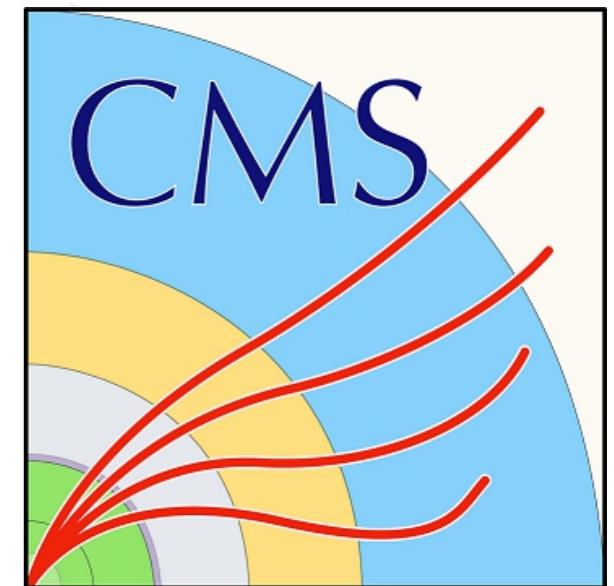


BSM Results at the LHC

APS DPF 2019

Maximilian Swiatlowski,
for the ATLAS and CMS Collaborations

Enrico Fermi Institute, University of Chicago



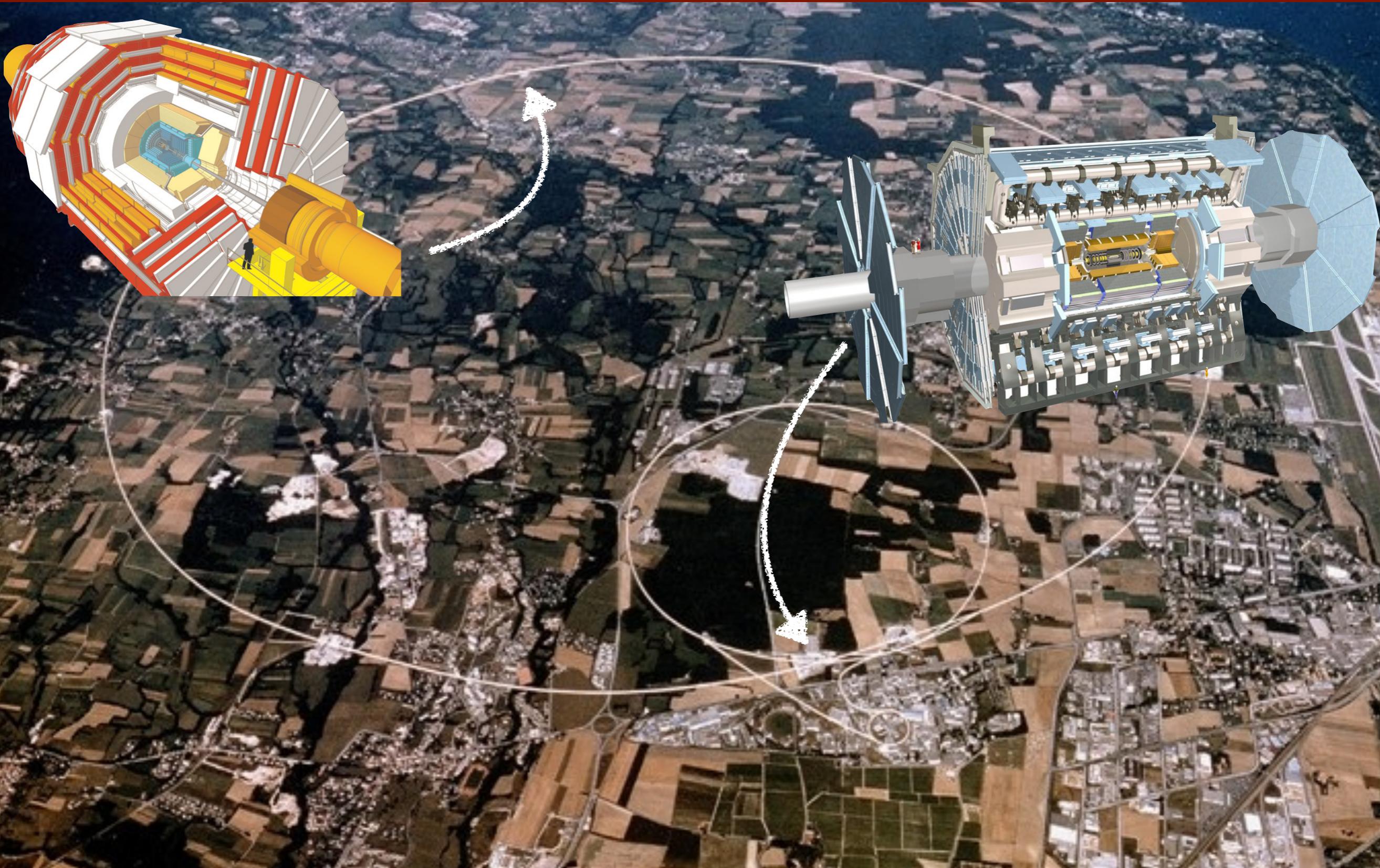


The LHC



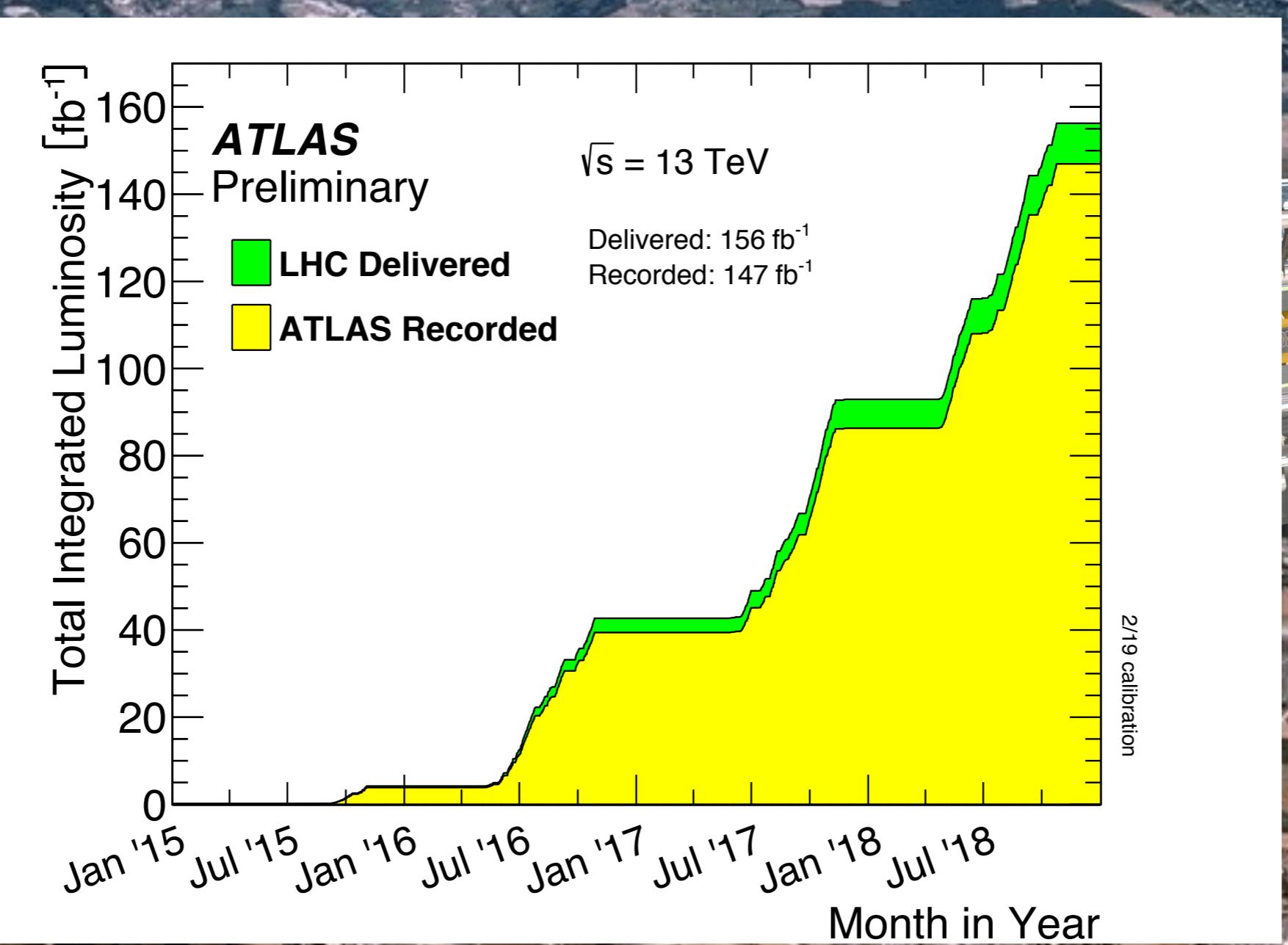


The LHC

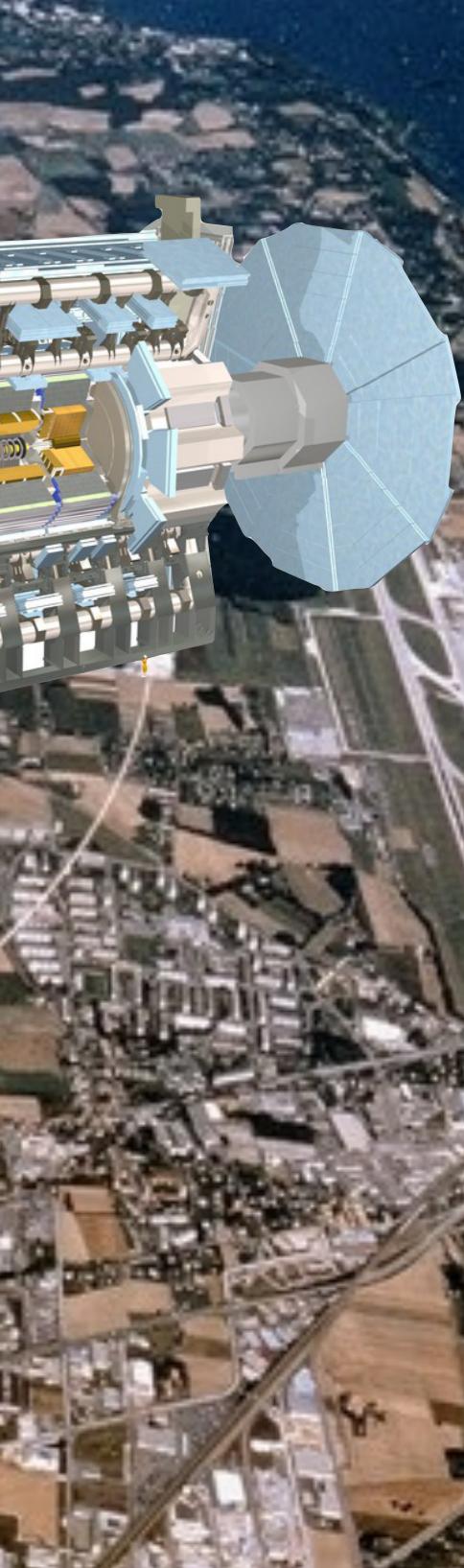
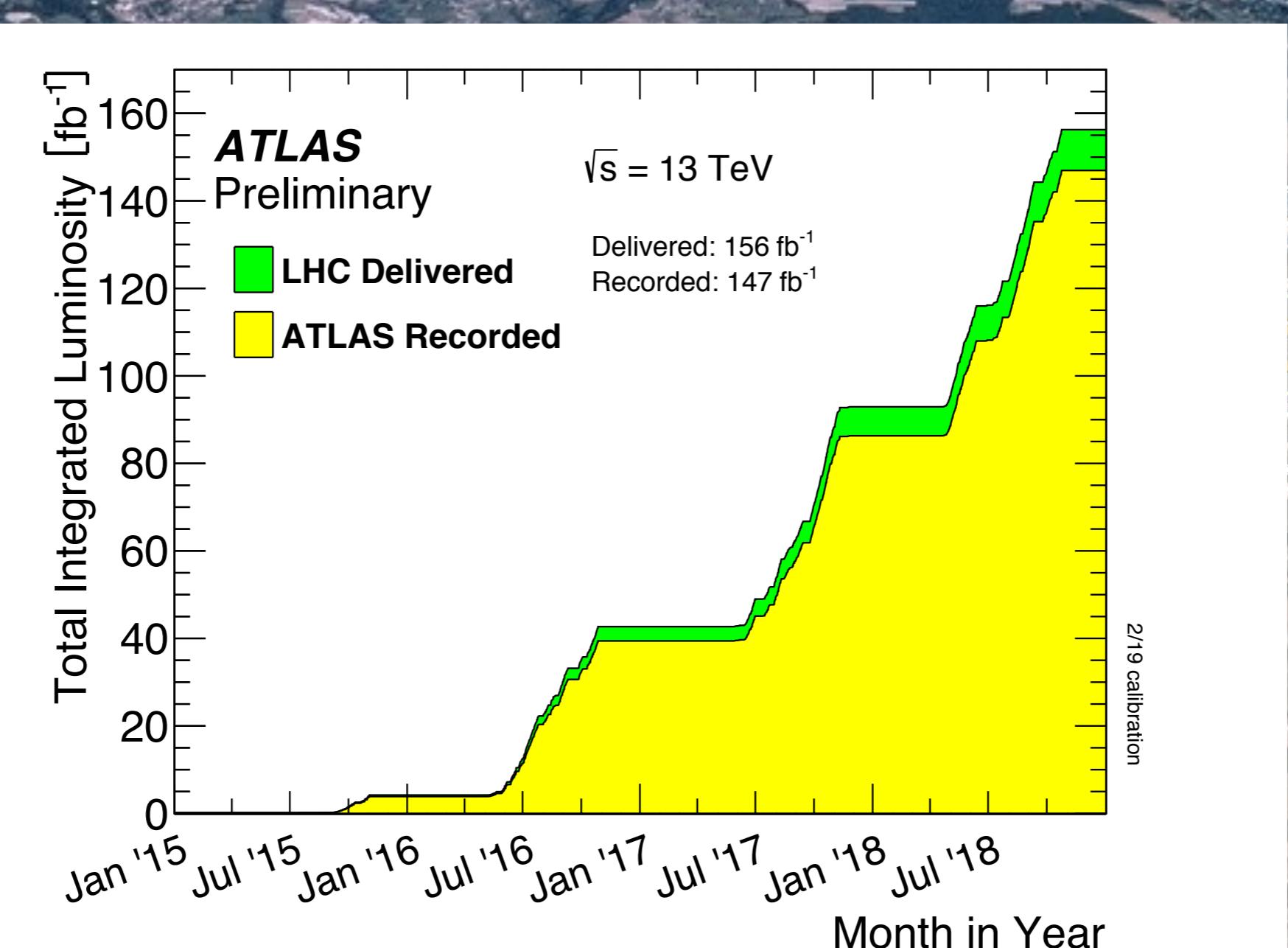




The LHC



The LHC



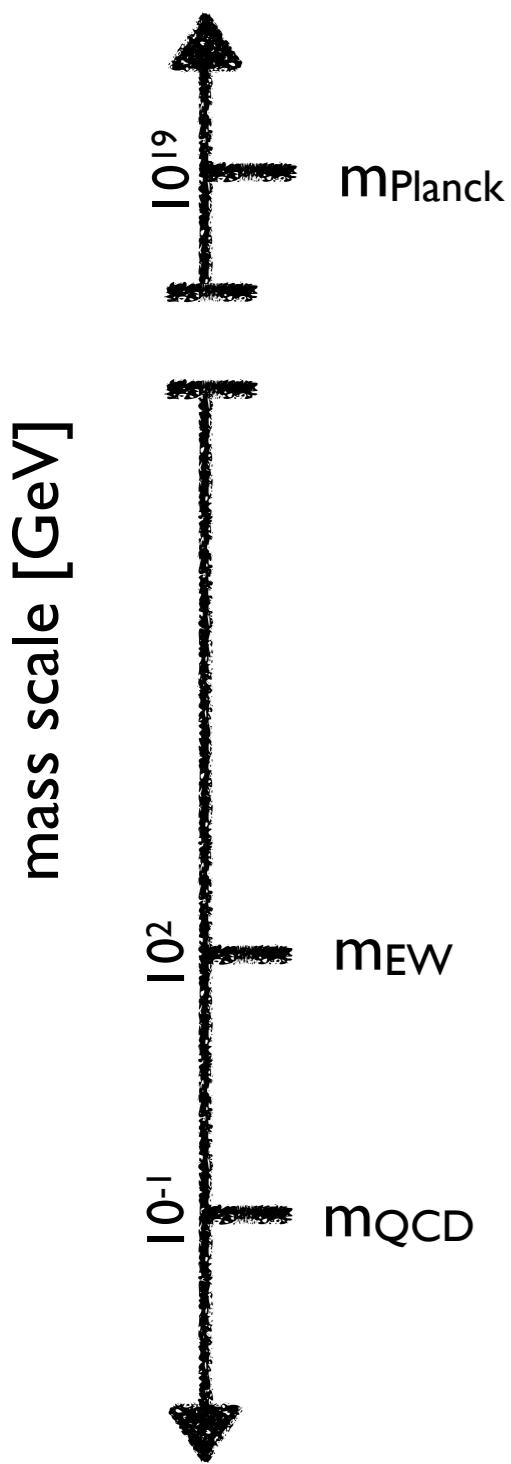
Huge recorded dataset from Run 2:
Thank you to the LHC team and machine!

Why BSM?



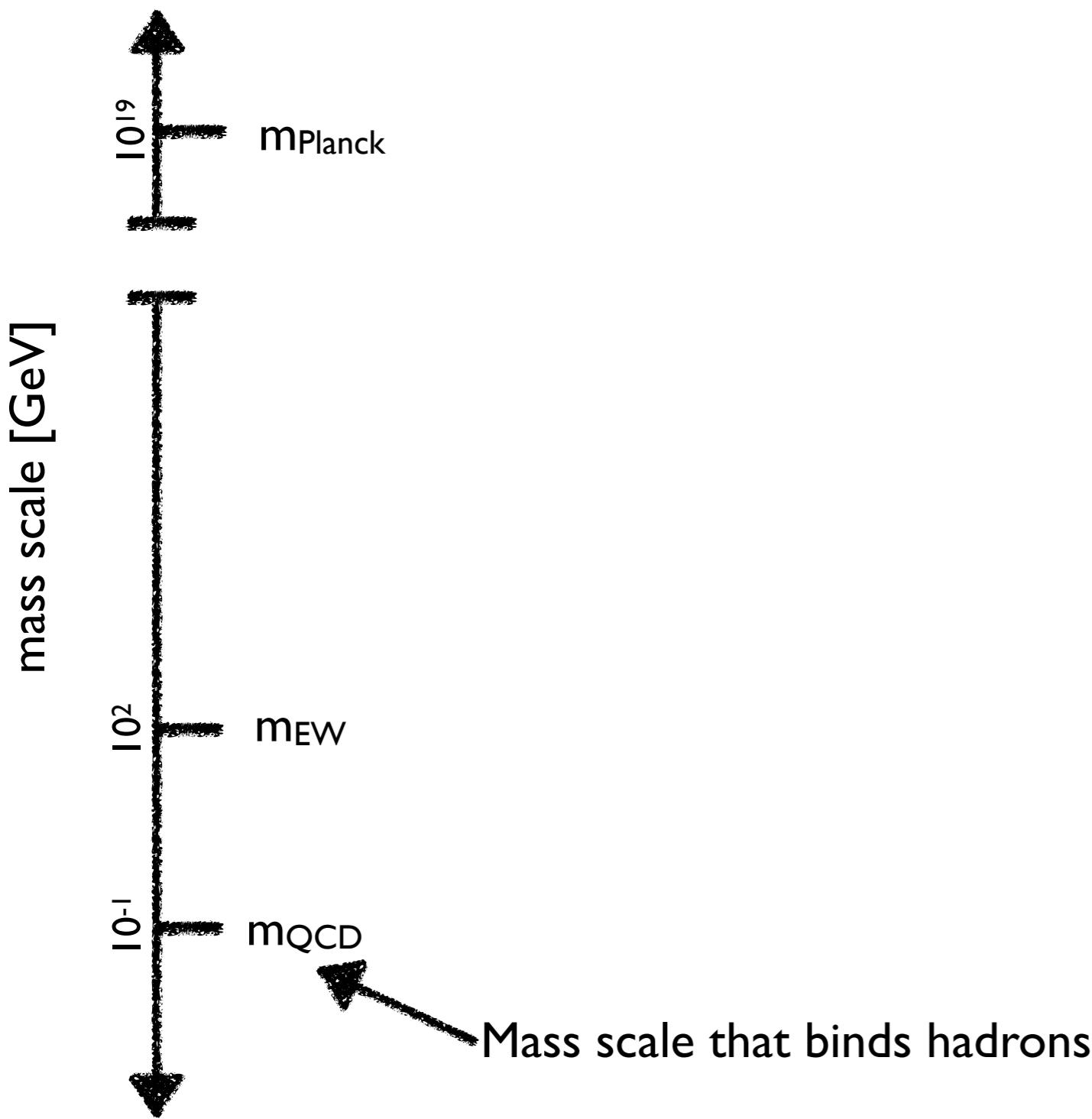


Why BSM?



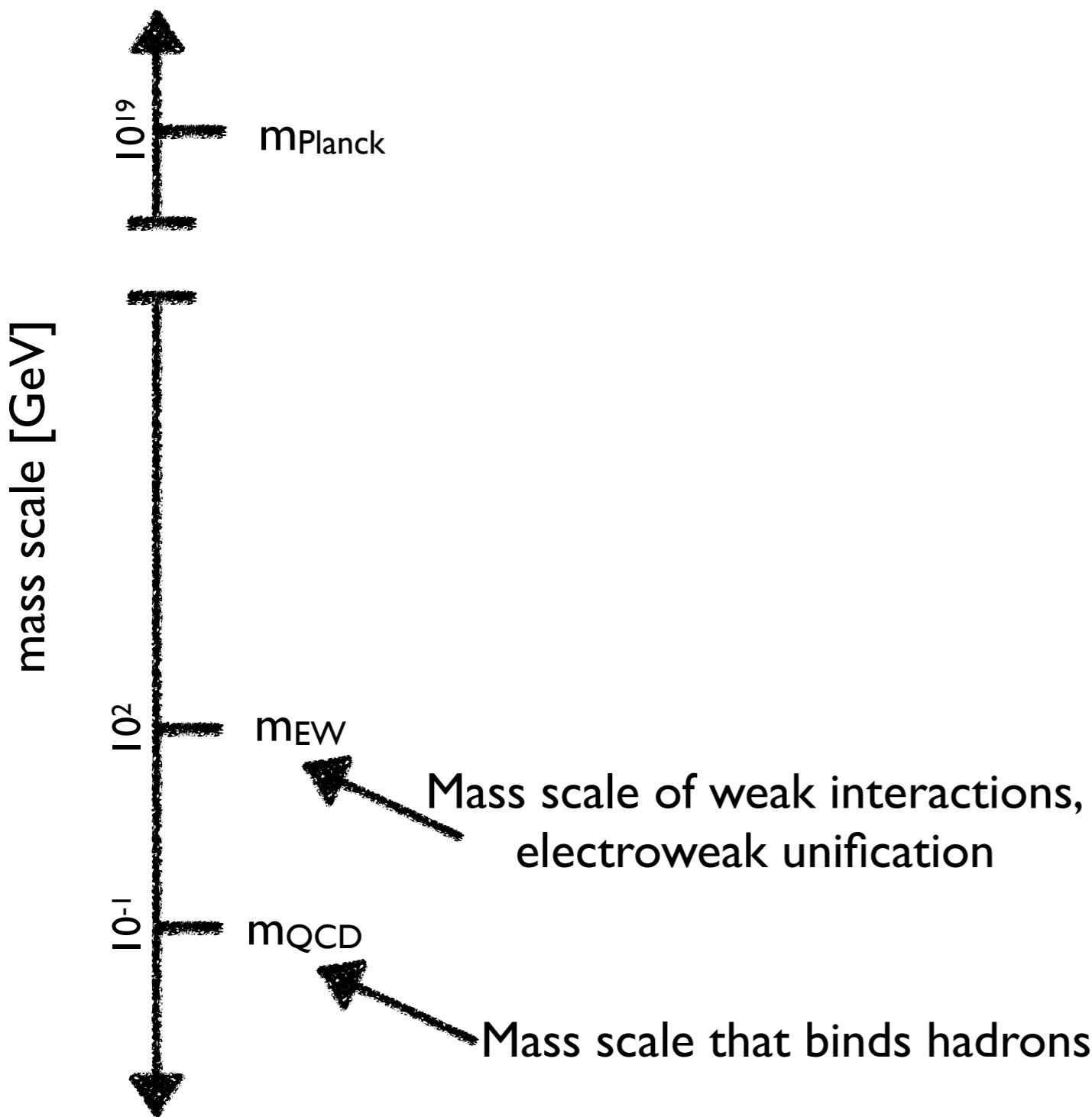


Why BSM?



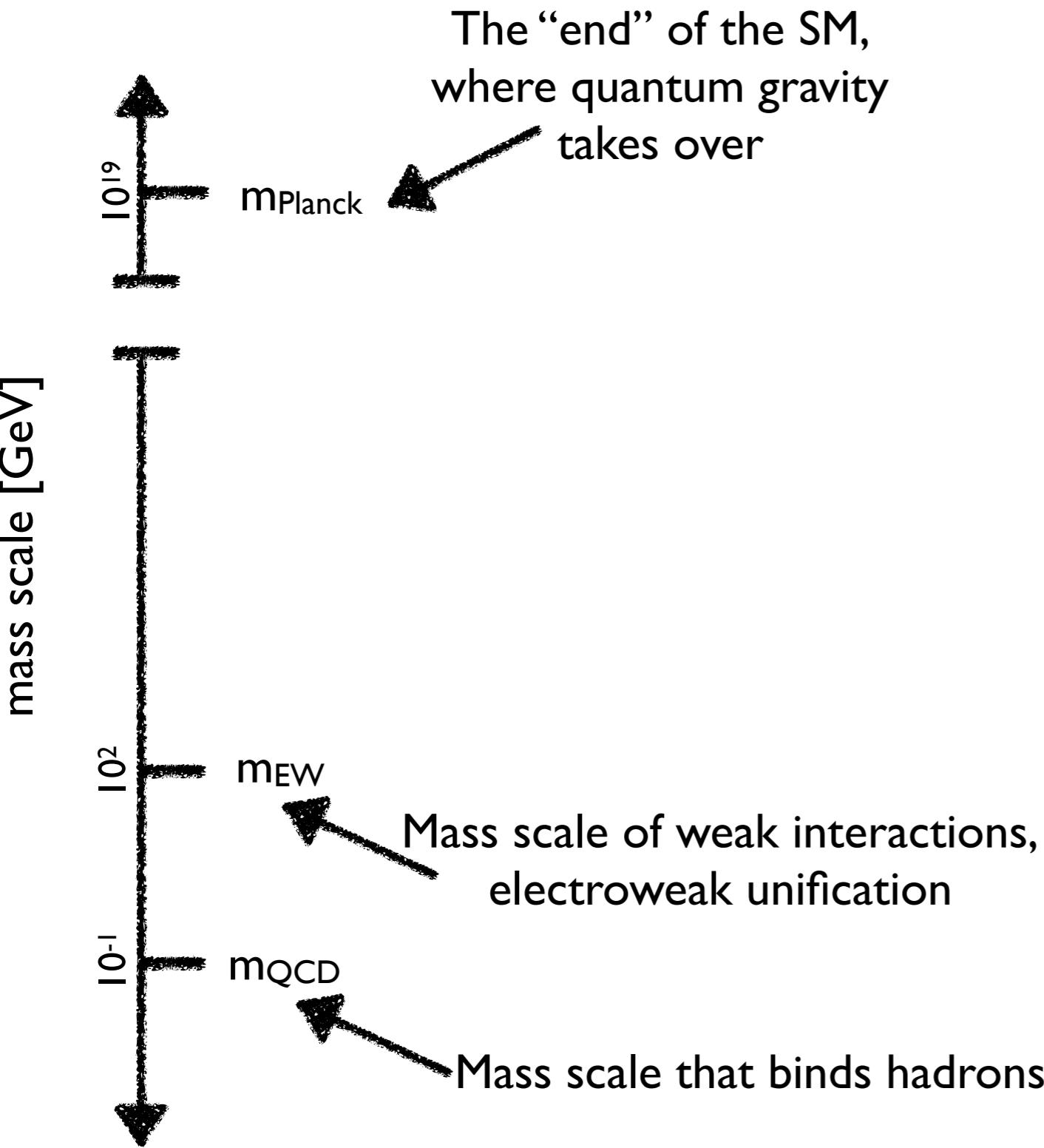


Why BSM?



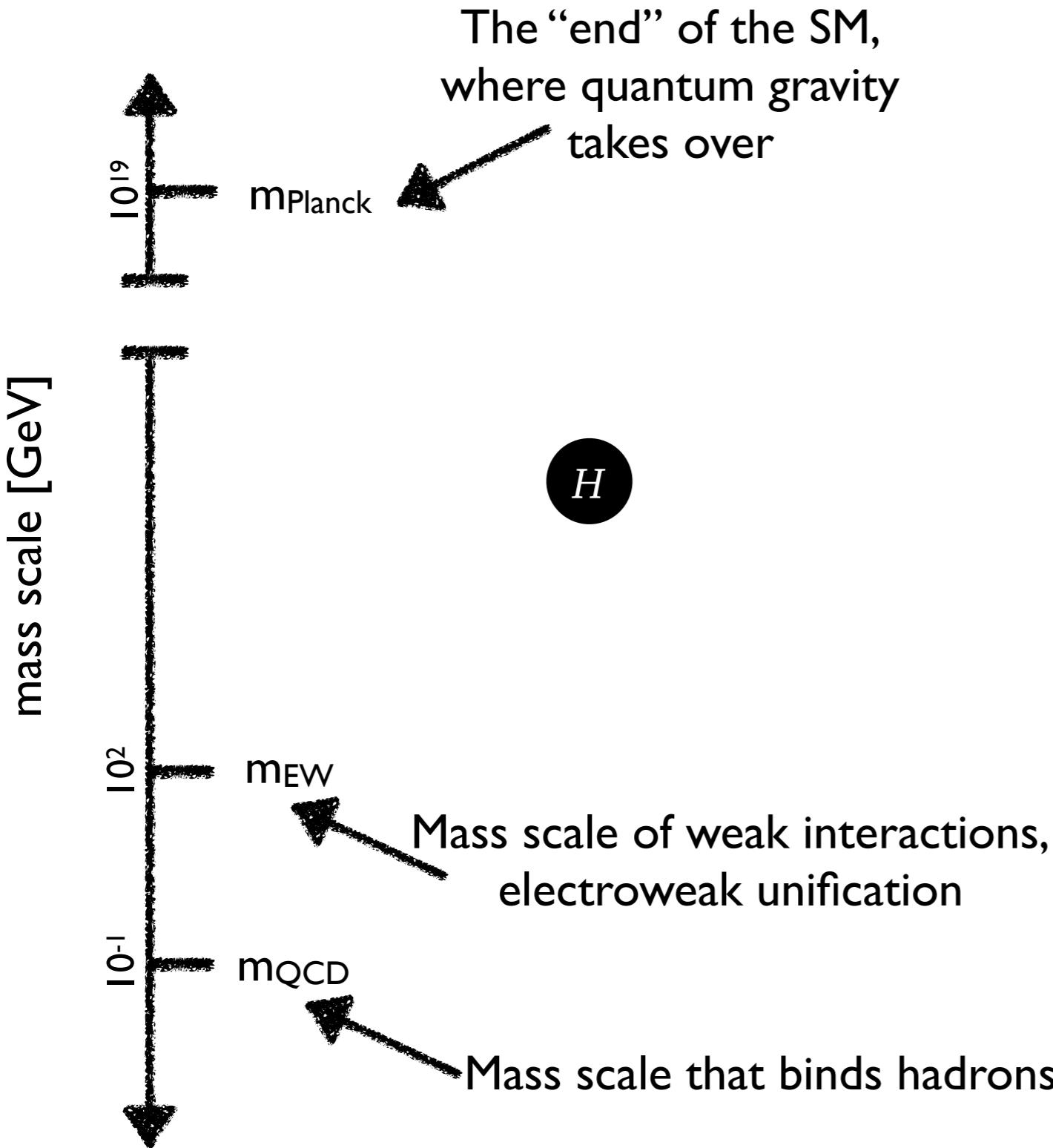


Why BSM?





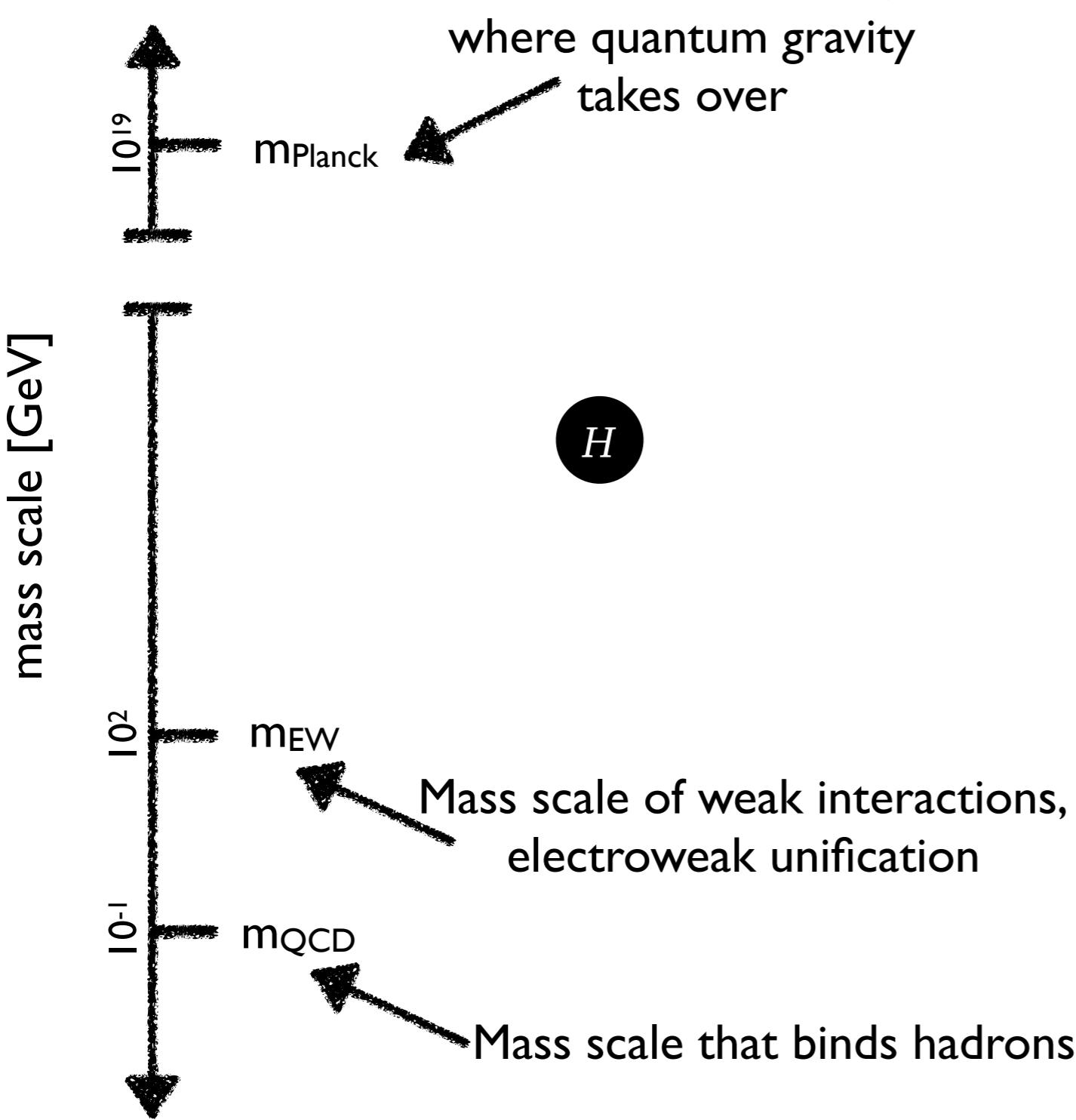
Why BSM?



To first order, the mass of the Higgs is a free parameter, and can be anywhere



Why BSM?

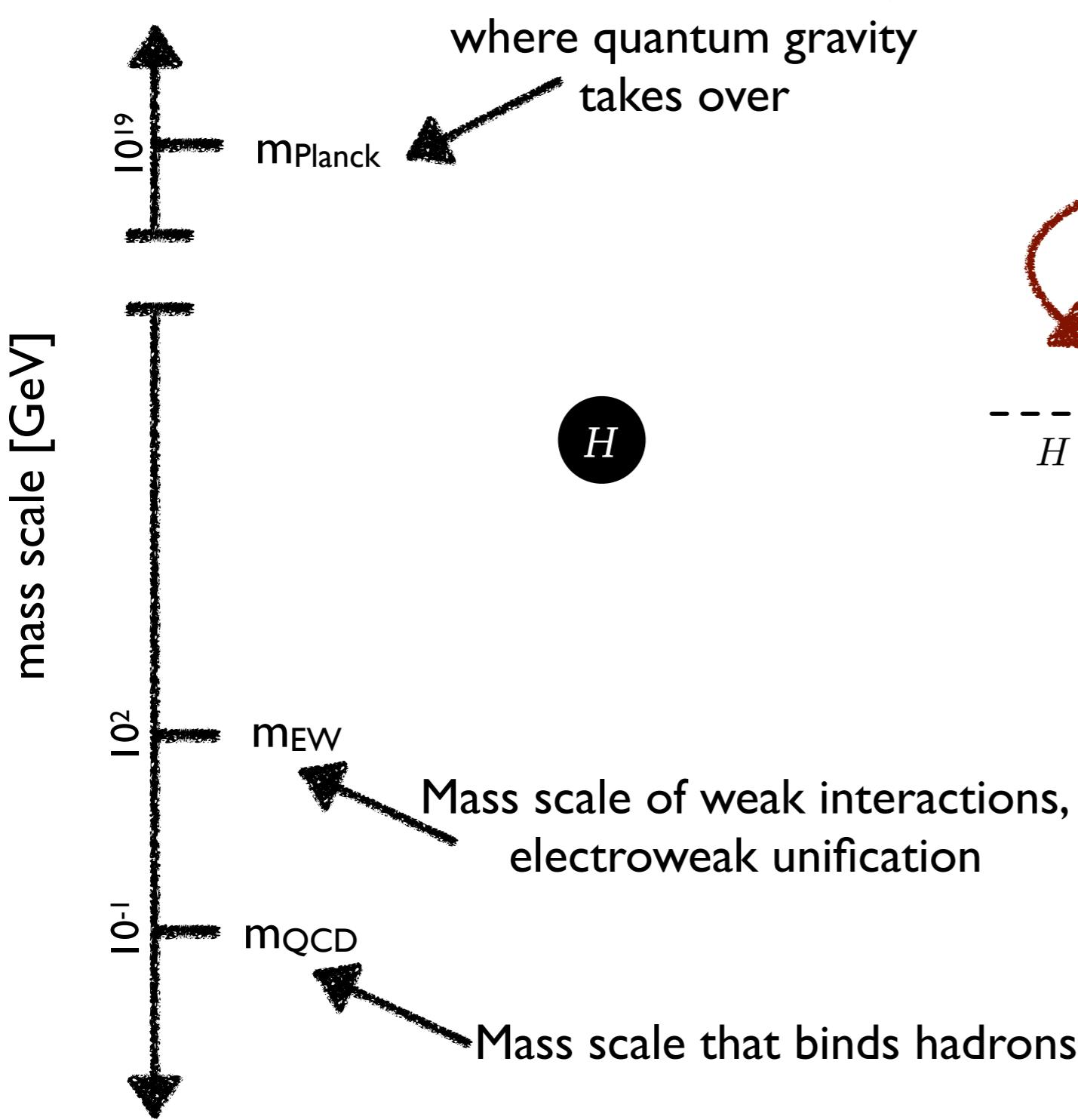


To first order, the mass of the Higgs is a free parameter, and can be anywhere

But the Higgs interacts with other particles, which affects its mass

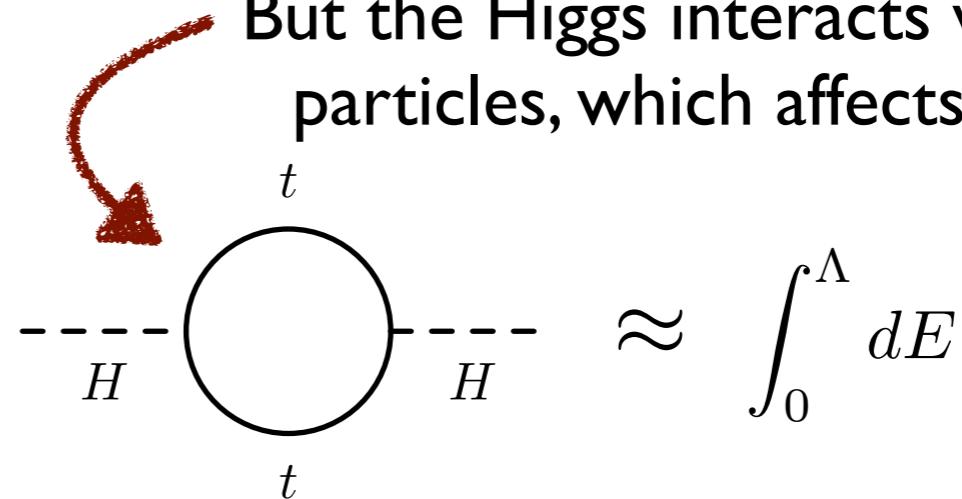


Why BSM?



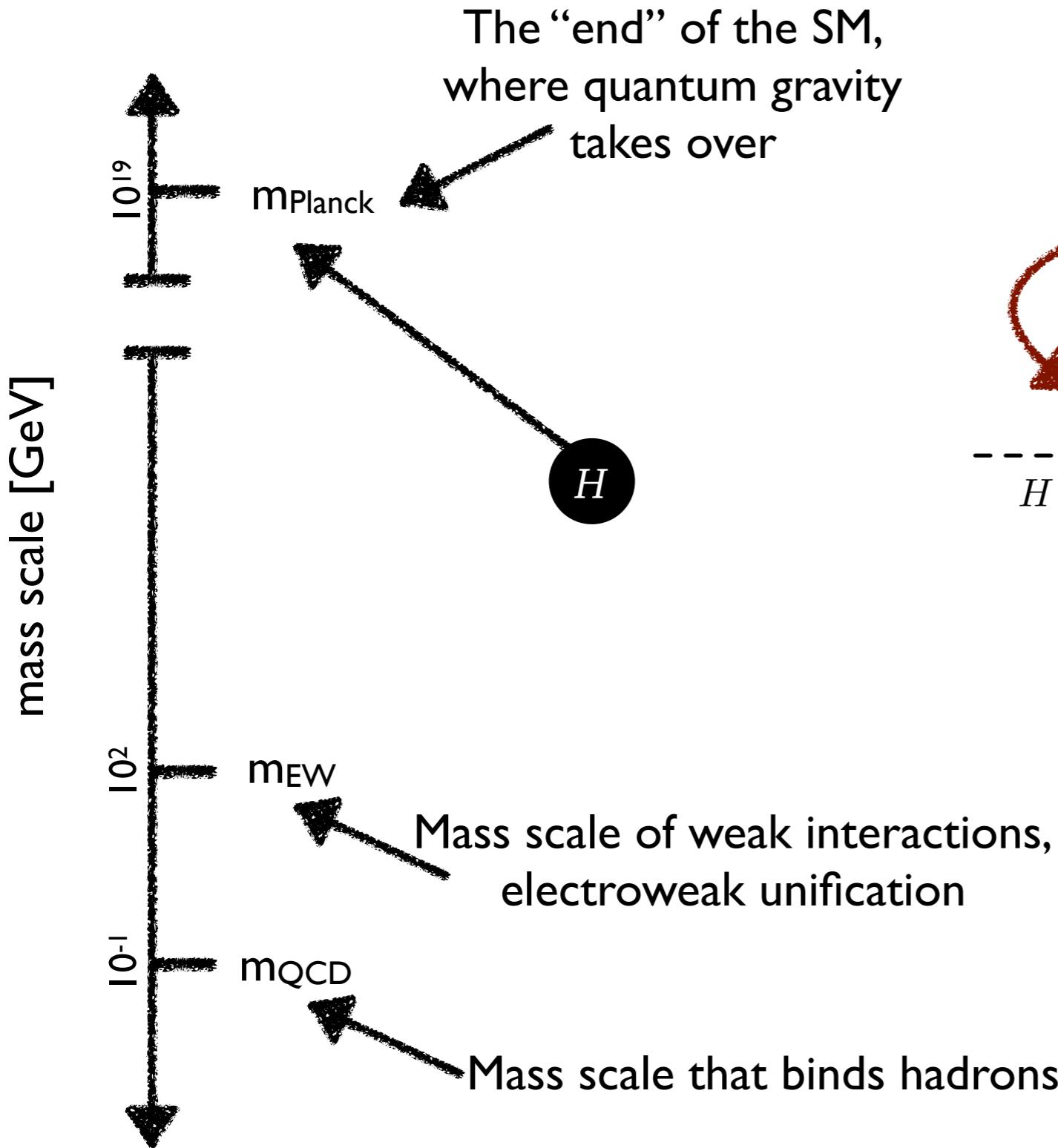
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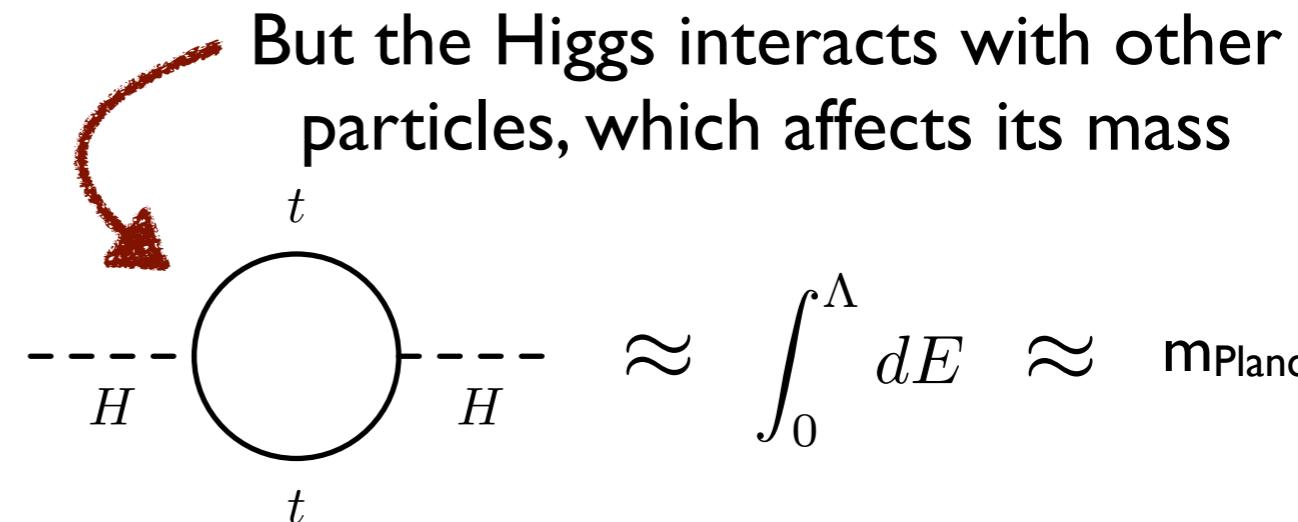




Why BSM?



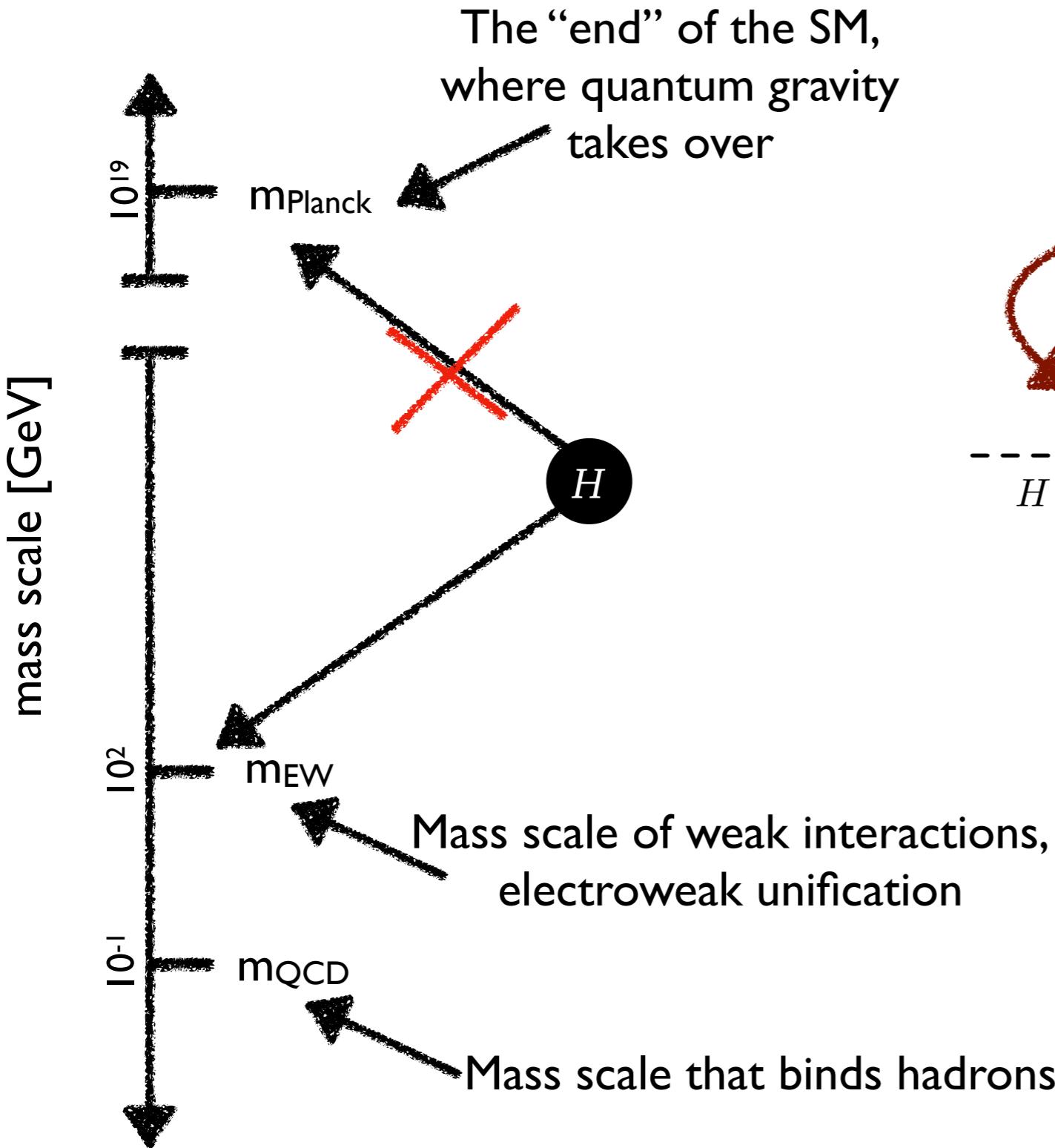
To first order, the mass of the Higgs is a free parameter, and can be anywhere



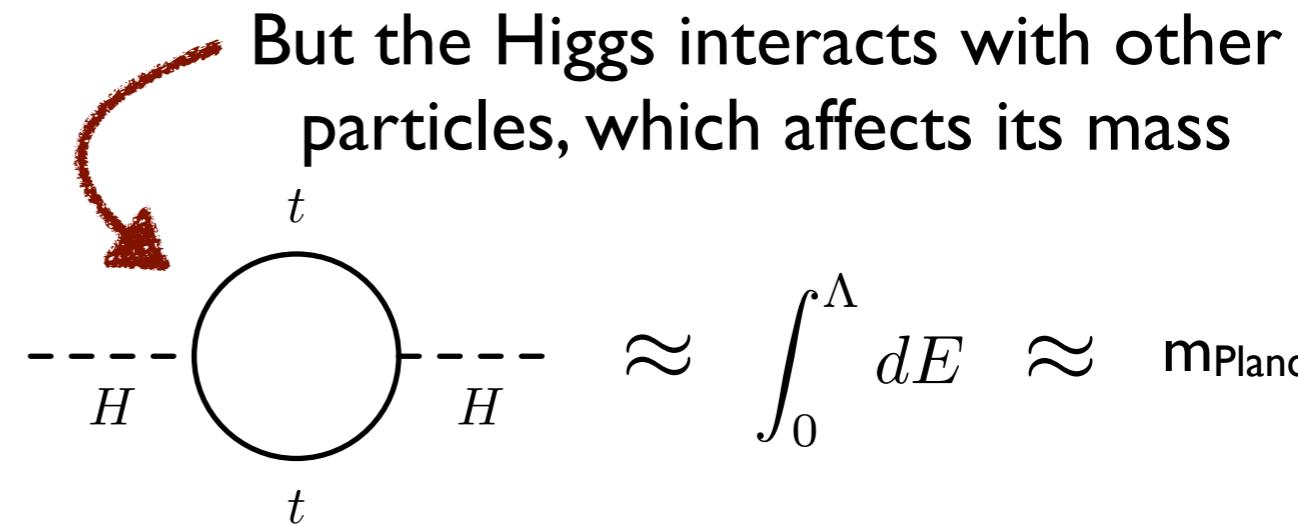
In the SM, this correction should set the Higgs mass to $\sim m_{\text{Planck}}$



Why BSM?



To first order, the mass of the Higgs is a free parameter, and can be anywhere

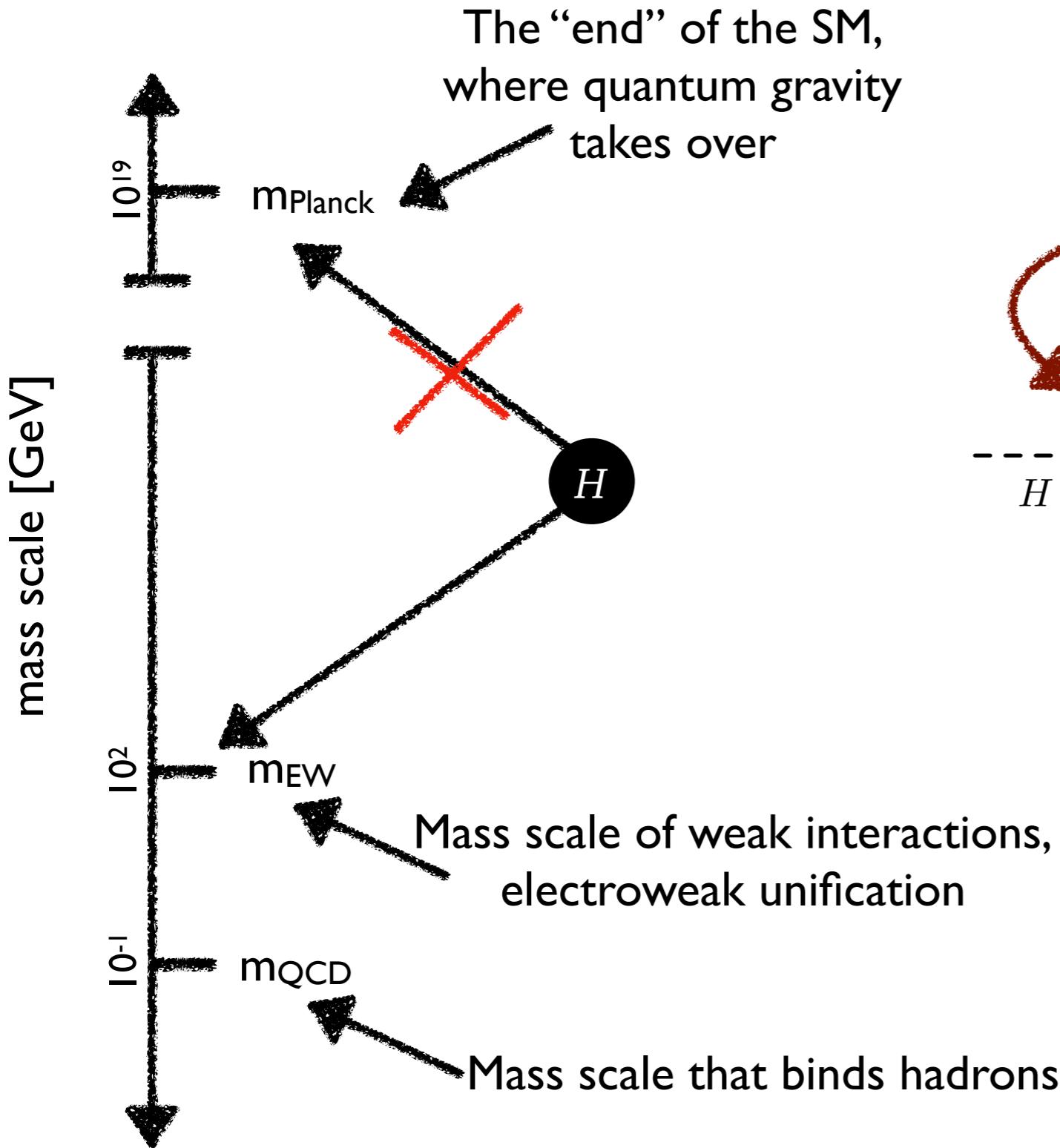


In the SM, this correction should set the Higgs mass to $\sim m_{\text{Planck}}$

But we observe it at 125 GeV!

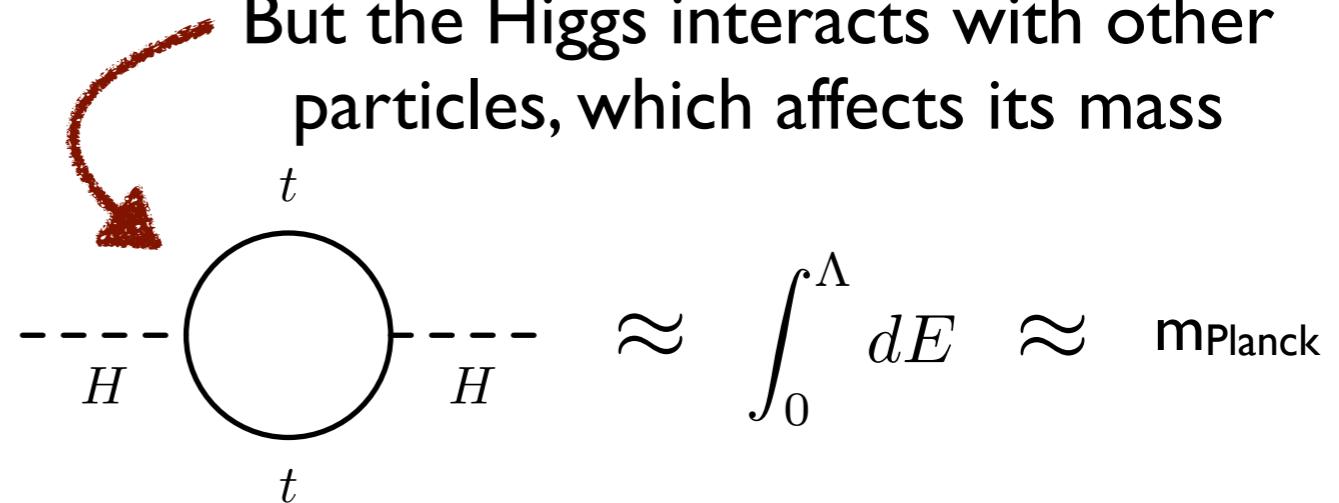


Why BSM?



To first order, the mass of the Higgs is a free parameter, and can be anywhere

But the Higgs interacts with other particles, which affects its mass



In the SM, this correction should set the Higgs mass to $\sim m_{\text{Planck}}$

But we observe it at 125 GeV!

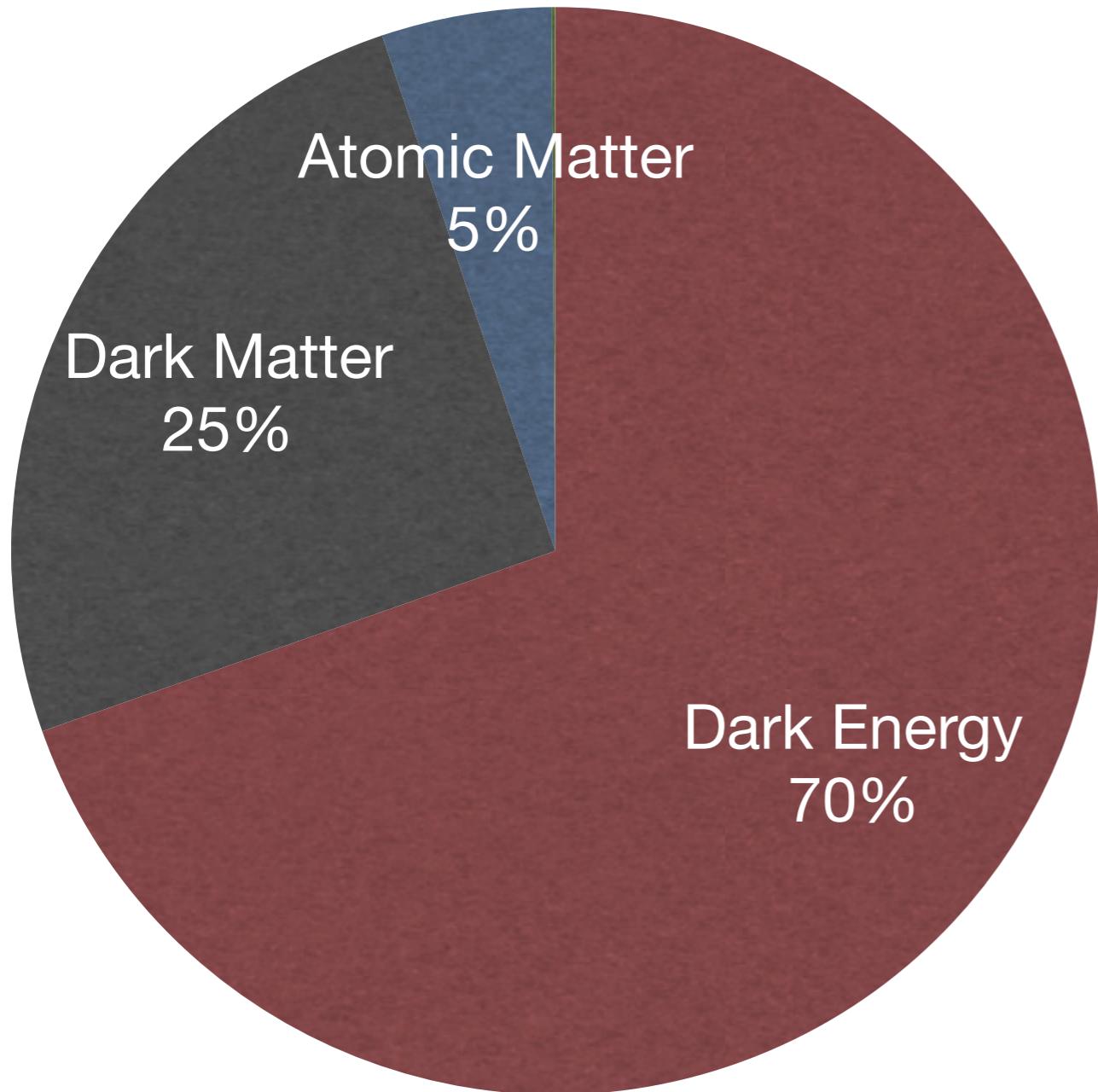
BSM can stabilize the Higgs mass

Why BSM?



