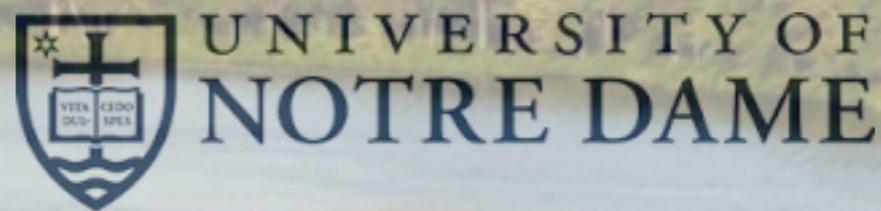


Theory Perspective on Particle Searches at the LHC

Adam Martin (amarti41@nd.edu)



DPF2019, Boston July 31st, 2019

Where do we stand:

Just a few years ago...

SUSY right around the corner! New W' , Z' , multiple Higgses!

Very non-standard Higgs, or even no Higgs!

Dark Matter @ LHC, extra dimensions!

ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: May 2019

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 139) \text{ fb}^{-1}$$

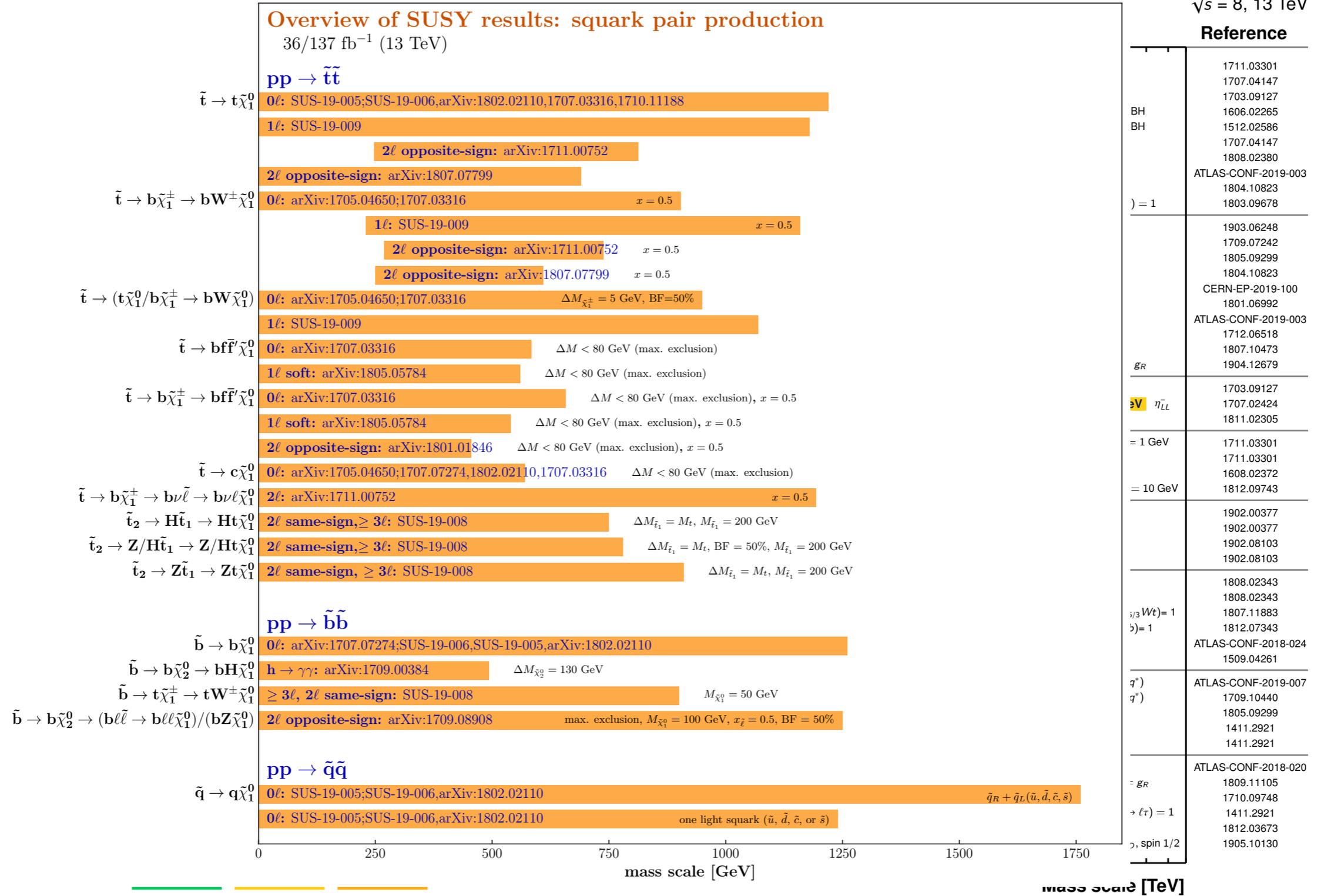
$$\sqrt{s} = 8, 13 \text{ TeV}$$

Model	ℓ, γ	Jets [†]	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit	Reference	
Extra dimensions	ADD $G_{KK} + g/q$	0 e, μ	1-4 j	Yes	36.1	M_D 7.7 TeV	$n = 2$ 1711.03301
	ADD non-resonant $\gamma\gamma$	2 γ	-	-	36.7	M_S 8.6 TeV	$n = 3$ HLZ NLO 1707.04147
	ADD QBH	-	2 j	-	37.0	M_{th} 8.9 TeV	$n = 6$ 1703.09127
	ADD BH high $\sum p_T$	$\geq 1 e, \mu$	$\geq 2 j$	-	3.2	M_{th} 8.2 TeV	$n = 6, M_D = 3 \text{ TeV}$, rot BH 1606.02265
	ADD BH multijet	-	$\geq 3 j$	-	3.6	M_{th} 9.55 TeV	$n = 6, M_D = 3 \text{ TeV}$, rot BH 1512.02586
	RS1 $G_{KK} \rightarrow \gamma\gamma$	2 γ	-	-	36.7	G_{KK} mass 4.1 TeV	$k/\overline{M}_{Pl} = 0.1$ 1707.04147
	Bulk RS $G_{KK} \rightarrow WW/ZZ$	multi-channel	-	-	36.1	G_{KK} mass 2.3 TeV	$k/\overline{M}_{Pl} = 1.0$ 1808.02380
	Bulk RS $G_{KK} \rightarrow WW \rightarrow qq\bar{q}\bar{q}$	0 e, μ	2 J	-	139	G_{KK} mass 1.6 TeV	$k/\overline{M}_{Pl} = 1.0$ ATLAS-CONF-2019-003
	Bulk RS $g_{KK} \rightarrow tt$	1 e, μ	$\geq 1 b, \geq 1J/2j$	Yes	36.1	g_{KK} mass 3.8 TeV	$\Gamma/m = 15\%$ 1804.10823
	2UED / RPP	1 e, μ	$\geq 2 b, \geq 3 j$	Yes	36.1	KK mass 1.8 TeV	Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow tt) = 1$ 1803.09678
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	2 e, μ	-	-	139	Z' mass 5.1 TeV	$\Gamma/m = 1\%$ 1903.06248
	SSM $Z' \rightarrow \tau\tau$	2 τ	-	-	36.1	Z' mass 2.42 TeV	1709.07242
	Leptophobic $Z' \rightarrow b\bar{b}$	-	2 b	-	36.1	Z' mass 2.1 TeV	1805.09299
	Leptophobic $Z' \rightarrow t\bar{t}$	1 e, μ	$\geq 1 b, \geq 1J/2j$	Yes	36.1	Z' mass 3.0 TeV	1804.10823
	SSM $W' \rightarrow \ell\nu$	1 e, μ	-	Yes	139	W' mass 6.0 TeV	CERN-EP-2019-100
	SSM $W' \rightarrow \tau\nu$	1 τ	-	Yes	36.1	W' mass 3.7 TeV	1801.06992
	HVT $V' \rightarrow WZ \rightarrow qq\bar{q}\bar{q}$ model B	0 e, μ	2 J	-	139	V' mass 3.6 TeV	$g_V = 3$ ATLAS-CONF-2019-003
	HVT $V' \rightarrow WH/ZH$ model B	multi-channel	-	-	36.1	V' mass 2.93 TeV	$g_V = 3$ 1712.06518
	LRSM $W_R \rightarrow tb$	multi-channel	-	-	36.1	W_R mass 3.25 TeV	1807.10473
	LRSM $W_R \rightarrow \mu N_R$	2 μ	1 J	-	80	W_R mass 5.0 TeV	$m(N_R) = 0.5 \text{ TeV}, g_L = g_R$ 1904.12679
CI	CI $qq\bar{q}\bar{q}$	-	2 j	-	37.0	Λ 21.8 TeV	η_{LL}^- 1703.09127
	CI $\ell\ell\bar{q}\bar{q}$	2 e, μ	-	-	36.1	Λ 40.0 TeV	η_{LL}^- 1707.02424
	CI $t\bar{t}\bar{t}\bar{t}$	$\geq 1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Λ 2.57 TeV	$ C_{4t} = 4\pi$ 1811.02305
DM	Axial-vector mediator (Dirac DM)	0 e, μ	1-4 j	Yes	36.1	m_{med} 1.55 TeV	$g_q = 0.25, g_\chi = 1.0, m(\chi) = 1 \text{ GeV}$ 1711.03301
	Colored scalar mediator (Dirac DM)	0 e, μ	1-4 j	Yes	36.1	m_{med} 1.67 TeV	$g = 1.0, m(\chi) = 1 \text{ GeV}$ 1711.03301
	$VV\chi\chi$ EFT (Dirac DM)	0 e, μ	1 J, $\leq 1 j$	Yes	3.2	M_* 700 GeV	$m(\chi) < 150 \text{ GeV}$ 1608.02372
	Scalar reson. $\phi \rightarrow t\chi$ (Dirac DM)	0-1 e, μ	1 b, 0-1 J	Yes	36.1	m_ϕ 3.4 TeV	$y = 0.4, \lambda = 0.2, m(\chi) = 10 \text{ GeV}$ 1812.09743
LQ	Scalar LQ 1 st gen	1,2 e	$\geq 2 j$	Yes	36.1	LQ mass 1.4 TeV	$\beta = 1$ 1902.00377
	Scalar LQ 2 nd gen	1,2 μ	$\geq 2 j$	Yes	36.1	LQ mass 1.56 TeV	$\beta = 1$ 1902.00377
	Scalar LQ 3 rd gen	2 τ	2 b	-	36.1	LQ_3^u mass 1.03 TeV	$\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1$ 1902.08103
	Scalar LQ 3 rd gen	0-1 e, μ	2 b	Yes	36.1	LQ_3^d mass 970 GeV	$\mathcal{B}(LQ_3^d \rightarrow t\tau) = 0$ 1902.08103
Heavy quarks	VLQ $TT \rightarrow Ht/Zt/Wb + X$	multi-channel	-	-	36.1	T mass 1.37 TeV	SU(2) doublet 1808.02343
	VLQ $BB \rightarrow Wt/Zb + X$	multi-channel	-	-	36.1	B mass 1.34 TeV	SU(2) doublet 1808.02343
	VLQ $T_{5/3} T_{5/3} T_{5/3} \rightarrow Wt + X$	2(SS) $\geq 3 e, \mu \geq 1 b, \geq 1 j$	Yes	36.1	$T_{5/3}$ mass 1.64 TeV	$\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$ 1807.11883	
	VLQ $Y \rightarrow Wb + X$	1 e, μ	$\geq 1 b, \geq 1 j$	Yes	36.1	Y mass 1.85 TeV	$\mathcal{B}(Y \rightarrow Wb) = 1, c_R(Wb) = 1$ 1812.07343
	VLQ $B \rightarrow Hb + X$	0 $e, \mu, 2 \gamma$	$\geq 1 b, \geq 1 j$	Yes	79.8	B mass 1.21 TeV	$\kappa_B = 0.5$ ATLAS-CONF-2018-024
VLQ $QQ \rightarrow WqVq$	1 e, μ	$\geq 4 j$	Yes	20.3	Q mass 690 GeV	1509.04261	
Excited fermions	Excited quark $q^* \rightarrow qg$	-	2 j	-	139	q^* mass 6.7 TeV	only u^* and d^* , $\Lambda = m(q^*)$ ATLAS-CONF-2019-007
	Excited quark $q^* \rightarrow q\gamma$	1 γ	1 j	-	36.7	q^* mass 5.3 TeV	only u^* and d^* , $\Lambda = m(q^*)$ 1709.10440
	Excited quark $b^* \rightarrow b\bar{g}$	-	1 b, 1 j	-	36.1	b^* mass 2.6 TeV	1805.09299
	Excited lepton ℓ^*	3 e, μ	-	-	20.3	ℓ^* mass 3.0 TeV	$\Lambda = 3.0 \text{ TeV}$ 1411.2921
	Excited lepton ν^*	3 e, μ, τ	-	-	20.3	ν^* mass 1.6 TeV	$\Lambda = 1.6 \text{ TeV}$ 1411.2921
Other	Type III Seesaw	1 e, μ	$\geq 2 j$	Yes	79.8	N^0 mass 560 GeV	$m(W_R) = 4.1 \text{ TeV}, g_L = g_R$ ATLAS-CONF-2018-020
	LRSM Majorana ν	2 μ	2 j	-	36.1	N_R mass 3.2 TeV	1809.11105
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$	2,3,4 e, μ (SS)	-	-	36.1	$H^{\pm\pm}$ mass 870 GeV	DY production 1710.09748
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$	3 e, μ, τ	-	-	20.3	$H^{\pm\pm}$ mass 400 GeV	DY production, $\mathcal{B}(H_L^{\pm\pm} \rightarrow \ell\tau) = 1$ 1411.2921
	Multi-charged particles	-	-	-	36.1	multi-charged particle mass 1.22 TeV	DY production, $ q = 5e$ 1812.03673
	Magnetic monopoles	-	-	-	34.4	monopole mass 2.37 TeV	DY production, $ g = 1g_D$, spin 1/2 1905.10130

