## Reevaluating Uncertainties in $\bar{B} \to X_s \gamma$ Decay

#### AYESH GUNAWARDANA

Wayne State University

July 30 th, 2019

Based on A.G and Gil Paz "Reducing Uncertainties in  $\bar{B} \to X_s \gamma$ " (To appear)

#### Introduction

# Why $ar{\mathcal{B}} ightarrow extbf{X}_{\!s} \gamma$

Why 
$$\bar{B} o X_s \gamma$$

•  $\bar{B} \to X_s \gamma$  decay is an important **New Physics** probe

## Why $\bar{B} o X_s \gamma$

- $\bar{B} \to X_s \gamma$  decay is an important **New Physics** probe
  - It is suppressed at tree level in SM
  - Can receive contributions from SM extensions.

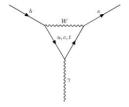


Figure:  $b \rightarrow s \gamma$  flavor changing neural current (FCNC) in SM

## Why $\bar{B} o X_s \gamma$

- $\bar{B} \to X_s \gamma$  decay is an important **New Physics** probe
  - It is suppressed at tree level in SM
  - Can receive contributions from SM extensions.

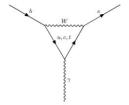


Figure:  $b \rightarrow s \gamma$  flavor changing neural current (FCNC) in SM

- SM extensions modify the  $C_{7\gamma}$  Wilson coefficient

## Why $\bar{B} \to X_s \gamma$

- $\bar{B} \to X_s \gamma$  decay is an important **New Physics** probe
  - It is suppressed at tree level in SM
  - Can receive contributions from SM extensions.

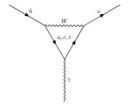


Figure:  $b \rightarrow s \gamma$  flavor changing neural current (FCNC) in SM

- SM extensions modify the  $C_{7\gamma}$  Wilson coefficient
- CP violation in  $\bar{B} o X_s \gamma$  can be enhanced by new physics

## Photon production

### Photon production

• Photon can be produced directly:

$$Q_{7\gamma}=rac{-e}{8\pi^2}m_bar{s}\sigma_{\mu
u}F^{\mu
u}(1+\gamma_5)b$$

#### Photon production

Photon can be produced directly:

$$Q_{7\gamma}=rac{-e}{8\pi^2}m_bar{s}\sigma_{\mu
u}F^{\mu
u}(1+\gamma_5)b$$

Also, gluon or quark pair can convert to photon

$$Q_{8g}=rac{-e}{8\pi^2}m_bar{s}\sigma_{\mu
u}G^{\mu
u}(1+\gamma_5)b$$

$$Q_1^q = (\bar{q}b)_{V-A}(\bar{s}q)_{V-A}$$

ullet The effective Lagrangian to describe  $ar{B} o X_{
m s}\gamma$ 

• The effective Lagrangian to describe  $ar{B} o X_s \gamma$ 

$$\mathcal{H}_{ ext{eff}} = rac{G_F}{\sqrt{2}} \sum_{q=u,c} V_{qb}^* V_{qs} \ (C_1 Q_1^q + \sum_{i=2}^6 C_i Q_i + C_{7\gamma} Q_{7\gamma} + C_{8g} Q_{8g}) + ext{h.c.}$$

• The effective Lagrangian to describe  $ar{B} o X_s \gamma$ 

$$\mathcal{H}_{ ext{eff}} = rac{G_F}{\sqrt{2}} \sum_{q=u,c} V_{qb}^* V_{qs} \ (C_1 Q_1^q + \sum_{i=2}^6 C_i Q_i + C_{7\gamma} Q_{7\gamma} + C_{8g} Q_{8g}) + ext{h.c.}$$

- Most important operators are  $Q_{7\gamma},\,Q_{8g}$  and  $Q_1^q.$
- $Q_{7\gamma}=rac{-e}{8\pi^2}m_bar{s}\sigma_{\mu
  u}F^{\mu
  u}(1+\gamma_5)b$
- $Q_1^q = (\bar{q}b)_{V-A}(\bar{s}q)_{V-A}$   $Q_{8g} = \frac{-e}{8\pi^2}m_b\bar{s}\sigma_{\mu\nu}G^{\mu\nu}(1+\gamma_5)b$

ullet The effective Lagrangian to describe  $ar{\mathcal{B}} o X_{m s}\gamma$ 

$$\mathcal{H}_{ ext{eff}} = rac{G_F}{\sqrt{2}} \sum_{q=u,c} V_{qb}^* V_{qs} \ (C_1 Q_1^q + \sum_{i=2}^6 C_i Q_i + C_{7\gamma} Q_{7\gamma} + C_{8g} Q_{8g}) + ext{h.c.}$$

