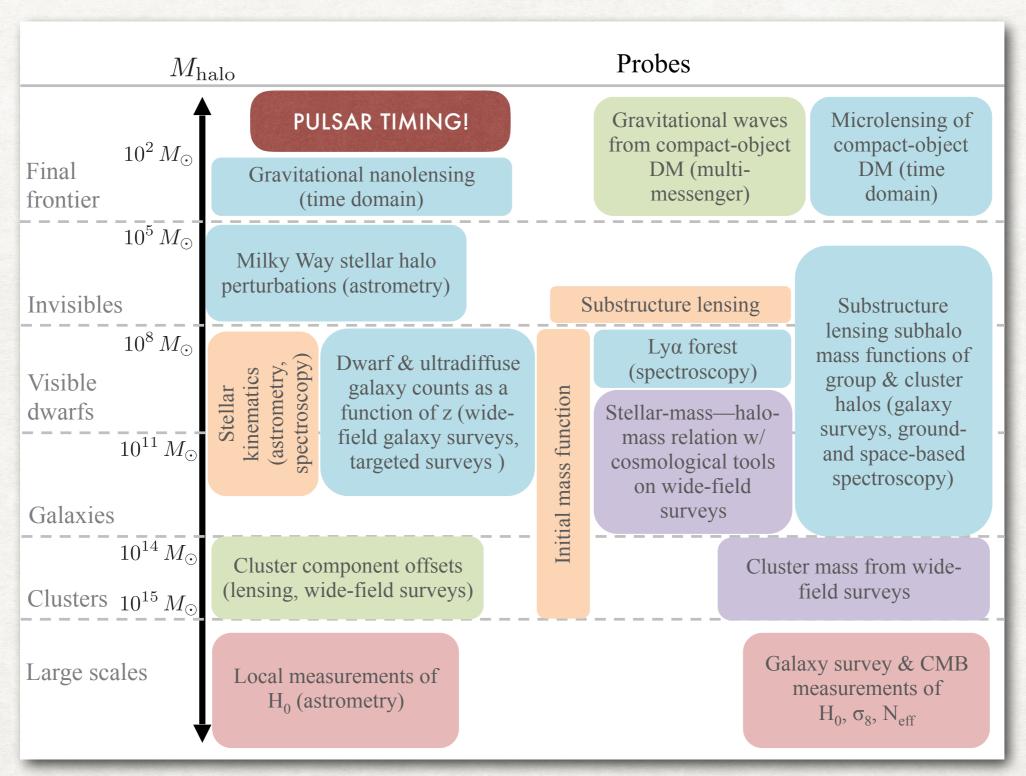
- Given an initial power spectrum, what is the substructure today?
- Well posed, hard to solve accurately
- Tidal stripping?
- How much of the DM is still in these miniclusters?
- Answers to these are important for direct detection too.

PROBES OF DIFFERENT MASSES



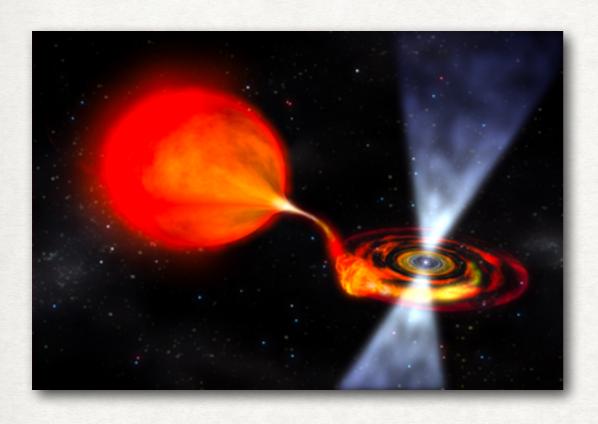
Source: [Buckley, Peter:1712.06615]

PROBES OF DIFFERENT MASSES



Source: [Buckley, Peter:1712.06615]

MILLI-SECOND PULSARS



- Neutron stars sped up through accretion.
- Fastest rotating pulsars have frequencies of a few kHz.
- Stable over remarkable timescales (T>20 years)
- Accurate timing models exist

PULSAR TIMING

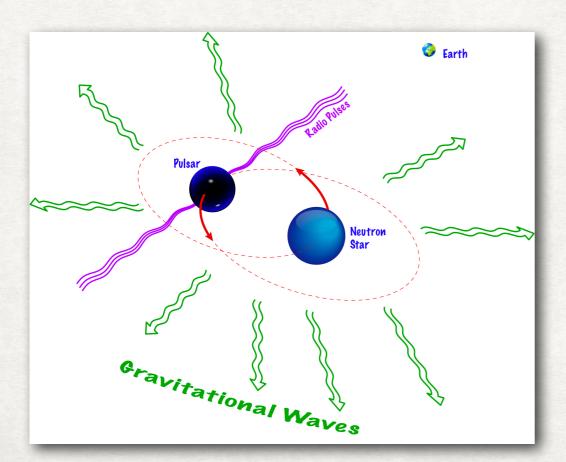
- Millisecond pulses ~ rotation period
- Phase: $\phi(t) = \phi_0 + \nu t + \frac{1}{2} \dot{\nu} t^2 + \frac{1}{6} \ddot{\nu} t^3 + \dots$
- ν ~ kHz
- $\dot{\nu}/\nu$ ~ 10-23 to 10-20 Hz
- $\ddot{\nu}/\nu$ < 10⁻³¹ Hz², not included in fits
- After fitting away the period and derivative, residuals are remarkably small* (and stable).

$$t_{
m RMS} \equiv \sqrt{rac{1}{N} \sum_n (t_n^{
m data} - t_n^{
m fit})^2}$$
 ~ 50 ns

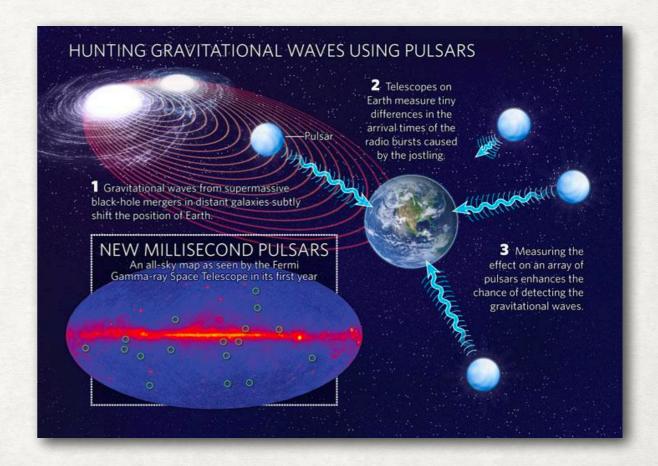
NEW PHYSICS FROM PULSAR TIMING

• Any new physics that predicts time dependent $\,\delta\phi\equiv\int dt\,\delta\nu(t)\,$ can possibly be constrained.

Hulse-Taylor binary used as a test of GR through its contracting orbit



Can be used as an extremely low frequency GW detector



PTA COLLABORATIONS



Today

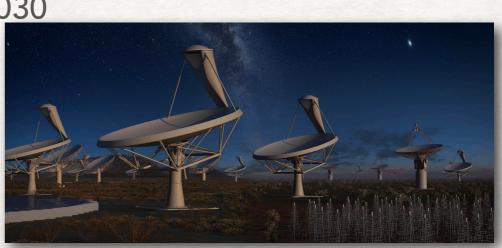
- Np ~ 73
- T = 10 to 20 years



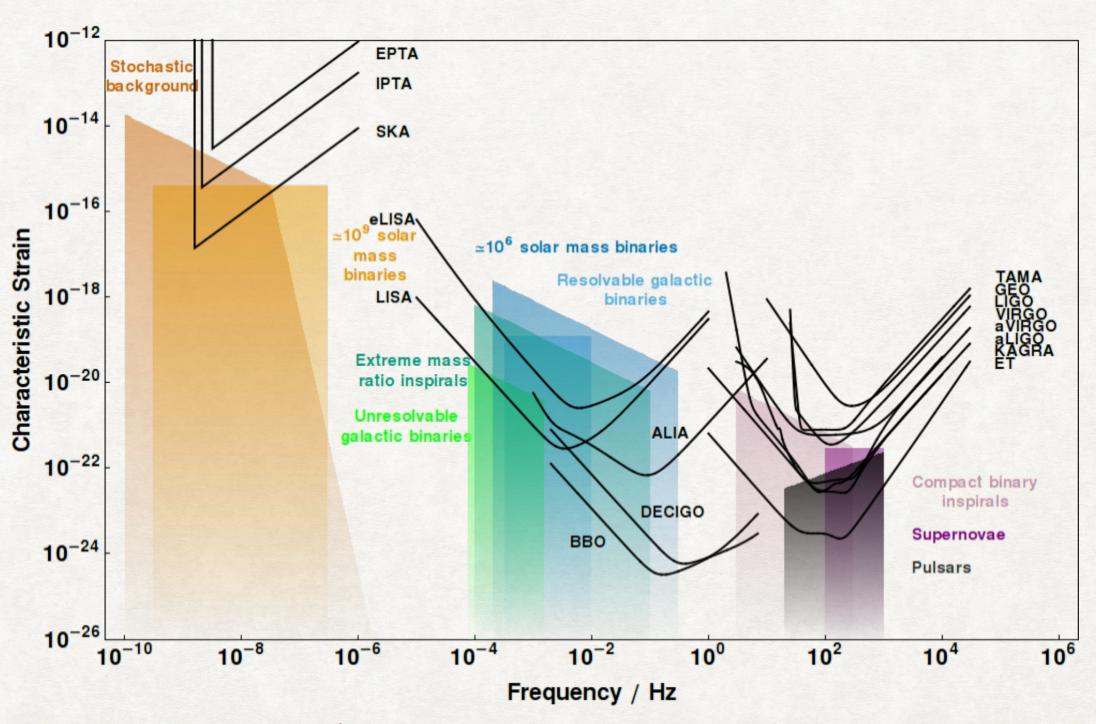
Future

- Several precursors currently running
- Np ~ 200-1000
- Projected to start ~ 2030
- T = 20 + years





PTA VS OTHER GRAVITY WAVE DETECTORS



Moore, Cole, Berry

PTA FOR DARK MATTER

- Gravitational probes are broadly of two varieties
- Either probe gravitational interaction of DM with some test mass.
 e.g. Carney, Ghosh, Krnjaic, Taylor. arXiv:1903.00492
- Or probe gravitational interaction between light and DM, e.g. Lensing/ LIGO etc
- PTAs have both kinds of signal (see also 1804.01991 van Tilburg, Taki, Weiner for larger masses with astrometry instead)

PTA FOR DARK MATTER

- Ultralight DM causing GW like delays
- [Khmelnitsky, Rubakov 1309.5888], [Graham, Kaplan, Mardon, Rajendran, Terrano 1512.06165]
- PTAs are sensitive accelerometers: Doppler Delay
- [Seto, Corray astro-ph/0702586], [Baghram, Afshordi, Zurek 1101.5487] [Kashiyama, Seto 1208.4101], [Kazumi, Oguri, Masamune 1801.07847]
- · Gravitational potential wells along the light path: Shapiro Delay
- [Siegel, 0801.3458], [Siegel, Hertzberk, Fry astro-ph/0702546], [Baghram, Afshordi, Zurek 1101.5487],
 [Clark, Lewis, Scott 1509.02938], [Schutz, Liu 1610.04234]
- This Work:
- Explicit calculations for SNR calculations for Doppler and Shapiro Delays
- Comprehensive analysis of both Doppler/Shapiro for all signal durations
- Extending analysis to diffuse halos

DOPPLER DELAY

- Recognize the ratio $\frac{\delta \nu}{\nu}$ is $v_{\rm rel}/c$
- Thus sensitive to tiny accelerations

$$\left(\frac{\delta\nu}{\nu}\right)_D = \hat{\mathbf{d}} \cdot \int \nabla\Phi \ dt,$$

 velocity shape for a point object transit looks like:

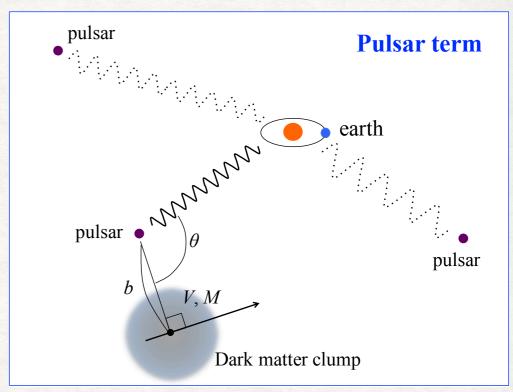
$$\left(rac{\delta
u}{
u}
ight)_D = rac{GM}{v^2 au_D} rac{1}{\sqrt{1+x_D^2}} \, \left(x_D \hat{\mathbf{b}} - \hat{\mathbf{v}}
ight) \cdot \hat{\mathbf{d}}$$
.