$\sqrt{s} = 8 \text{ TeV}$ $\sqrt{s} = 13 \text{ TeV}$ partial data $\sqrt{s} = 13 \text{ TeV}$ full data

10⁻¹ 1 10 Mass scale [TeV]

Remarkable success of SM, through an incredible variety of tests/searches



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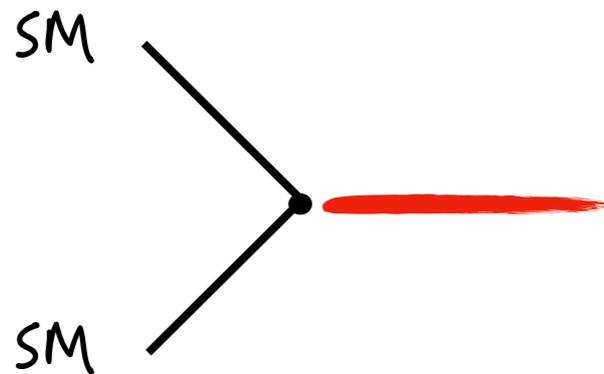
Still possible that one of these scenarios is lying just out of reach & will be dug up with more data. However, mass limits will increase very slowly in next phase of LHC

- **LHC as a precision machine, look for indirect effects of new physics**
- **leave no stone unturned**

Go heavy: Effective Field Theory (EFT) Approach

New physics may be too heavy to make on-shell at the LHC

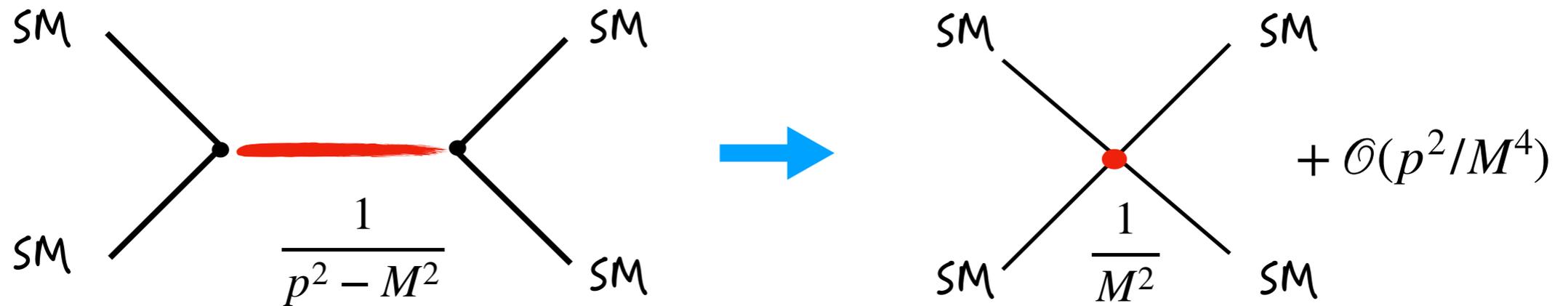
instead, look for its indirect effects on SM processes



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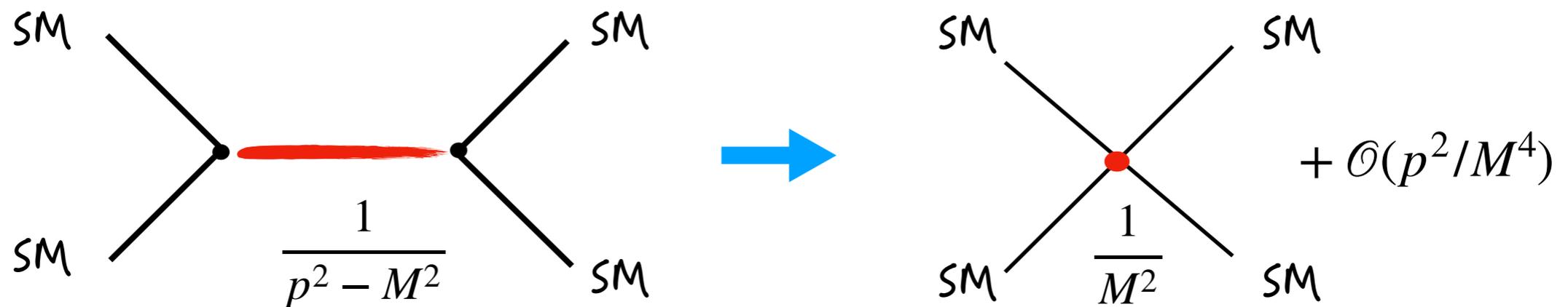
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Go heavy: Effective Field Theory (EFT) Approach

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Effective vertices with only SM particles, expansion in derivatives = SMEFT

$$\mathcal{L} = \mathcal{L}_{SM} + \sum \frac{c_i}{\Lambda^2} \mathcal{O}_i^{d=6} + \sum \frac{c_i}{\Lambda^4} \mathcal{O}_i^{d=8} + \dots$$

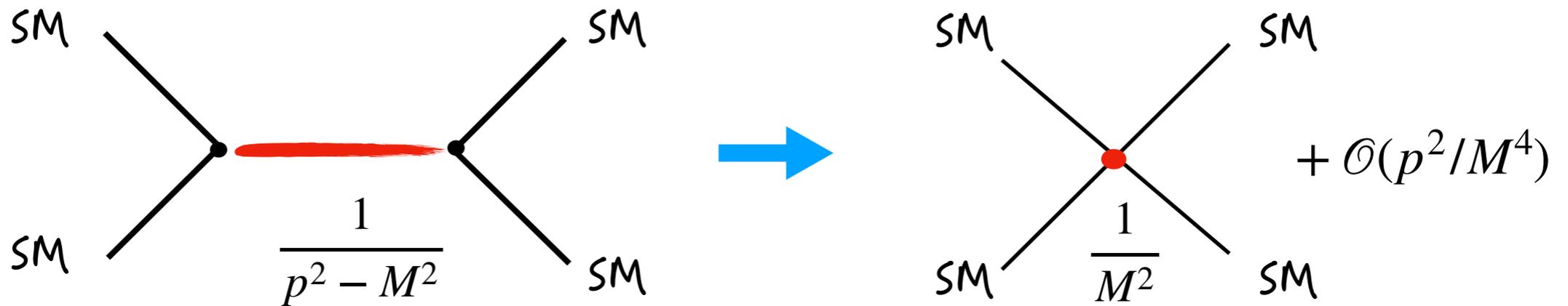
BSM SM particles

Program: study SM final states, look for deviations. Clean processes usually the most sensitive: $h \rightarrow \gamma\gamma$, $pp \rightarrow 4\ell$, dibosons...

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instead, look for its indirect effects on SM processes



Including SMEFT terms:

$$|A|^2 = |A_{SM}|^2 + \frac{2E^2}{\Lambda^2} \text{Re}(A_{SM}^* A_6) + \frac{E^4}{\Lambda^4} |A_6|^2 + \frac{2E^4}{\Lambda^4} \text{Re}(A_{SM}^* A_8) \dots$$

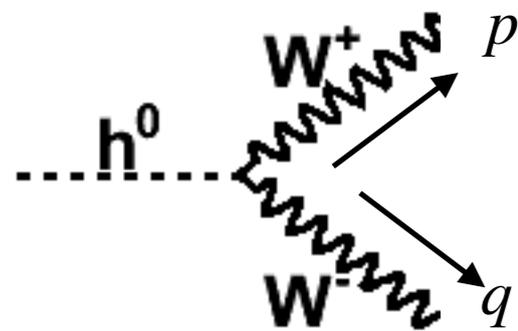
↑ interference piece
 ↑ (new physics)²

Go heavy: SMEFT

(* also CP structure)

Higher dimensional operators can have different Lorentz/derivative structure than SM!

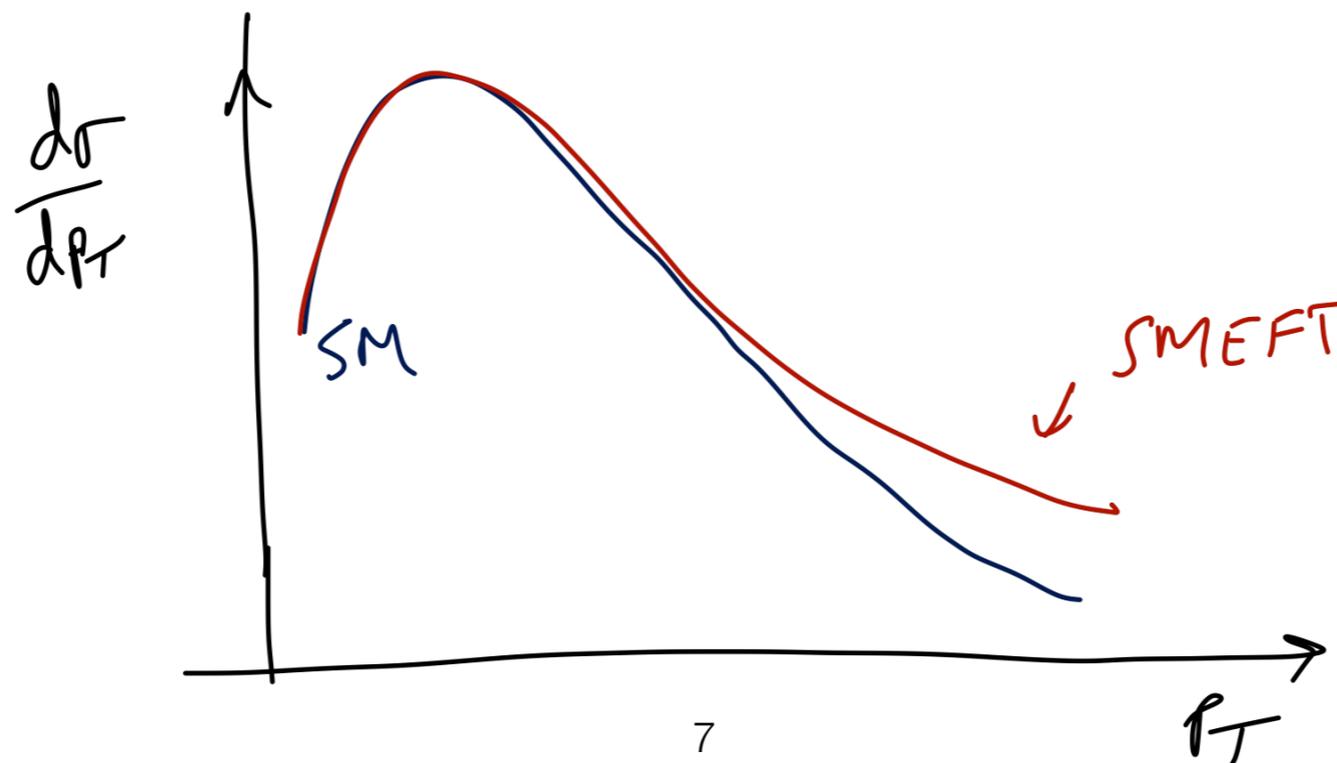
$$\text{ex.) } L_{EFT} \supset \frac{g_2^2 c_{HW}}{\Lambda^2} H^\dagger H W_{\mu\nu}^a W^{a,\mu\nu} + \frac{c_{H\Box}}{\Lambda^2} H^\dagger H \Box (H^\dagger H)$$



new momentum dependence

$$A_{hW(p)W(q)} \propto \frac{2m_W^2}{v} \left(g_{\mu\nu} \left(\overset{\text{SM term}}{1} + \frac{c_{H\Box} v^2}{4\Lambda^2} \right) + \underbrace{(p \cdot q g_{\mu\nu} - q_\mu p_\nu)}_{\text{new momentum dependence}} \frac{8c_{HW}}{\Lambda^2} \right)$$