- Most important operators are  $Q_{7\gamma}, Q_{8g}$  and  $Q_1^q$ .
- $Q_{7\gamma}=rac{-e}{8\pi^2}m_bar{s}\sigma_{\mu
  u}F^{\mu
  u}(1+\gamma_5)b$
- $Q_1^q=(ar q b)_{V-A}(ar s q)_{V-A}$   $Q_{8g}=rac{-e}{8\pi^2}m_bar s\sigma_{\mu
  u}G^{\mu
  u}(1+\gamma_5)b$
- At leading power: Only  $Q_{7\gamma}-Q_{7\gamma}$  contributes to decay rate
- At  $1/m_b$ :  $\Gamma$  get  $Q_1-Q_{7\gamma}$ ,  $Q_{8g}-Q_{8g}$  and  $Q_{7\gamma}-Q_{8g}$  contributions

#### Decay rate

World average for experimental value:

$$\mathcal{B}(B \to X_s \gamma) (E_{\gamma} > 1.6 \text{ GeV}) = (3.32 \pm 0.15) \times 10^{-4}$$

[ Y. Amhis et. al. EPJC 77, 895 (2017)]

NNLO prediction

$$\Gamma\left(\overline{B} \to X_q \gamma\right) = \underbrace{\Gamma\left(b \to X_q^p \gamma\right)}_{\text{Perturbatively calculable}} + \underbrace{\delta\Gamma_{\text{nonp}}}_{\mathcal{O}(\frac{\Lambda_{\text{QCD}}}{m_b})}$$

• SM prediction (2015) [Misiak et. al. PRL 114, 221801 (2015)]

$$\mathcal{B}_{s\gamma}^{\rm SM} = (3.36 \pm 0.23) \times 10^{-4}$$

for 
$$E_{\gamma} > 1.6 \; {\rm GeV}$$

#### Decay rate

World average for experimental value:

$$\mathcal{B}(B \to X_s \gamma) (E_{\gamma} > 1.6 \text{ GeV}) = (3.32 \pm 0.15) \times 10^{-4}$$

[ Y. Amhis et. al. EPJC 77, 895 (2017)]

NNLO prediction

$$\Gamma\left(\overline{B} \to X_q \gamma\right) = \underbrace{\Gamma\left(b \to X_q^p \gamma\right)}_{\text{Perturbatively calculable}} + \underbrace{\delta\Gamma_{\text{nonp}}}_{\mathcal{O}(\frac{\Lambda_{\text{QCD}}}{m_b})}$$

• SM prediction (2015) [Misiak et. al. PRL 114, 221801 (2015)]

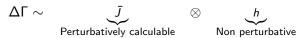
$$\mathcal{B}_{s\gamma}^{\rm SM} = (3.36 \pm 0.23) \times 10^{-4}$$

for  $E_{\gamma} > 1.6 \text{ GeV}$ 

- $\delta\Gamma_{\rm nonp} \equiv \text{Non-perturbative contribution}$ 
  - The largest contribution to the error 5% from  $\mathcal{O}(\frac{\Lambda_{\text{QCD}}}{m_b})$

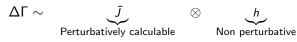
## Order $1/m_b$ power corrections to $\Gamma(\bar{B} \to X_s \gamma)$

Non-perturbative effects arise from Resolved Photon Contributions

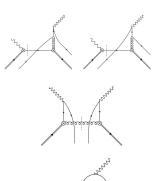


## Order $1/m_b$ power corrections to $\Gamma(\bar{B} \to X_s \gamma)$

Non-perturbative effects arise from Resolved Photon Contributions



-  $Q_{7\gamma}-Q_{8g}$ 



-  $Q_{8g}-Q_{8g}$ 





- 2010 estimates for non-perturbative contribution to error
  - From  $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
  - From  $Q_{8g} Q_{8g} \in [-0.3, +1.9]\%$
  - From  $Q_{7\gamma} Q_{8g} \in [-4.4, +5.6]\%$