Impact parameter

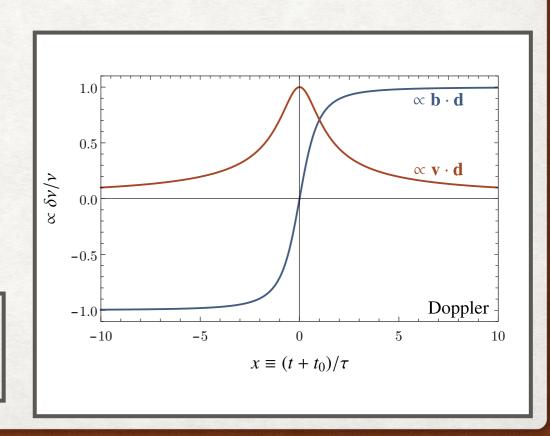
Signal period

Dimensionless time variable

 $\left|\mathbf{b}\right| = au v$



Source: Kashiyama, Seto - 1208.4101



SHAPIRO DELAY

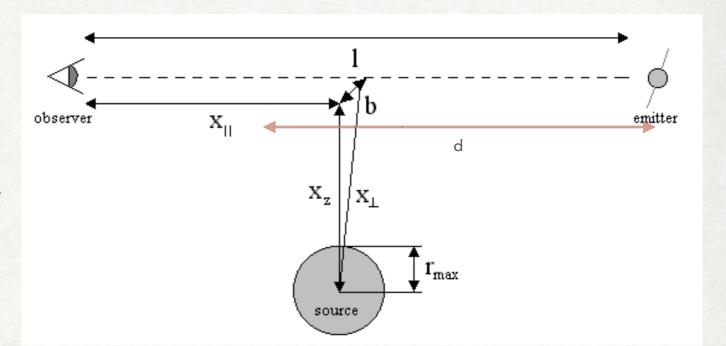
- Similar to Sachs-Wolfe effect
- · In frequency domain given by,

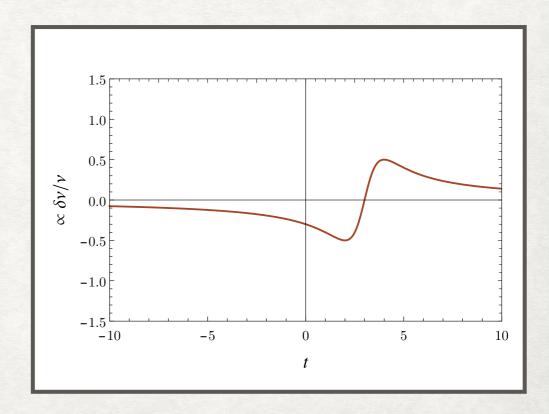
Dimensionless time variable

$$\left(\frac{\delta\nu}{\nu}\right)_S = -2\int \mathbf{v} \cdot \nabla\Phi \ dz$$

· For a point object,

$$\left(\frac{\delta\nu}{\nu}\right)_S = \frac{4GM}{\tau_S} \frac{x_S}{1+x_S^2}$$
 Duration of signal

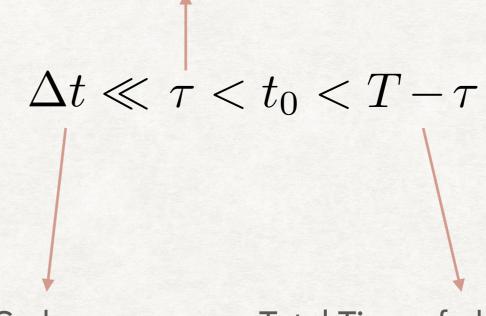




DYNAMIC VS STATIC

Characteristic signal period

Dynamic if



Cadence

Total Time of observation

Static otherwise

$$\tau \gtrsim T$$

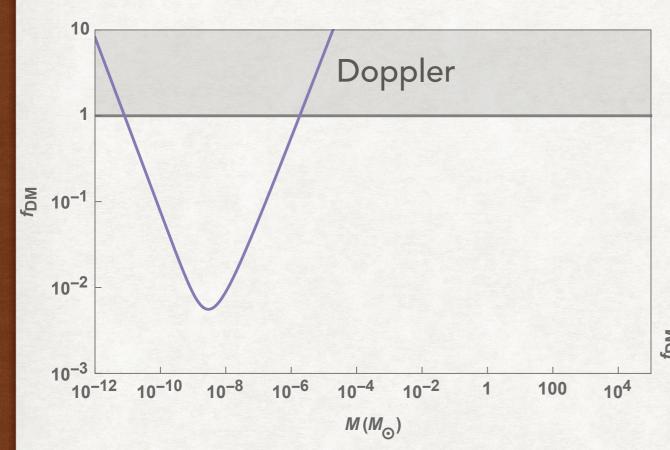
DETECTING SIGNALS

- Dynamic: Similar to a bump hunt / LIGO signal / Microlensing signal
- Single event shape and amplitude, straightforward filtering.
- Static: Works only if a second derivative continues to not be measured.

START WITH MONOCHROMATIC POINT MASSES

BOUNDS FROM DYNAMIC SIGNALS

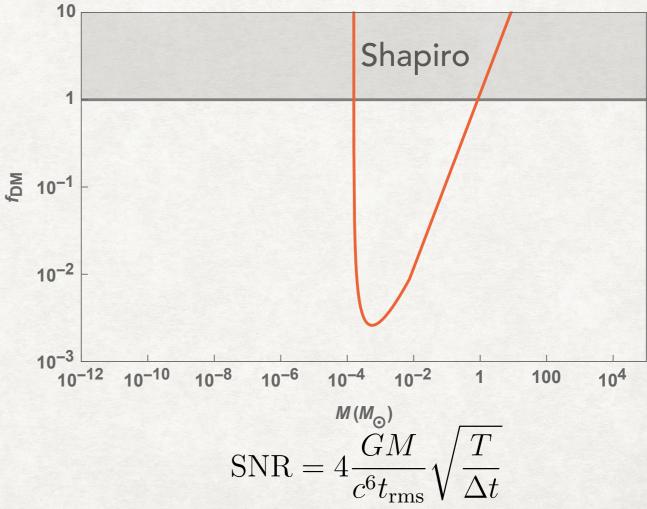
$$SNR_D = \frac{1}{2\sqrt{3}} \frac{GMT^{\frac{3}{2}}}{c t_{rms} v^2 \sqrt{\Delta t} \tau}$$



Small sampling volume localized around pulsar/earth

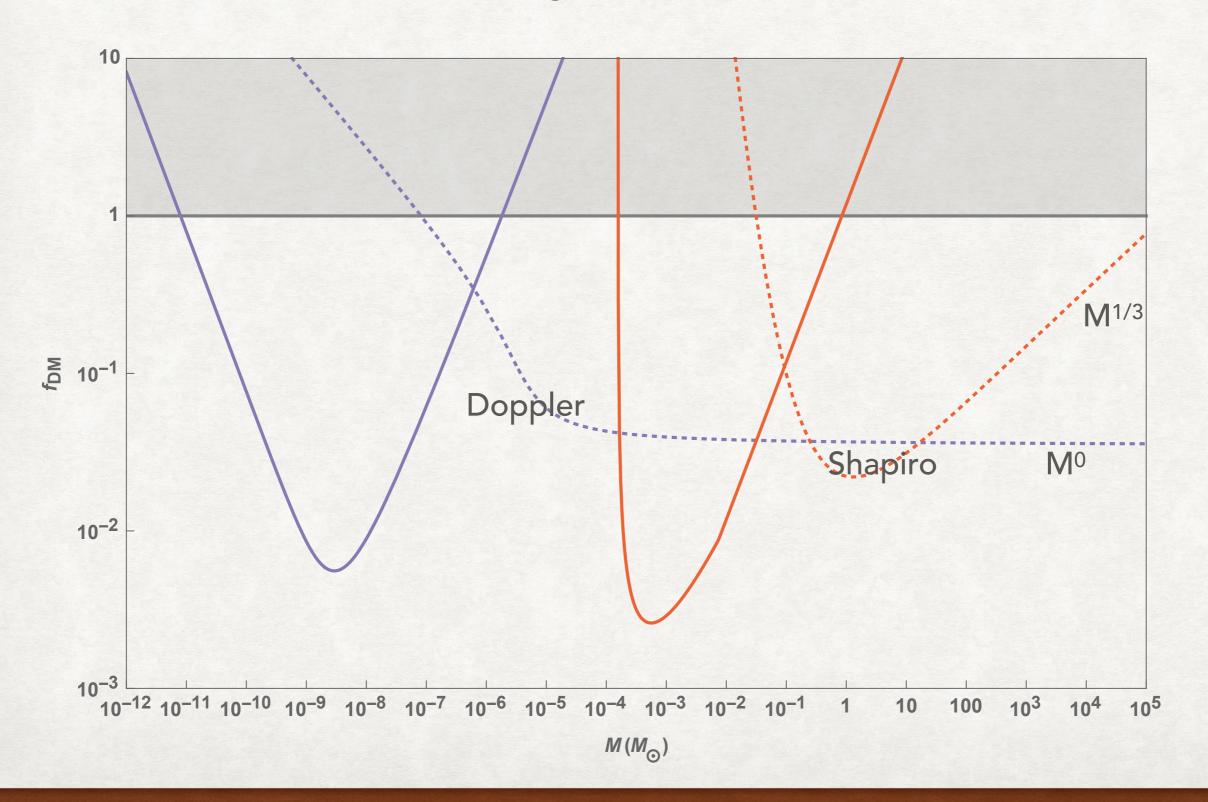
- Just like any direct detection experiment
- LHS: closest event below threshold
- RHS: closest event takes too long

