

# Nucleon Binding Energy and Transverse Momentum Imbalance in Neutrino-Nucleus Reactions

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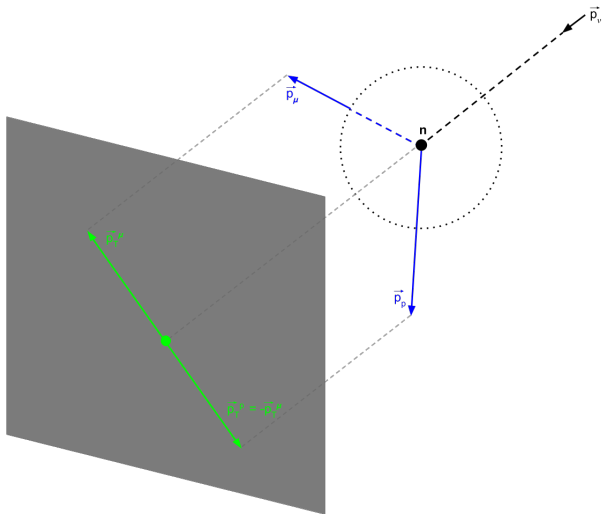
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University of Minnesota Duluth

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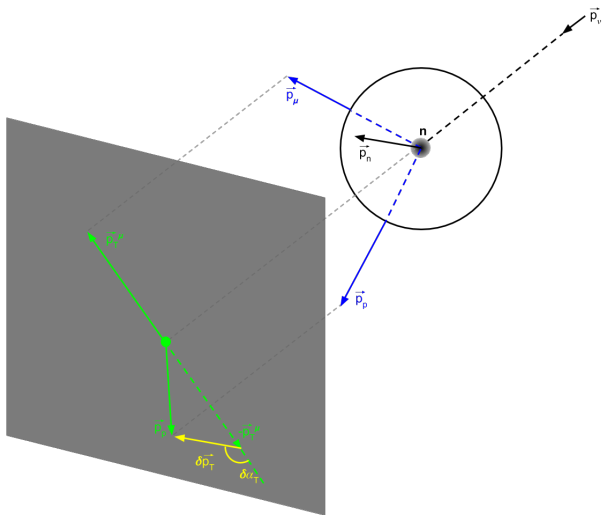
## Outline

- 1 The Single Transverse Kinematic Imbalances
- 2 Measurements of Single-TKI Projections
- 3 Results and Sensitivities to Binding Energy

We have developed a powerful tool to diagnose FSI



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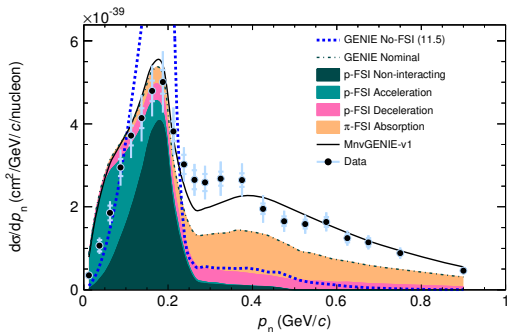


# Powerful tool to diagnose FSI

MINER $\nu$ A has had a fruitful journey on the transverse path:

Inferred initial neutron momentum:

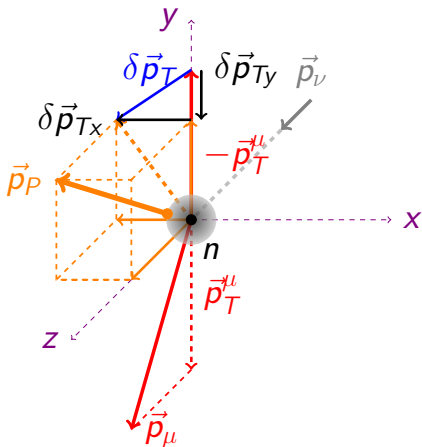
$$p_n = \sqrt{\delta p_T^2 + \delta p_L^2}$$



Pile up of MC due to “accelerated” proton after going through FSI at low  $p_n$  bins. This turns out to be a bug.