For more handles to tease out new physics effects, look to differential distributions:
Higgs distributions are a natural target



Go heavy: SMEFT

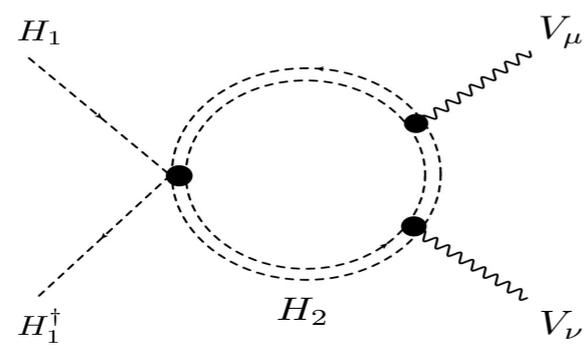
Well defined mapping from UV theory \rightarrow EFT \rightarrow comparison with data

EFT

$$\begin{aligned}
 c_{3G} &= \frac{g_s^2}{(4\pi)^2} \frac{1}{20} \\
 c_{3W} &= \frac{g^2}{(4\pi)^2} \frac{1}{20} \\
 c_{2G} &= \frac{g_s^2}{(4\pi)^2} \frac{1}{20} \\
 c_{2W} &= \frac{g^2}{(4\pi)^2} \frac{1}{20} \\
 c_{2B} &= \frac{g'^2}{(4\pi)^2} \frac{1}{20}
 \end{aligned}$$

$$\begin{aligned}
 c_H &= \frac{h_t^4}{(4\pi)^2} \frac{3}{4} \left[\left(1 + \frac{1}{3} \frac{g'^2 c_{2\beta}}{h_t^2} + \frac{1}{12} \frac{g'^4 c_{2\beta}^2}{h_t^4} \right) - \frac{7}{6} \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 + \frac{1}{14} \frac{(g^2 + 2g'^2) c_{2\beta}}{h_t^2} \right) + \frac{7}{30} \frac{X_t^4}{m_{\tilde{t}}^4} \right] \\
 c_T &= \frac{h_t^4}{(4\pi)^2} \frac{1}{4} \left[\left(1 + \frac{1}{2} \frac{g^2 c_{2\beta}}{h_t^2} \right)^2 - \frac{1}{2} \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 + \frac{1}{2} \frac{g^2 c_{2\beta}}{h_t^2} \right) + \frac{1}{10} \frac{X_t^4}{m_{\tilde{t}}^4} \right] \\
 c_R &= \frac{h_t^4}{(4\pi)^2} \frac{1}{2} \left[\left(1 + \frac{1}{2} \frac{g^2 c_{2\beta}}{h_t^2} \right)^2 - \frac{3}{2} \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 + \frac{1}{12} \frac{(3g^2 + g'^2) c_{2\beta}}{h_t^2} \right) + \frac{3}{10} \frac{X_t^4}{m_{\tilde{t}}^4} \right] \\
 c_D &= \frac{h_t^2}{(4\pi)^2} \frac{1}{20} \frac{X_t^2}{m_{\tilde{t}}^2}
 \end{aligned}$$

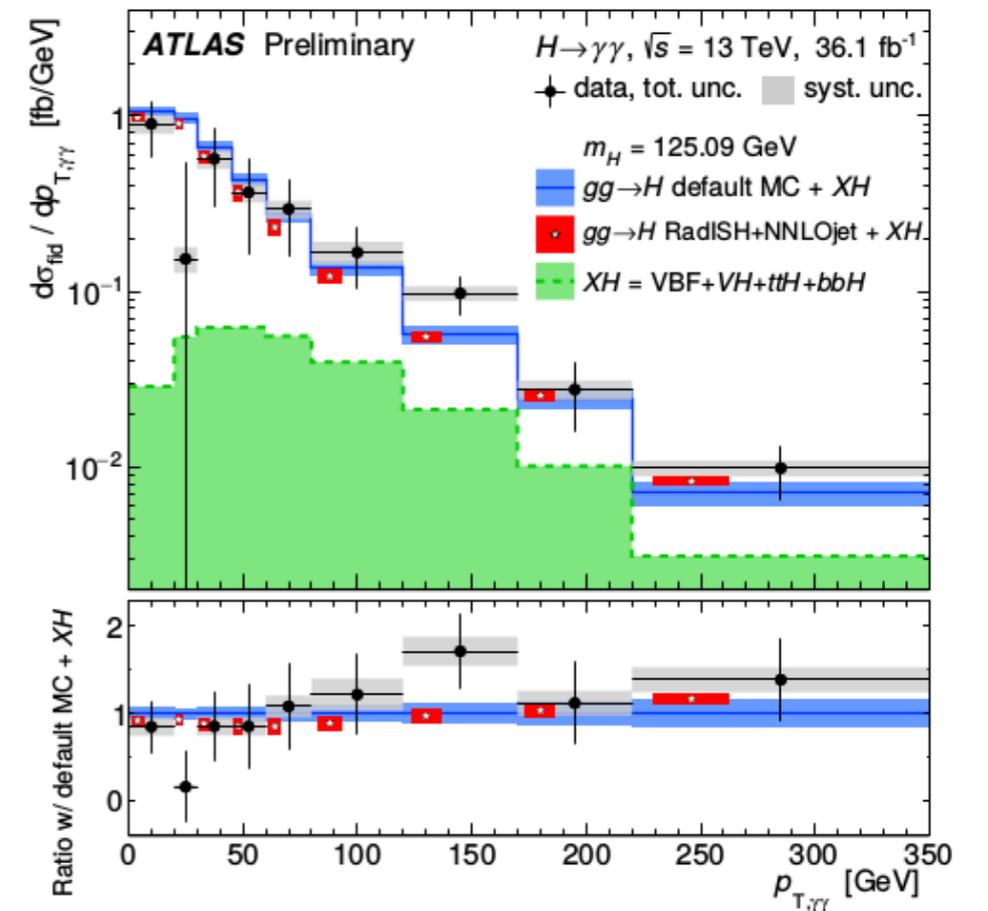
Models



...

[Gaillard '86, Henning et al 1412.1837, 1604.01019, de Blas et al 1711.10391]

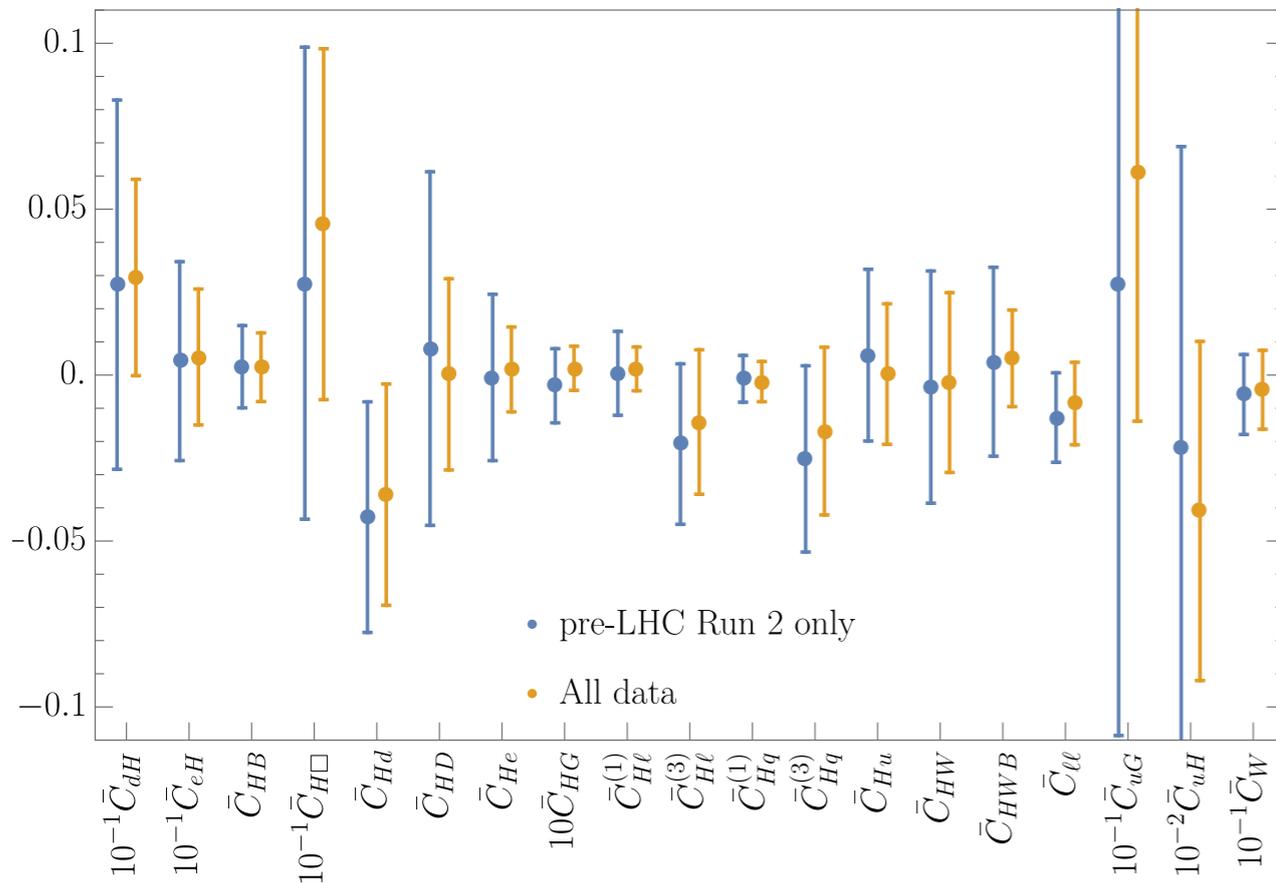
Observables



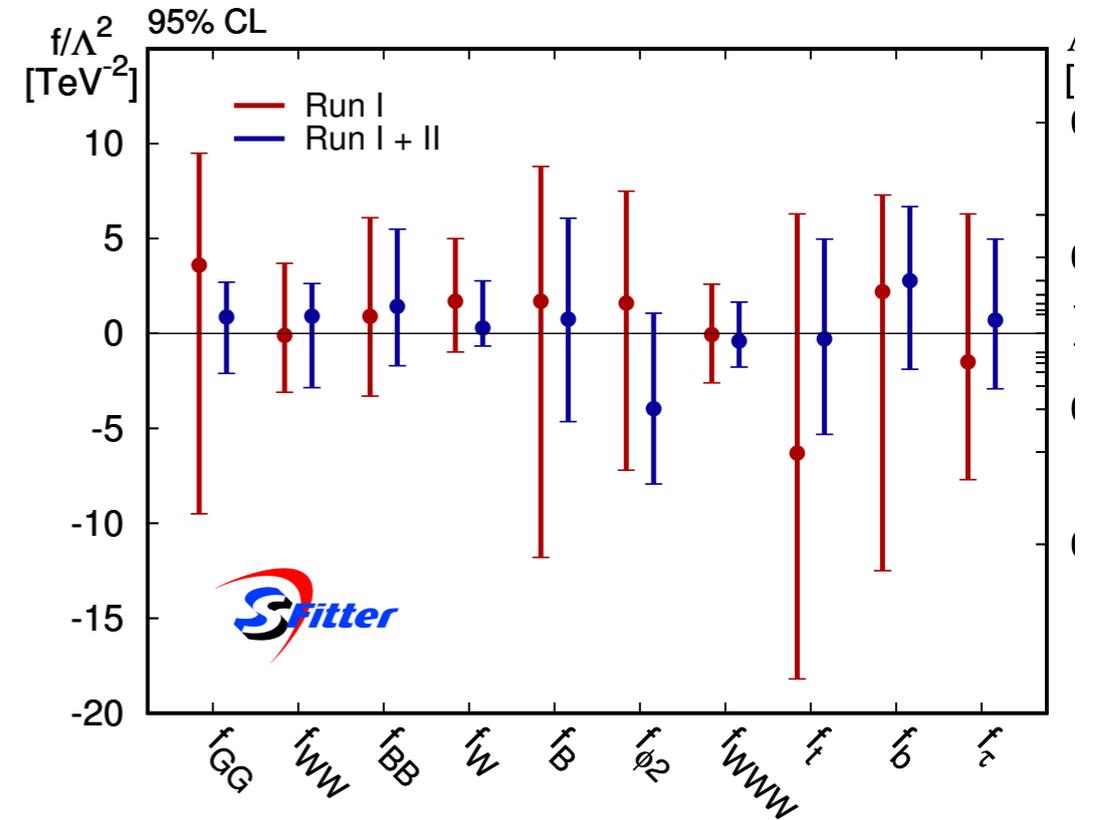
Go heavy: SMEFT

Recent results: global fits to Higgs data + LEP I/II

[Ellis et al 1803.03252]



[Plehn et al 1812.07587]



What's next?