Large sampling volume
Along the whole light path



FRACTION VS M SCALING -STATIC

Could also look for a static signal: i.e. a non-zero second derivative



THE COSMIC ENERGY INVENTORY

Parameter	Components ^a	Totals ^a
Dark sector:		0.954 ± 0.003
Dark energy	0.72 ± 0.03	
Dark matter	0.23 ± 0.03	
Primeval gravitational waves	$\lesssim 10^{-10}$	
Primeval thermal remnants:		0.0010 ± 0.0005
Electromagnetic radiation	$10^{-4.3\pm0.0}$	
Neutrinos	$10^{-2.9 \pm 0.1}$	
Prestellar nuclear binding energy	$-10^{-4.1 \pm 0.0}$	
Baryon rest mass:		0.045 ± 0.003
Warm intergalactic plasma	0.040 ± 0.003	3
Virialized regions of galaxies	0.024 ± 0.005	
Intergalactic	0.016 ± 0.005	
Intracluster plasma	0.0018 ± 0.000)7
Main-sequence stars: spheroids and bulges	0.0015 ± 0.000)4
Main-sequence stars: disks and irregulars	0.00055 ± 0.000)14
White dwarfs	0.00036 ± 0.000	800
Neutron stars	0.00005 ± 0.000	002
Black holes	0.00007 ± 0.000	002
Substellar objects	0.00014 ± 0.000	007
H 1 + He 1	0.00062 ± 0.000)10
Molecular gas	0.00016 ± 0.000	006
Planets	10^{-6}	
Condensed matter	$10^{-5.6 \pm 0.3}$	
Sequestered in massive black holes	$10^{-5.4}(1+\epsilon_n)$	

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Intergalactic	0.016 ± 0.005	
Intracluster plasma	0.0018 ± 0.00	007
Main-sequence stars: spheroids and bulges	0.0015 ± 0.00	004
Main-sequence stars: disks and irregulars	0.00055 ± 0.00	014
White dwarfs	0.00036 ± 0.00	0008
Neutron stars	0.00005 ± 0.00	0002
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Intergalactic	0.016 ± 0.005	C
Intracluster plasma	0.0018 ± 0.0	Static
Main-sequence stars: spheroids and bulges	0.0015 ± 0.0	0004
Main-sequence stars: disks and irregulars	$0.00055 \pm 0.$	00014
White dwarfs	0.00036 ± 0.0	0008
Neutron stars	0.00005 ± 0	00002
Black holes	0.00007 ± 0.0	00002
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Main-sequence stars: disks and irregulars		0.00055 ± 0.00014		
White dwarfs		0.00036 ± 0.0008		
Neutron stars		0.00005 ± 0.00002		
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THE COSMIC ENERGY INVENTORY

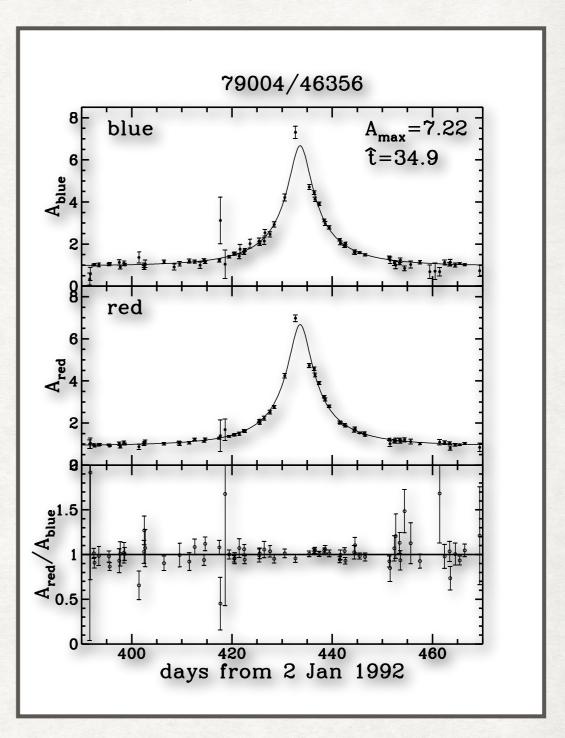
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OTHER SOURCES OF BACKGROUND

- Glitches: Sudden increase in frequency, followed by a slow relaxation (days-year).
- Dispersion through interstellar medium
- Pulsars may suffer from intrinsic red noise
- Beyond the scope of a collaboration-independent analysis
- Noise subtractions do not reduce signal when running it through a simulated PTA analysis. (We thank Stephen Taylor of the Nanograv collaboration for this).

DYNAMIC BACKGROUNDS

- Dynamic signal more spectacular than static signal.
- Shape differences could help differentiate from glitches etc.
- Dynamic signals are non-dispersive
- Will not be limited soon by baryonic background.



Dispersion used in Microlensing to differentiate lensing blip from a dispersive blip

PTA TIMELINE

	T [yr]	$t_{ m RMS} \ [m ns]$	$\Delta t \; [\mathrm{wk}]$	$d [\mathrm{kpc}]$	N_P
Current SKA Optimistic	5 - 30 20 20	$50 - 10^4$ 50 25	$egin{array}{c} 1-4\ 2\ 1 \end{array}$	0	

With this we consider 4 scenarios:

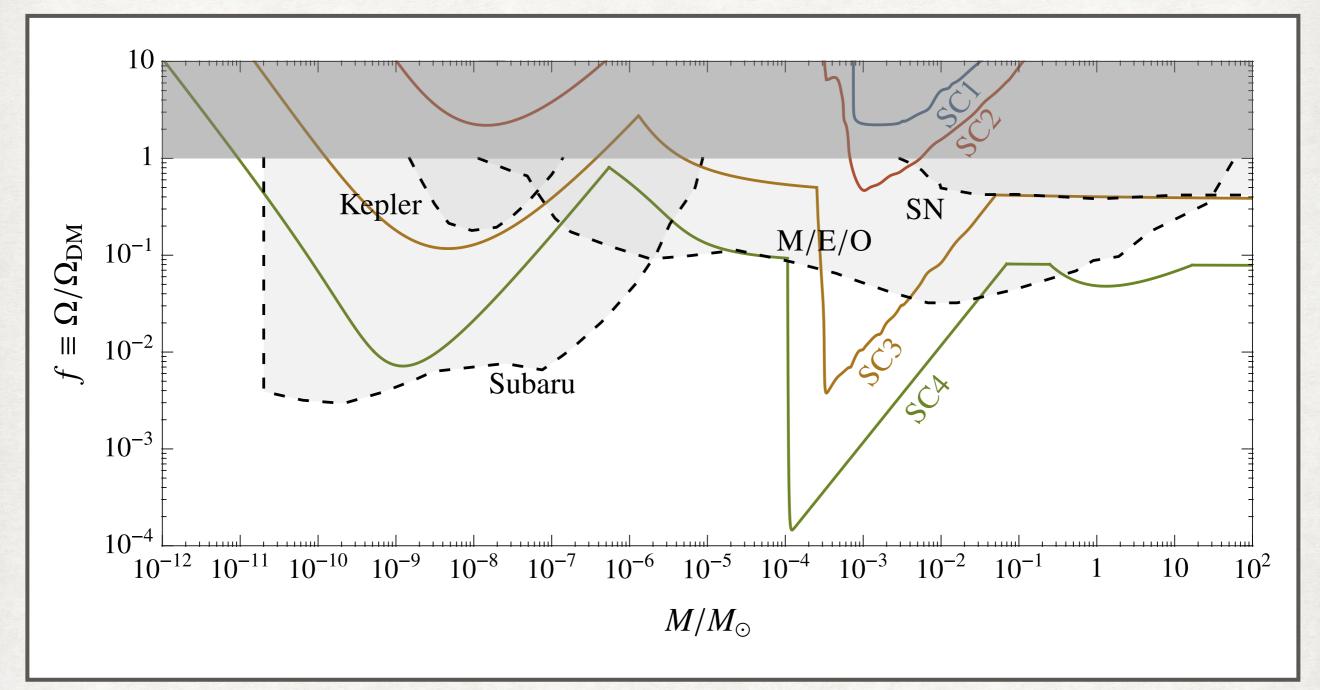
SC1: Estimated limits from current data only

SC2: Current pulsars running for 10 more years

SC3: Assume SKA starts 10 year from now and current pulsars continue

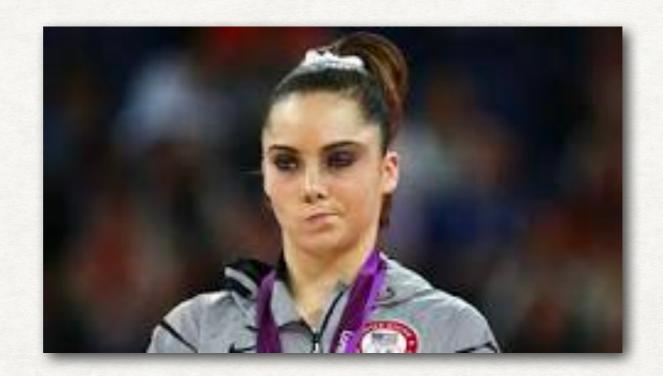
SC4: Optimistic

RESULTS FOR PBH



Lensing constraint from Subaru, Machos, Eros, Ogle (MEO) and SN lensing

IS THIS A SILVER MEDAL?



Limits comparable but subdominant to lensing for the most part

MORE DIFFUSE OBJECTS

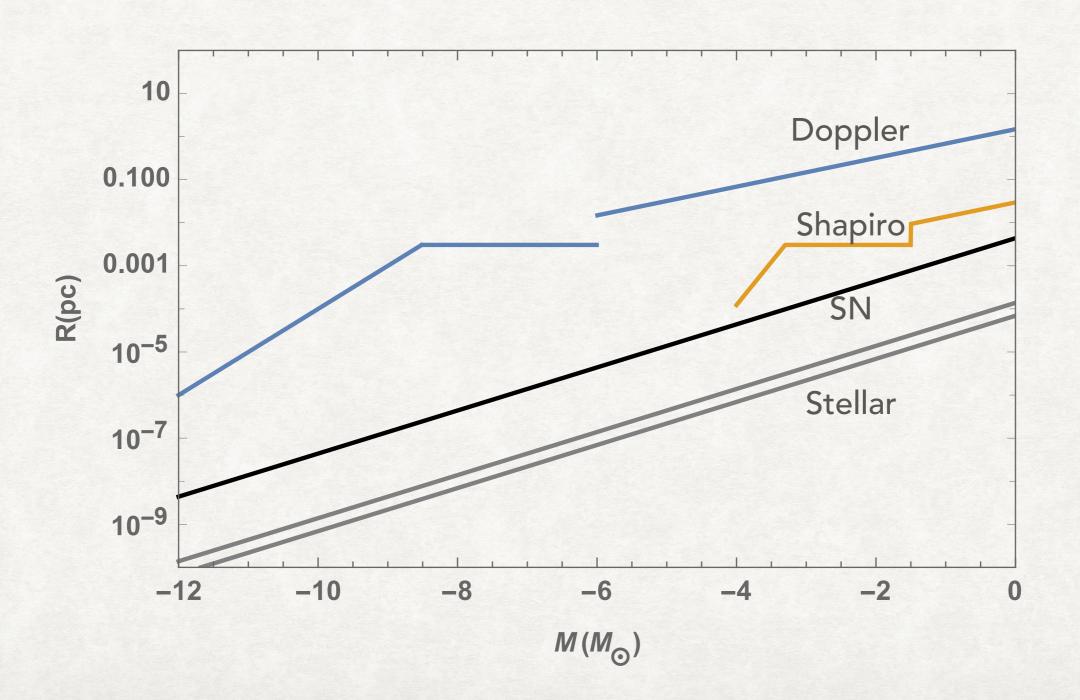
- We have seen point-like objects till now.
- If size of the object < impact parameter, Gauss' law: treat object as point like
- Signal loss if object size > impact parameter.
- Can get conservative estimate with M_{enc}(b).
- For illustrative purposes, NFW halo with clumpiness parametrized by the concentration parameter $c=r_{vir}/r_{c.}$
- retrieve PBH in the large c limit