Phys.Rev.Lett. 121 (2018) no.2, 022504

## Single-TKI Transverse projections



We can project  $\delta p_T$  into components parallel:

$$\hat{y} = -\hat{p}_T^\mu$$

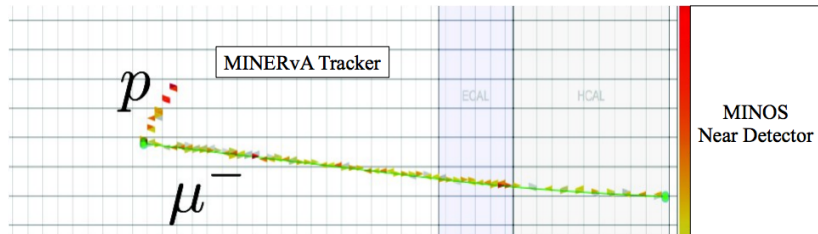
and perpendicular:

$$\hat{x} = \frac{\vec{p}_\nu \times \vec{p}_\mu}{|\vec{p}_\nu \times \vec{p}_\mu|}$$

to the muon transverse momentum  $\vec{p}_T^\mu$

## Signal Definition

The new variables are natural extensions of the Single-TKI measurements in Phys.Rev.Lett. 121 (2018) no.2, 022504.



- *CCQELike*: 1 muon, 0 mesons and at least 1 proton
- Muon and proton kinematic requirements to avoid regions that MINERνA has no acceptance:
  - $1.5 \text{ GeV}/c < p_\mu < 10 \text{ GeV}/c, \theta_\mu < 20^\circ$
  - $0.45 \text{ GeV}/c < p_p < 1.2 \text{ GeV}/c, \theta_p < 70^\circ$

## MC Setup

Cross section extracted with GENIE 2.8.4 and MnvGENIE-v1:

Expert level information for record only

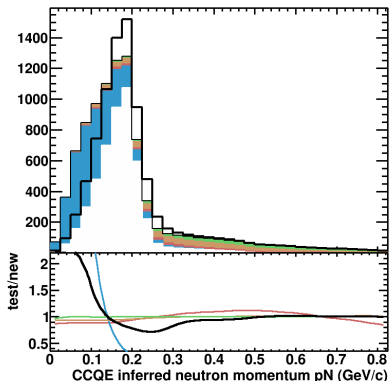
- FMBodekRitchie Fermi gas model, Rein-Seghal RES model, Bodek-Yang DIS parametrization and GENIE hA FSI model.
- Valencia 2p2h model with MINER $\nu$ A tune
- RPA modification to QE
- Non-resonant pion reduction to 43% of nominal

And finally, MnvGENIE-v1.0.1 removes the elastic FSI in GENIE to correct the GENIE FSI bug:

- removes elastic FSI,
- increases no-FSI by 50%,
- shown to effectively account for the GENIE FSI bug.



## A GENIE FSI Fix

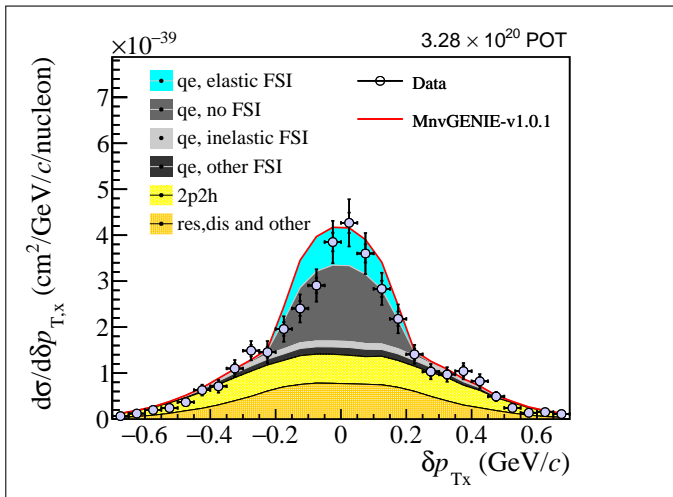


Elastic hadron-nucleus scattering in neutrino-nucleus reactions and transverse kinematics measurements (arXiv 1906.10576)

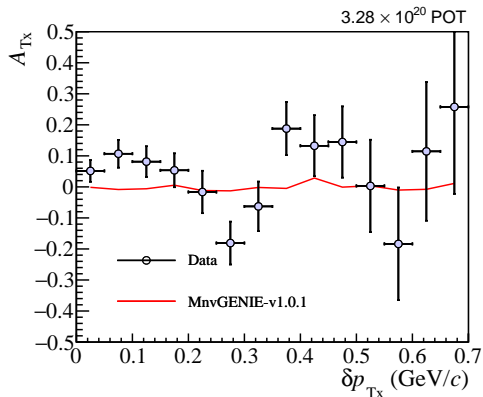
**Black:** a distribution with corrected elastic FSI

**Blue:** the bugged version of elastic FSI in GENIE

Replacing elastic FSI with no-FSI effectively reproduces the corrected distribution without having to reprocess the large number of existing MC.

