

New Tools  
Lattice **QCD**: New Ideas  
New Results

---



Andreas S. Kronfeld  
Fermilab & IAS TU München



2019 Meeting of the APS Division of Particles and Fields  
Northeastern University, Boston | July 29, 2019



# Basic Idea

---

- Discretize space to have Kronecker  $\delta_{xy}$  instead of Dirac  $\delta(x-y)$  so the degrees of freedom are countable [Z. Phys. 56 (1929) 1].
- Set up path integral for these degrees of freedom, which introduces discrete time steps [RMP 20 (1948) 367].
- Figure out how to enforce gauge invariance [PRD 10 (1974) 2445].
- End up with mathematically well-posed definition of QFT:

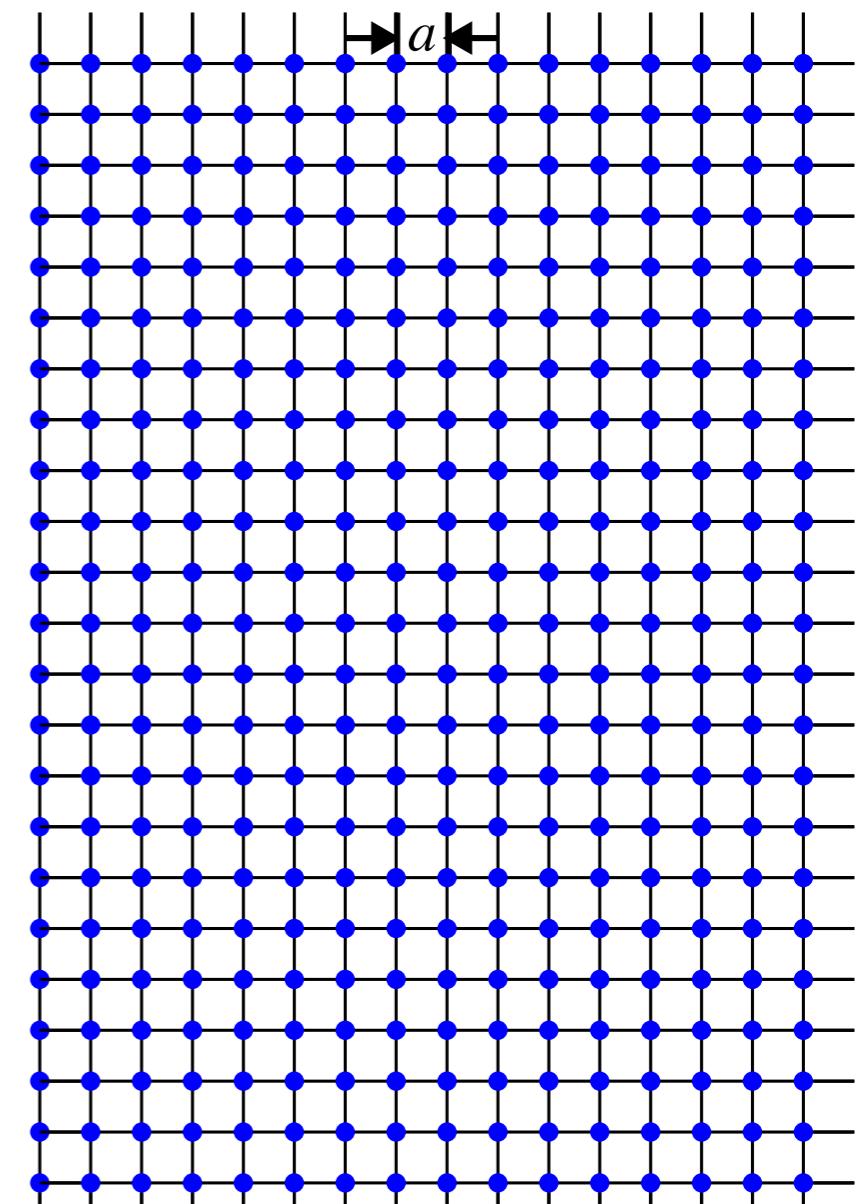
$$\langle \bullet \rangle = \frac{1}{Z} \int \boxed{\mathcal{D}U \mathcal{D}\Psi \mathcal{D}\bar{\Psi}}_{\text{MC}} \exp(-S) [\bullet]_{\text{hand}}$$

- In finite volume, the number of degrees of freedom is finite (albeit large).



- With a finite number of degrees of freedom, a computer can integrate.
- Random number generator (Markov chain Monte Carlo) produces **lattice data** varying  

$$(a, L, L_4, \{m_{\text{sea}}\}, m_q, m_Q)$$
- Extract masses, matrix elements, etc., from various correlation functions.
- Use **effective field theories** to fit the lattice data to obtain results with
  - $a \rightarrow 0$ ,
  - $L, L_4 \rightarrow \infty$ ,
  - quark masses tuned to  $n_f$  hadron masses.



# Resources

---

- USQCD Whitepapers:
  - Quark and lepton flavor, Lehner, Meinel *et al.* [arXiv:1904.09479](#);
  - Beyond the standard model, Neil *et al.* [arXiv:1904.09964](#);
  - Neutrino-nucleus scattering, ASK, Richards *et al.* [arXiv:1904.09931](#);
  - Fundamental symmetries, Cirigliano, Davoudi *et al.* [arXiv:1904.09704](#);
  - Hadrons and nuclei, Detmold, Edwards *et al.* [arXiv:1904.09512](#);
  - Hot, dense QCD, Karsch, Mukherjee *et al.*, [arXiv:1904.09951](#);
  - Calculations at the exascale, Joó, Jung *et al.*, [arXiv:1904.09725](#);
- Flavor Lattice Averaging Group, [arXiv:1902.08191](#), updates at **FLAG2019**.

# Outline

focus on material not covered in || sessions

---

- Introduction
- Something Cool
- Quark Flavor Physics: Form Factors and Decay Constants
- Higgs Physics: Quark Masses;  $\alpha_s$ ; PDFs
- Lepton-Flavor Physics: Muon  $g-2$ ; Nucleonic Ingredients for  $\nu A$  Scattering
- Outlook

Something Cool

# Work and Fields

arXiv:1803.05656

- The work done against an electric force is

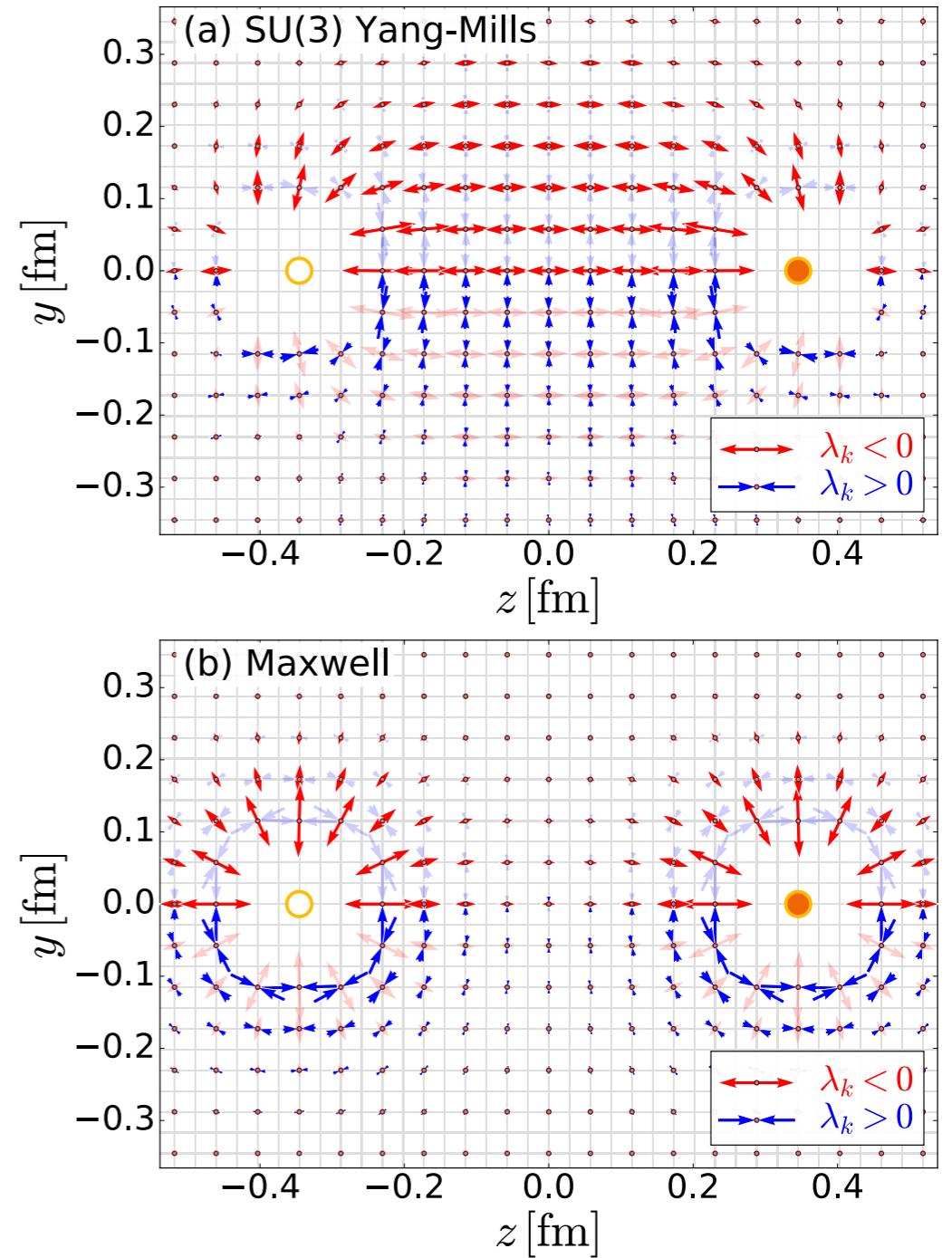
$$W = \mathbb{P} \int_{x_1}^{x_2} \mathbf{F}(x) \cdot d\mathbf{x}$$

where  $\mathbf{F} = q\mathbf{E}$  and  $\mathbf{E}$  is the field of the charge distribution apart from  $q$ .