- 2010 estimates for non-perturbative contribution to error
  - From  $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
  - From  $Q_{8g} Q_{8g} \in [-0.3, +1.9]\%$
  - From  $Q_{7\gamma} Q_{8g} \in [-4.4, +5.6]\%$
- The contribution from  $Q_{7\gamma} Q_{8g}$ 
  - Obtained on experiment with 95% confidence level range

[M. Benzke, S. J. Lee, M. Neubert and G. Paz JHEP 1008, 099(2010)]

- 2010 estimates for non-perturbative contribution to error
  - From  $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
  - From  $Q_{8g} Q_{8g} \in [-0.3, +1.9]\%$
  - From  $Q_{7\gamma} Q_{8g} \in [-4.4, +5.6]\%$
- The contribution from  $Q_{7\gamma} Q_{8g}$ 
  - Obtained on experiment with 95% confidence level range
  - [M. Benzke, S. J. Lee, M. Neubert and G. Paz JHEP 1008, 099(2010)]
    - New Belle result for  $Q_{7\gamma}-Q_{8g}$  contribution  $\sim 2\%$
  - [S. Watanuki et. al. PRD 99, 032012(2019)]

- 2010 estimates for non-perturbative contribution to error
  - From  $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
  - From  $Q_{8g} Q_{8g} \in [-0.3, +1.9]\%$
  - From  $Q_{7\gamma} Q_{8g} \in [-4.4, +5.6]\%$
- The contribution from  $Q_{7\gamma} Q_{8g}$ 
  - Obtained on experiment with 95% confidence level range
  - [M. Benzke, S. J. Lee, M. Neubert and G. Paz JHEP 1008, 099(2010)]
    - New Belle result for  $Q_{7\gamma}-Q_{8g}$  contribution  $\sim 2\%$
  - [S. Watanuki et. al. PRD 99, 032012(2019)]
- Now  $Q_1^c Q_{7\gamma}$  is the largest contribution to the error! Can we reduce it?

## $Q_1^c - Q_{7\gamma}$ contribution

• The contribution to the error from  $Q_1^c - Q_{7\gamma}$  is given by

$$\frac{C_1}{C_{7\gamma}}\frac{\Lambda_{17}}{m_b}$$

### $Q_1^c - Q_{7\gamma}$ contribution

• The contribution to the error from  $Q_1^c-Q_{7\gamma}$  is given by

$$\frac{C_1}{C_{7\gamma}}\frac{\Lambda_{17}}{m_b}$$

where

$$\Lambda_{17} = e_c \operatorname{Re} \int_{-\infty}^{\infty} \frac{d\omega_1}{\omega_1} \left| 1 - \underbrace{F\left(\frac{m_c^2 - i\varepsilon}{m_b\omega_1}\right)}_{\text{perturbative}} + \frac{m_b\omega_1}{12m_c^2} \right| \underbrace{h_{17}\left(\omega_1\right)}_{\text{non-perturbative}}$$

### $Q_1^c - Q_{7\gamma}$ contribution

• The contribution to the error from  $Q_1^c-Q_{7\gamma}$  is given by

$$\frac{C_1}{C_{7\gamma}}\frac{\Lambda_{17}}{m_b}$$

where

$$\Lambda_{17} = e_c \operatorname{Re} \int_{-\infty}^{\infty} \frac{d\omega_1}{\omega_1} \left[ 1 - \underbrace{F\left( \frac{m_c^2 - i\varepsilon}{m_b \omega_1} \right)}_{\operatorname{perturbative}} + \frac{m_b \omega_1}{12 m_c^2} \right] \underbrace{h_{17}\left(\omega_1\right)}_{\operatorname{non-perturbative}}$$

- Need a new model for h<sub>17</sub> to reduce the error
  - New information on moments of  $h_{17}$ : constrain **new model**
  - What can we learn from moments?