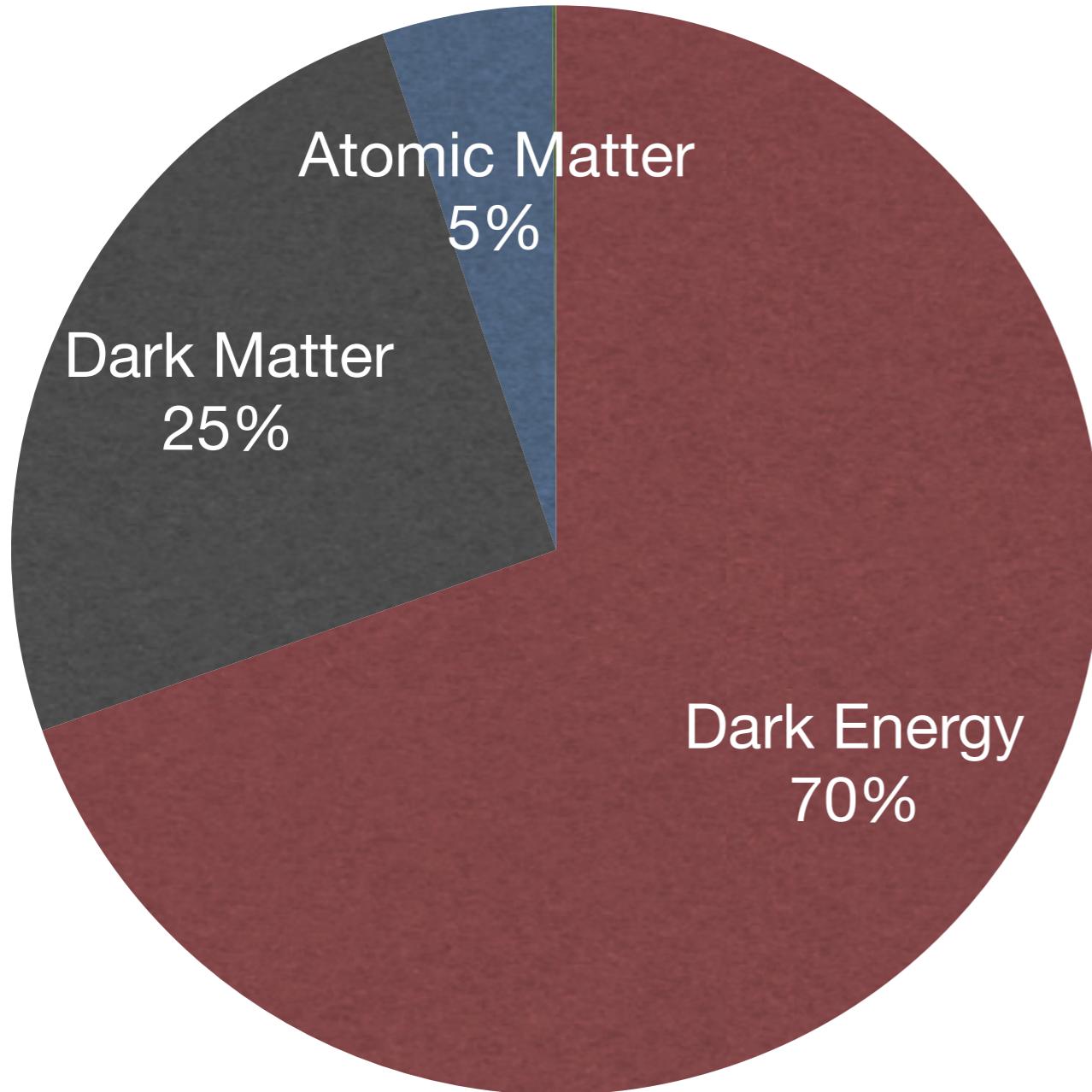


Why BSM?





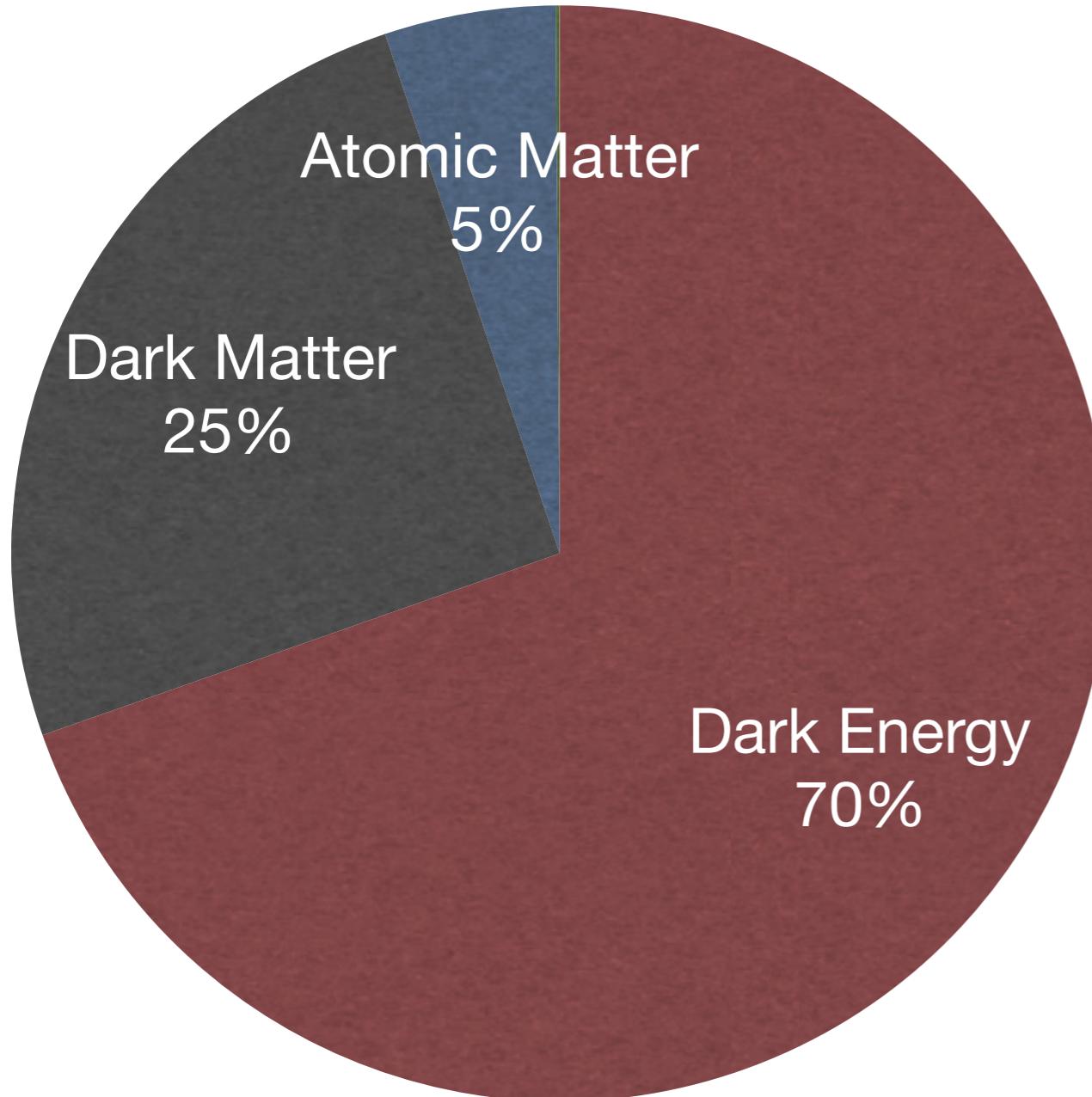
Why BSM?



We don't understand the majority of the universe!



Why BSM?



We don't understand the majority of the universe!

BSM can provide Dark Matter candidates, which we could produce at the LHC

Why BSM?





Why BSM?

10,000,000,001

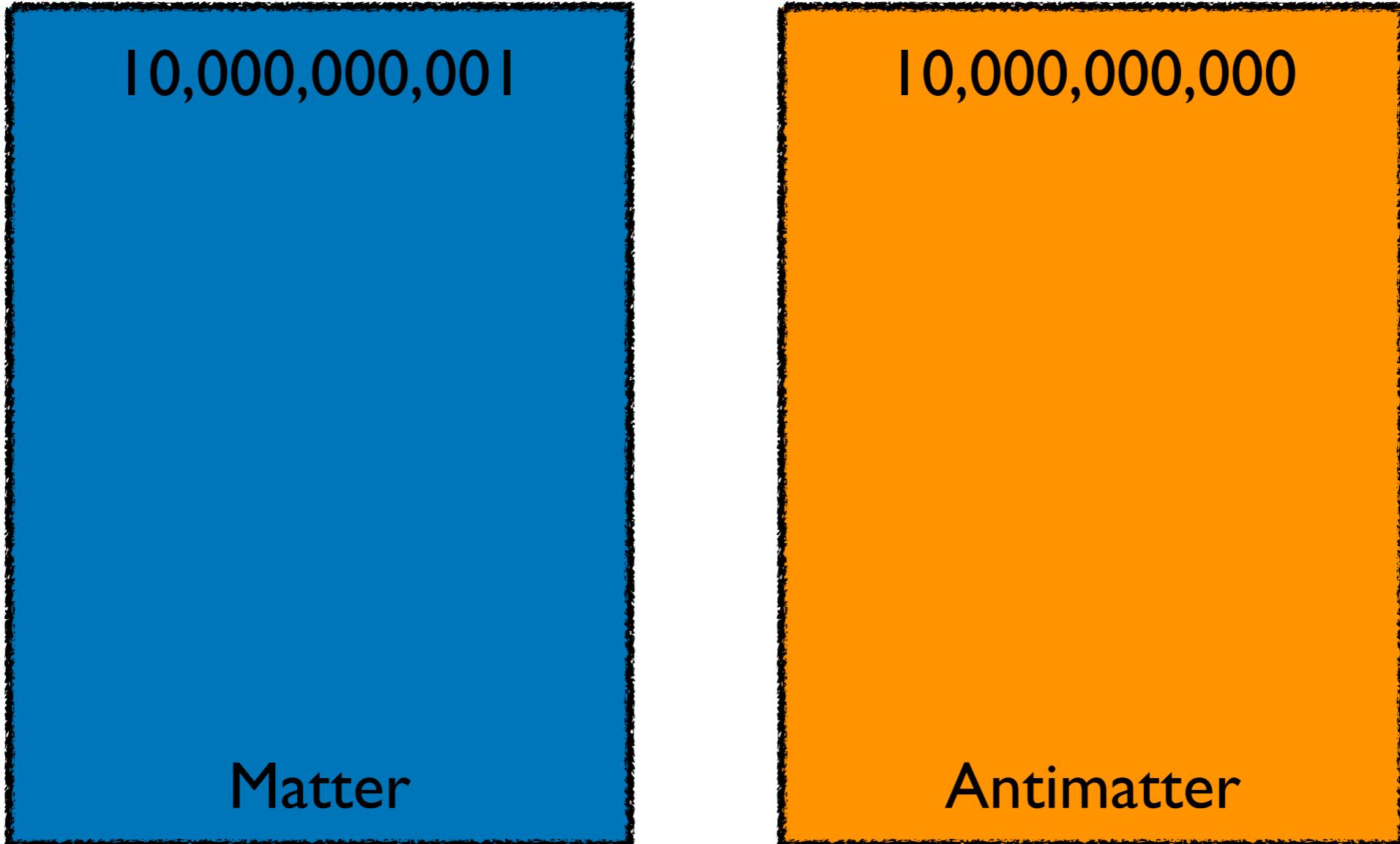
Matter

10,000,000,000

Antimatter



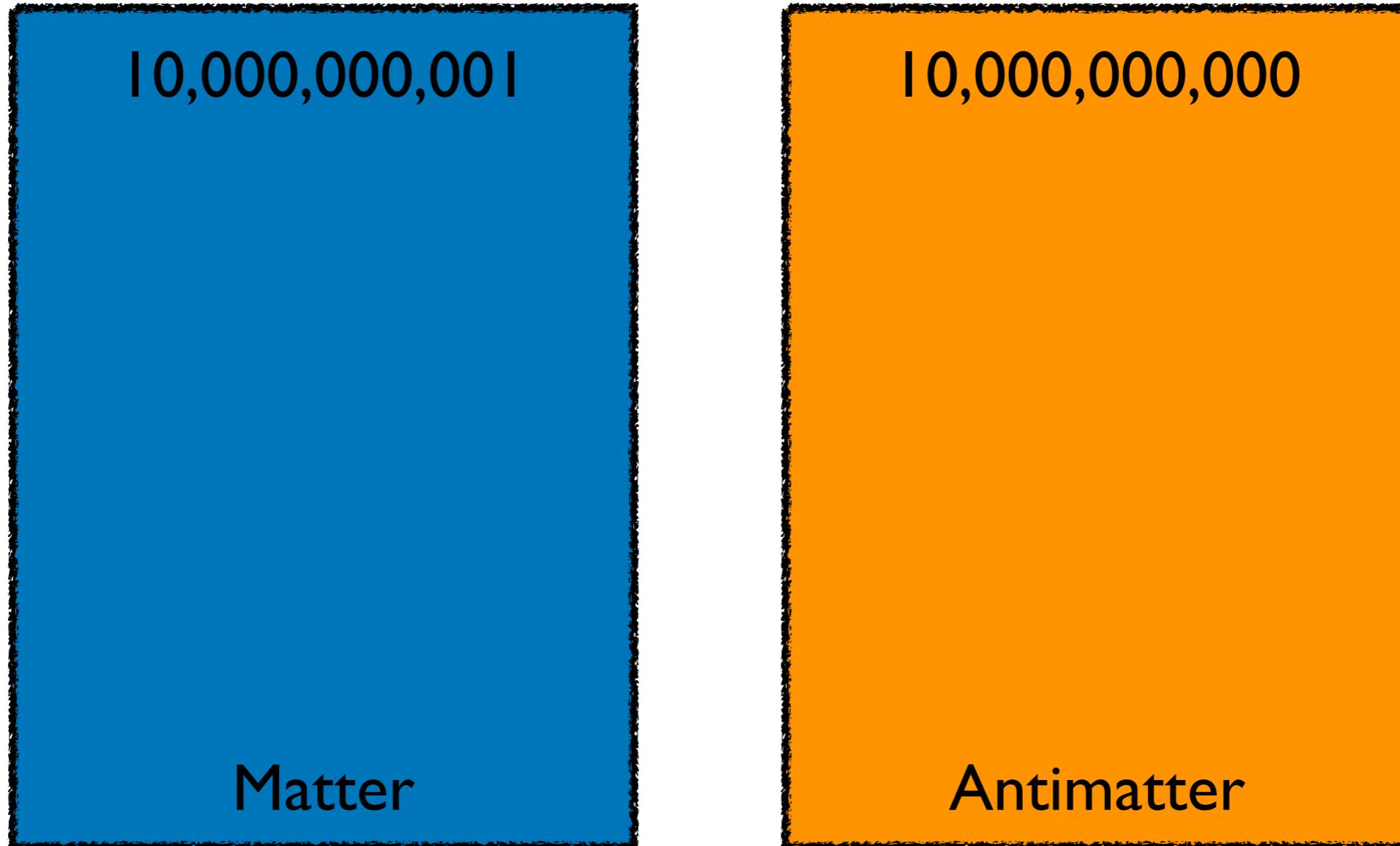
Why BSM?



Where did the tiny matter/antimatter asymmetry that leads to a matter dominated universe arise?



Why BSM?

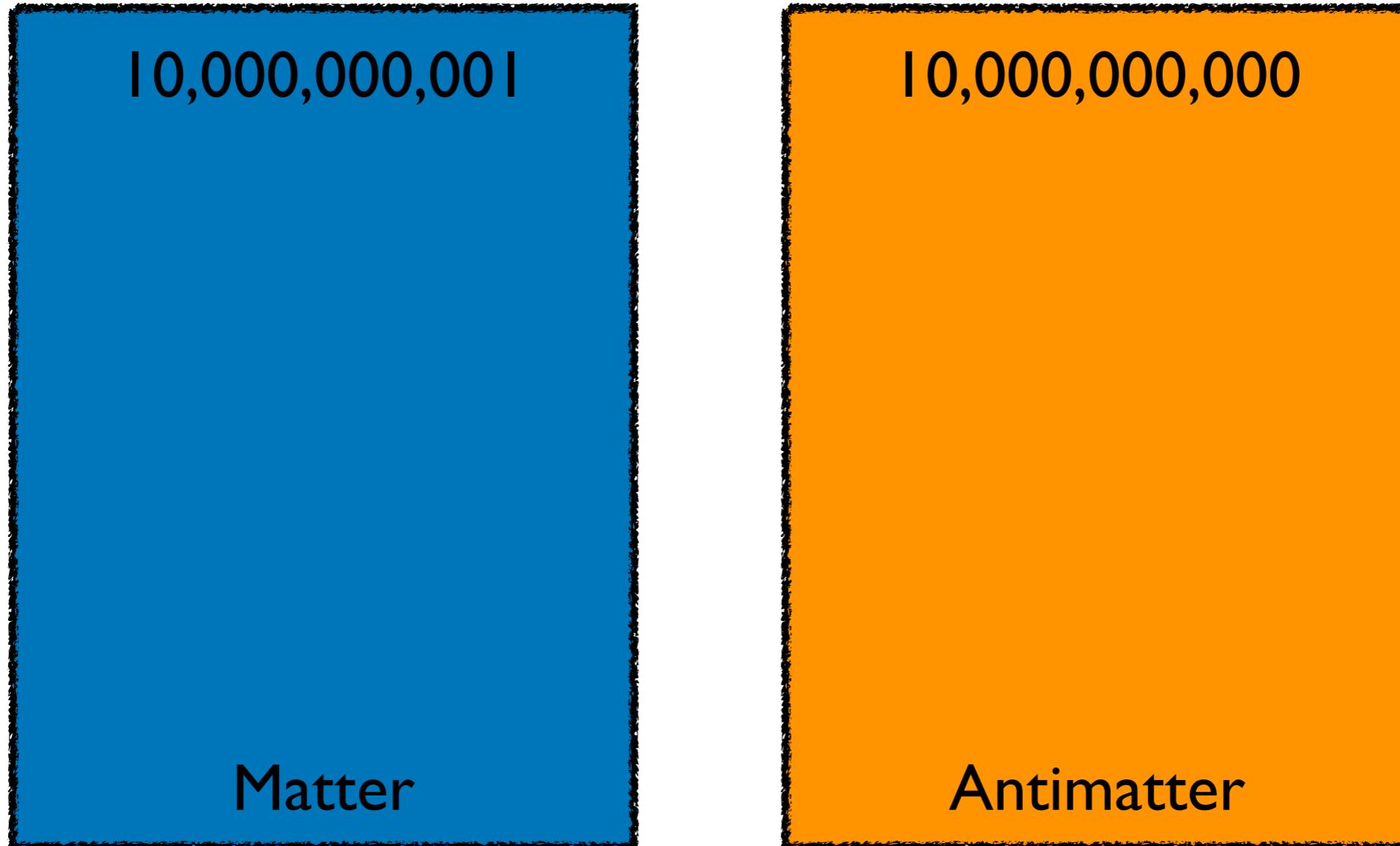


Where did the tiny matter/antimatter asymmetry that leads to a matter dominated universe arise?

CP violation in the SM is not enough



Why BSM?



Where did the tiny matter/antimatter asymmetry
that leads to a matter dominated universe arise?

CP violation in the SM is not enough

BSM can explain the matter dominated universe

How do we search for BSM physics at the LHC?

BSM = Beyond Standard Model

BSM physics = Physics beyond the Standard Model

BSM particles = Particles beyond the Standard Model

BSM interactions = Interactions beyond the Standard Model

BSM processes = Processes beyond the Standard Model

BSM signals = Signals beyond the Standard Model

BSM backgrounds = Backgrounds beyond the Standard Model

BSM constraints = Constraints beyond the Standard Model

BSM limits = Limits beyond the Standard Model

BSM discovery = Discovery beyond the Standard Model

BSM exclusion = Exclusion beyond the Standard Model

BSM interpretation = Interpretation beyond the Standard Model

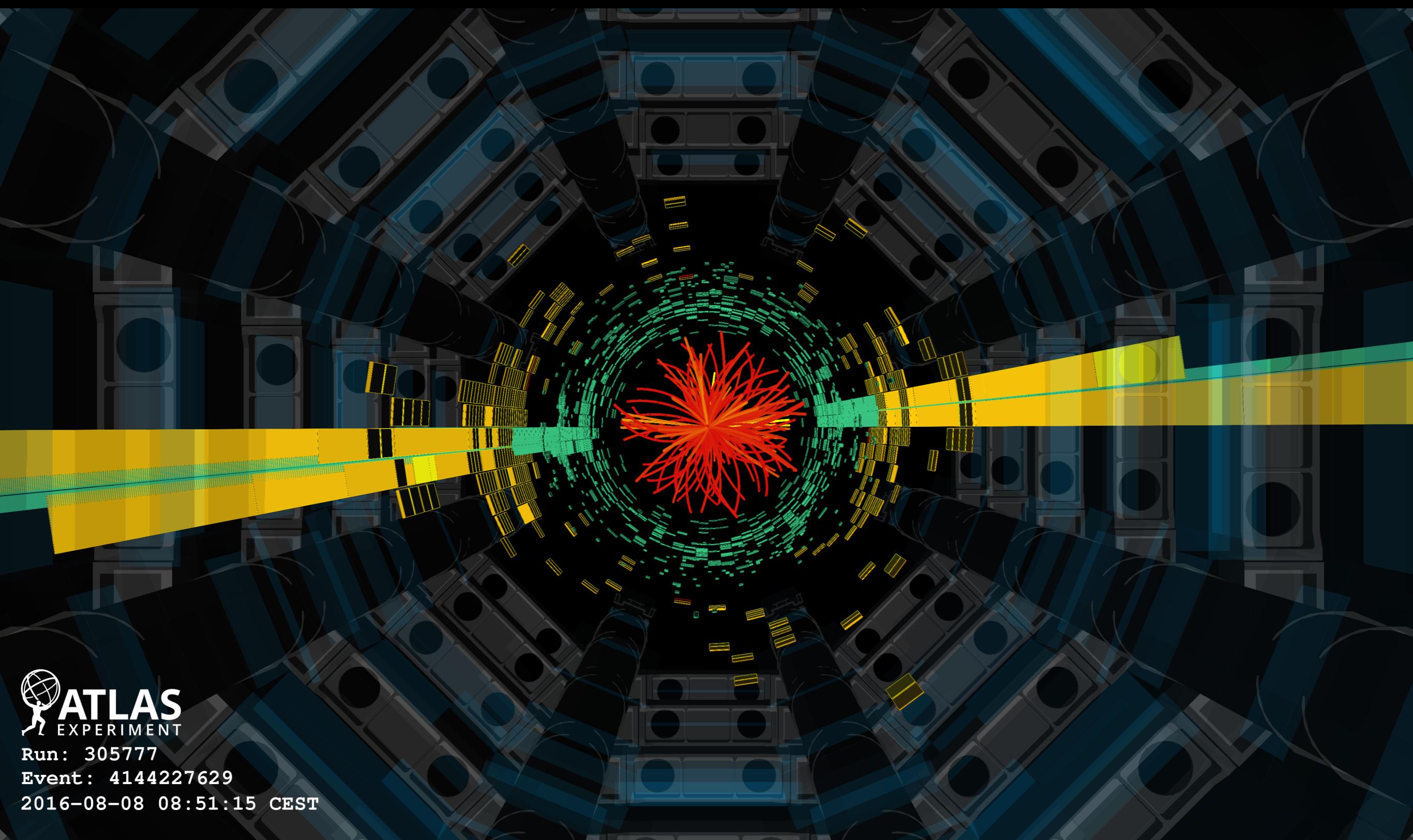
BSM theory = Theory beyond the Standard Model

BSM model = Model beyond the Standard Model

BSM mechanism = Mechanism beyond the Standard Model

BSM mechanism = Mechanism beyond the Standard Model

How do we search for BSM physics at the LHC?



ATLAS
EXPERIMENT
Run: 305777
Event: 4144227629
2016-08-08 08:51:15 CEST

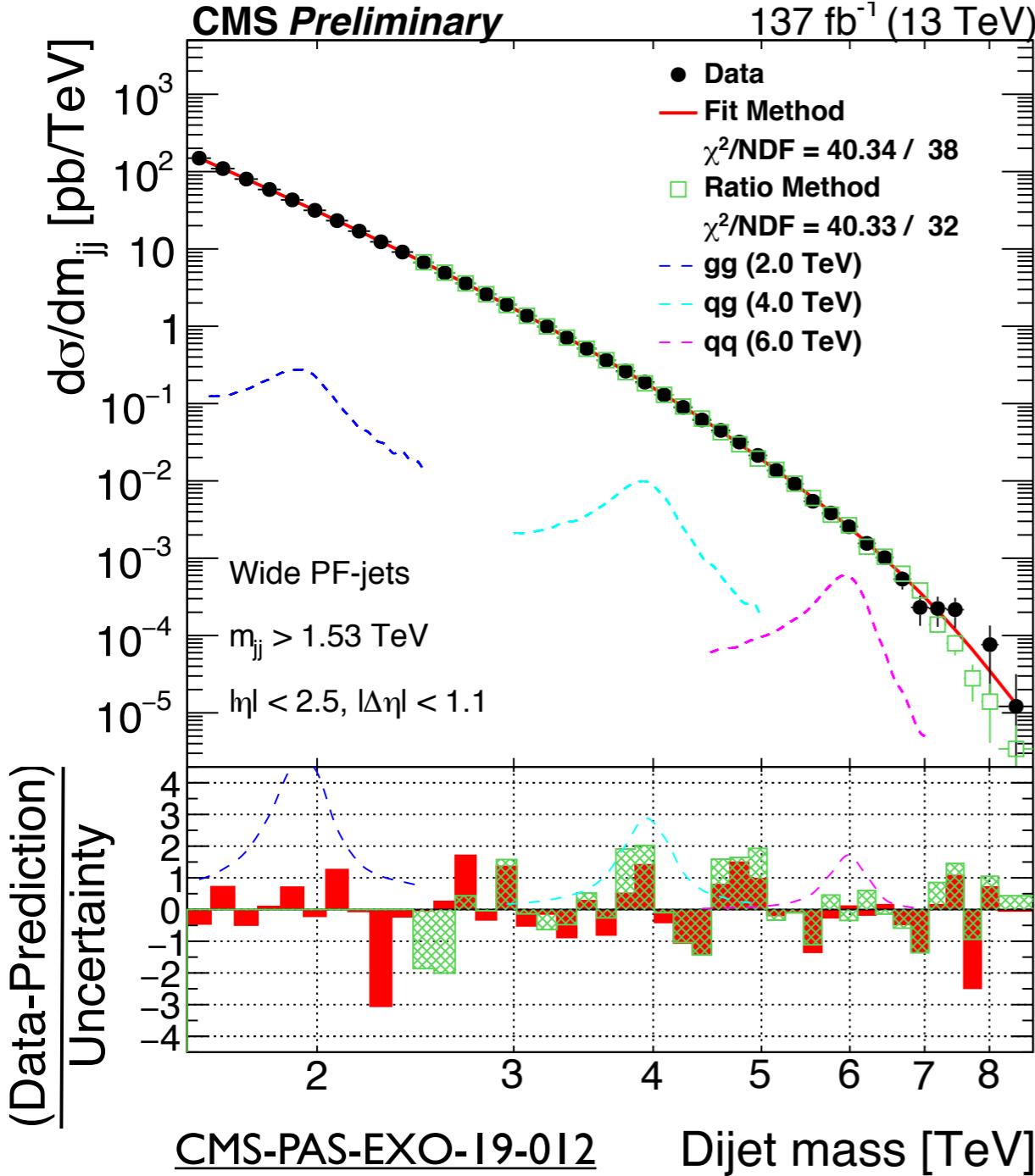
Searching with Dijets





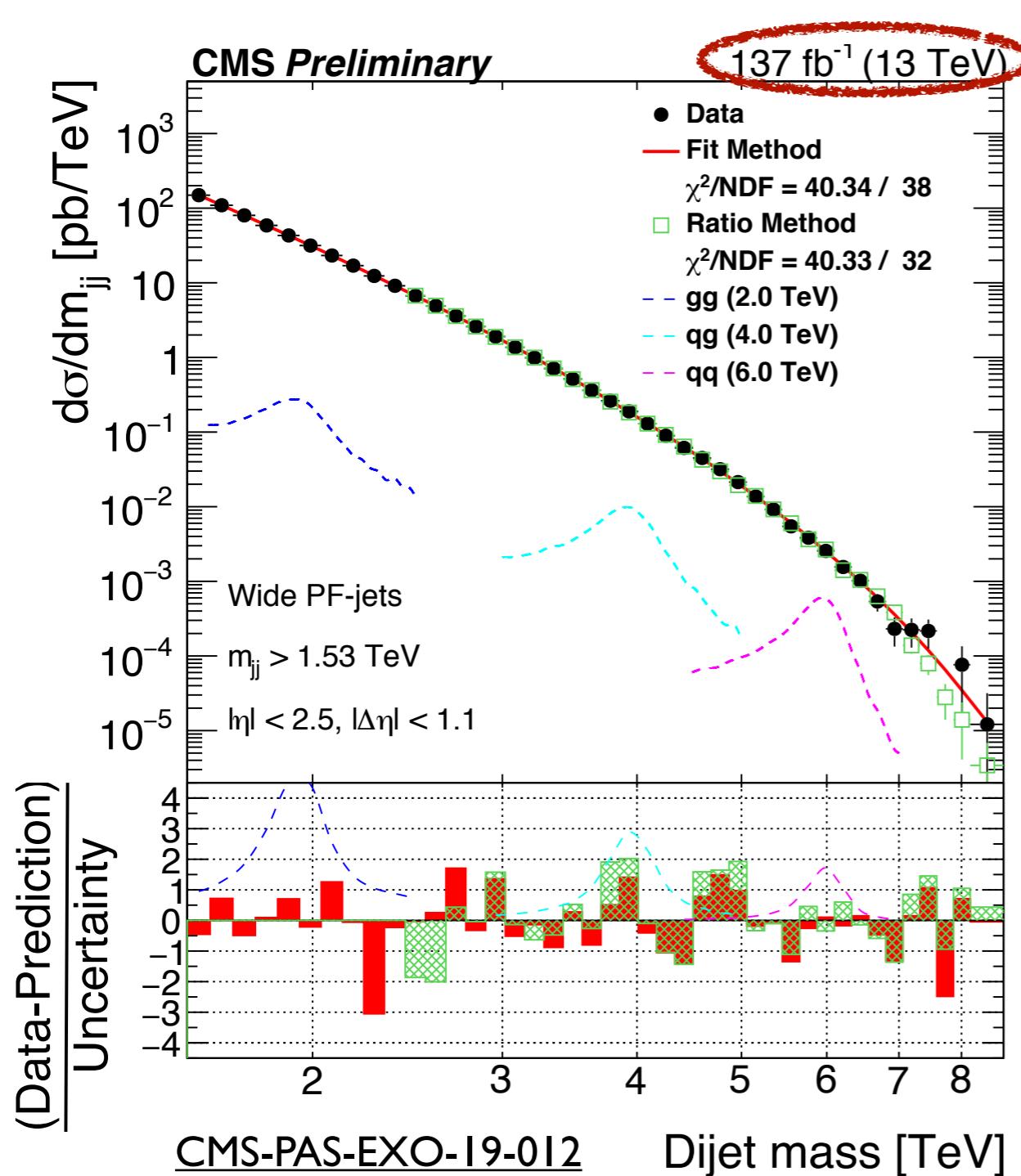
Searching with Dijets

One of the best BSM signatures:
pairs of jets





Searching with Dijets

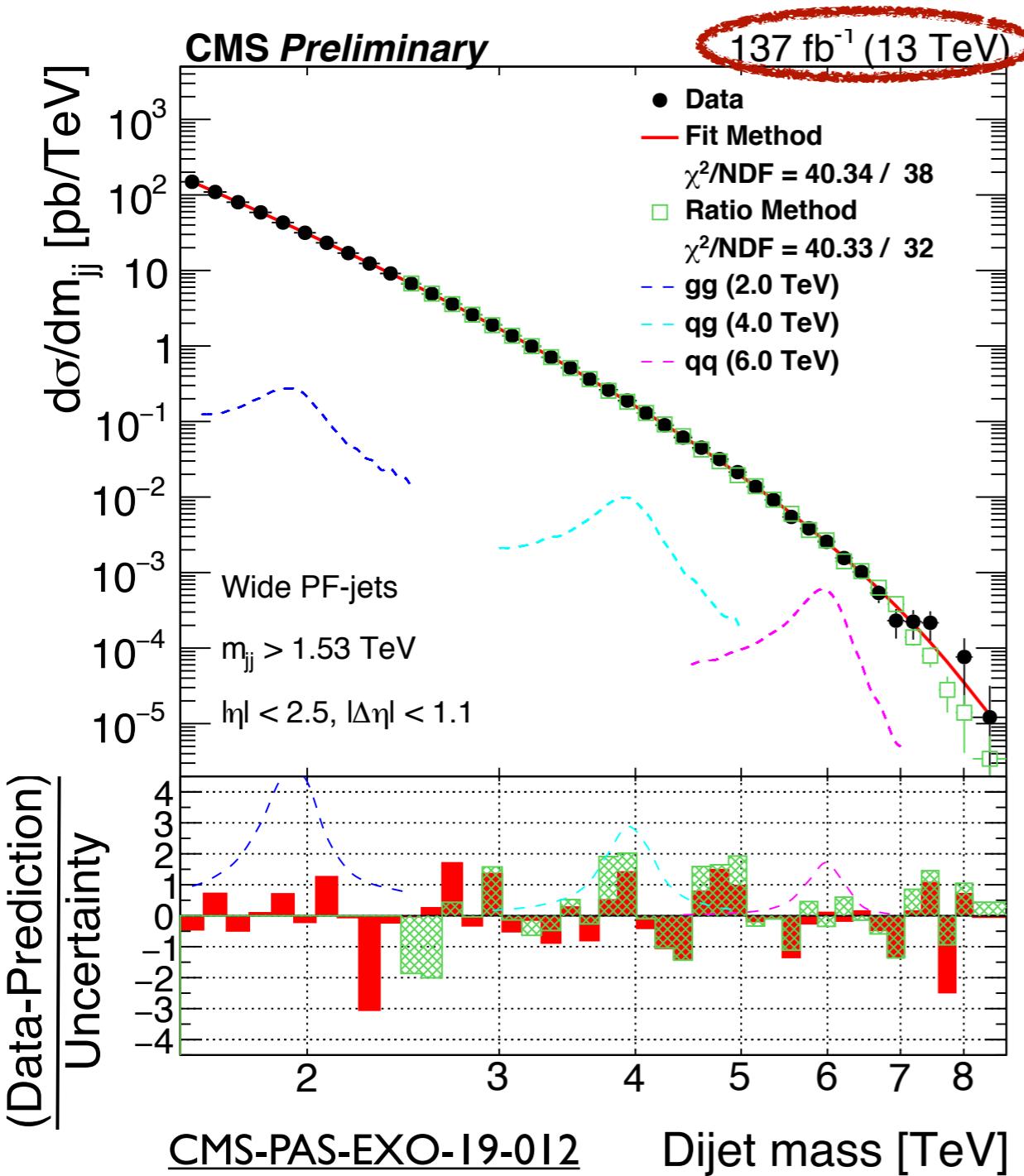


One of the best BSM signatures:
pairs of jets

Huge dataset enables some of
the best sensitivity yet!



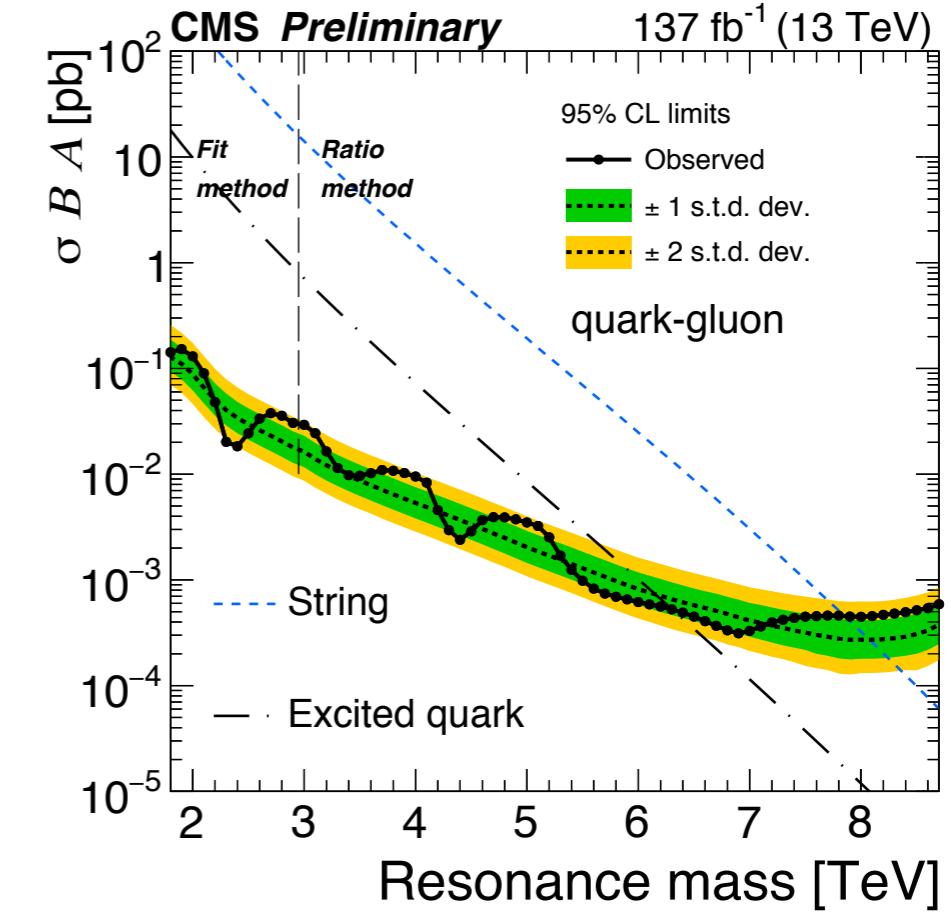
Searching with Dijets



One of the best BSM signatures:
pairs of jets

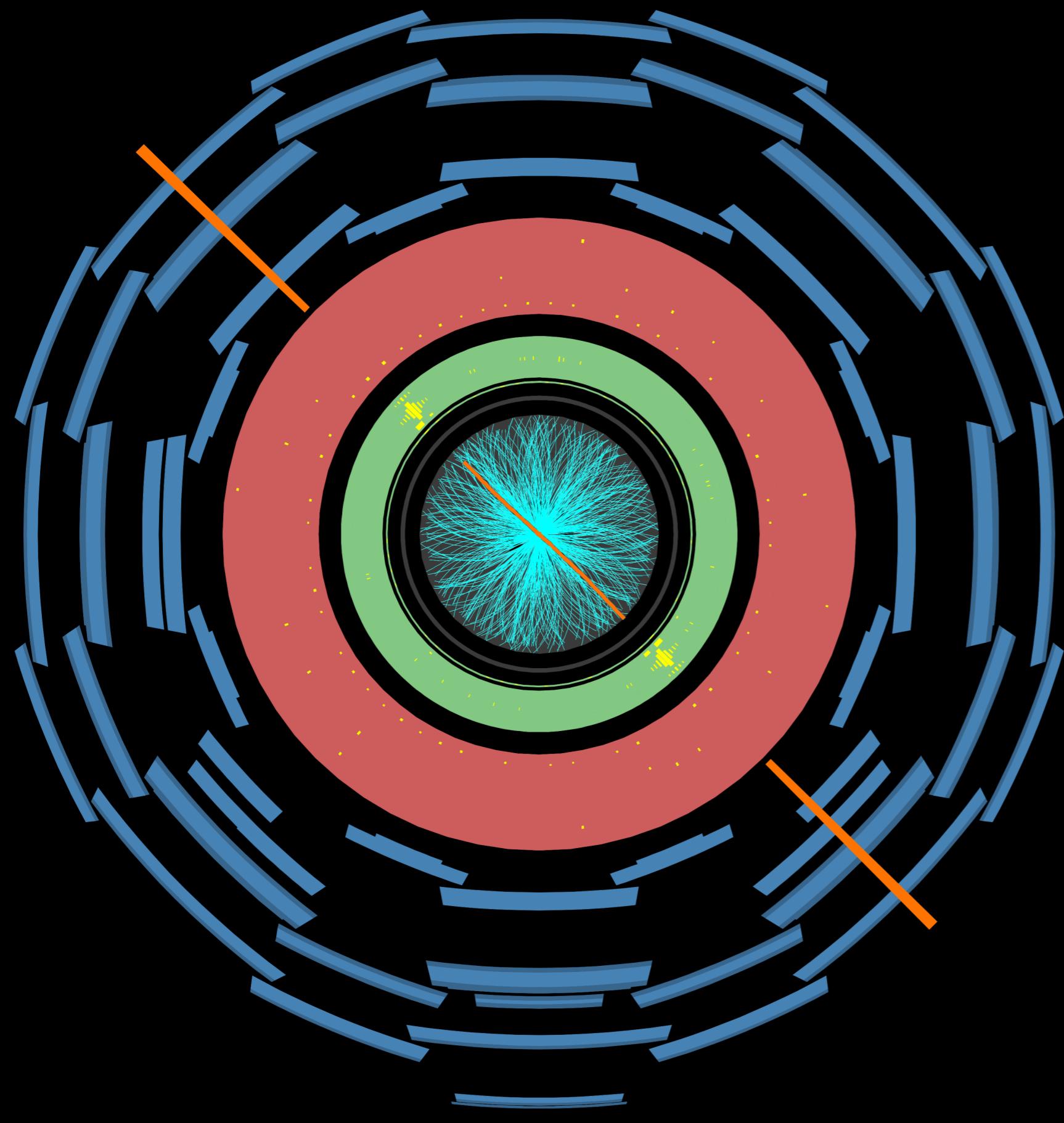
Huge dataset enables some of
the best sensitivity yet!

But no hints of new physics...



What about other particles?

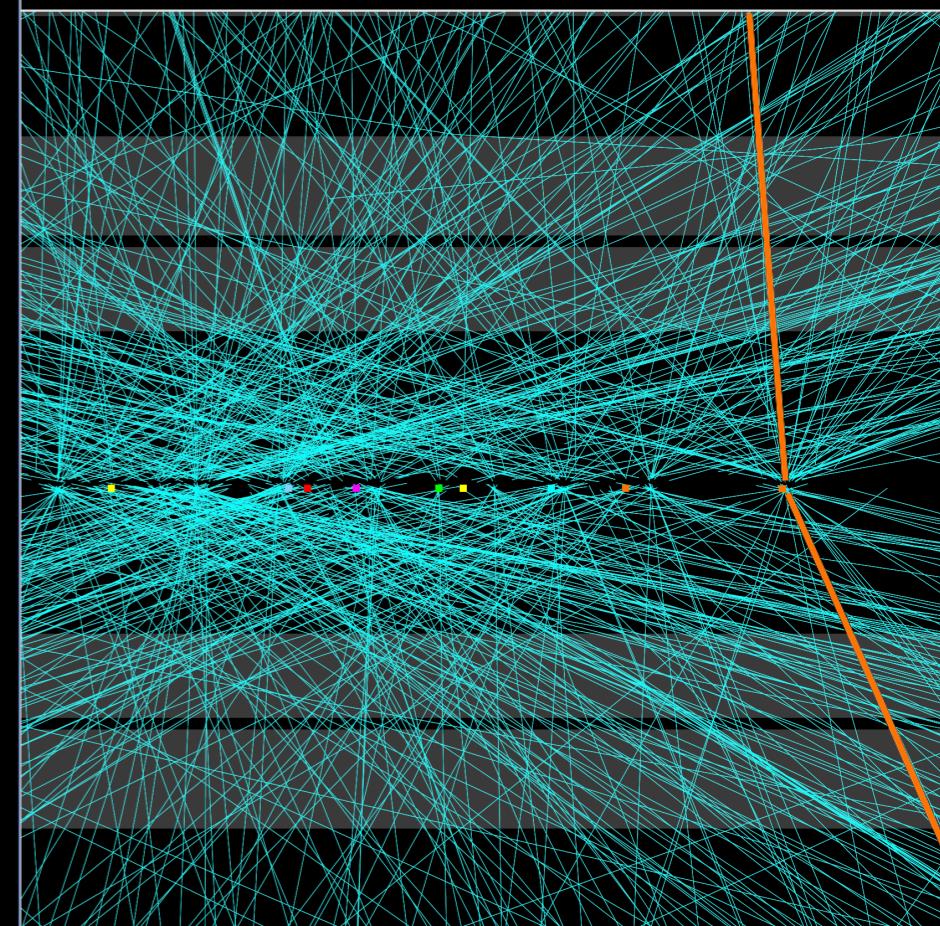
What about other particles?



ATLAS
EXPERIMENT

Run Number: 336852, Event Number: 1440436043

Date: 2017-09-29 11:44:35 CEST

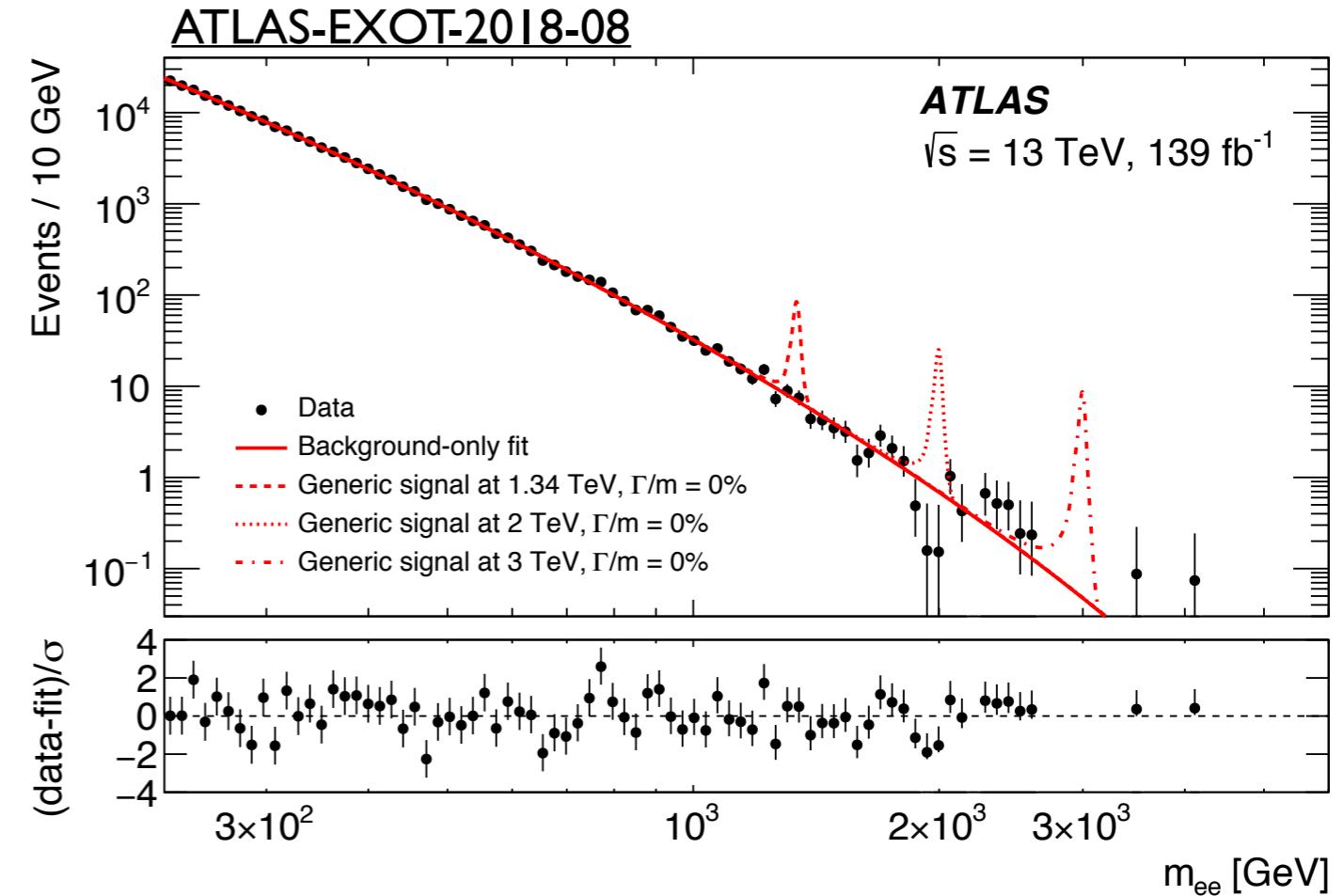
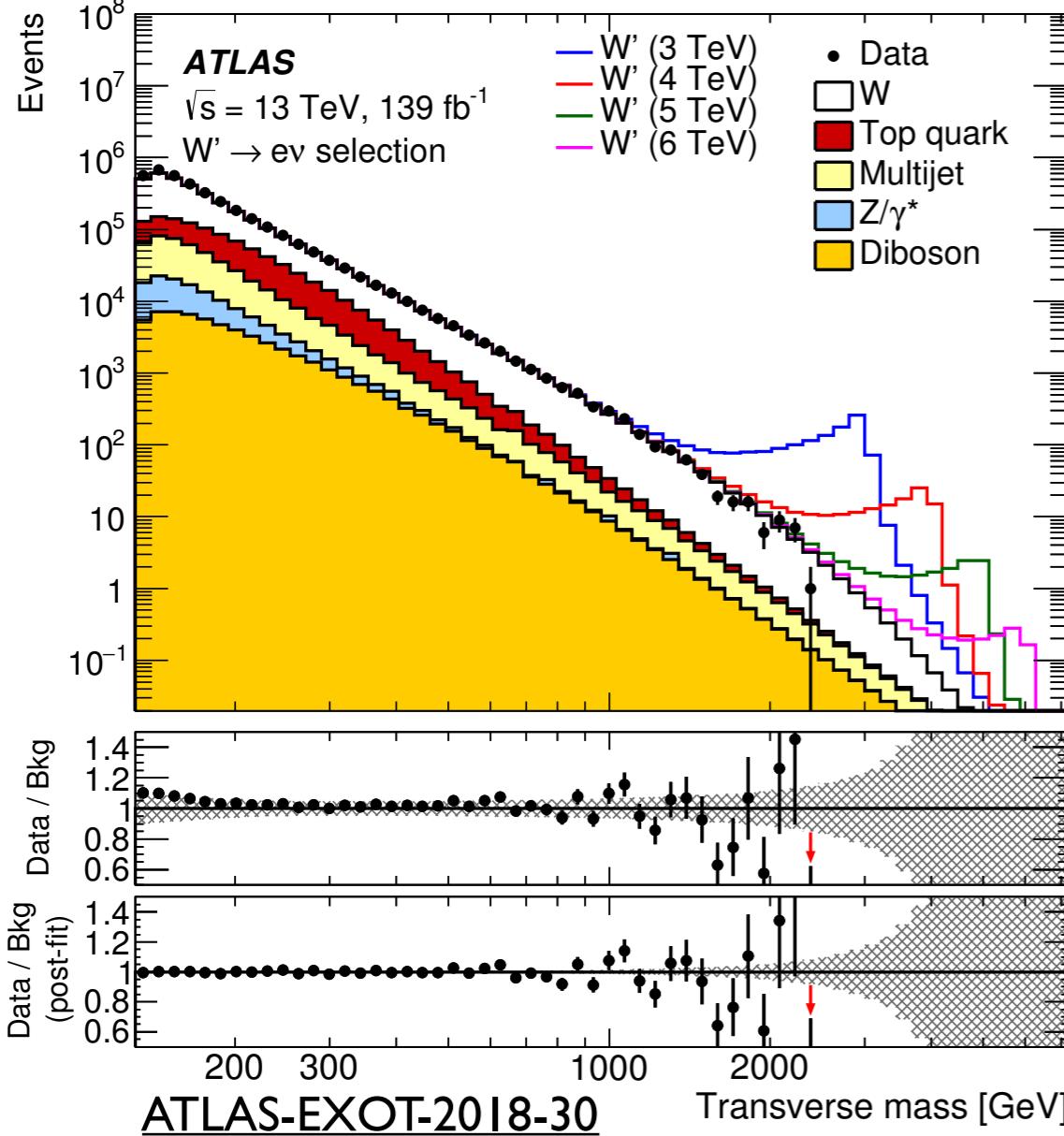


Searching with Leptons





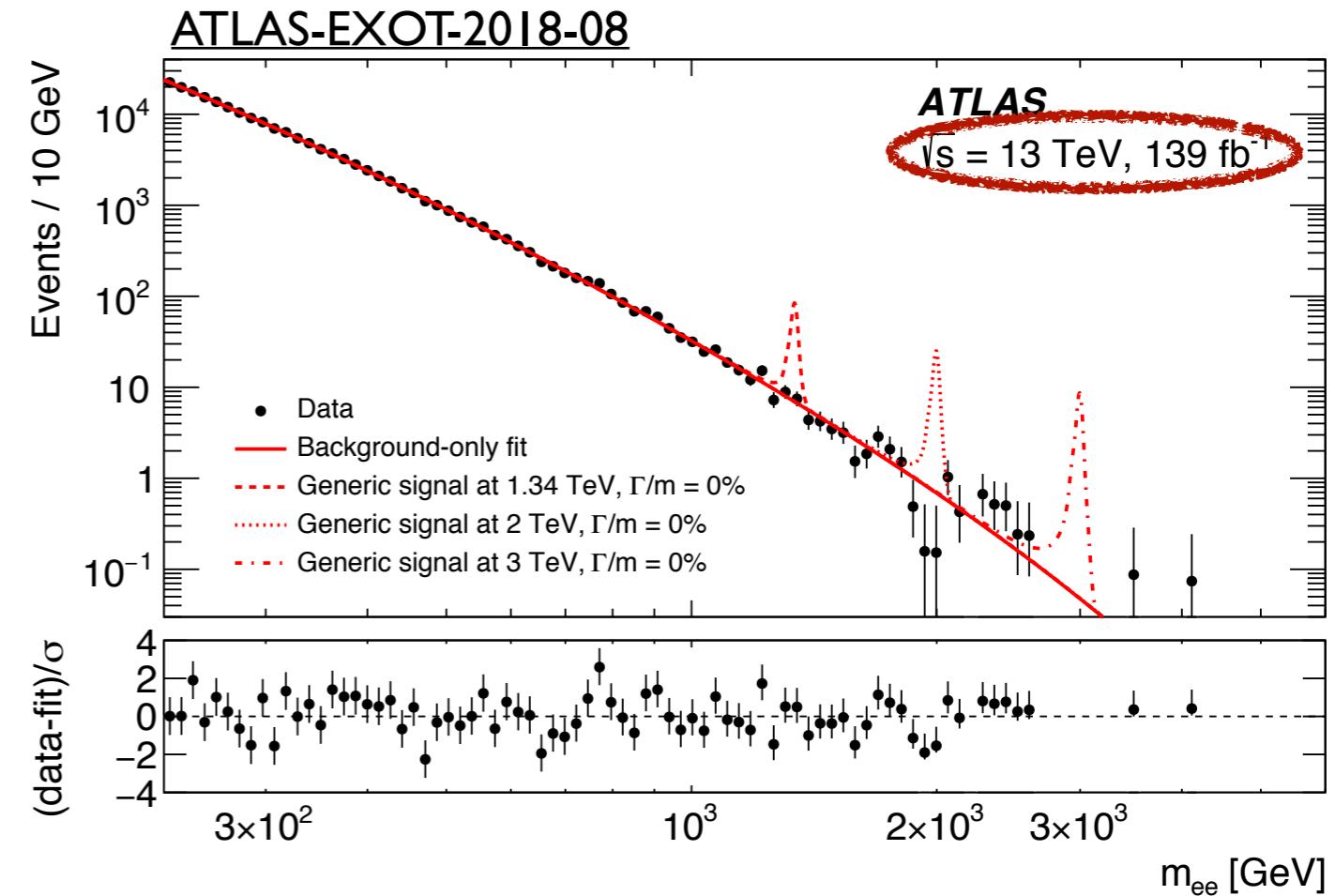
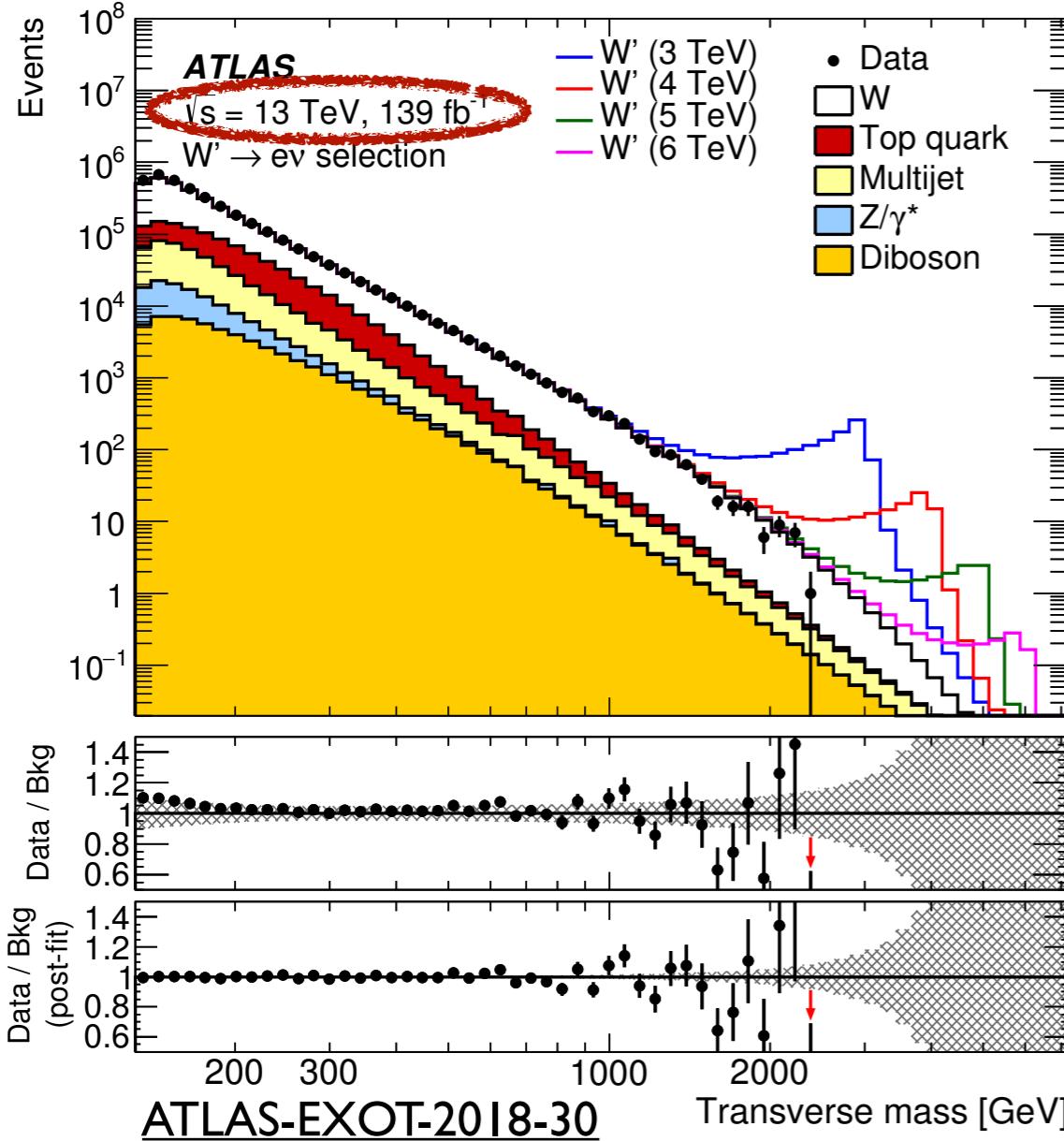
Searching with Leptons



You can also search for BSM with lepton resonances!



Searching with Leptons



You can also search for BSM with lepton resonances!

Not better luck here,
even with the full datasets

What about
more
complicated
signatures?