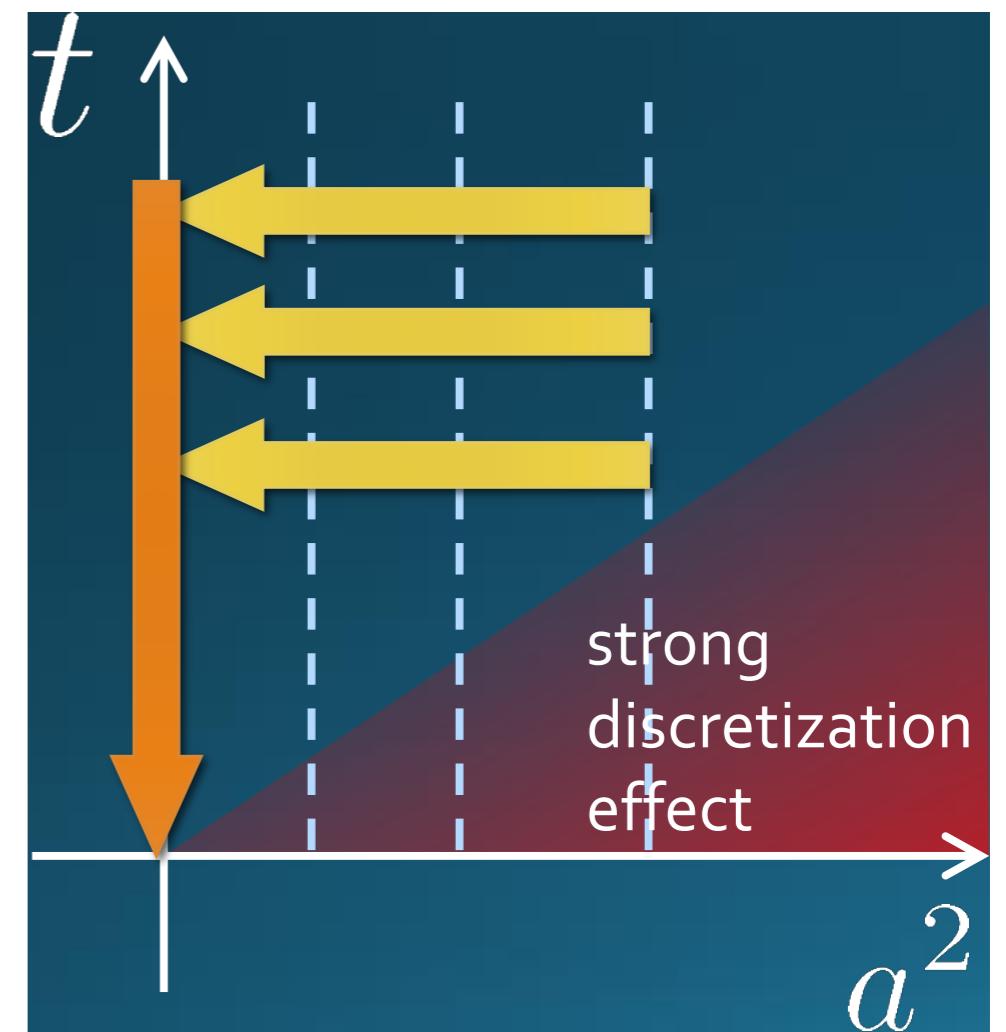
- On the other hand, it should also be the change in energy of the total field,  $\mathbf{E}_{\text{tot}}$ , including that of  $q$ :

$$\mathcal{E} = \frac{1}{2} \int d^3x \left[ |\mathbf{E}_{\text{tot}}|_q^2 \text{ at } \mathbf{x}_2 - |\mathbf{E}_{\text{tot}}|_q^2 \text{ at } \mathbf{x}_1 \right]$$

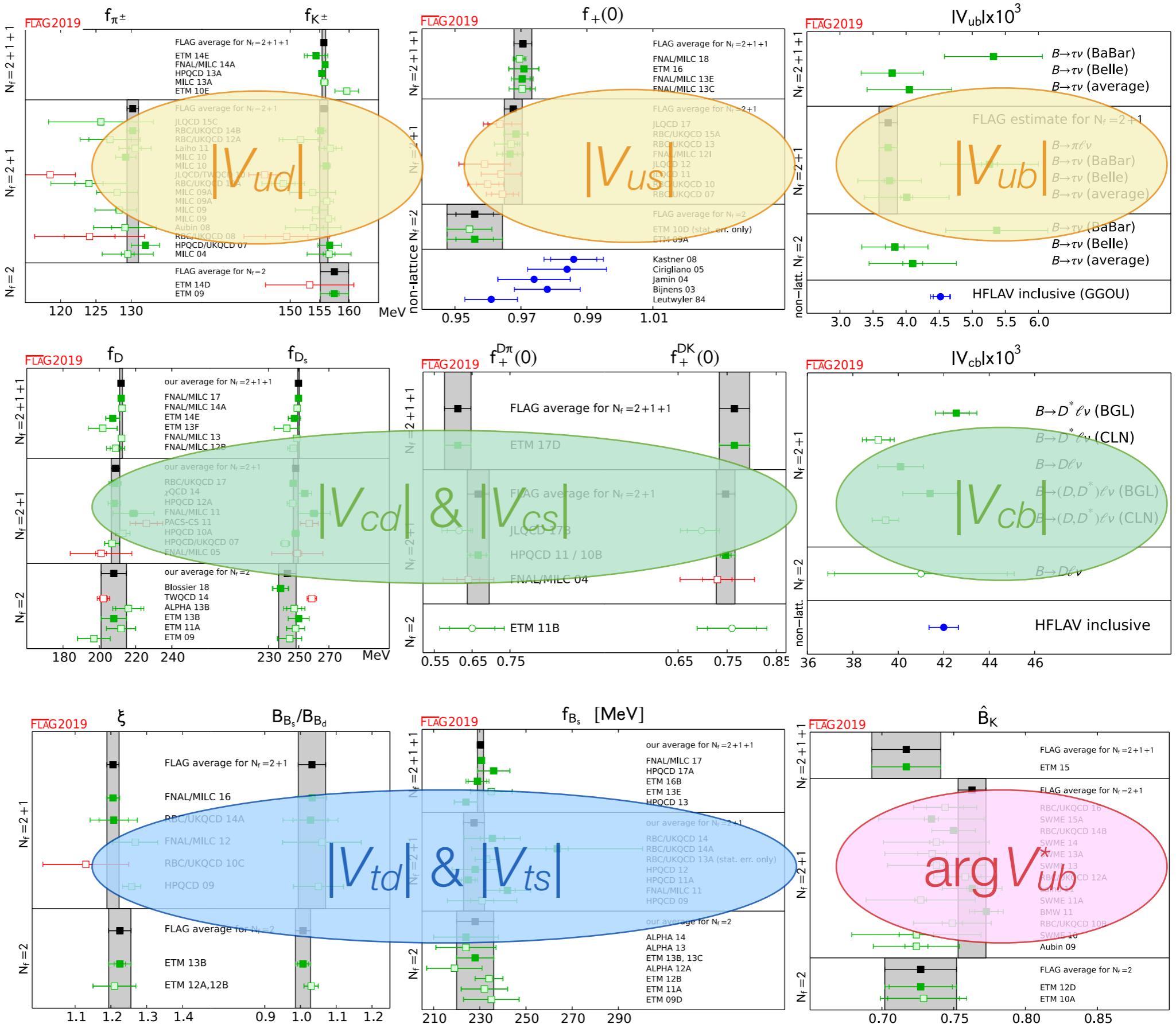
- Yanagihara et al. show this works in YM.



- Yanagihara *et al.* [arXiv:1803.05656] did more, showing how the stress part of the stress-energy tensor produces the force even in a confining quantum system.
- To do so, they had to overcome a problem with the lattice-UV-regulated stress-energy tensor: power-law mixing with the unit operator.
- Old idea, more recently developed into a powerful tool: "**gradient flow**".
- Smears gauge field over distance  $\sqrt{8t}$ ; flow time  $t$  has dimension distance<sup>2</sup>.
- **No power law divergences**, if you take  $a \rightarrow 0$  first, then  $t \rightarrow 0$ .
- Use to renormalize any operator.