•  $h_{17}$  can be thought of as a gluon PDF of a B meson

- $h_{17}$  can be thought of as a gluon PDF of a B meson
  - Non-local operator matrix element
  - Describe the hadronic effects of the process

- $h_{17}$  can be thought of as a gluon PDF of a B meson
  - Non-local operator matrix element
  - Describe the hadronic effects of the process

$$\begin{split} &h_{17}(\omega_1) = \\ &= \int \frac{dr}{2\pi} e^{-i\omega_1 r} \frac{\langle \bar{B}|(\bar{h}S_{\bar{n}})(0)\not n(1+\gamma_5)i\gamma^{\perp}\bar{n}_{\beta}(S_{\bar{n}}gG^{\alpha\beta}S_{\bar{n}})(r\bar{n})(S_{\bar{n}}^{\dagger}h)(0)|\bar{B}\rangle}{2M_B} \end{split}$$

-  $S_n(x)$  is the Wilson line

- $h_{17}$  can be thought of as a gluon PDF of a B meson
  - Non-local operator matrix element
  - Describe the hadronic effects of the process

$$\begin{split} &h_{17}(\omega_1) = \\ &= \int \frac{dr}{2\pi} e^{-i\omega_1 r} \frac{\langle \bar{B}|(\bar{h}S_{\bar{n}})(0)\not n(1+\gamma_5)i\gamma^{\perp}\bar{n}_{\beta}(S_{\bar{n}}gG^{\alpha\beta}S_{\bar{n}})(r\bar{n})(S_{\bar{n}}^{\dagger}h)(0)|\bar{B}\rangle}{2M_B} \end{split}$$

-  $S_n(x)$  is the Wilson line

$$S_n(x) = \mathbf{P} \exp \left( ig \int_{-\infty}^0 du n \cdot A_s(x + un) \right)$$

$$n^{\mu}\equiv (1,0,0,1)$$
 and  $\overline{n}^{\mu}\equiv (1,0,0,-1)$ 

• k th moment of  $h_{17}$ ; Obtained using  $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$ 

• k th moment of  $h_{17}$ ; Obtained using  $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$ 

$$\langle \omega_{1}^{k} \textit{h}_{17} \rangle = (-1)^{k} \frac{1}{2 \textit{M}_{\textit{B}}} \left\langle \overline{\textit{B}} \left| \left( \overline{\textit{h}} \textit{S}_{\bar{\textit{n}}} \right) (0) \cdots \left[ \left( i \overline{\textit{n}} \cdot \partial \right)^{k} \right] \left( \textit{S}_{\bar{\textit{n}}}^{\dagger} \textit{g} \textit{G}_{\textit{s}}^{\alpha \beta} \textit{S}_{\bar{\textit{n}}} \right) (r \overline{\textit{n}}) \left( \textit{S}_{\bar{\textit{n}}}^{\dagger} \textit{h} \right) (0) \right| \overline{\textit{B}} \right\rangle \right|_{r=0}$$

• k th moment of  $h_{17}$ ; Obtained using  $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$ 

$$\langle \omega_{1}^{k} h_{17} \rangle = (-1)^{k} \frac{1}{2M_{B}} \left\langle \overline{B} \left| \left( \overline{h} S_{\overline{n}} \right) (0) \cdots \overline{\left( i \overline{n} \cdot \partial \right)^{k}} \right| \left( S_{\overline{n}}^{\dagger} g G_{s}^{\alpha \beta} S_{\overline{n}} \right) (r \overline{n}) \left( S_{\overline{n}}^{\dagger} h \right) (0) \right| \overline{B} \right\rangle \bigg|_{r=0}$$

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

• k th moment of  $h_{17}$ ; Obtained using  $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$ 

$$\langle \omega_{1}^{k} h_{17} \rangle = (-1)^{k} \frac{1}{2M_{B}} \left\langle \overline{B} \left| \left( \overline{h} S_{\overline{n}} \right) (0) \cdots \overline{\left( i \overline{n} \cdot \partial \right)^{k}} \right| \left( S_{\overline{n}}^{\dagger} g G_{s}^{\alpha \beta} S_{\overline{n}} \right) (r \overline{n}) \left( S_{\overline{n}}^{\dagger} h \right) (0) \right| \overline{B} \right\rangle \bigg|_{r=0}$$