PROBES VS OBJECT SIZE

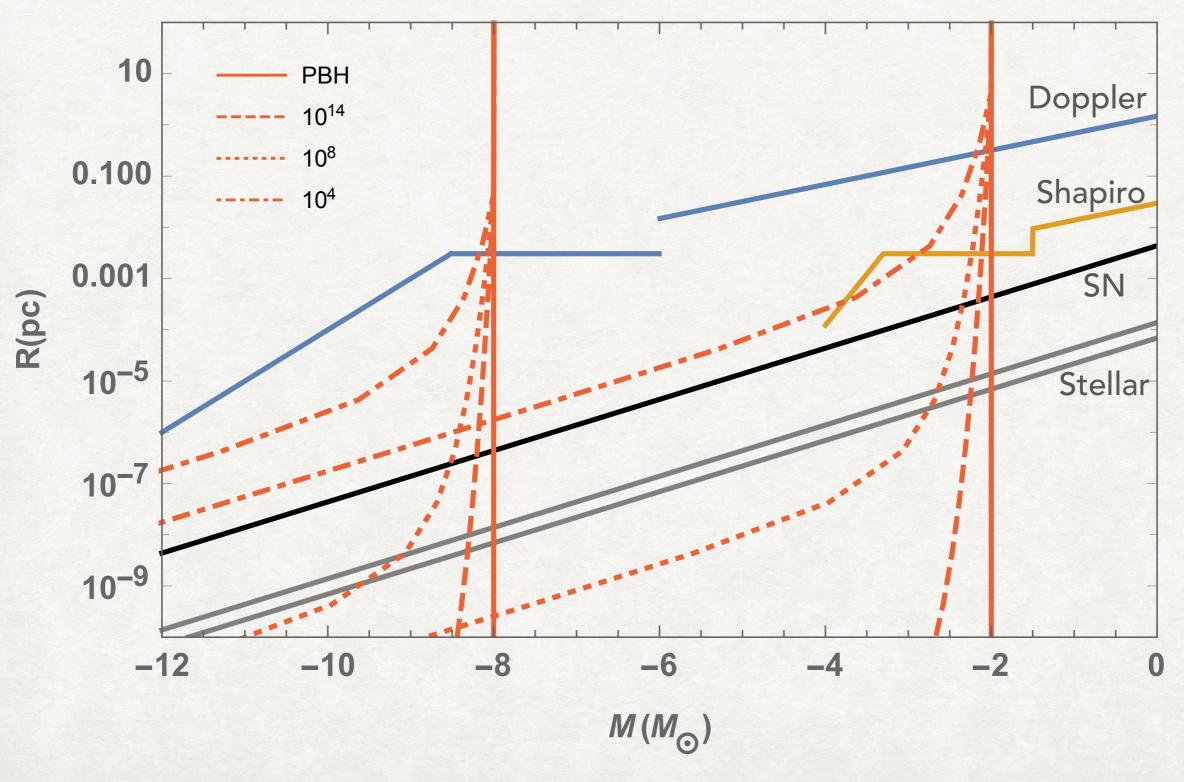
Lensing: (Billion Stars x few hours)

• PTA: (100-1000 pulsars x few years)

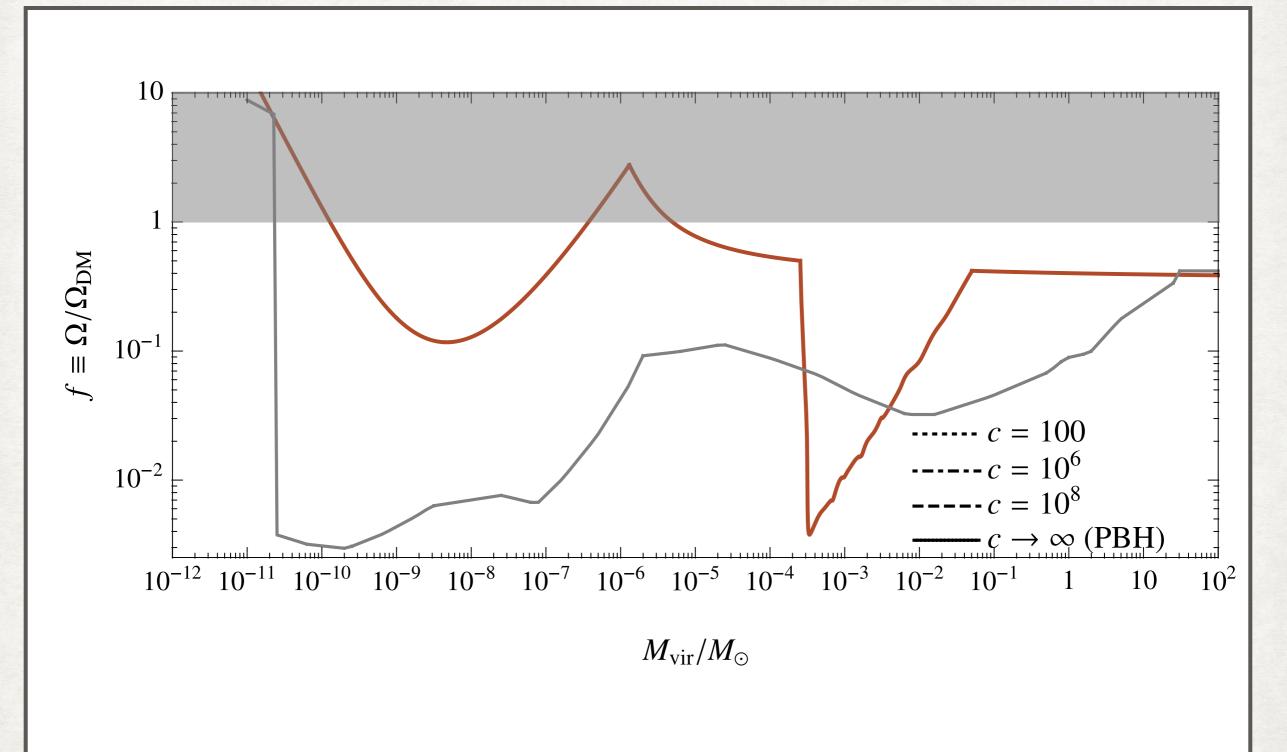
small impact parameter enormous impact parameter