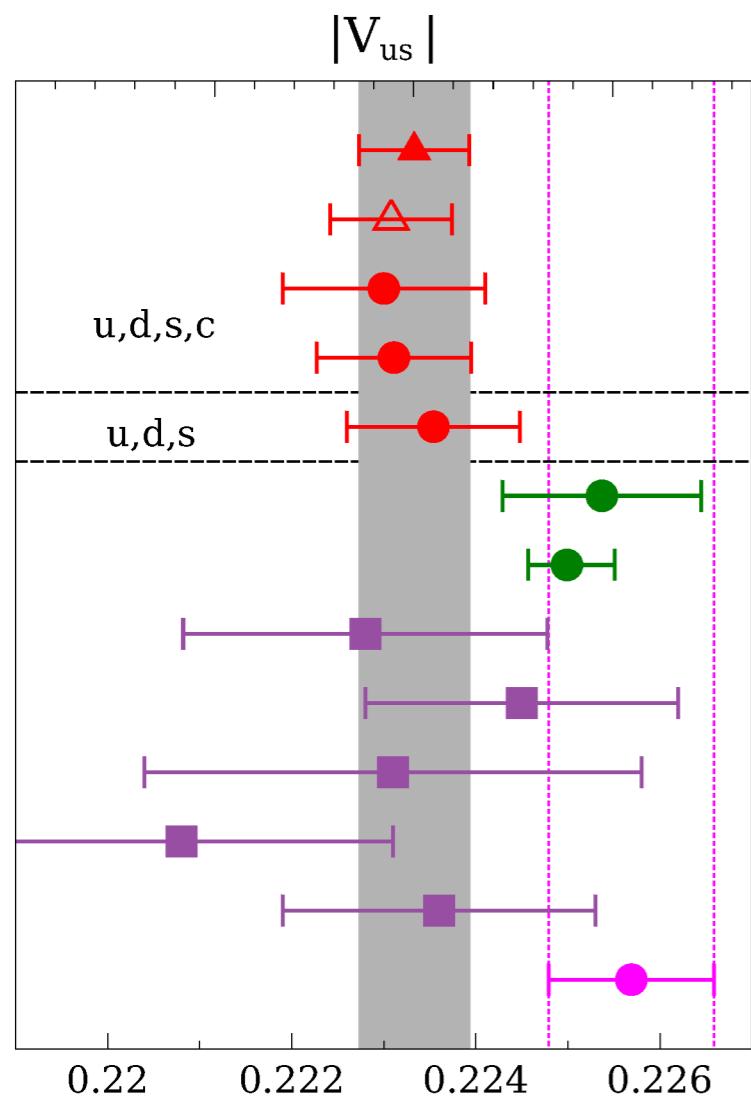


# Quark-Flavor Physics



# $|V_{us}|$ : Semileptonic

- QCD uncertainty in  $|V_{us}|$  now the same as experiment [arXiv:1809.02827].



This work

$$|V_{us}| = 0.22333(42)_{\text{expt}}(43)_{f_{+}(0)}$$

This work (only neutral kaon exp. data)

$K_{l3}$  ETMC 2016

$K_{l3}$  Fermilab Lattice/MILC 2014

$K_{l3}$  RBC/UKQCD 2014

$K_{l2}$  FLAG 2016 +  $f_K$  FLAG  $N_f=2+1$

$K_{l2} + f_K/f_\pi$  Fermilab Lattice/MILC 2017

$\tau \rightarrow s$  inclusive, Boyle et al. 2018

$\tau \rightarrow s$  inclusive +  $K_{l2}$  input, Boyle et al. 2018

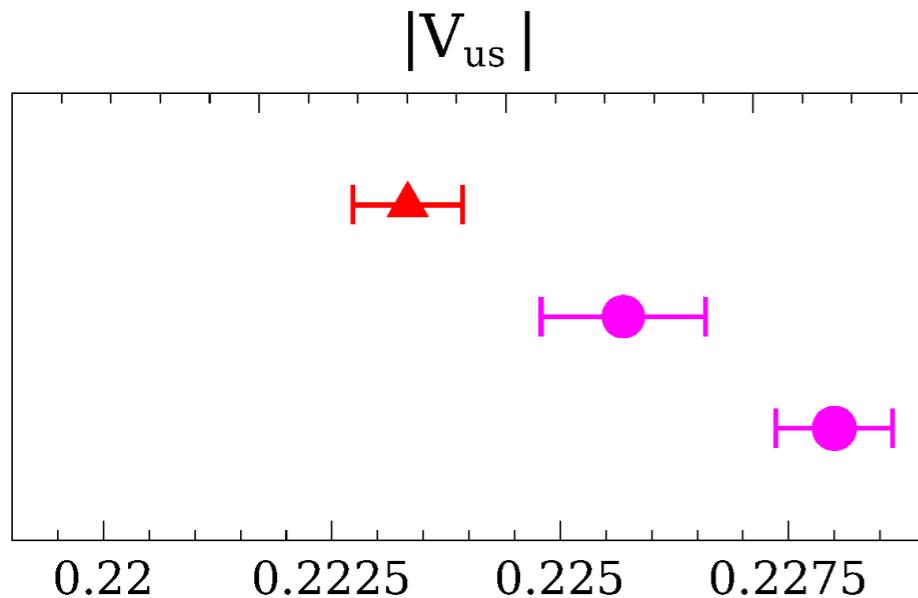
$\tau \rightarrow s$  inclusive, Hudspith et al. 2017

$\tau \rightarrow K \ell \nu / \tau \rightarrow \pi \ell \nu$  HFLAV2017 +  $f_K/f_\pi$  Fermilab Lattice/MILC 2017

Unitarity with  $|V_{ud}| = 0.97420(21)$ , RC from Marciano & Sirlin 2005

# $|V_{us}|$ : Semileptonic

- QCD uncertainty in  $|V_{us}|$  now the same as experiment [arXiv:1809.02827].

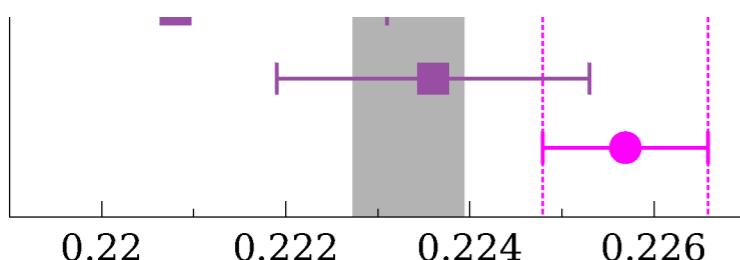


This work

Unitarity with  $|V_{ud}|=0.97420(21)$ , RC from Marciano & Sirlin 2005

Unitarity with  $|V_{ud}|=0.97366(15)$ , RC from Seng et al. 2018

$\sim 2\sigma$  violation of unitarity becomes  $>5\sigma$  with new RC

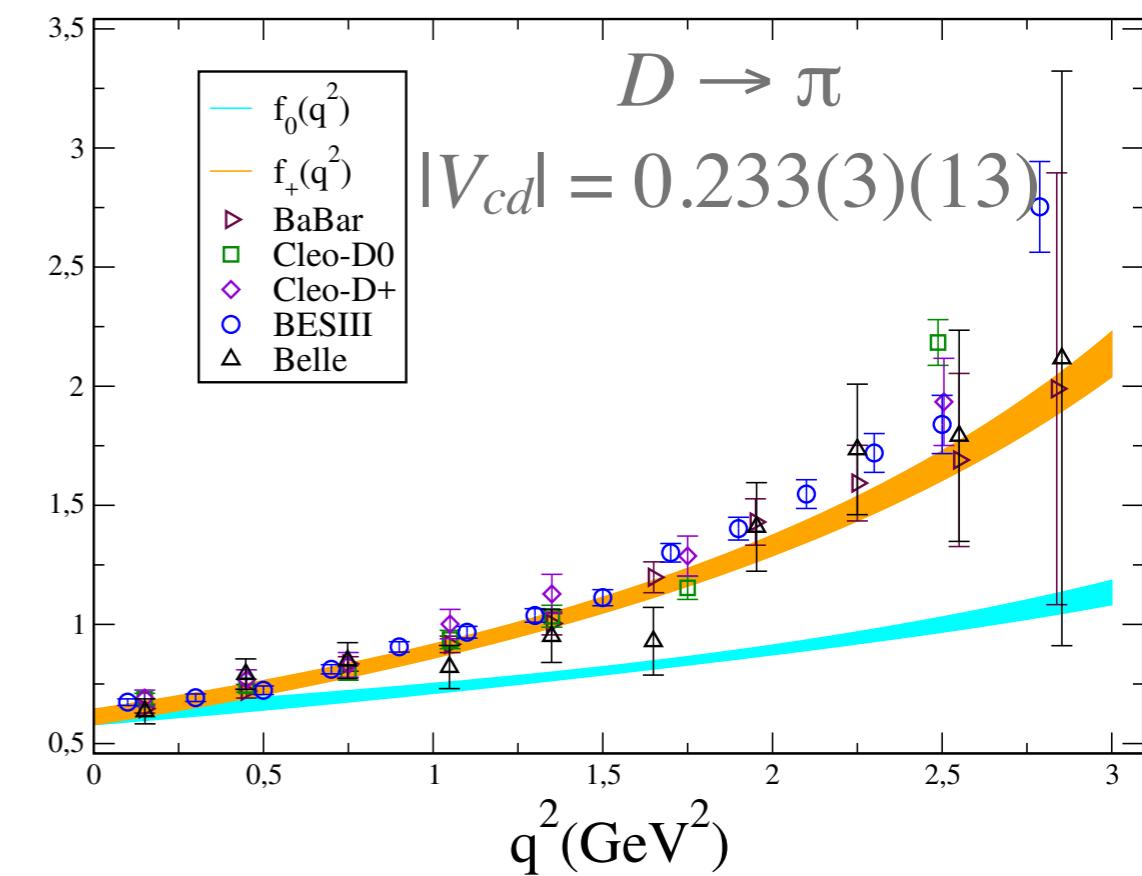
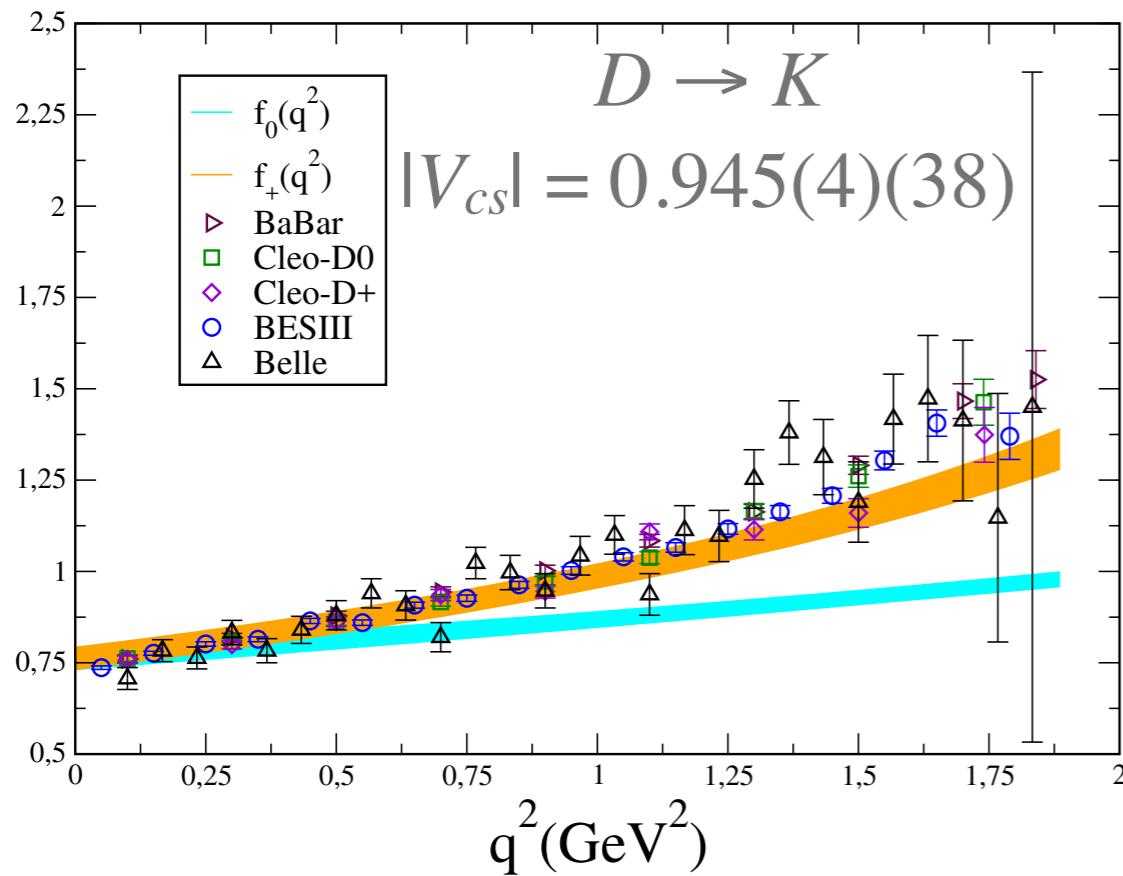