$\delta p_{Tx}$ 

The distribution has hints of asymmetry and there are small structures at -0.3 and 0.35

$\delta p_{Tx}$  asymmetry

We calculate the asymmetry in each  $\delta p_{Tx}$  bin:

$$A_{Tx} = \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-},$$

for  $\delta p_{Tx} > 0$ .

And calculated  $1.3\sigma$  effect.

Figment of imagination or something real?

## Some old single-pion production results

Asymmetries in  $\Delta$  events exists. We have an irreducible pion absorption background.

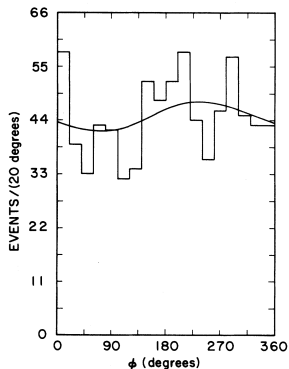
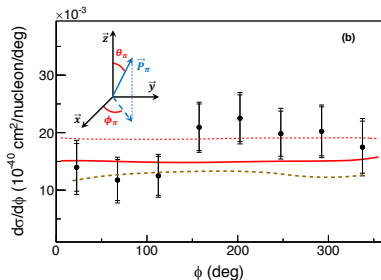
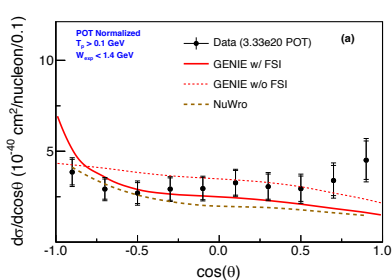


FIG. 16. Distribution of events in the pion azimuthal angle  $\phi$  for the final state  $\mu^- p \pi^+$ , with  $M(p \pi^+) < 1.4$  GeV. The curve is the area-normalized prediction of the Adler model.

Study of single-pion production by weak charged currents in low-energy  $\nu d$  interactions (Phys. Rev. D 26, 3297 (1982))

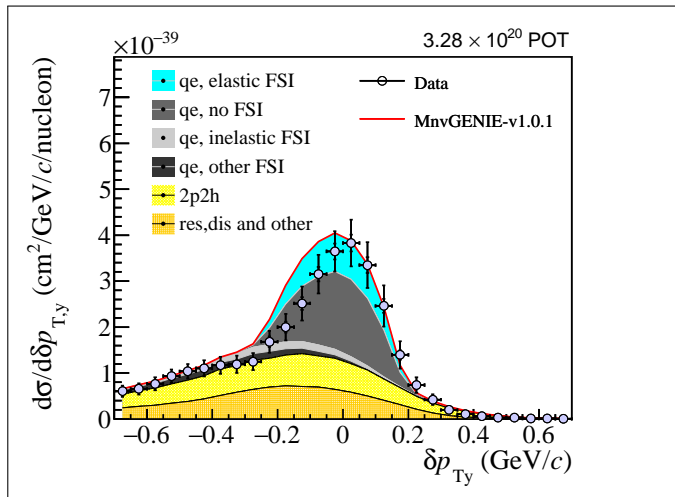
# MINER $\nu$ A's very own $\pi^0$ production results

Measurement of  $\nu_\mu$  charged-current single  $\pi^0$  production on hydrocarbon in the few-GeV region using MINER $\nu$ A (Phys.Rev. D96 (2017) no.7, 072003)



Future high-statistics Single-TKI measurements in the Medium Energy (ME) will shed light on the significance of the asymmetry.

It might be possible to constrain the fraction of pion absorption events with this asymmetry ( arXiv:1907.11212 )

$\delta p_{Ty}$ 

A clear peak shift.  $\delta p_{Ty}$  is sensitive to the binding energy implementation of the models.

## GENIE implementation and corrections

In order to investigate if the peak shift originate from mishandling of binding energy, we considered two adhoc corrections: **Correction common to both:**

- Add 25 MeV to outgoing proton. This is the removal energy arbitrarily subtracted from final state proton in Carbon.
- Subtract 10.1 MeV from the muon to account for the lack of excitation energy in particle generation.