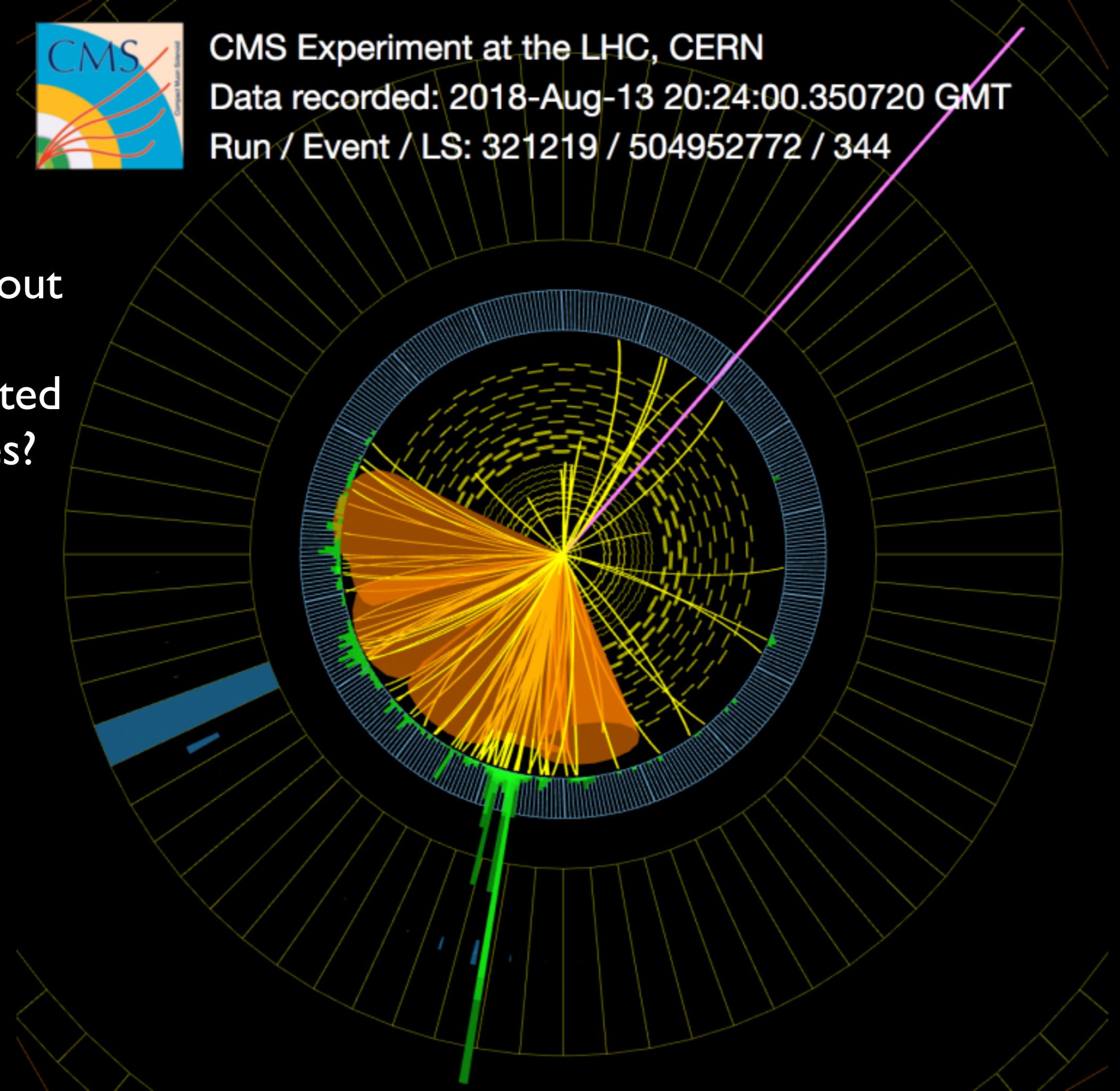


CMS Experiment at the LHC, CERN

Data recorded: 2018-Aug-13 20:24:00.350720 GMT

Run / Event / LS: 321219 / 504952772 / 344

What about
more
complicated
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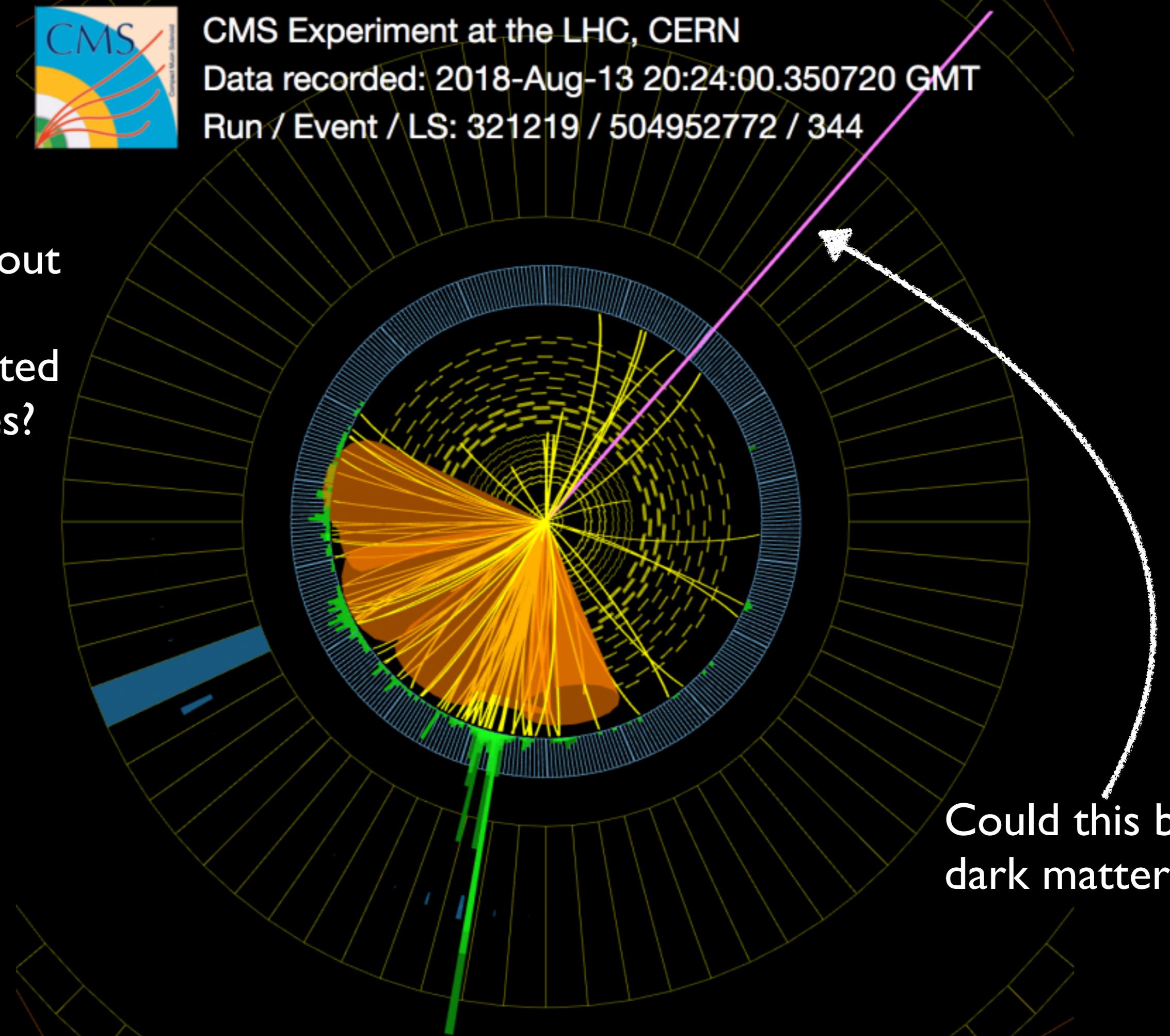


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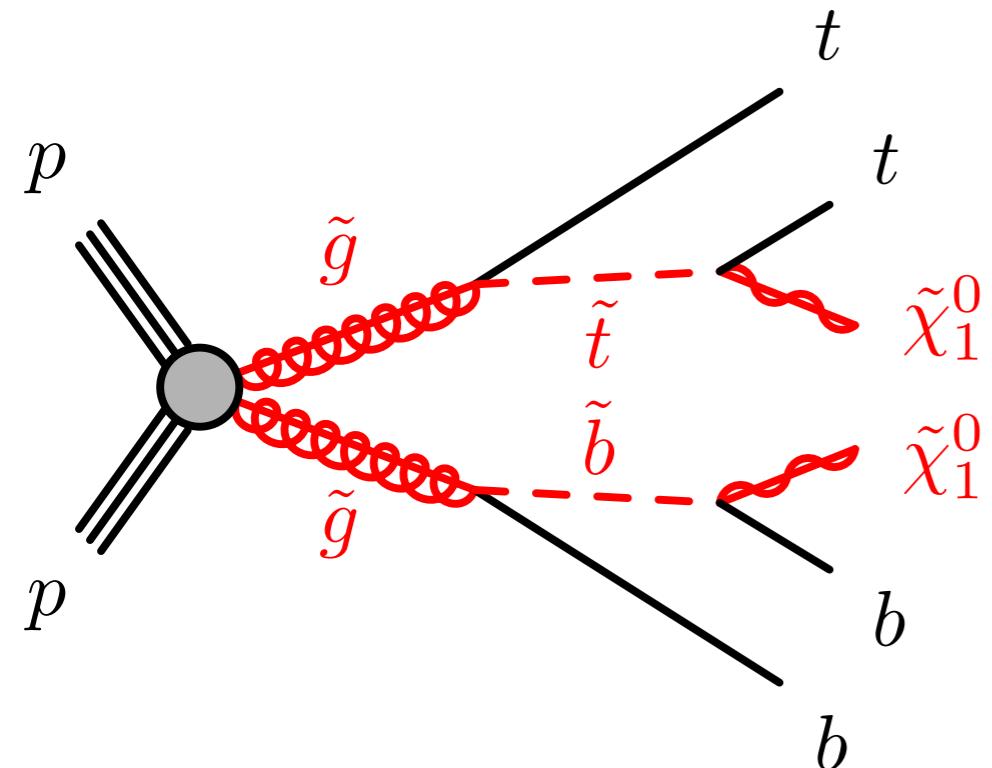


SUSY Signatures





SUSY Signatures

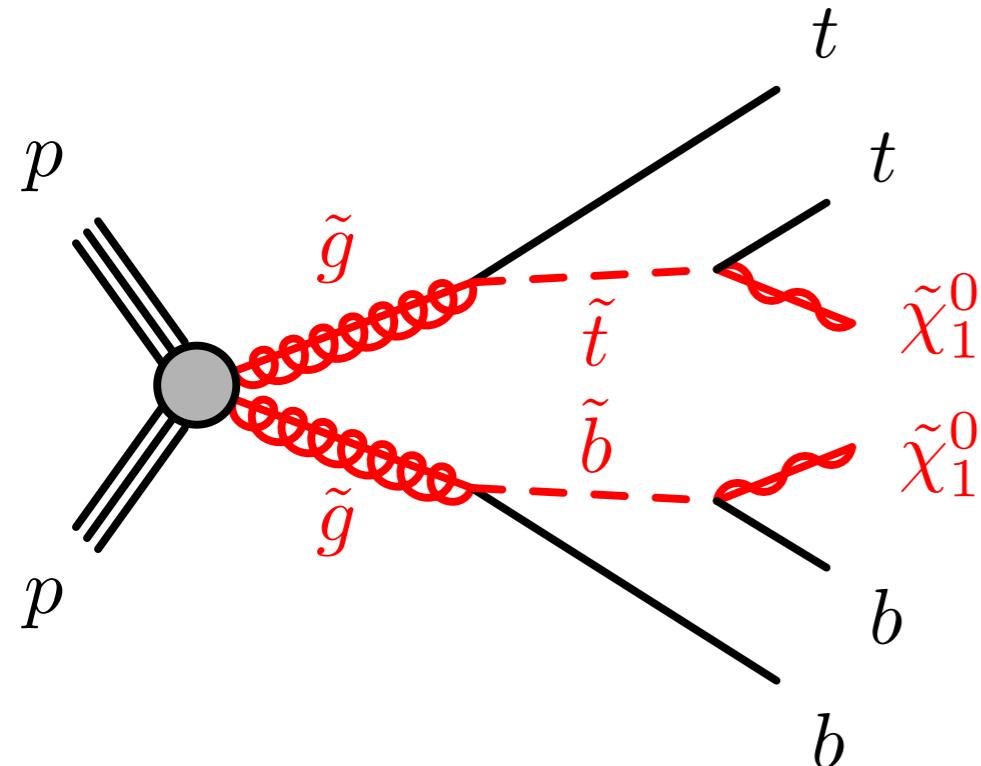


Supersymmetry predicts

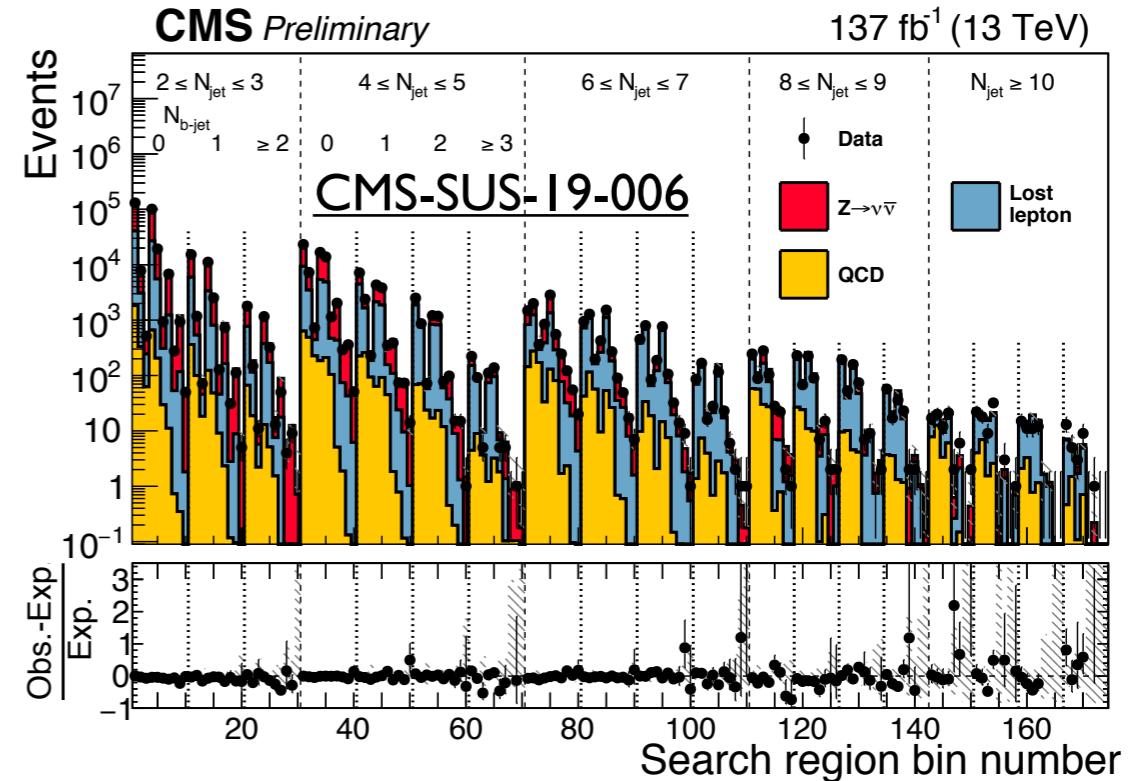
more complicated final states...



SUSY Signatures

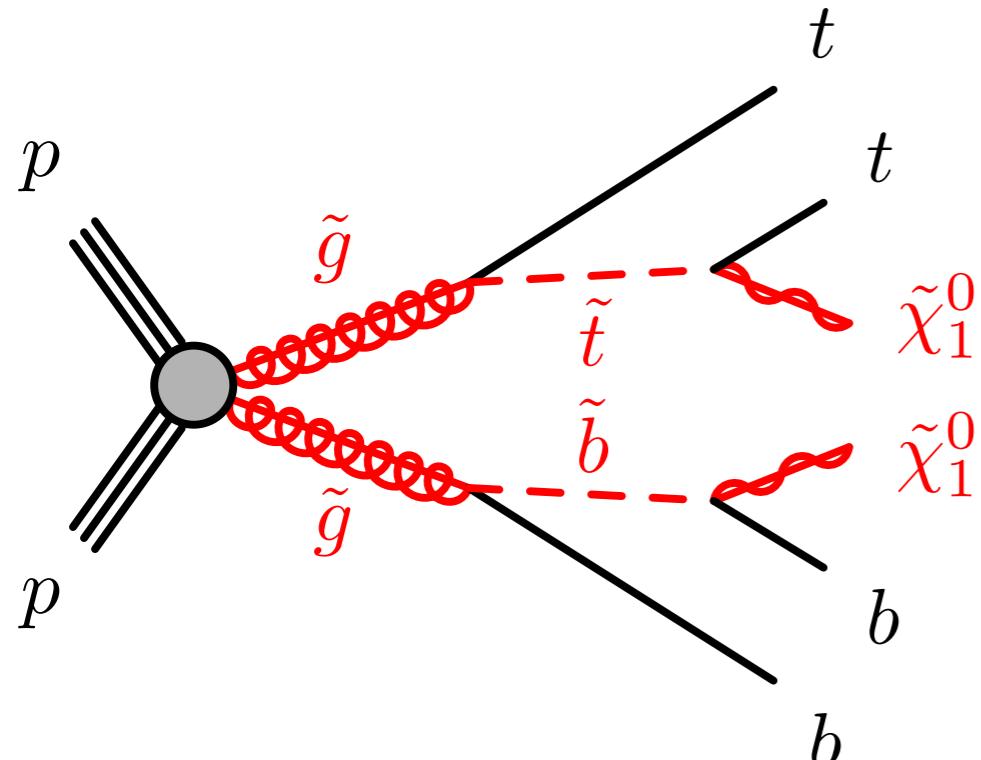


Supersymmetry predicts
more complicated final states...

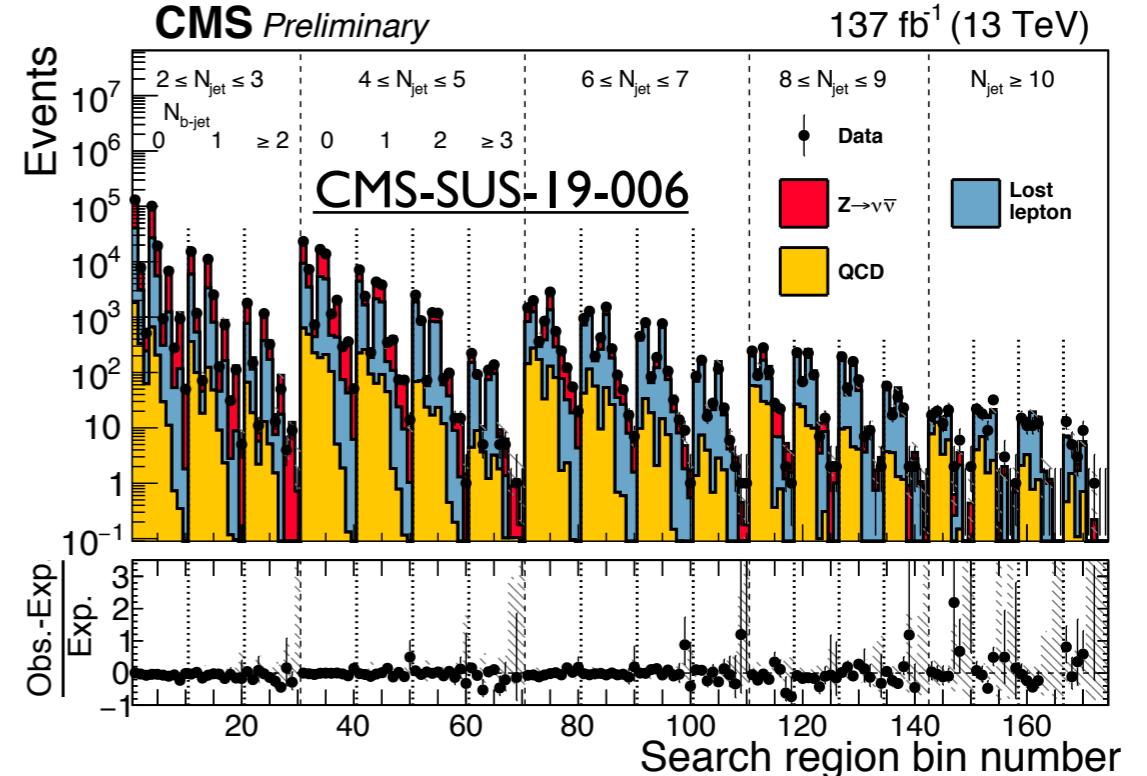
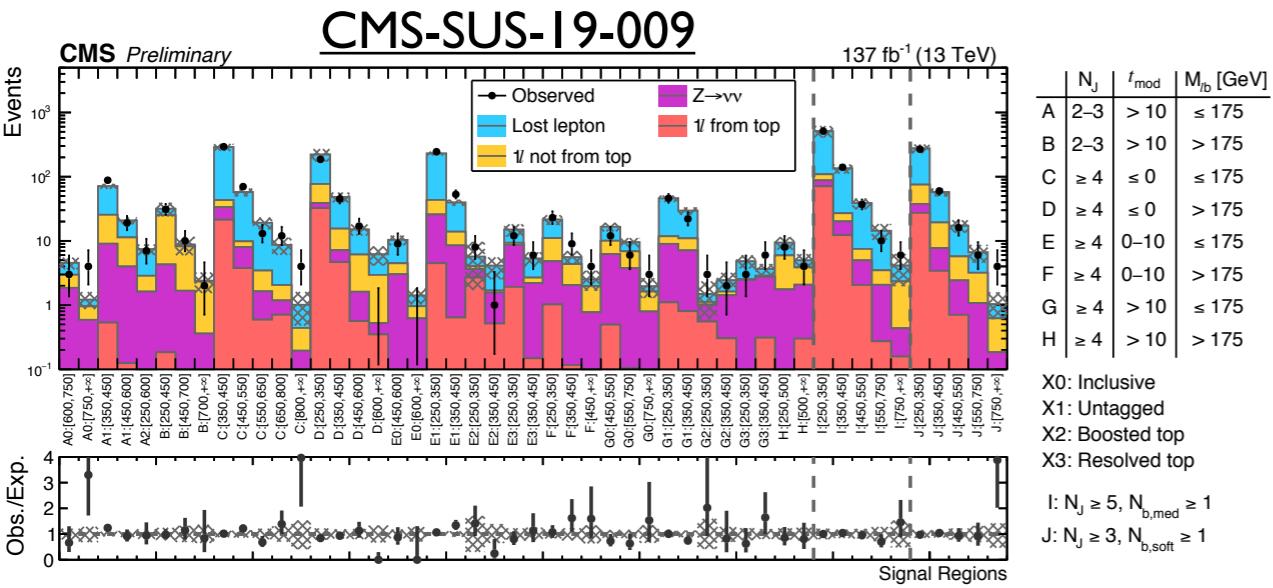




SUSY Signatures

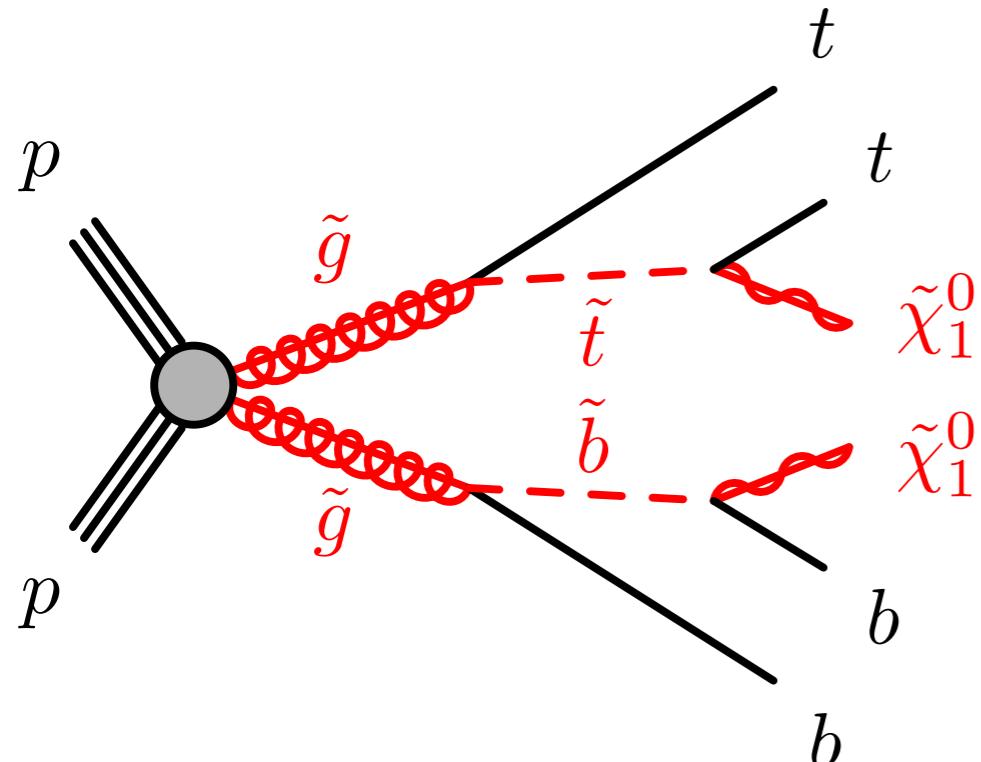


Supersymmetry predicts
more complicated final states...

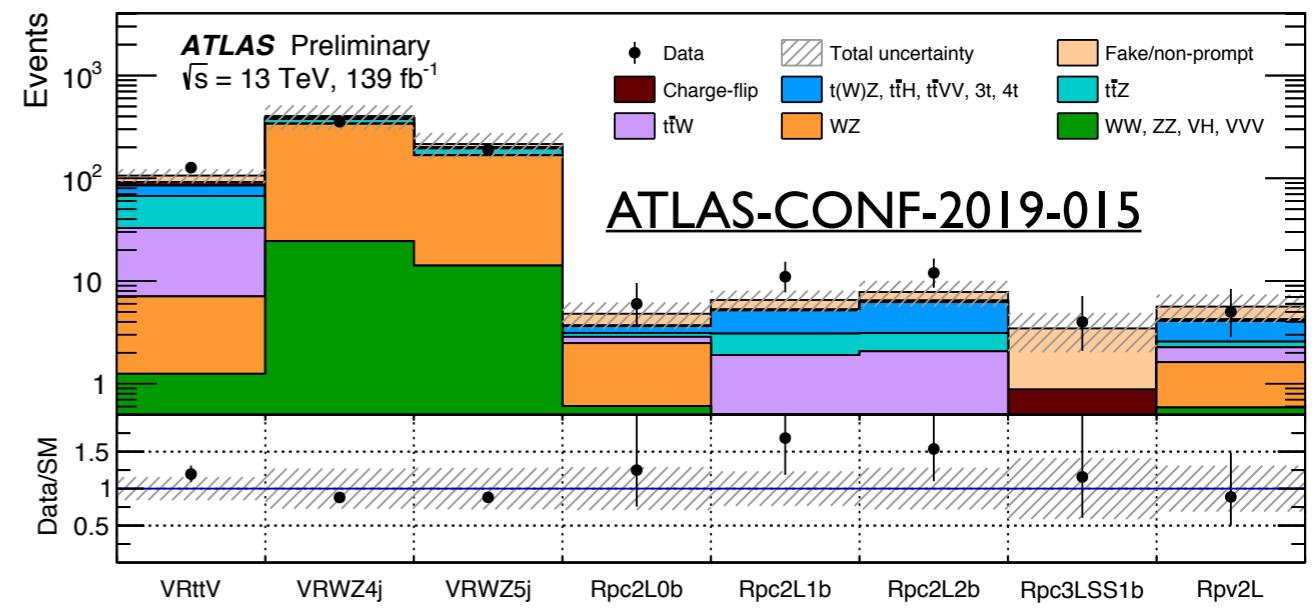
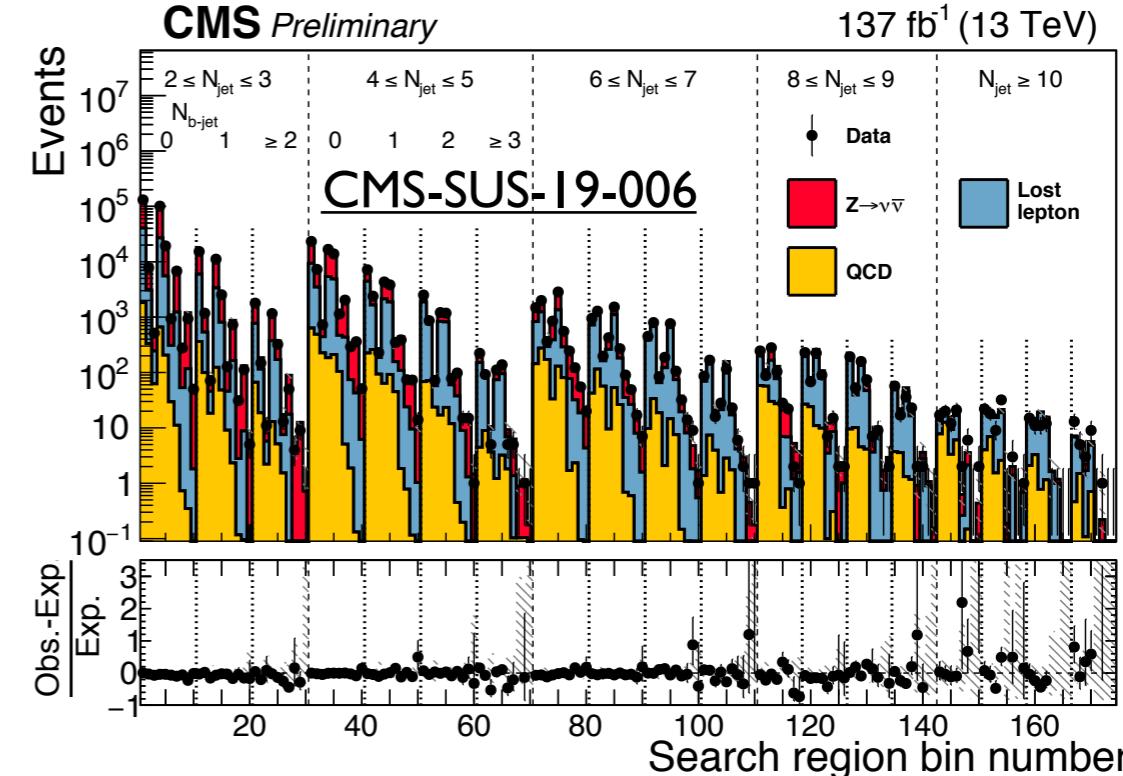
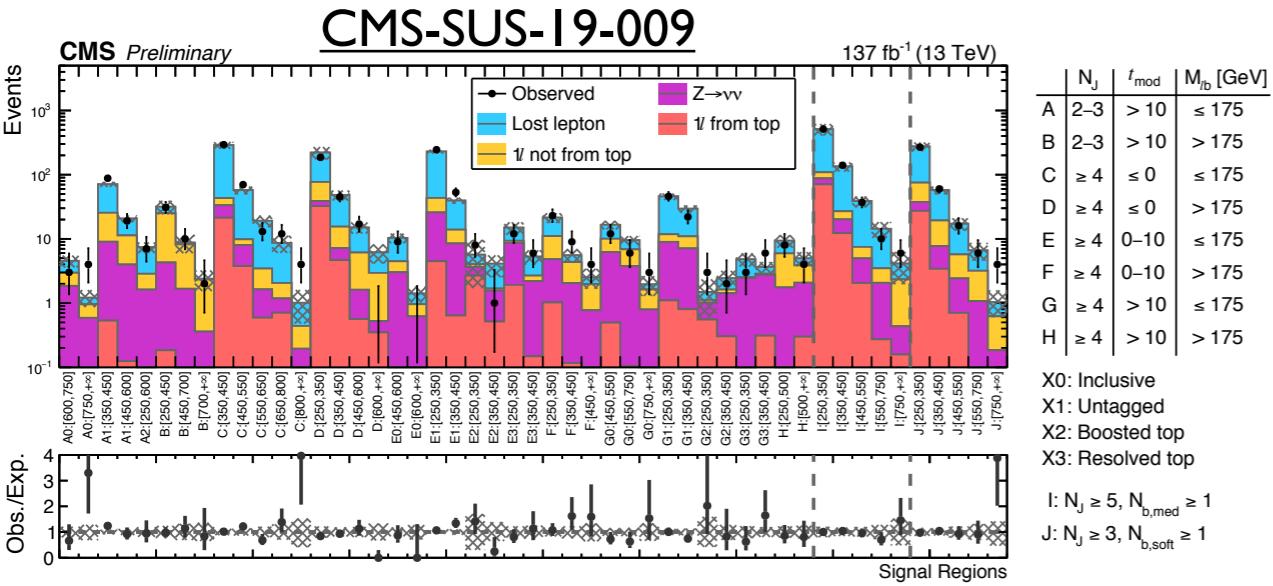




SUSY Signatures

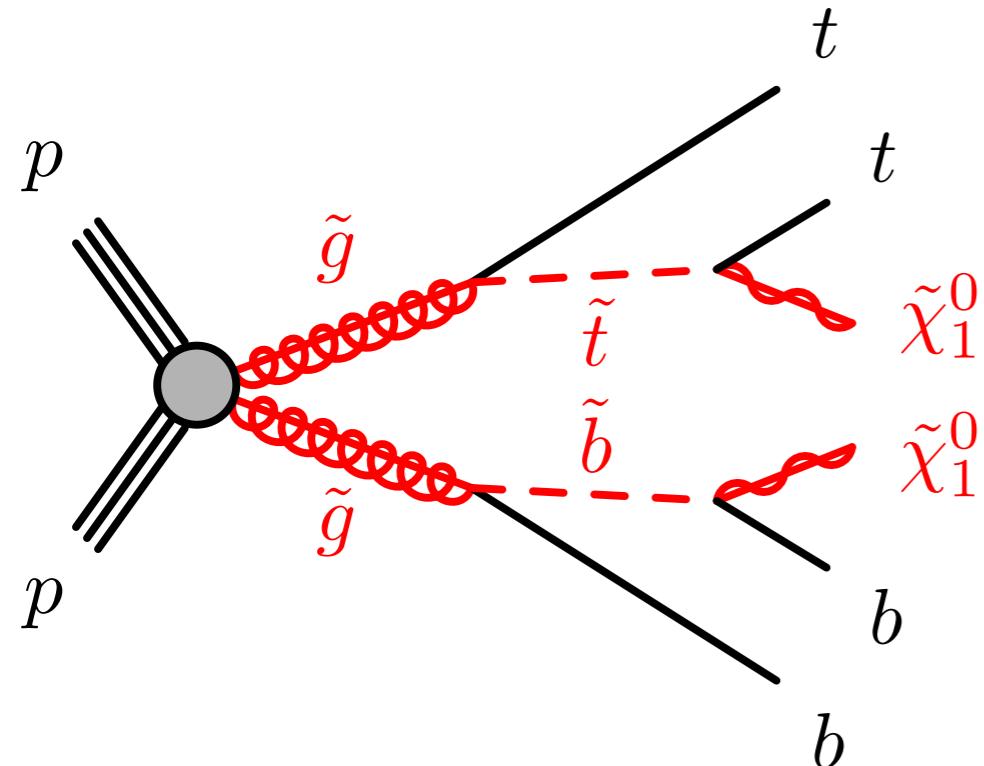


Supersymmetry predicts
more complicated final states...

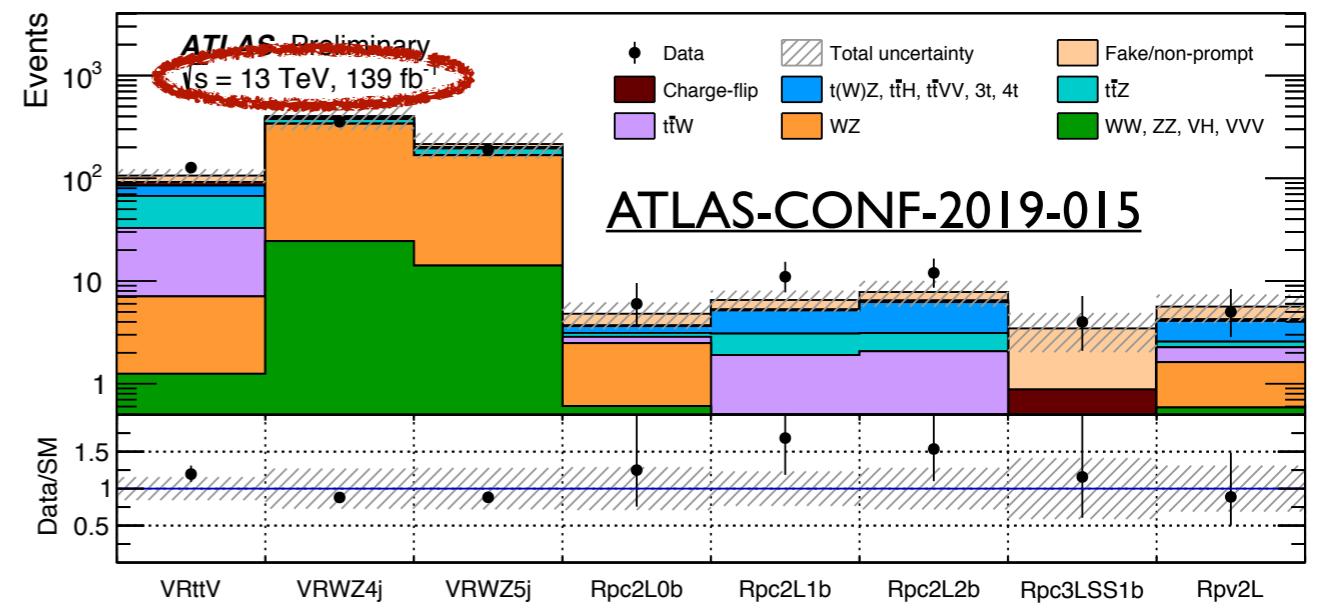
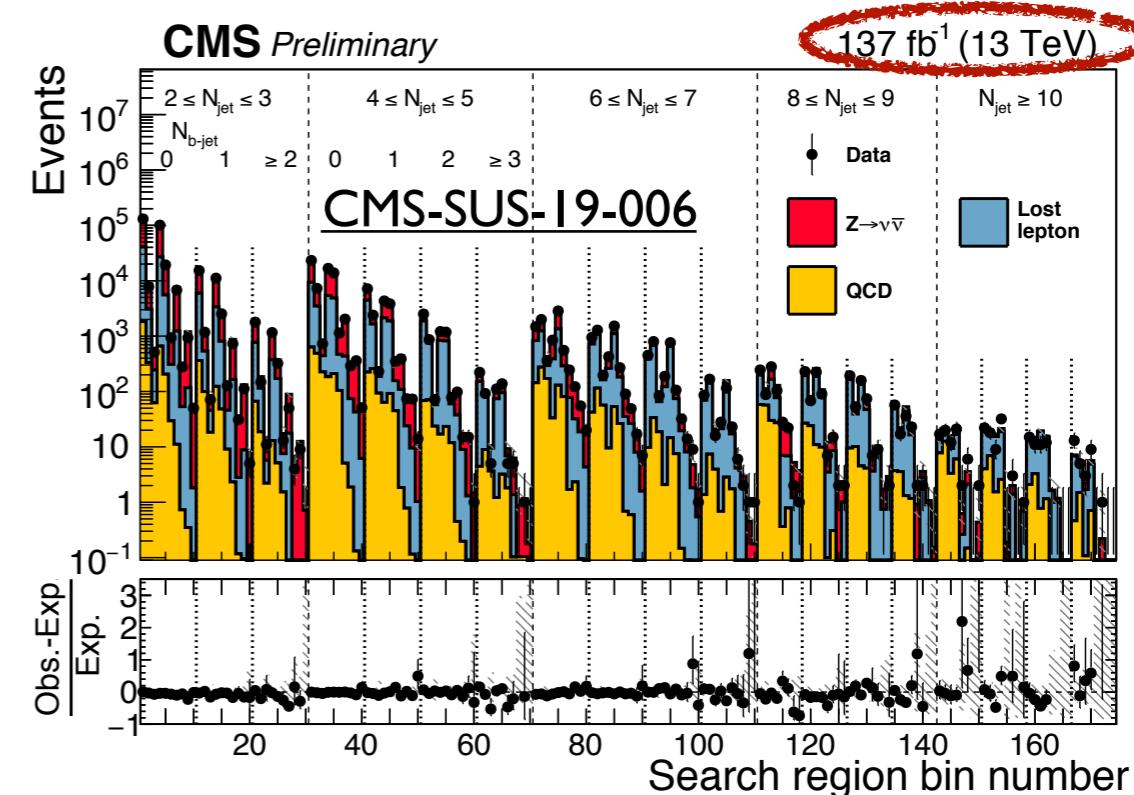
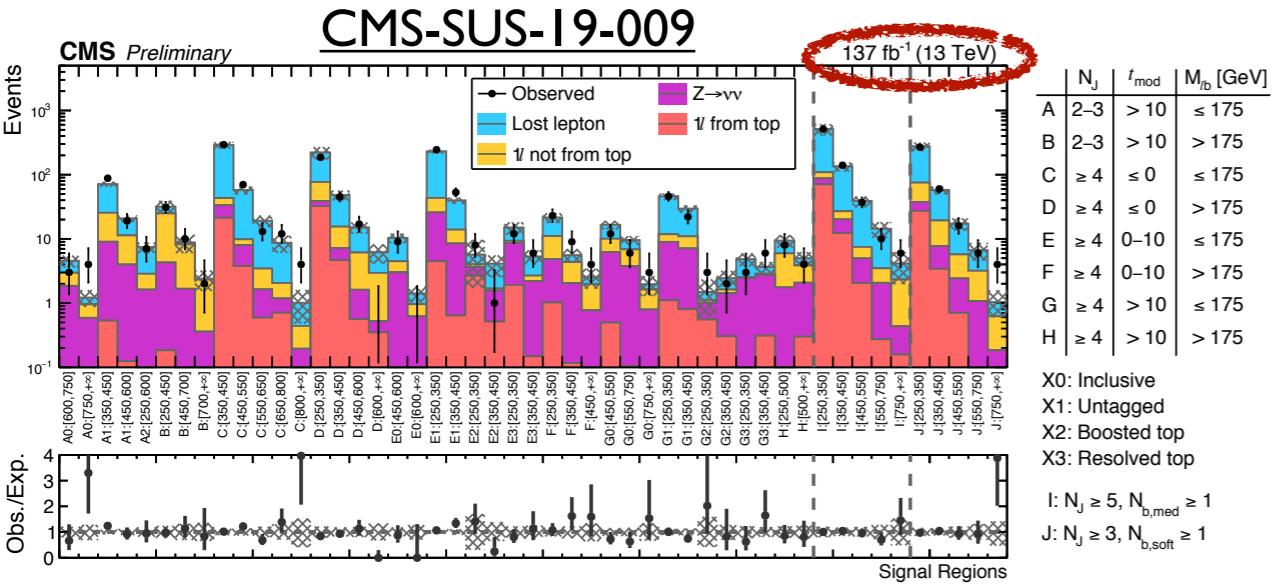




SUSY Signatures



Supersymmetry predicts
more complicated final states...



But still no hints, even with full data!

Our Sensitivity is Better Than Ever



Our Sensitivity is Better Than Ever



ATLAS SUSY Searches* - 95% CL Lower Limits
July 2019

ATLAS Preliminary
 $\sqrt{s} = 13 \text{ TeV}$

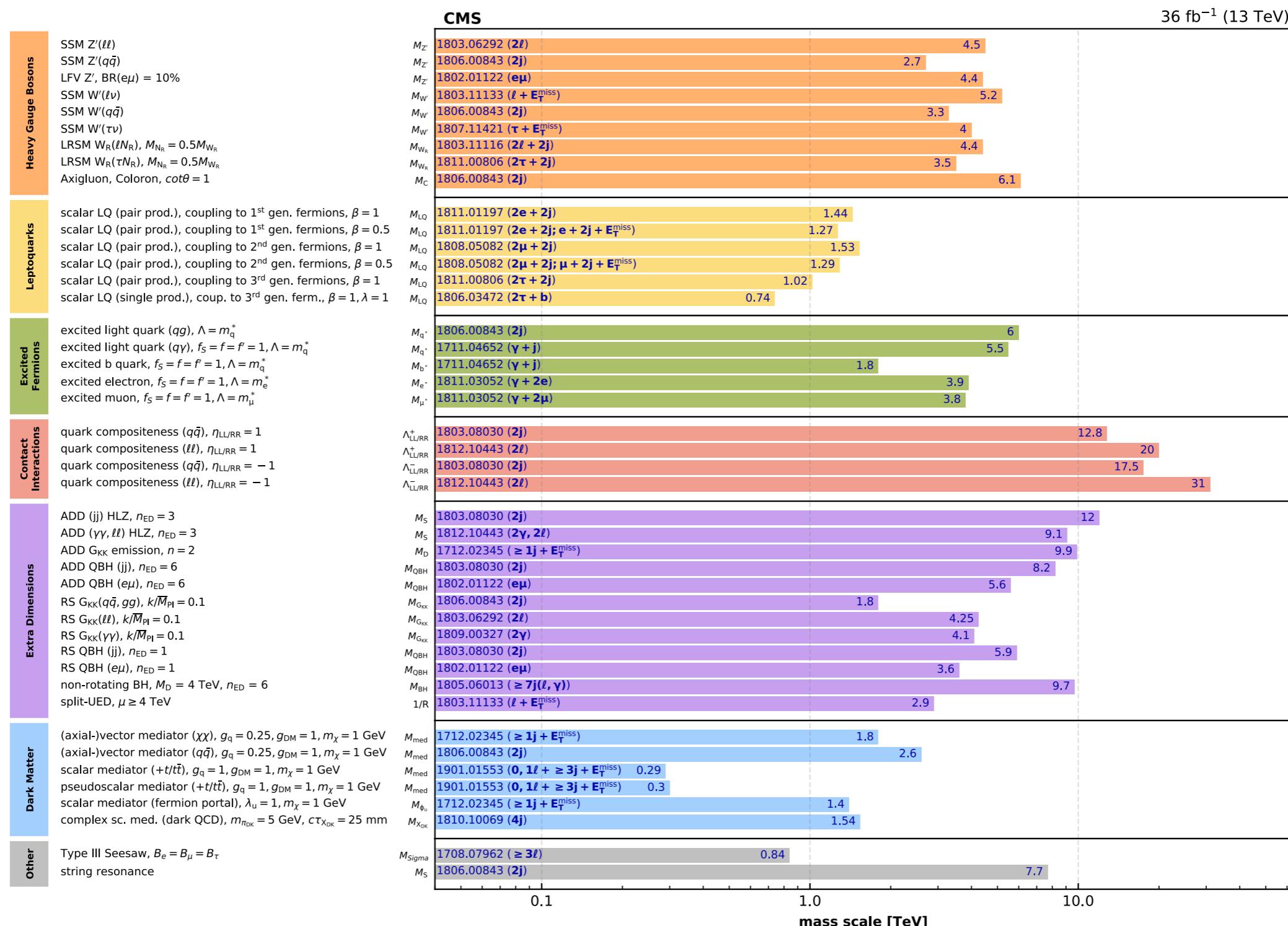
Model	Signature	$\int \mathcal{L} dt [fb^{-1}]$	Mass limit					Reference	
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{g}\rightarrow q\tilde{\chi}_1^0$ mono-jet	0 e, μ E_T^{miss}	2-6 jets 1-3 jets	36.1 36.1	\tilde{q} [2x, 8x Degen.] \tilde{q} [1x, 8x Degen.]	0.43 0.71	0.9 1.55	$m(\tilde{\chi}_1^0) < 100 \text{ GeV}$ $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1712.02332 1711.03301
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow g\tilde{\chi}_1^0$	0 e, μ	2-6 jets	E_T^{miss}	36.1	\tilde{g} \tilde{g}	ForbIDDEN	2.0 0.95-1.6	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 900 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow g\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	3 e, μ $ee, \mu\mu$	4 jets 2 jets	E_T^{miss}	36.1 36.1	\tilde{g} \tilde{g}	1.2	1.85	$m(\tilde{\chi}_1^0) < 800 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow qqWZ\tilde{\chi}_1^0$	0 e, μ SS e, μ	7-11 jets 6 jets	E_T^{miss}	36.1 139	\tilde{g} \tilde{g}	1.15	1.8	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow t\bar{t}\tilde{\chi}_1^0$	0-1 e, μ SS e, μ	3 b 6 jets	E_T^{miss}	79.8 139	\tilde{g} \tilde{g}	1.25	2.25	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300 \text{ GeV}$
3^{rd} gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{\chi}_1^0/\tilde{\nu}_1^\pm$	Multiple		36.1	\tilde{b}_1	ForbIDDEN	0.9		1708.09266, 1711.03301
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$	Multiple		36.1	\tilde{b}_1	ForbIDDEN	0.58-0.82		1708.09266
		Multiple		139	\tilde{b}_1	ForbIDDEN	0.74		ATLAS-CONF-2019-015
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{\chi}_2^0$	0 e, μ	6 b	E_T^{miss}	139	\tilde{b}_1	ForbIDDEN	0.23-0.48	SUSY-2018-31
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow Wb\tilde{\chi}_1^0 \text{ or } \tilde{\chi}_1^0$	0-2 e, μ	0-2 jets/1-2 b	E_T^{miss}	36.1	\tilde{t}_1		1.0	1506.08616, 1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow Wb\tilde{\chi}_1^0$	1 e, μ	3 jets/1 b	E_T^{miss}	139	\tilde{t}_1		0.44-0.59	ATLAS-CONF-2019-017
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1\rightarrow \tilde{G}$	1 $\tau + 1 e, \mu, \tau$	2 jets/1 b	E_T^{miss}	36.1	\tilde{t}_1		1.16	1803.10178
EW direct	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c}\rightarrow c\tilde{\chi}_1^0$	0 e, μ	2 c	E_T^{miss}	36.1	\tilde{c}	0.46	0.85	1805.01649
		0 e, μ	mono-jet	E_T^{miss}	36.1	\tilde{t}_1	0.43		1805.01649
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2\rightarrow \tilde{t}_1 + h$	1-2 e, μ	4 b	E_T^{miss}	36.1	\tilde{t}_2		0.32-0.88	1706.03986
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2\rightarrow \tilde{t}_1 + Z$	3 e, μ	1 b	E_T^{miss}	139	\tilde{t}_2	ForbIDDEN	0.86	ATLAS-CONF-2019-016
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \text{ via } WZ$	2-3 e, μ $ee, \mu\mu$	≥ 1	E_T^{miss} E_T^{miss}	36.1 139	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.205	0.6	1403.5294, 1806.02293
Long-lived particles	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp \text{ via } WW$	2 e, μ		E_T^{miss}	139	$\tilde{\chi}_1^\pm$		0.42	ATLAS-CONF-2019-008
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^0 \text{ via } Wh$	0-1 e, μ	2 b/2 γ	E_T^{miss}	139	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	ForbIDDEN	0.74	ATLAS-CONF-2019-019, ATLAS-CONF-2019-XYZ
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp \text{ via } \tilde{\ell}/\tilde{\nu}$	2 e, μ		E_T^{miss}	139	$\tilde{\chi}_1^\pm$		1.0	ATLAS-CONF-2019-008
	$\tilde{\tau}\tilde{\tau}, \tilde{\tau}\rightarrow \tau\tilde{\chi}_1^0$	2 τ		E_T^{miss}	139	$\tilde{\tau}$ [$\tilde{\tau}_L, \tilde{\tau}_R, L$]	0.16-0.3	0.12-0.39	ATLAS-CONF-2019-018
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell}\rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0 jets	E_T^{miss}	139	$\tilde{\ell}$		0.256	ATLAS-CONF-2019-008
		2 e, μ	≥ 1	E_T^{miss}	139	$\tilde{\ell}$			ATLAS-CONF-2019-014
	$\tilde{H}\tilde{H}, \tilde{H}\rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ	≥ 3 b	E_T^{miss}	36.1	\tilde{H}	0.13-0.23	0.29-0.88	1806.04030
		4 e, μ	0 jets	E_T^{miss}	36.1	\tilde{H}	0.3		1804.03602
RPV	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	E_T^{miss}	36.1	$\tilde{\chi}_1^\pm$ $\tilde{\chi}_1^\pm$	0.15	0.46	Pure Wino Pure Higgsino
	Stable \tilde{g} R-hadron		Multiple		36.1	\tilde{g}		2.0	1712.02118 ATL-PHYS-PUB-2017-019
	Metastable \tilde{g} R-hadron, $\tilde{g}\rightarrow qq\tilde{\chi}_1^0$		Multiple		36.1	\tilde{g} [$\tau(\tilde{g})=10 \text{ ns}, 0.2 \text{ ns}$]		2.05 2.4	1902.01636, 1808.04095 1710.04901, 1808.04095
	$L\text{F} V pp\rightarrow \tilde{v}_\tau + X, \tilde{v}_\tau\rightarrow e\mu/e\tau/\mu\tau$	$e\mu, e\tau, \mu\tau$			3.2	\tilde{v}_τ		1.9	$\lambda'_{311}=0.11, \lambda_{132/133/233}=0.07$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\ell\ell\nu\nu$	4 e, μ	0 jets	E_T^{miss}	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [$\lambda_{133}\neq 0, \lambda_{12k}\neq 0$]	0.82	1.33	$m(\tilde{\chi}_1^0)=100 \text{ GeV}$
RPV	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0\rightarrow qqq$	4-5 large- R jets		E_T^{miss}	36.1	\tilde{g} [$m(\tilde{\chi}_1^0)=200 \text{ GeV}, 1100 \text{ GeV}$]	1.05	1.3	Large λ'_{112}
		Multiple		E_T^{miss}	36.1	\tilde{g} [$\lambda''_{112}=2e-4, 2e-5$]		1.9 2.0	$m(\tilde{\chi}_1^0)=200 \text{ GeV}$, bino-like
	$\tilde{\tau}, \tilde{\tau}\rightarrow \tilde{\ell}\tilde{\chi}_1^0, \tilde{\chi}_1^0\rightarrow tbs$	Multiple			36.1	\tilde{g} [$\lambda'_{323}=2e-4, 1e-2$]	0.55	1.05	$m(\tilde{\chi}_1^0)=200 \text{ GeV}$, bino-like
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow bs$	2 jets + 2 b			36.7	\tilde{t}_1 [qq, bs]	0.42	0.61	ATLAS-CONF-2018-003
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow q\ell$	2 e, μ	2 b		36.1	\tilde{t}_1 [$1e-10 < \lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9$]	1.0	0.4-1.45	1710.07171 1710.05544
		1 μ	DV		136	\tilde{t}_1 [$1e-10 < \lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9$]		1.6	ATLAS-CONF-2019-006

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Our Sensitivity is Better Than Ever



Overview of CMS EXO results

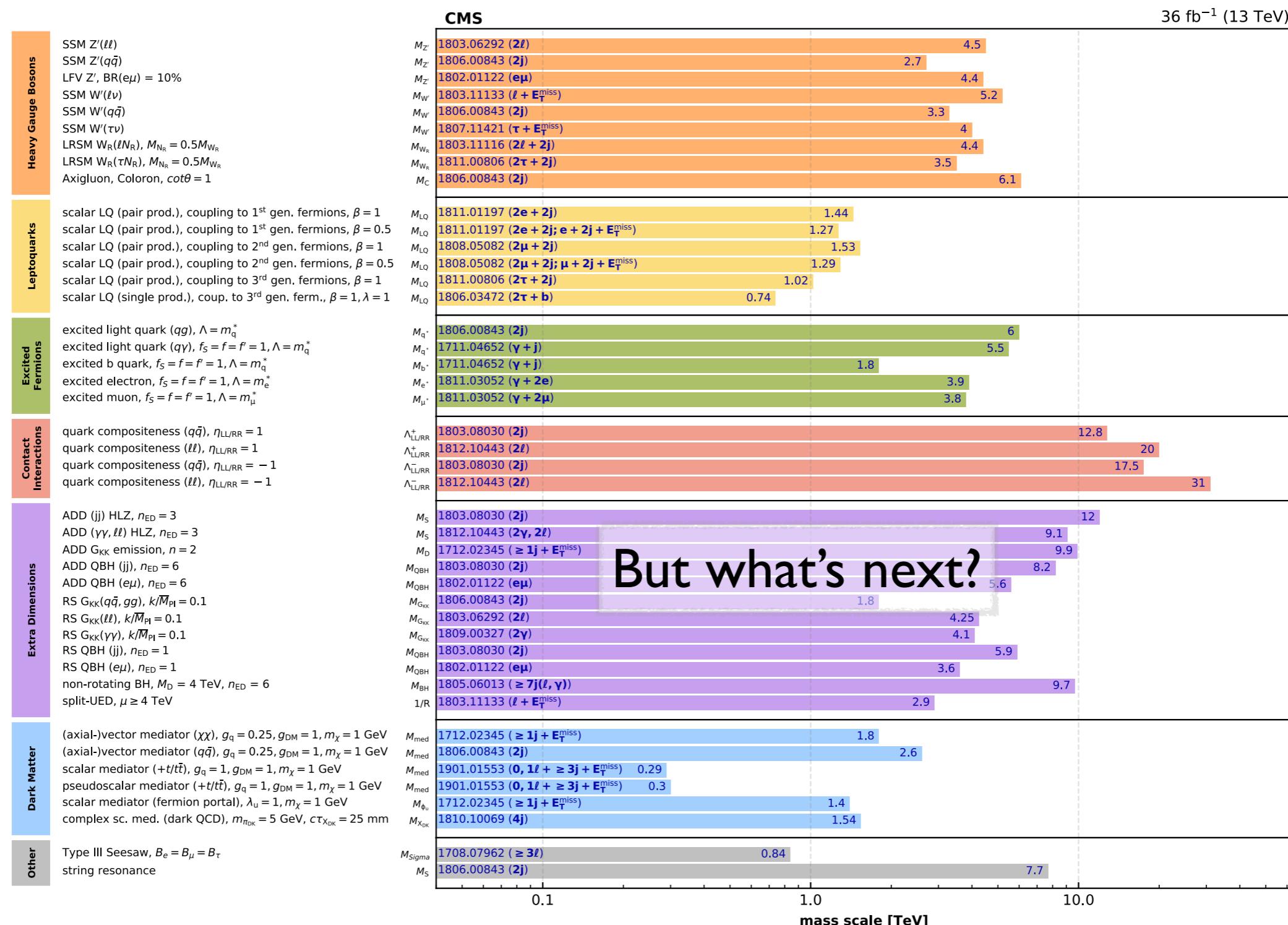


Selection of observed exclusion limits at 95% C.L. (theory uncertainties are not included).

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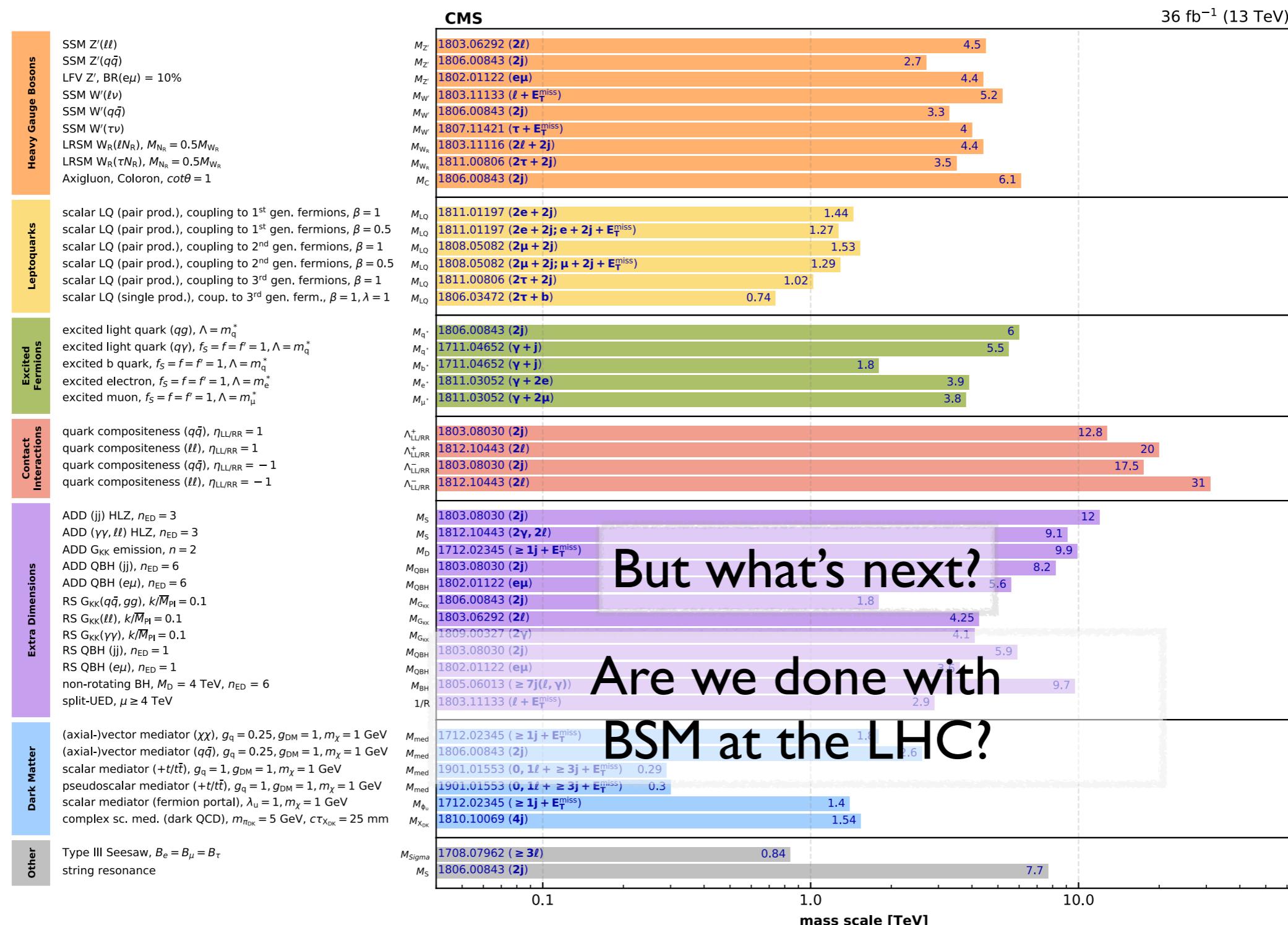
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January 2019

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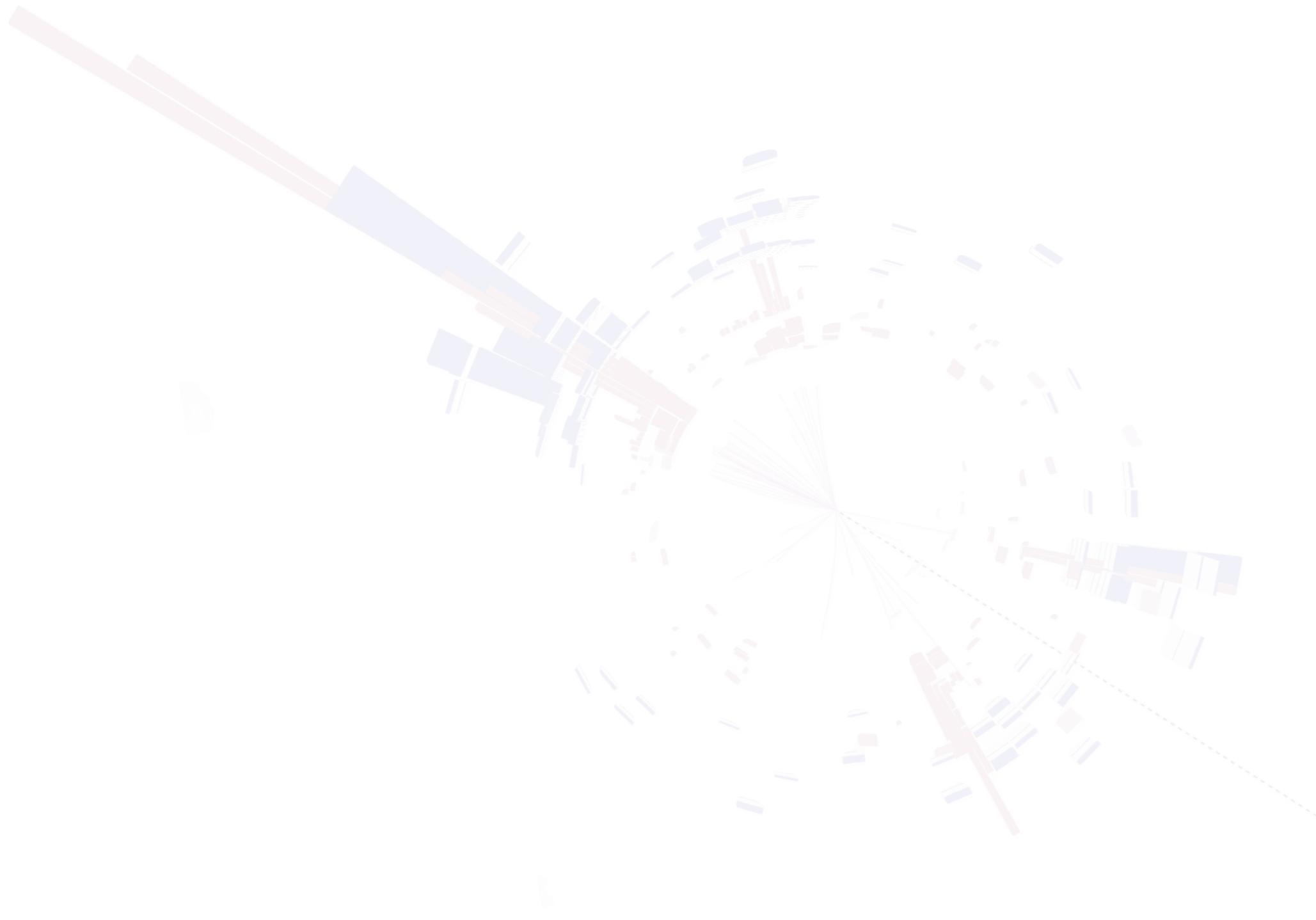
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Low Hanging Fruit Records



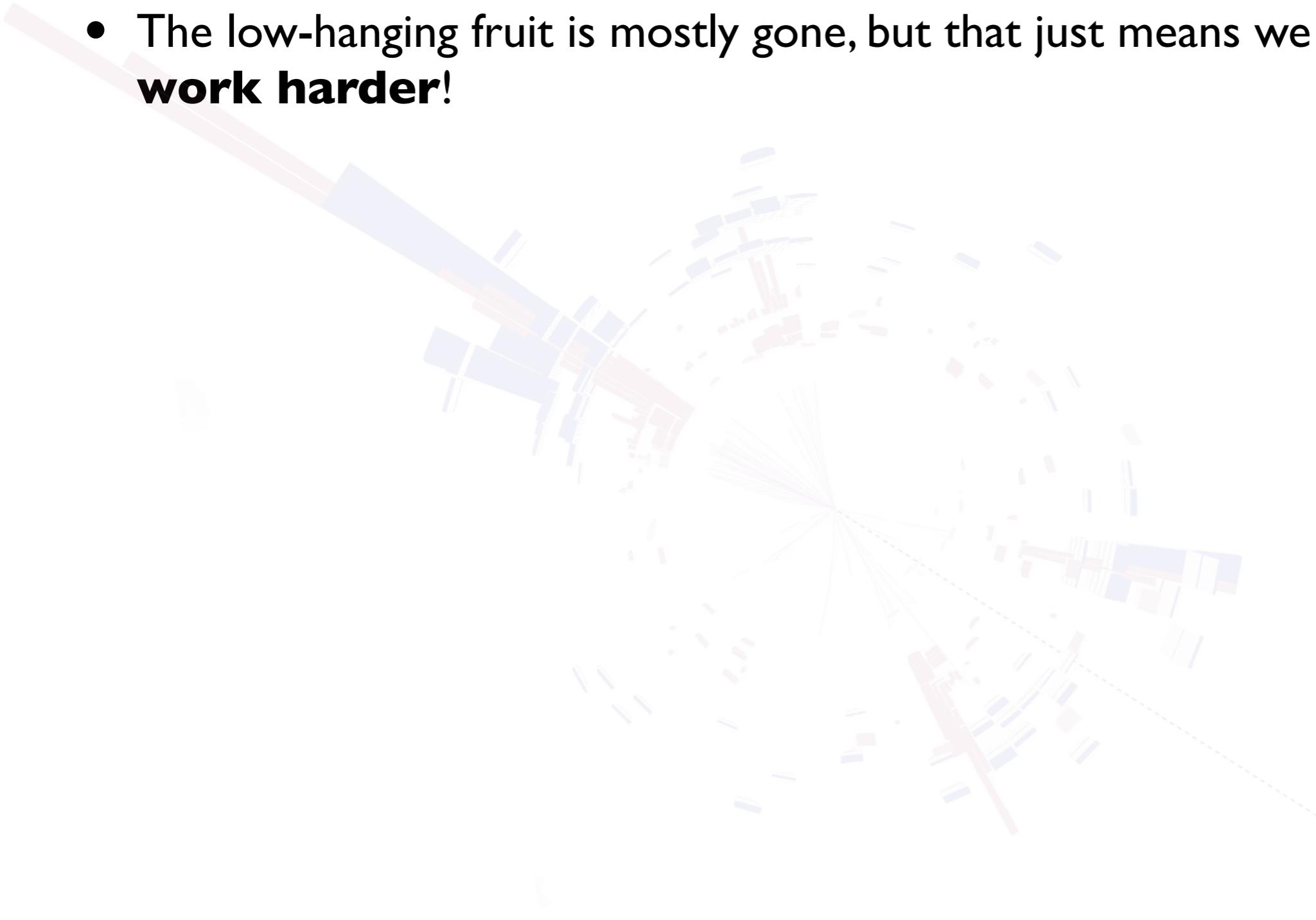
BSM is More Vibrant Than Ever!





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- The low-hanging fruit is mostly gone, but that just means we have to **work harder!**





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 - My apologies for omitting many interesting and exciting results!