[see talk by Tian]

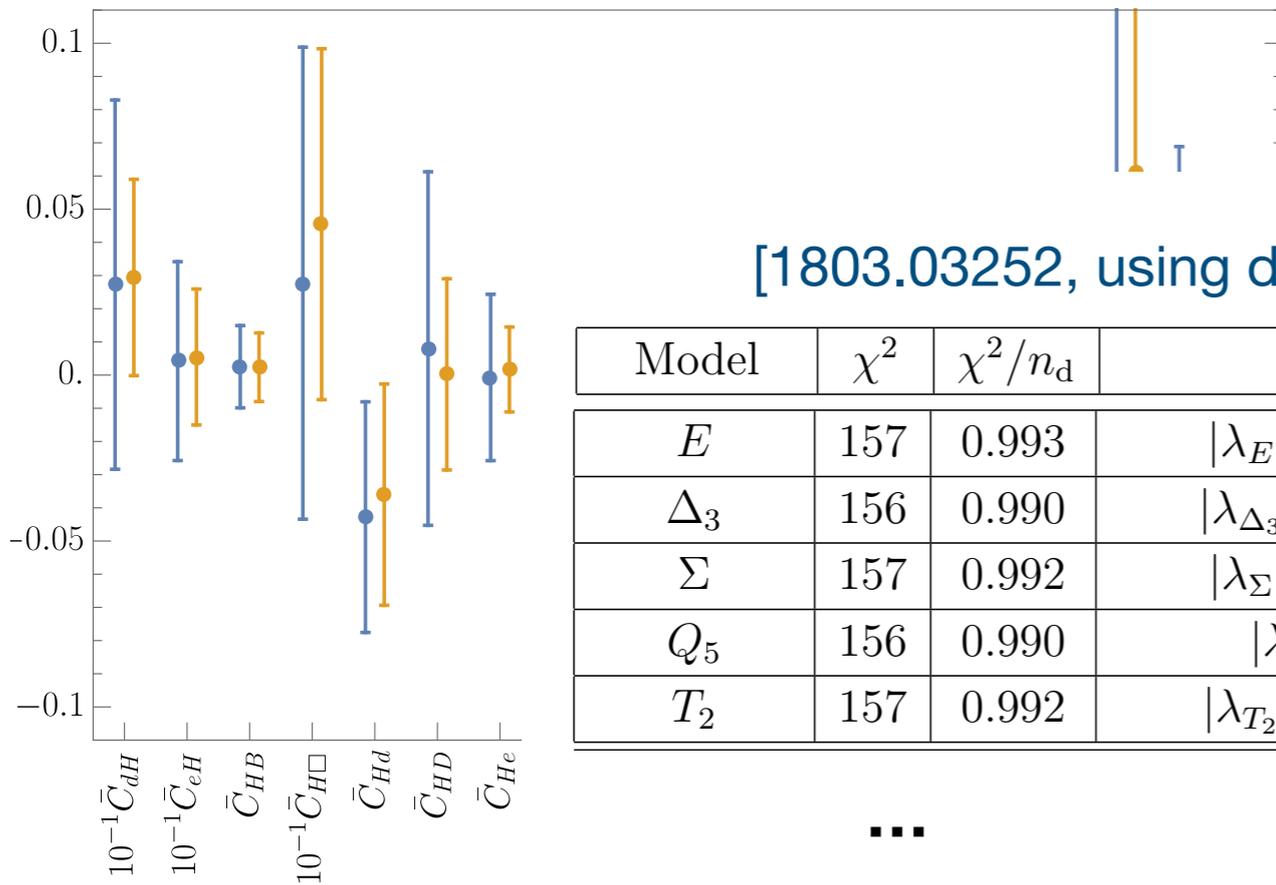
$\Lambda = \text{NP scale} = \text{where you want to put the next collider, so knowing the uncertainty is important}$

May require EFT + NLO SM, or dim-8 EFT

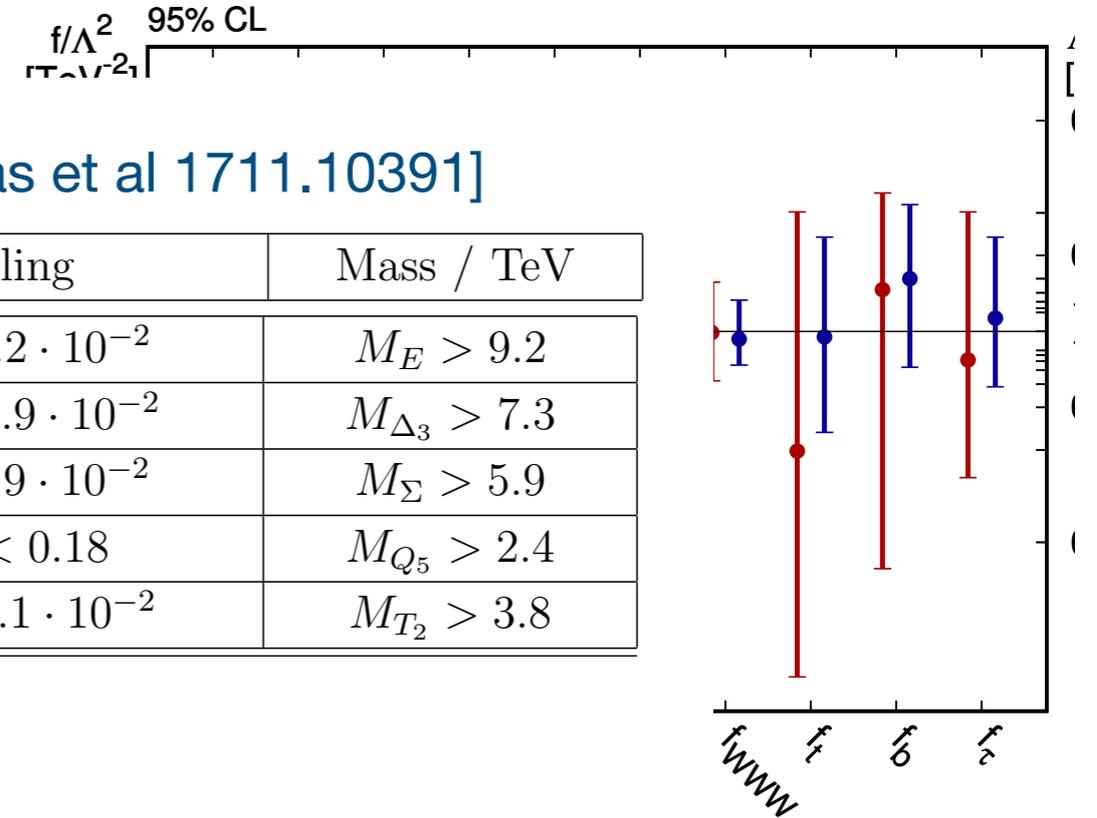
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Recent results: global fits to Higgs data + LEP I/II

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[Plehn et al 1812.07587]



[1803.03252, using de Blas et al 1711.10391]

Model	χ^2	χ^2/n_d	Coupling	Mass / TeV
E	157	0.993	$ \lambda_E ^2 < 1.2 \cdot 10^{-2}$	$M_E > 9.2$
Δ_3	156	0.990	$ \lambda_{\Delta_3} ^2 < 1.9 \cdot 10^{-2}$	$M_{\Delta_3} > 7.3$
Σ	157	0.992	$ \lambda_\Sigma ^2 < 2.9 \cdot 10^{-2}$	$M_\Sigma > 5.9$
Q_5	156	0.990	$ \lambda_{Q_5} ^2 < 0.18$	$M_{Q_5} > 2.4$
T_2	157	0.992	$ \lambda_{T_2} ^2 < 7.1 \cdot 10^{-2}$	$M_{T_2} > 3.8$

...

What's next?

[see talk by Tian]

$\Lambda = \text{NP scale} = \text{where you want to put the next collider, so knowing the uncertainty is important}$

May require EFT + NLO SM, or dim-8 EFT

Go color neutral:

LHC is great at producing of colored states = fit well with naturalness solutions with colored 'top-partners', etc. close to TeV scale

Much weaker limits for uncolored states: squark limits ~ 1.5 TeV vs. electroweakino limits ~ 300 GeV)

Go color neutral:

Much weaker limits for uncolored states: squark limits ~ 1.5 TeV vs. electroweakino limits ~ 300 GeV)

Lots of model building recently working to solve hierarchy problem using color neutral partners = 'neutral naturalness'.

	Scalar Top Partner	Fermionic Top Partner
Colored Top Partner	Supersymmetry	Compositeness/Warped Extra Dimensions
EW-charged Top Partner	Folded SUSY [Chacko, Harnik, Goh, Burdman]	Quirky Little Higgs [Cai, Cheng, Terning]
Neutral Top Partner	Hyperbolic Higgs [Cohen, Craig, Giudice, McCullough] Tripled Top [Cheng, Li, Salvioni, Verhaaren]	Twin Higgs [Chacko, Harnik, Goh] [Barbieri, Gregoire, Hall]

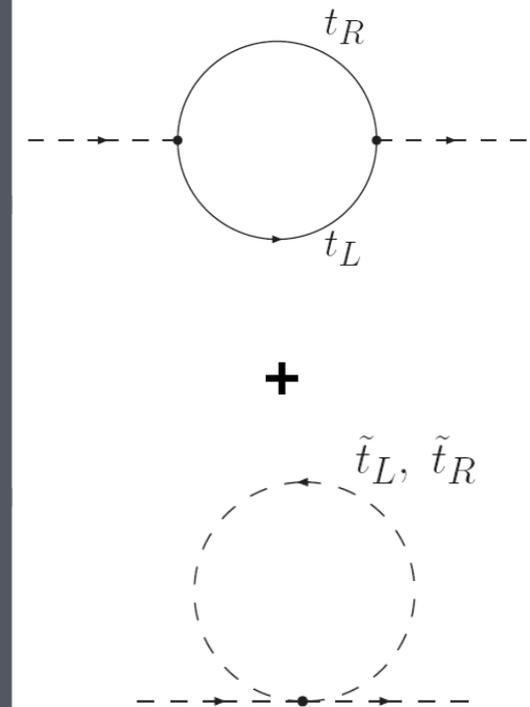


Table from B. Batell

[Table inspiration from Curtin, Verhaaren]

Go color neutral:

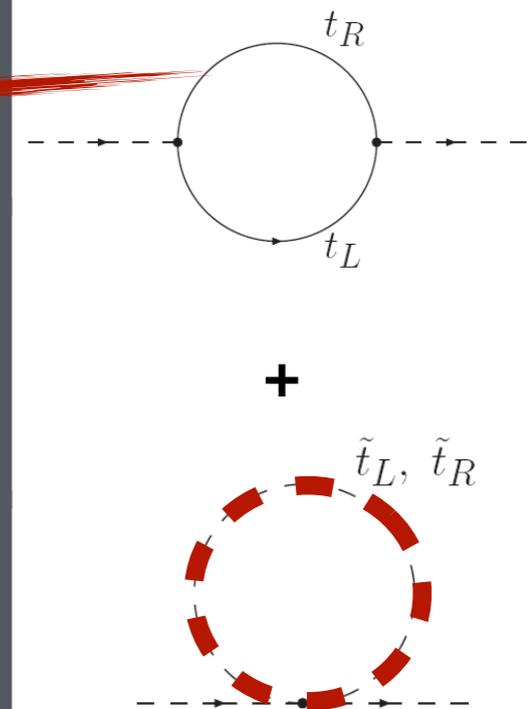
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Lots of model building recently working to solve hierarchy problem using color neutral partners = 'neutral naturalness'.