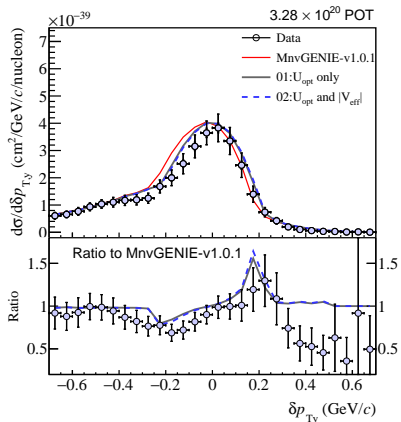
Then we have:

- 1  $U_{\text{opt}}$ : Subtract  $|U_{\text{opt}}|$  from the proton and add to the muon
- 2  $U_{\text{opt}}$  and  $V_{\text{eff}}$ : Also add  $|V_{\text{eff}}^{\text{P}}|$  to the proton and subtract  $|V_{\text{eff}}|$  from the muon

$U_{\text{opt}}$  was a fit obtained from *Eur.Phys.J. C79 (2019) no.4, 293* and averages to be 2.3 MeV due to proton minimum momentum cuts.

$$V_{\text{eff}} = V_{\text{eff}}^{\text{P}} = 3.1 \text{ MeV}$$

## Comparing GENIE corrections



Generator peaks better conform to data.

$U_{opt}$  and  $V_{eff}$  have small effects.

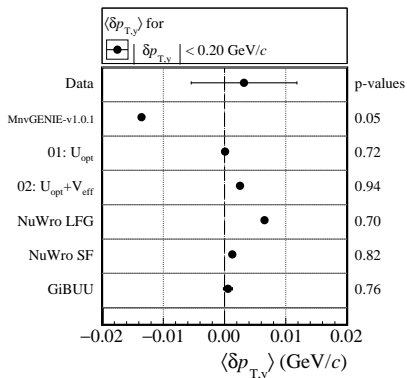
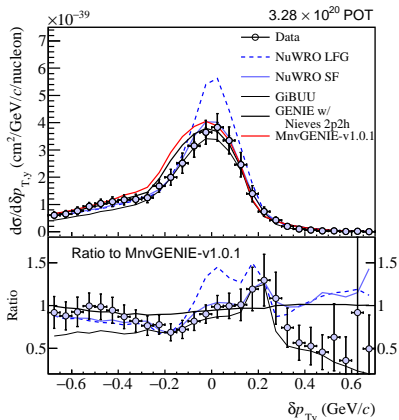
Undo the 25 MeV removal energy while simulate  $E_x$  is by far the most significant.

The take home message:

the adhoc corrections produce more sensible peak positions. We need more careful treatment of the GENIE binding energy.



## Model comparisons



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the adhoc corrections produce more sensible peak positions. We need more careful treatment of the GENIE binding energy.

## Conclusion

- Extended the Single-TKI measurements to  $\delta p_{Tx}$  and  $\delta p_{Ty}$ .
- Applied corrections to GENIE FSI model.
- Hints of asymmetry in  $\delta p_{Tx}$  not modelled by current generators.
- $\delta p_{Ty}$  implies binding energy needs better treatment.

The Single-TKI variables can be measured for other elements such as Argon, providing constraints on nuclear effects such as the binding energy for current and future experiments.

Future MINER $\nu$ A measurements of the Single-TKI variables in the high-statistics ME samples will further our understanding of nuclear effects and aid in precise nuclear modelling.

# Backup

## A simple case without FSI

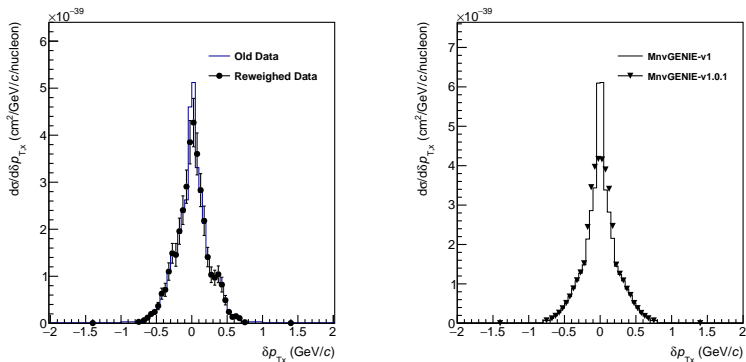
For a bound nucleon scattering without FSI:

$$\begin{pmatrix} E_\nu \\ 0 \\ 0 \\ p_\nu \end{pmatrix} + \begin{pmatrix} M_N - \epsilon^N \\ k_x \\ k_y \\ k_z \end{pmatrix} = \begin{pmatrix} E_{l'} \\ 0 \\ p_{l'y} \\ p_{l'z} \end{pmatrix} + \begin{pmatrix} E_p \\ p_{px} \\ p_{py} \\ p_{pz} \end{pmatrix} \quad (1)$$