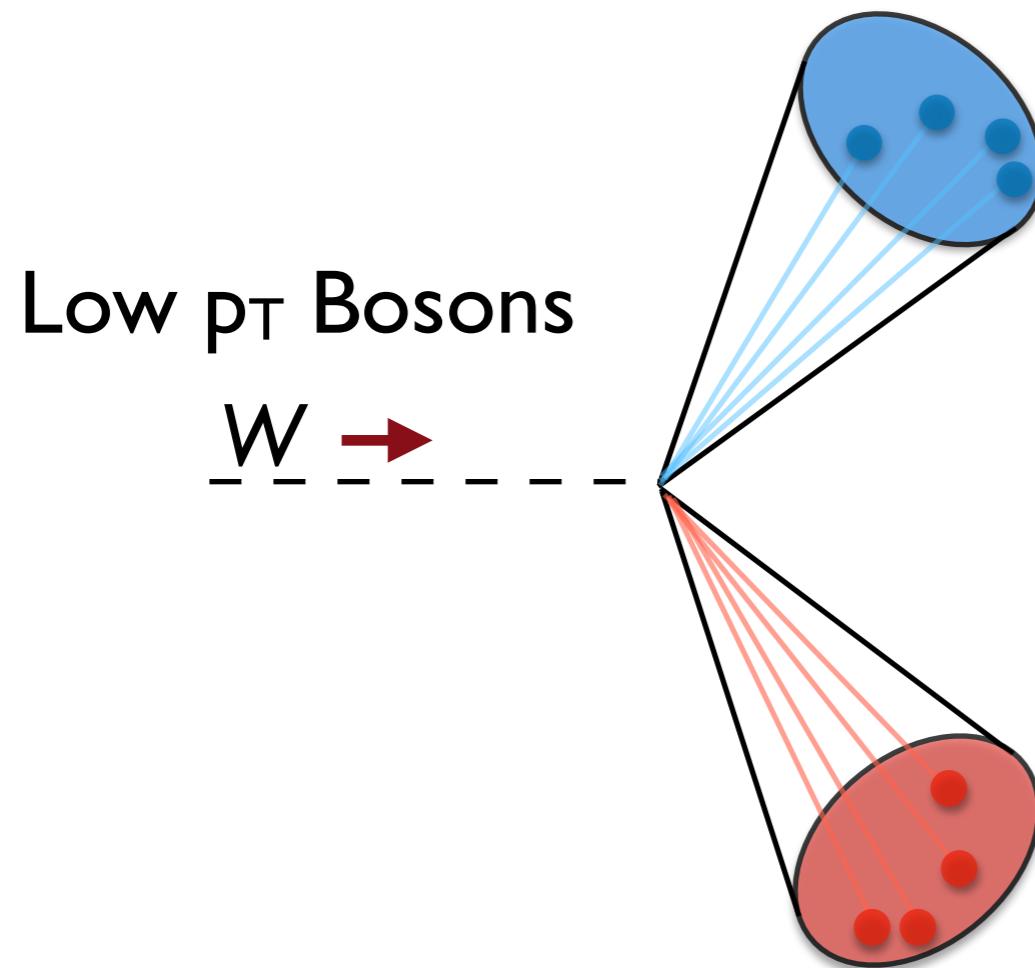
Squeeze Every Drop

*Or: how advances in reconstruction,
triggering, and machine learning are unlocking
new insights into BSM*

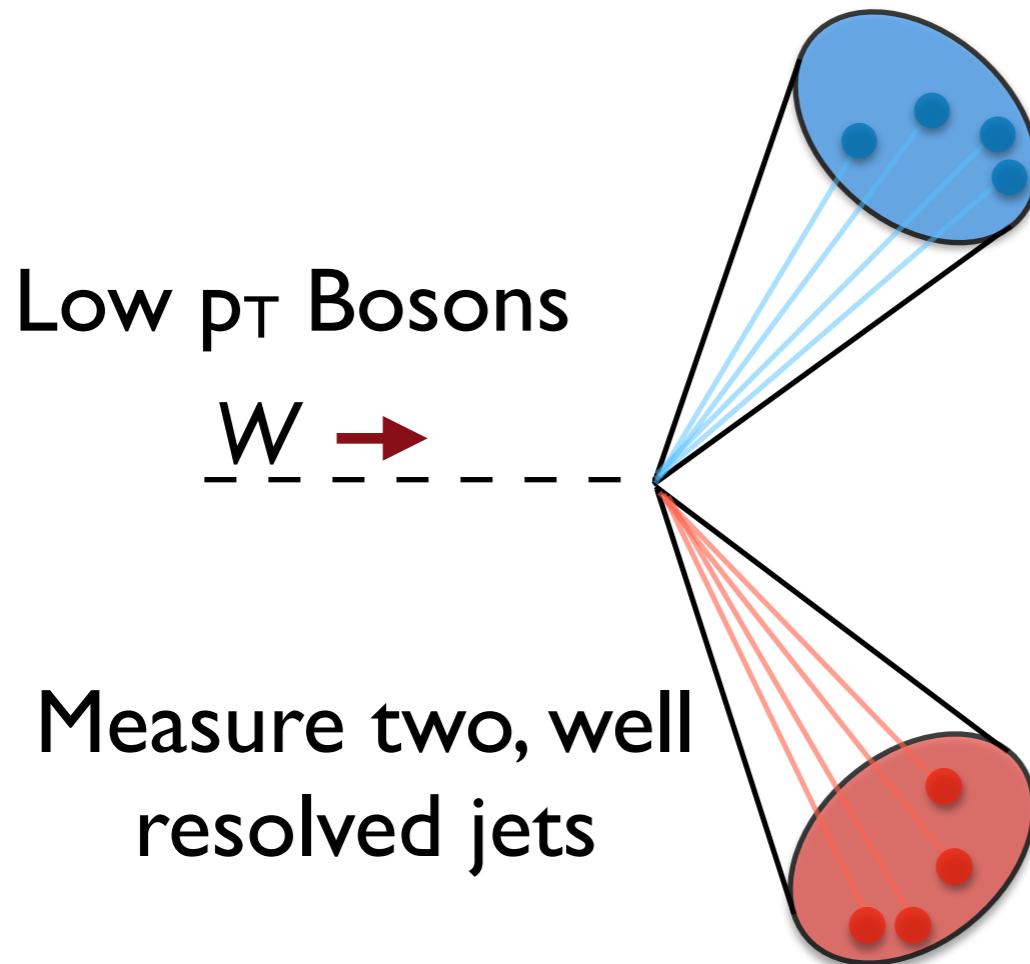
The Challenge of High PT



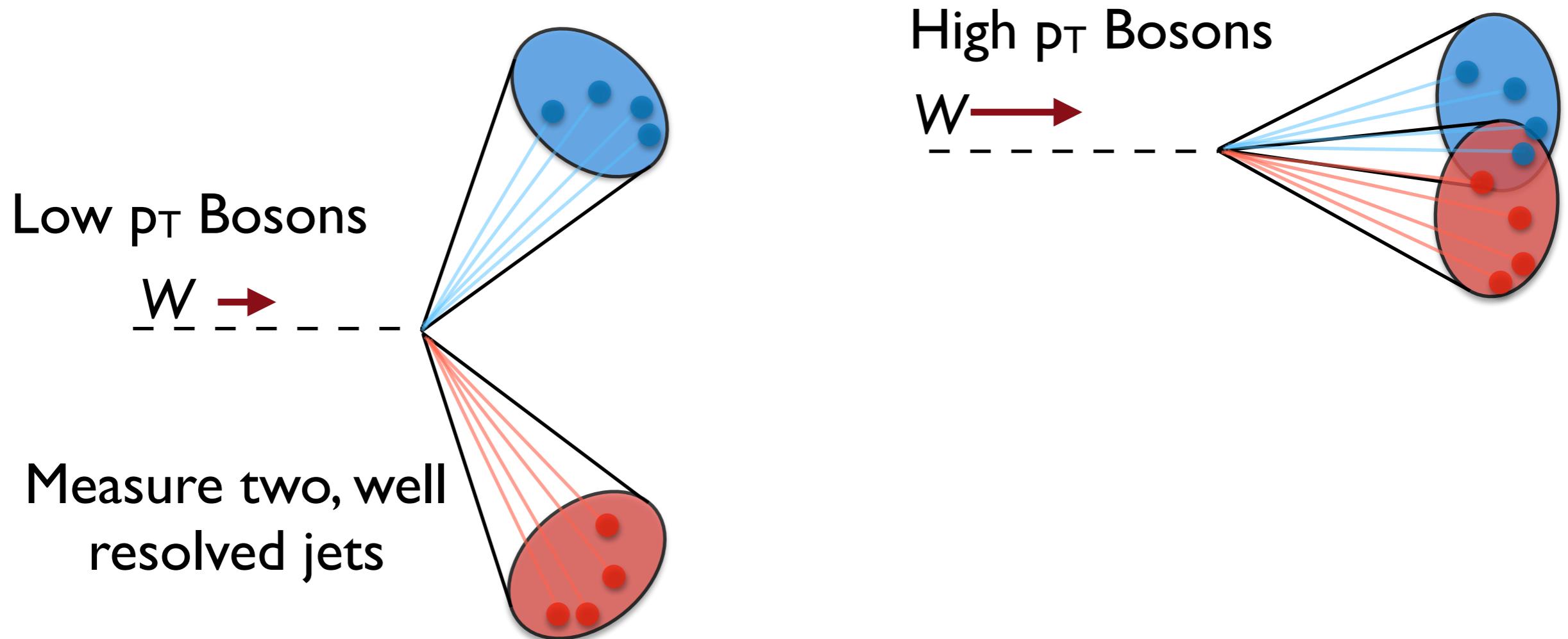
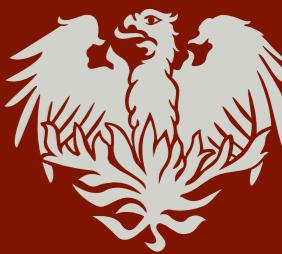
The Challenge of High p_T



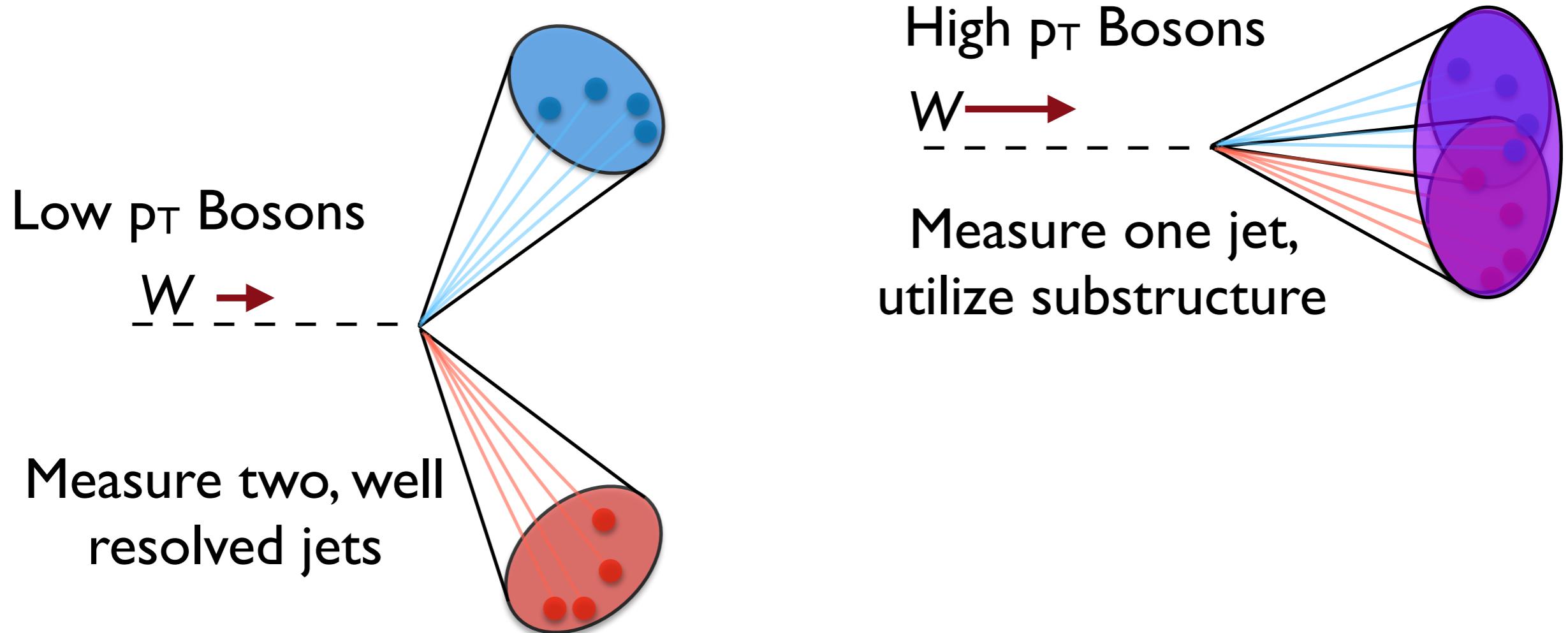
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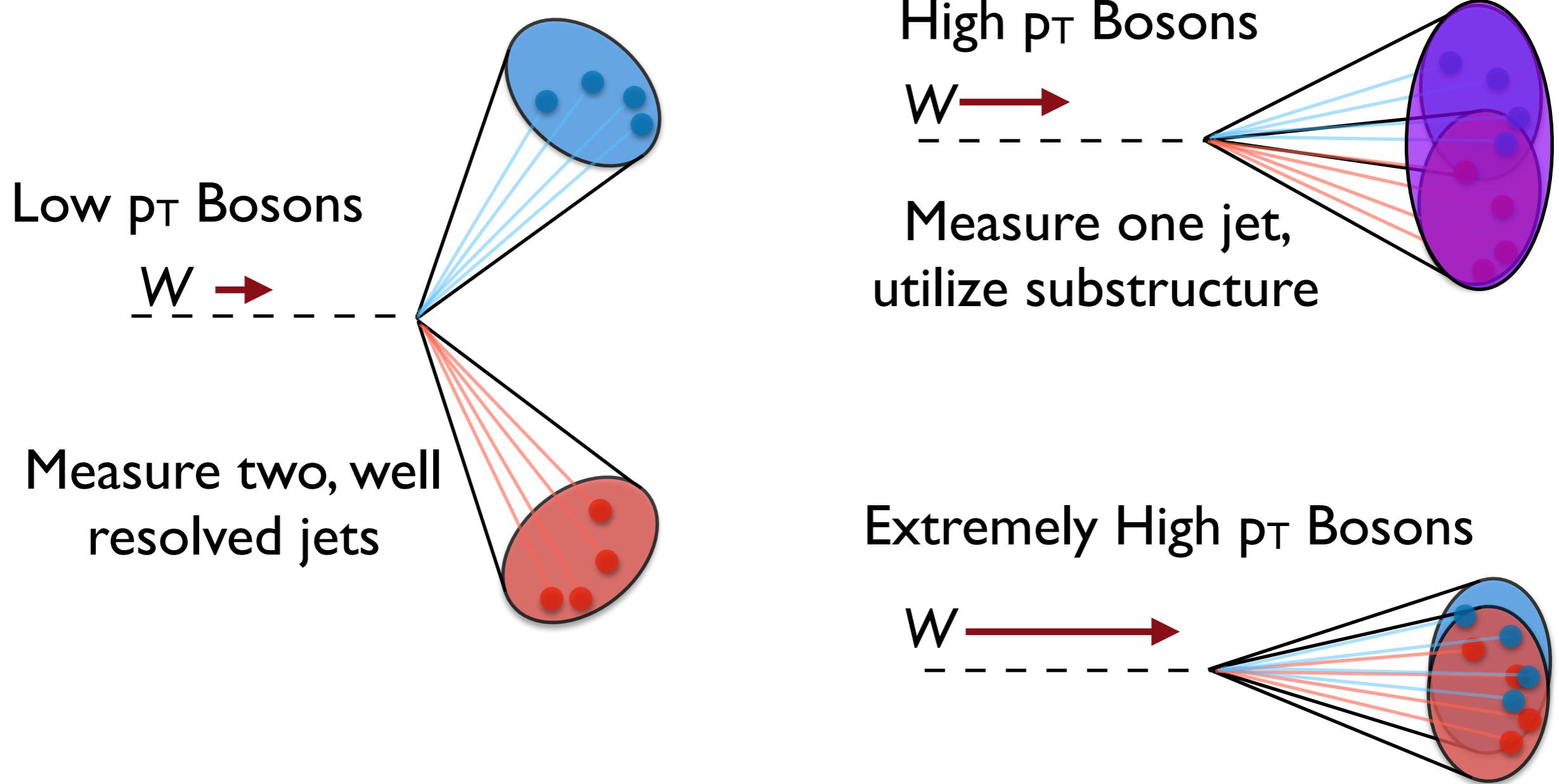
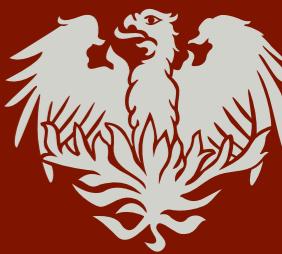
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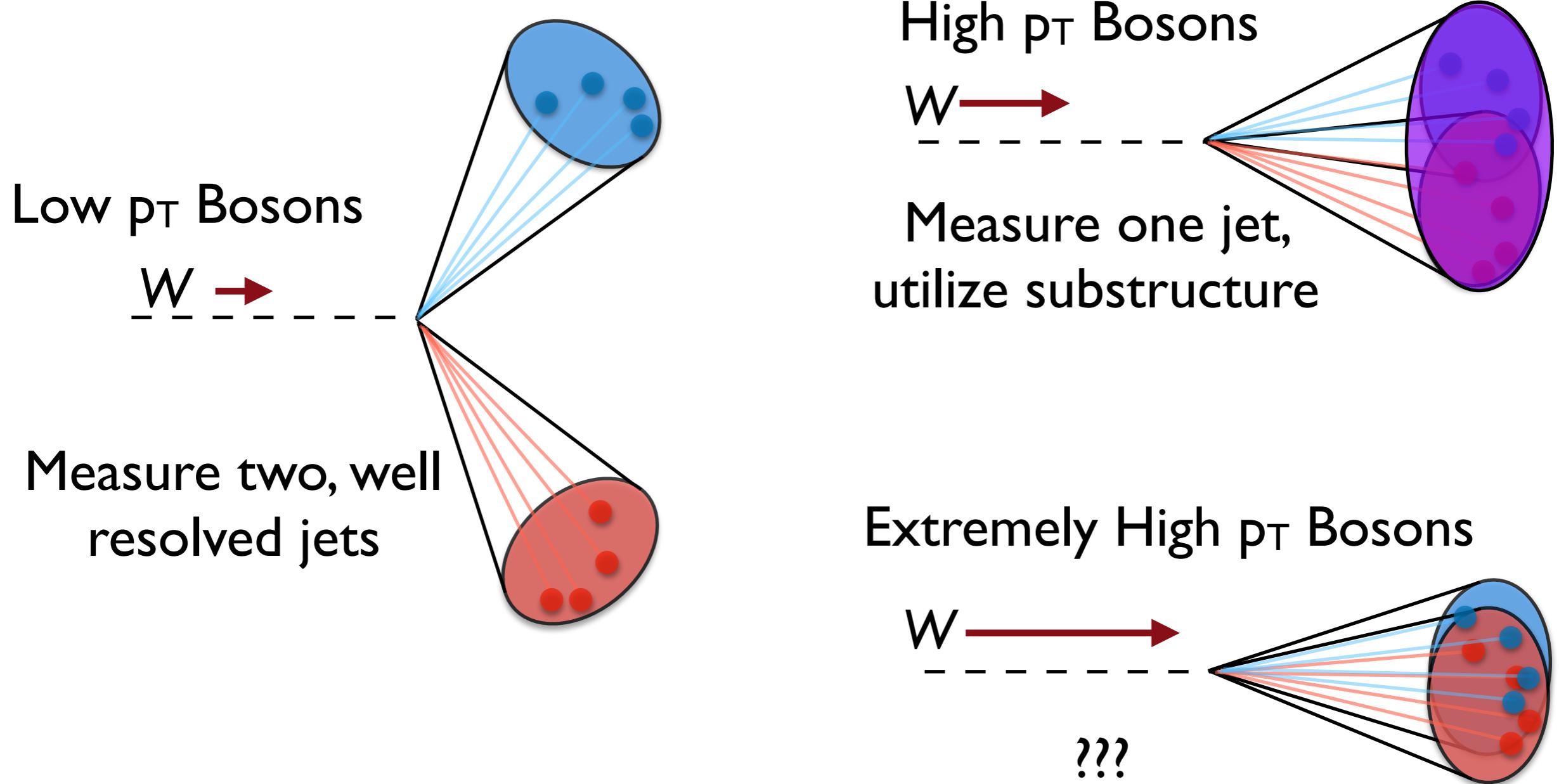
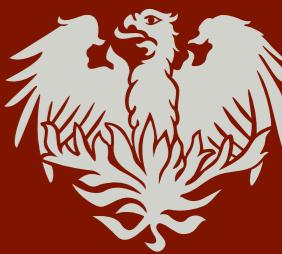
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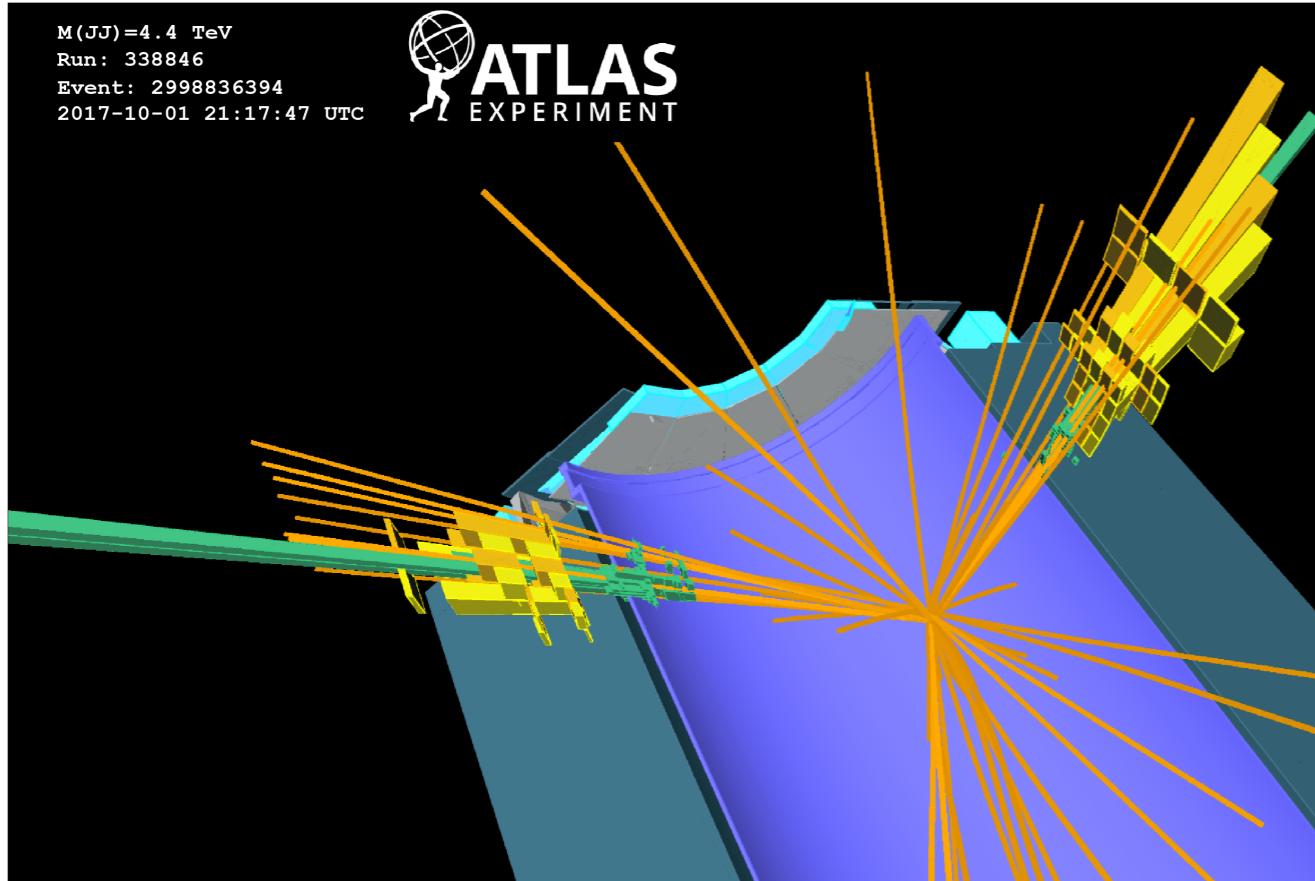




Tagging at High pT

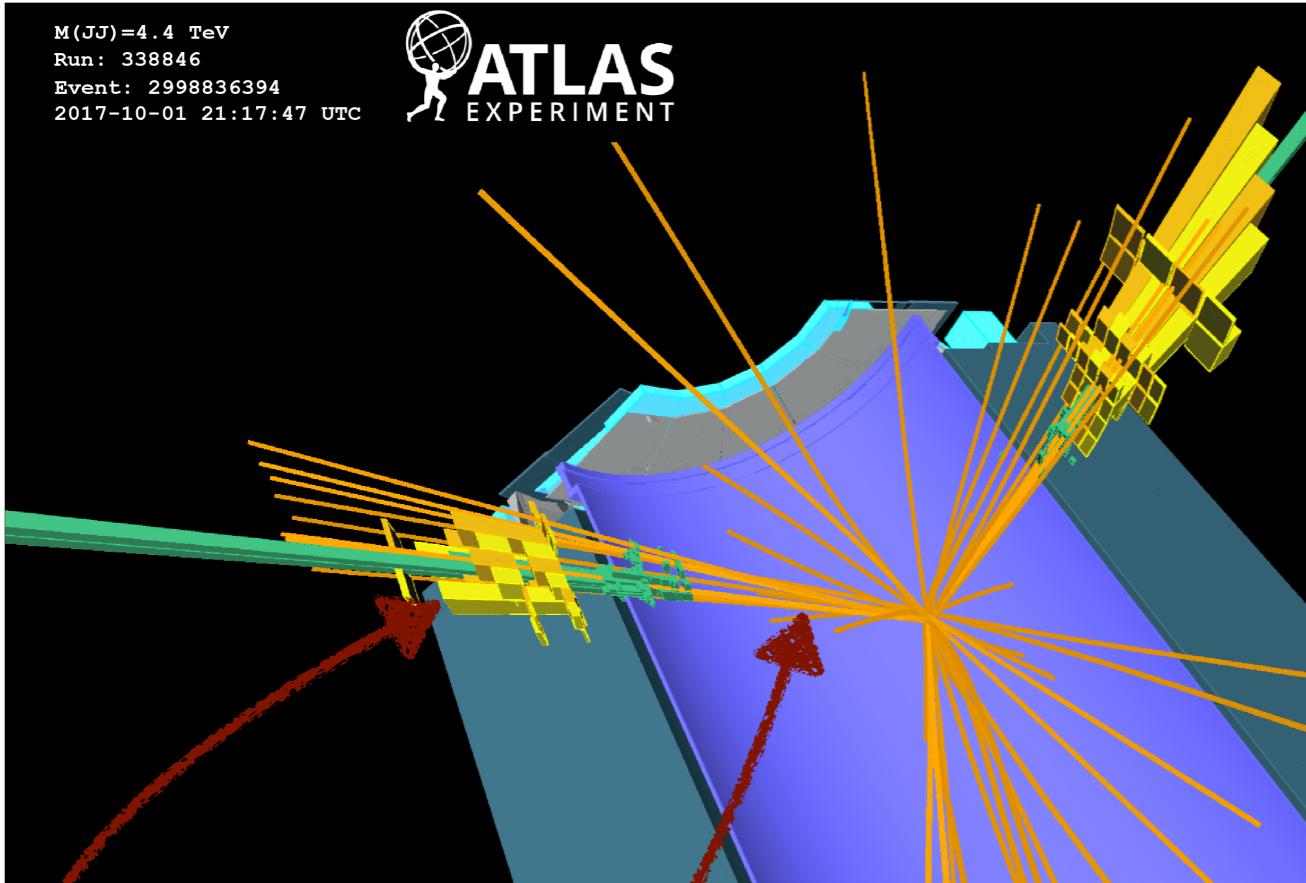


Tagging at High PT





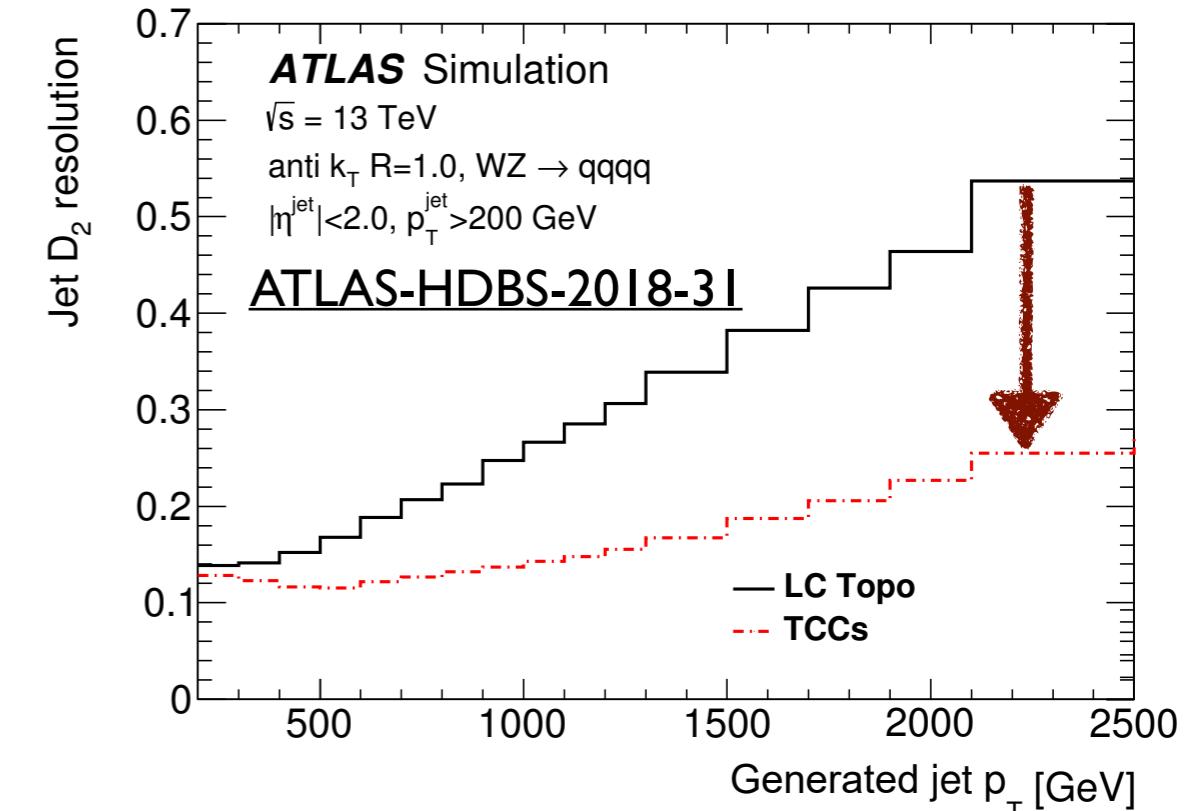
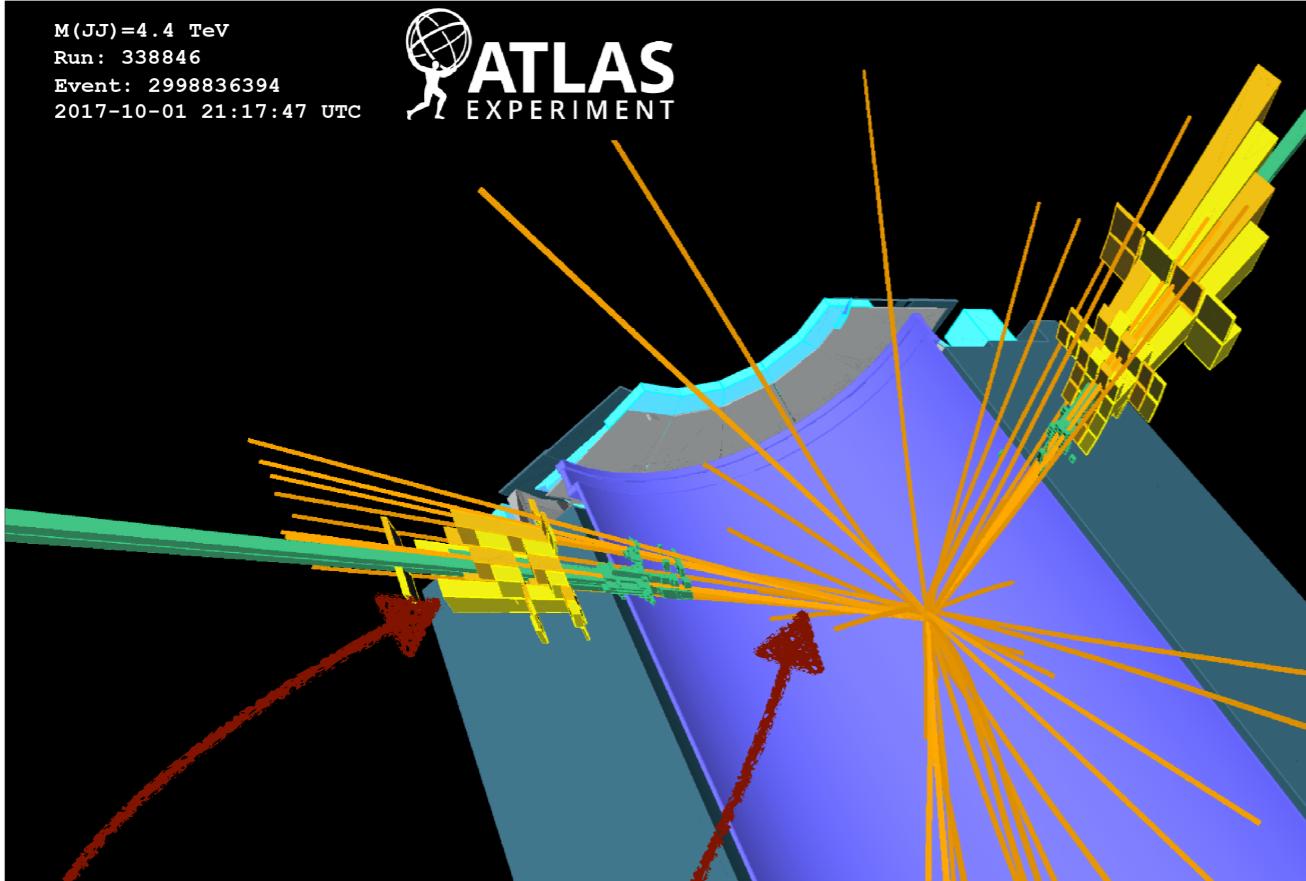
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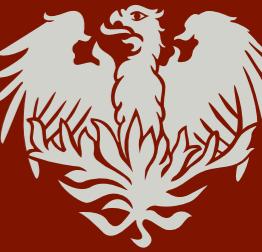
Utilize better spatial resolution
from tracker to separate energy
deposits in the calorimeter!



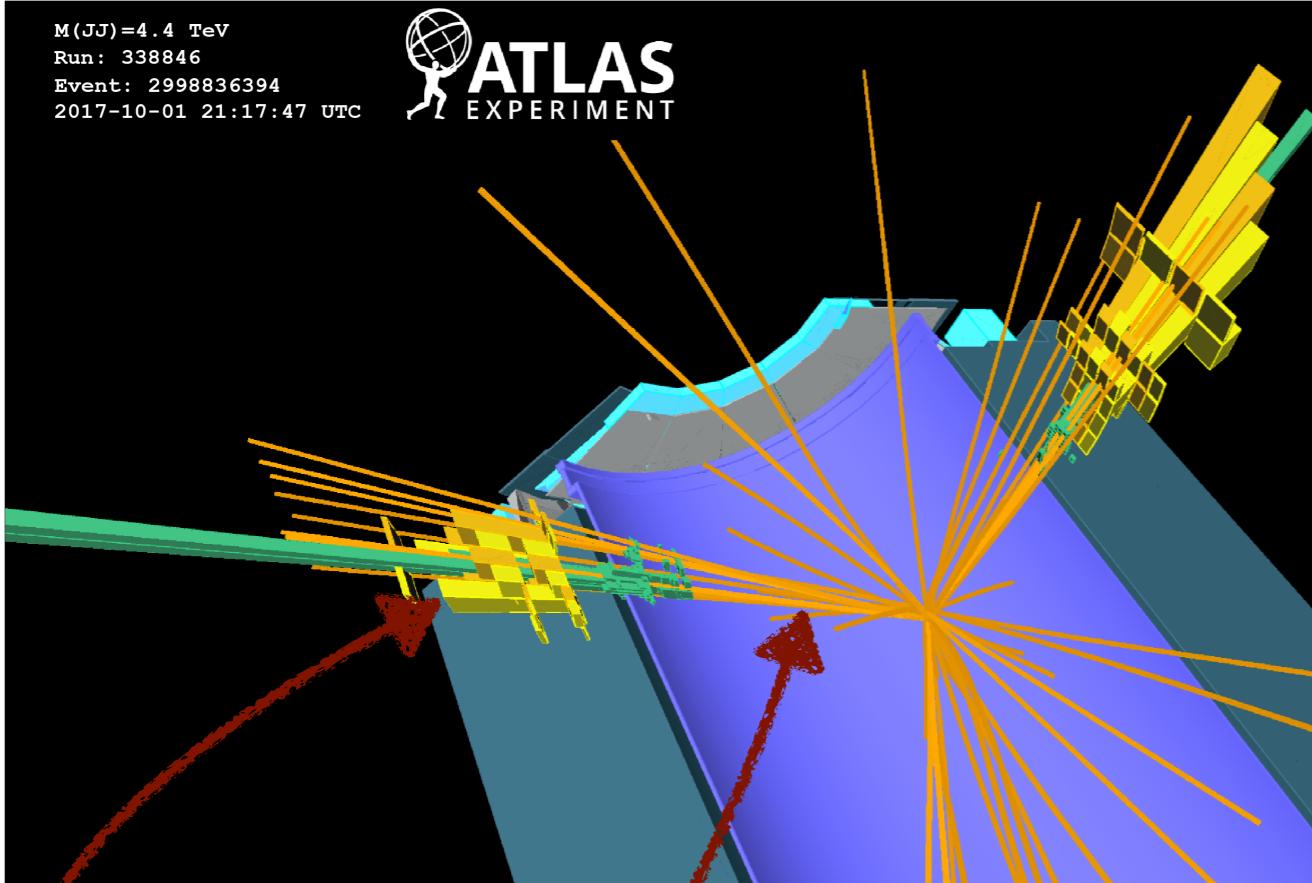
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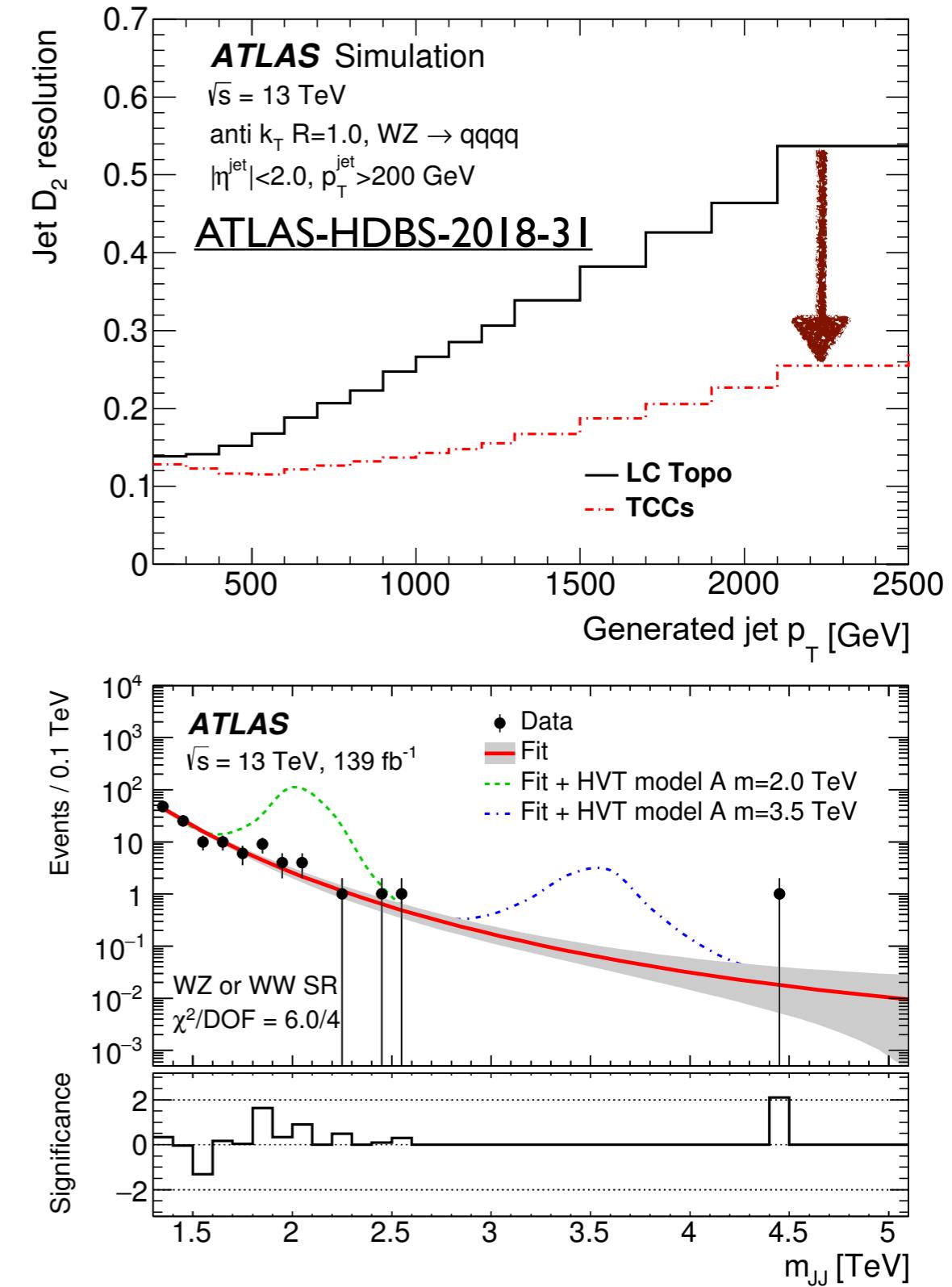


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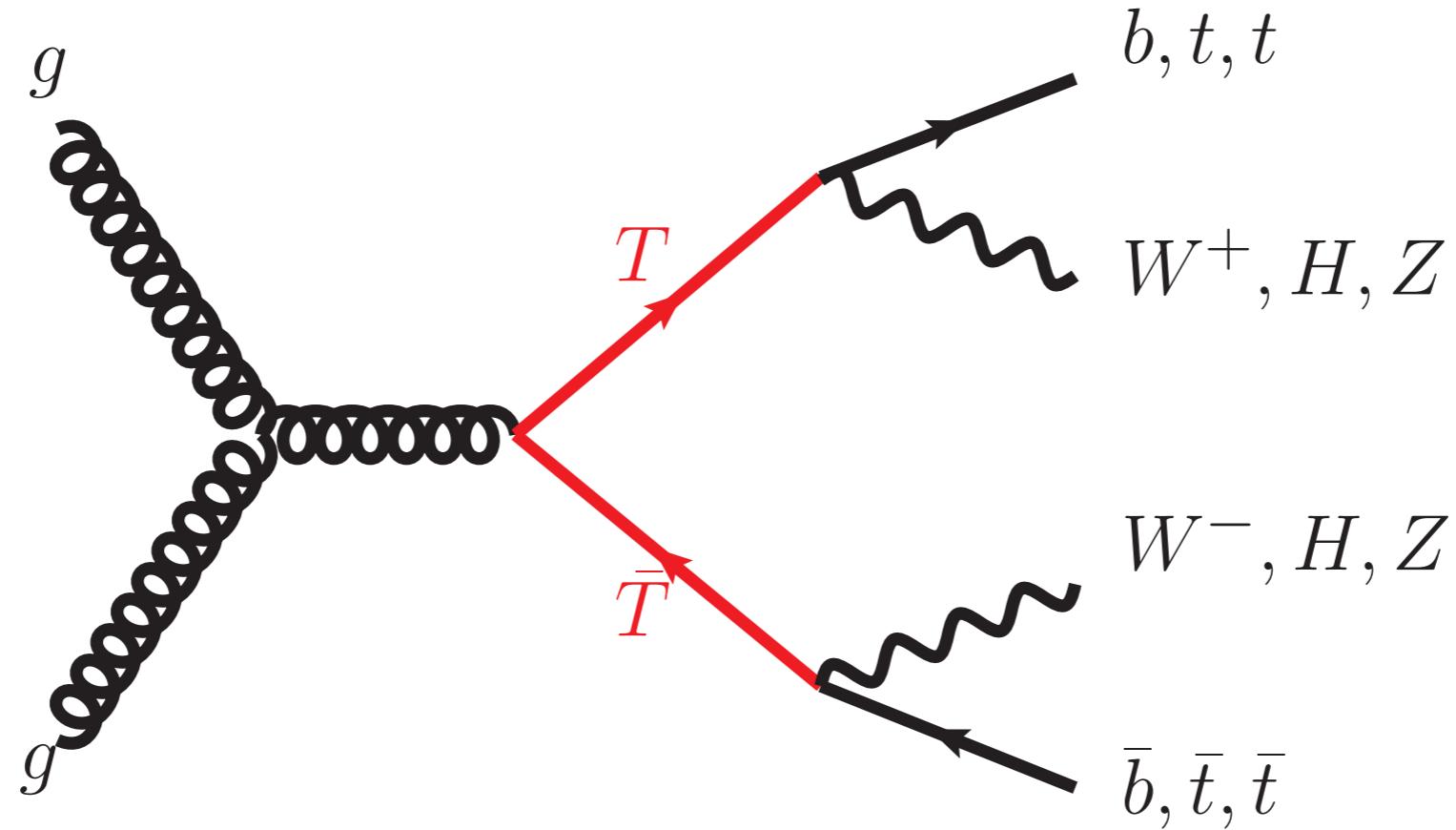
Enables strongest sensitivity yet
to boosted all-hadronic final states



The Challenge of Branching

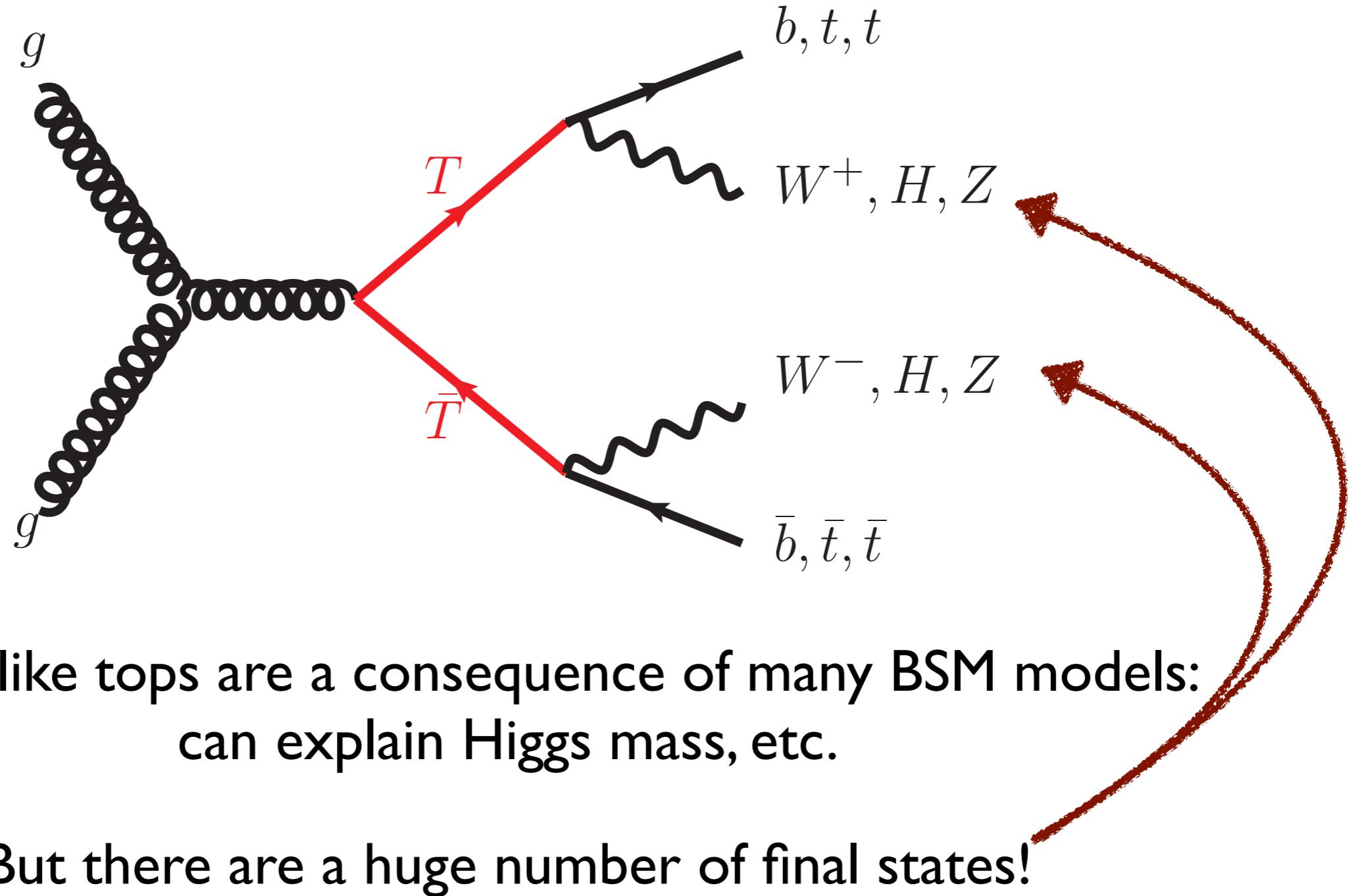


The Challenge of Branching

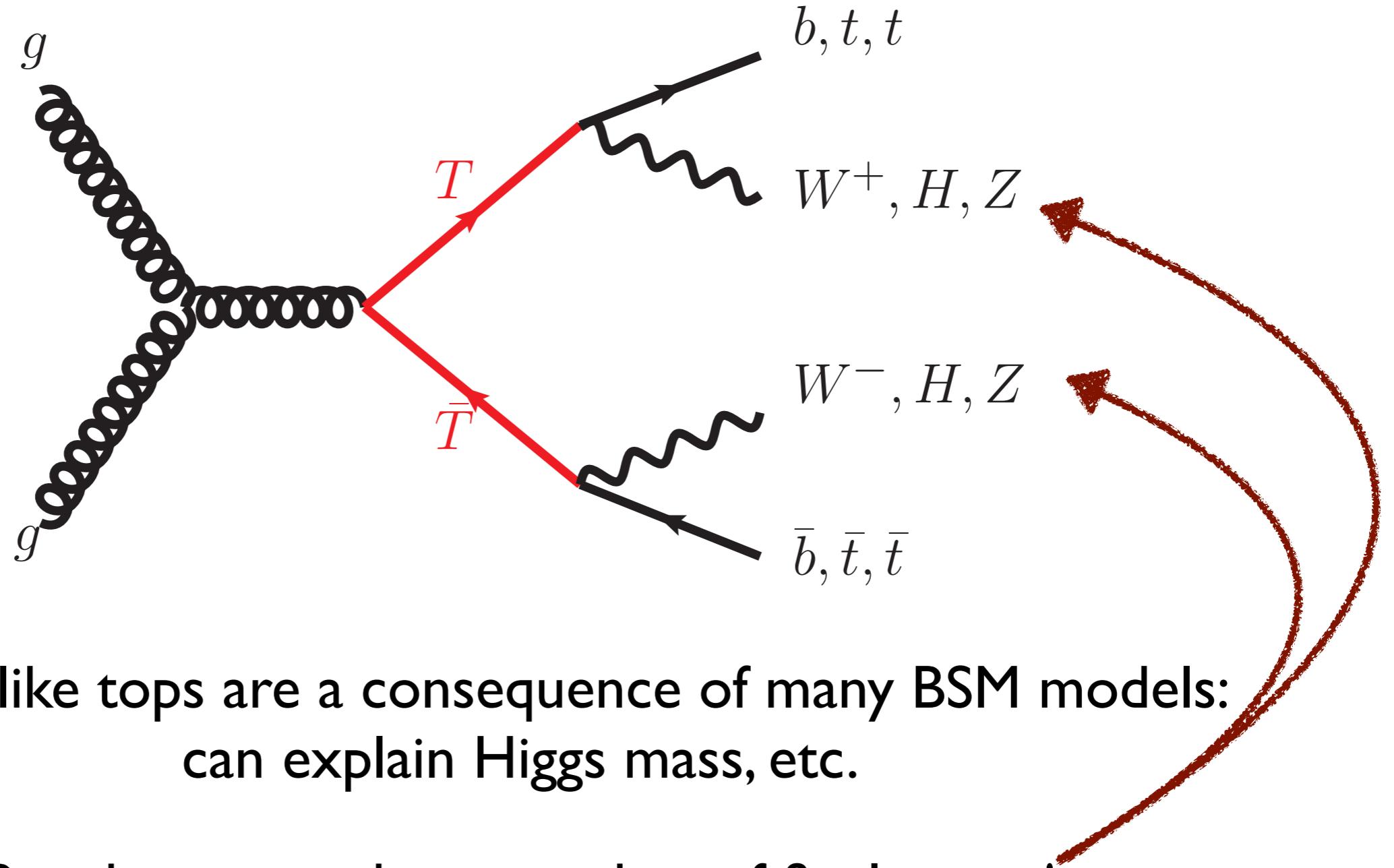


Vector-like tops are a consequence of many BSM models:
can explain Higgs mass, etc.

The Challenge of Branching



The Challenge of Branching



Vector-like tops are a consequence of many BSM models:
can explain Higgs mass, etc.

But there are a huge number of final states!

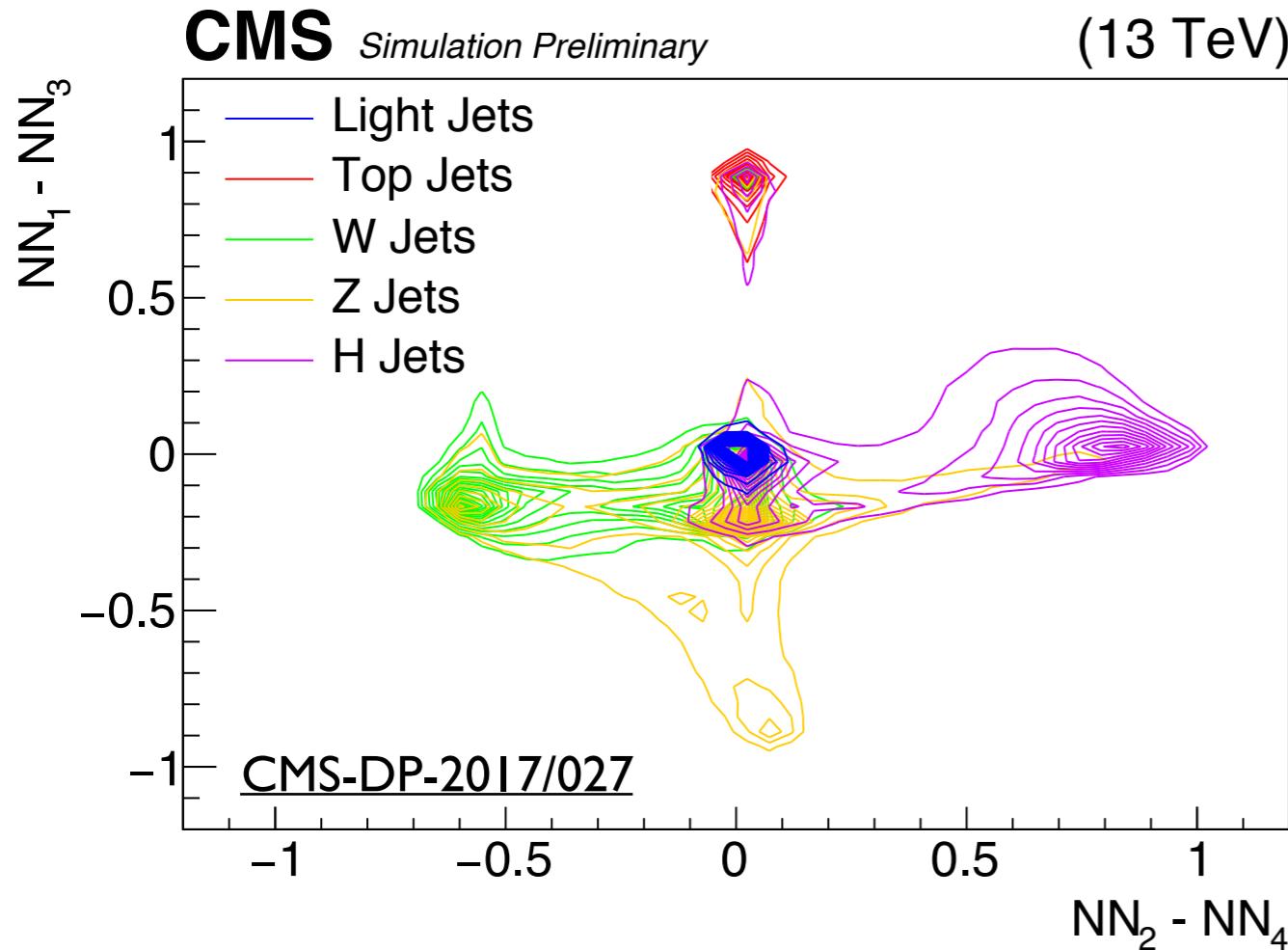
How do you search for all of these efficiently?

Multi-Node Tagging





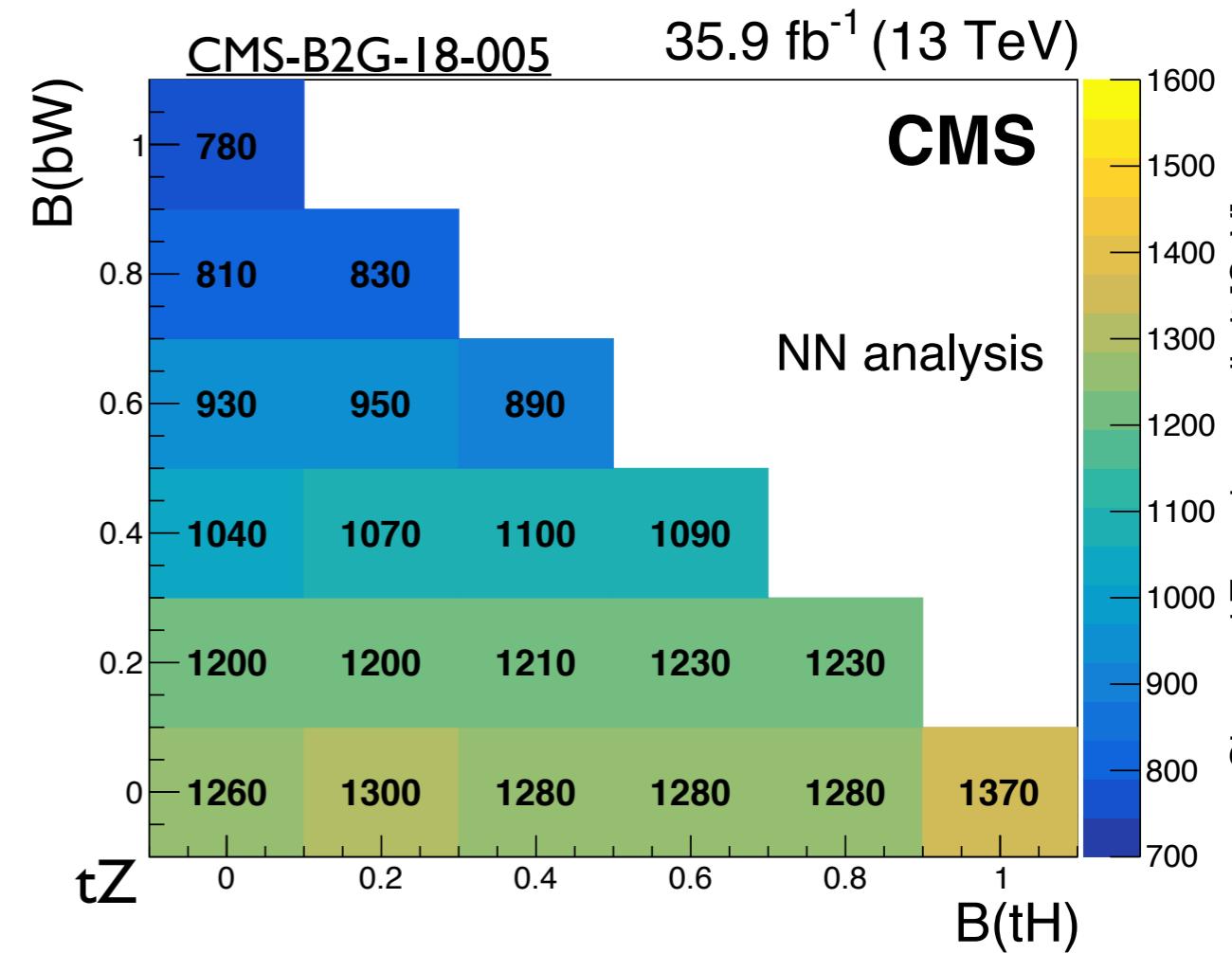
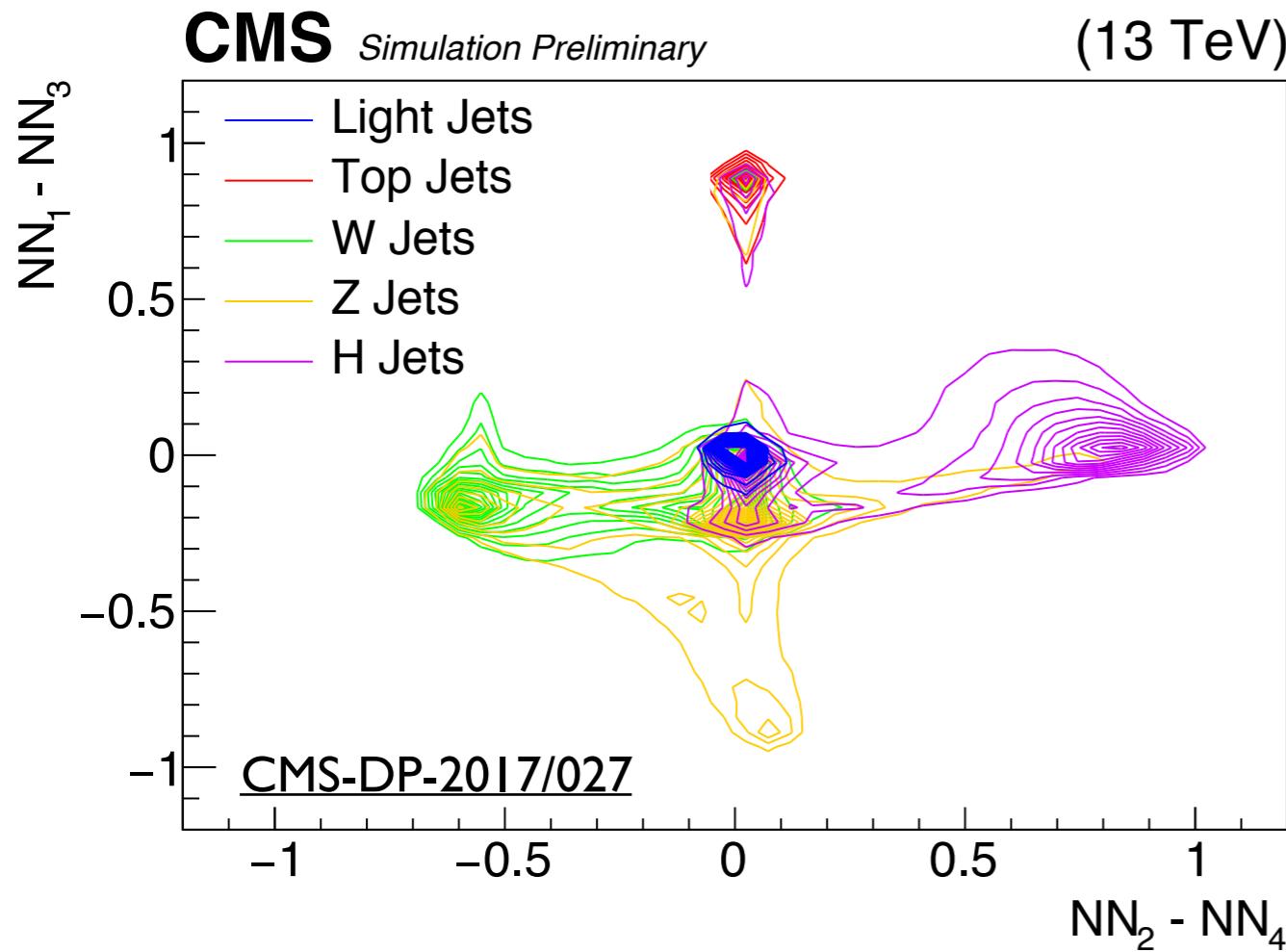
Multi-Node Tagging



Train a neural network to
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Multi-Node Tagging



Train a neural network to distinguish all signal classes from background

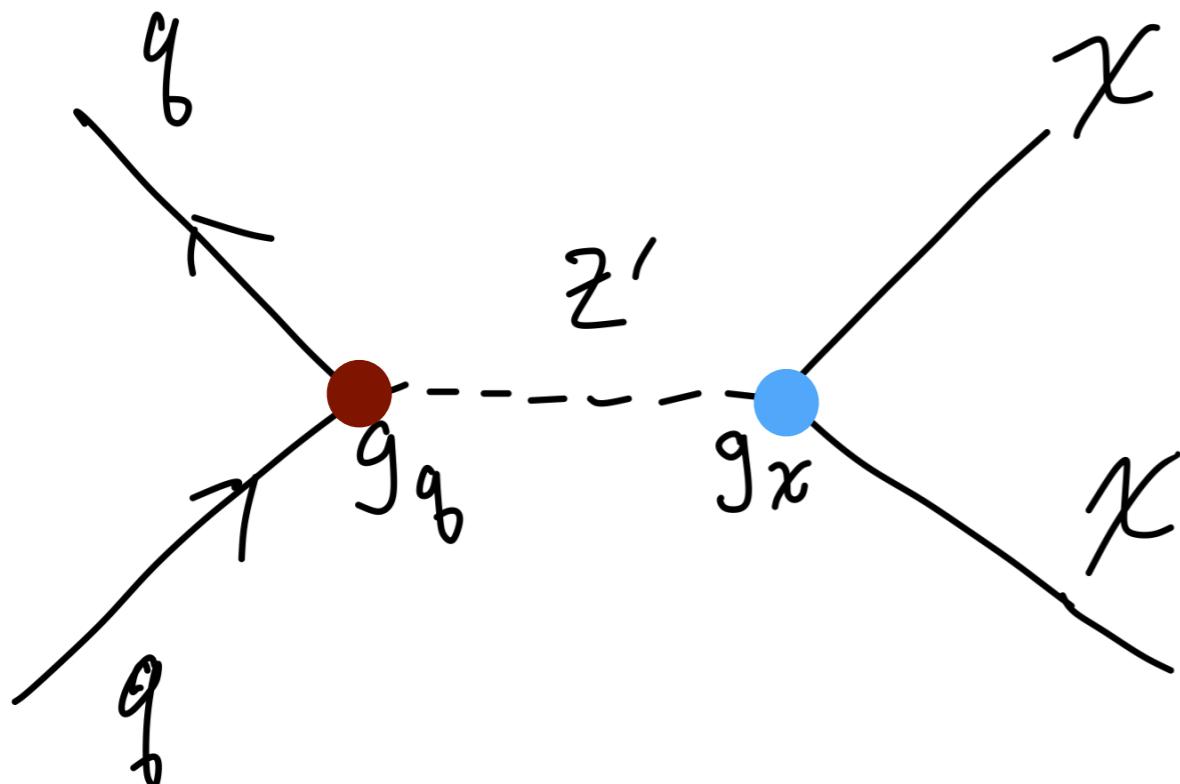
Use the NN to search for new physics in many final states at once

The Challenges of Dijets





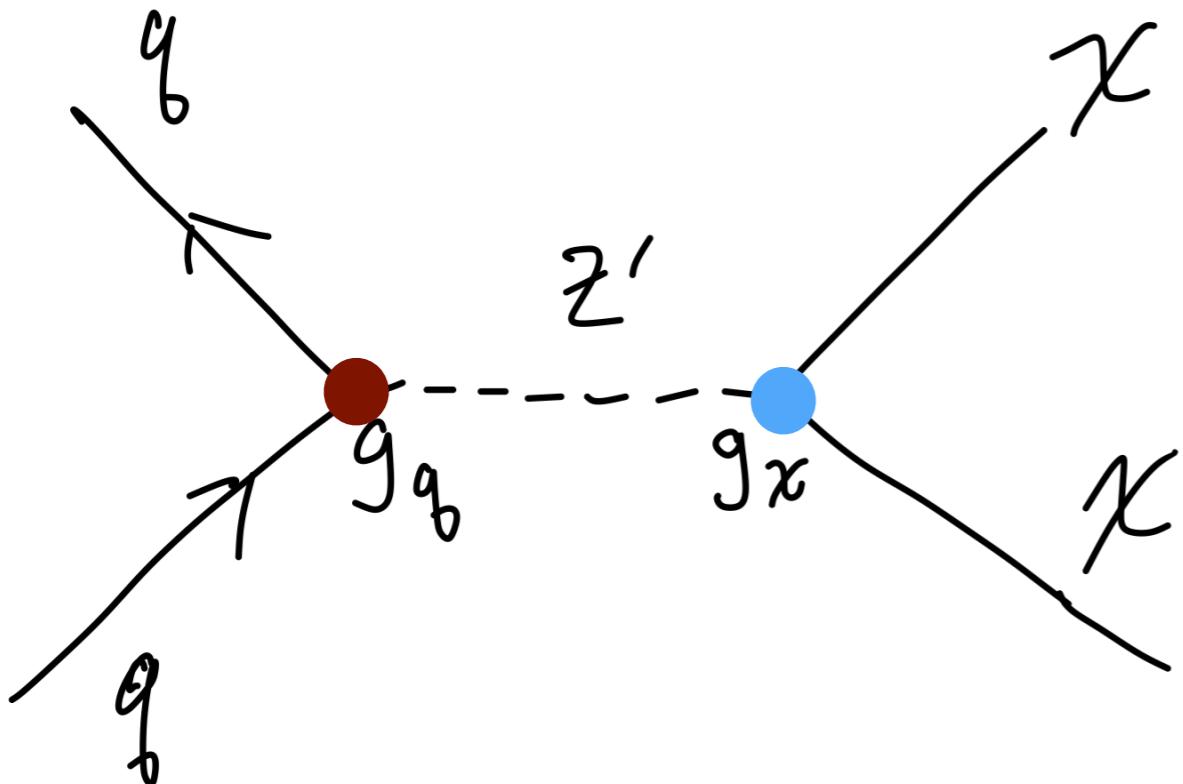
The Challenges of Dijets



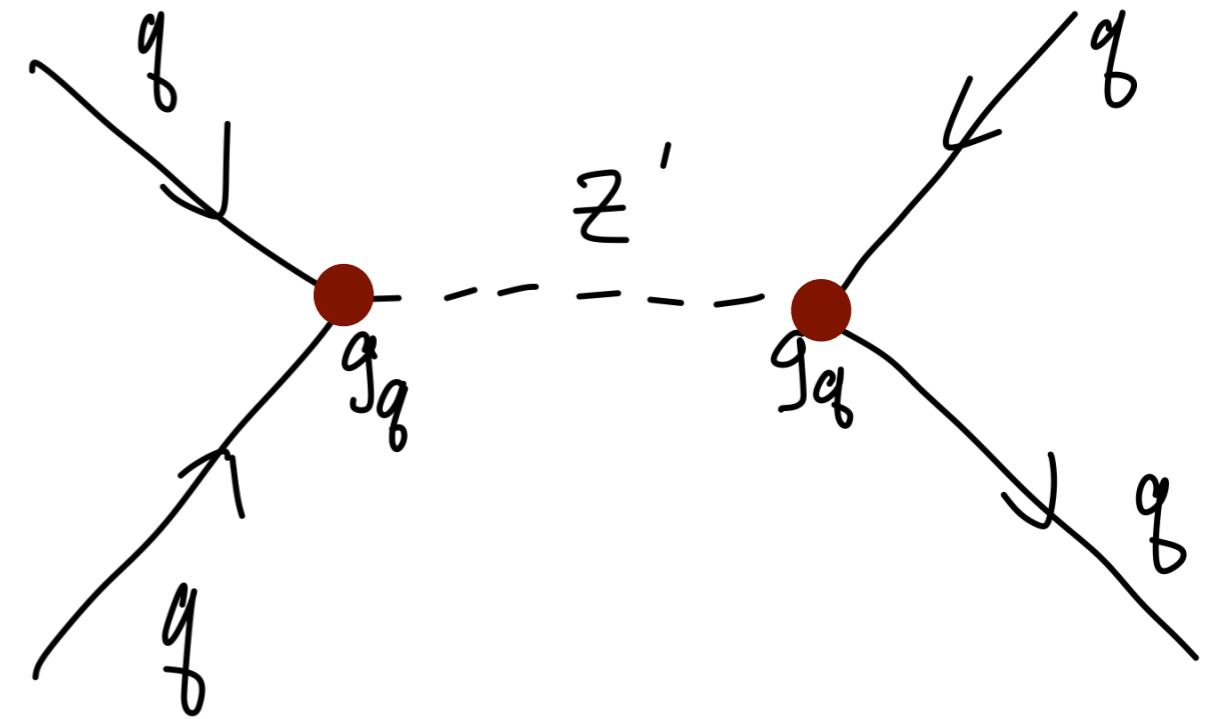
If we can produce DM
at the LHC like this...