	Scalar Top Partner	Fermionic Top Partner	
Colored Top Partner	Supersymmetry	Compositeness/Warped Extra Dimensions	<p>+ extra Higgs, mixes with SM Higgs</p>
EW-charged Top Partner	Folded SUSY [Chacko, Harnik, Goh, Burdman]	Quirky Little Higgs [Cai, Cheng, Terning]	
Neutral Top Partner	Hyperbolic Higgs [Cohen, Craig, Giudice, McCullough] Tripled Top [Cheng, Li, Salvioni, Verhaaren]	Twin Higgs [Chacko, Harnik, Goh] [Barbieri, Gregoire, Hall]	

Table from B. Batell [Table inspiration from Curtin, Verhaaren]

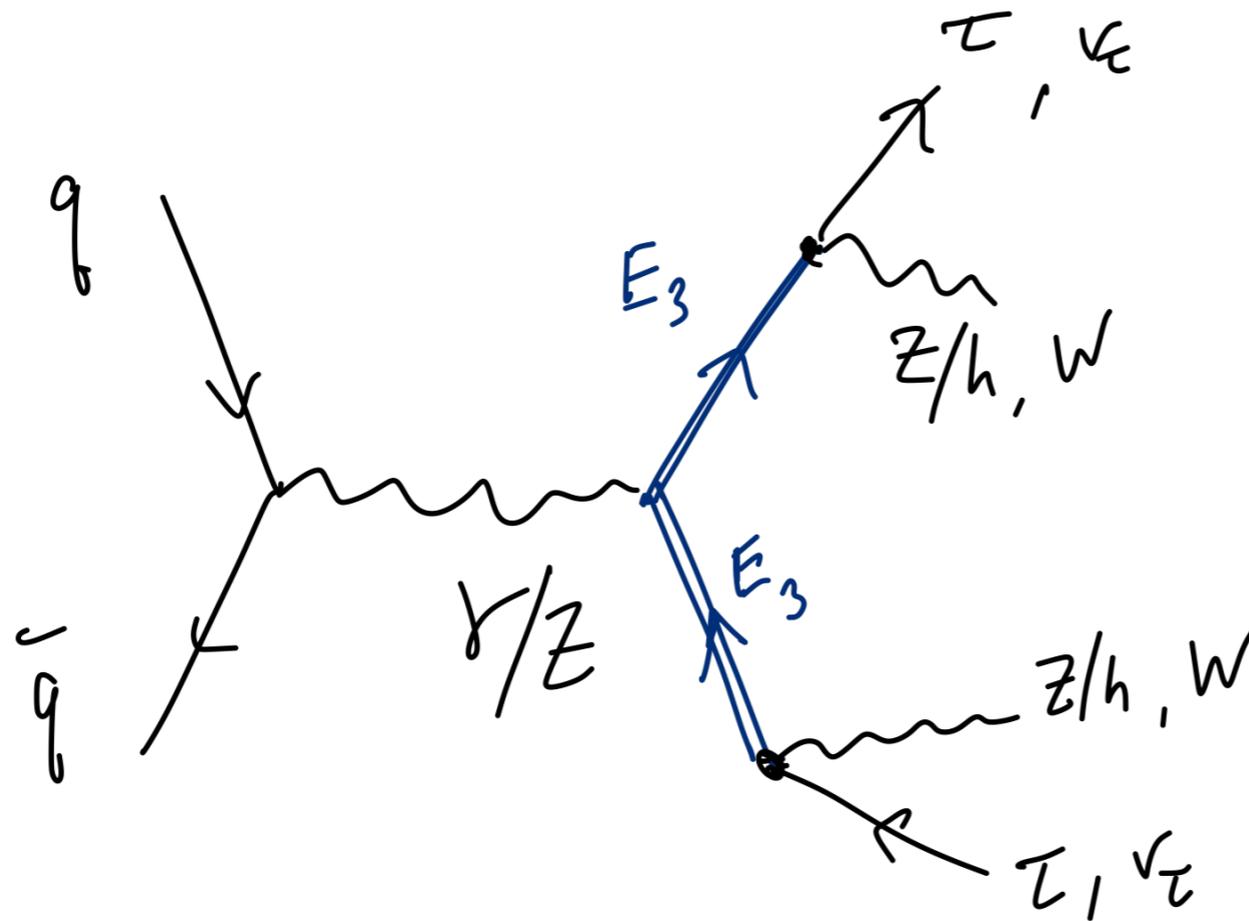
$$y_A Q_A H_A u_A + y_B Q_B H_B u_B + \lambda |H|_A^2 |H_B|^2$$

Go color neutral:

LHC limits weaken even further if we have **no BSM MET**

ex.) new lepton with vector-like charge assignments under $U(1)_Y$, talks primarily to tau

[Bhattiprolu, S. Martin 1905.00498]



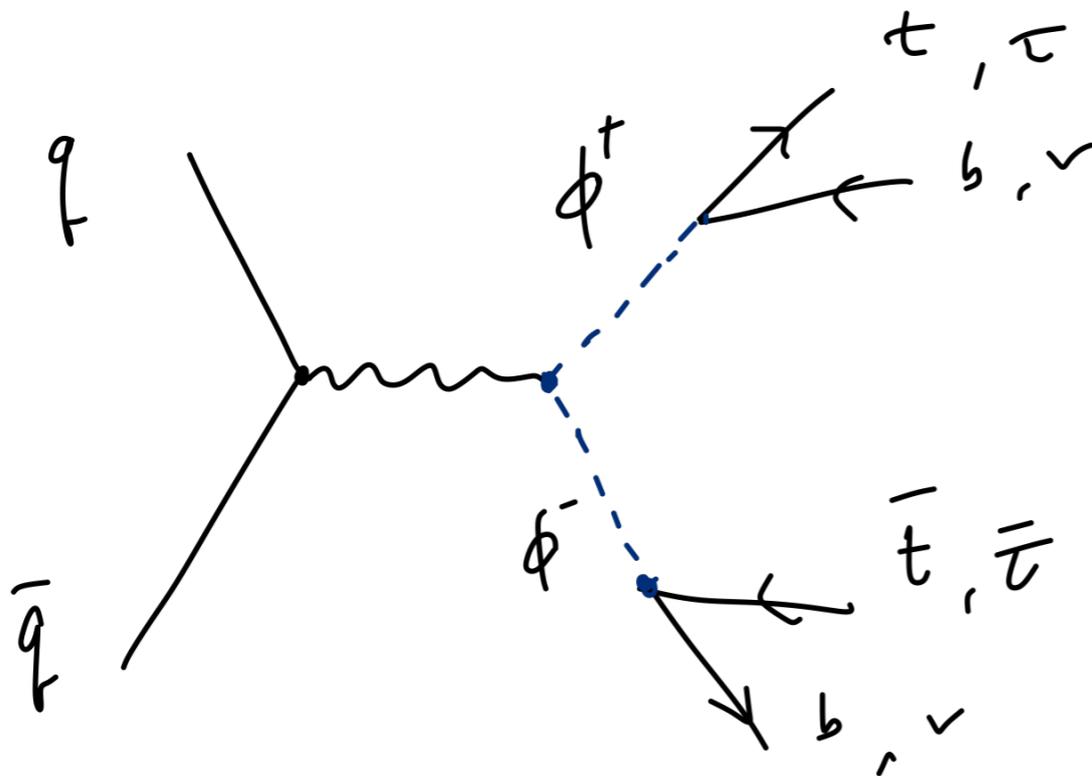
Best limit: LEP II

Could be: the lightest stuff from heavy SUSY/compositeness

Go color neutral:

ex .) new $SU(2)_w$ triplet scalar, talks to SM fermions via higher dimensional operators

[Kribs, AM, Ostdiek, Tong 1809.10184]



$$\mathcal{L}_{decay} \supset \frac{y_u Q \tau^A H u_c \phi^A}{\Lambda} + \frac{y_\ell L \tau^A H e_c \phi^A}{\Lambda} + \dots$$

* ($pp \rightarrow \phi^\pm \phi^0$ too)

Could be... a remnant from solutions to other SM shortcomings (DM)

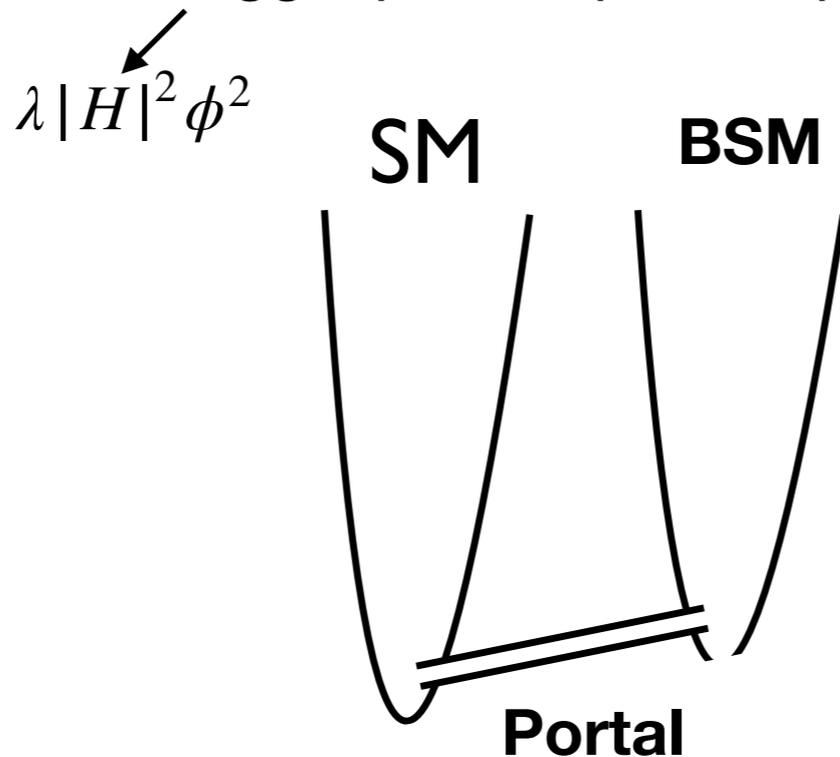
8 TeV limits > 13 TeV limits, as 13 TeV analyses focus on heavier stuff with BSM MET

Model independent, multi-lepton searches at 8 TeV (ATLAS 1411.2921, CMS 1404.5801) are great & extremely powerful — would be great to have 13 TeV versions

Go long lived:

$$\epsilon F^{\mu\nu} X_{\mu\nu} \quad \frac{(SM)(BSM)}{M_{mediator}}$$

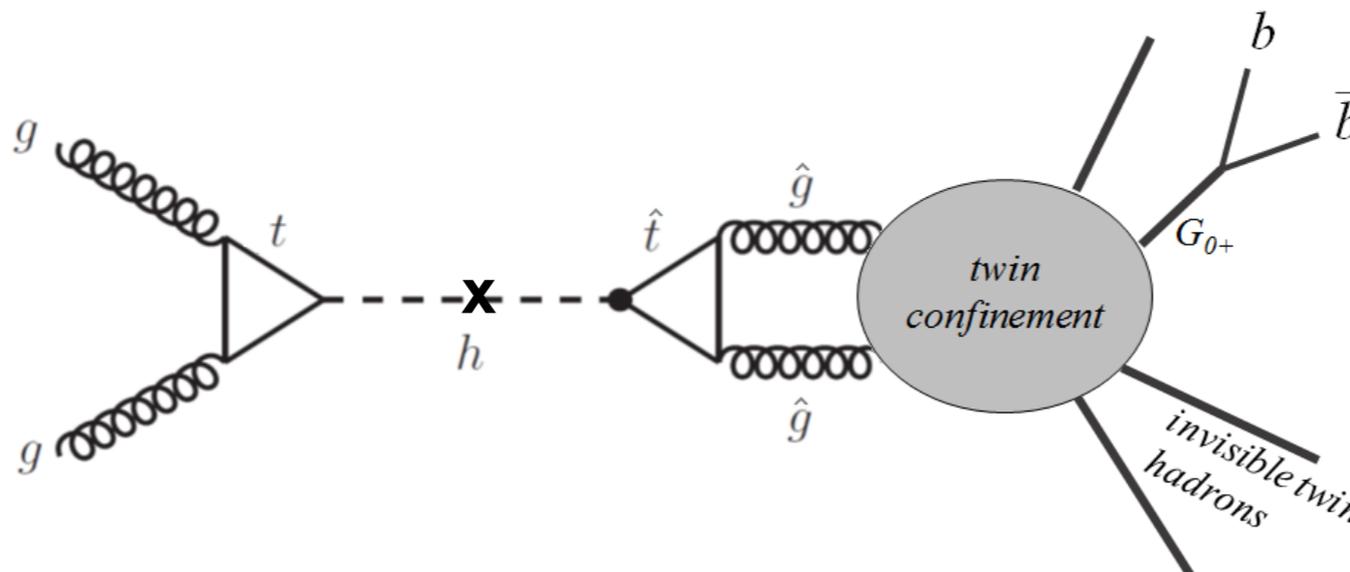
Even weaker connections: Higgs 'portal', photon portal, heavy mediator portal, etc.



hidden by weak coupling rather than high mass

If SM \leftrightarrow BSM connection is sufficiently weak \rightarrow **long lived particles (LLP)**!
 Show up as exotic decays of SM particles and/or displaced vertices.