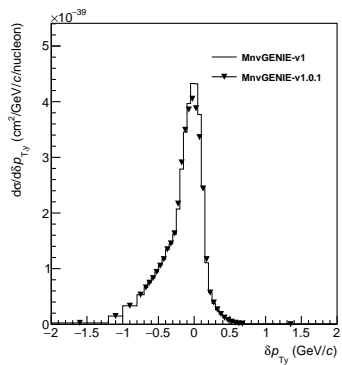
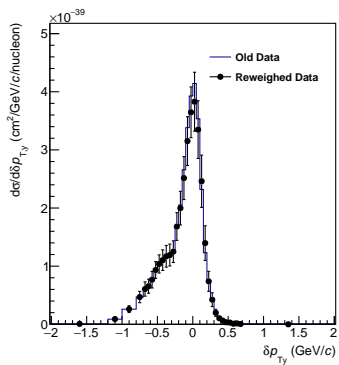
and rearranging:

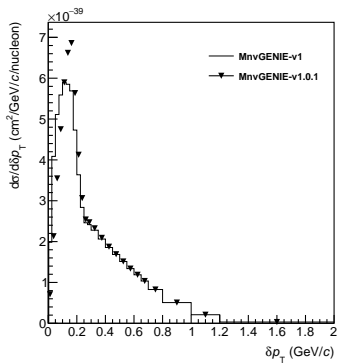
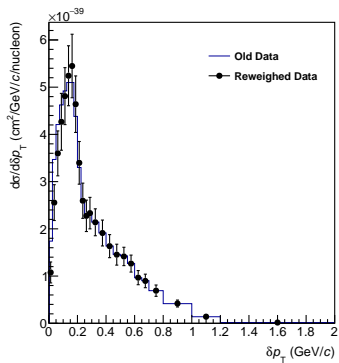
$$\begin{pmatrix} \nu \\ 0 \\ q_T \\ q_l \end{pmatrix} + \begin{pmatrix} M_N - \epsilon^N \\ k_x \\ k_y \\ k_z \end{pmatrix} = \begin{pmatrix} E_p \\ p_{px} = \delta p_{Tx} \\ q_T + \delta p_{Ty} \\ p_{pz} \end{pmatrix} \quad (2)$$

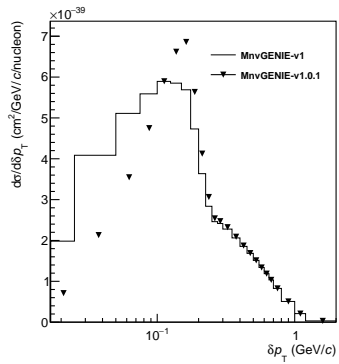
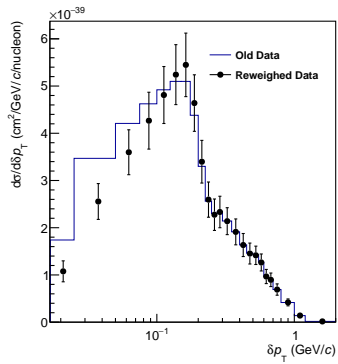
$\delta p_{Tx}$  is precisely the measure of Fermi motion, while the effects of  $\epsilon$  occurs in y and z axis only.