$\tau \rightarrow K \ell \nu / \tau \rightarrow \pi \ell \nu$  HFLAV2017+  $f_K/f_\pi$  Fermilab Lattice/MILC 2017

Unitarity with  $|V_{ud}|=0.97420(21)$ , RC from Marciano & Sirlin 2005

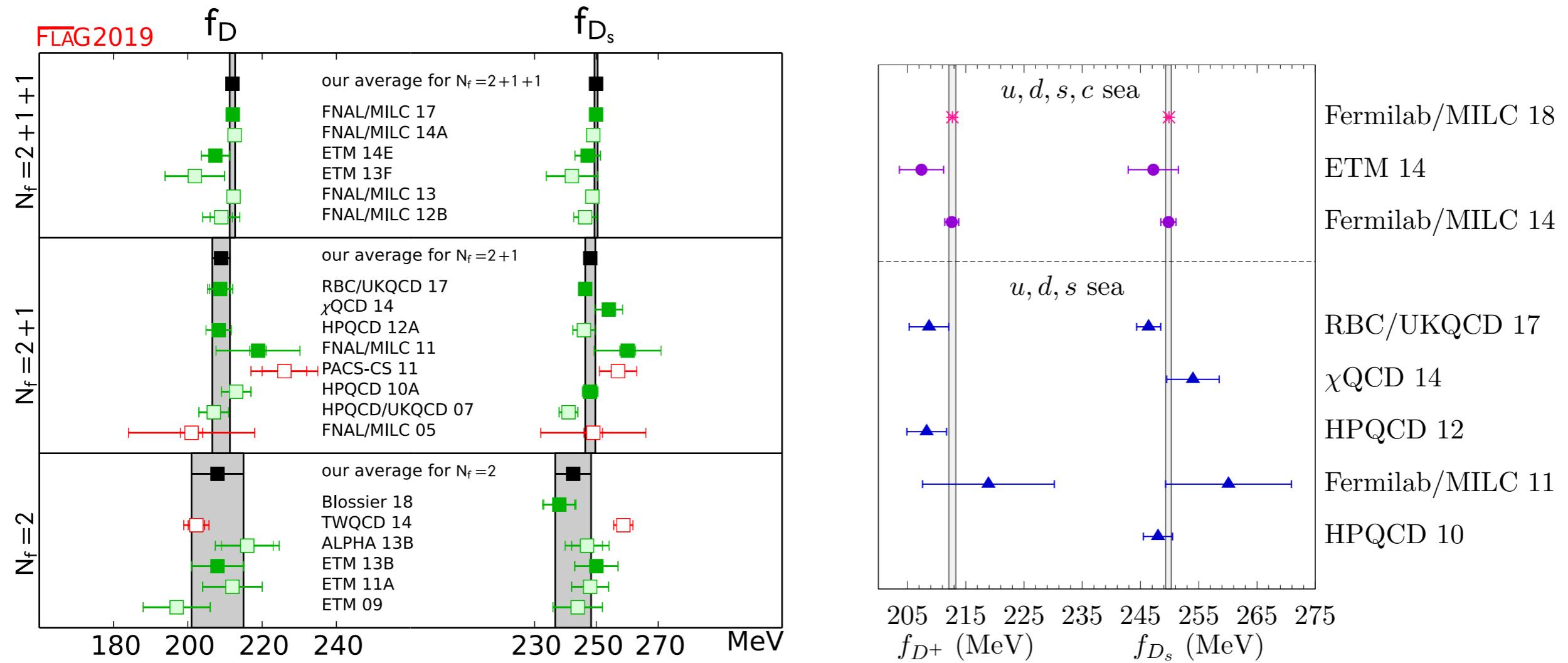
# $|V_{cd}|$ and $|V_{cs}|$ : Semileptonic

- Most recent entry in FLAG 2016 was HPQCD 2010 ( $D \rightarrow K$ ), 2011 ( $D \rightarrow \pi$ ); see also [arXiv:1305.1462](#) (unpublished).
- Newer results from ETM removes hypercubic artifacts [[arXiv:1706.03017](#)] and determines CKM [[arXiv:1706.03657](#)].