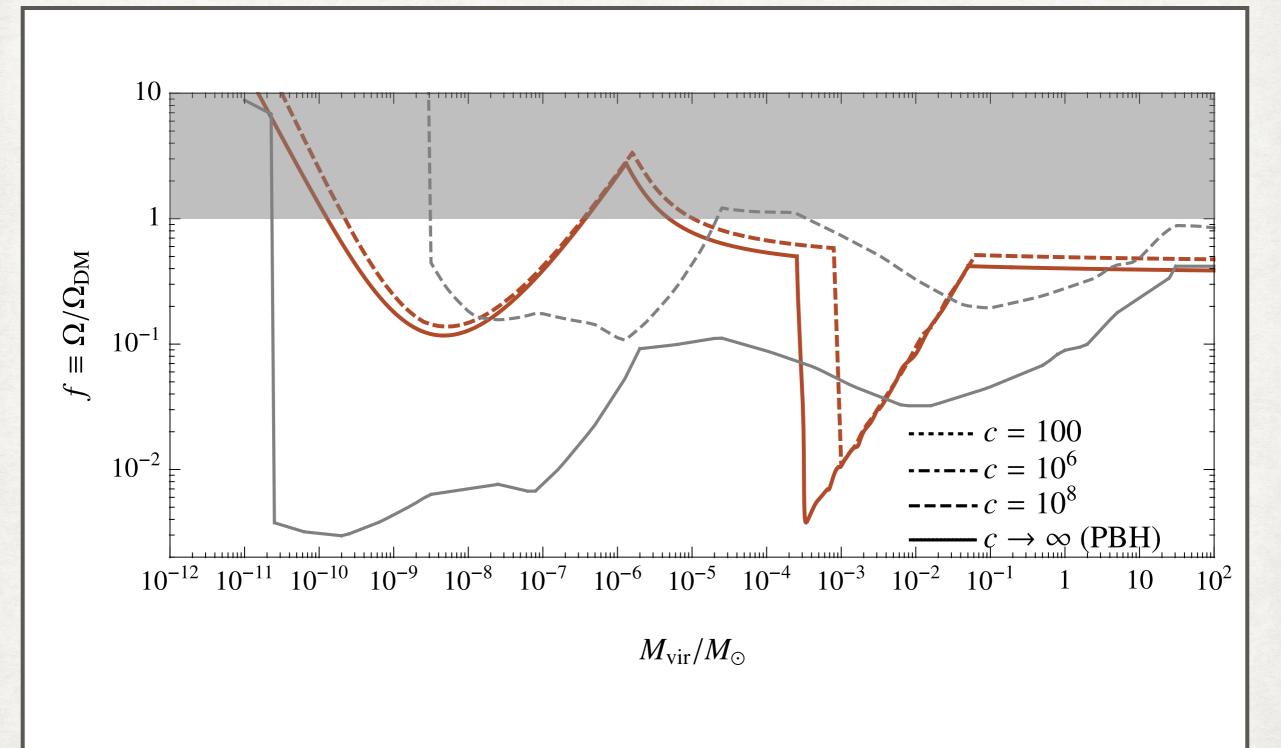


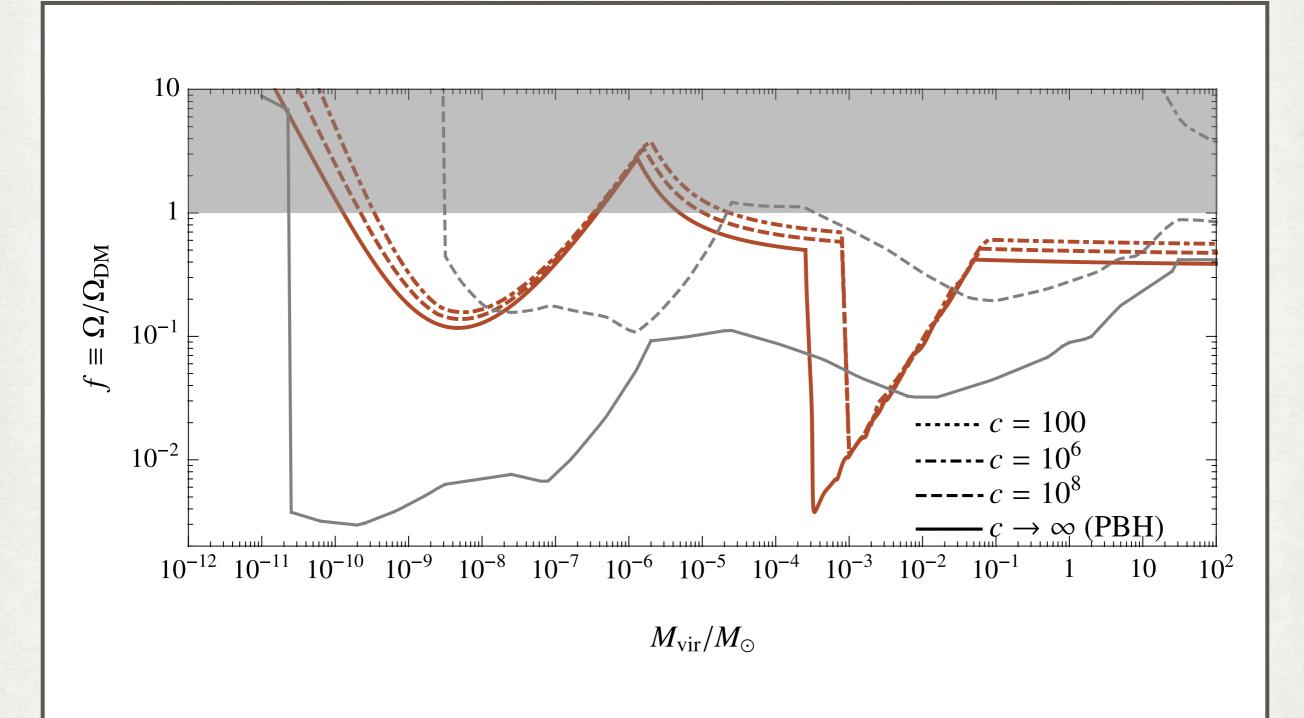
SENSITIVITY TO DIFFUSE HALOS

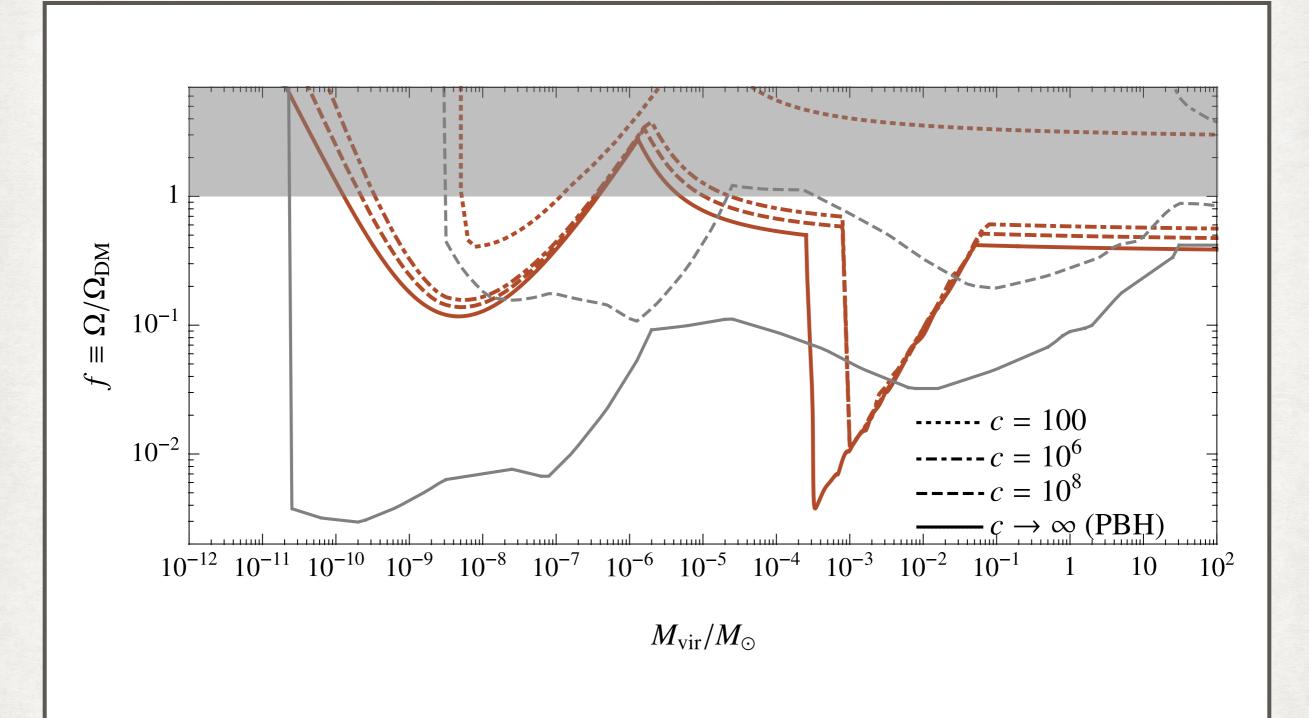


Limits iff red line intersects a probe radius







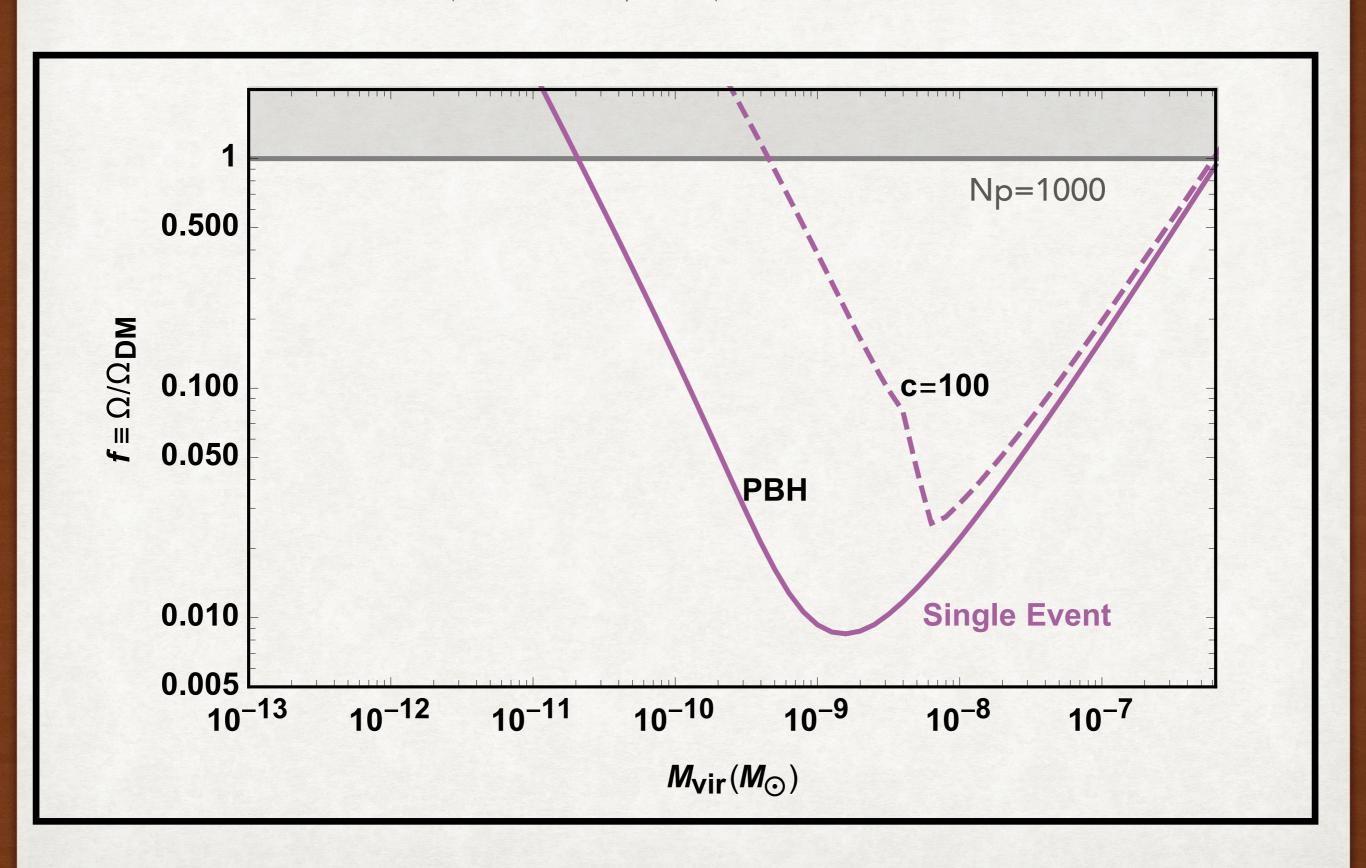


WORK IN PROGRESS

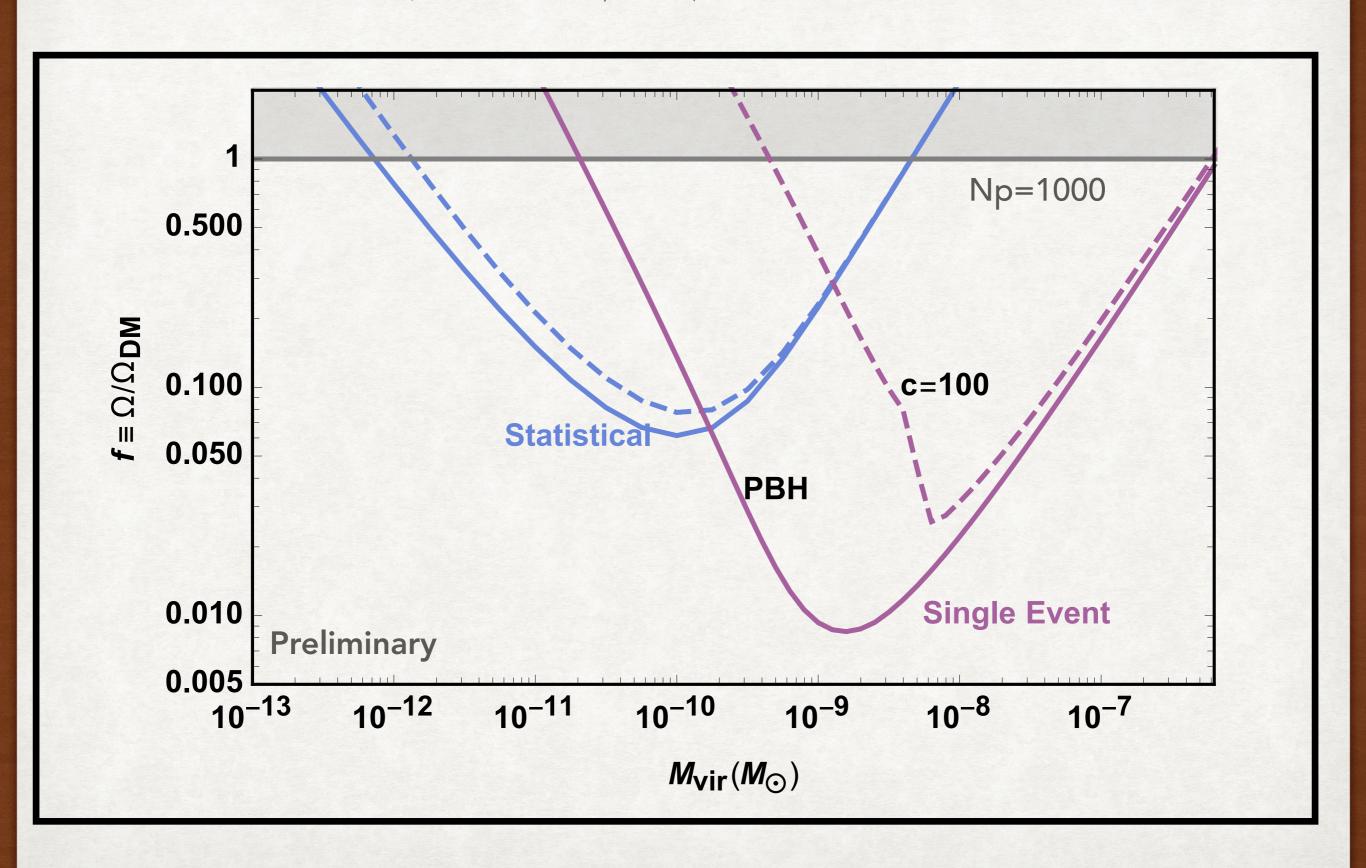
A MULTI EVENT STATISTICAL SIGNAL SEARCH

- Till now only single blip:
- advantage: Shape based filter search with very good sensitivity
- · At small masses, can we instead do a statistical search?
- Yes: Compare a Binary merger GW signal to searching for a stochastic signal.
- Doppler: look for a random walk (memory effect)
- Shapiro: Statistics of multiple blips