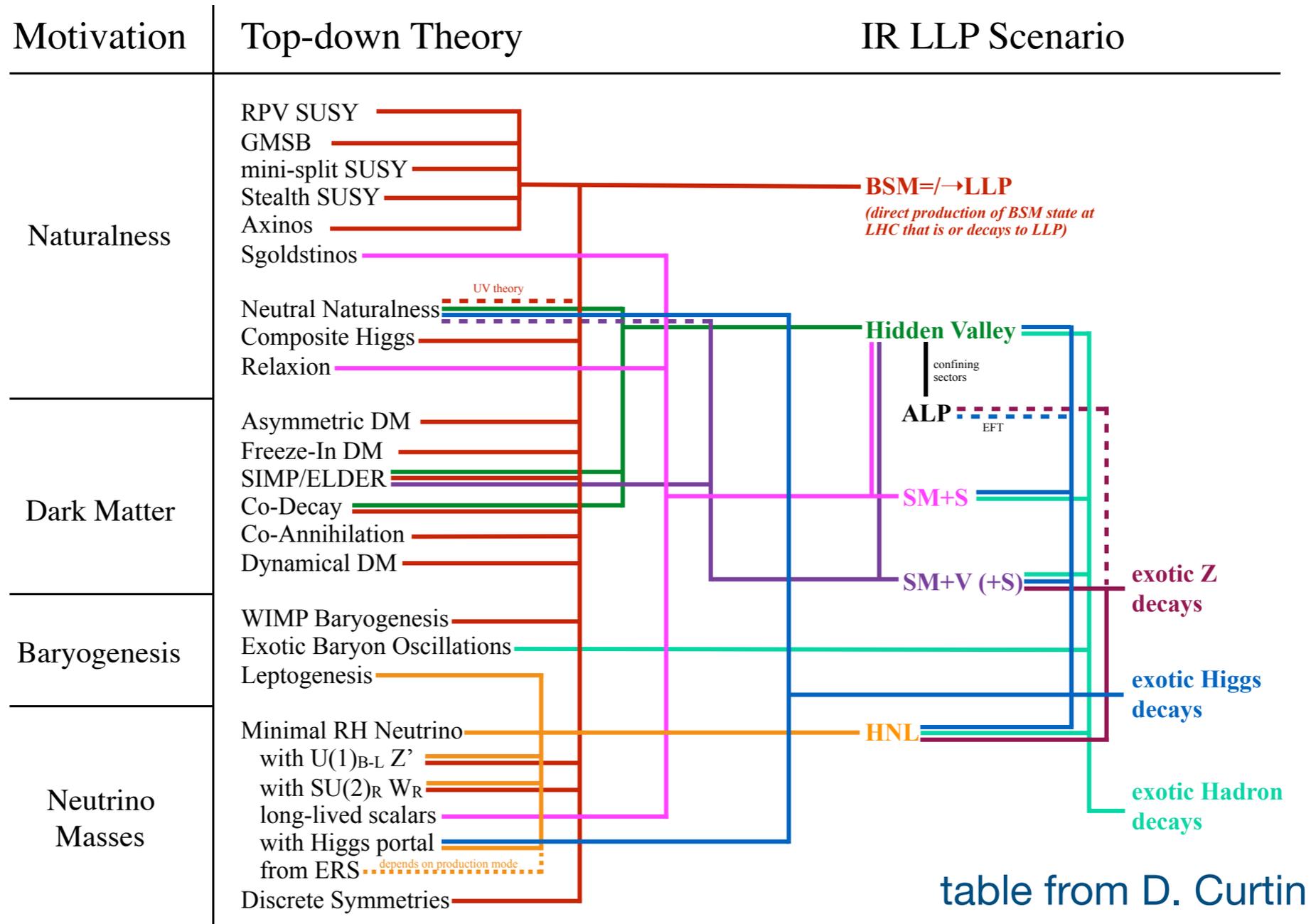
ex.)



Neutral naturalness signal!

Go long lived:

LLPs are an ingredient in many 'motivated' UV scenarios



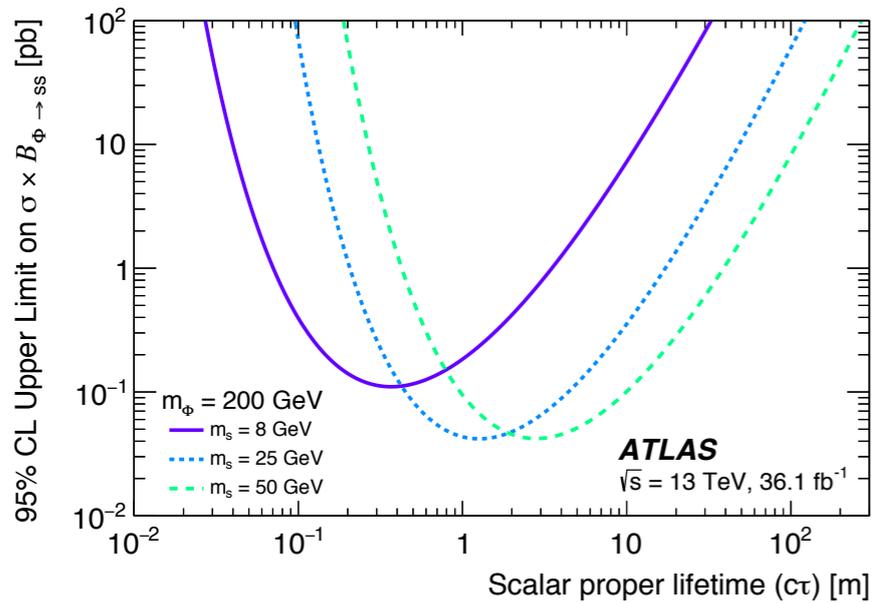
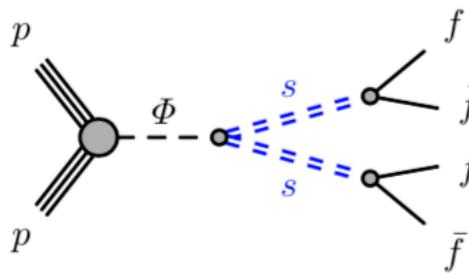
Go long lived:

Neutral LLPs: often difficult to trigger on and with hard to predict backgrounds, but lots of attention recently

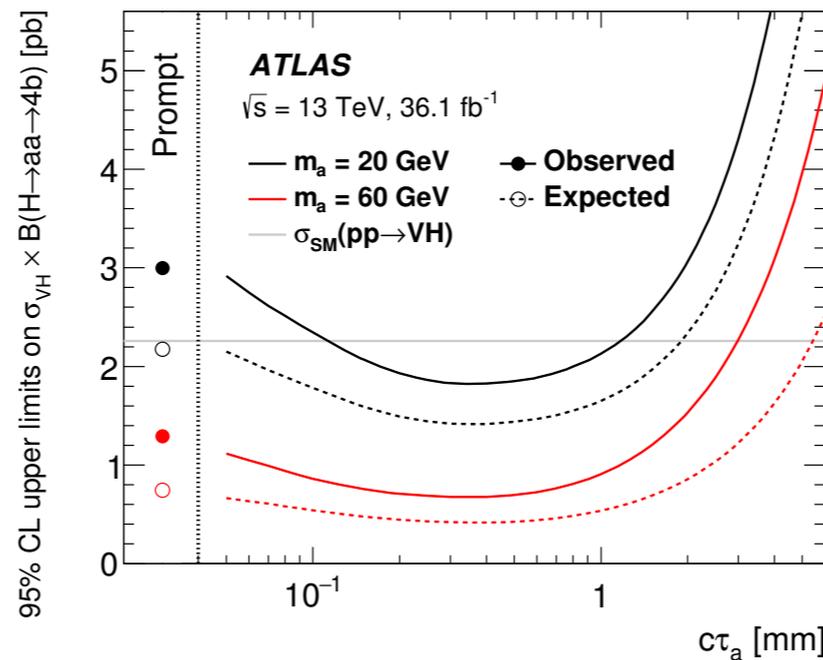
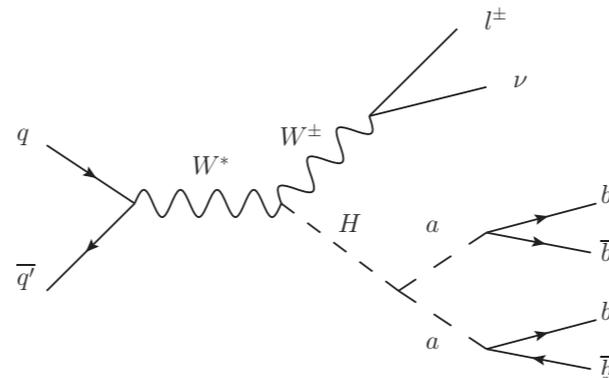
(see talks by Proffitt, Adams, Rifki, Gan,..)

ex.)

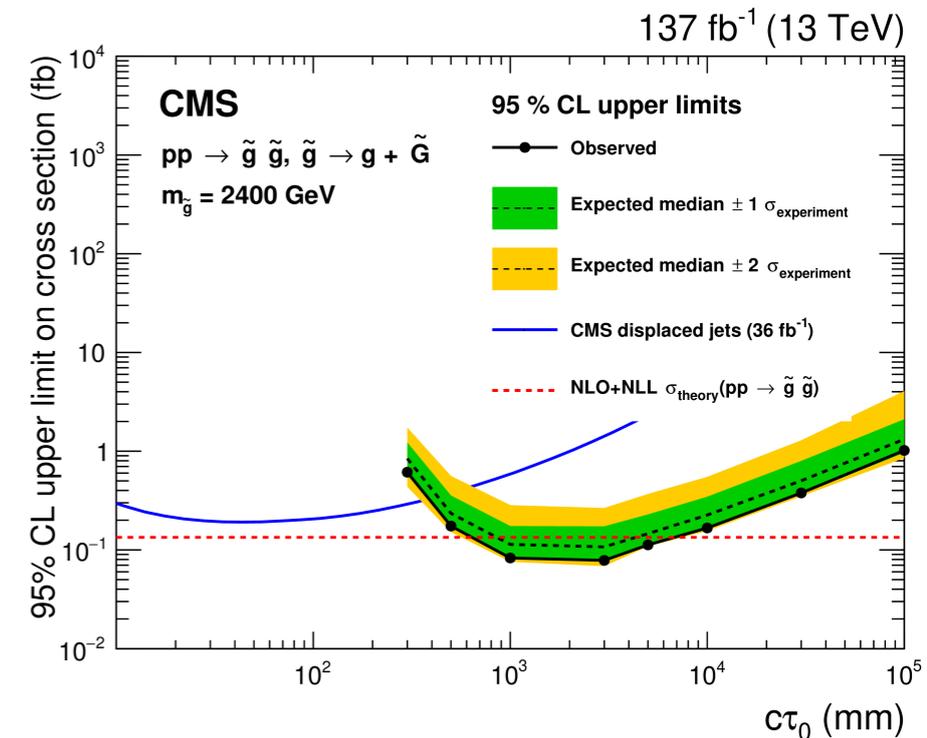
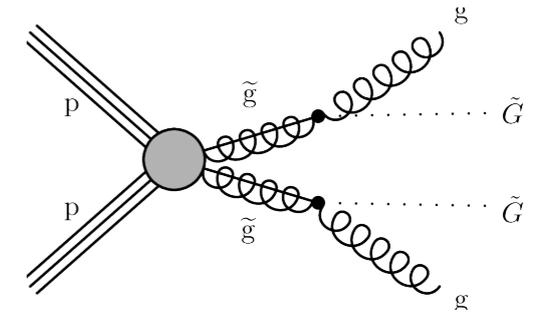
ATLAS displaced jets 1811.07370