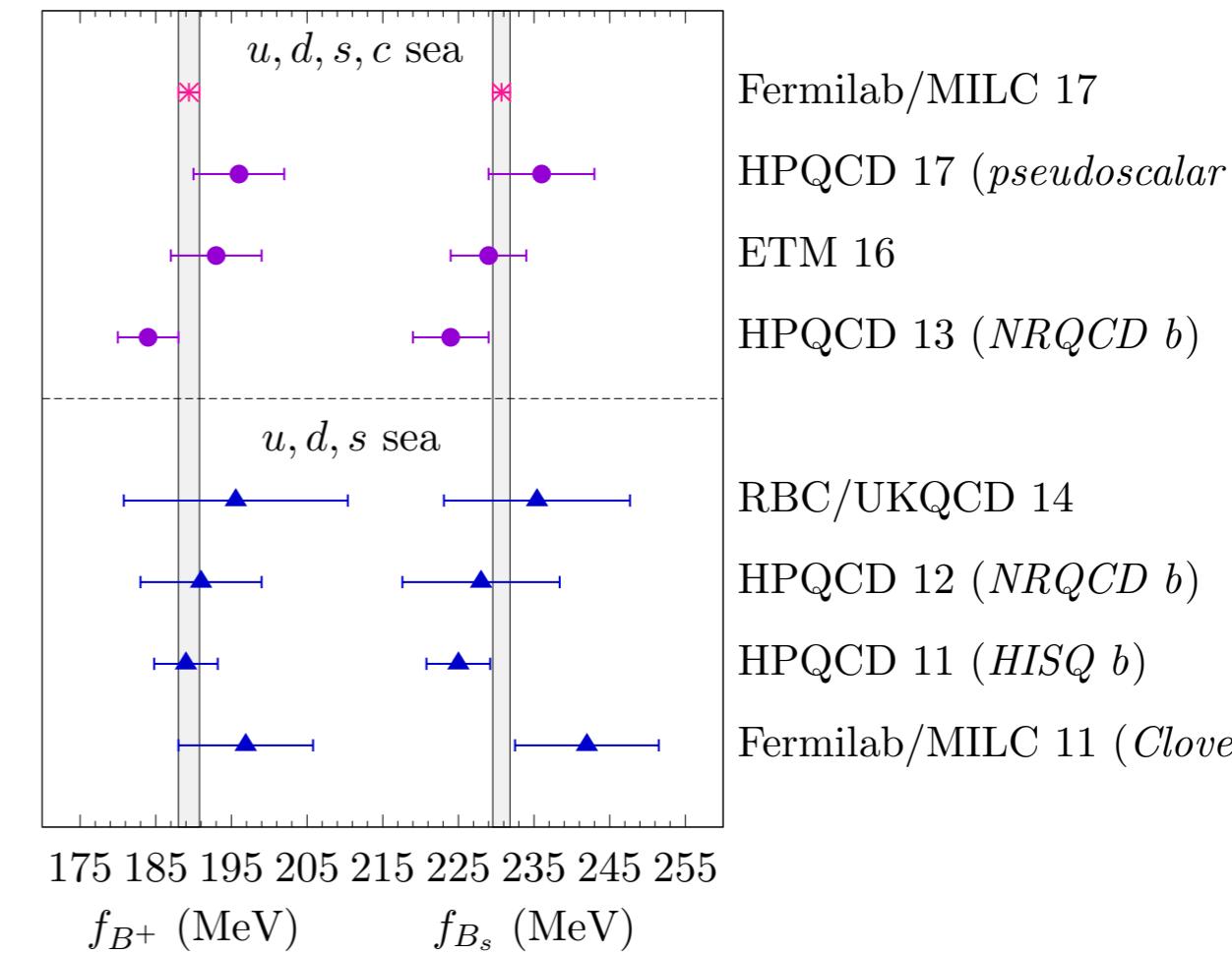
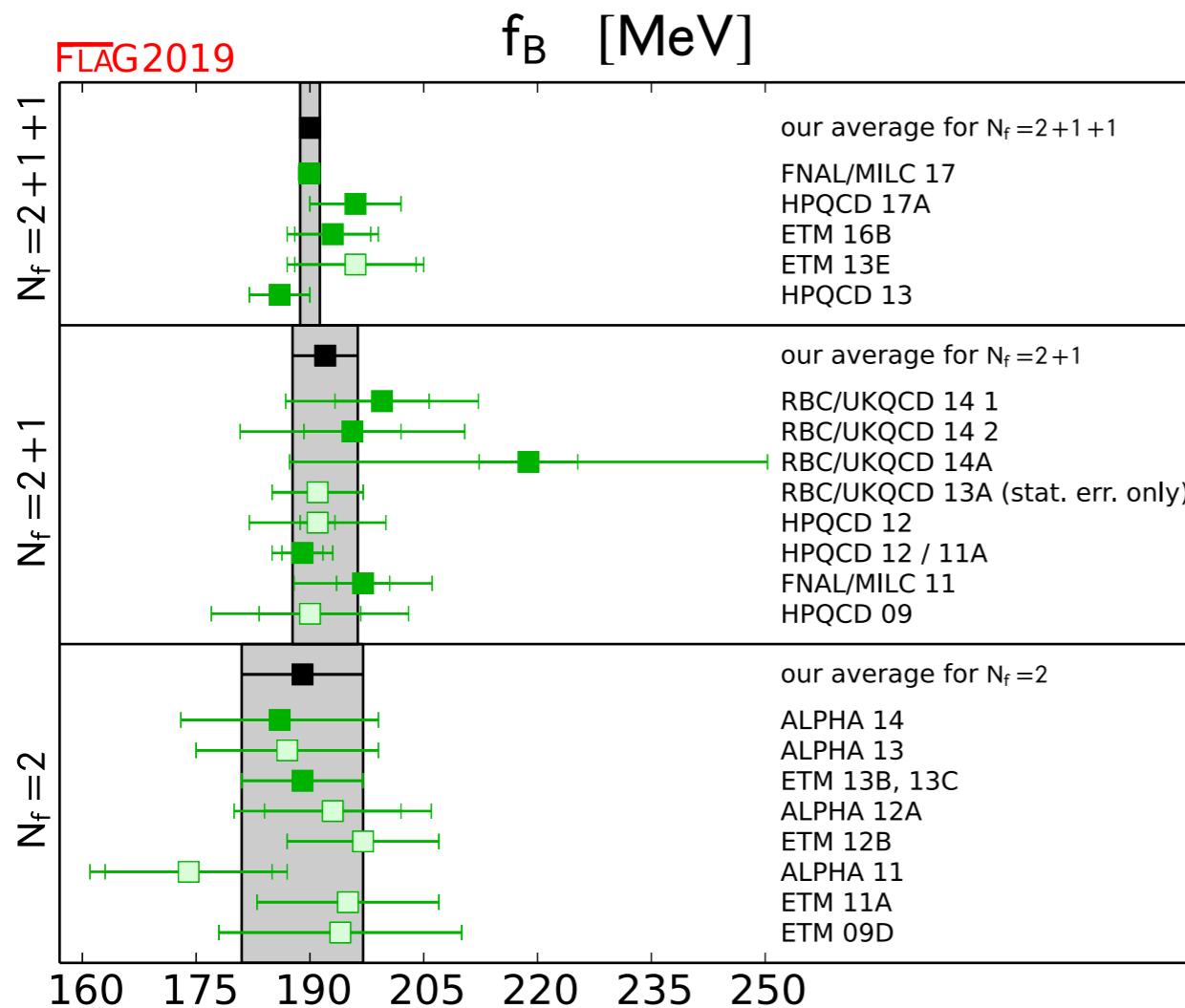
# $|V_{cd}|$ and $|V_{cs}|$ : Leptonic

- Major update from Fermilab Lattice/MILC [arXiv:1712.09262]: ~0.2%
- Also RBC/UKQCD [arXiv:1701.02644].



# $|V_{ub}|$ : Leptonic

- Updates from Fermilab Lattice/MILC [[arXiv:1712.09262](#)]: ~0.7%
- Other updates from HPQCD [[arXiv:1711.09981](#)], ETM [[arXiv:1603.04306](#)].



# $|V_{ub}|$ and $|V_{cb}|$

---

- Older  $B \rightarrow \pi l \nu$  from RBC/UKQCD [[arXiv:1501.05373](#)] and Fermilab Lattice/MILC [[arXiv:1503.07839](#)].
- Baryon decay  $\Lambda_b \rightarrow p l \nu$ : Detmold, Lehner, Meinel [[arXiv:1503.01421](#) and **S. Meinel, Weds. 14:00**].
- $B_s \rightarrow K l \nu$  done [[arXiv:1808.09285](#), [arXiv:1901.02561](#)] or underway.
- Updates to  $B \rightarrow \pi l \nu$  coming [[arXiv:1710.09442](#) and **W. Jay, Tues. 17:20**].
- Aim to keep QCD uncertainties similar to Belle 2 and LHCb uncertainties.
- Explore prospects for, e.g.,  $B \rightarrow K^* ll$  [**S. Meinel, Weds. 14:00**.  
 $\hookrightarrow K\pi$ ]

# Lepton-Flavor Physics

# Muon Anomalous Magnetic Moment

- Fermilab experiment aims to improve precision by a factor of 4.

- By-far dominant theoretical uncertainties stem from hadronic contributions, which can be computed with lattice QCD.

- Hadronic vacuum polarization (HVP); cf., S. Gottlieb, Mon. 17:40.

- Hadronic light-by-light (HLbL):

- first calculation at physics light-quark masses [[arXiv:1610.04603](https://arxiv.org/abs/1610.04603)];
- suggests "Glasgow consensus" is not way off.

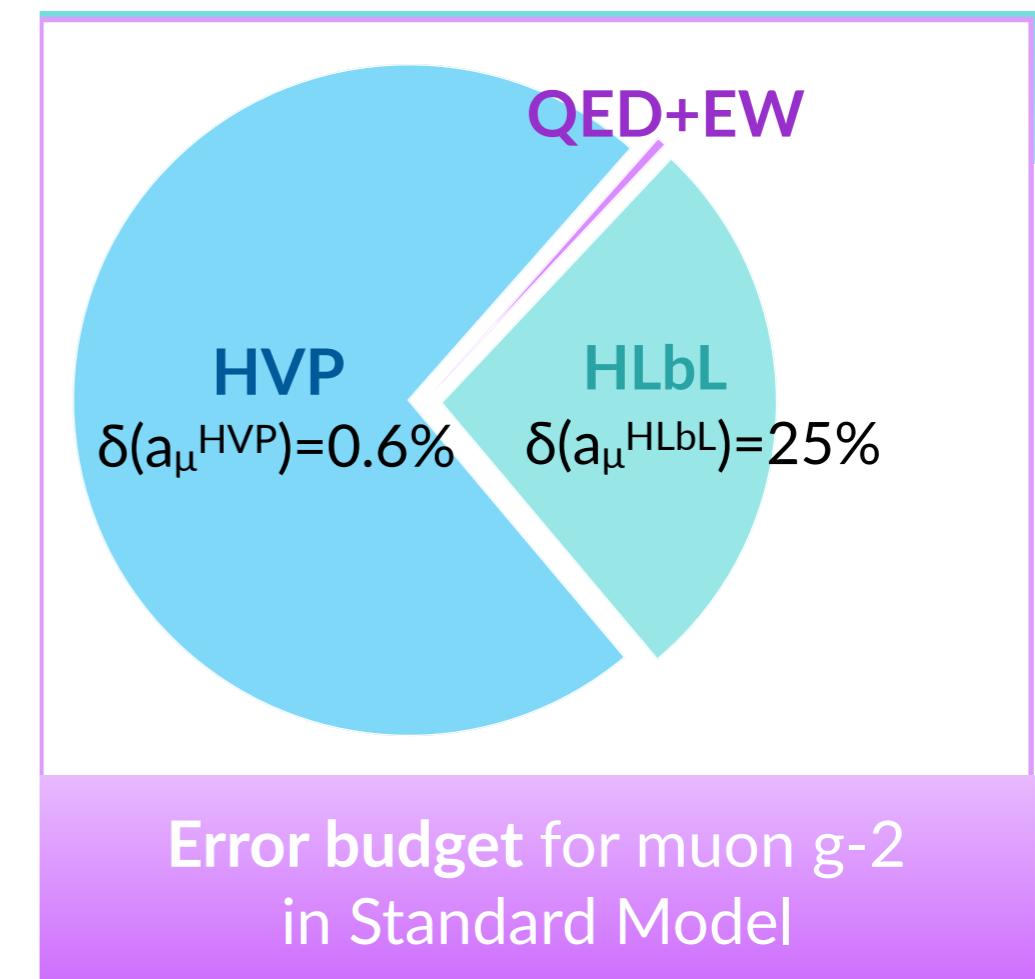


image from Ruth Van de Water

# Neutrino Physics— $\nu N$ Scattering

---

- Key issue in neutrino oscillations is to reconstruct the neutrino energy, but the nuclear remnant (and possibly other particles) are not detected.
- Therefore, nuclear models are needed.
- The models require nucleonic ingredients which should come from QCD:
- Therefore, aim to get scattering amplitudes from first principles:
  - $\bar{\nu}p$  or  $\nu d$  scattering experiments in the Tokai, NuMI, or LBNF beam;
  - lattice QCD with error budgets as comprehensive as those for CKM.