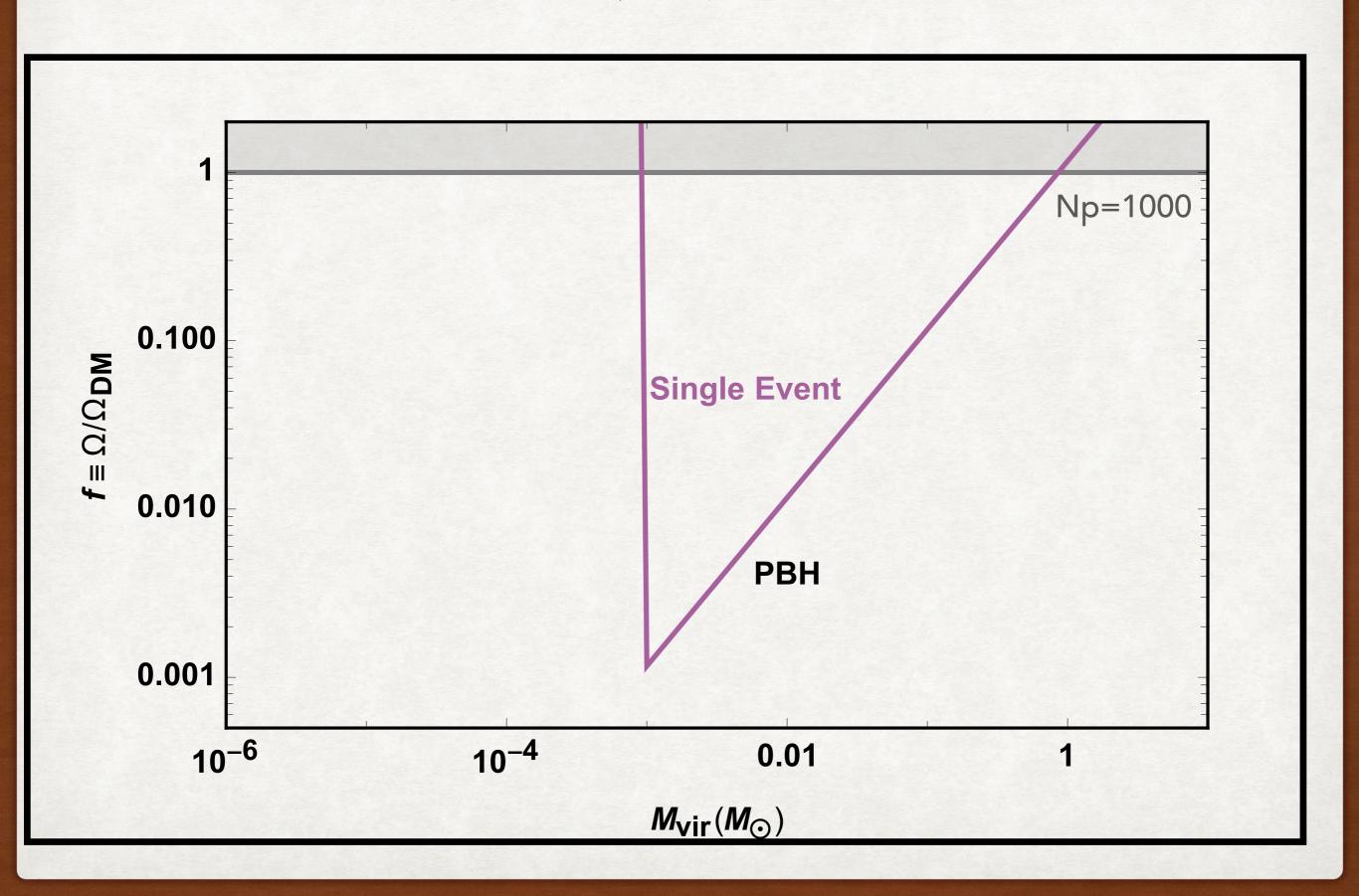
SINGLE EVENT DOPPLER



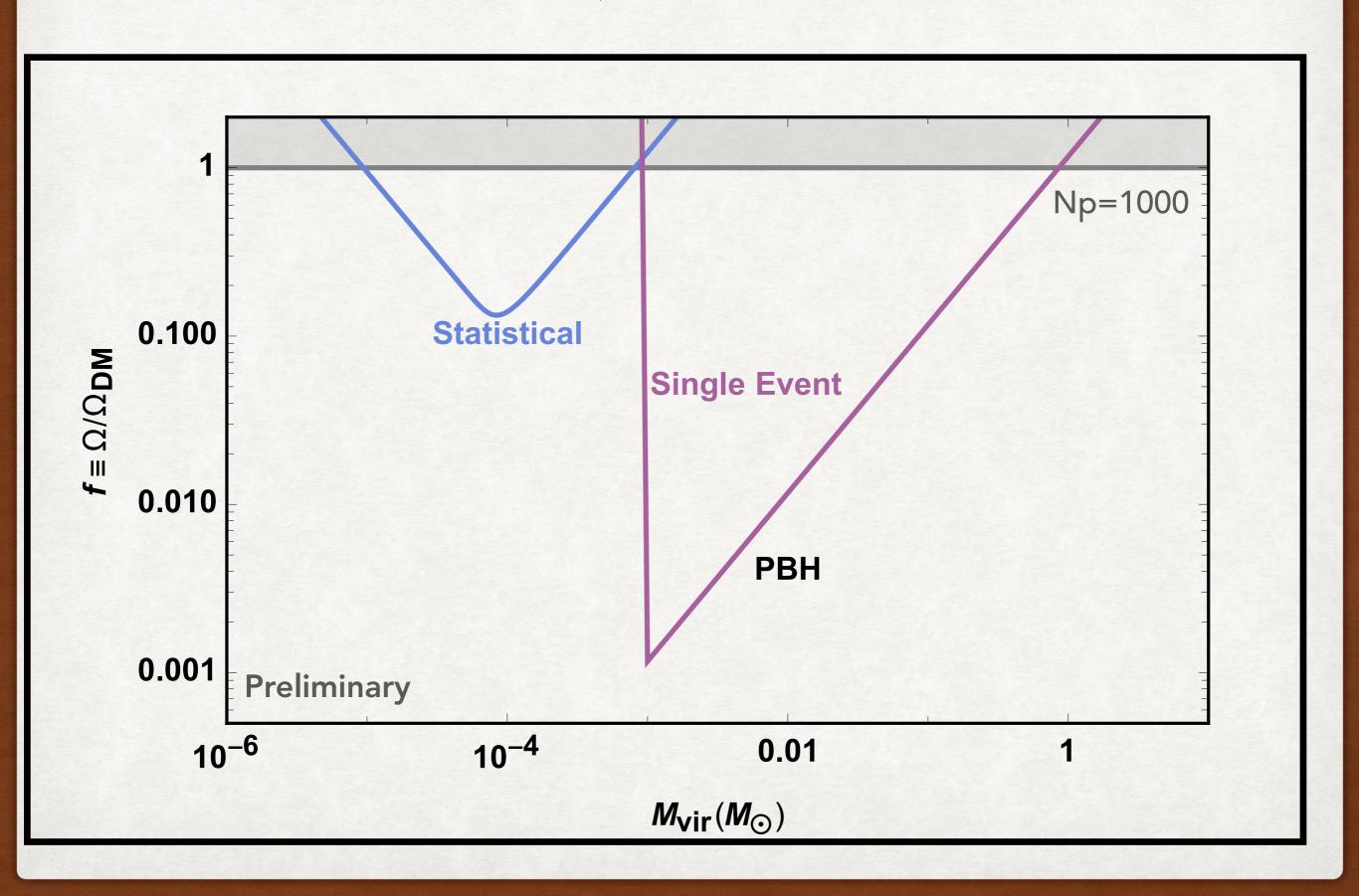
MULTI-EVENT DOPPLER



SINGLE-EVENT SHAPIRO



MULTI-EVENT SHAPIRO



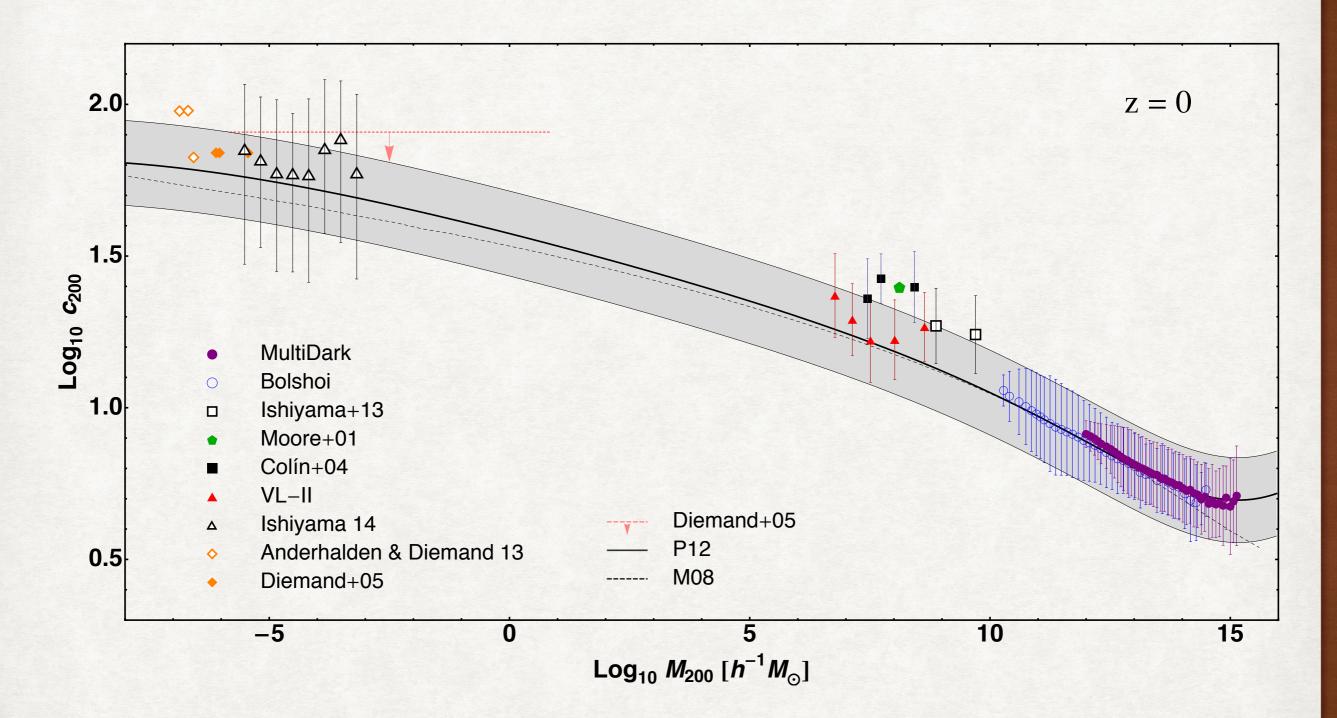
EXTENDED HALO MASS FUNCTION

- Assume typical scale-free Halo mass function from Press-Schechter.
- dn/dM ~ M-2
- M_{min}, the IR cutoff
- M_{max}, the UV cutoff
- Equal amount of DM in every decade of masses,
- Even large $M_{\text{max}}/M_{\text{min}}$ can be probed using sensitivity solely in a small subset window.

LIMIT SETTING PARAMETERS

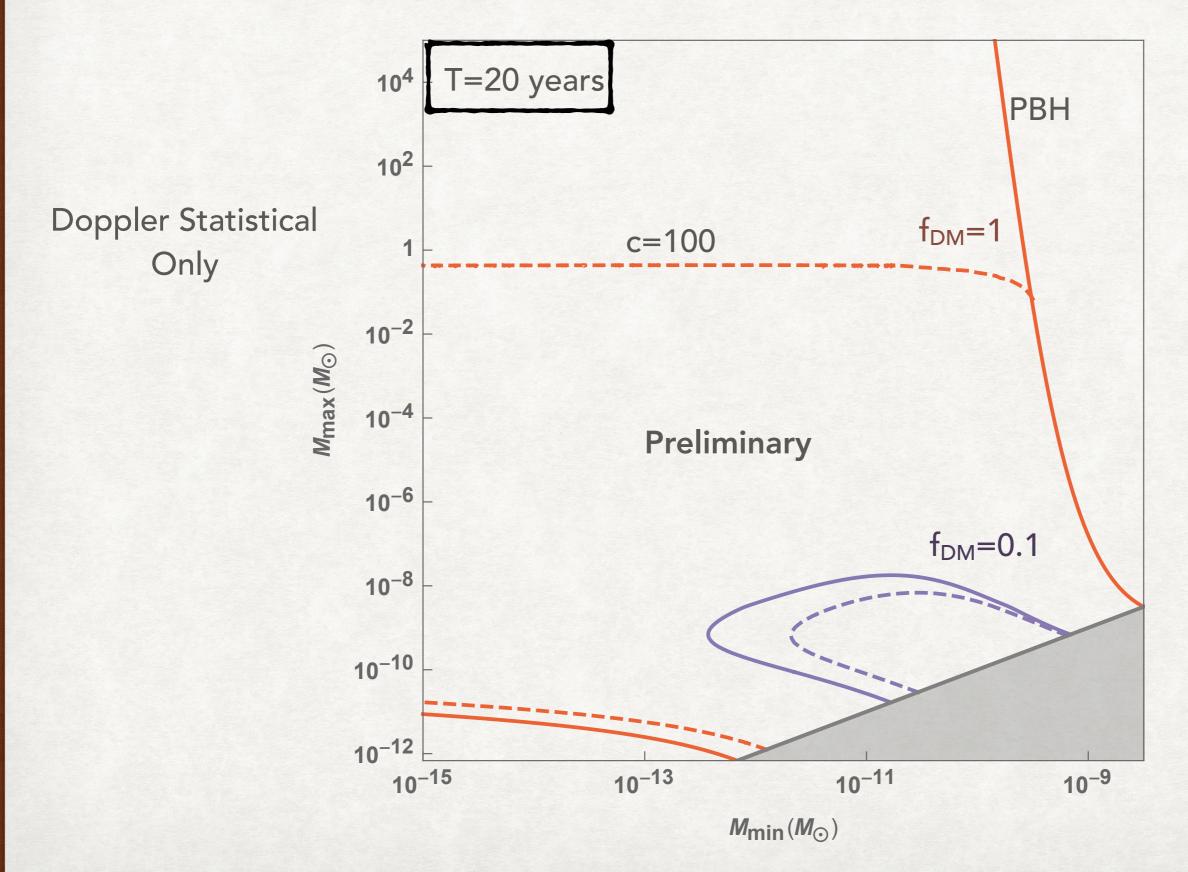
- Set Limits for
- c, the concentration parameter
- f_{DM} the fraction of dark matter that has not disrupted
- Ignoring tidal disruption and sweeping it into c and f_{DM}

GOAL: C=100

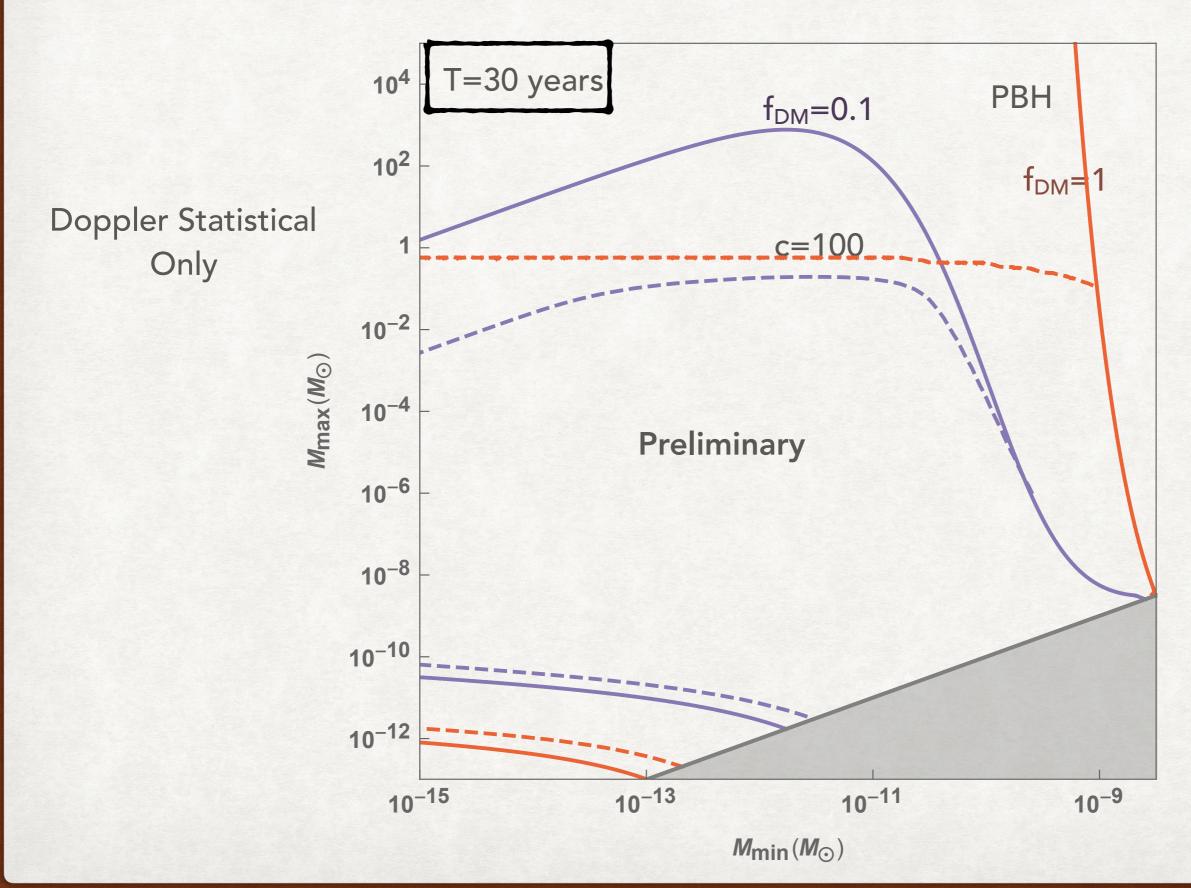


M usually cut off at 10-6 because WIMPs wash out small-scale structure...

LIMIT PROJECTIONS ON EXTENDED MASS



LIMIT PROJECTIONS ON EXTENDED MASS



TO DO LIST

- MSPs across the GC?
- in DM rich environments?
- Extra galactic MSPs?

BACKUP

PROBES VS OBJECT SIZE

$$r_{\rm PTA} \sim 10^{-3} \ {\rm pc} \times \left\{ \begin{array}{l} \frac{M}{10^{-9} M_{\odot}} & ({\rm Doppler \ Dynamic}) \\ \left(\frac{M}{10^{-3} M_{\odot}}\right)^2 & ({\rm Shapiro \ Dynamic}) \end{array} \right.$$