The Challenges of Dijets



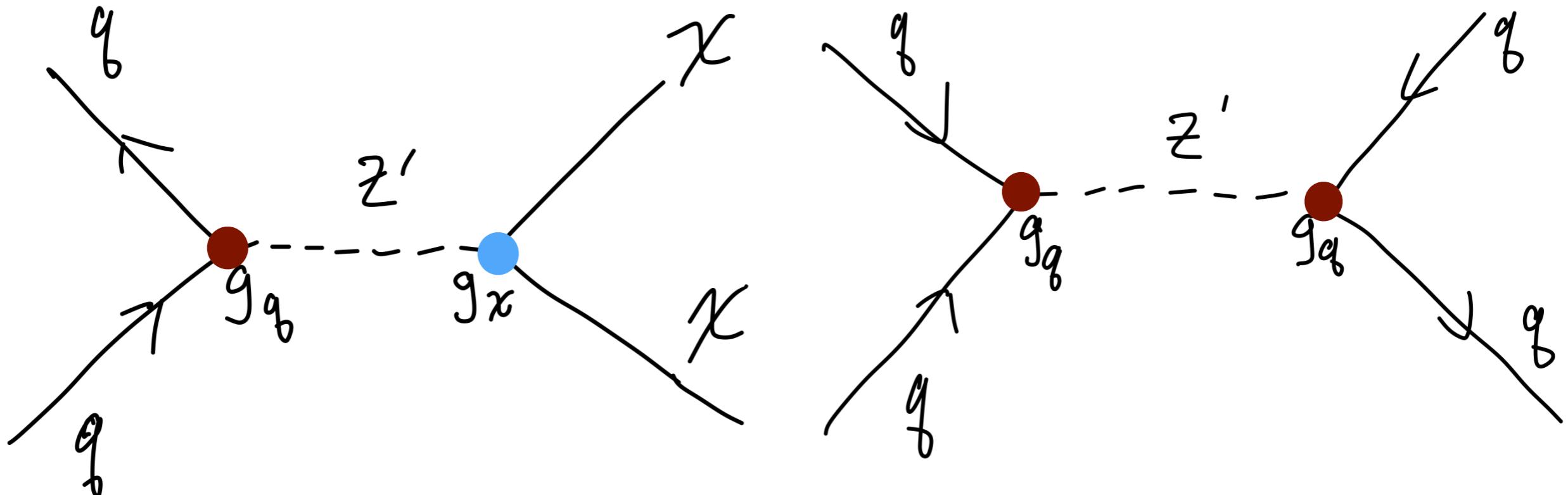
If we can produce DM
at the LHC like this...



Then the mediator should decay
back to jets!



The Challenges of Dijets



If we can produce DM
at the LHC like this...

Then the mediator should decay
back to jets!

But if the Z' is low enough mass,
we won't trigger the event from the jets!

Data Scouting





Data Scouting

Trigger limitations
come from
total bandwidth



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$$\text{Event Size} \times \text{Event Rate} = \text{Total Bandwidth}$$



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Normally, we record
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But we can reco
events in the trigger,
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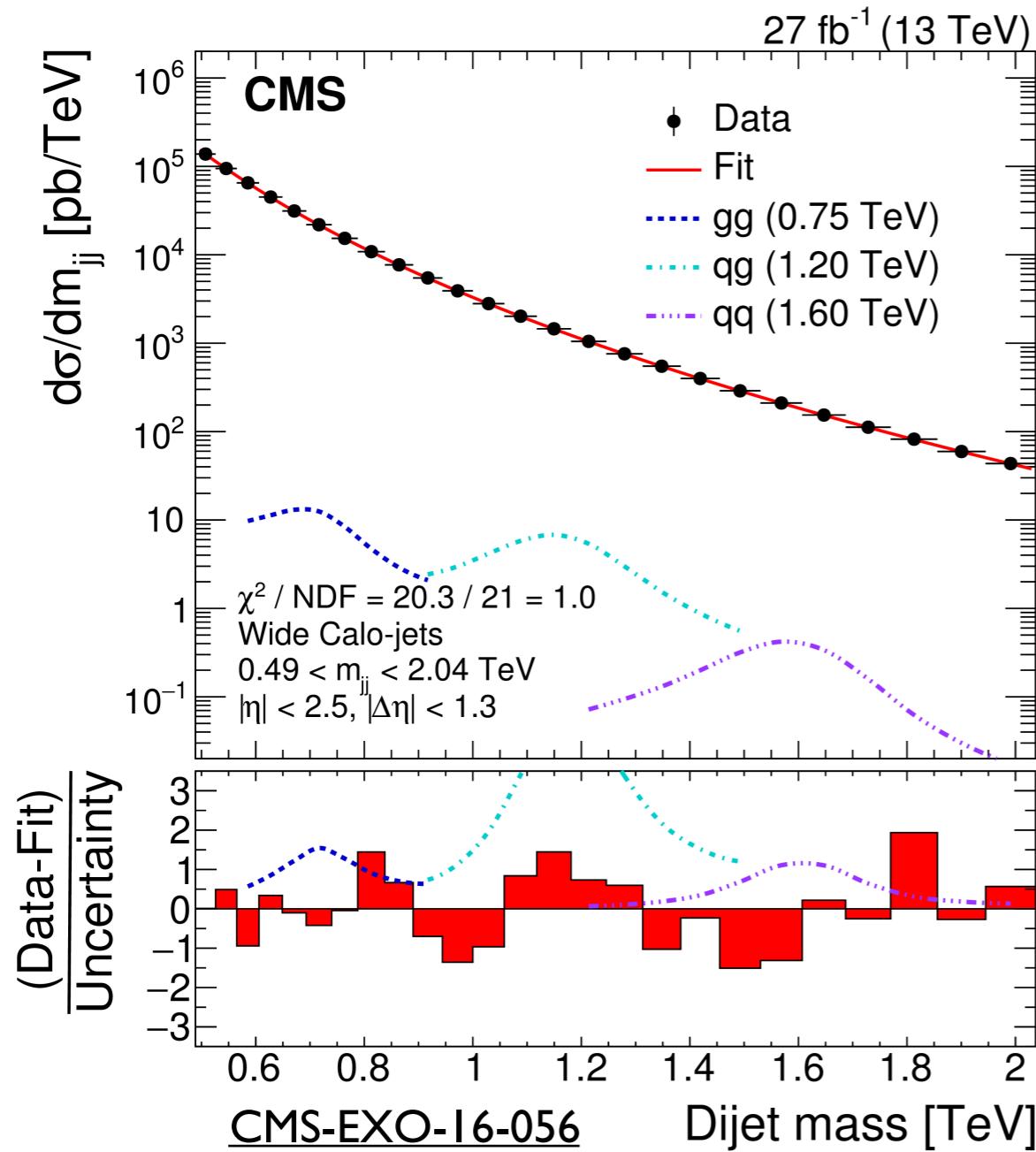
$$\text{Size} \times \text{Event Rate} = \text{Total Bandwidth}$$

If you save smaller events, you can save more of them!

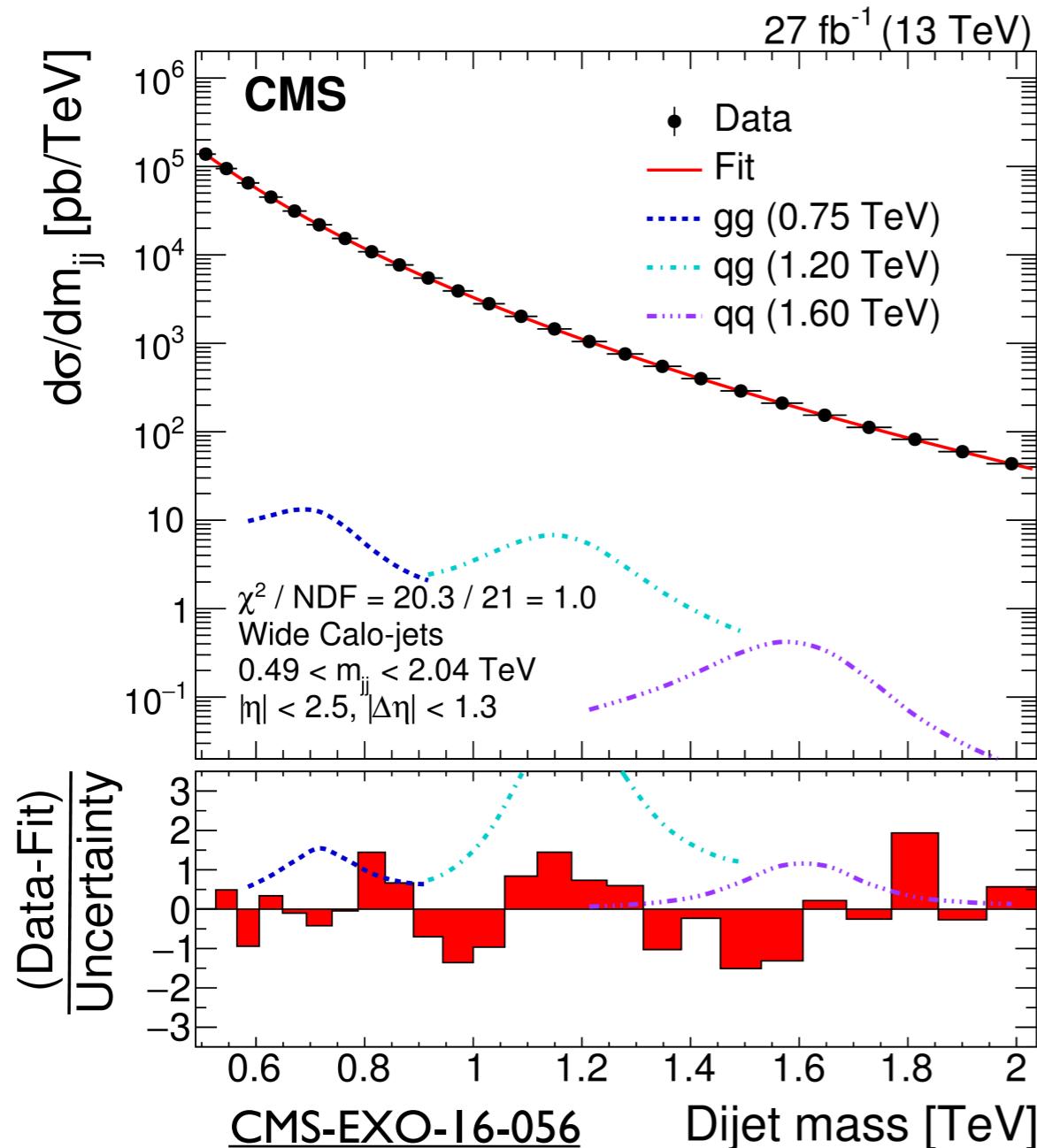
What You Can Accomplish



What You Can Accomplish

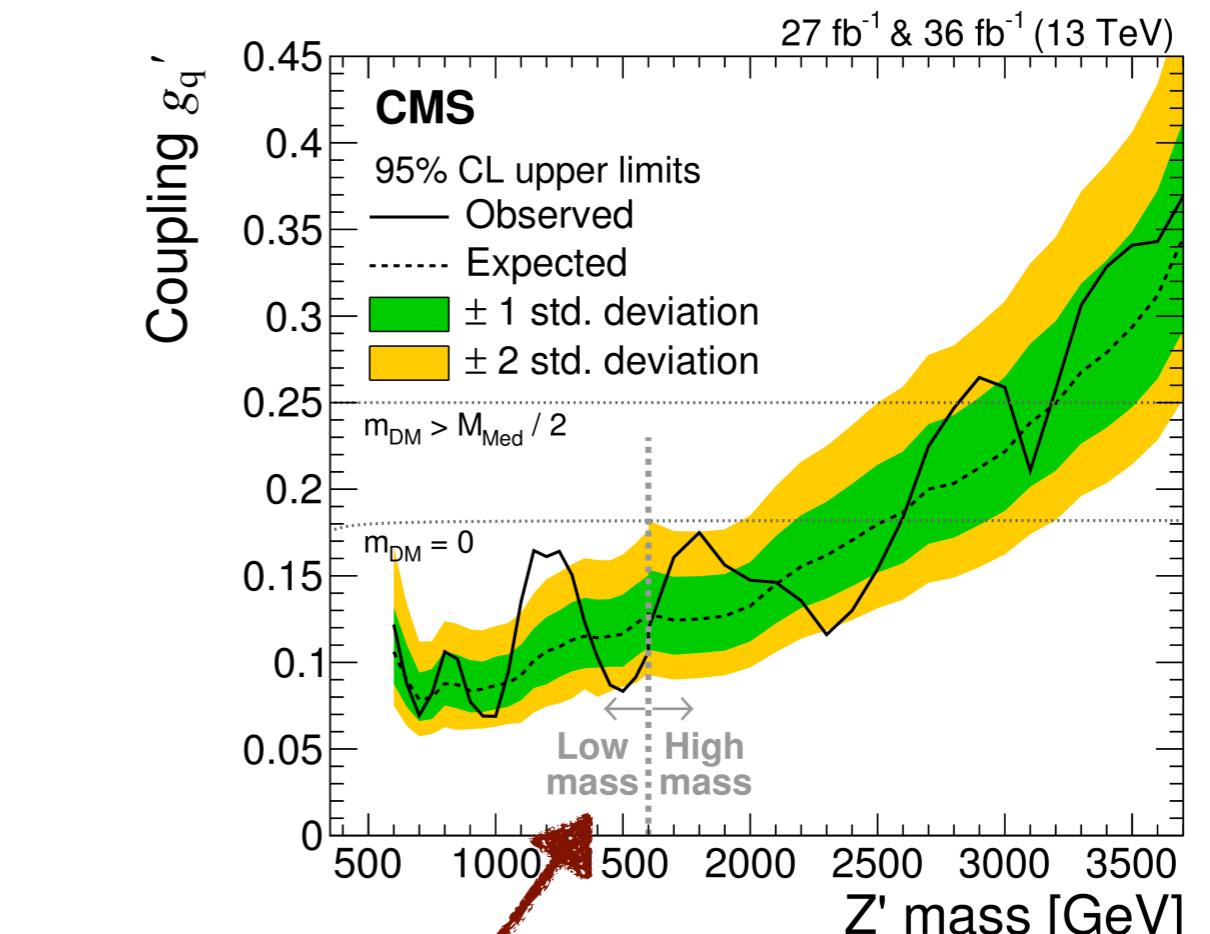
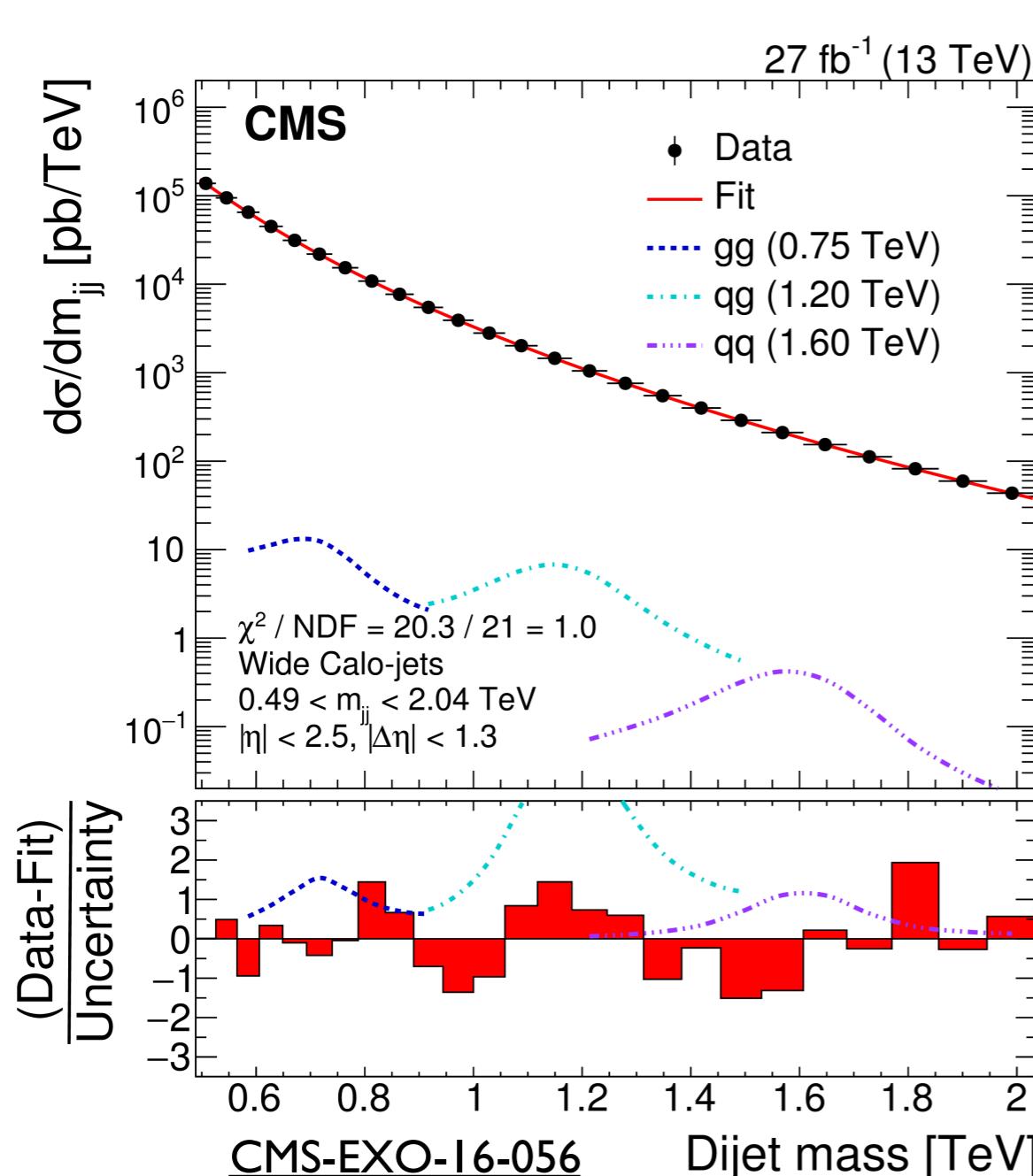


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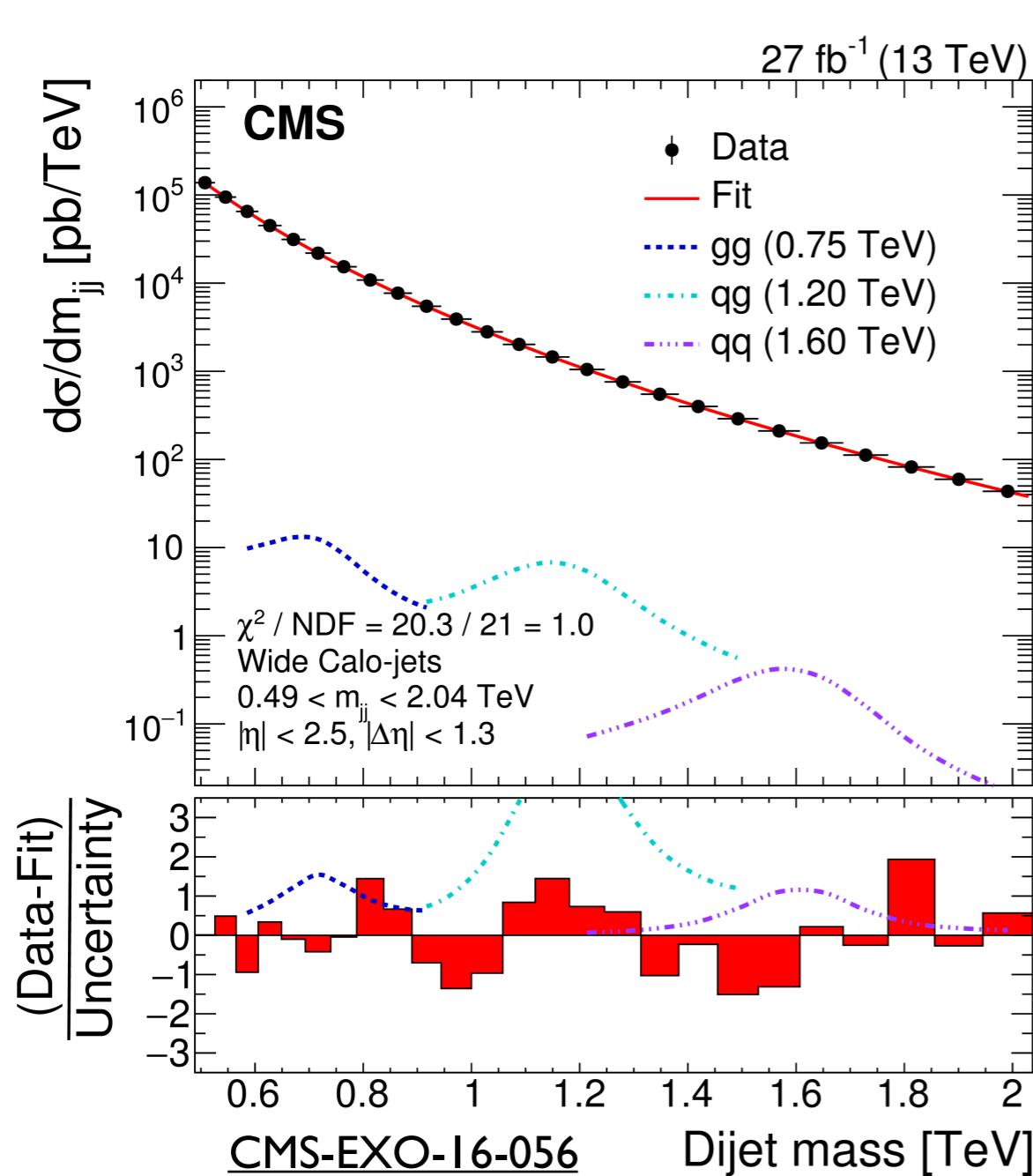
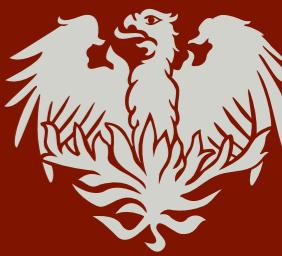
Can collect huge datasets
at low mass using this technique!

What You Can Accomplish

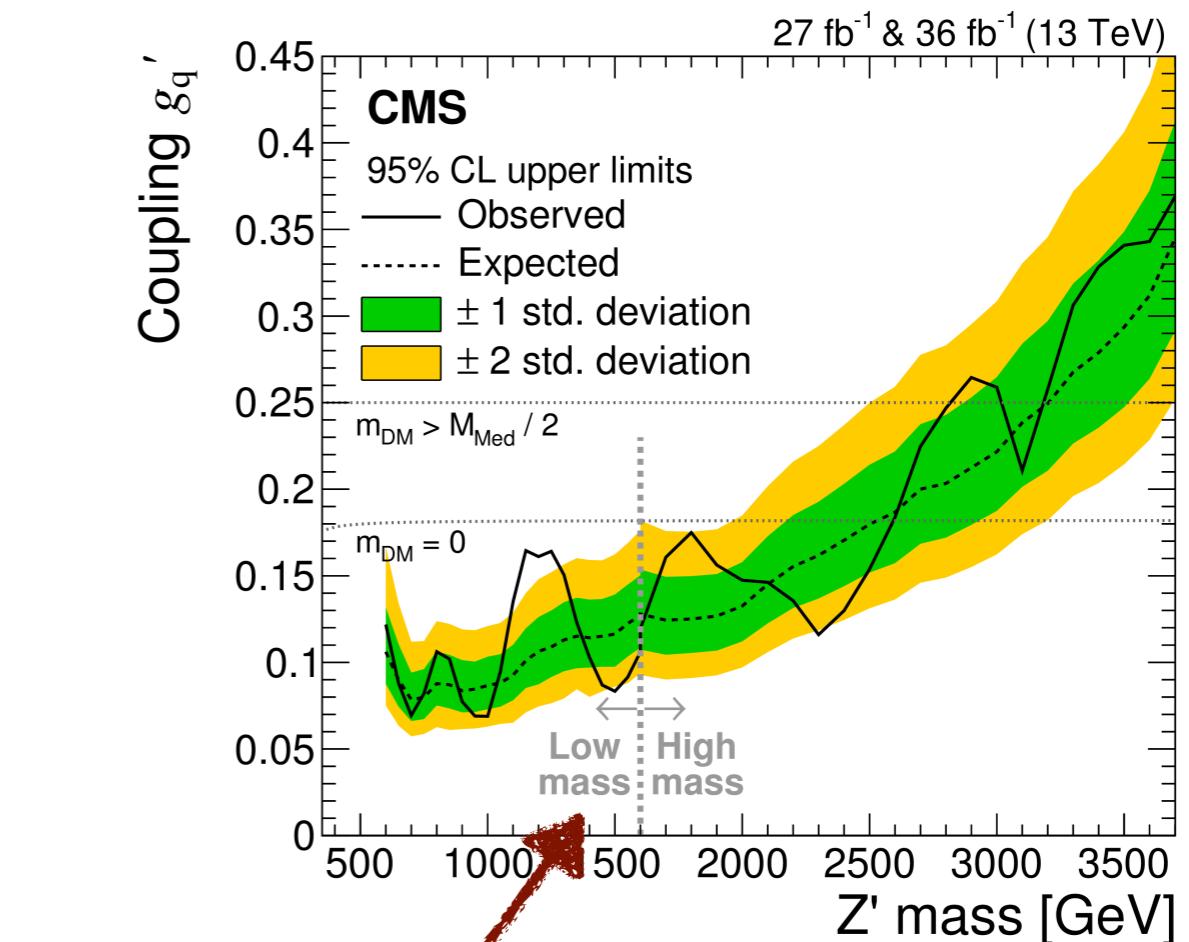


Can collect huge datasets
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What You Can Accomplish



Can collect huge datasets
at low mass using this technique!



Can probe unique physics
phase space by recording
only a portion of the events!

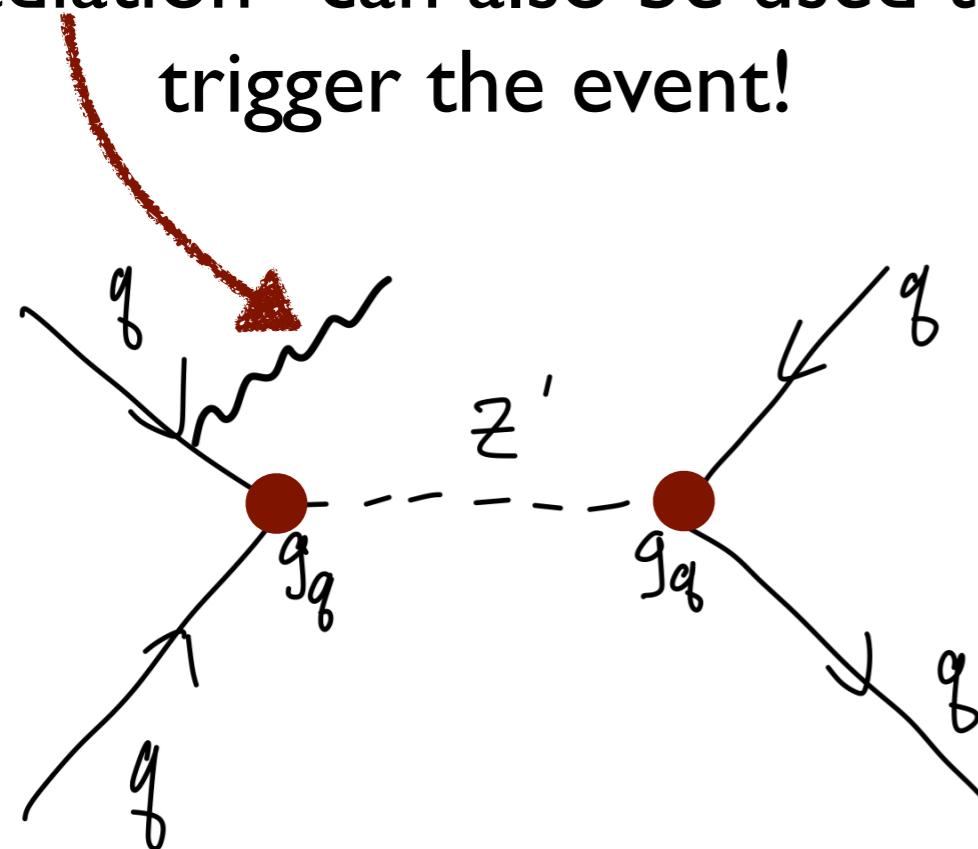
Extra Radiation





Extra Radiation

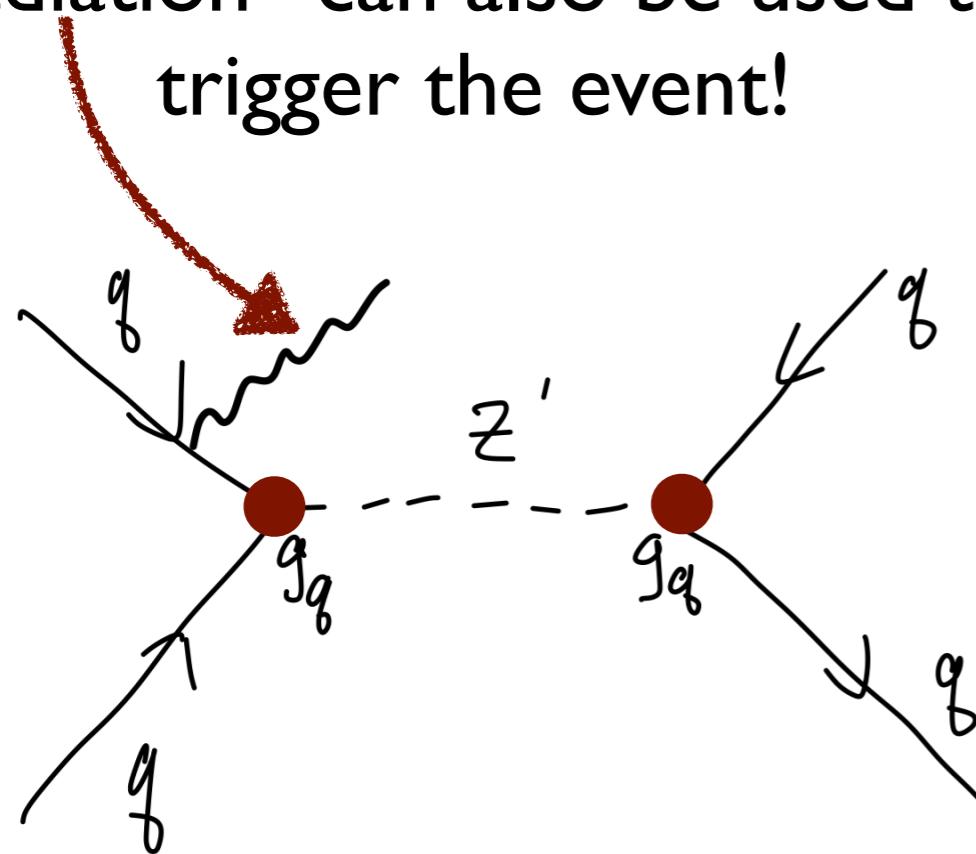
Photon or jet “initial state radiation” can also be used to trigger the event!



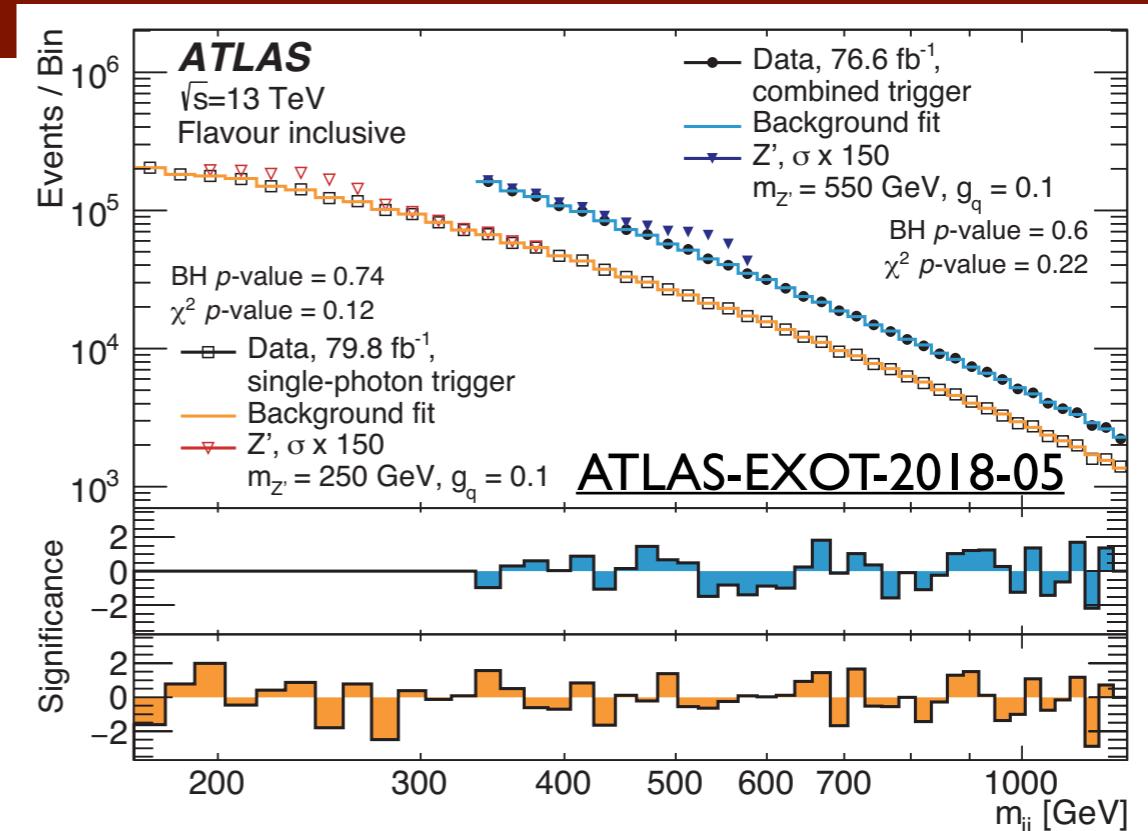


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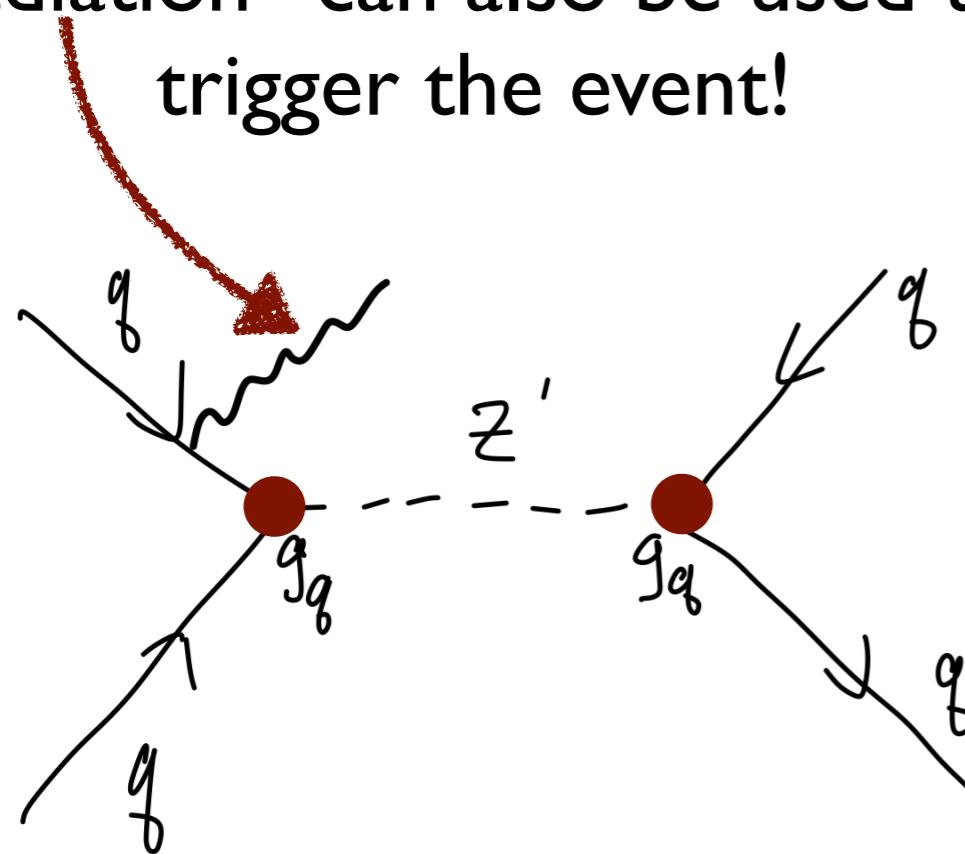
Even lower regions of mediator mass can be explored with these techniques





Extra Radiation

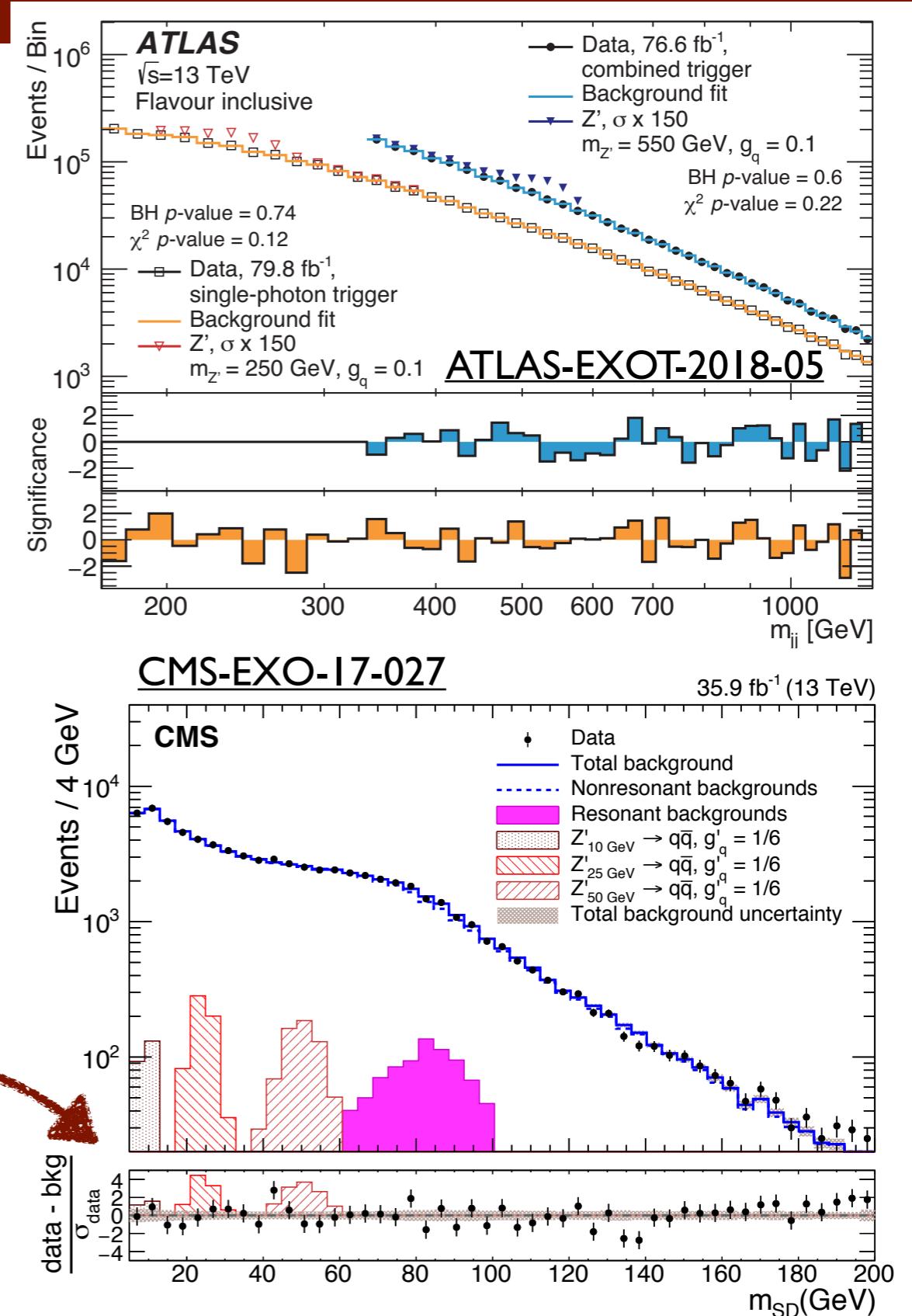
Photon or jet “initial state radiation” can also be used to trigger the event!



Even lower regions of mediator mass can be explored with these techniques



Can even search as low as 10 GeV!

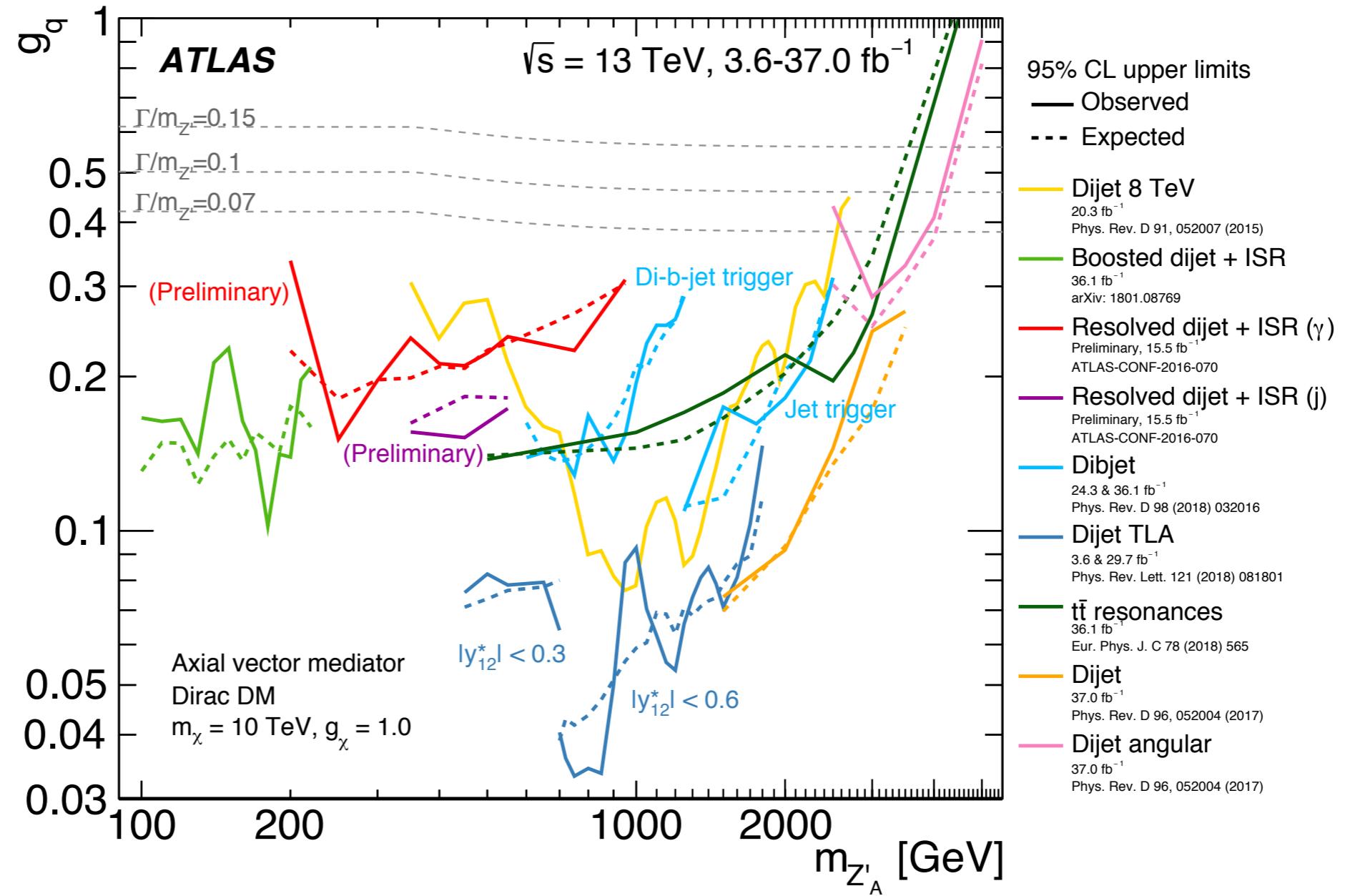


The Big Picture for Z'



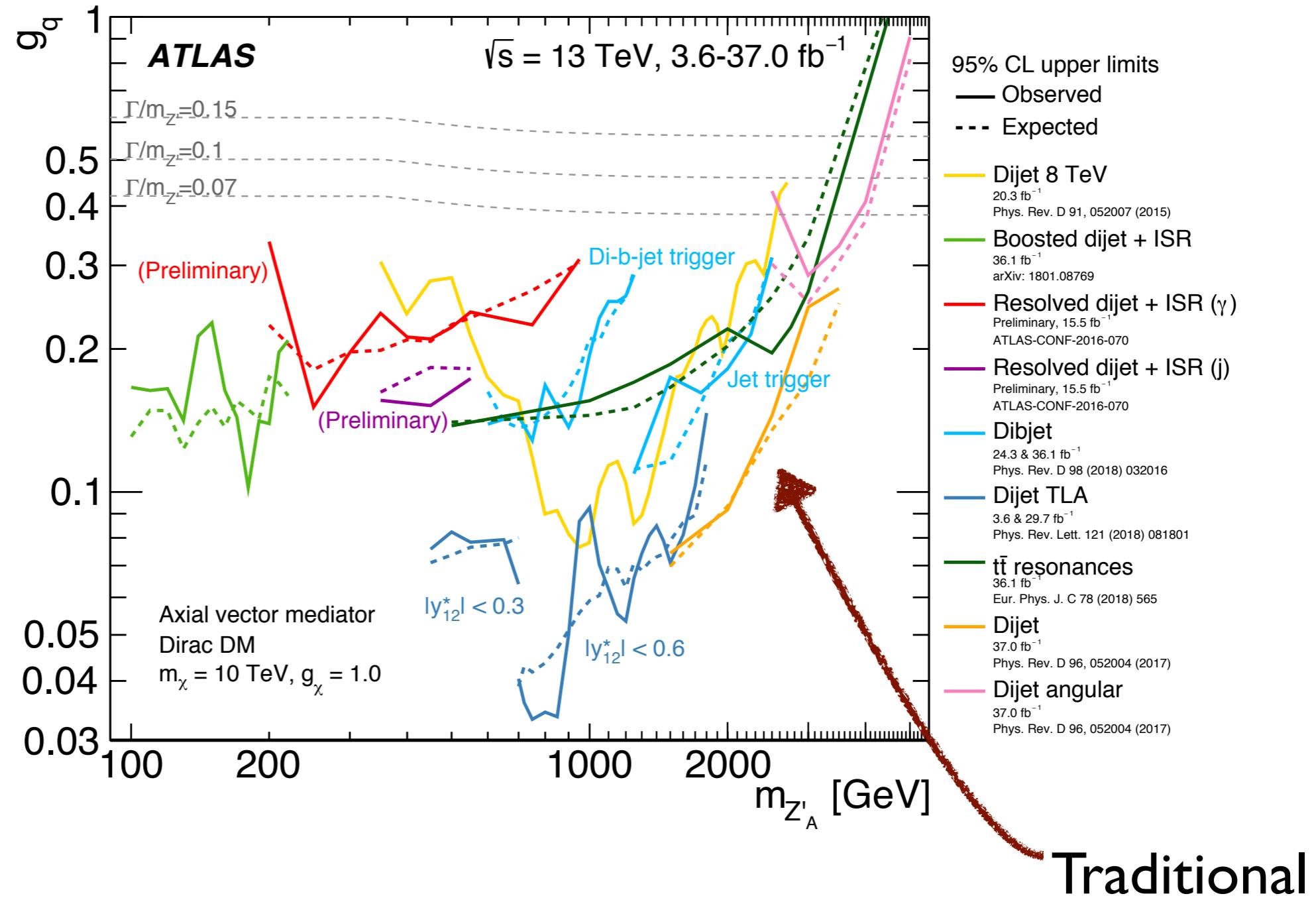


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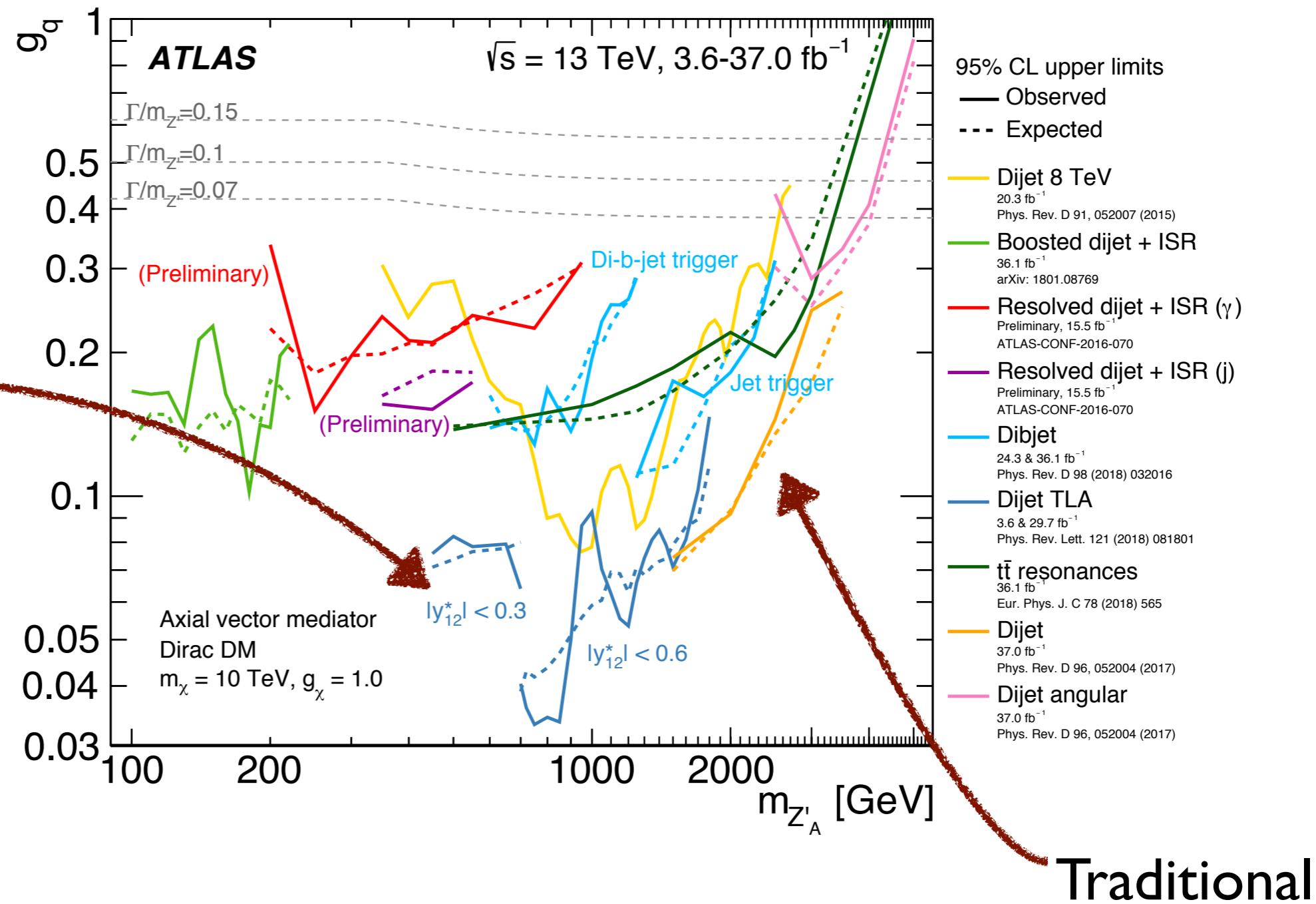
The Big Picture for Z'





The Big Picture for Z'

Trigger
Only



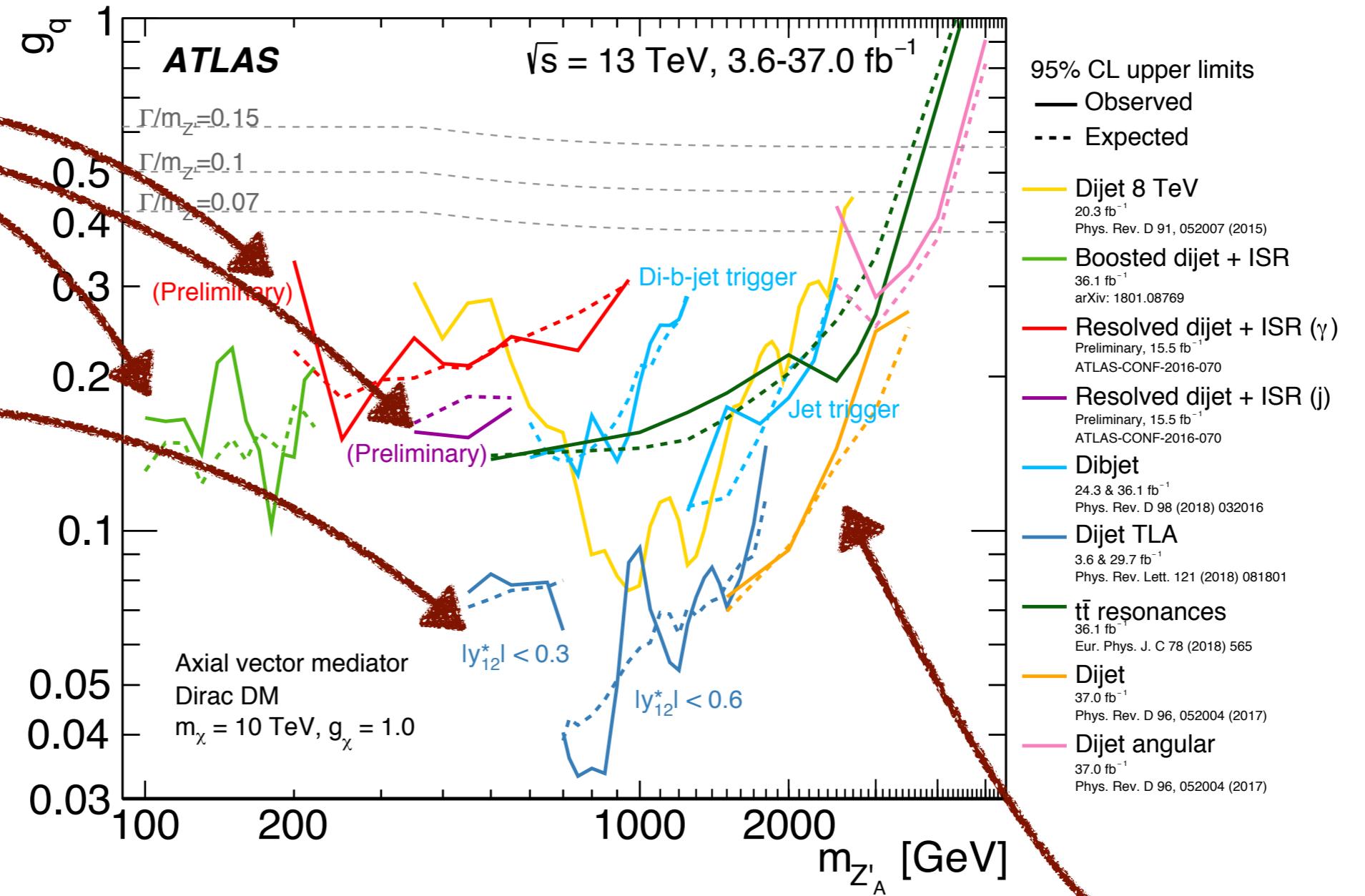


The Big Picture for Z'

ISR

Trigger
Only

Traditional

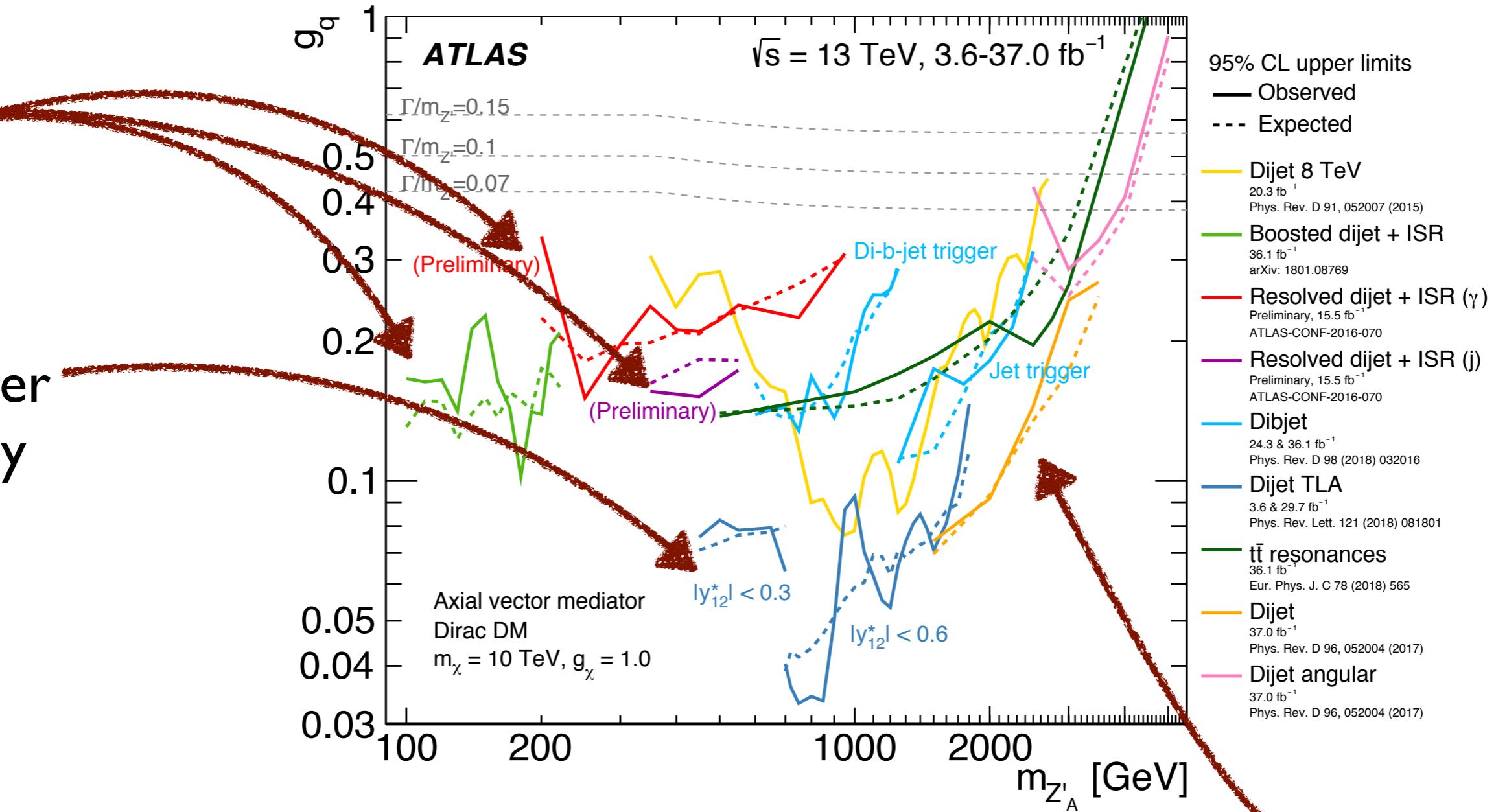




The Big Picture for Z'

ISR

Trigger
Only



In the hunt for Z', new techniques
are taking center stage!

Traditional

Rarer than Rare

*Or: how huge datasets and clever strategies
are enabling searches for vanishingly
small signals*

Why Search for Rare Signals?



Why Search for Rare Signals?