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

• k th moment of  $h_{17}$ ; Obtained using  $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$ 

$$\langle \omega_{1}^{k} h_{17} \rangle = (-1)^{k} \frac{1}{2M_{B}} \left\langle \overline{B} \left| \left( \overline{h} S_{\overline{n}} \right) (0) \cdots \overline{\left( i \overline{n} \cdot \partial \right)^{k}} \right| \left( S_{\overline{n}}^{\dagger} g G_{s}^{\alpha \beta} S_{\overline{n}} \right) (r \overline{n}) \left( S_{\overline{n}}^{\dagger} h \right) (0) \right| \overline{B} \right\rangle \bigg|_{r=0}$$

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

- Apply this for k derivatives  $\Rightarrow k$  commutators of  $i\bar{n} \cdot D$
- $[iD^{\mu}, iD^{\nu}] = igG^{\mu\nu}$

• k th moment of  $h_{17}$ ; Obtained using  $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$ 

$$\langle \omega_{1}^{k} h_{17} \rangle = (-1)^{k} \frac{1}{2M_{B}} \left\langle \overline{B} \left| \left( \overline{h} S_{\overline{n}} \right) (0) \cdots \overline{\left( i \overline{n} \cdot \partial \right)^{k}} \right| \left( S_{\overline{n}}^{\dagger} g G_{s}^{\alpha \beta} S_{\overline{n}} \right) (r \overline{n}) \left( S_{\overline{n}}^{\dagger} h \right) (0) \right| \overline{B} \right\rangle \bigg|_{r=0}$$

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

- Apply this for k derivatives  $\Rightarrow k$  commutators of  $i\bar{n} \cdot D$
- $[iD^{\mu}, iD^{\nu}] = igG^{\mu\nu}$
- New result Moments over  $\omega_1$

$$\langle \omega_1^k h_{17} \rangle = (-1)^k \frac{1}{2M_B} \langle \bar{B} | \bar{h} \cdots \underbrace{[i\bar{n} \cdot D, [i\bar{n} \cdot D, \cdots [i\bar{n} \cdot D, D]] D^{\alpha}, i\bar{n} \cdot D] \cdots]}_{k \text{ times}} ]s^{\lambda} h | \bar{B} \rangle$$

## Moments of the $g_{17}$

 Procedure to obtain these HQET matrix elements derived in [A. Gunawardana and G. Paz, JHEP 07(2017)137 [arXiv:1702.08904]]

$$\langle h_{17} 
angle = 2\lambda_2 = 2\mu_G^2/3$$
  $\langle \omega_1^2 h_{17} 
angle = rac{2}{15} \left(5m_5 + 3m_6 - 2m_9
ight)$  New result

m<sub>i</sub> were extracted from data for the first time in 2016
 [P. Gambino, K. J Healey, S. Turczyk PLB 763, 60 (2016)]

$$\mu_G^2 = 0.355 \pm 0.060 \text{ GeV}^2$$
  $m_5 = 0.072 \pm 0.045 \text{ GeV}^4$   $m_6 = 0.060 \pm 0.164 \text{ GeV}^4$   $m_9 = -0.280 \pm 0.352 \text{ GeV}^4$ 

• Relative errors are large:

• Relative errors are large: Numerical error is 17% for  $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for  $\langle \omega_1^2 h_{17} \rangle$ 

- Relative errors are large: Numerical error is 17% for  $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for  $\langle \omega_1^2 h_{17} \rangle$
- These moments still give useful information

- Relative errors are large: Numerical error is 17% for  $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for  $\langle \omega_1^2 h_{17} \rangle$
- These moments still give useful information
  - 2019 estimate  $\langle \omega_1^2 h_{17} \rangle \in (0.03, 0.27) \; \text{GeV}^4$
  - 2010 models provide  $\langle \omega_1^2 h_{17} \rangle \in (-0.31, 0.49) \text{ GeV}^4$ .
  - These older models were constructed before  $m_i$  were extracted
  - New estimate is significantly smaller than old estimate.