# Axial Form Factor from Lattice QCD

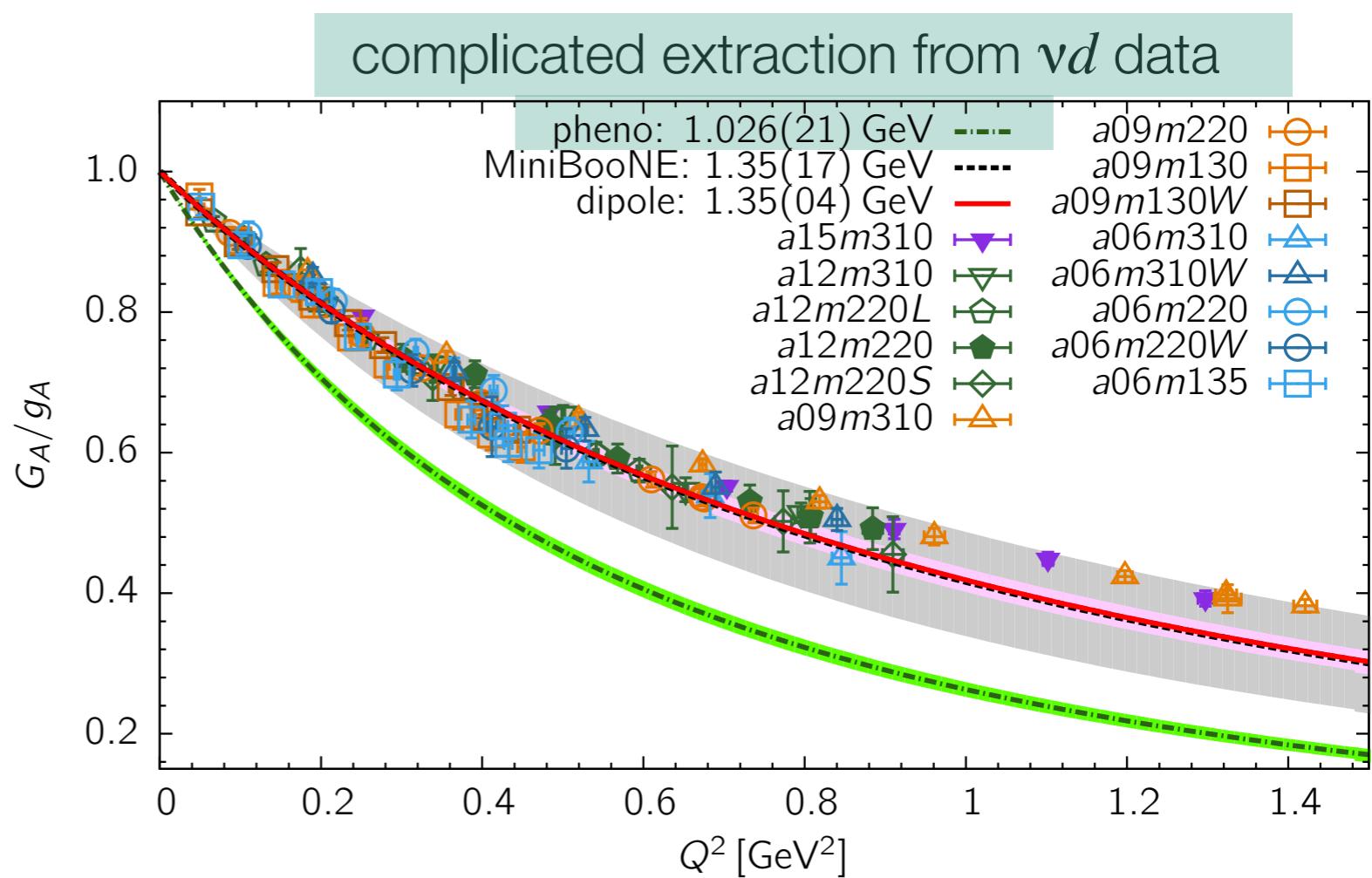
- PNDME c USQCD, [arXiv:1901.00060](#) and [Lattice '19 talk](#):
- Lattice slopes smaller than that extracted from  $\nu d$ .
- Even continuum limit:

$r_A = 0.481(85)$  fm,

vs Meyer *et al.*:

$r_A = 0.68(16)$  fm.

- *Caveat emptor* extraction from  $\nu d$  data.



# Quark Masses



# Four Methods

---

- Mass renormalization [ $\bar{m} = m_{\overline{\text{MS}}}(m_{\overline{\text{MS}}})$ ]:
  - $\bar{m} = Z_m(m_0 - m_{\text{cr}})$ , with  $Z_m$  either  $\leq 2$  loops or **nonperturbative**. 
- Ward identities:
  - $m_{\text{AW}} Z_P \langle P \rangle = Z_A \langle \partial \cdot A \rangle$ , run  $m_{\text{AW}}$  to high scale  $\mu$  & convert to  $m_{\overline{\text{MS}}}(\mu)$ .
- Continuum limit  $\otimes$  continuum pQCD:
  - $\lim_{a \rightarrow 0} G_n(Qa) = 1 + \sum_{k=1}^{K \geq 3} G_n^{(k)} \alpha_s^k(Q)$ , e.g., quarkonium moments. 
- Continuum limit  $\otimes$  HQET:
  - $M = m_{\text{MRS}} + \bar{\Lambda}_{\text{MRS}} + \dots$ , e.g., heavy-light meson masses. 

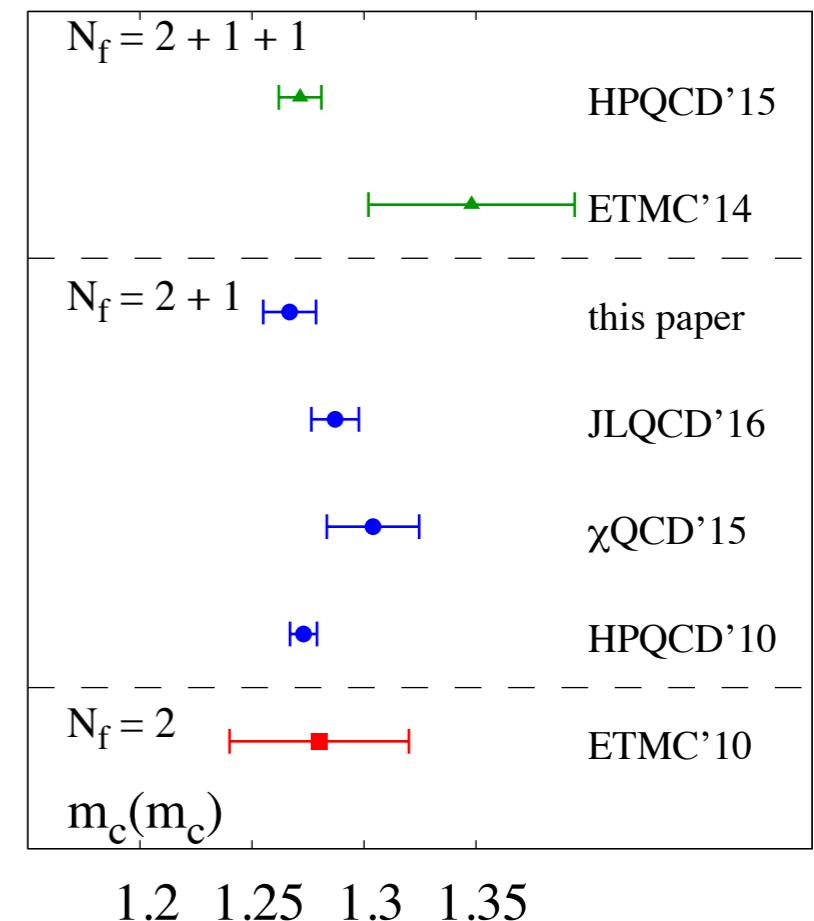
# Charmonium Correlators

- Lately, the most precise determinations of heavy-quarks come from:

$$\lim_{a \rightarrow 0} \frac{G_n}{G_n^{\text{tree}}} = g_n(\alpha_s(\mu), m_{c,\overline{\text{MS}}}/\mu)$$
$$G_n = \sum_t t^n m_0^2 \left\langle \bar{c} \gamma^5 c(t) \bar{c} \gamma^5 c(0) \right\rangle, \quad \begin{matrix} \text{even} \\ n \geq 4 \end{matrix}$$

Bochkarev & de Forcrand [[hep-lat/9505025](#)],  
Allison et al. [HPQCD, [arXiv:0805.2999](#)].