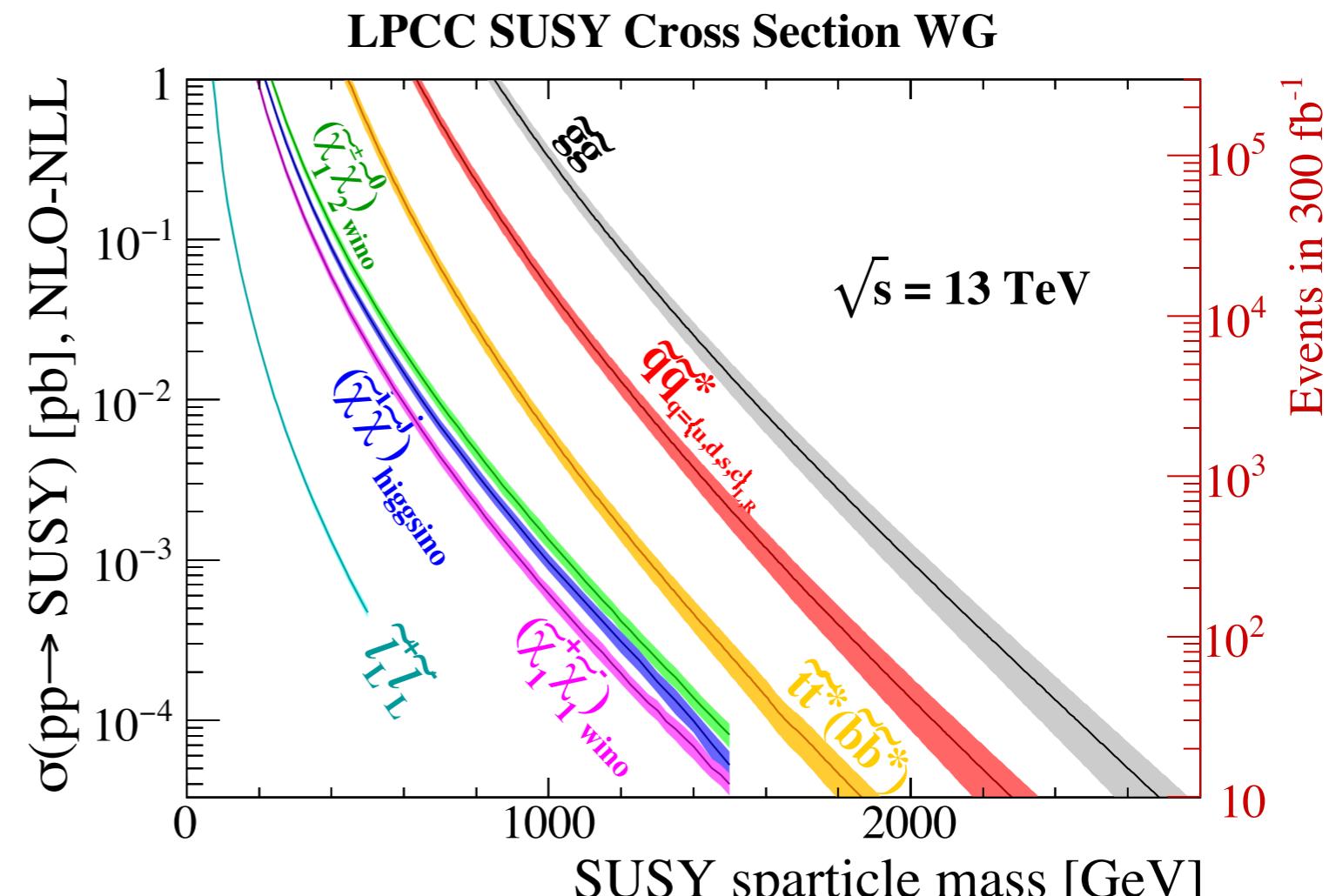


Because many theories predict low cross-section signals!



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Because many theories predict low cross-section signals!



<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections>

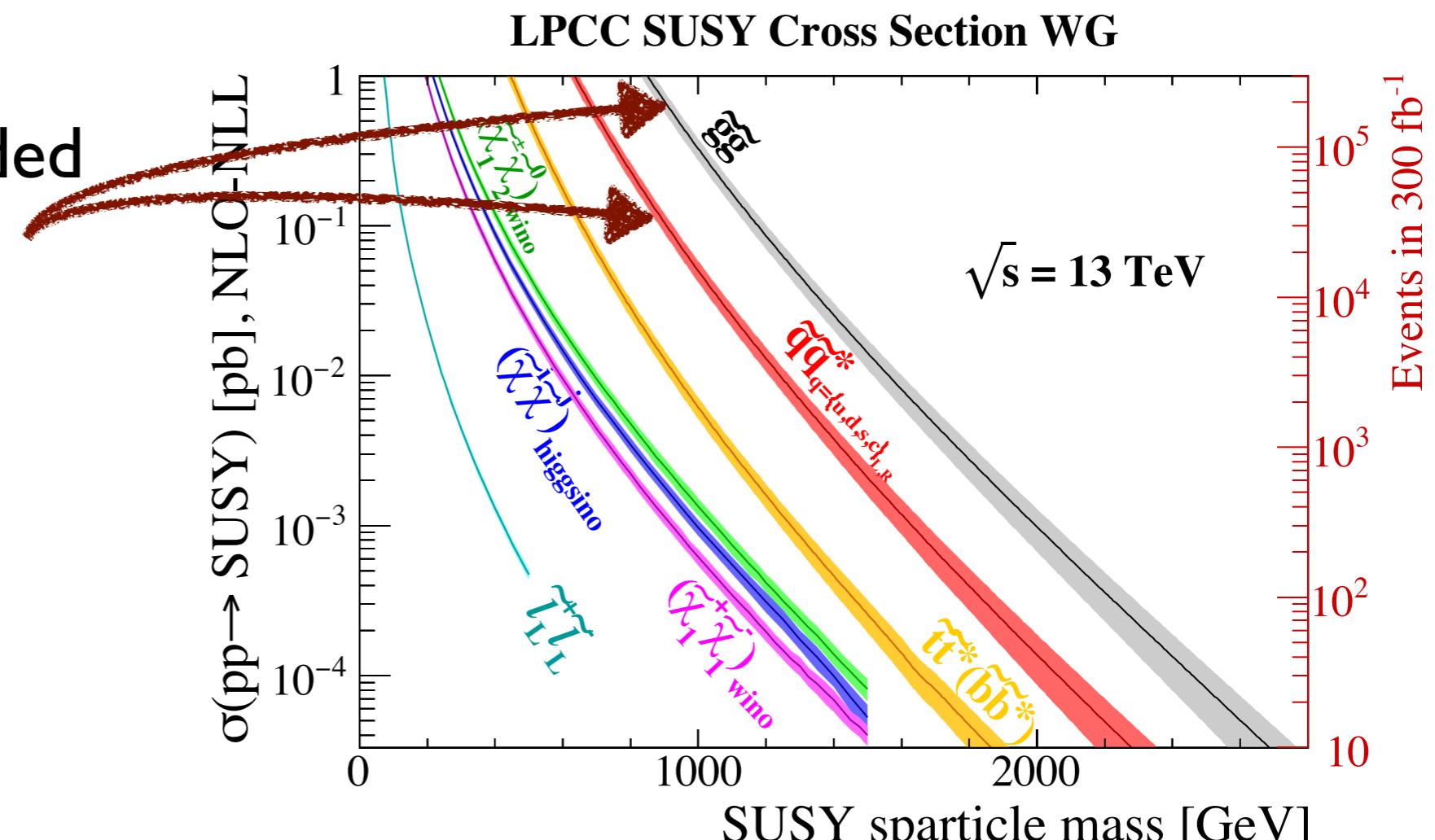
arXiv:1407.5066



Why Search for Rare Signals?

Because many theories predict low cross-section signals!

If you've already excluded
gluinos and squarks
and stops...



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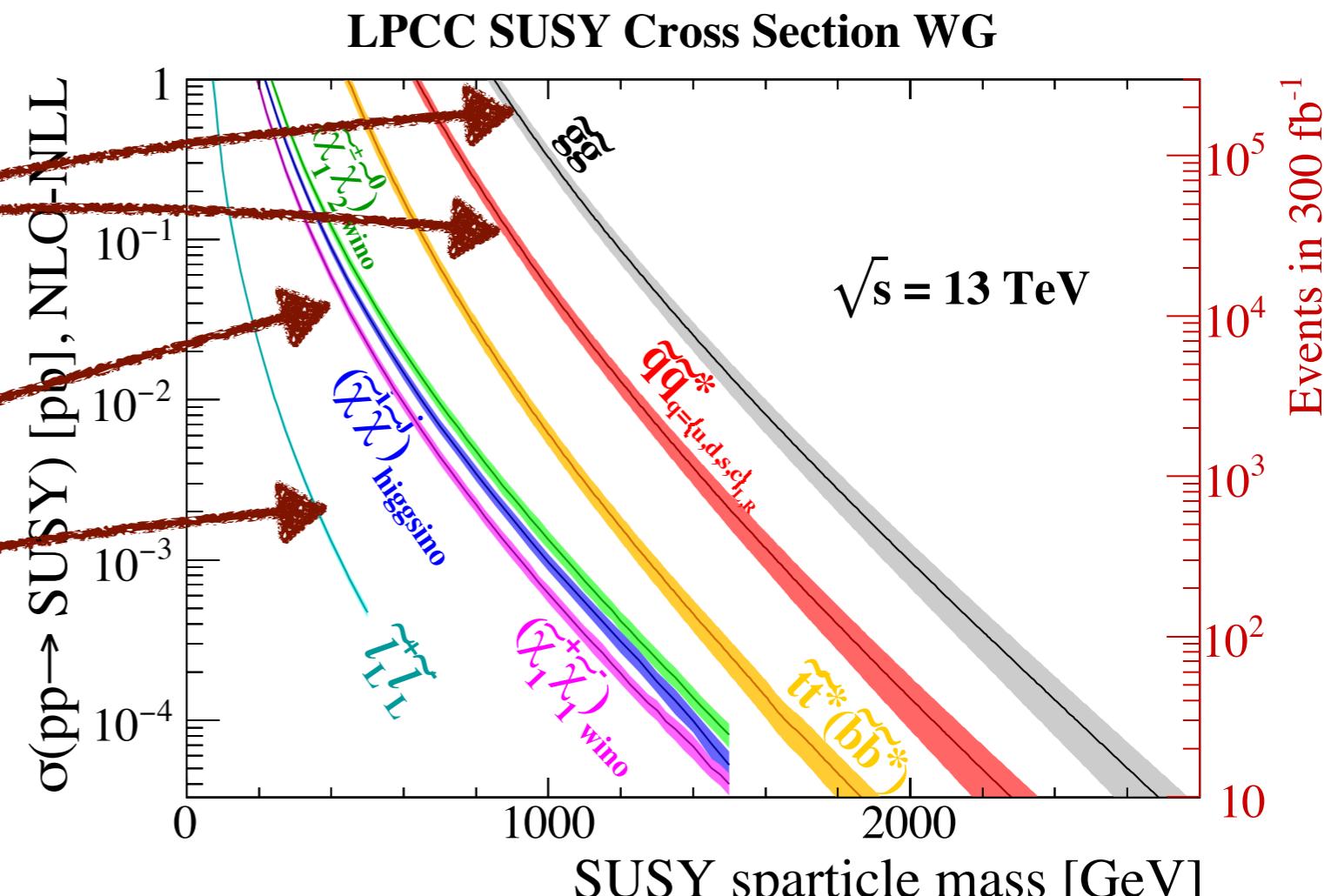


Why Search for Rare Signals?

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Higgsinos and sleptons
might still be in sight!



<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections>

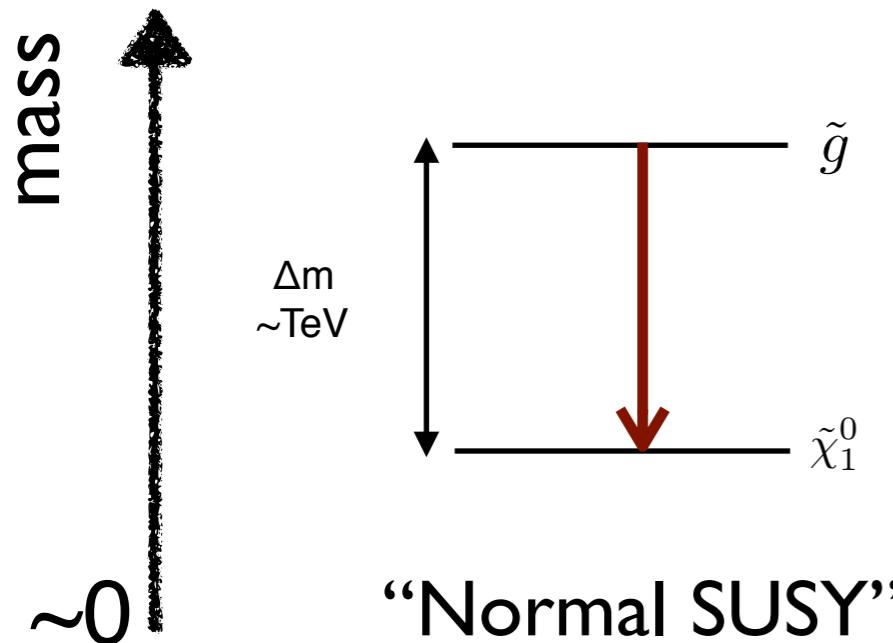
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Higgsinos: The Challenge





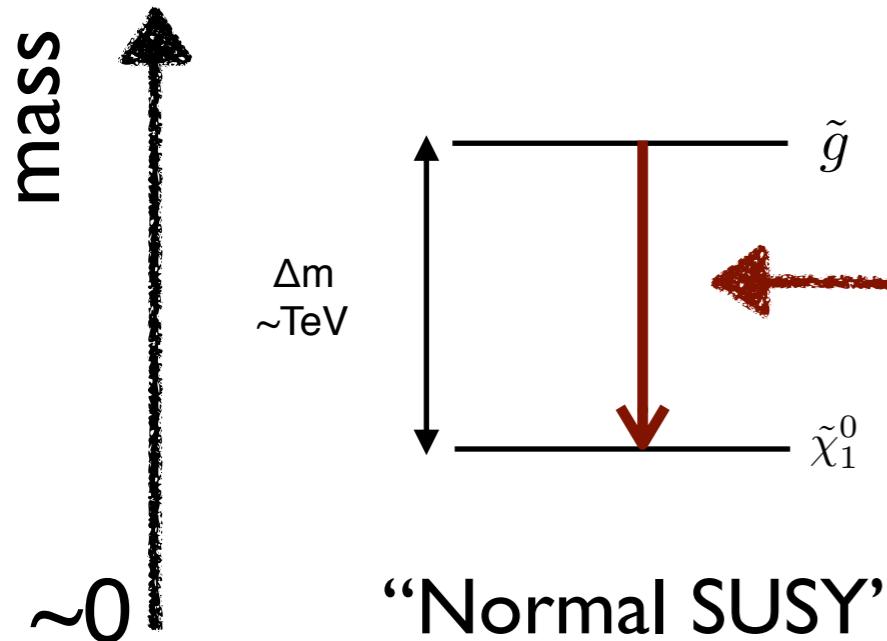
Higgsinos: The Challenge



**“Classic” SUSY searches
exploit large mass splittings**



Higgsinos: The Challenge

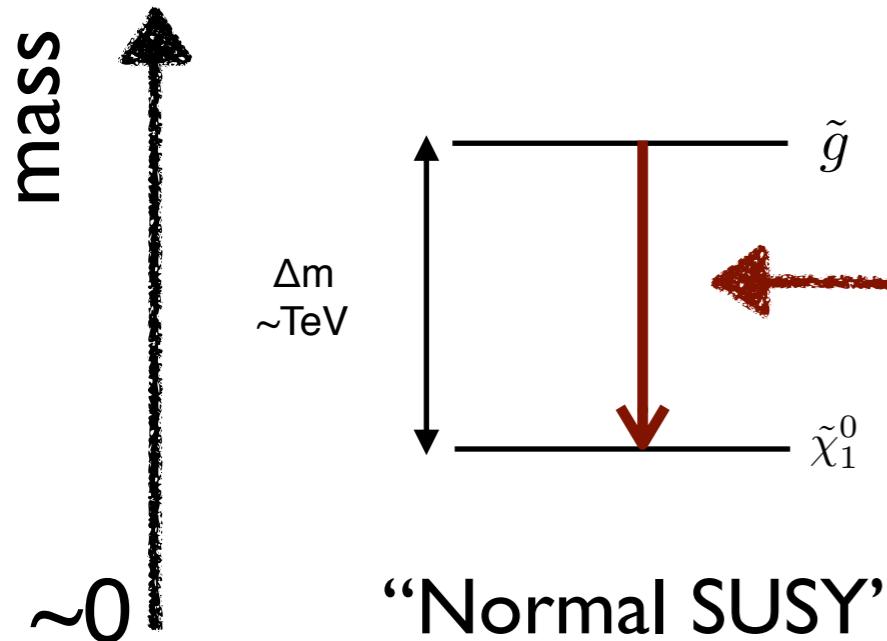


“Classic” SUSY searches
exploit large mass splittings

Cascades from initial sparticles
will be high energy,
lots of missing energy

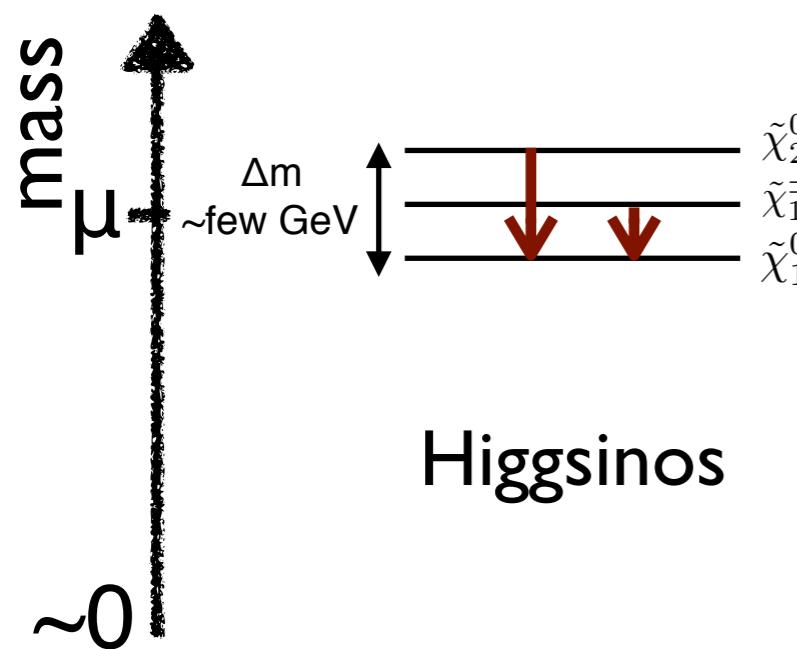


Higgsinos: The Challenge



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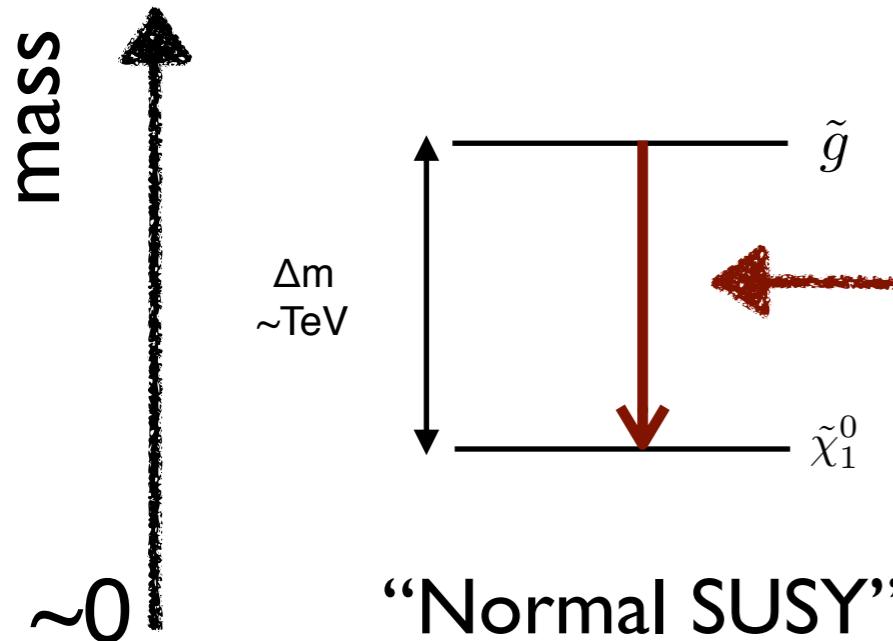
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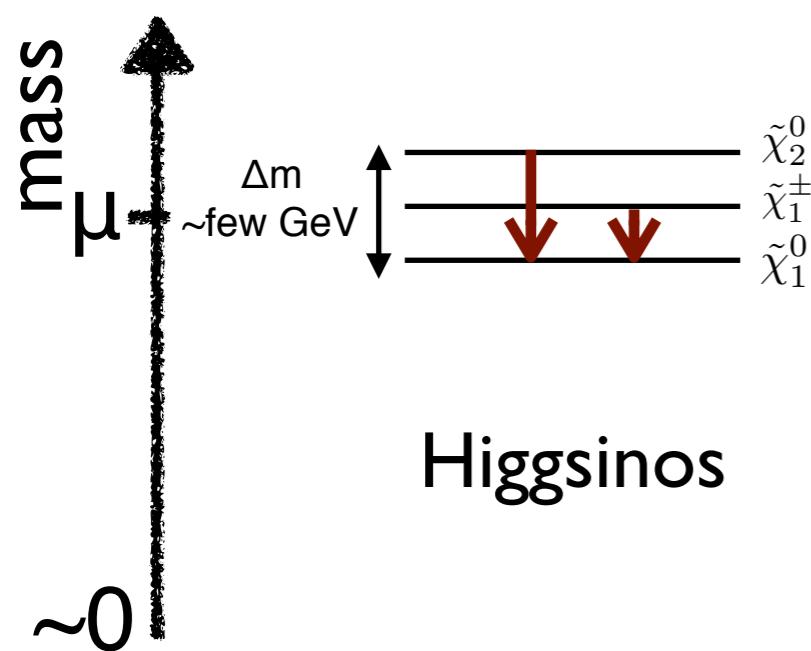
Higgsinos are another story...



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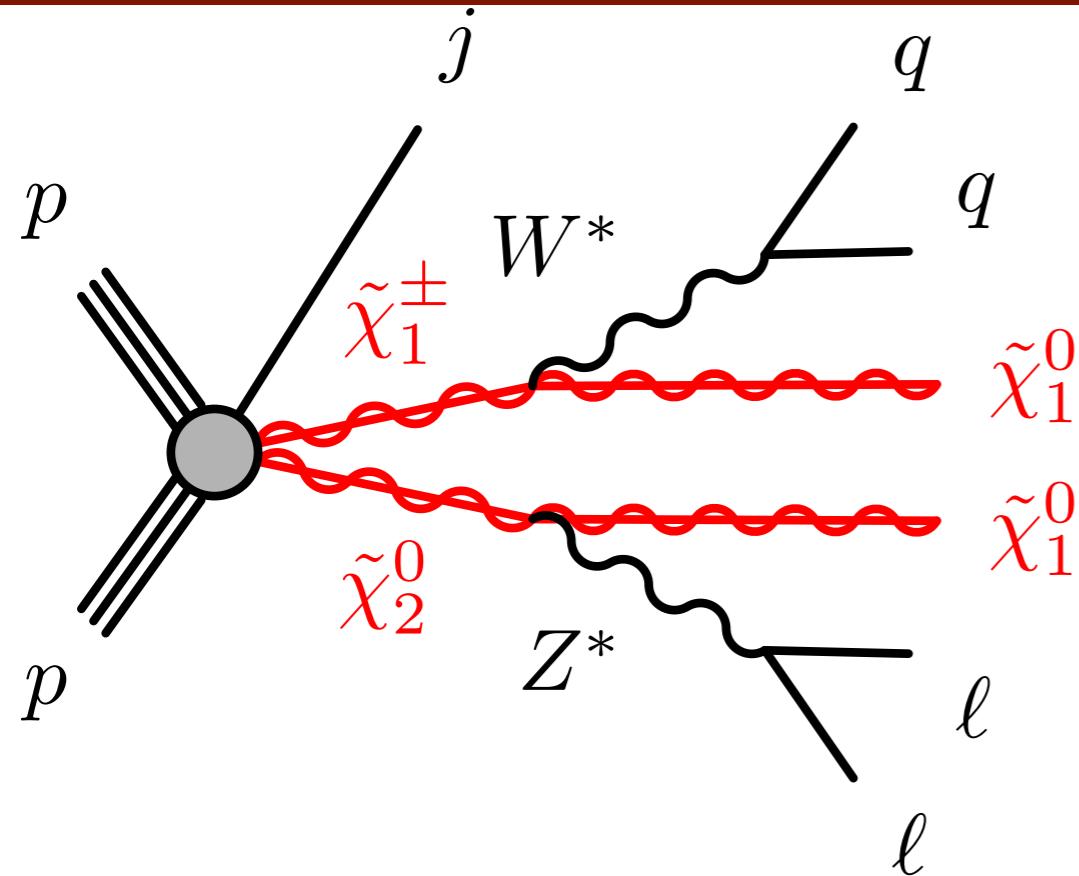
Higgsinos are another story...

No large mass splitting:
very soft visible particles,
no obvious missing energy

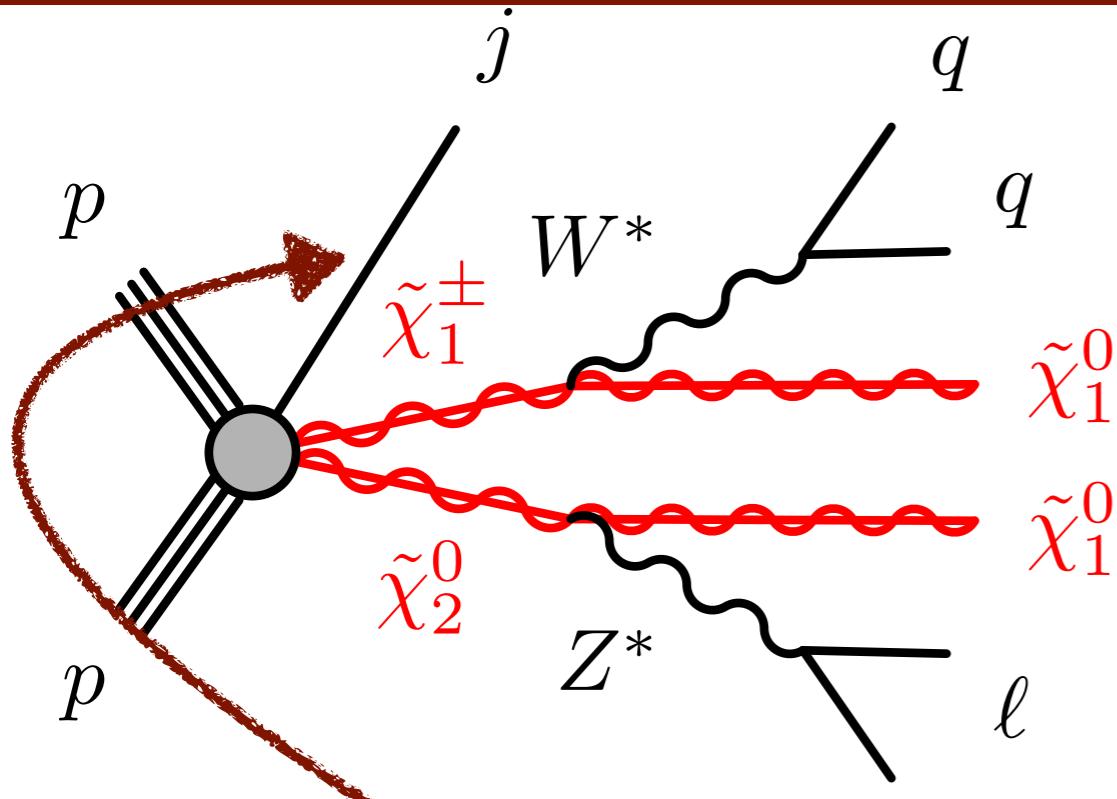
Searching for Rare Higgsinos



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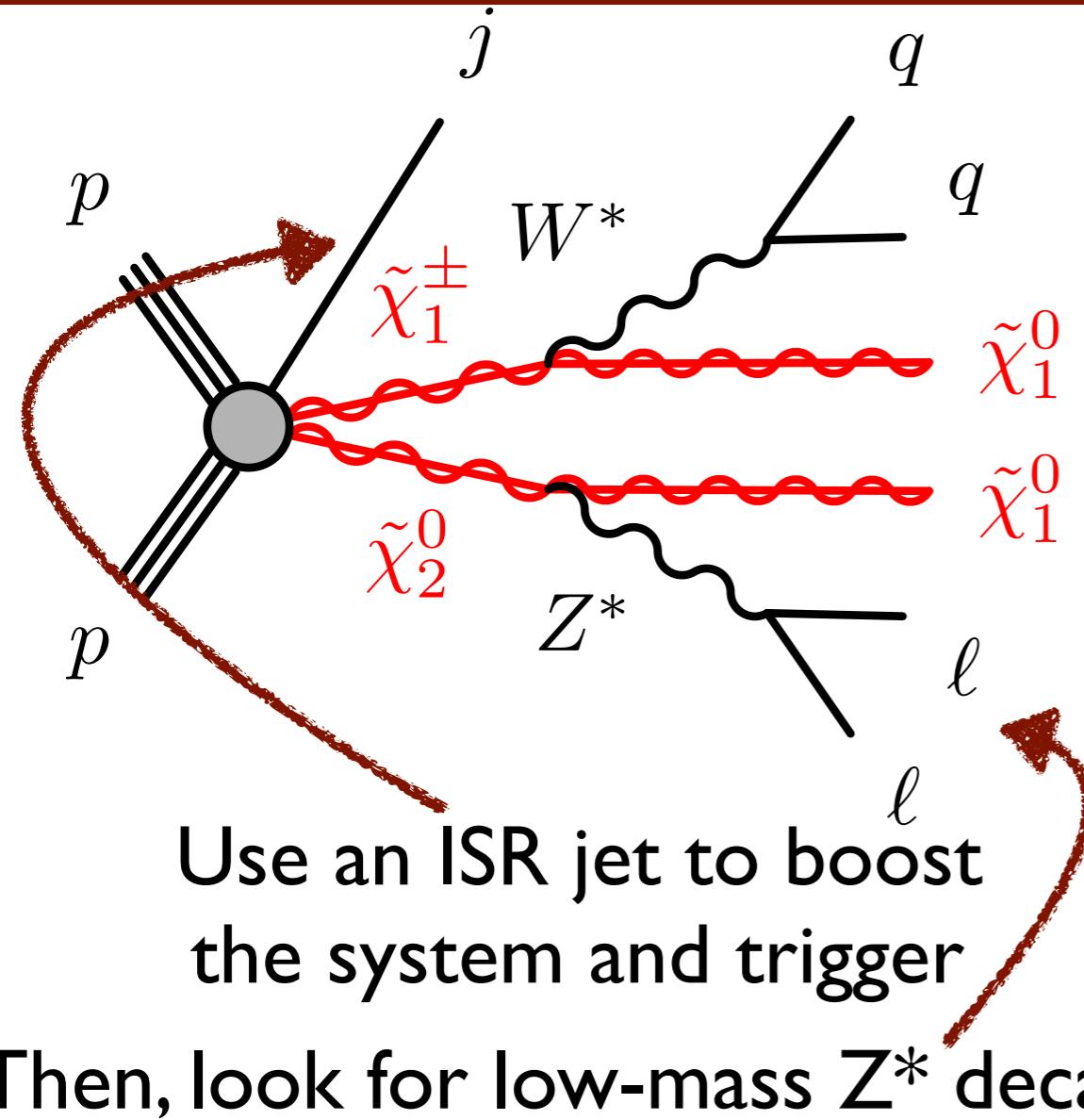
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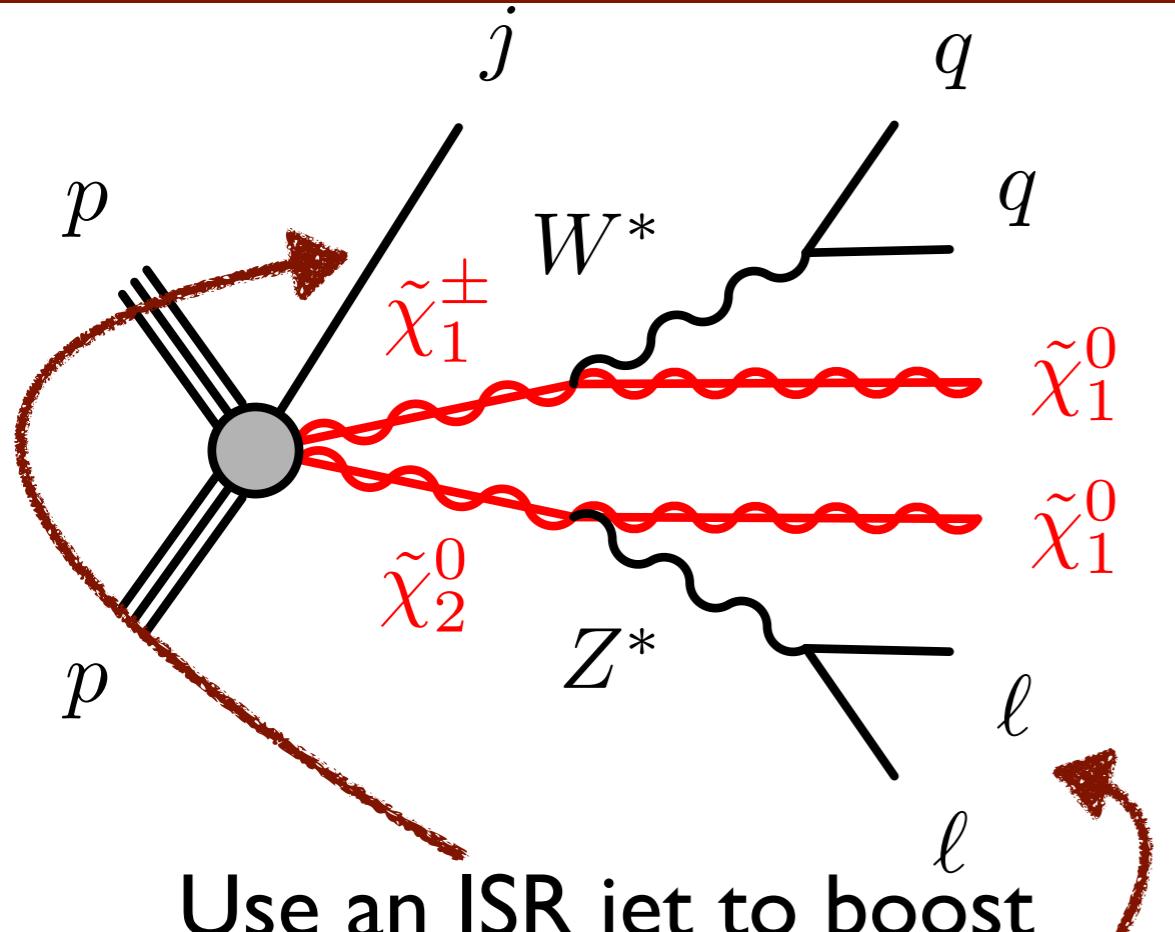
Use an ISR jet to boost
the system and trigger



Searching for Rare Higgsinos



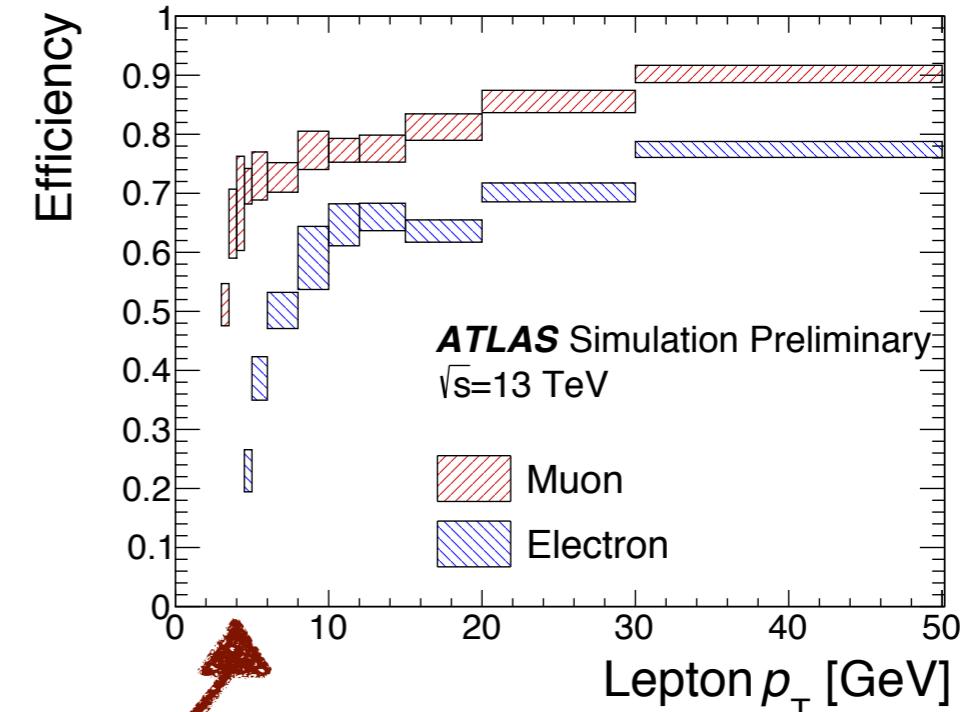
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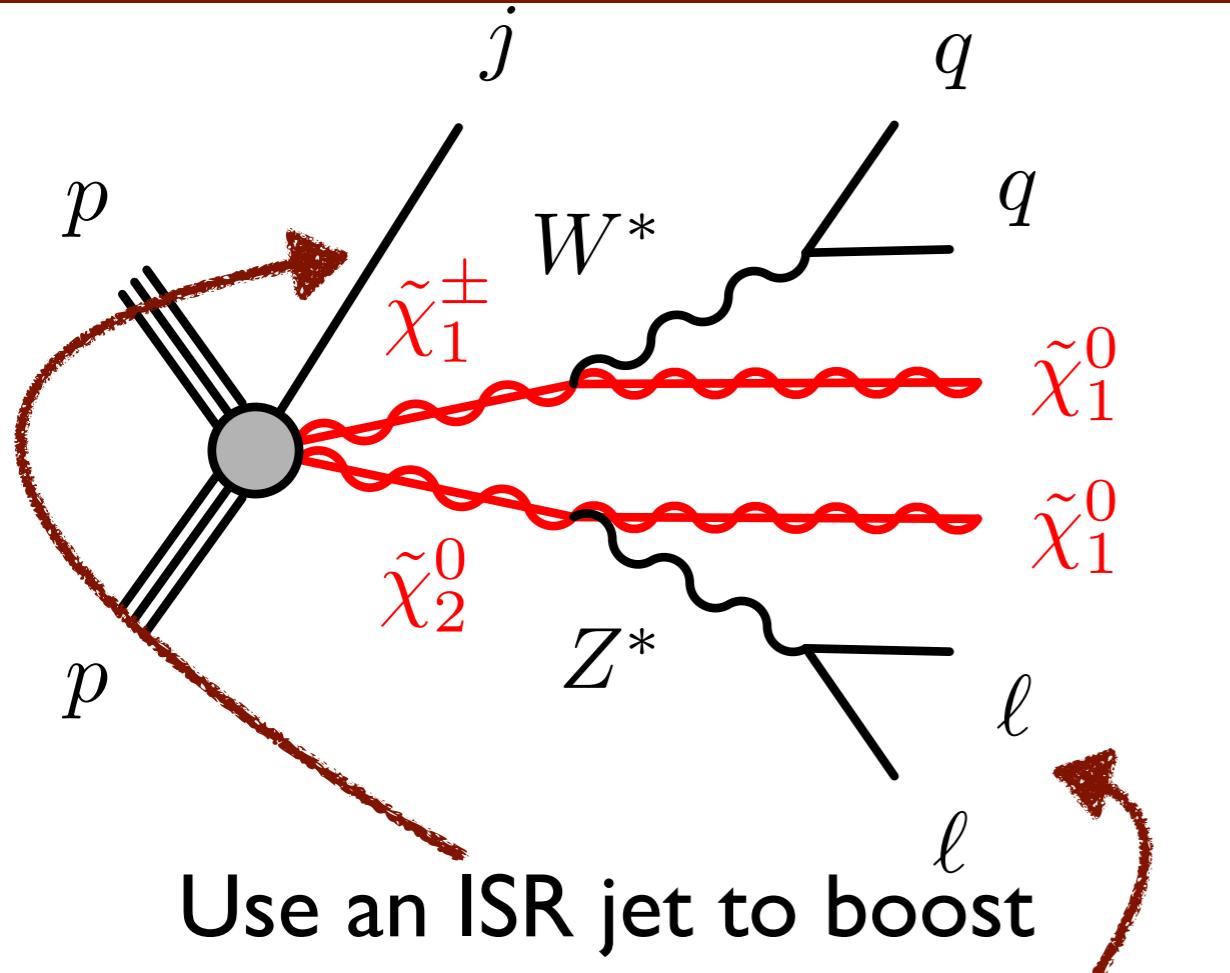
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Then, look for low-mass Z^* decays

Critical to reconstruct leptons
at extremely low p_T



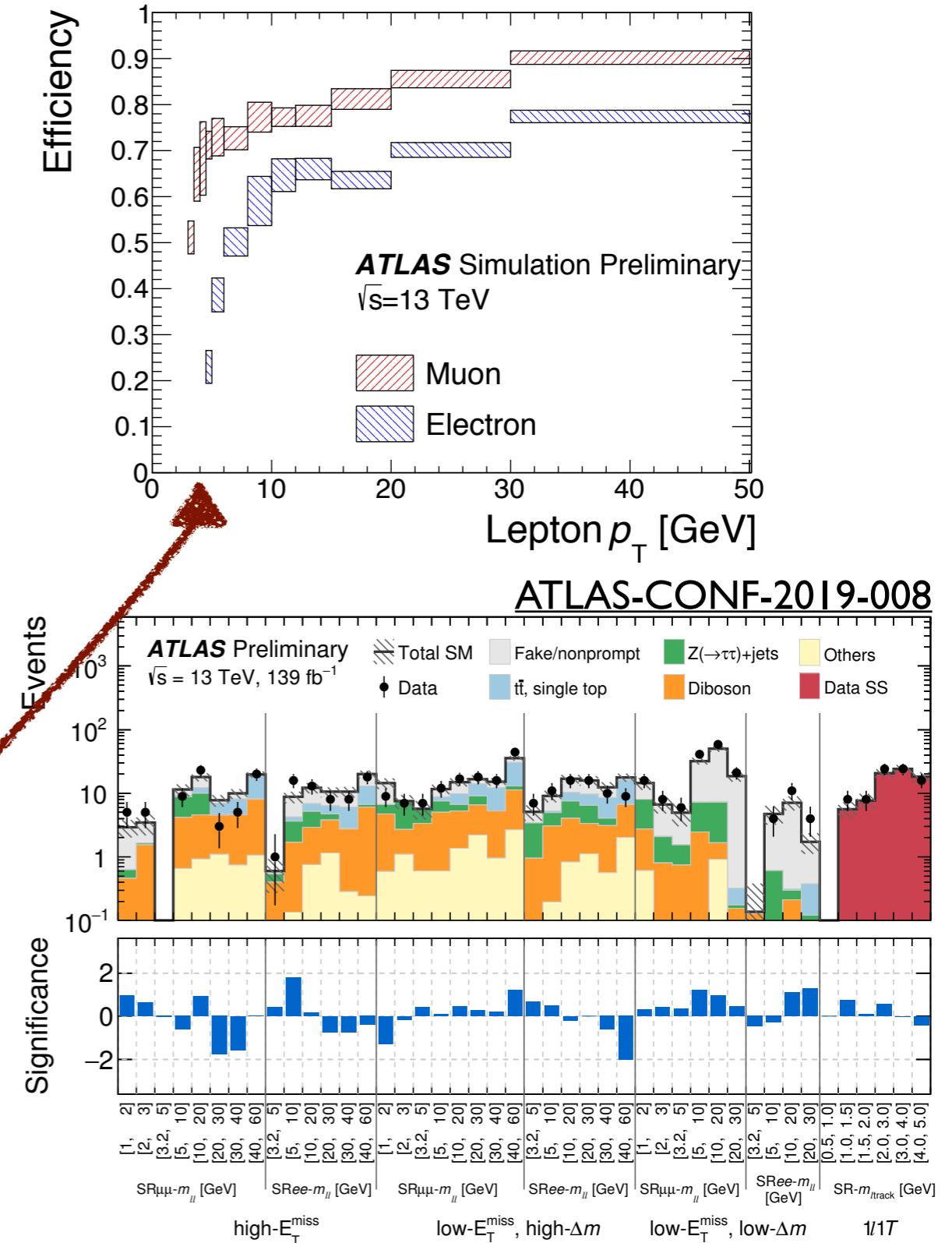
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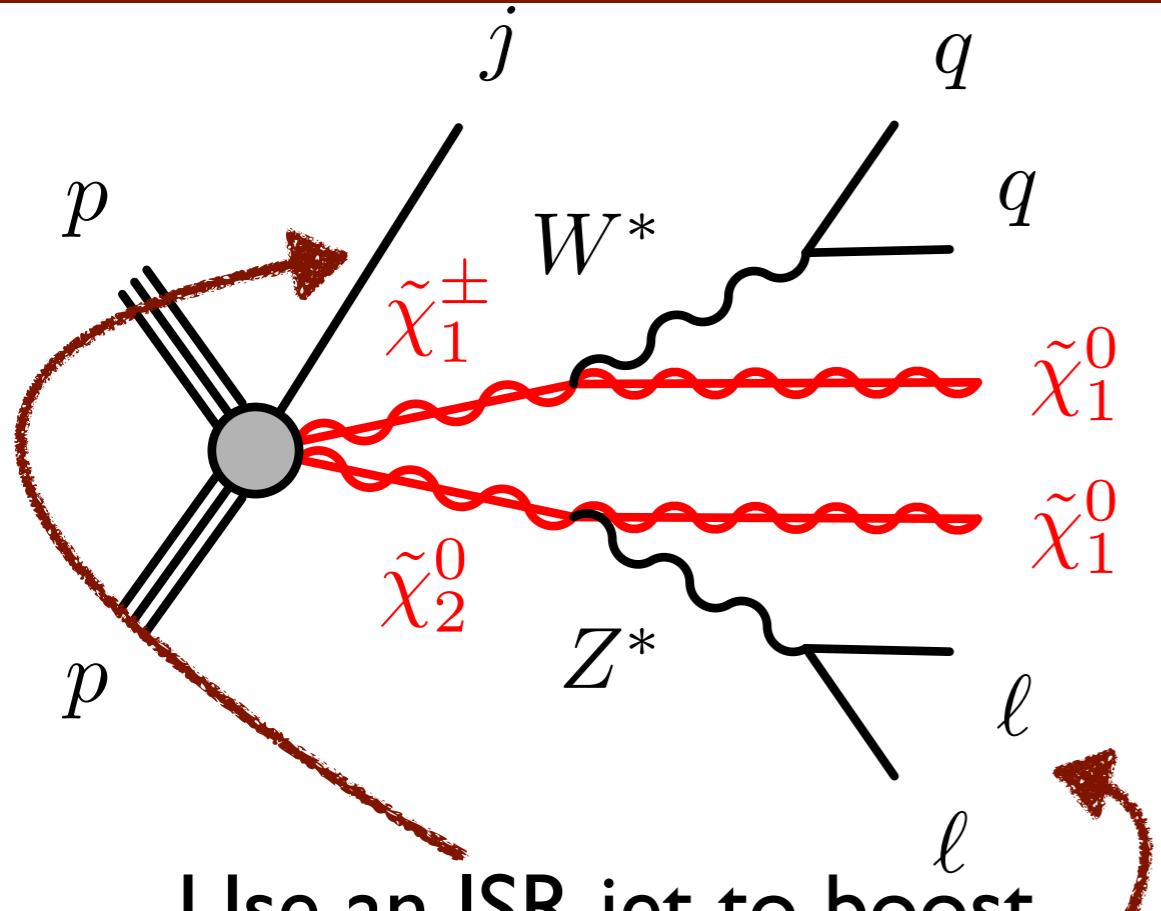
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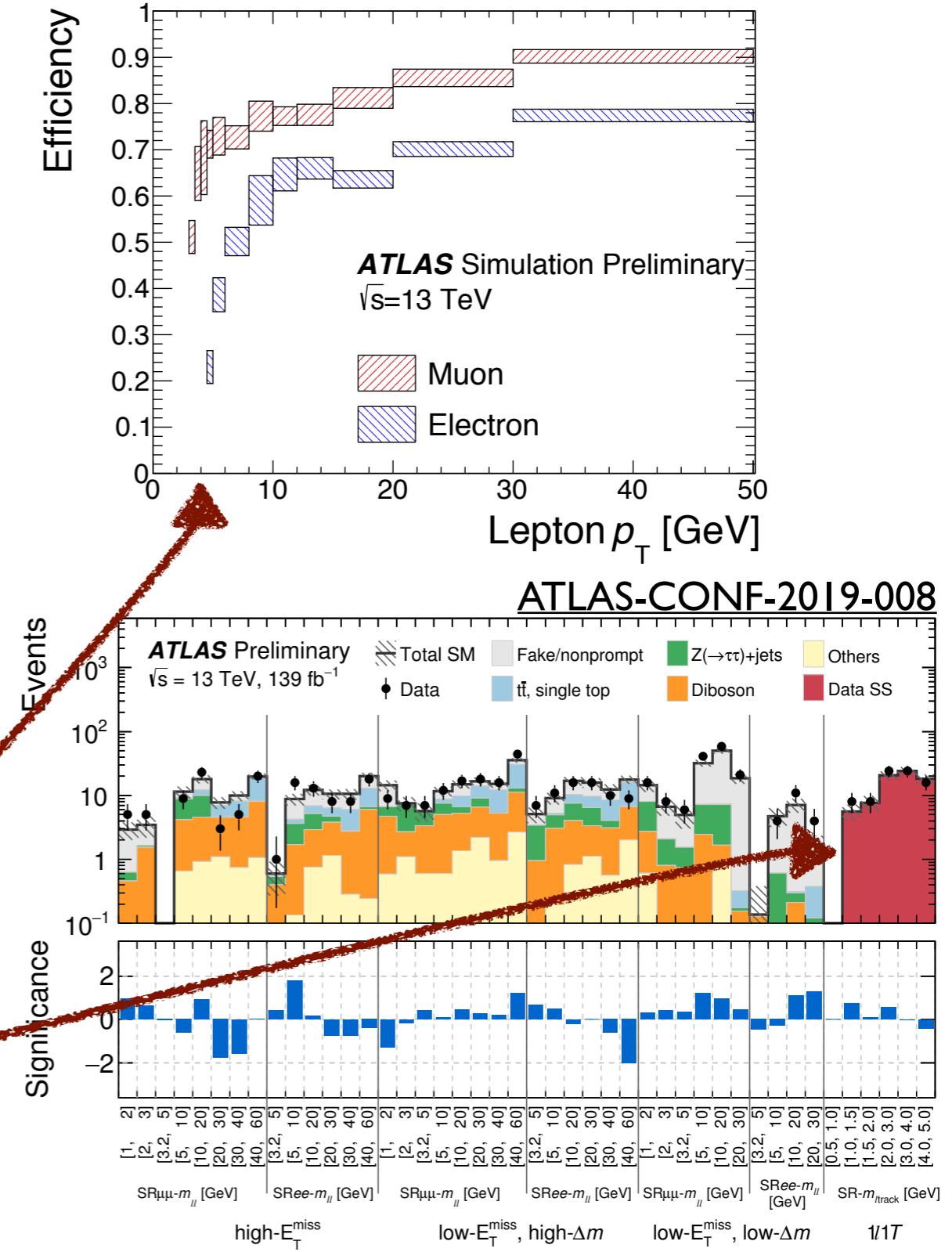


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Critical to reconstruct leptons at extremely low p_T

New lepton + track region extends sensitivity to even lower masses!

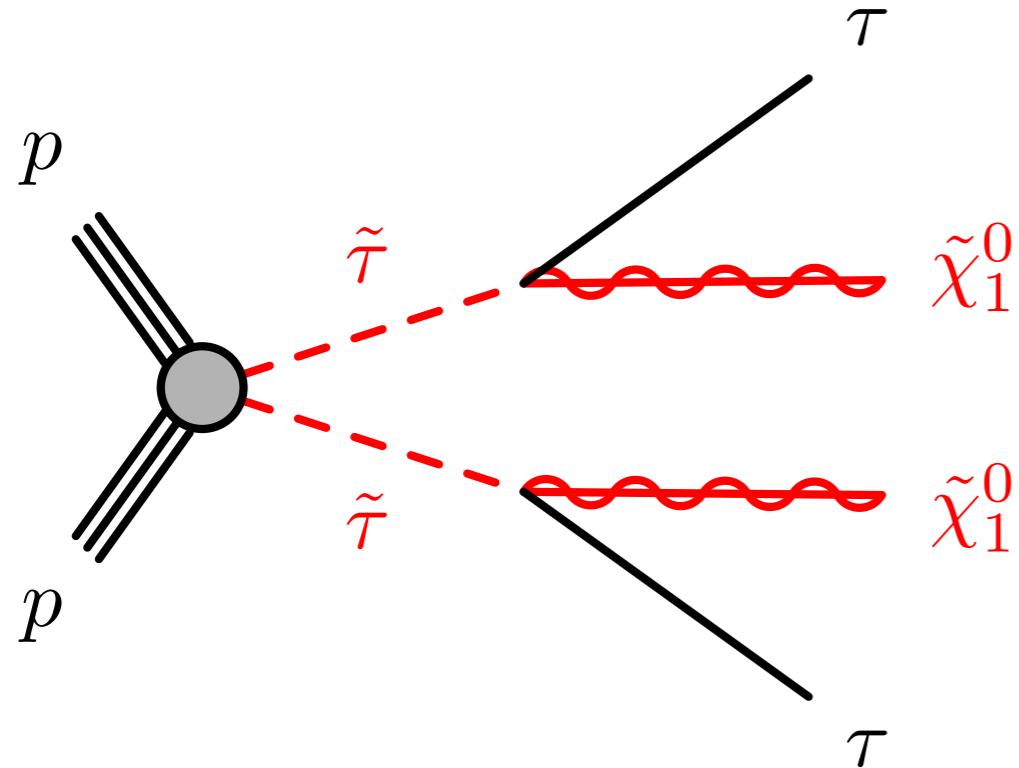




Even Rarer: Staus

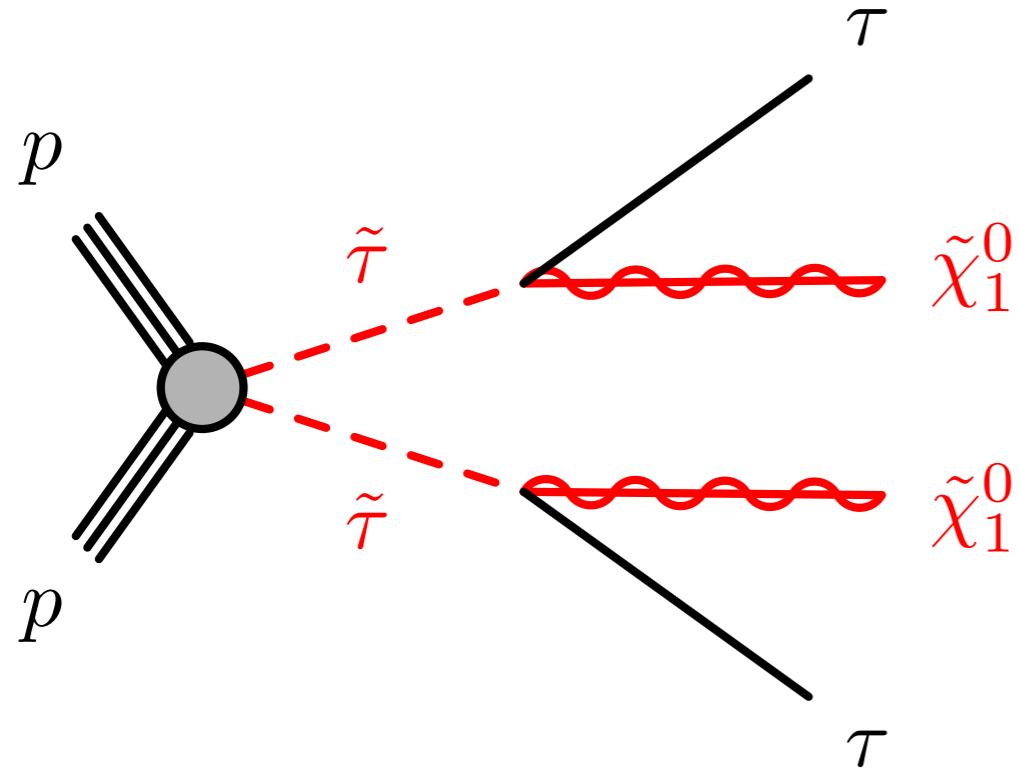


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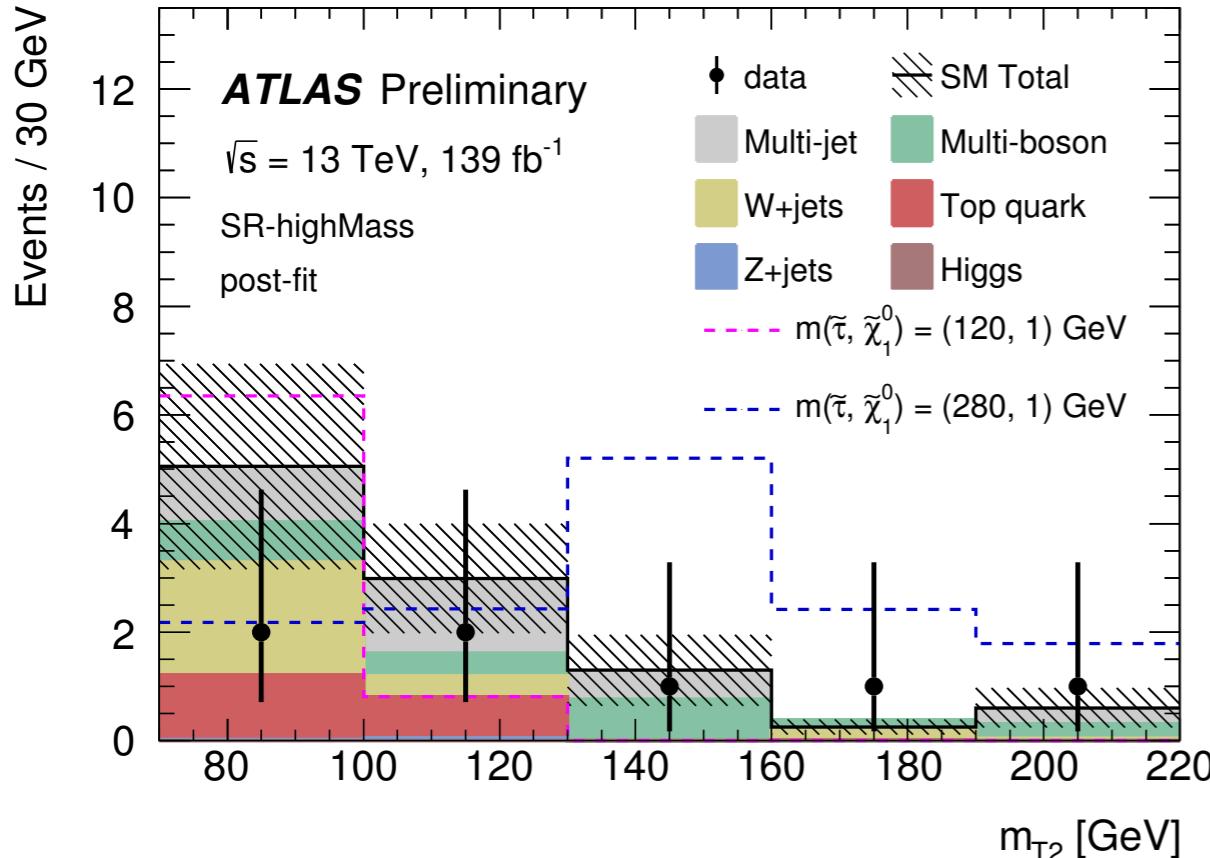
Even Rarer: Staus



Staus are an even larger challenge:
incredibly low cross-section!

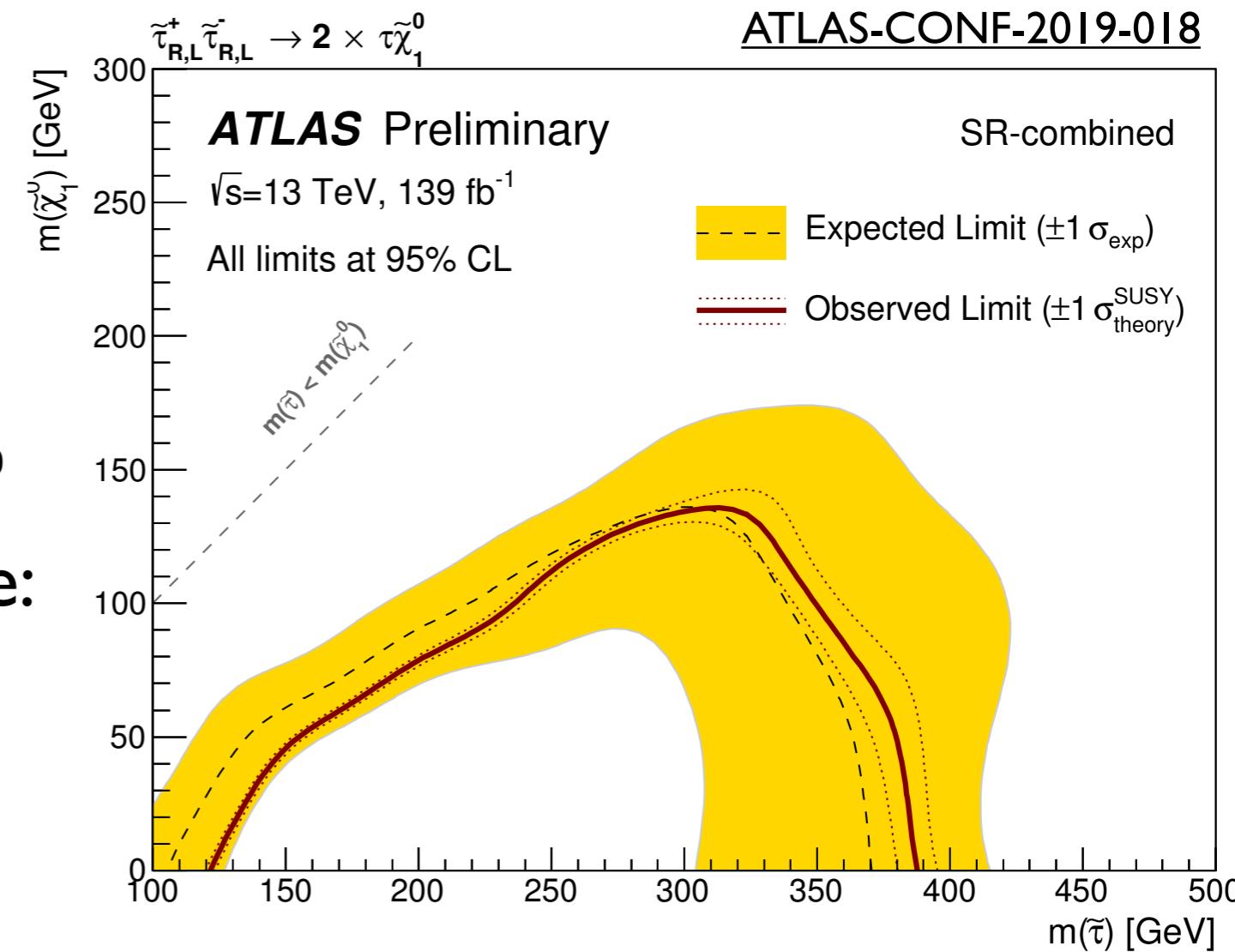


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Effective triggers and the large
dataset enable first substantial
sensitivity at the LHC!