- Relative errors are large: Numerical error is 17% for  $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for  $\langle \omega_1^2 h_{17} \rangle$
- These moments still give useful information
  - 2019 estimate  $\langle \omega_1^2 h_{17} \rangle \in (0.03, 0.27) \text{ GeV}^4$
  - 2010 models provide  $\langle \omega_1^2 h_{17} \rangle \in (-0.31, 0.49) \text{ GeV}^4$ .
  - These older models were constructed before  $m_i$  were extracted
  - New estimate is significantly smaller than old estimate.
- Expect in future
  - Further improvements on HQET matrix elements
  - Belle II or LQCD data ⇒ Better constrains on moments

# **Applications**

• Properties of  $h_{17}$ 

- Properties of h<sub>17</sub>
  - Real and even function over  $\omega_1$
  - $\langle \omega_1^k h_{17}(\omega_1) \rangle = 0$  for  $k = 1, 3, 5, \cdots$
  - $h_{17}$  has a dimension of mass
  - Range of  $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$

- Properties of h<sub>17</sub>
  - Real and even function over  $\omega_1$
  - $\langle \omega_1^k h_{17}(\omega_1) \rangle = 0$  for  $k=1,3,5,\cdots$
  - $h_{17}$  has a dimension of mass
  - Range of  $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$
- We use Hermite polynomials  $H_n(x)$

- Properties of h<sub>17</sub>
  - Real and even function over  $\omega_1$
  - $-\langle \omega_1^k h_{17}(\omega_1) \rangle = 0 \text{ for } k = 1, 3, 5, \cdots$
  - h<sub>17</sub> has a dimension of mass
  - Range of  $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$
- We use Hermite polynomials  $H_n(x)$
- Our model:  $h_{17}(\omega_1) = \sum_n a_{2n} H_{2n}(\frac{\omega_1}{\sqrt{2}\sigma}) e^{\frac{-\omega_1^2}{2\sigma}}$

- where

$$a_0 = rac{\left\langle \omega_1^0 h_{17} 
ight
angle}{\sqrt{2\pi} |\sigma|}, \quad a_2 = rac{\left\langle \omega_1^2 h_{17} 
ight
angle - \sigma^2 \left\langle \omega_1^0 h_{17} 
ight
angle}{4\sqrt{2\pi} |\sigma|^3}, \quad a_4 = \cdots$$

- Properties of h<sub>17</sub>
  - Real and even function over  $\omega_1$
  - $-\langle \omega_1^k h_{17}(\omega_1) \rangle = 0 \text{ for } k = 1, 3, 5, \cdots$
  - h<sub>17</sub> has a dimension of mass
  - Range of  $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$
- We use Hermite polynomials  $H_n(x)$
- Our model:  $h_{17}(\omega_1) = \sum_n a_{2n} H_{2n}(\frac{\omega_1}{\sqrt{2}\sigma}) e^{\frac{-\omega_1^2}{2\sigma}}$ 
  - where

$$a_0=rac{\left\langle \omega_1^0 h_{17} 
ight
angle}{\sqrt{2\pi}|\sigma|}, \quad a_2=rac{\left\langle \omega_1^2 h_{17} 
ight
angle -\sigma^2 \left\langle \omega_1^0 h_{17} 
ight
angle}{4\sqrt{2\pi}|\sigma|^3}, \quad a_4=\cdots$$

 $|h_{17}| < 1$  GeV and no peaks beyond  $\omega_1 = 1$  GeV

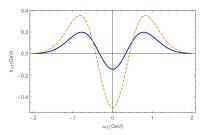


Figure: 2019 model vs 2010 model for  $h_{17}$ 

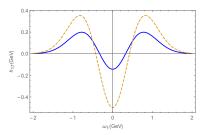


Figure: 2019 model vs 2010 model for  $h_{17}$ 

• Orange dashed line: 2010 model  $h_{17}\left(\omega_1,\mu\right)=\frac{2\lambda_2}{\sqrt{2\pi}\sigma}\frac{\omega_1^2-\Lambda^2}{\sigma^2-\Lambda^2}e^{-\frac{\omega_1^2}{2\sigma^2}}$ 

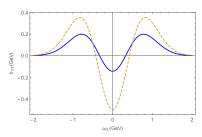


Figure: 2019 model vs 2010 model for  $h_{17}$ 

- Orange dashed line: 2010 model  $h_{17}(\omega_1,\mu) = \frac{2\lambda_2}{\sqrt{2\pi}\sigma} \frac{\omega_1^2 \Lambda^2}{\sigma^2 \Lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$ 
  - $\sigma = 0.5 \text{ GeV}, \Lambda = 0.425 \text{ GeV} \text{ and } \Rightarrow \langle \omega_1^2 h_{17} \rangle = 0.49 \text{ GeV}^4$