ATLAS VH, H-> 2a -> 4b
1806.07355



CMS non-prompt jets 1906.06441



Go long lived:

Neutral LLPs: often difficult to trigger on and with hard to predict backgrounds, but lots of attention recently

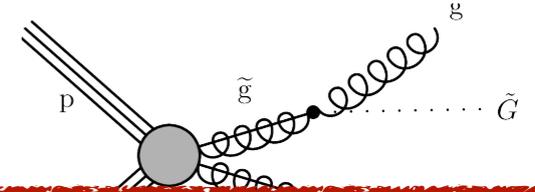
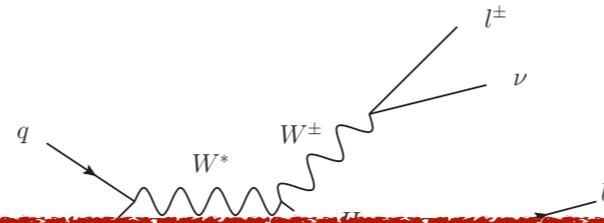
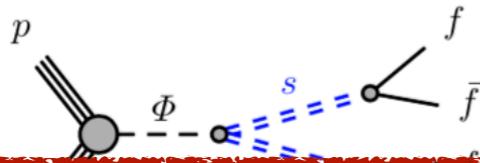
(see talks by Proffitt, Adams, Rifki, Gan,..)

ex.)

ATLAS displaced jets 1811.07370

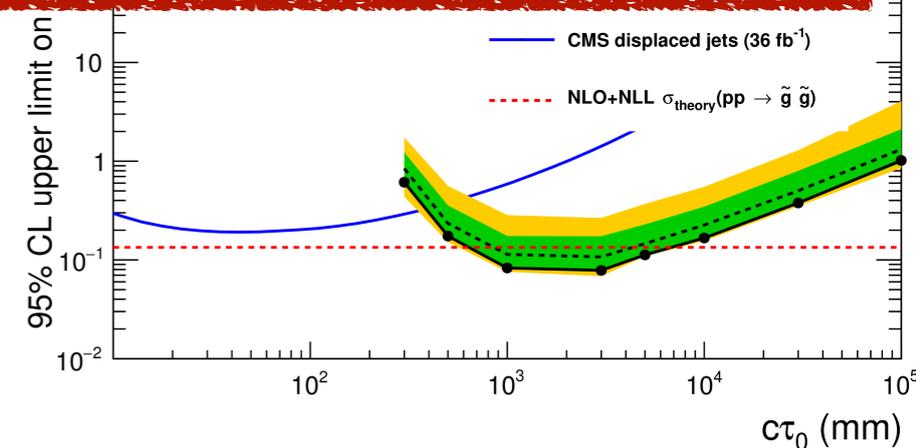
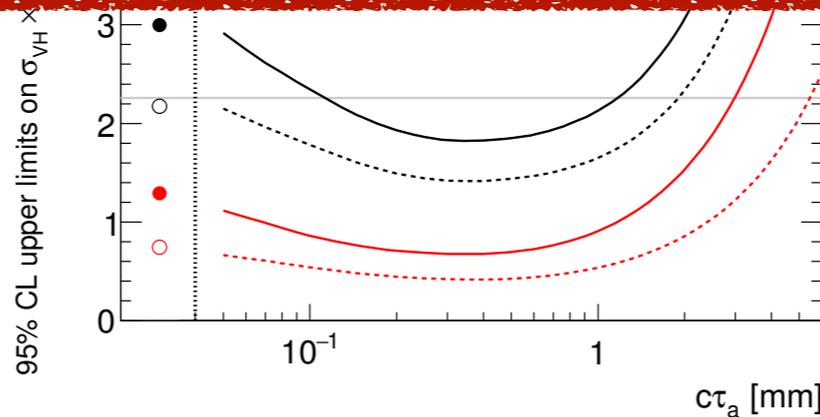
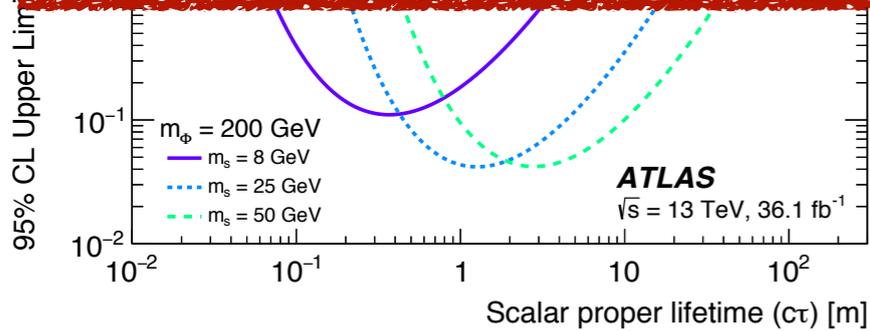
ATLAS VH, $H \rightarrow 2a \rightarrow 4b$
1806.07355

CMS non-prompt jets 1906.06441



LHC upgrades to trigger/timing could have a big impact — see 1903.04497
(talk by Flowers)

Complementarity between ATLAS/CMS detectors & new dedicated LLP experiments/proposals: MATHUSLA, CODEX-B, FASER...



Conclusions

- SUSY/Composite Higgs extensions of SM still well motivated, worth pursuing with full force .. however increasingly strained by null results
- **SMEFT**: model independent parameterization of physics too heavy to make on shell. Well defined, systematic. Lots of directions to pursue, power in differential distributions
- **Keep the net wide:**

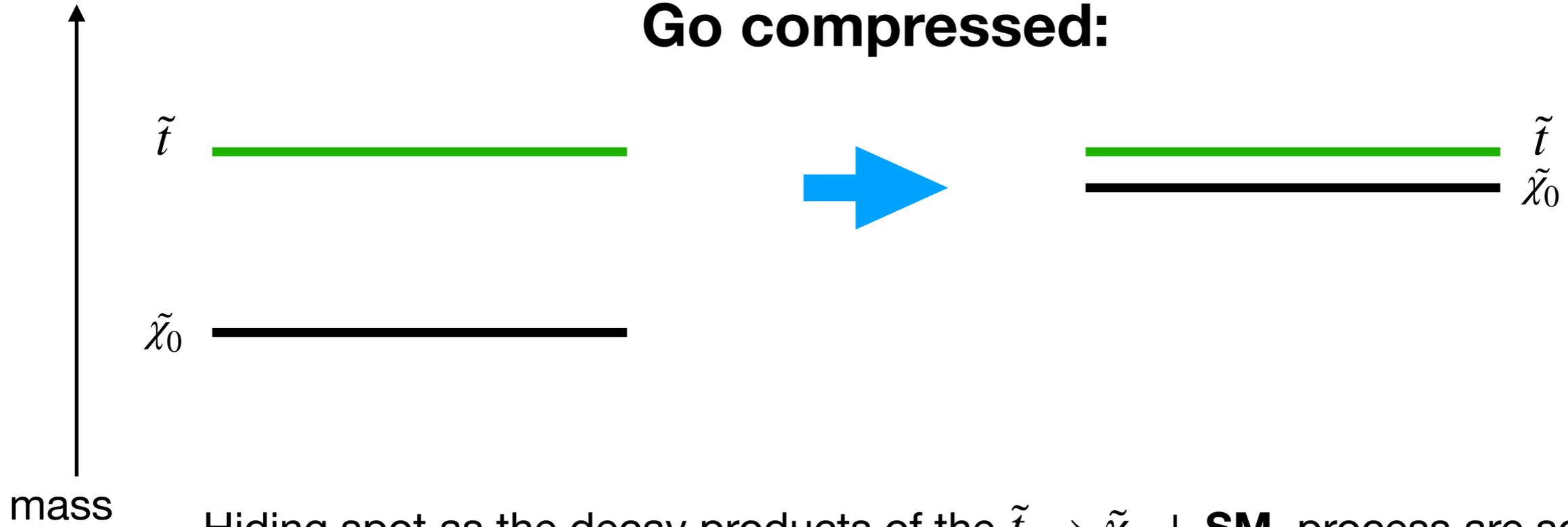
Color neutral states, especially with no BSM MET could still be lurking at light scales.
Keep the model independent searches coming!

Long lived particles: generic feature in many models, including 'neutral natural' setups. New frontier! Many innovative ideas but lots of space left to explore

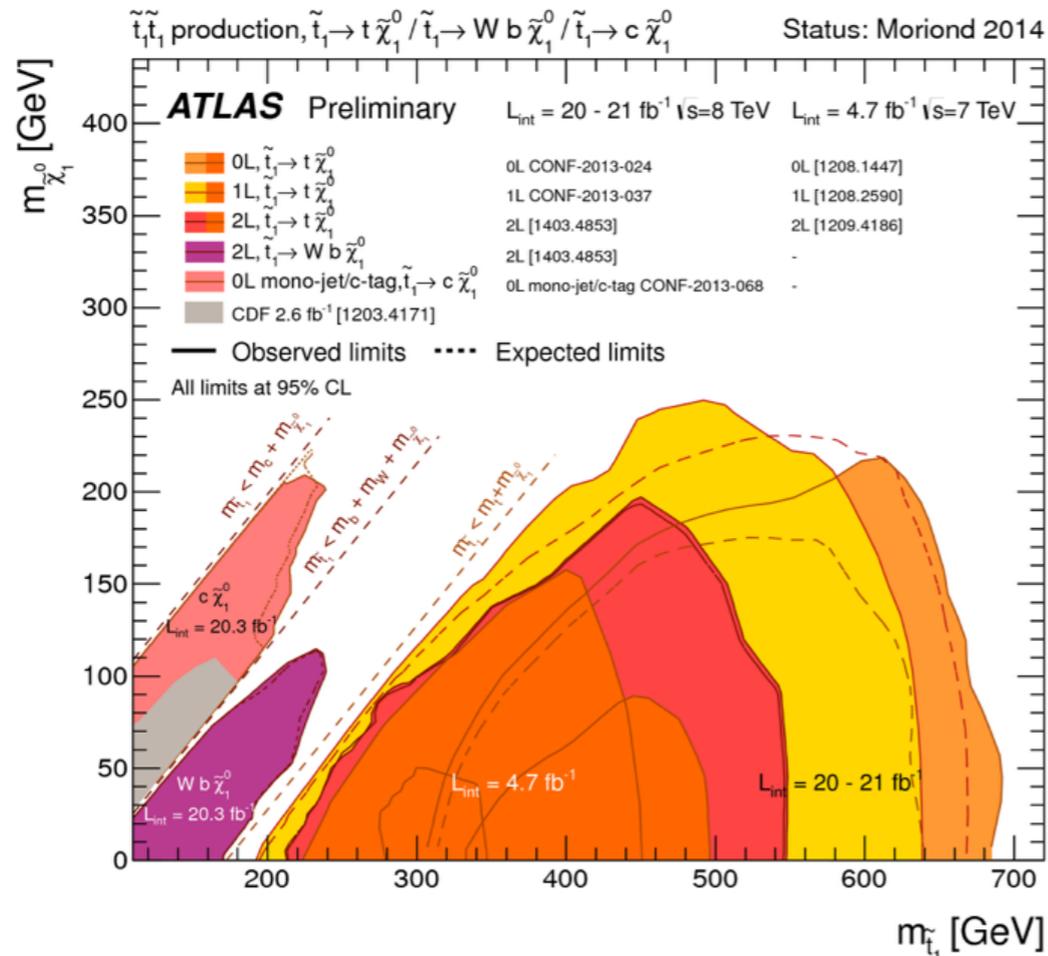
Looking forward to more data & results! Thank you!