- New analysis on 11 ensembles from hotQCD ( $n_f = 2+1$  HISQ) by Maezawa & Petreczky [[arXiv:1606.08798](#)] compatible with earlier HPQCD results [[arXiv:1408.4169](#)].



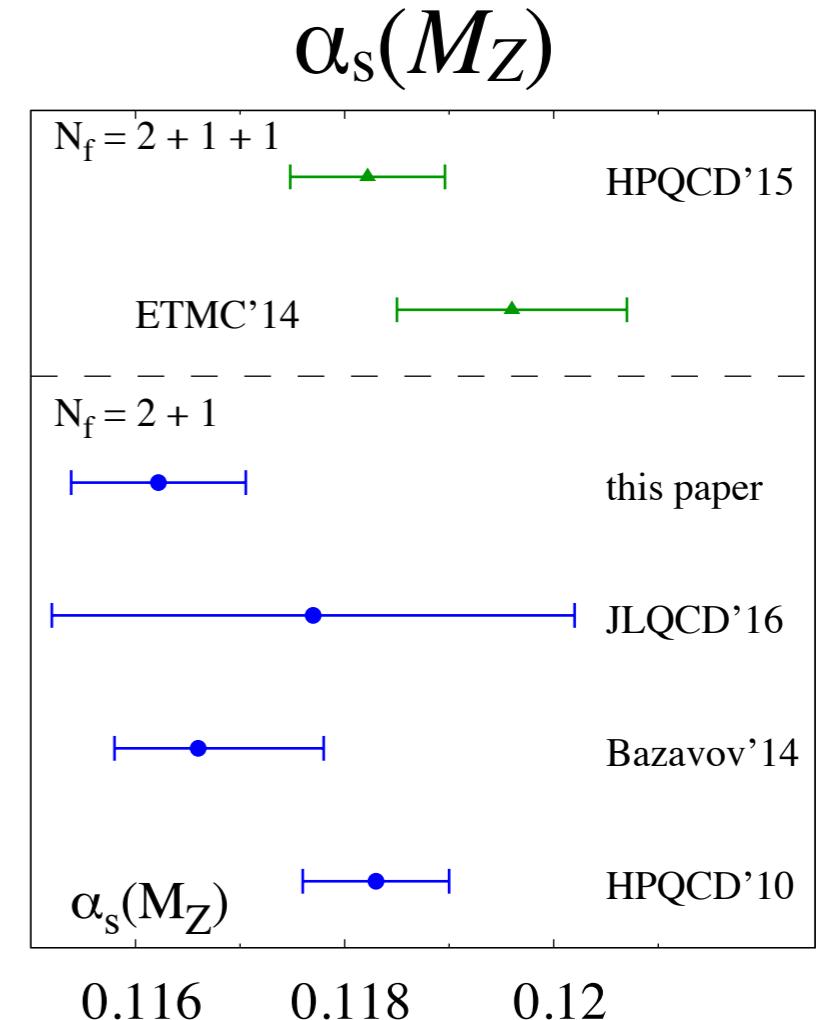
# Charmonium Correlators

- Lately, the most precise determinations of heavy-quarks come from:

$$\lim_{a \rightarrow 0} \frac{G_n}{G_n^{\text{tree}}} = g_n(\alpha_s(\mu), m_{c,\overline{\text{MS}}}/\mu)$$
$$G_n = \sum_t t^n m_0^2 \left\langle \bar{c} \gamma^5 c(t) \bar{c} \gamma^5 c(0) \right\rangle, \quad \begin{matrix} \text{even} \\ n \geq 4 \end{matrix}$$

Bochkarev & de Forcrand [[hep-lat/9505025](#)],  
Allison et al. [HPQCD, [arXiv:0805.2999](#)].

- New analysis on 11 ensembles from hotQCD ( $n_f = 2+1$  HISQ) by Maezawa & Petreczky [[arXiv:1606.08798](#)] compatible with earlier HPQCD results [[arXiv:1408.4169](#)].



# RI-SMOM Method

arXiv:1805.06225

---

- Fix gauge and use quark vertex function:

$$\Lambda_P = \langle S^{-1}(p_1) \rangle \langle \bar{\psi} \gamma^5 \psi(p_1, p_2) \rangle \langle S^{-1}(p_2) \rangle$$

$$|p_1 + p_2| = |p_1| = |p_2| \equiv \mu$$

$$Z_P^{\text{SMOM}}(\mu) \equiv \frac{\Lambda_P(\mu)}{\Lambda_P^{(0)}}$$

$$Z_m^{\text{SMOM}}(\mu) = 1/Z_P^{\text{SMOM}}(\mu)$$

- Trickiest part of the analysis is that this gauge-fixed quantity has “gluon-mass condensate” contribution  $\sim \langle A^2 \rangle / \mu^2$ .
- Run  $m^{\text{SMOM}}(\mu) = Z_m^{\text{SMOM}}(\mu)m_0$  to a high scale and convert to  $m_{\overline{\text{MS}}}(\mu)$ .

# Heavy-light Meson Masses

---

- From HQET (or other approaches to the  $1/m_h$  expansion):

$$M_{H_J} = m_h + \bar{\Lambda} + \frac{\mu_\pi^2}{2m_h} - d_J \frac{\mu_G^2(m_h)}{2m_h}$$

- This formula used to determine  $\bar{\Lambda}$ ,  $\mu_\pi^2$ , and  $\mu_G^2(m_b)$  from lattice QCD in [hep-ph/0006345](#); new, improved theory [[arXiv:1712.04983](#)].

# Heavy-light Meson Masses

- From HQET (or other approaches to the  $1/m_h$  expansion):

$$M_{HJ} = m_h + \bar{\Lambda} + \frac{\mu_\pi^2}{2m_h} - d_J \frac{\mu_G^2(m_h)}{2m_h}$$

mass of spin- $J$  meson

mass of heavy quark

energy of gluons and light quarks

kinetic energy of heavy quark

1 for  $B$ ,  $-\frac{1}{3}$  for  $B^*$

spin-orbit interaction

- This formula used to determine  $\bar{\Lambda}$ ,  $\mu_\pi^2$ , and  $\mu_G^2(m_b)$  from lattice QCD in [hep-ph/0006345](#); new, improved theory [[arXiv:1712.04983](#)].

# Minimal Renormalon Subtraction

[arXiv:1701.00347](https://arxiv.org/abs/1701.00347), [arXiv:1712.04983](https://arxiv.org/abs/1712.04983)

---

- New idea how to isolate the leading  $\mathcal{O}(\Lambda)$  ambiguity in the pole mass, yielding an unambiguous definition of  $\bar{\Lambda}$ .
- For the bottom quark (tree, 1-loop, 2-loop, 3-loop, 4-loop):

$$m_{b,\text{pole}}/\bar{m}_b = (1, 1.093, 1.143, 1.183, 1.224)$$

$$m_{b,\text{MRS}}/\bar{m}_b = (1.157, 1.133, 1.131, 1.132, 1.132)$$

- For the top quark (tree, 1-loop, 2-loop, 3-loop, 4-loop):

$$m_{t,\text{MRS}}/\bar{m}_t = (1.0687, 1.0576, 1.0573, 1.0574, 1.0574)$$

could have implications for top physics at LHC.

# Results from MRS/HQET

$\alpha_s$  parametric not  
PT truncation

- Masses in numerical form:

$$m_{l,\overline{\text{MS}}}(2 \text{ GeV}) = 3.404(14)_{\text{stat}}(08)_{\text{syst}}(19)\alpha_s(04)_{f_\pi,\text{PDG}} \text{ MeV}$$

$$m_{u,\overline{\text{MS}}}(2 \text{ GeV}) = 2.118(17)_{\text{stat}}(32)_{\text{syst}}(12)\alpha_s(03)_{f_\pi,\text{PDG}} \text{ MeV}$$

$$m_{d,\overline{\text{MS}}}(2 \text{ GeV}) = 4.690(30)_{\text{stat}}(36)_{\text{syst}}(26)\alpha_s(06)_{f_\pi,\text{PDG}} \text{ MeV}$$

$$m_{s,\overline{\text{MS}}}(2 \text{ GeV}) = 92.52(40)_{\text{stat}}(18)_{\text{syst}}(52)\alpha_s(12)_{f_\pi,\text{PDG}} \text{ MeV}$$

$$m_{c,\overline{\text{MS}}}(3 \text{ GeV}) = 984.3(4.2)_{\text{stat}}(1.6)_{\text{syst}}(3.2)\alpha_s(0.6)_{f_\pi,\text{PDG}} \text{ MeV}$$

$$m_{b,\overline{\text{MS}}}(m_{b,\overline{\text{MS}}}) = 4203(12)_{\text{stat}}(1)_{\text{syst}}(8)\alpha_s(1)_{f_\pi,\text{PDG}} \text{ MeV}$$