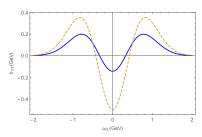


Figure: 2019 model vs 2010 model for  $h_{17}$ 

- Orange dashed line: 2010 model  $h_{17}(\omega_1, \mu) = \frac{2\lambda_2}{\sqrt{2\pi}\sigma} \frac{\omega_1^2 \Lambda^2}{\sigma^2 \Lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$ 
  - $\sigma=0.5~{\sf GeV}, \Lambda=0.425~{\sf GeV}$  and  $\Rightarrow \langle \omega_1^2 h_{17} \rangle=0.49~{\sf GeV}^4$
- Blue line; 2019 model:  $\sigma=0.5$  GeV and  $\langle \omega_1^2 h_{17} \rangle=0.27$  GeV<sup>4</sup>

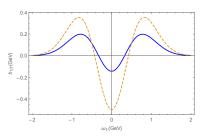


Figure: 2019 model vs 2010 model for  $h_{17}$ 

- Orange dashed line: 2010 model  $h_{17}(\omega_1, \mu) = \frac{2\lambda_2}{\sqrt{2\pi}\sigma} \frac{\omega_1^2 \Lambda^2}{\sigma^2 \Lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$ 
  - $\sigma=0.5~{\sf GeV}, \Lambda=0.425~{\sf GeV}$  and  $\Rightarrow \langle \omega_1^2 h_{17} \rangle=0.49~{\sf GeV}^4$
- Blue line; 2019 model:  $\sigma=0.5$  GeV and  $\langle \omega_1^2 h_{17} \rangle=0.27$  GeV<sup>4</sup>
- New function is 50% smaller than the 2010

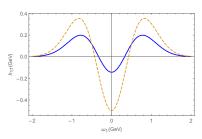


Figure: 2019 model vs 2010 model for  $h_{17}$ 

- Orange dashed line: 2010 model  $h_{17}\left(\omega_1,\mu\right)=rac{2\lambda_2}{\sqrt{2\pi}\sigma}rac{\omega_1^2-\Lambda^2}{\sigma^2-\Lambda^2}e^{-rac{\omega_1^2}{2\sigma^2}}$ 
  - $\sigma=0.5$  GeV,  $\Lambda=0.425$  GeV and  $\Rightarrow \langle \omega_1^2 h_{17} \rangle=0.49$  GeV<sup>4</sup>
- Blue line; 2019 model:  $\sigma=0.5$  GeV and  $\langle \omega_1^2 h_{17} \rangle=0.27$  GeV<sup>4</sup>
- New function is 50% smaller than the 2010
  - New model give better constraints on  $Q_1^c-Q_{7\gamma}$  contribution

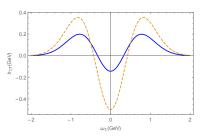


Figure: 2019 model vs 2010 model for  $h_{17}$ 

- Orange dashed line: 2010 model  $h_{17}\left(\omega_1,\mu\right)=rac{2\lambda_2}{\sqrt{2\pi}\sigma}rac{\omega_1^2-\Lambda^2}{\sigma^2-\Lambda^2}e^{-rac{\omega_1^2}{2\sigma^2}}$ 
  - $\sigma = 0.5 \text{ GeV}, \Lambda = 0.425 \text{ GeV} \text{ and } \Rightarrow \langle \omega_1^2 h_{17} \rangle = 0.49 \text{ GeV}^4$
- Blue line; 2019 model:  $\sigma=0.5$  GeV and  $\langle \omega_1^2 h_{17} \rangle=0.27$  GeV<sup>4</sup>
- New function is 50% smaller than the 2010
  - New model give better constraints on  $\mathit{Q}_{1}^{c}-\mathit{Q}_{7\gamma}$  contribution
- Consider also unknown higher moments, up to 6 Hermite polynomials

• Direct CP Asymmetry experimental bound:

$$A_{CP} = (1.5 \pm 2.0) \%$$

[ Y. Amhis et. al. EPJC 77, 895 (2017)]