EXTRA

Go compressed:

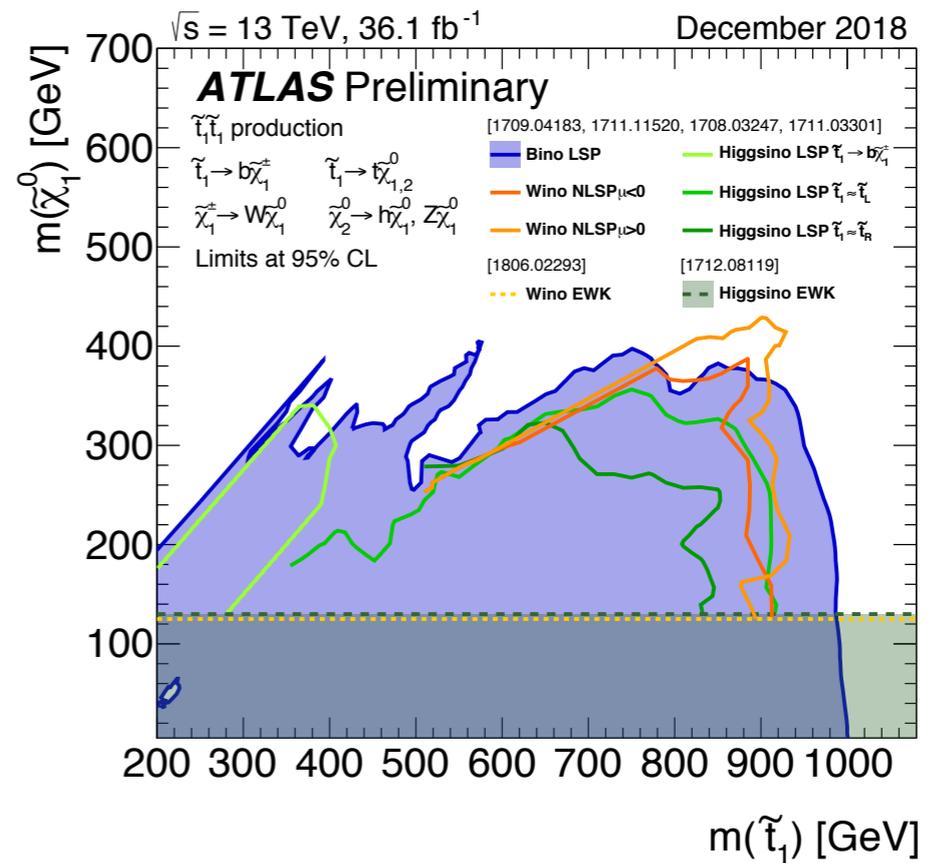
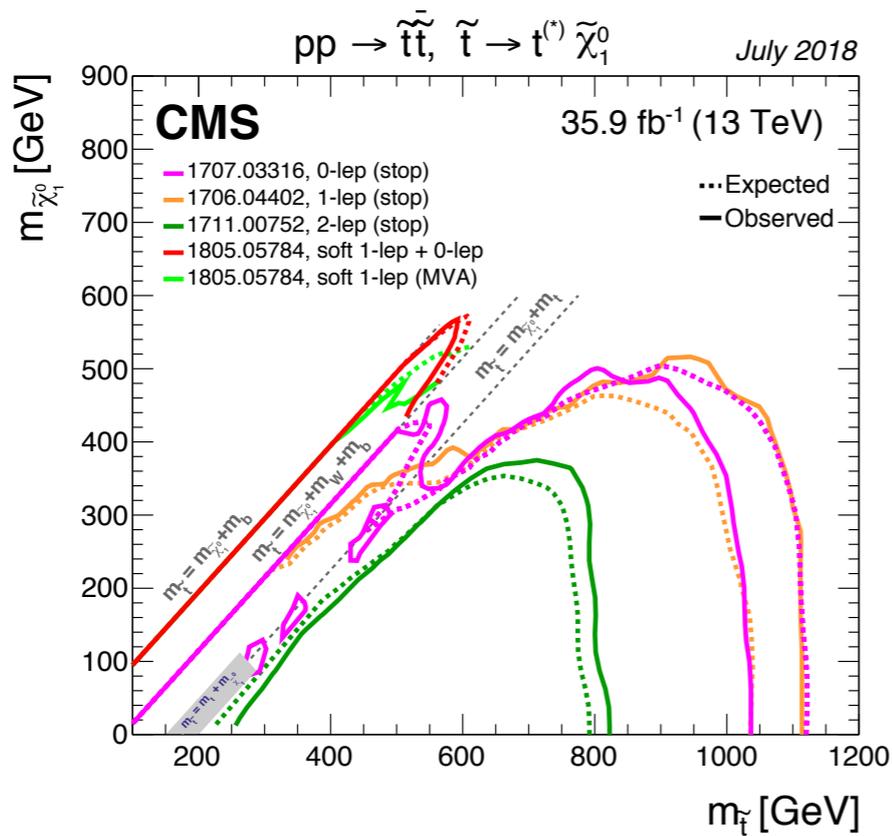
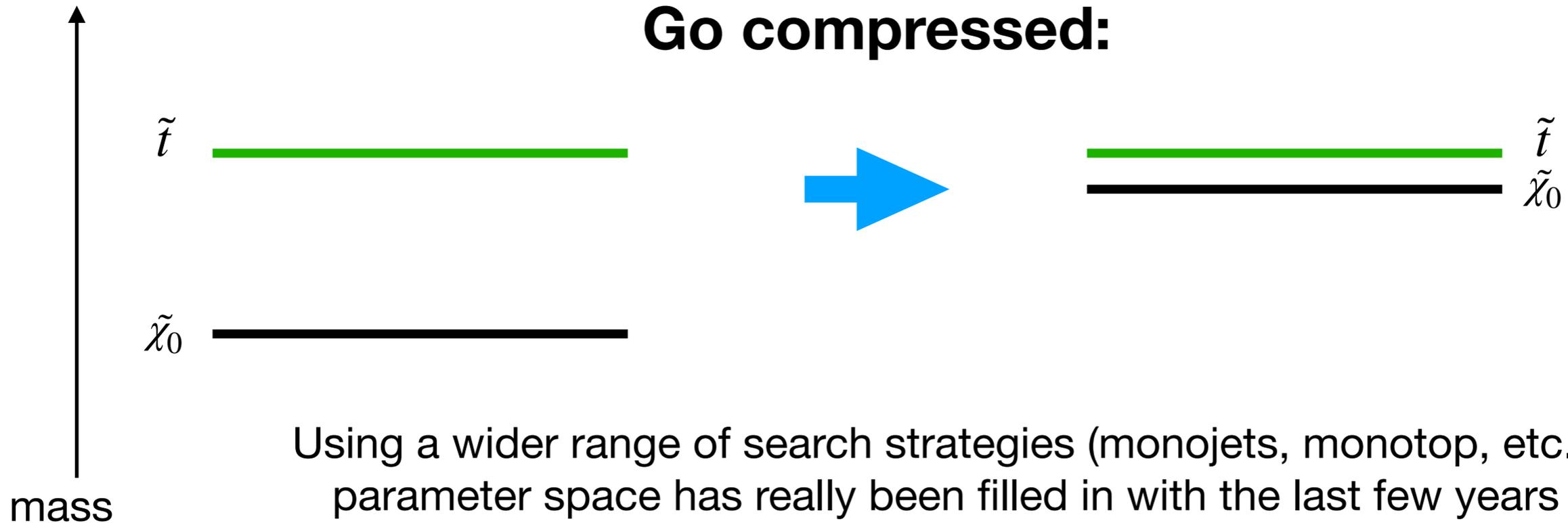


Hiding spot as the decay products of the $\tilde{t} \rightarrow \tilde{\chi}_0 + \mathbf{SM}$ process are soft



Plenty of other examples, e.g. degenerate chargino-neutralino

Go compressed:



Hiding spots remain, but it's getting a lot tougher

(talk by Chen)

Twin Higgs done more correctly...

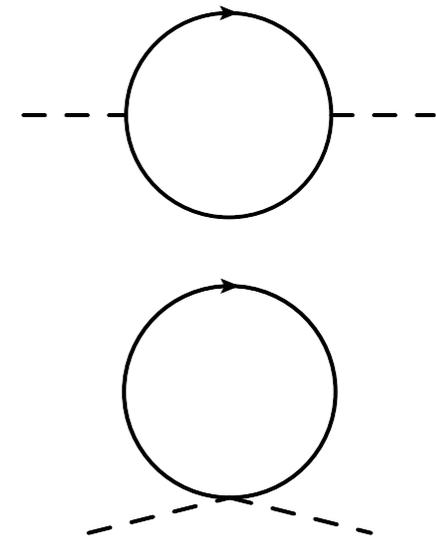
← not SM colored

$$y_t Q_A H_A u_A + y_t Q_B H_B u_B \quad A \leftrightarrow B \text{ related by } Z_2$$

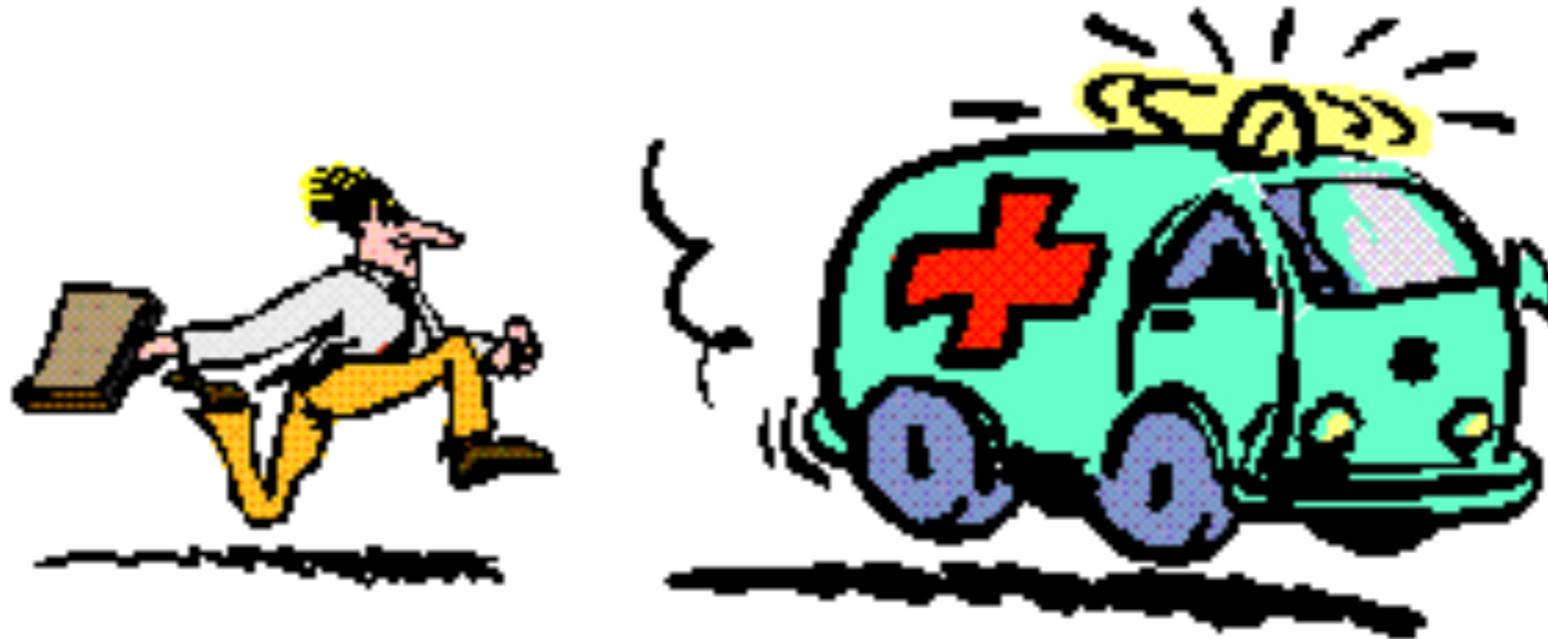
Combine H_A, H_B , they have a $SU(4)$ global symmetry that breaks at f

SM h is Goldstone from $SU(4)$ breaking

$$y_t Q_A h u_A - y_t Q_B \frac{h^2}{f} u_B$$



Chase the anomalies!



- **For the theorists:** 'motivated' models haven't shown up, so keep broadening our scope
- **For the experimenters:** Don't let "no theory paper exists with that particle/final state" stop you

Even if anomaly fades, creative efforts lead to new models/revitalize old models/
point out issues in search strategies!