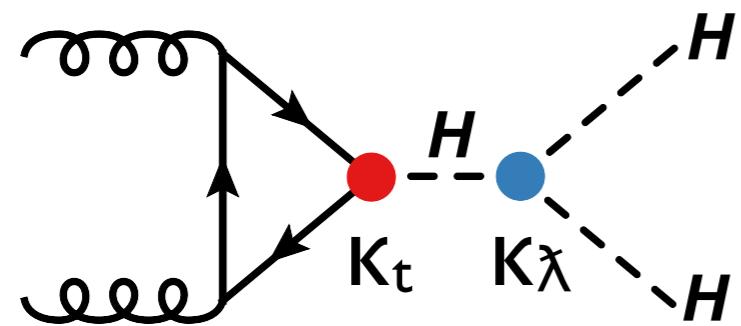


Di-Higgs at the LHC





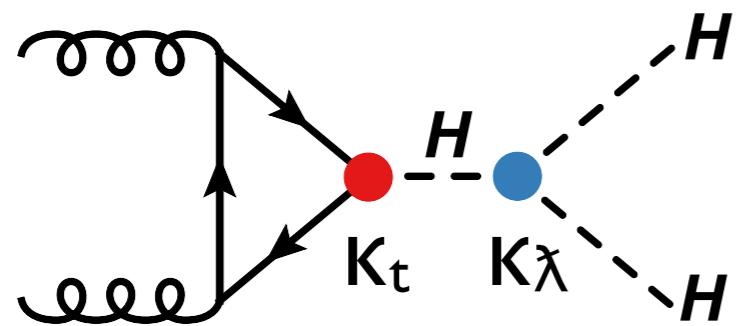
Di-Higgs at the LHC



Di-Higgs is another exciting new target for the LHC



Di-Higgs at the LHC

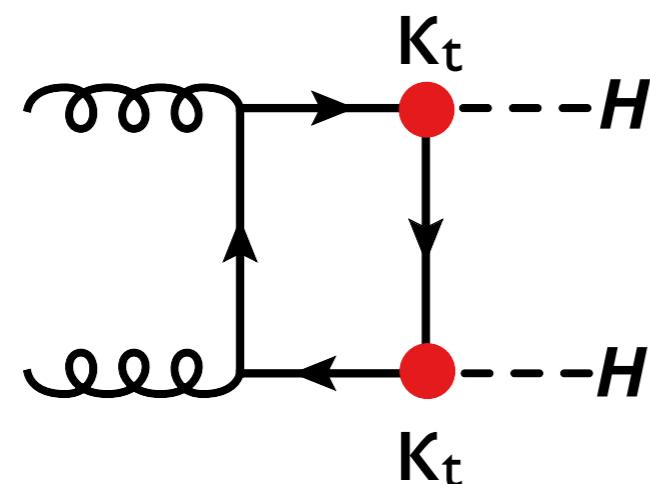
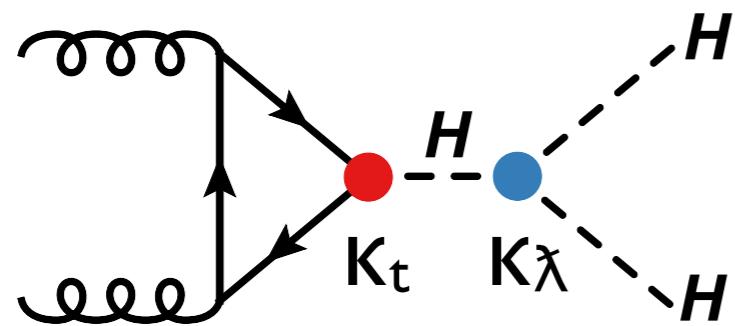


Di-Higgs is another exciting new target for the LHC

Exciting signal: can reveal the shape of the Higgs potential!



Di-Higgs at the LHC



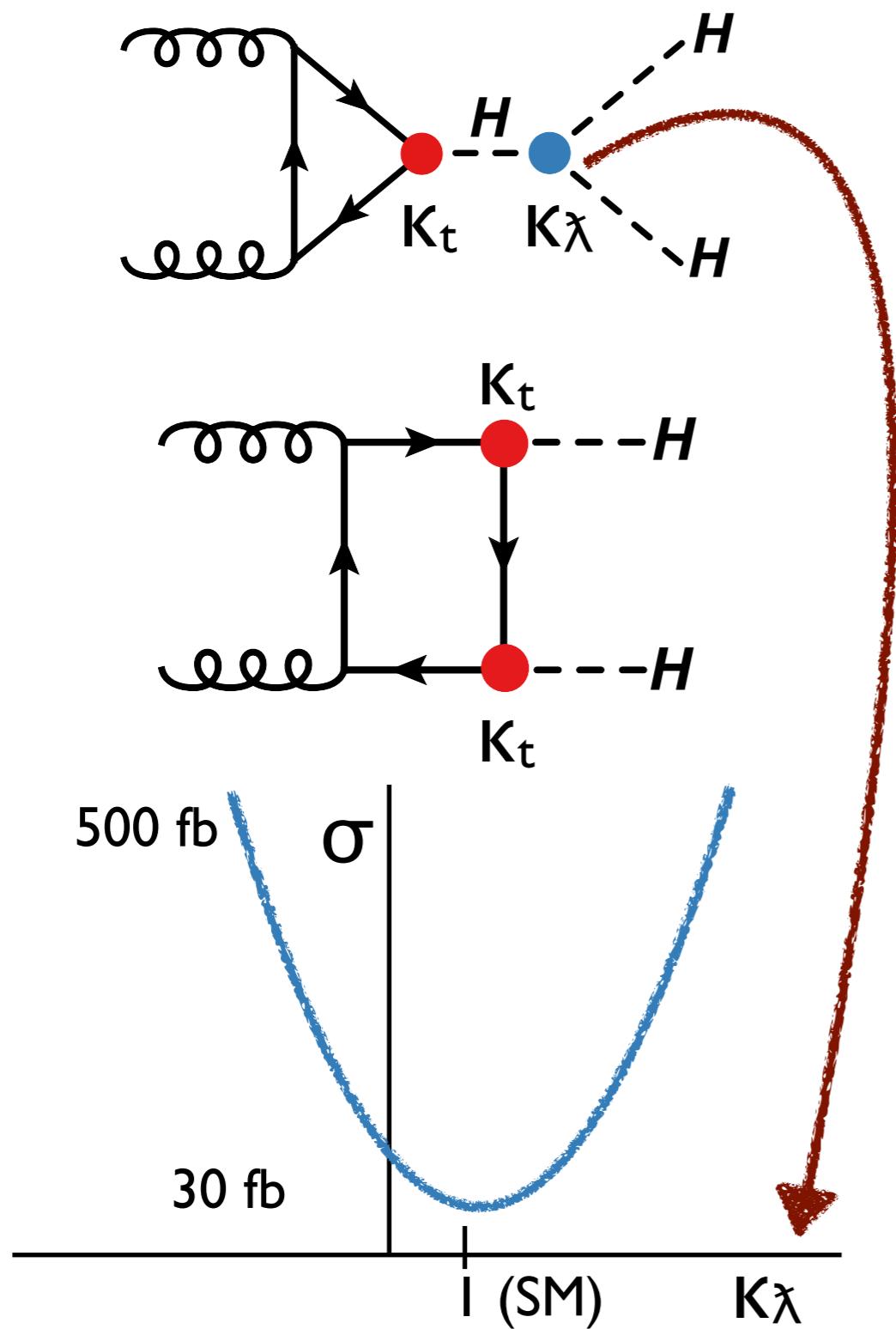
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Interference between SM diagrams leads to very low x-sec



Di-Higgs at the LHC



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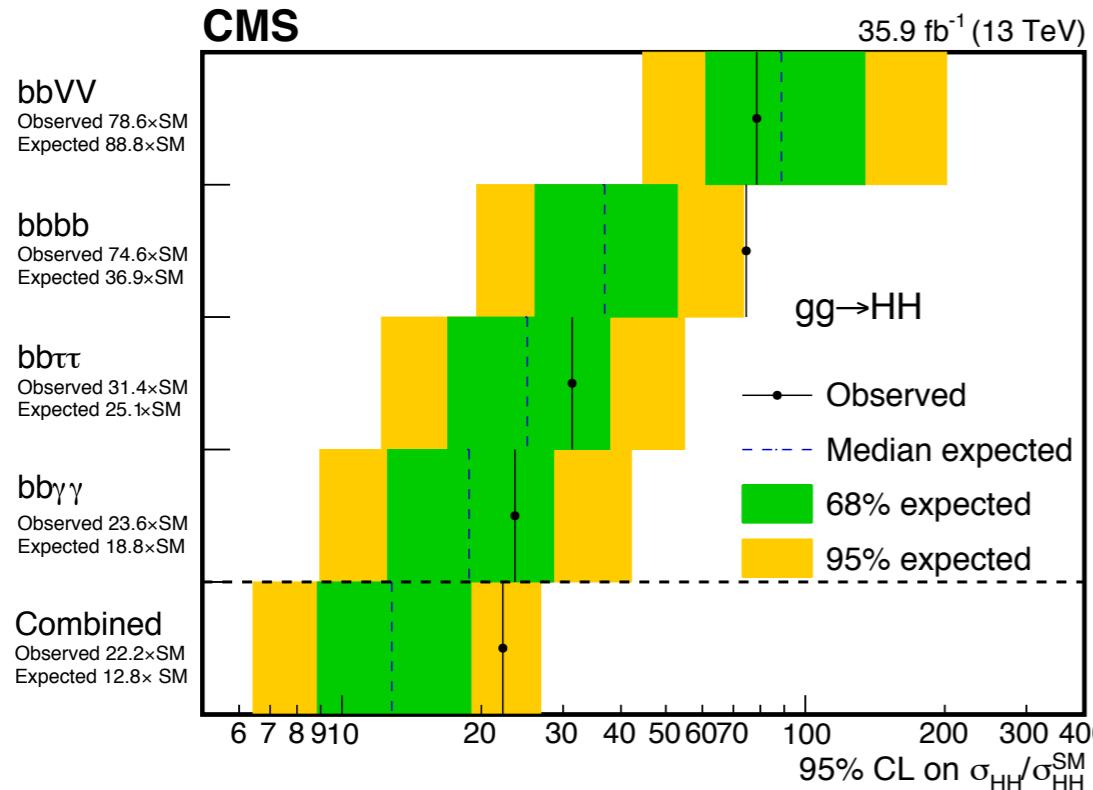
But small deviations from the SM can lead to huge x-sec increases!

Results on Di-Higgs





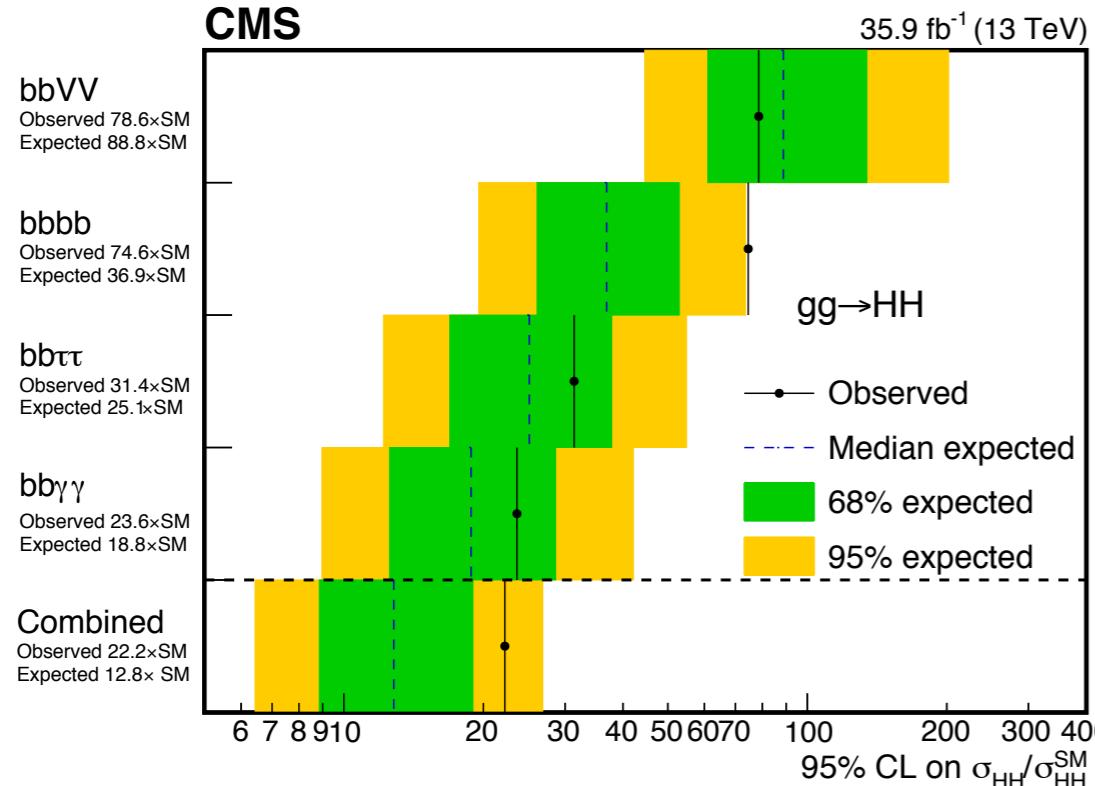
Results on Di-Higgs



Many orthogonal channels
can be combined to
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Results on Di-Higgs

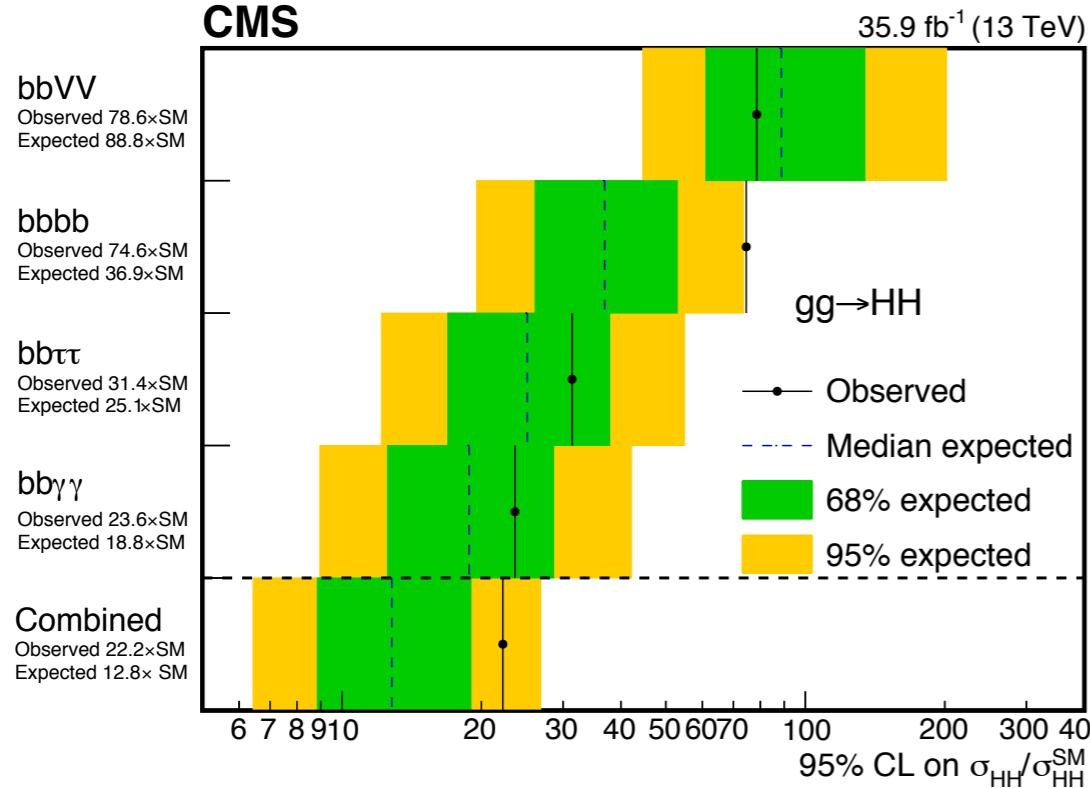


Many orthogonal channels can be combined to increase sensitivity!

Note the use of b-jets: rare signals mean that high-BR decays of the Higgs need to be used

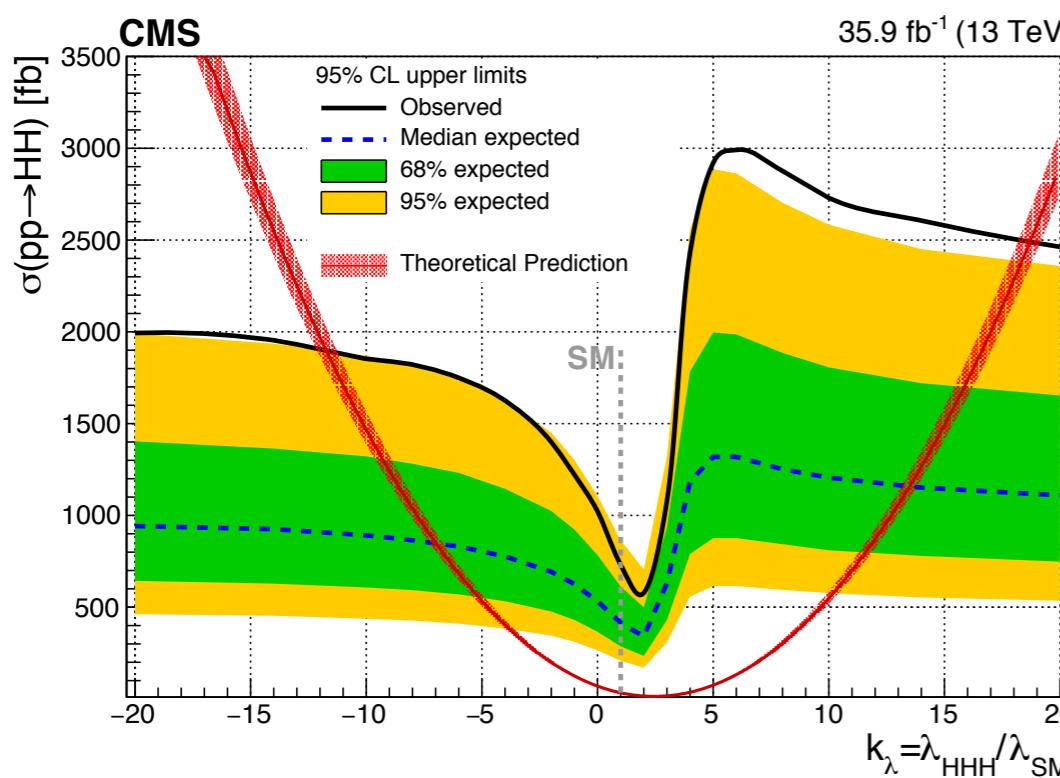


Results on Di-Higgs



Many orthogonal channels can be combined to increase sensitivity!

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Wide range of couplings still allowed!
Even more data still necessary

Hunting Rare SM Deviations



Hunting Rare SM Deviations

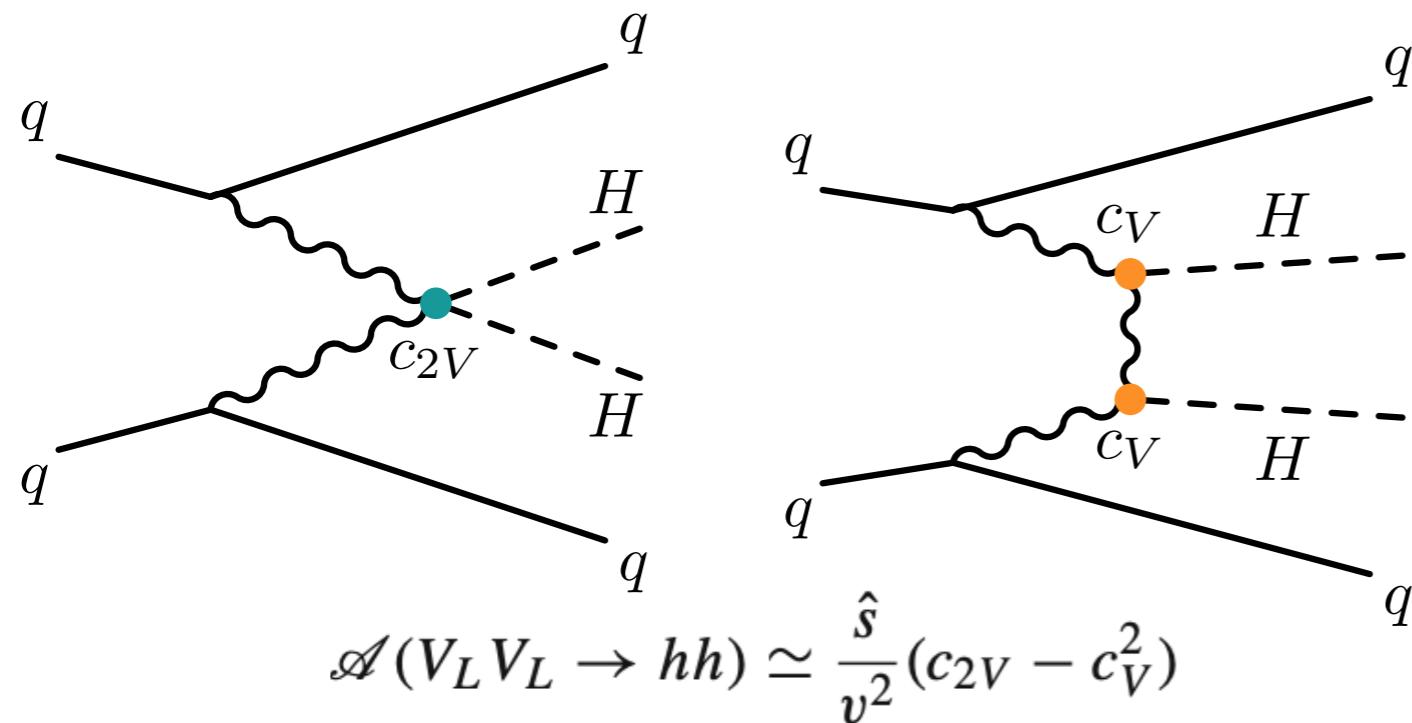


Interference in other SM
diagrams lead to similar
vanishingly small x-sec



Hunting Rare SM Deviations

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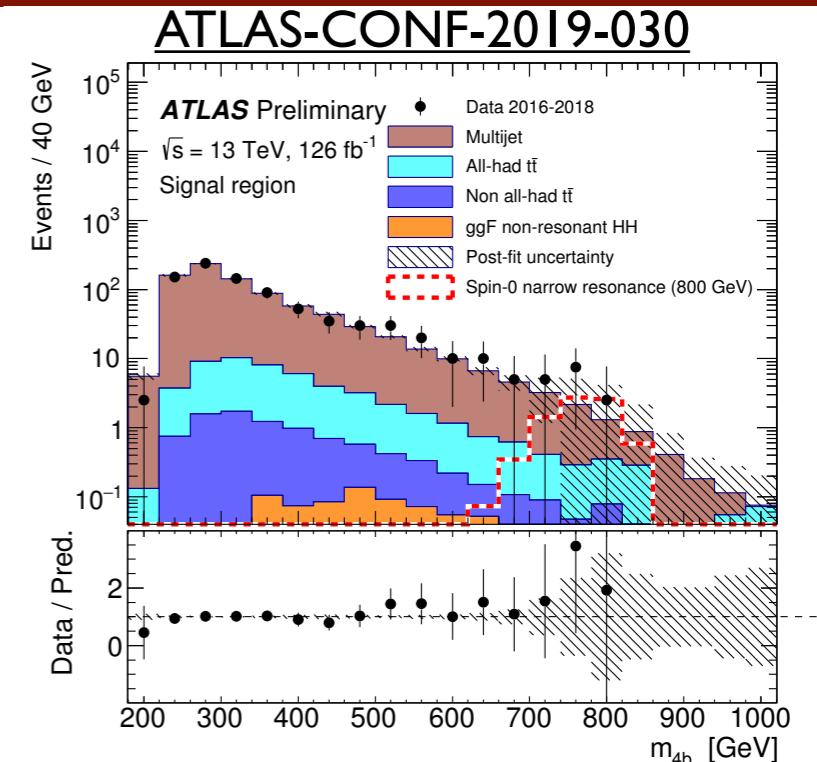
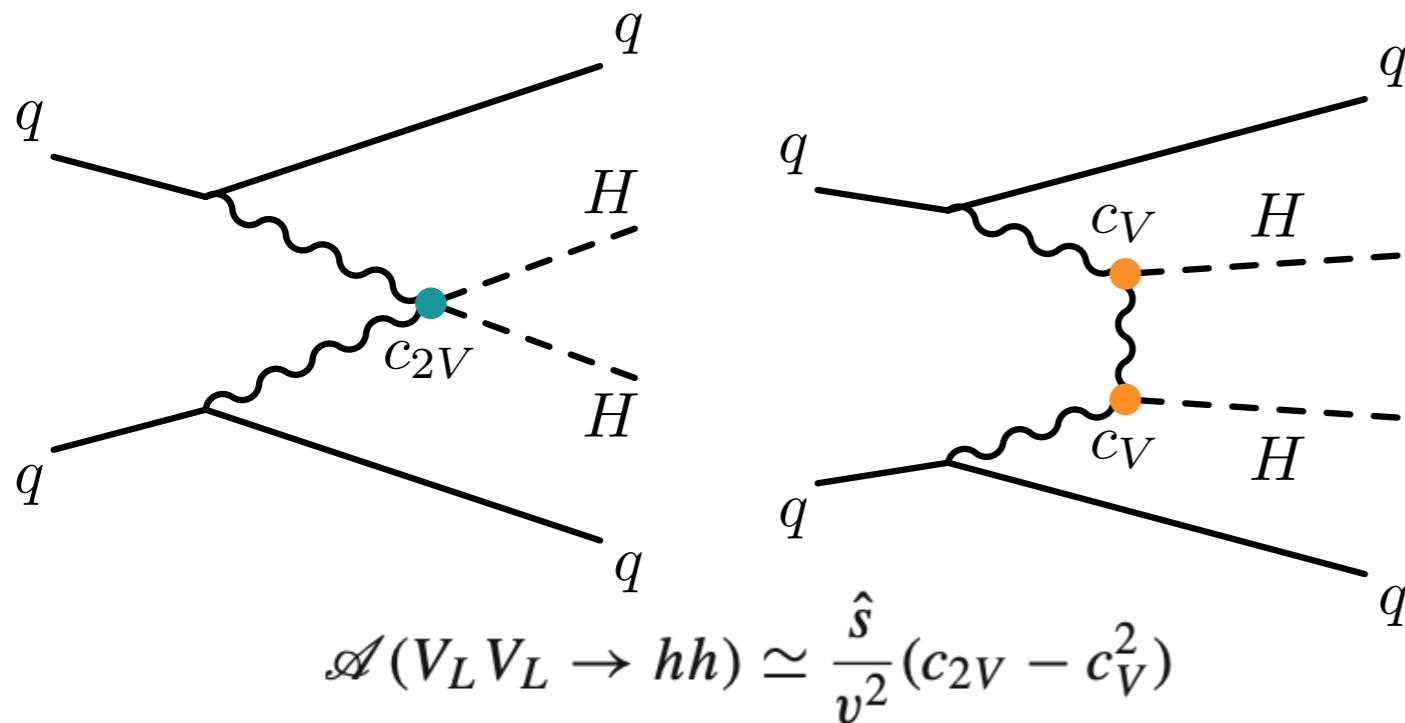


VBF di-Higgs production
is one example!



Hunting Rare SM Deviations

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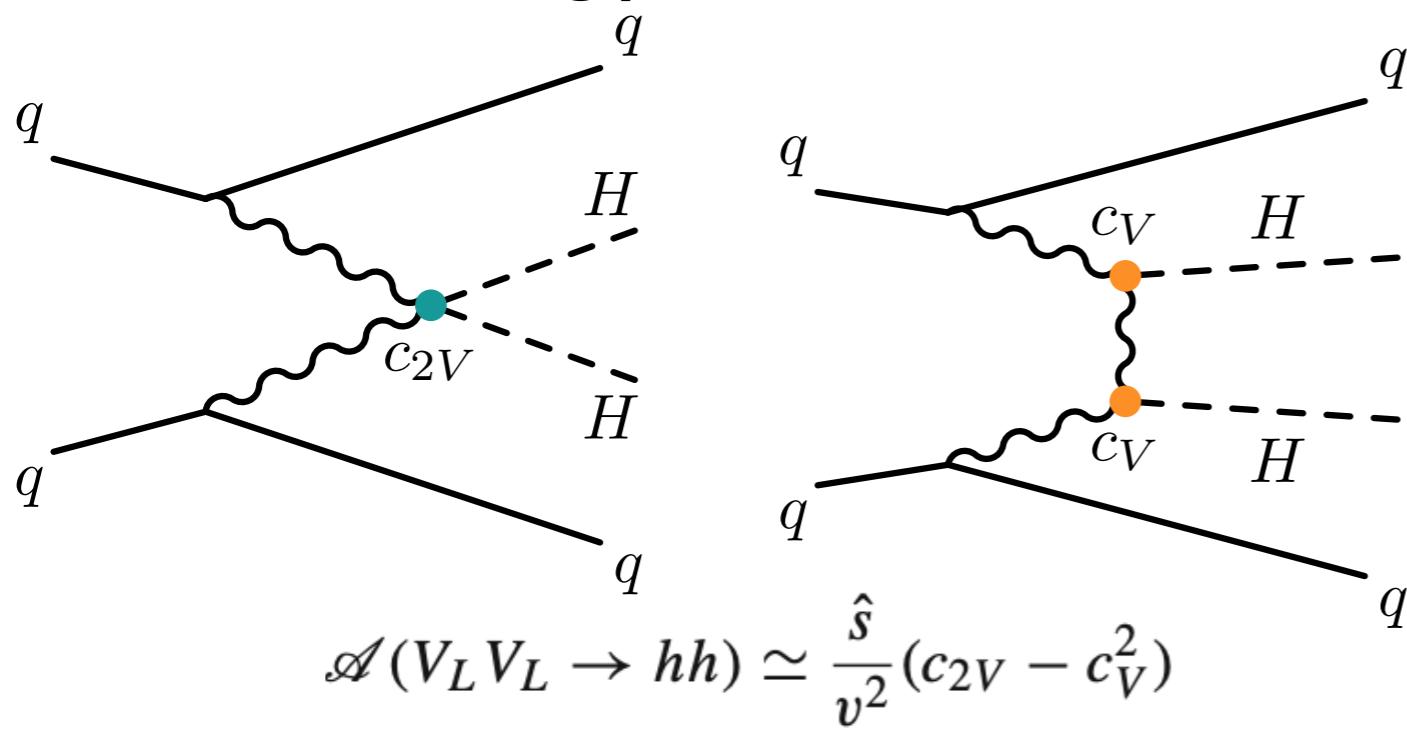


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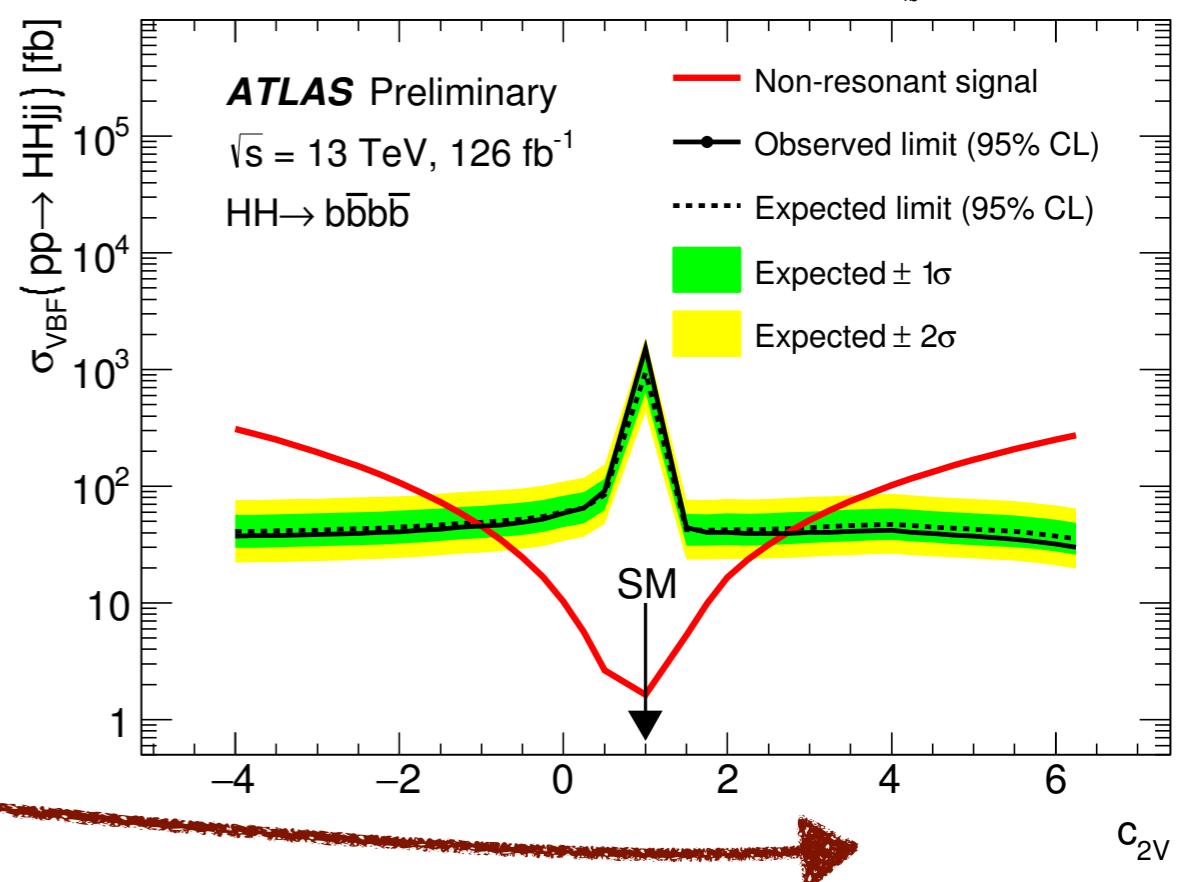
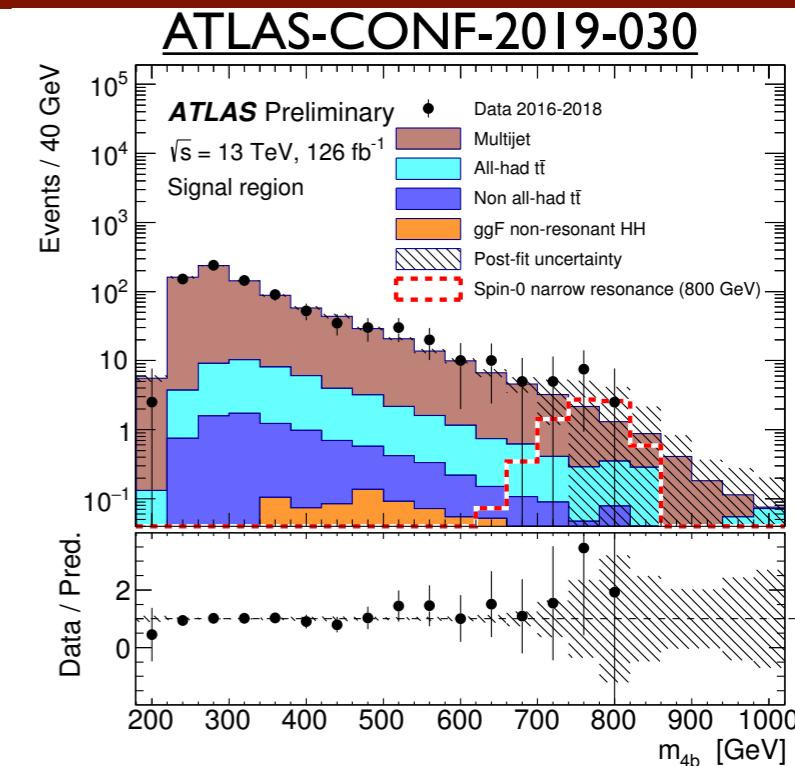
Hunting Rare SM Deviations

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First limits on the 4-point
VVHH coupling

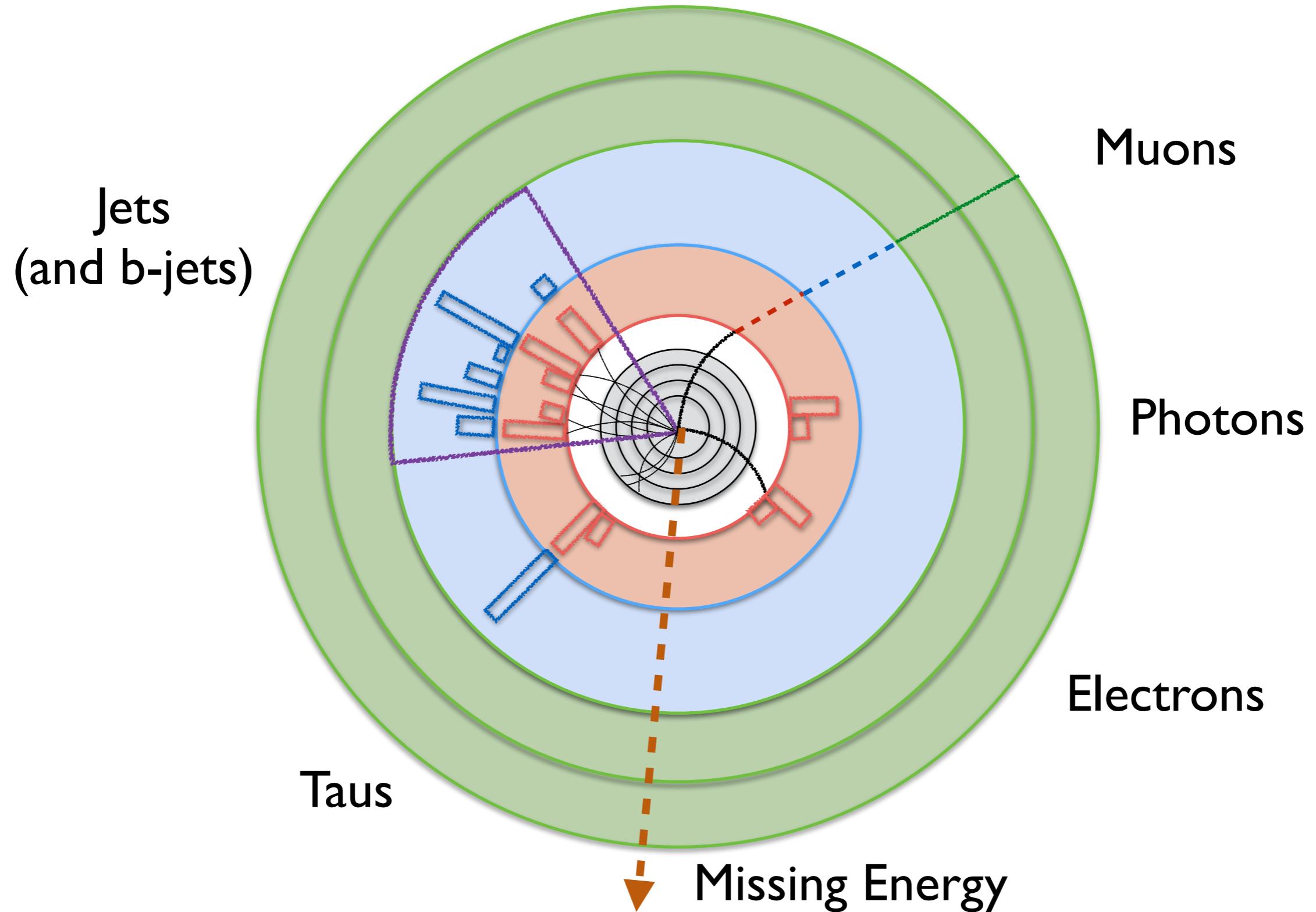


Not Your Advisor's Signals

*Or: how new reconstruction strategies are expanding
our sensitivity to signals our detectors weren't
built to measure*

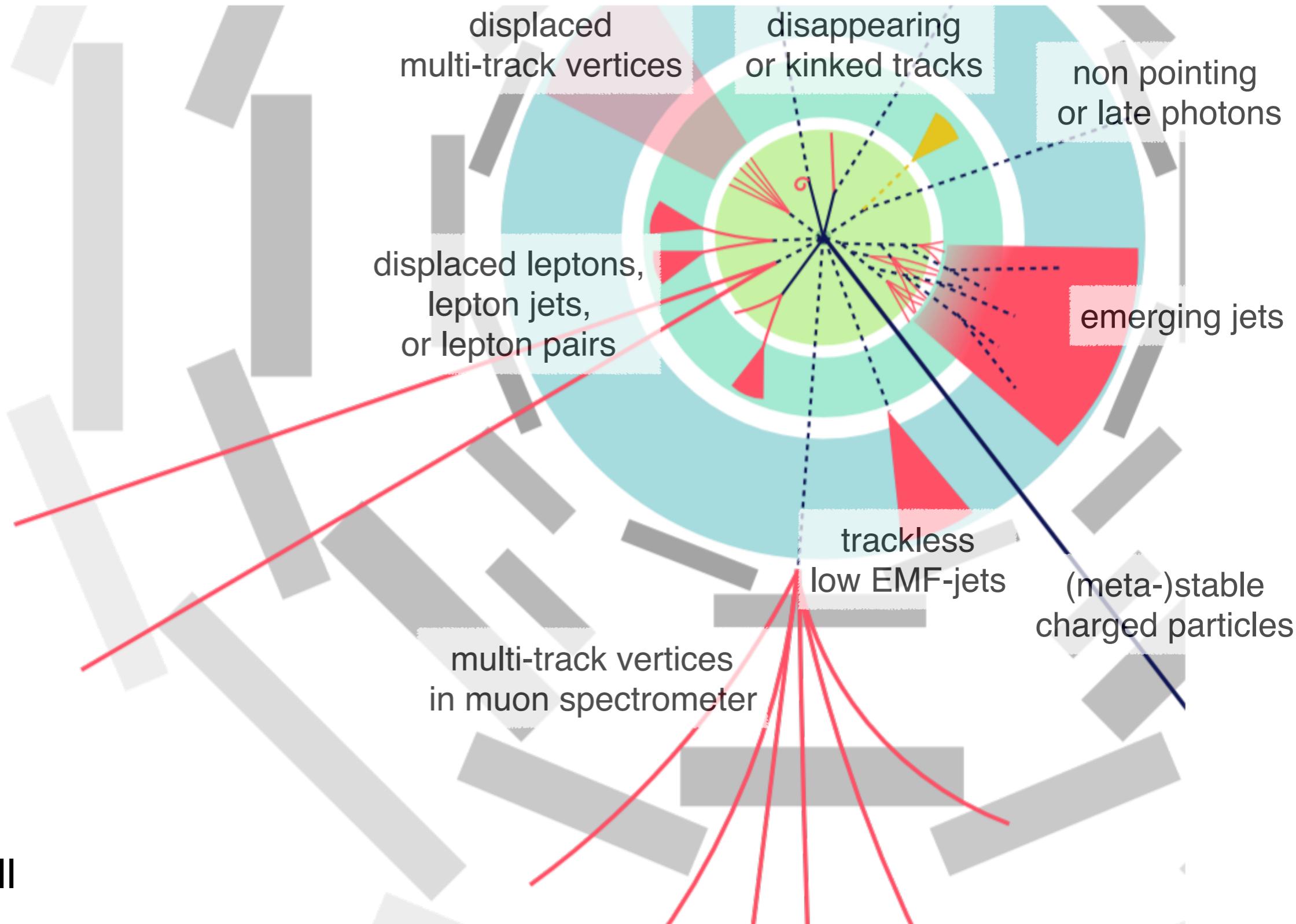


Your Advisor's Signals



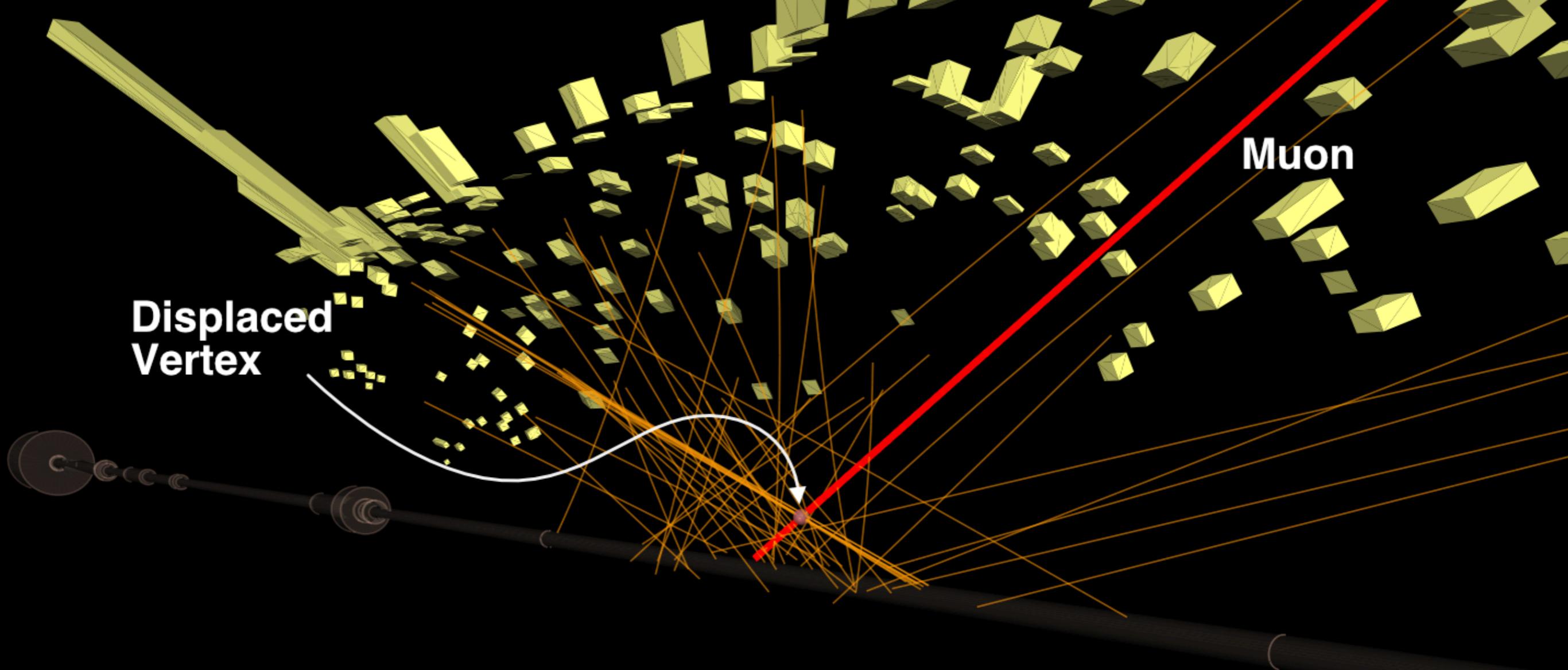


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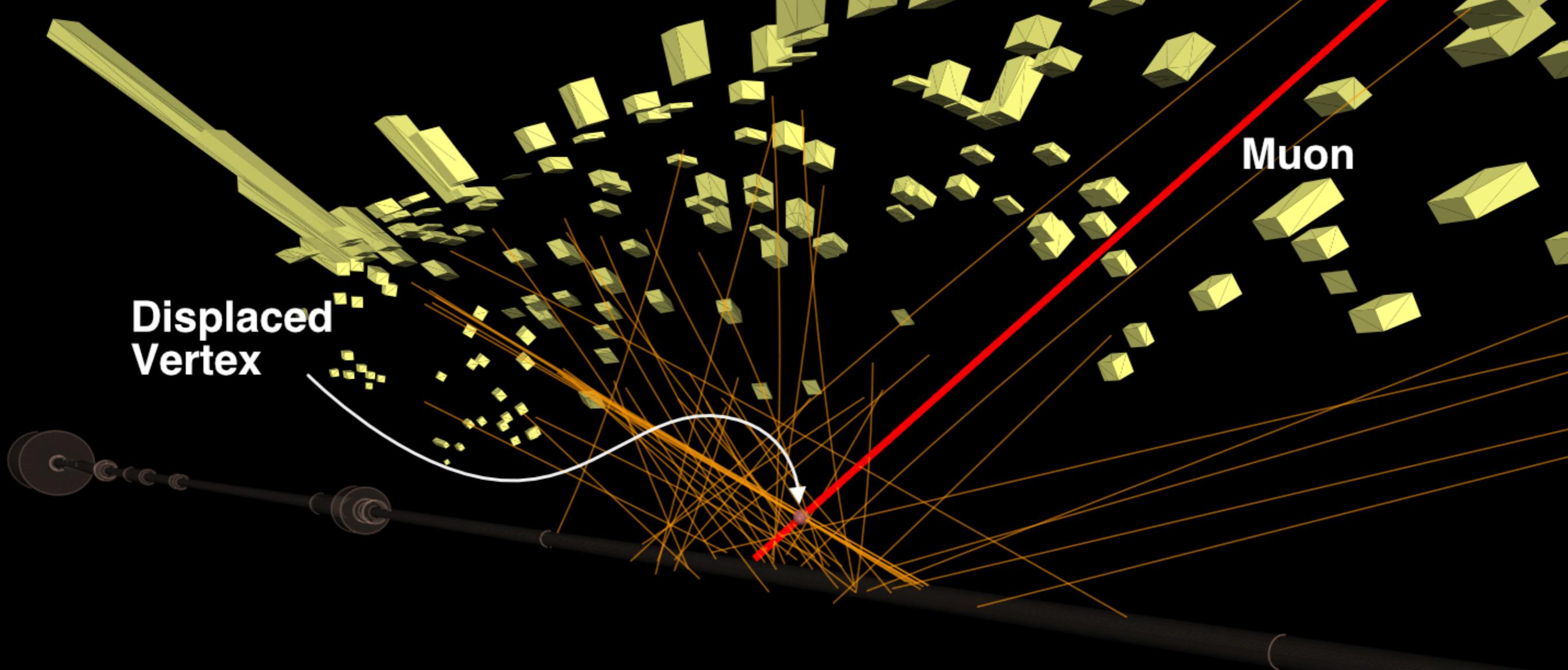
Heather Russell

Long-lived particles will
travel an appreciable distance
before they decay



Long-lived particles will
travel an appreciable distance
before they decay

Simulated Signal Event
Top Squark Pair Production
 $m(\tilde{t}) = 1.5 \text{ TeV}, \tau(\tilde{t}) = 1 \text{ ns}$
 $\tilde{t} \rightarrow \mu j$



Long-lived particles will travel an appreciable distance before they decay

Special algorithms can reconstruct these “displaced tracks” and look for displaced vertices

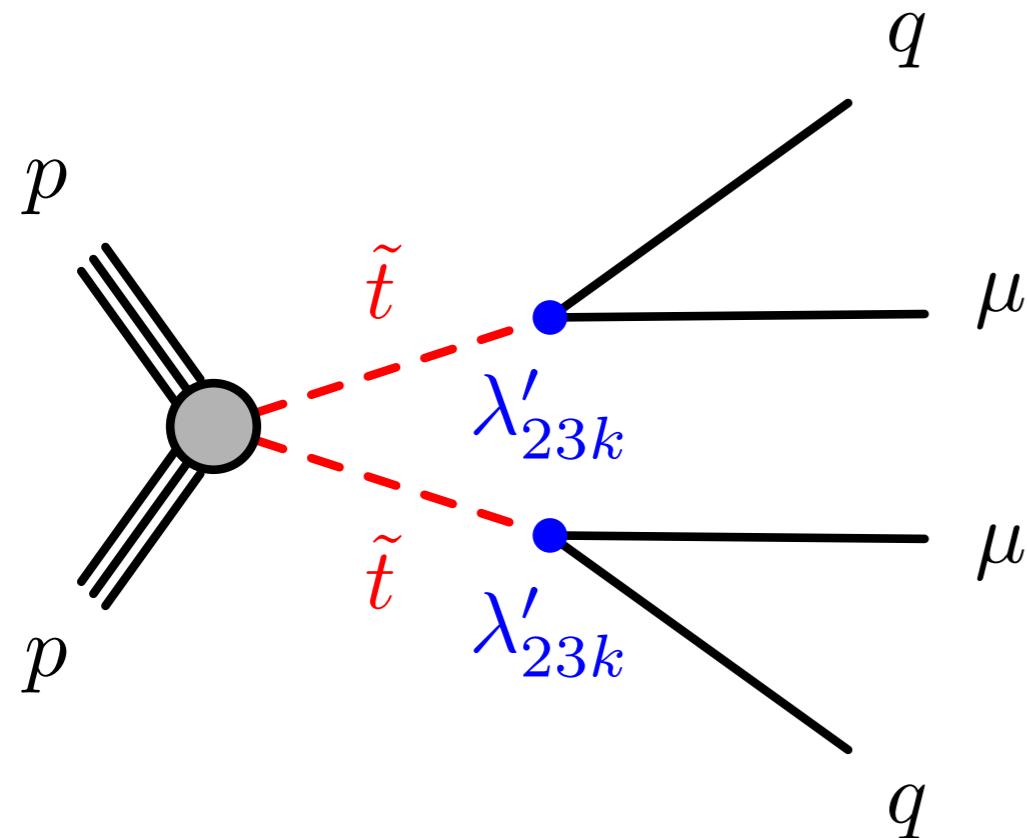
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Displaced Vertex + μ



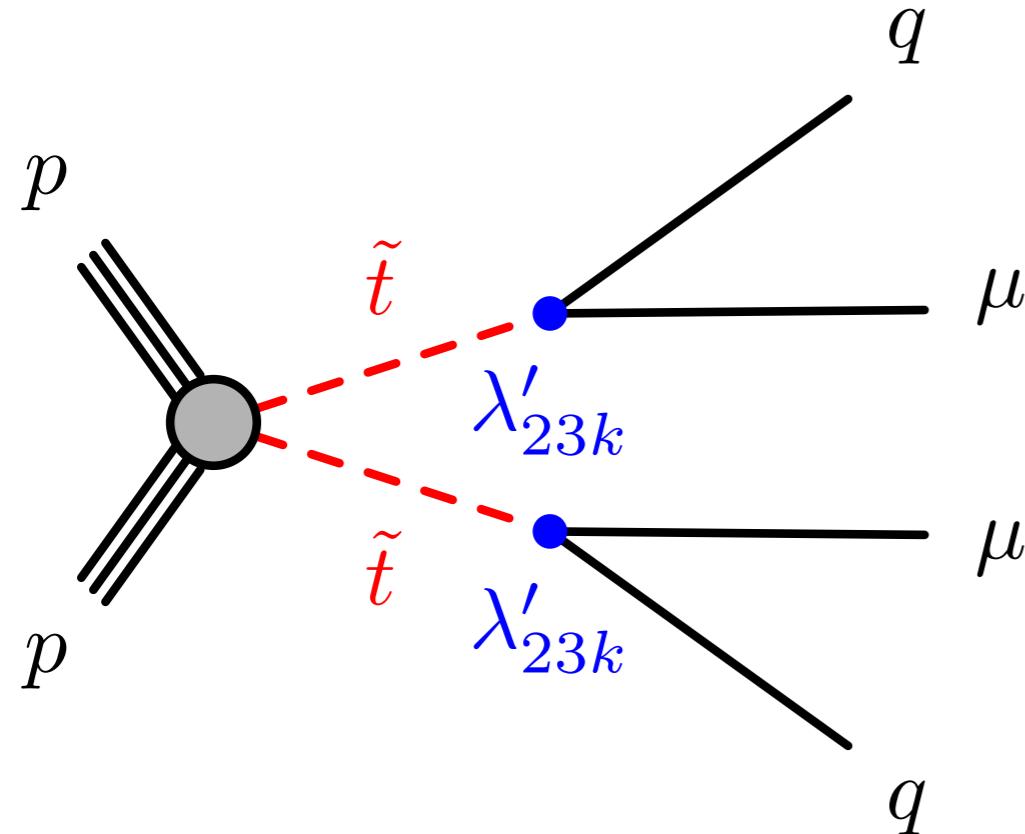


Displaced Vertex + μ





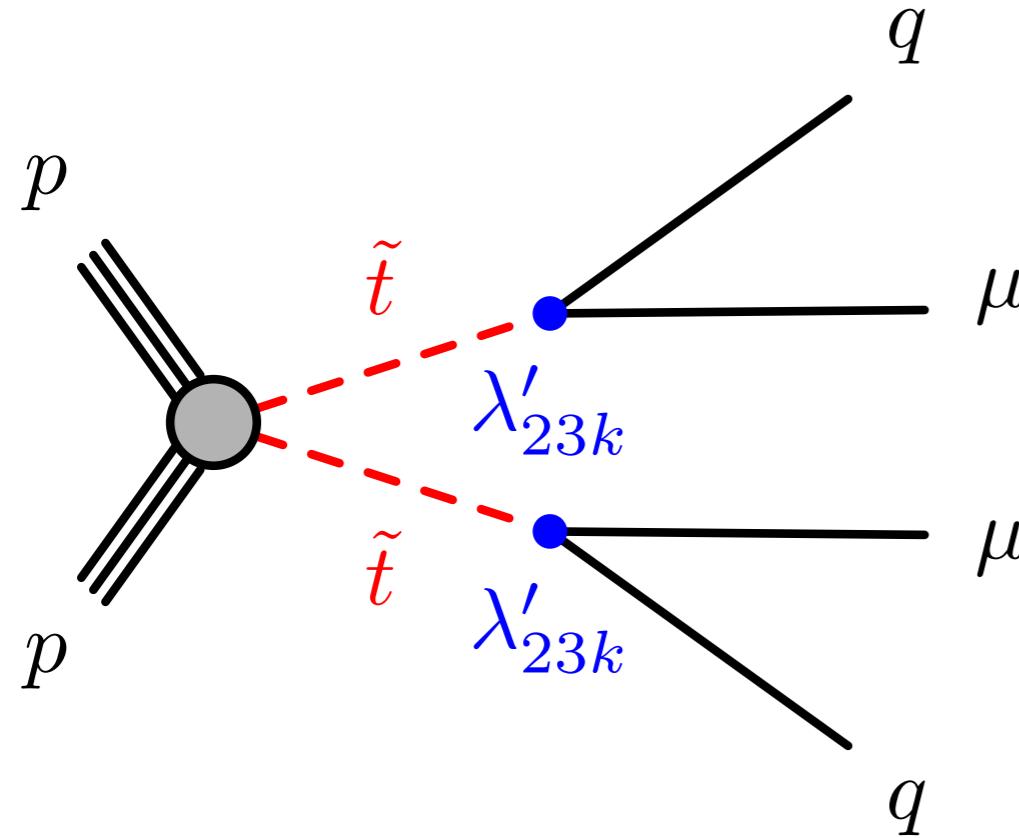
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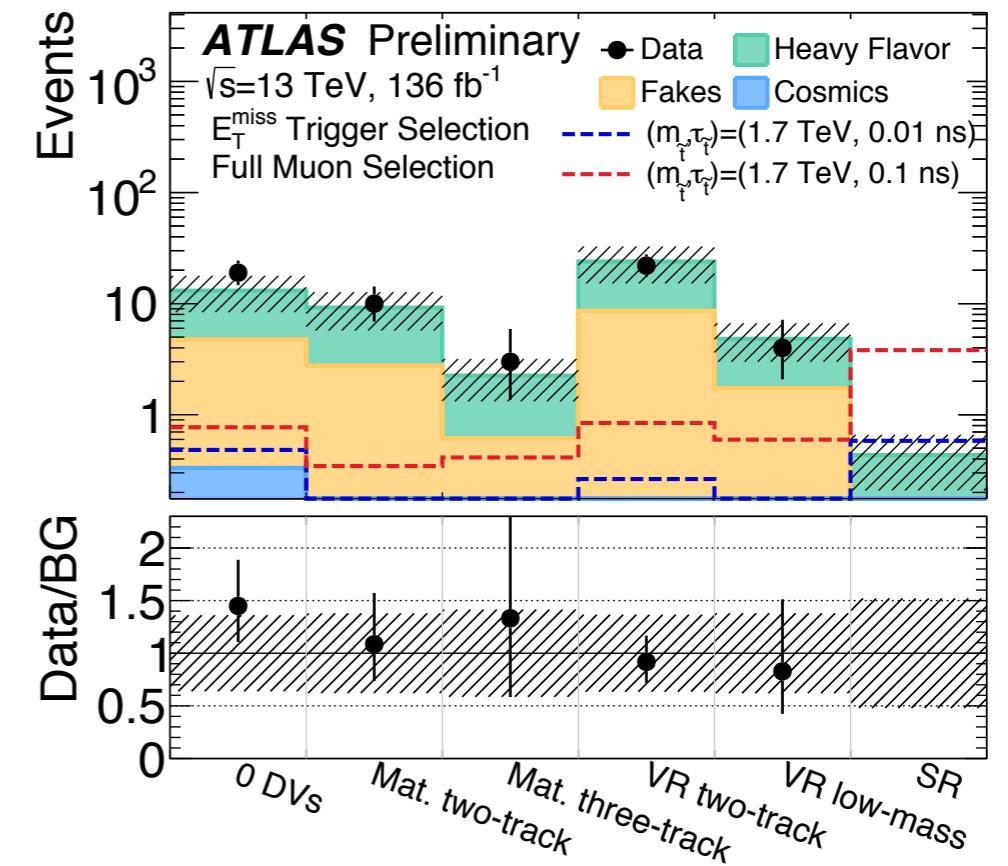
Stops can be long-lived in RPV
models with small couplings



Displaced Vertex + μ



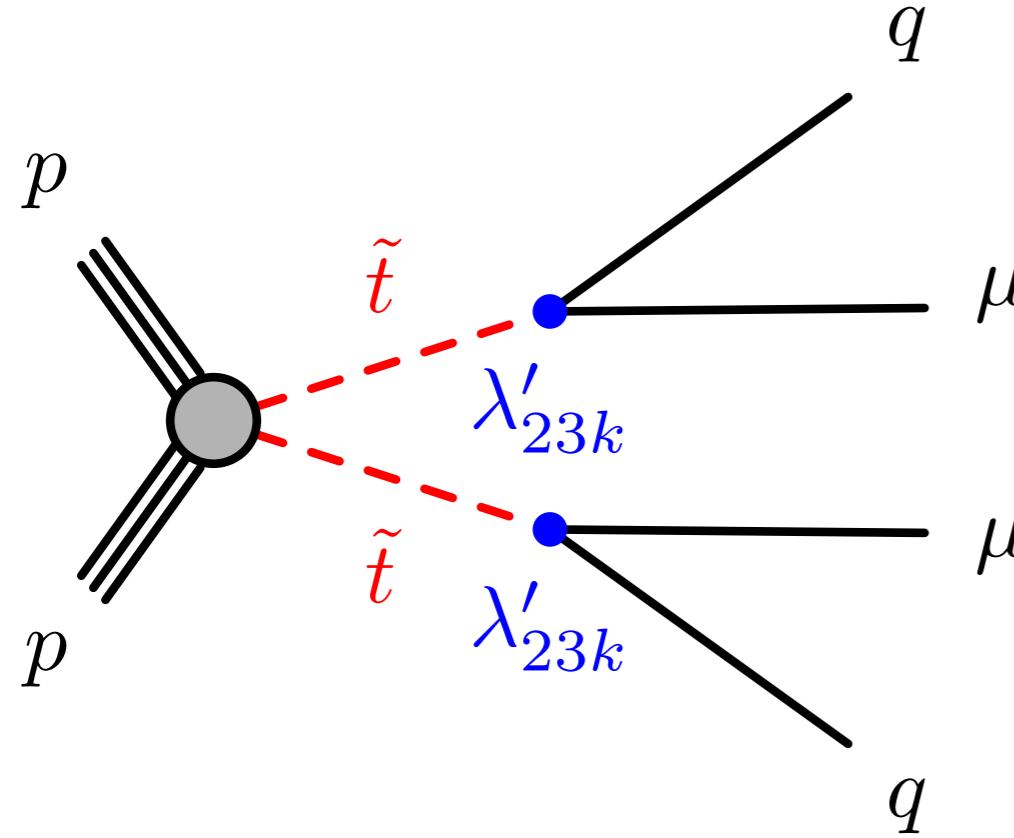
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Data-driven backgrounds estimate
~0 background, with no signal observed



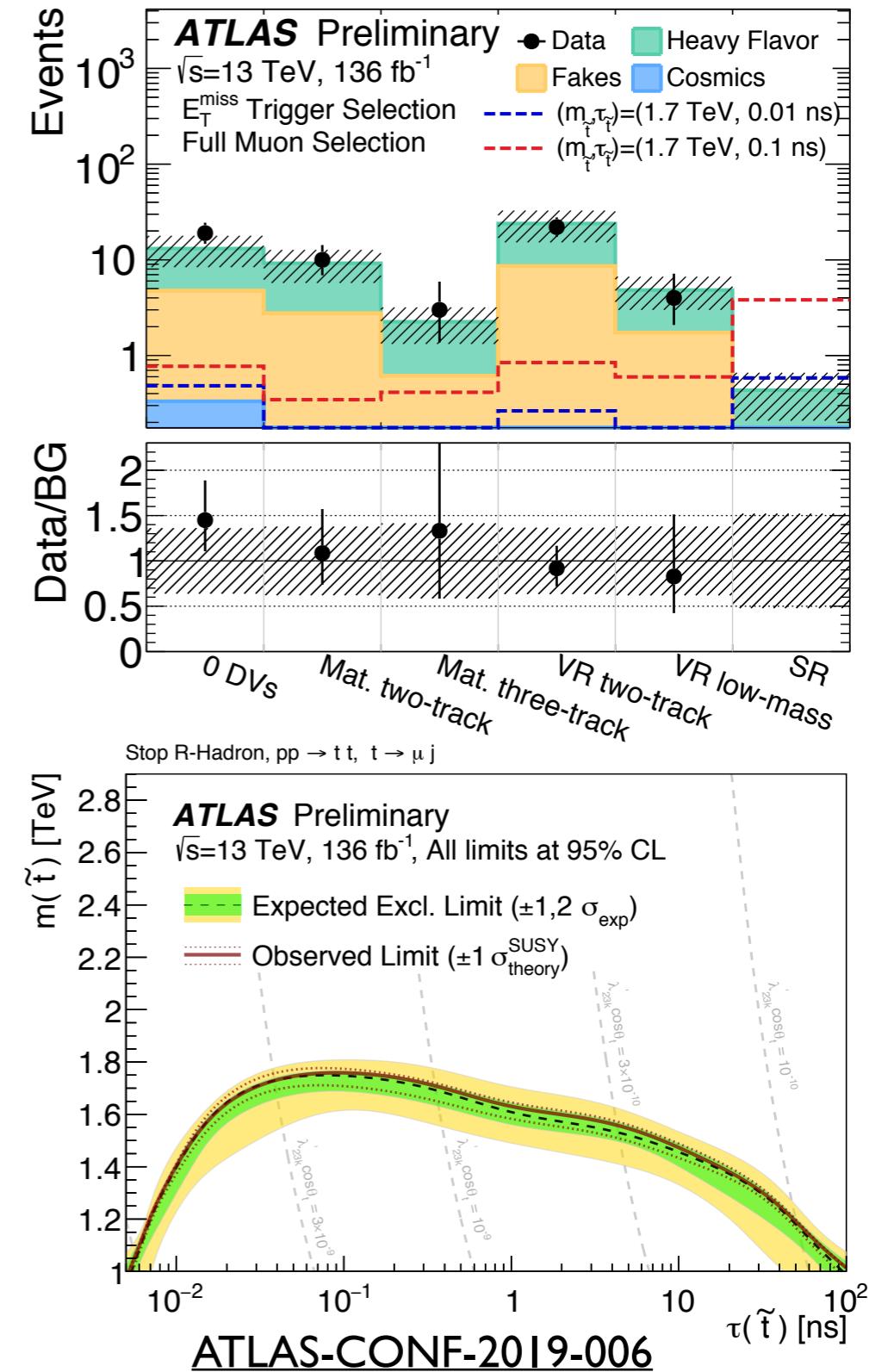
Displaced Vertex + μ



Stops can be long-lived in RPV models with small couplings

Data-driven backgrounds estimate
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Strongest limits yet on stops!