Direct CP Asymmetry experimental bound:

$$A_{CP} = (1.5 \pm 2.0) \%$$

[ Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\begin{split} \mathcal{A}_{X_s\gamma}^{\mathrm{SM}} &= \left(1.15 \times \frac{\tilde{\Lambda}_{17}^u - \tilde{\Lambda}_{17}^c}{300 \mathrm{MeV}} + 0.71\right)\% \text{ CP asymmetry} \\ \tilde{\Lambda}_{17}^u &= \frac{2}{3} h_{17}(0) \\ \tilde{\Lambda}_{17}^c &= \frac{2}{3} \int_{4m_c^2/m_b}^{\infty} \frac{d\omega_1}{\omega_1} \underbrace{f\left(\frac{m_c^2}{m_b\omega_1}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_1)}_{\text{Non-perturbative}} \end{split}$$

Direct CP Asymmetry experimental bound:

$$A_{CP} = (1.5 \pm 2.0) \%$$

[ Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\begin{split} \mathcal{A}_{X_s\gamma}^{\mathrm{SM}} &= \left(1.15 \times \frac{\tilde{\Lambda}_{17}^u - \tilde{\Lambda}_{17}^c}{300 \mathrm{MeV}} + 0.71\right) \% \text{ CP asymmetry} \\ \tilde{\Lambda}_{17}^u &= \frac{2}{3} h_{17}(0) \\ \tilde{\Lambda}_{17}^c &= \frac{2}{3} \int_{4m_c^2/m_b}^{\infty} \frac{d\omega_1}{\omega_1} \underbrace{f\left(\frac{m_c^2}{m_b\omega_1}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_1)}_{\text{Non-perturbative}} \end{split}$$

Previously known values:

• Direct CP Asymmetry experimental bound:

$$A_{CP} = (1.5 \pm 2.0) \%$$

[ Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\begin{split} \mathcal{A}_{X_s\gamma}^{\mathrm{SM}} &= \left(1.15 \times \frac{\tilde{\Lambda}_{17}^u - \tilde{\Lambda}_{17}^c}{300 \mathrm{MeV}} + 0.71\right) \% \text{ CP asymmetry} \\ \tilde{\Lambda}_{17}^u &= \frac{2}{3} h_{17}(0) \\ \tilde{\Lambda}_{17}^c &= \frac{2}{3} \int_{4m_c^2/m_b}^{\infty} \frac{d\omega_1}{\omega_1} \underbrace{f\left(\frac{m_c^2}{m_b\omega_1}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_1)}_{\text{Non-perturbative}} \end{split}$$

Previously known values:

$$-330 {
m MeV} < \tilde{\Lambda}_{17}^{u} < +525 {
m MeV} -9 {
m MeV} < \tilde{\Lambda}_{17}^{c} < +11 {
m MeV}$$

[M. Benzke, S. J. Lee, M. Neubert and G. Paz PRL 106, 141801(2011)]

• Direct CP Asymmetry experimental bound:

$$A_{CP} = (1.5 \pm 2.0) \%$$

[ Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\begin{split} \mathcal{A}_{X_s\gamma}^{\mathrm{SM}} &= \left(1.15 \times \frac{\tilde{\Lambda}_{17}^u - \tilde{\Lambda}_{17}^c}{300 \mathrm{MeV}} + 0.71\right) \% \text{ CP asymmetry} \\ \tilde{\Lambda}_{17}^u &= \frac{2}{3} h_{17}(0) \\ \tilde{\Lambda}_{17}^c &= \frac{2}{3} \int_{4m_c^2/m_b}^{\infty} \frac{d\omega_1}{\omega_1} \underbrace{f\left(\frac{m_c^2}{m_b\omega_1}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_1)}_{\text{Non-perturbative}} \end{split}$$

Previously known values:

$$-330 {
m MeV} < \tilde{\Lambda}_{17}^{u} < +525 {
m MeV} -9 {
m MeV} < \tilde{\Lambda}_{17}^{c} < +11 {
m MeV}$$

[M. Benzke, S. J. Lee, M. Neubert and G. Paz PRL 106, 141801(2011)]

• We plan to improve these estimates

AYESH GUNAWARDANA Wayne State University Reeva

### Conclusion

- $ar{B} o X_s \gamma$  is a important New Physics probe
- Non perturbative error of the decay rate is 5%
- $Q_1^c-Q_{7\gamma}$  is the largest contribution to the error
- Better estimates for  $Q_1^c-Q_{7\gamma}$  obtained from moments of  $h_{17}$
- New estimates for CP asymmetry
- Reduce non-perturbative error on rate and CP asymmetry