$$r_{\rm PTA} \sim 10^{-3} \ {
m pc} \times \left\{ \begin{array}{l} \left(\frac{M}{10^{-8} M_{\odot}}\right)^{\frac{1}{3}} & {
m (Doppler Static)} \\ \left(\frac{M}{10^{-3} M_{\odot}}\right)^{\frac{1}{3}} & {
m (Shapiro Static)} \end{array} \right.$$

$$r_E \sim \begin{cases} 10^{-6} \text{ pc} \left(\frac{M}{10^{-4} M_{\odot}}\right)^{\frac{1}{2}} & \text{(Stellar Lensing)} \\ 10^{-2} \text{ pc} \left(\frac{M}{10 M_{\odot}}\right)^{\frac{1}{2}} & \text{(Supernovae Lensing)} \end{cases}$$

TOWARDS A STATISTICAL DESCRIPTION OF SIGNALS

- Gravitational effects can be related to the local density
- Correlations in delay are measuring correlations in density and hence measuring the matter power spectrum
- Can handle multiple lower threshold events
- Need to be careful about how much is sampled cuts in momentum space
- Thus not only Power, but halo mass function important
- Doppler Random walk signal
- Shapiro Statistics of Multiple Blips
- More careful with background

TO DO

- Limits on current power spectrum.
- Can then project that to limits on particle models, by assuming a rough clustering model (or N-body simulations?)
- Pulsars near Galactic center to increase local DM density.
- Extra-galactic pulsars to increase baseline.