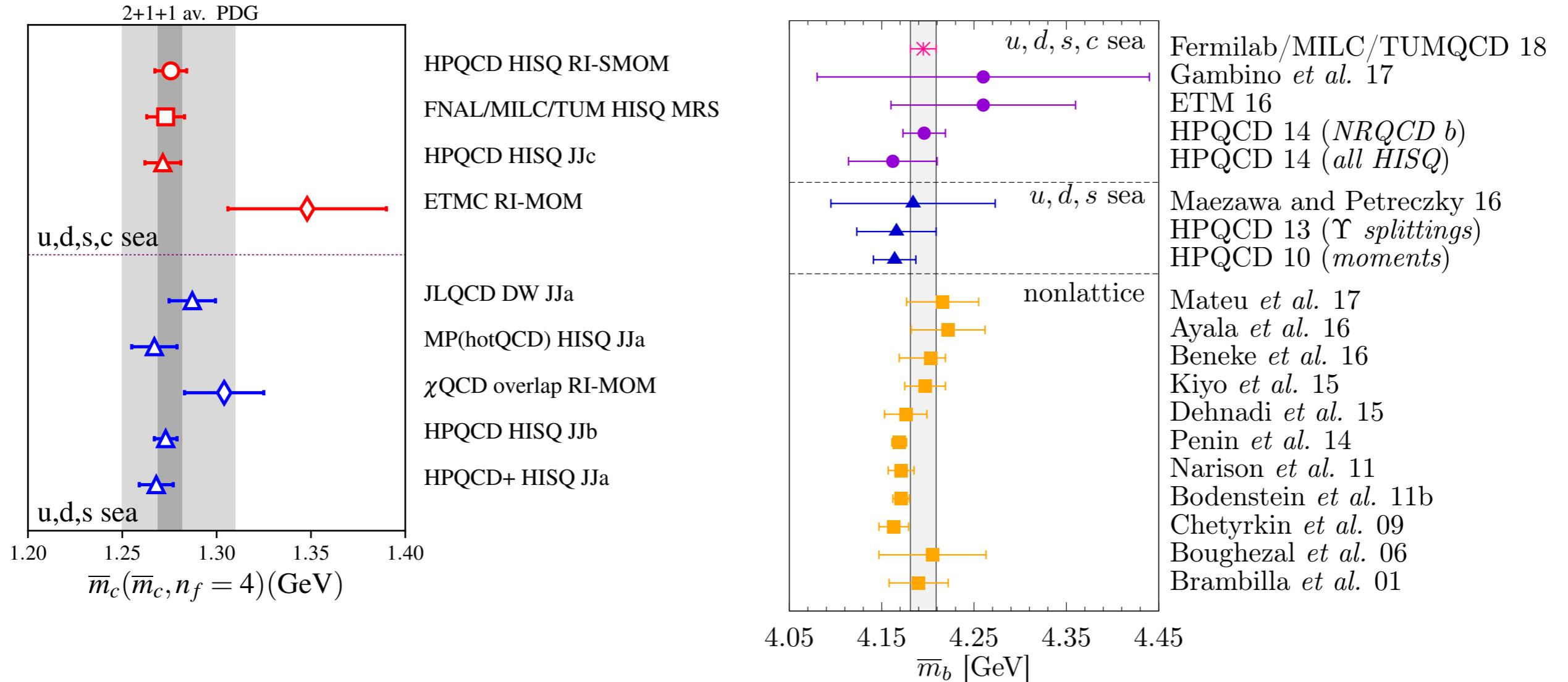
- Mass ratios:

$$m_c/m_s = 11.784(11)_{\text{stat}}(17)_{\text{syst}}(00)\alpha_s(08)_{f_\pi,\text{PDG}}$$

$$m_b/m_s = 53.93(7)_{\text{stat}}(8)_{\text{syst}}(1)\alpha_s(5)_{f_\pi,\text{PDG}}$$

$$m_b/m_c = 4.577(5)_{\text{stat}}(7)_{\text{syst}}(0)\alpha_s(1)_{f_\pi,\text{PDG}}$$

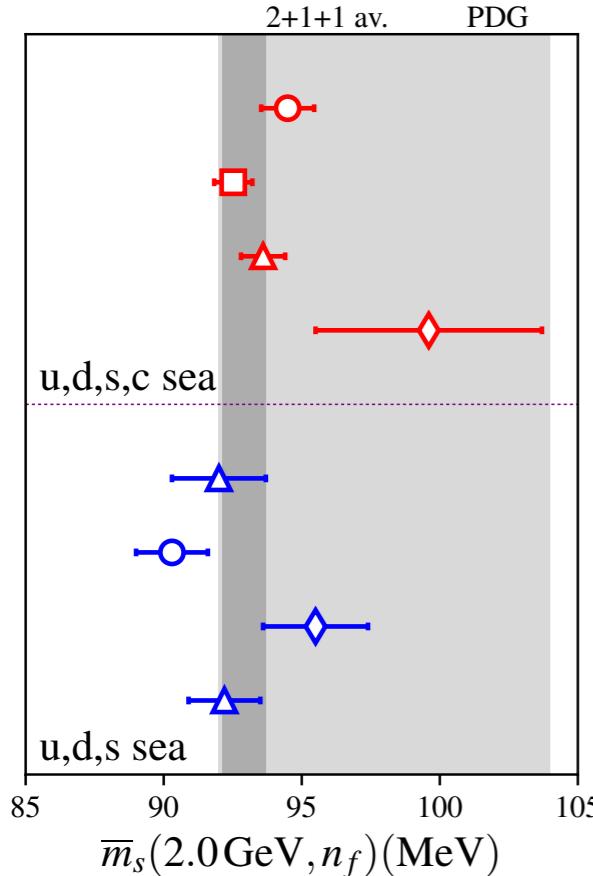
# Heavy Comparisons



- Precision: 0.3% for bottom to 0.5% for charm.

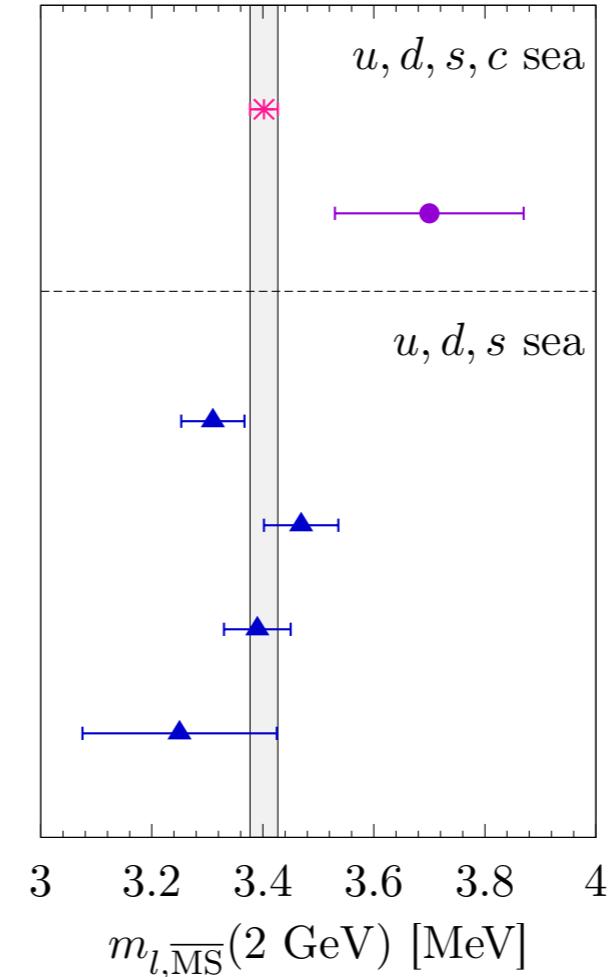
plots from [arXiv:1802.04248](https://arxiv.org/abs/1802.04248), [arXiv:1805.06225](https://arxiv.org/abs/1805.06225)

# Light Comparisons



2+1+1 av.      PDG

HPQCD HISQ RI-SMOM  
FNAL/MILC/TUM HISQ MRS +  $m_h/m_s$   
HPQCD HISQ JJ +  $m_c/m_s$   
ETMC RI-MOM  
  
u,d,s,c sea  
  
MP(hotQCD) HISQ JJ +  $m_c/m_s$   
RBC/UKQCD DW RI-SMOM  
BMW clover RI-MOM  
HPQCD HISQ JJ +  $m_c/m_s$   
  
u,d,s sea



$u, d, s, c$  sea  
  
 $u, d, s$  sea  
  
Fermilab/MILC/TUMQCD 18  
ETM 14  
RBC/UKQCD 14  
BMW 10  
HPQCD 10  
MILC 09

plots from [arXiv:1802.04248](https://arxiv.org/abs/1802.04248)  
[arXiv:1805.06225](https://arxiv.org/abs/1805.06225)

- Precision: 2% for up quark.

Consistent picture: all quarks but top

# Higgs Physics

---

- Quark masses and  $\alpha_s$  relevant for Higgs physics.
- Anything else?
- Moments of parton distribution functions (PDFs):
  - tame power-law divergences with gradient flow [[arXiv:1612.01584](#)].
- Dependence on Bjorken  $x$ :
  - numerous techniques (quasi-pdf, pseudo-pdf, "lattice cross sections", Ioffe-time matrix elements, ...) stemming from Ji [[arXiv:1305.1539](#)].
- Community whitepaper [[arXiv:1711.07916](#)]  $\Leftarrow$  lattice-pheno synergy.

# Outlook

- Lattice QCD now has had huge impact in quark-flavor physics:
  - results are increasingly precise: sub-% threshold has been passed for  $f_{B(s)}$ ,  $f_{D(s)}$ ,  $m_b$ ,  $m_c$  (and good enough for Belle 2, BES III);
  - quark masses already more than good enough for ILC Higgs studies.
- Excellent consistency of different quark-mass methods suggests that PDG estimates for quark masses must be updated.
- Lattice QCD is on track to reduce hadronic uncertainties in muon  $g-2$ .
- Bright future for nucleon matrix elements (form factors for neutrino physics, PDFs for LHC, many others for the precision frontier).

Thank you!