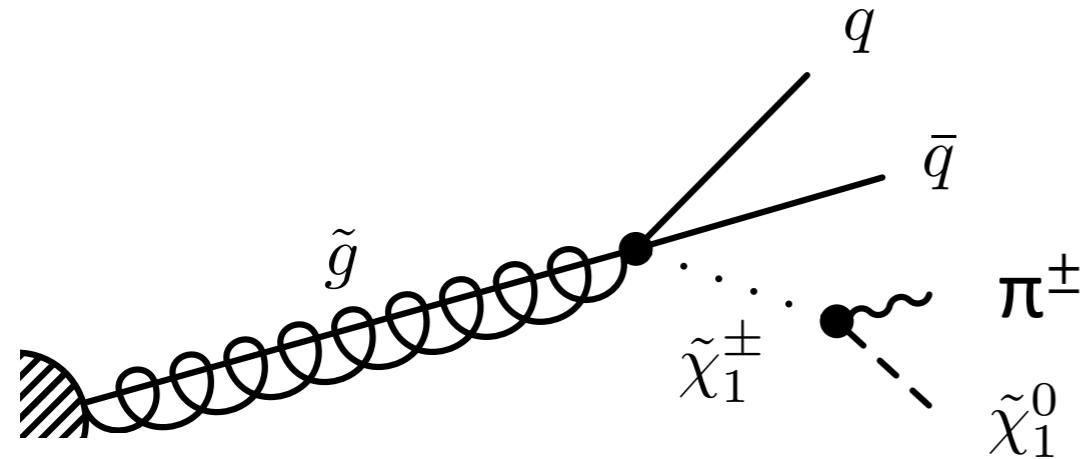


Disappearing Tracks





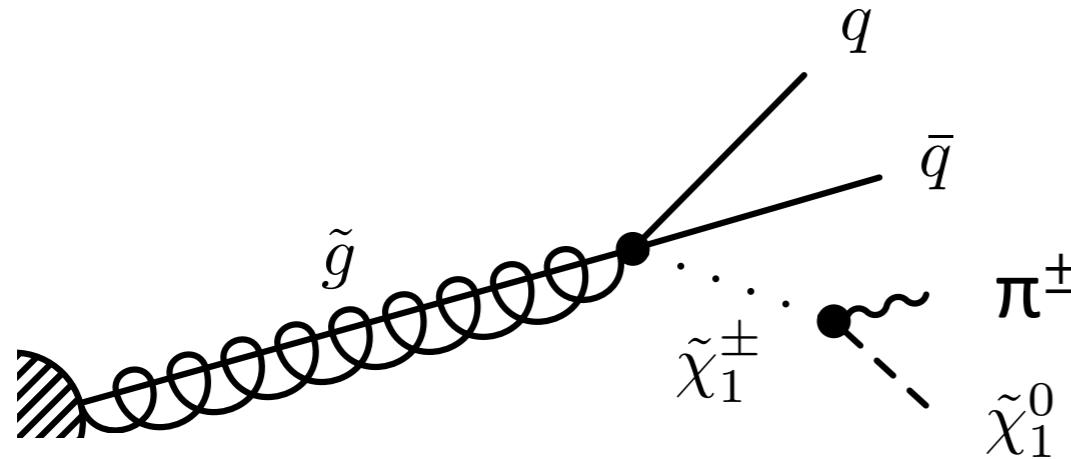
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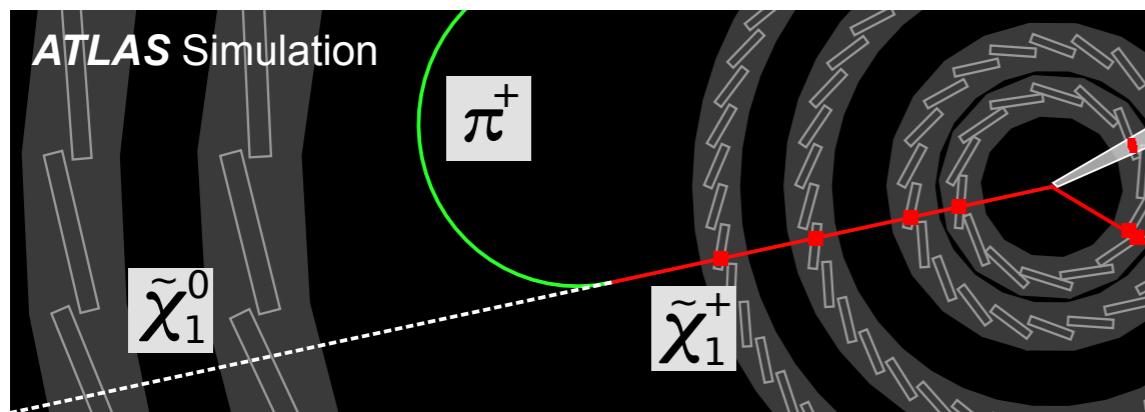
With extremely small mass splittings between SUSY particles, charginos can become long-lived!



Disappearing Tracks



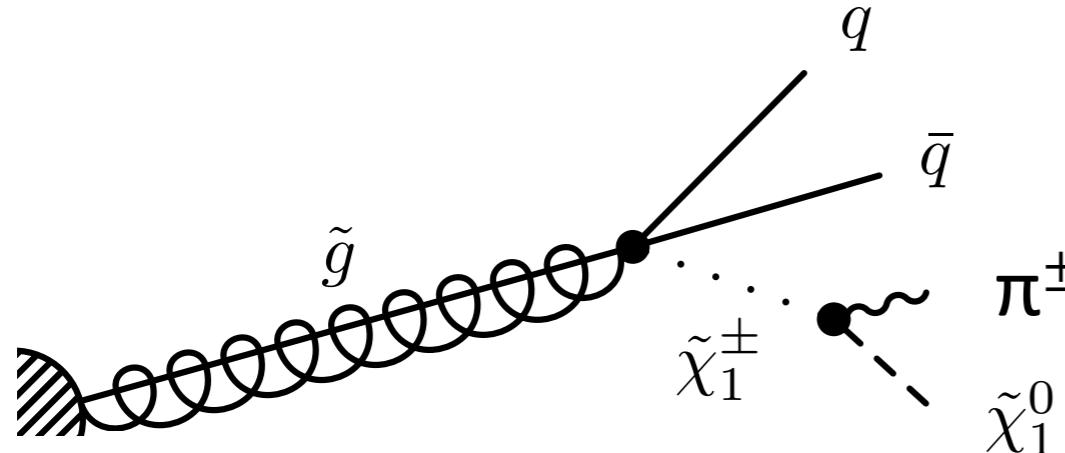
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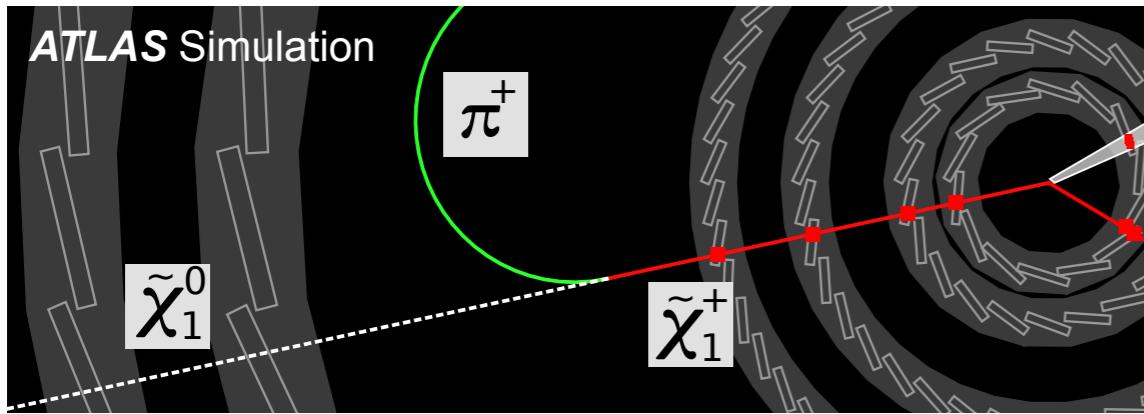
Can interact with the first few layers of the tracker, and then “disappear”



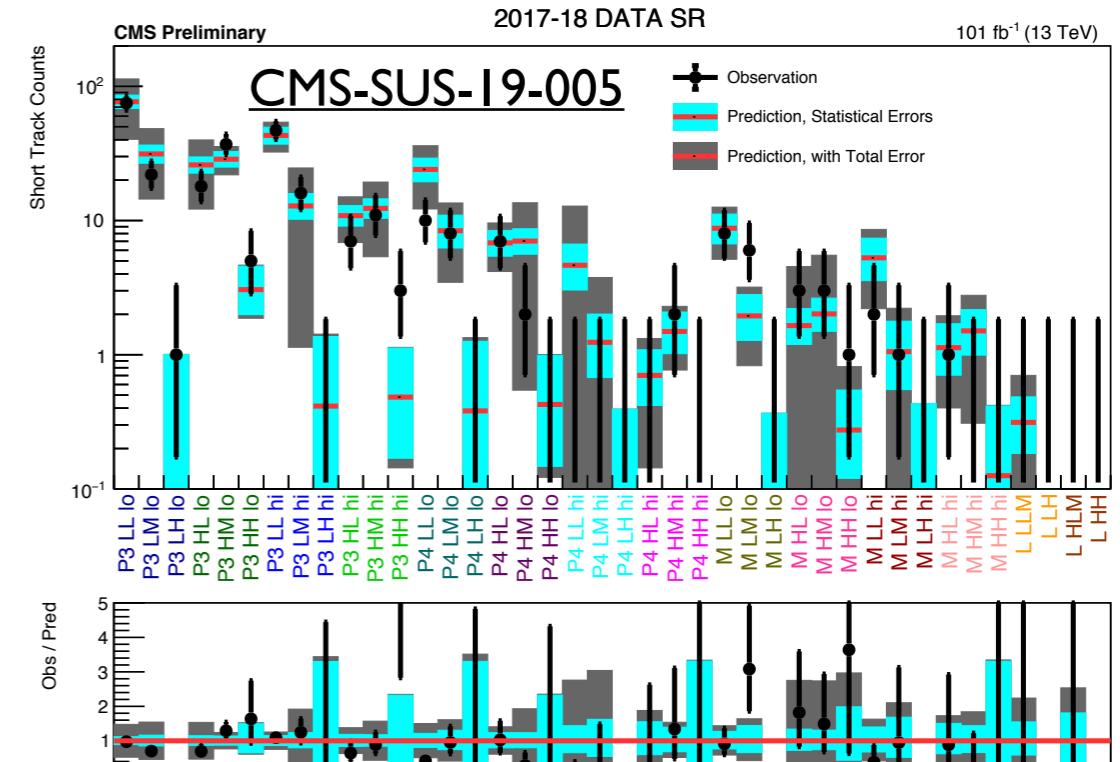
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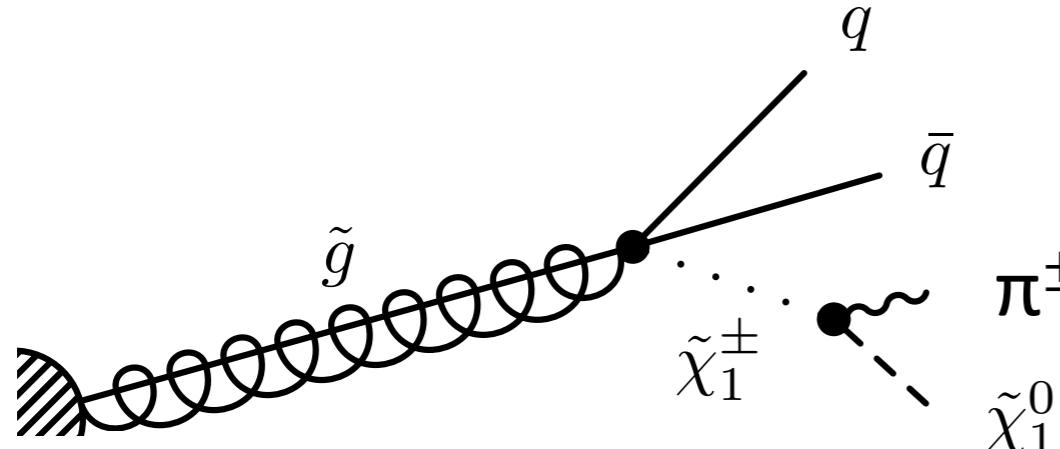


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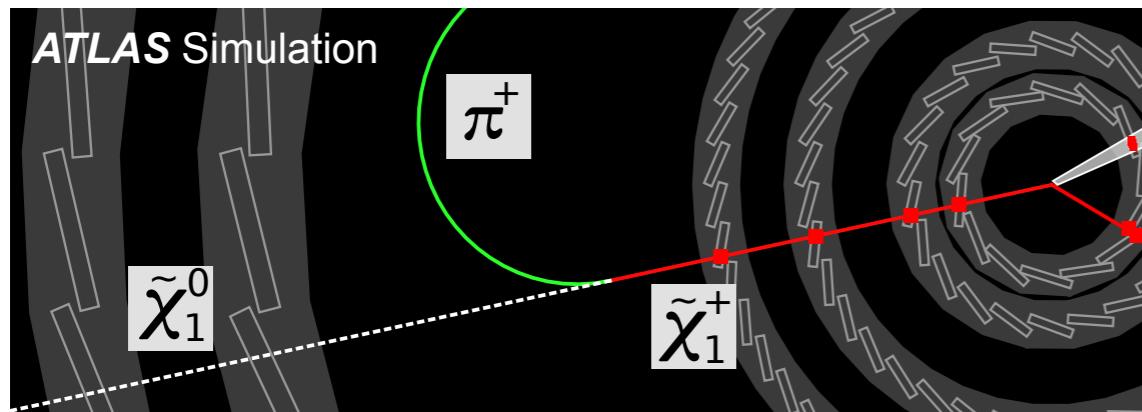




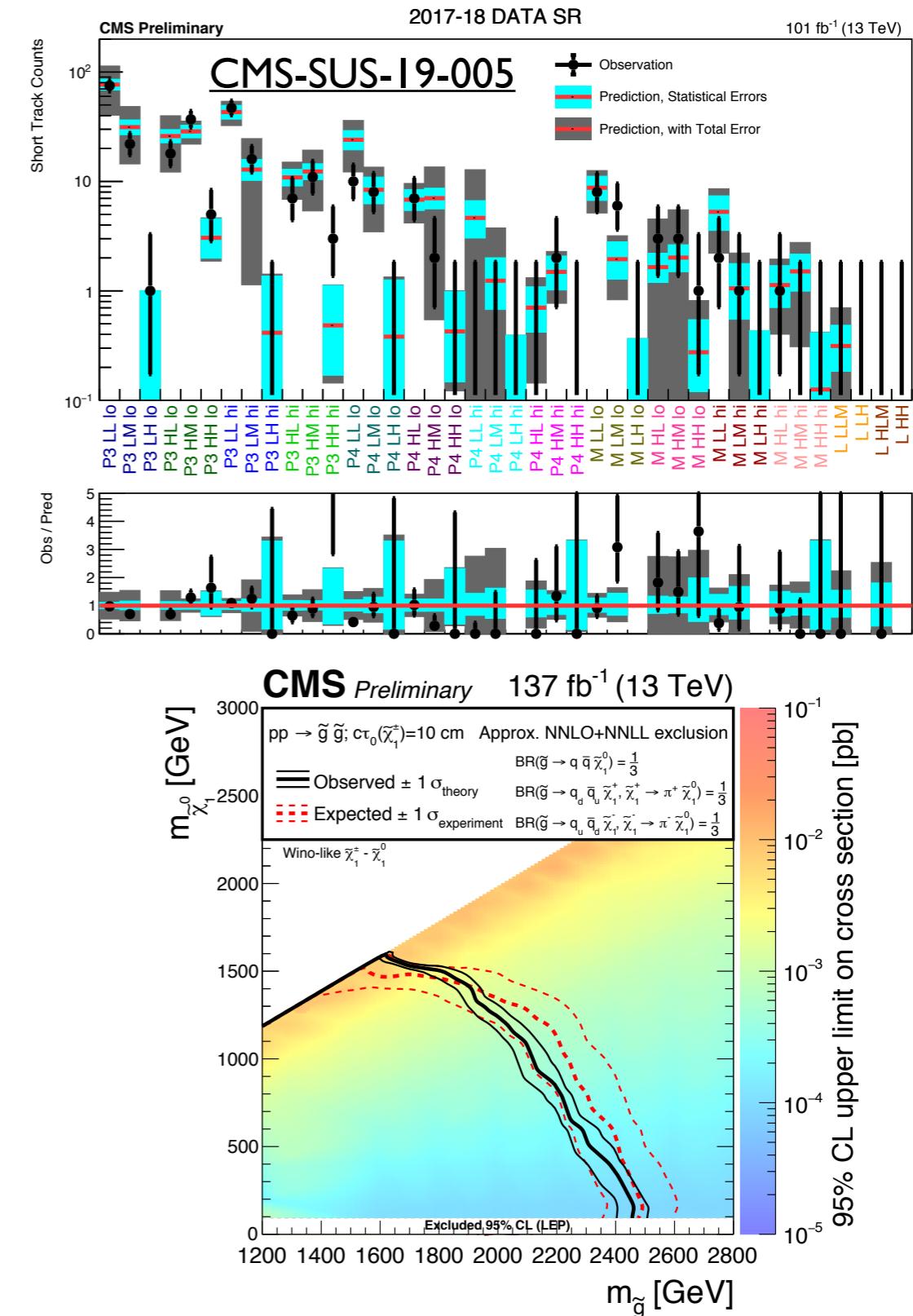
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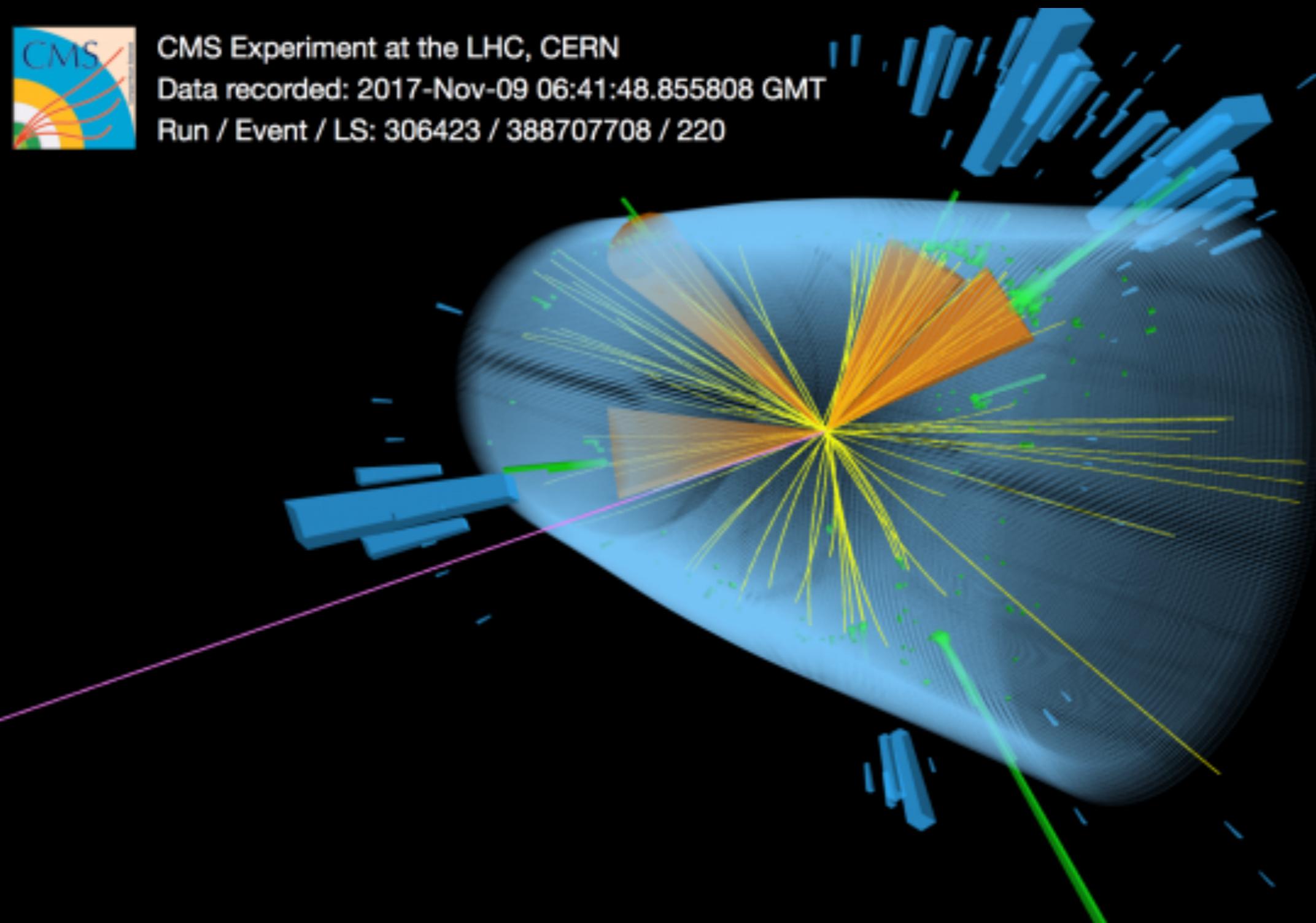


Can interact with the first few layers of the tracker, and then “disappear”



Can also extend reconstruction using additional information: **timing**

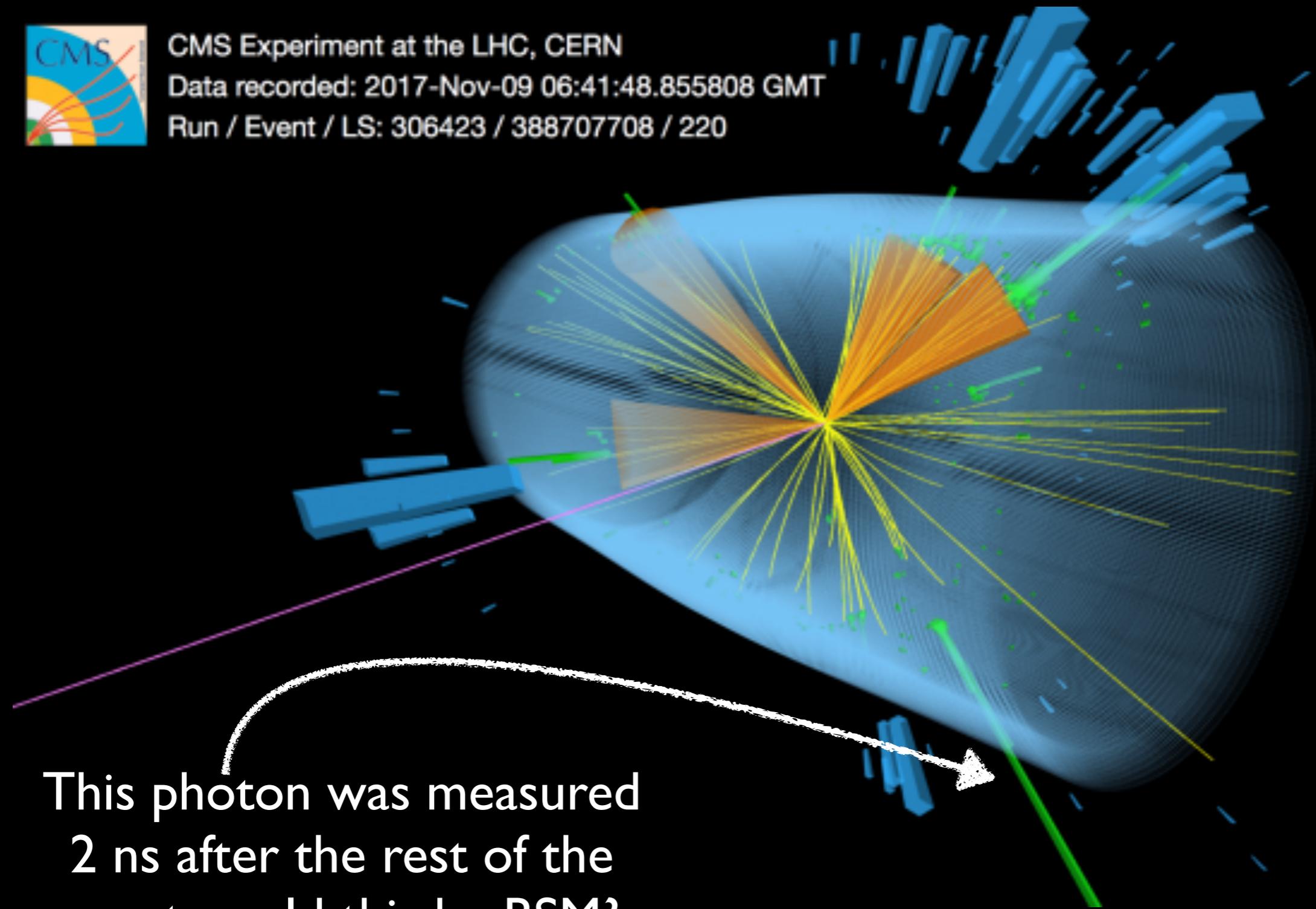
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CMS Experiment at the LHC, CERN
Data recorded: 2017-Nov-09 06:41:48.855808 GMT
Run / Event / LS: 306423 / 388707708 / 220



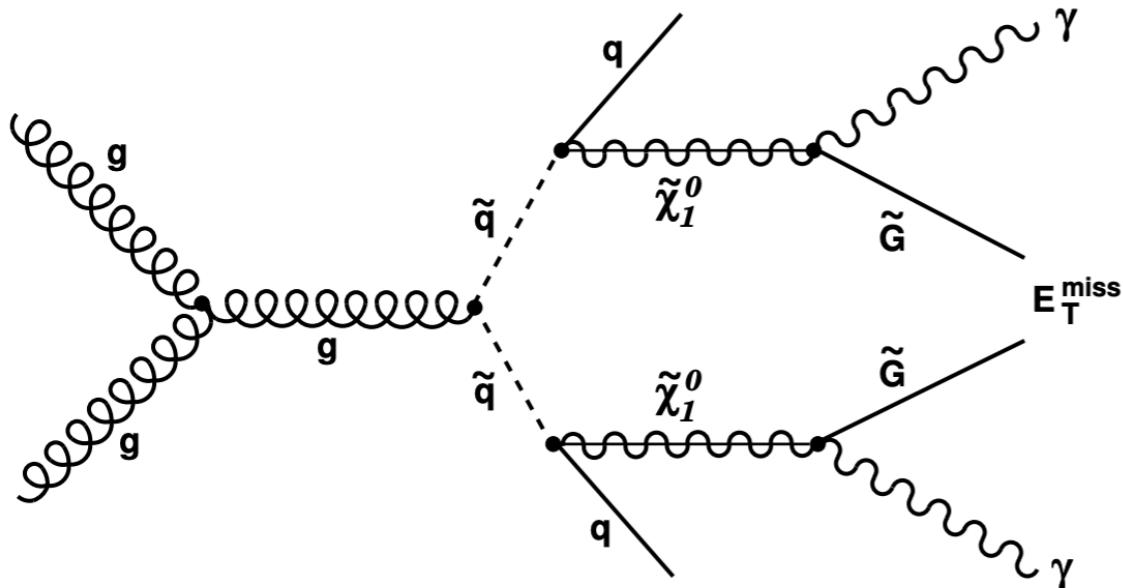
This photon was measured
2 ns after the rest of the
event: could this be BSM?

Delayed Photons





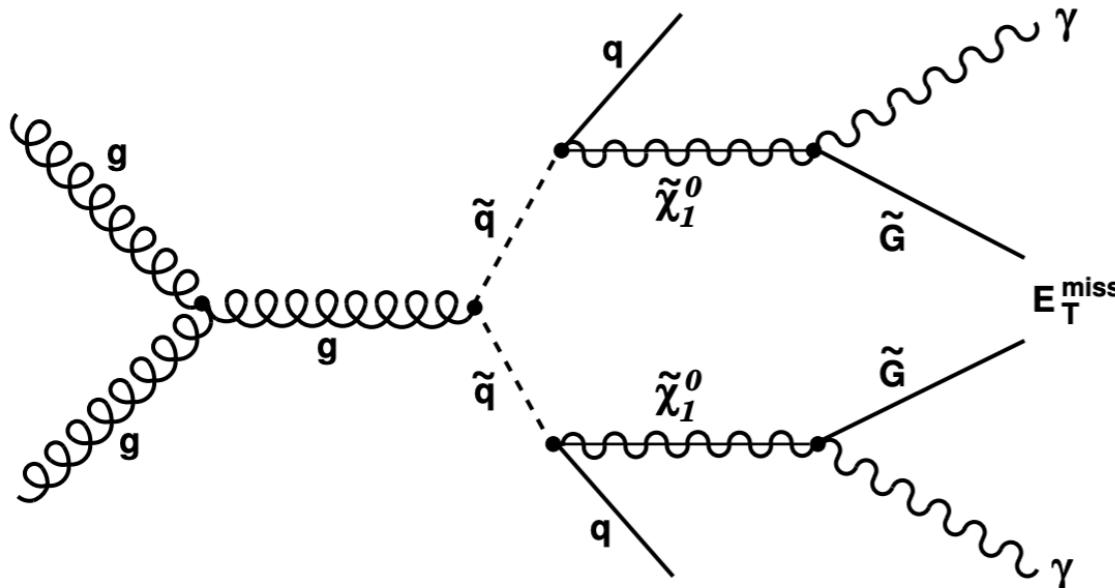
Delayed Photons



Many models (GMSB SUSY, etc.) have long-lived particles decaying to photons



Delayed Photons

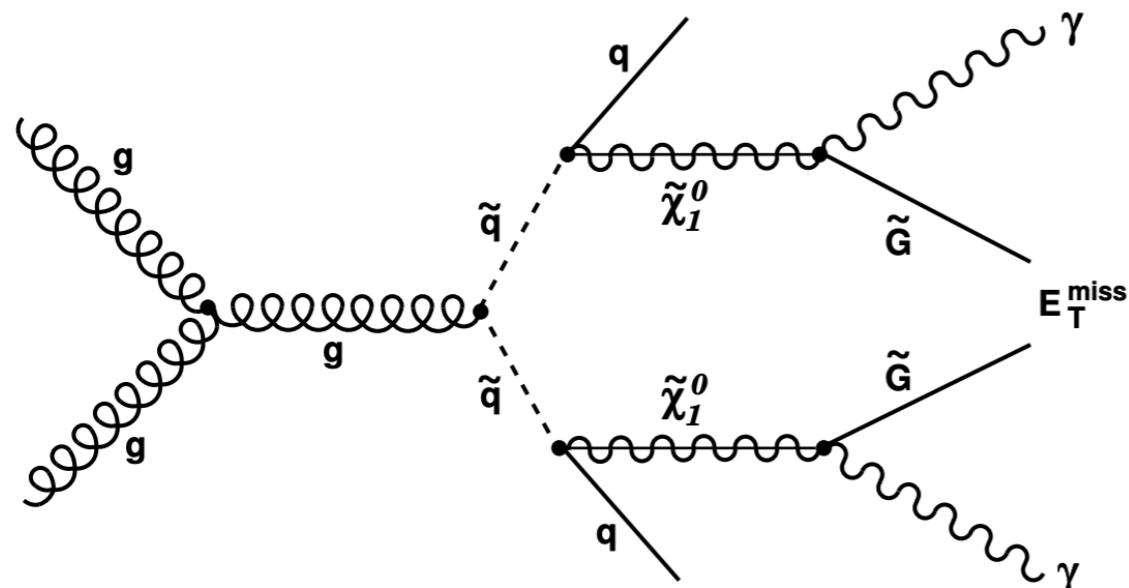


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These appear “late” compared to other energy in the event: use crystal timing in CMS to search

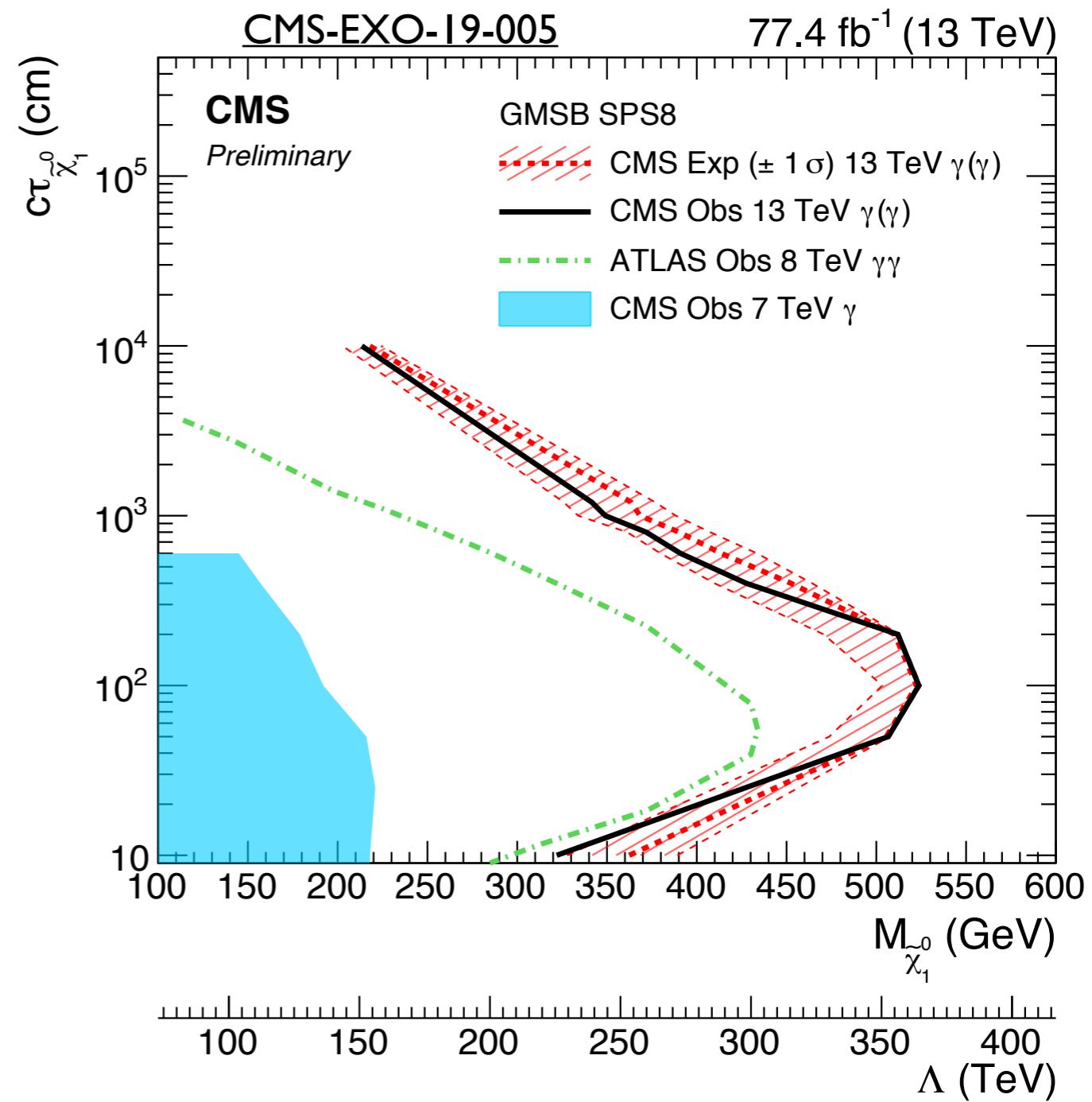


Delayed Photons



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Conclusions

Keep Looking!





Keep Looking!

**BSM is extremely well motivated
at the LHC!**



Keep Looking!

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Naturalness, dark matter, and
baryogenesis all need answers,
and the LHC could find them

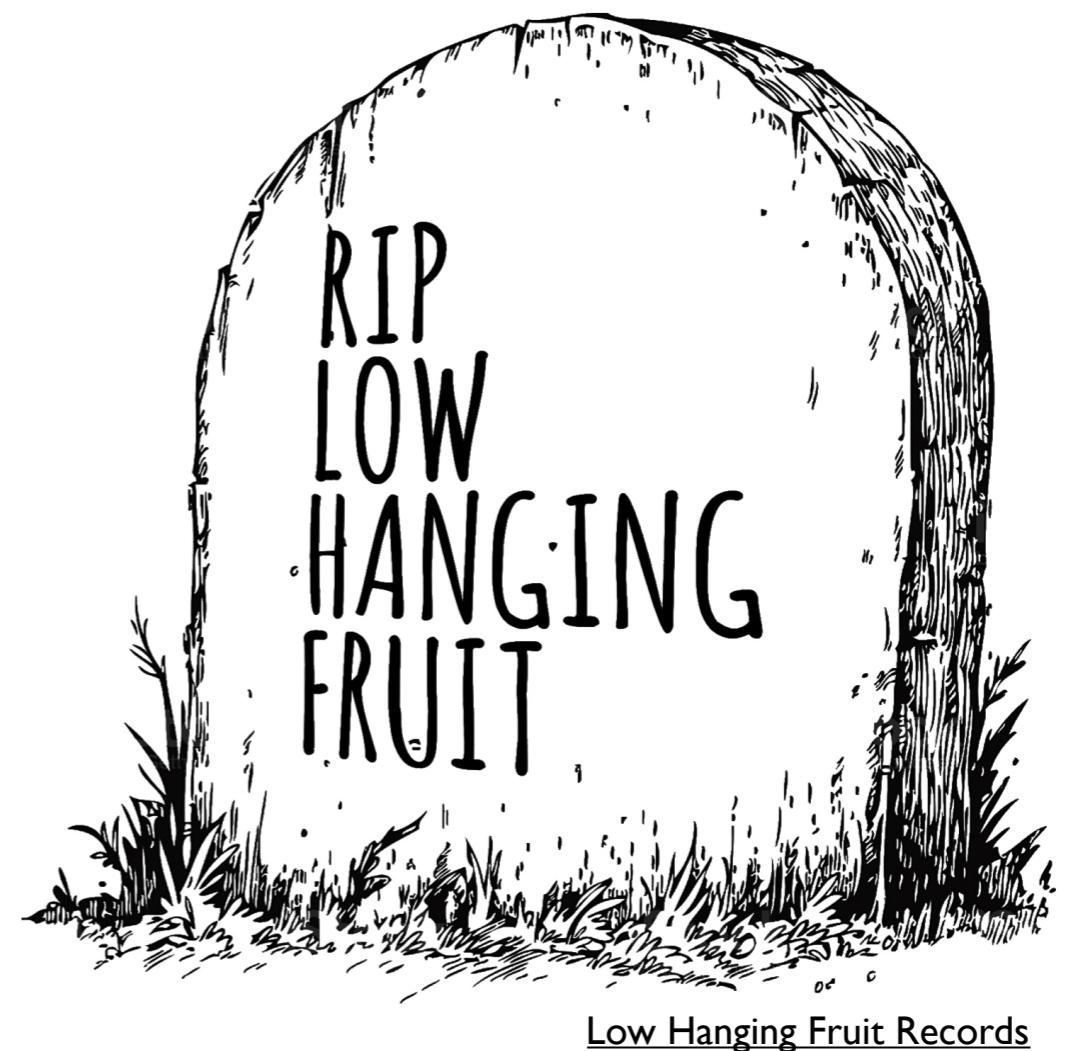


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Searches are moving
into a new era: the low-
hanging fruit has been
picked, and the
challenge is increasing



Low Hanging Fruit Records



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Naturalness, dark matter, and
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Searches are moving
into a new era: the low-
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picked, and the
challenge is increasing

But new tools, large datasets, and creative reconstruction
mean we can rise to the challenge!



Low Hanging Fruit Records

Thank you!

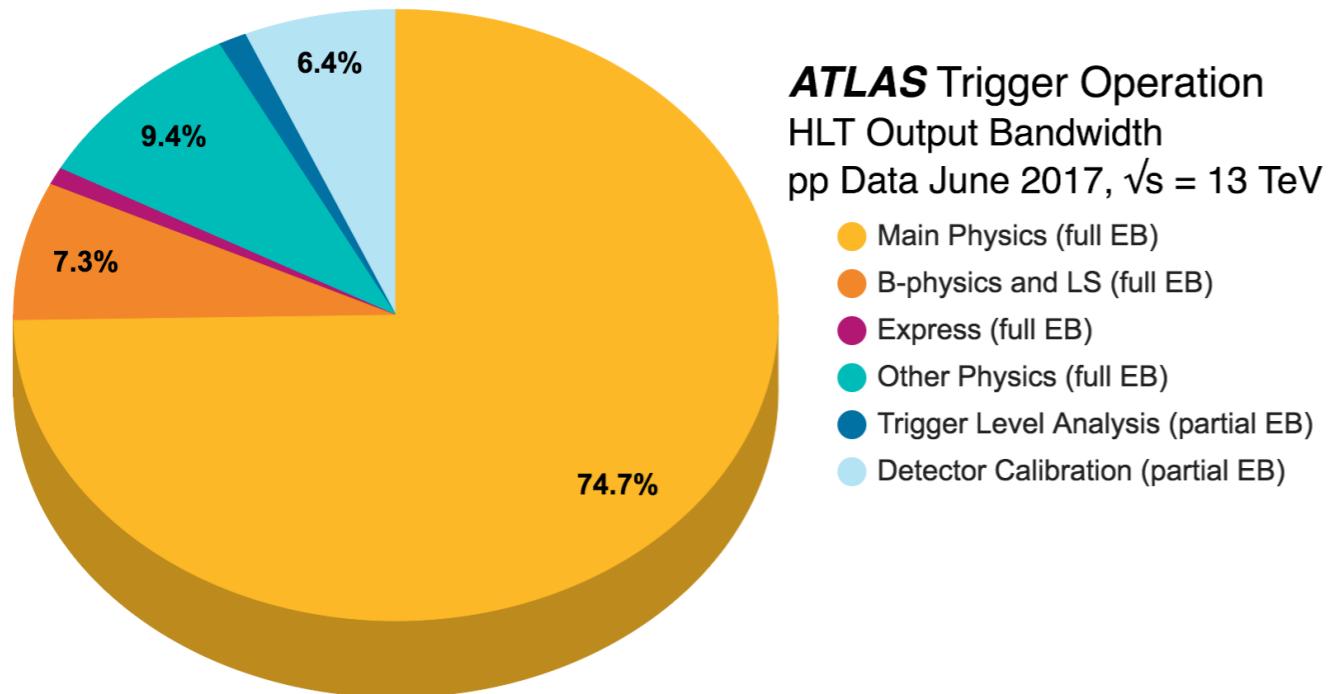
Backup

Data Scouting





Data Scouting

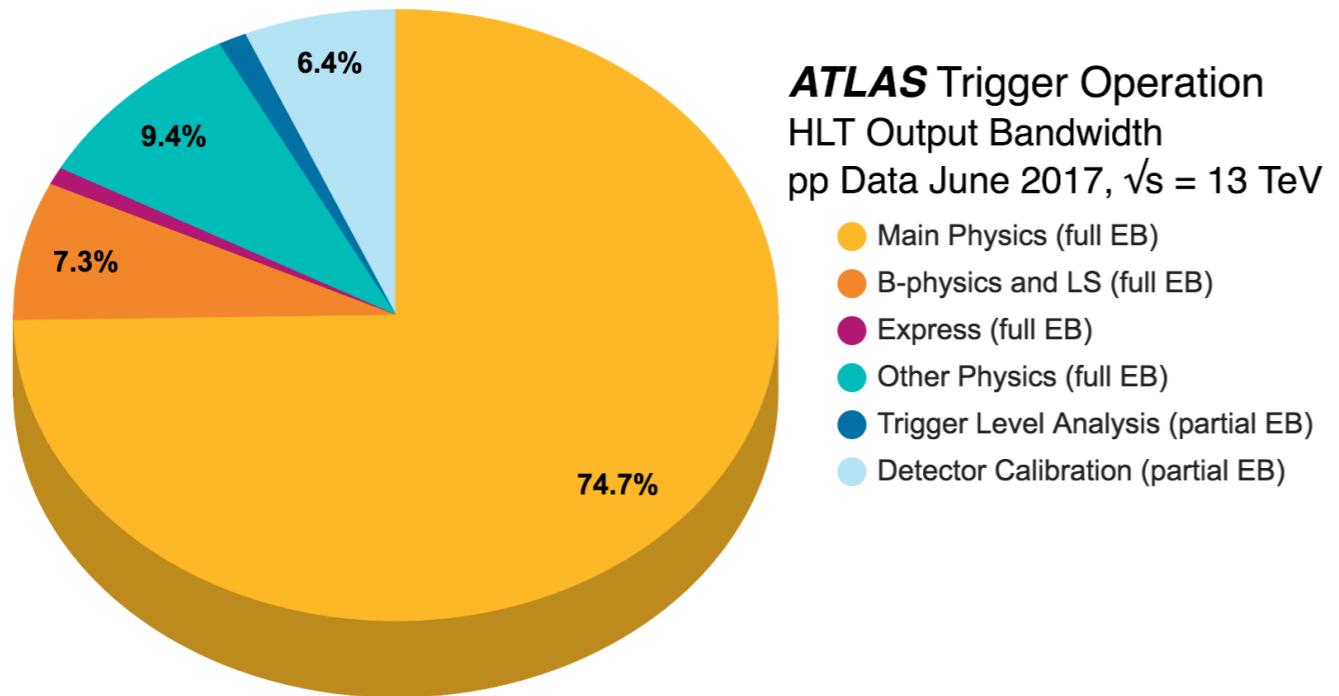


ATLAS Trigger Operation
HLT Output Bandwidth
pp Data June 2017, $\sqrt{s} = 13$ TeV

- Main Physics (full EB)
- B-physics and LS (full EB)
- Express (full EB)
- Other Physics (full EB)
- Trigger Level Analysis (partial EB)
- Detector Calibration (partial EB)



Data Scouting



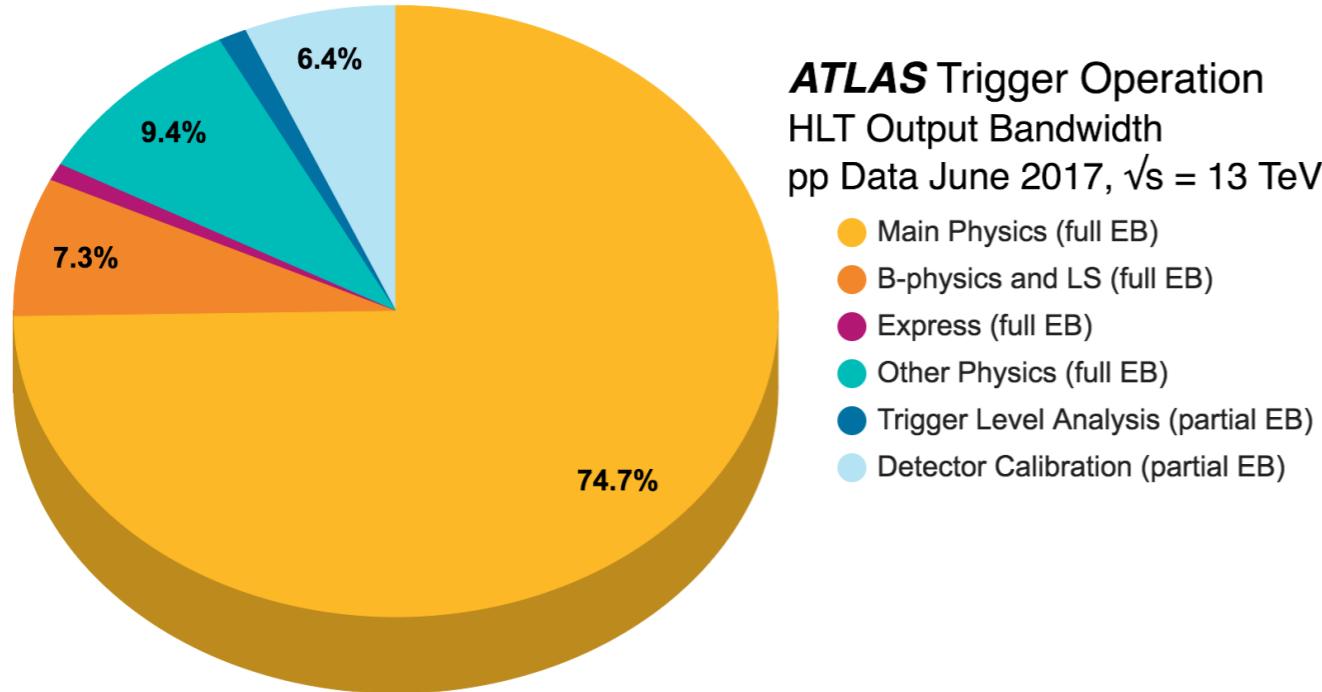
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Data Scouting



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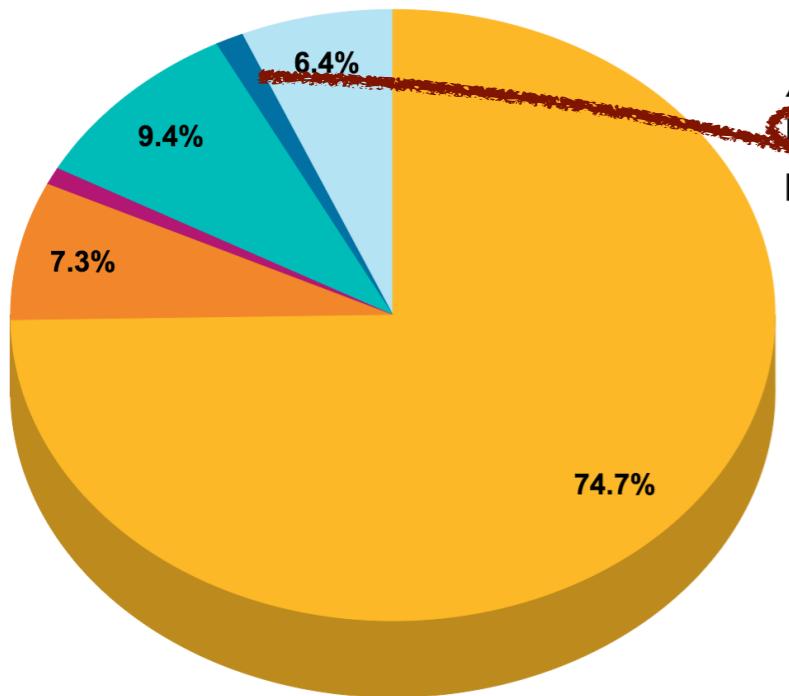
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Most of the time, record the “full event” when triggered

But for some analyses, we don’t need the full event:
just the jets can be enough to do physics!

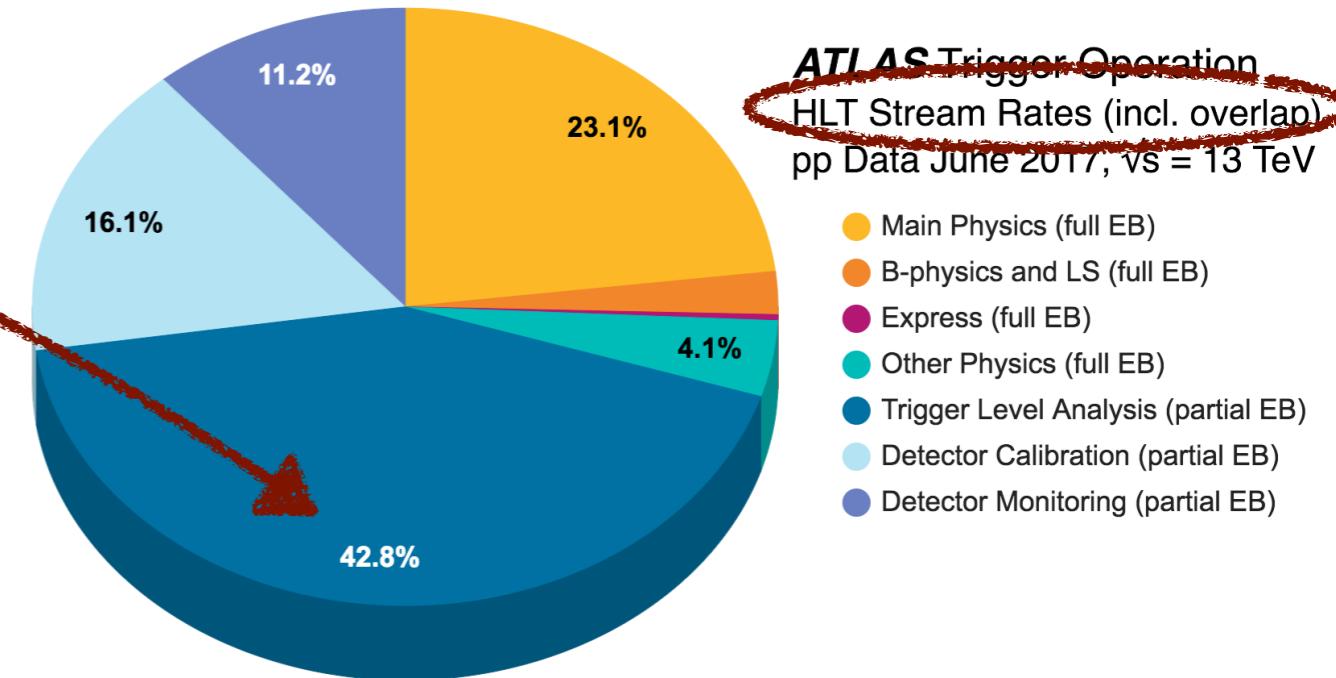


Data Scouting



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ATLAS Trigger Operation
HLT Stream Rates (incl. overlap)
pp Data June 2017, $\sqrt{s} = 13$ TeV

- Main Physics (full EB)
- B-physics and LS (full EB)
- Express (full EB)
- Other Physics (full EB)
- Trigger Level Analysis (partial EB)
- Detector Calibration (partial EB)
- Detector Monitoring (partial EB)

Most of the time, record the “full event” when triggered

But for some analyses, we don’t need the full event:
just the jets can be enough to do physics!

If you make the event small, then
you can record a lot more data!



Machine Learning for di-Higgs



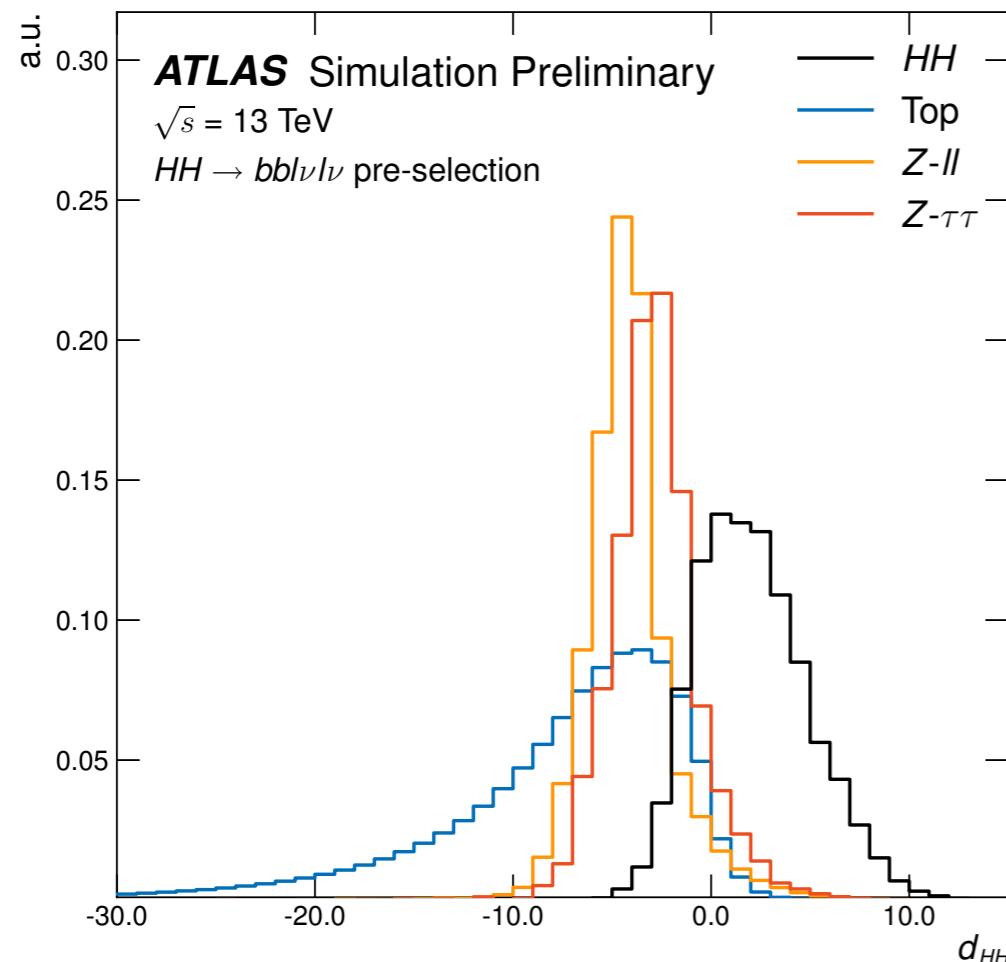
Machine Learning for di-Higgs

Rare signals are an excellent target for machine learning!



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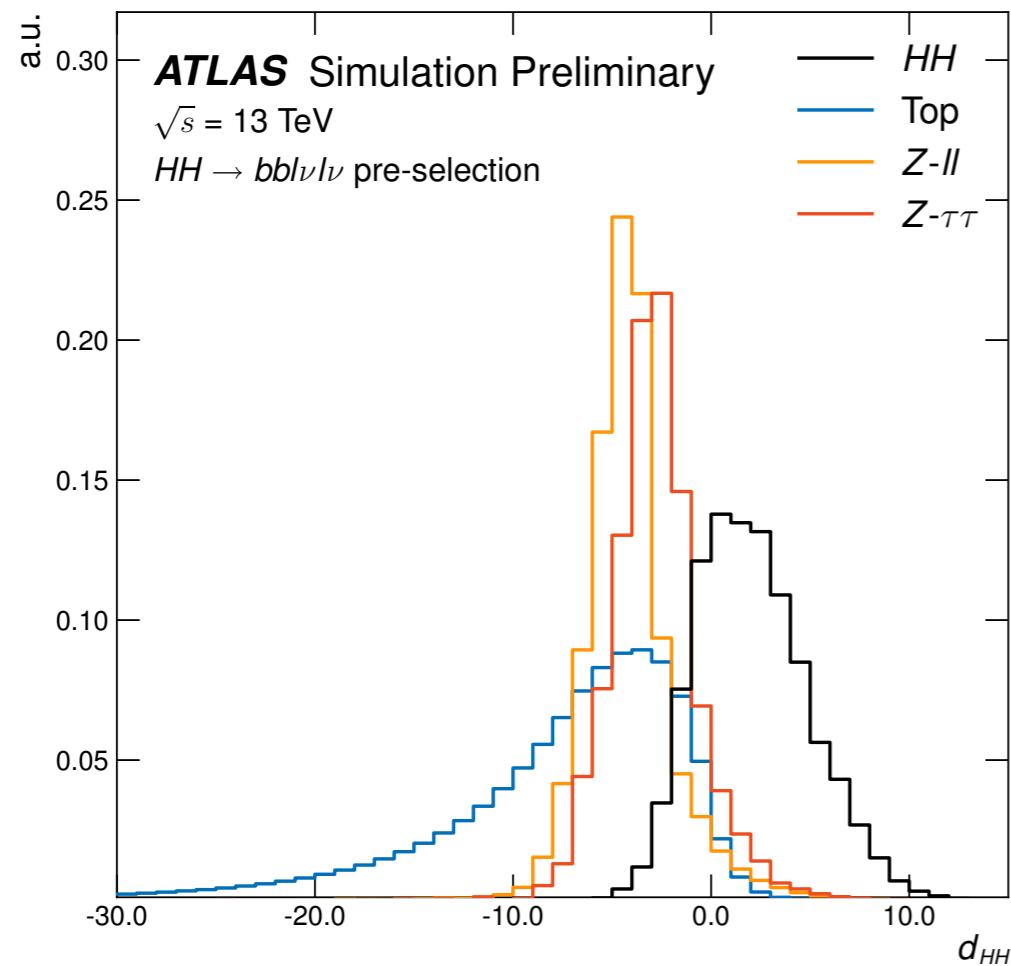


ATLAS uses a DNN to search for di-Higgs in $bbll$ final state

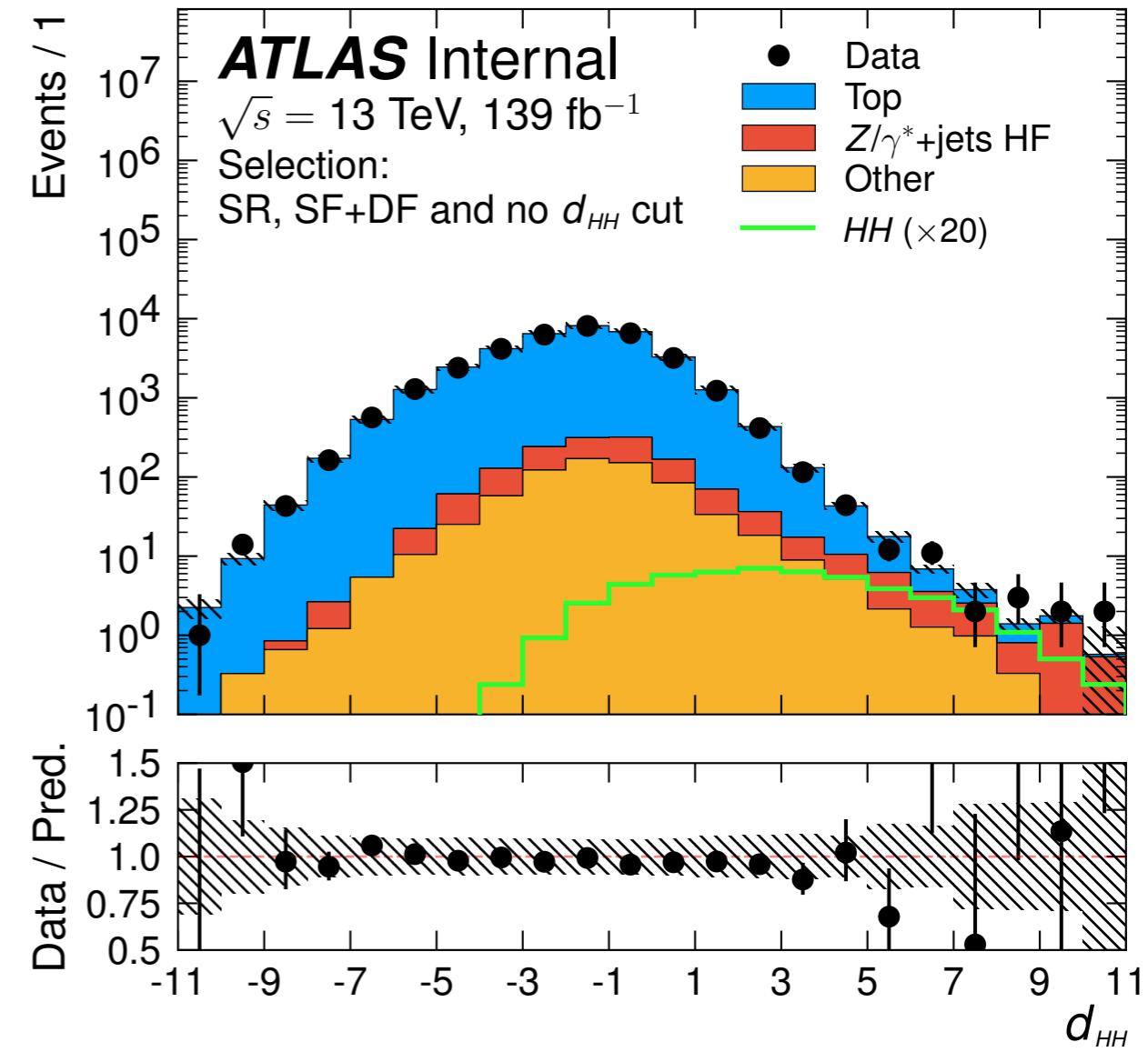


Machine Learning for di-Higgs

Rare signals are an excellent target for machine learning!



ATLAS uses a DNN to search for di-Higgs in $b\bar{b}ll$ final state



Machine learning can make this rarer channel competitive in the hunt for di-Higgs!

CMS Jet Timing



ATLAS WH results