CONCLUSIONS

- Pulsar timing can probe structure at a wide range of small scales.
- Doppler and Shapiro delays, especially in the dynamic regime, can provide a compelling discovery signal for DM subhalos
- Limits on PBH provide a second probe of the lensing region
- · Limits on more diffuse objects could do much better than lensing

COMPACT OBJECTS

Parametrize the profile as NFW or UCMH.

$$\rho(r, M_{\text{vir}}) = \frac{\rho_s}{(r/r_s)^{\alpha} (1 + r/r_s)^{\beta}}$$

$$\alpha=1,\ \beta=2$$

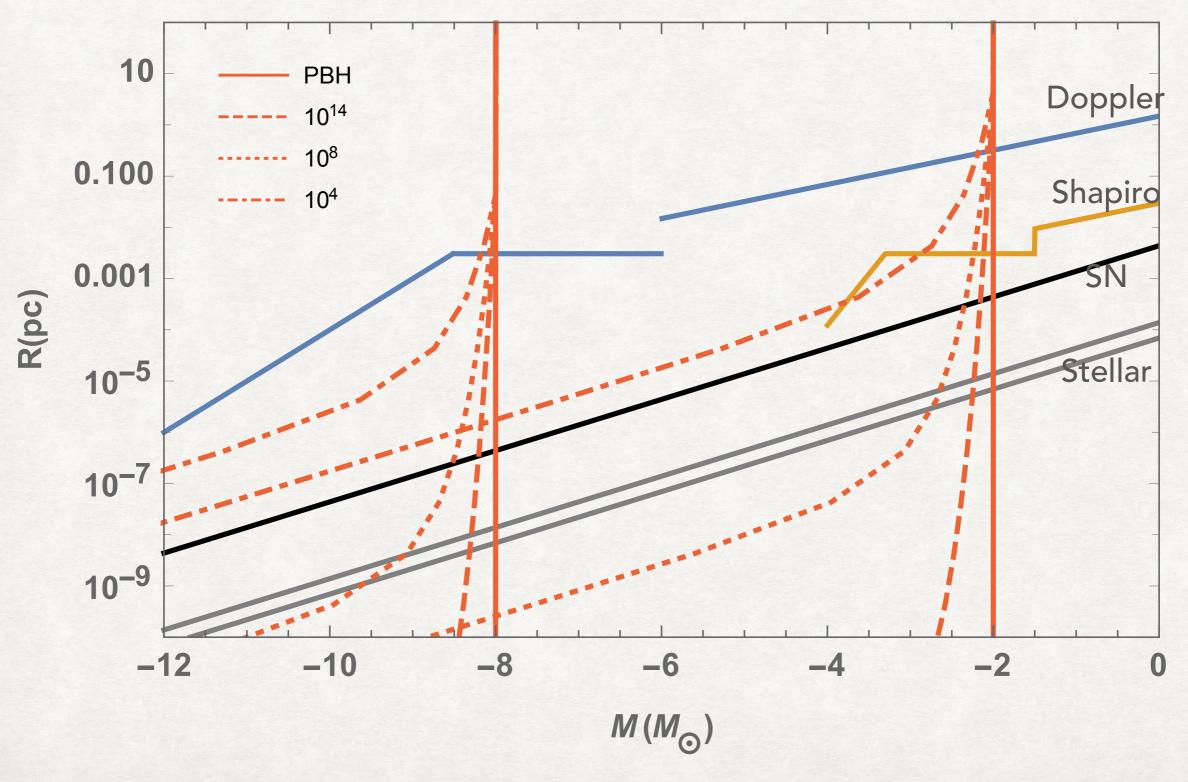
$$\alpha=9/4,\ \beta=0$$
 UCMH

$$r_{\rm vir} \equiv (3M_{\rm vir}/800\pi \dot{\rho}_c)^{1/3}$$

NFW

$$c \equiv r_{\rm vir}/r_s$$

SENSITIVITY TO CONCENTRATION



Limits iff red line intersects a probe radius

SIGNAL TO NOISE RATIO (STATIC SIGNALS)

- In the limit that you don't see the whole signal, Taylor expand.
- A constant first derivative i.e. spin-down or sometimes even spinup is already observed (and hard to ascertain from first principles).
- Subtracted as part of the fitting procedure.
- Second derivative much less common.
- Non-observation of second derivative can be used to set constraints.

Doppler

$$\frac{\ddot{\nu}}{\nu} \simeq \frac{2GMv}{r_{\min}^3} \sim 3 \times 10^{-32} \left(\frac{N_P f}{200}\right) \text{ Hz}^2$$

Shapiro

$$\frac{\ddot{\nu}}{\nu} \simeq \frac{16GMv^3}{r_{\times,\text{min}}^3}$$

$$\sim 8 \times 10^{-33} \left(\frac{N_P f}{200}\right)^{\frac{3}{2}} \left(\frac{M_{\odot}}{M}\right)^{\frac{1}{2}} \left(\frac{d}{\text{kpc}}\right)^{\frac{3}{2}} \text{ Hz}^2$$

STATIC SIGNAL SENSITIVITY

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Uncertainty in second derivative purely from rms fluctuations

$$\sigma_{\ddot{\nu}/\nu} = 6\sqrt{\frac{2800\Delta t}{T}} \frac{t_{\rm RMS}}{T^3}$$

$$f_{\rm D, \ stat} \lesssim 0.4 \left(\frac{200}{N_P}\right) \left(\frac{20 \ \rm yr}{T}\right)^{\frac{7}{2}}$$

$$f_{\rm S, \ stat} \lesssim \left(\frac{200}{N_P}\right) \left(\frac{M}{M_{\odot}}\right)^{\frac{1}{3}} \left(\frac{20 \ \rm yr}{T}\right)^{\frac{7}{3}} \left(\frac{\rm kpc}{d}\right)$$

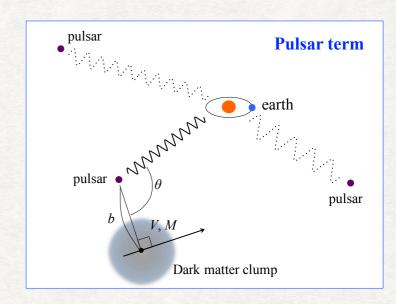
Notice no M dependence here

Timescales

- To determine typical timescale, let us determine object of closest approach
- Cross-section for Doppler, is a circle.
- Remembering $|\mathbf{b}| = \tau v$

$$\tau_{\min} \simeq \frac{1}{v} \sqrt{\frac{M}{N_P f \rho_{\rm DM} v T}}$$

$$\sim \frac{20 \text{ yr}}{\sqrt{N_P f}} \left(\frac{M}{10^{-9} M_{\odot}}\right)^{\frac{1}{2}} \left(\frac{20 \text{ yr}}{T}\right)^{\frac{1}{2}}$$



 N_P pulsars N_P x cross-section

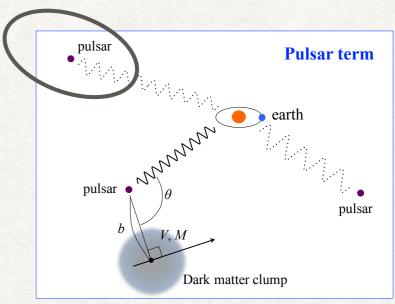
Number of pulsars

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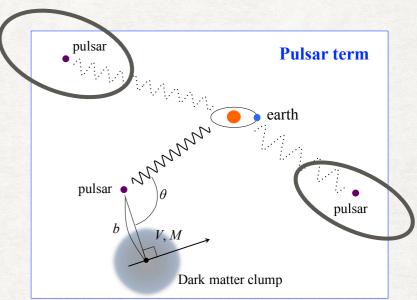
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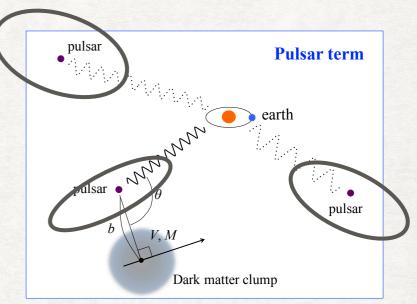
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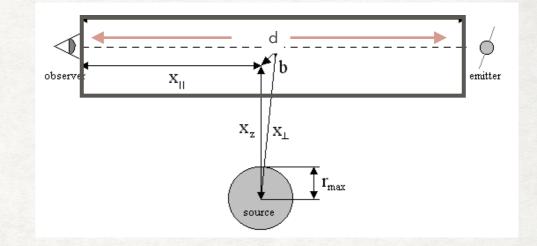
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Timescales

Cross-section for Shapiro is a rectangle

$$\tau_{\rm min} \simeq \frac{2}{v} \frac{M}{N_P f \rho_{\rm DM} v T d},$$

$$\sim \frac{20 \text{ yr}}{N_P f} \left(\frac{M}{10^{-4} M_{\odot}}\right) \left(\frac{20 \text{ yr}}{T}\right) \left(\frac{\text{kpc}}{d}\right)$$



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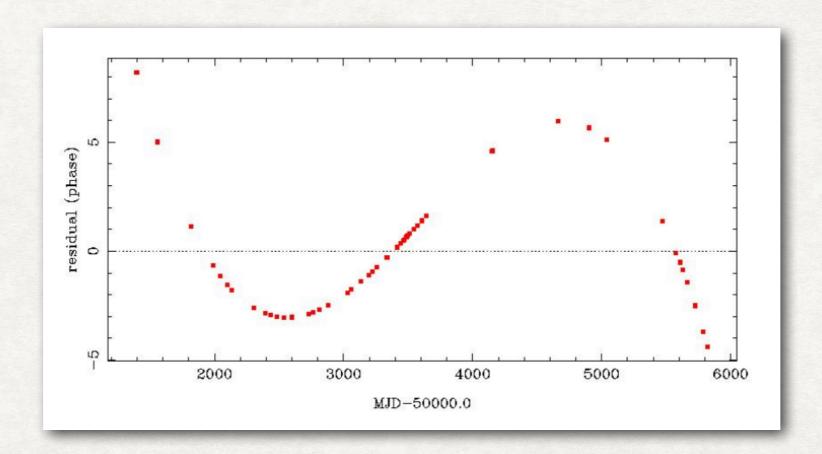
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Notice no M dependence here

STATIC BACKGROUNDS

- A few pulsars already display non-zero second derivative.
- Will need to supplement with E&M observations to subtract known nearby objects.



MONTECARLO SIMULATION

- Assume PBHs randomly distributed
- Isotropic Maxwell distribution in velocity truncated at vesc.
- Simulate N_p randomly distributed pulsars at appropriate distances.
- Simulate order O(10⁴) universes and require more than 95% universes pass SNR>4.

PTA FOR DARK MATTER

- Ultralight DM causing GW like delays
- [Khmelnitsky, Rubakov 1309.5888], [Graham, Kaplan, Mardon, Rajendran, Terrano 1512.06165]
- PTAs are sensitive accelerometers: Doppler Delay
- [Seto, Corray astro-ph/0702586] , [Baghram, Afshordi, Zurek 1101.5487] [Kashiyama, Seto 1208.4101], [Kazumi, Oguri, Masamune 1801.07847]
- Gravitational potential wells along the light path: Shapiro Delay
- [Siegel, 0801.3458], [Siegel, Hertzberk, Fry astro-ph/0702546],
 [Baghram, Afshordi, Zurek 1101.5487], [Clark, Lewis, Scott 1509.02938],
 [Schutz, Liu 1610.04234]
- Our Work
 Explicit calculations for SNR calculations for Doppler and Shapiro Delays
 Correcting several mistakes in literature
 Extending analysis to diffuse halos