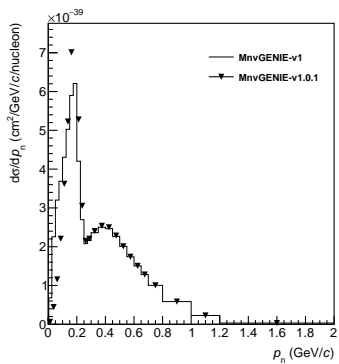
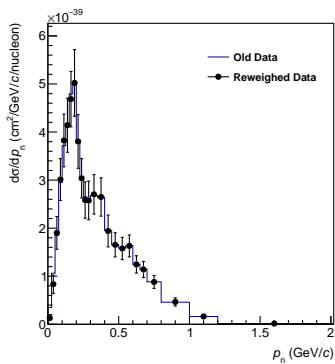
**Figure:**  $\delta p_{Tx}$ , data on the left, MC on the right. Elastic FSI preferentially produce nucleons close to the reaction plane, therefore producing an unphysical peak near 0. The updated MC does not experience the bug and therefore produces a smoother distribution. There is a corresponding reduction in the peak in data as a result of increased MC efficiencies at the peak

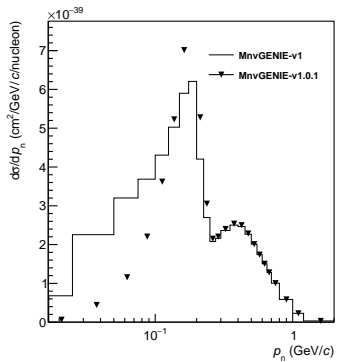
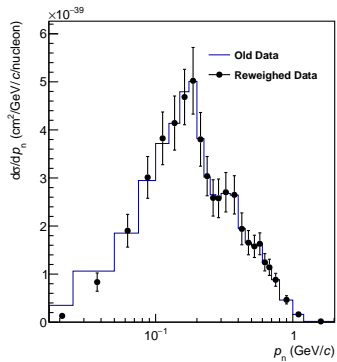
Figure:  $\delta p_{Ty}$

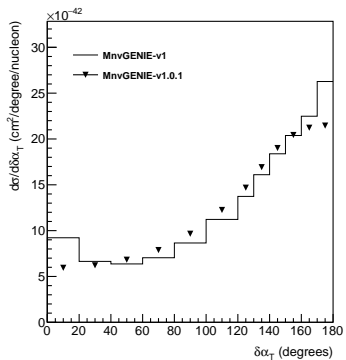
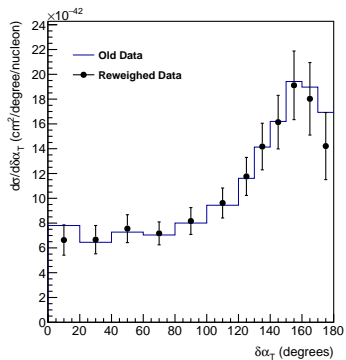
Figure:  $\delta p_T$

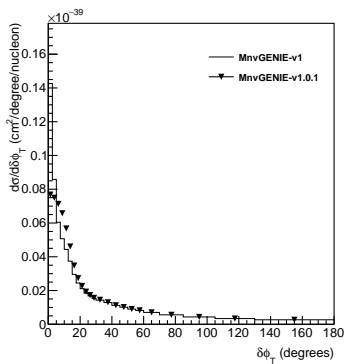
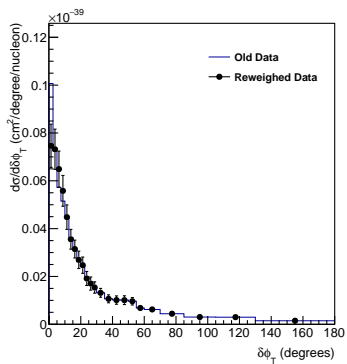
Figure:  $\delta p_T$

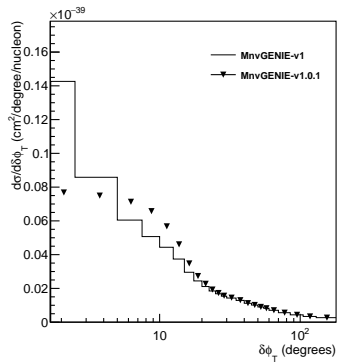
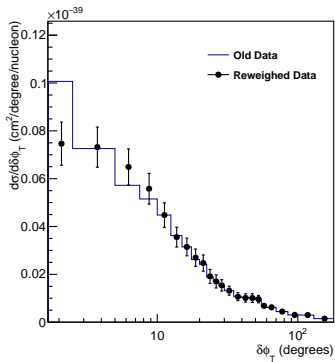


Figure:  $p_n$

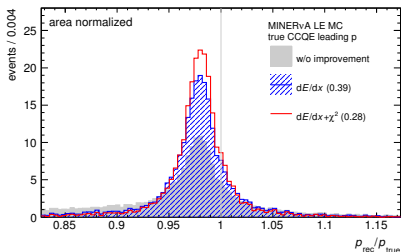
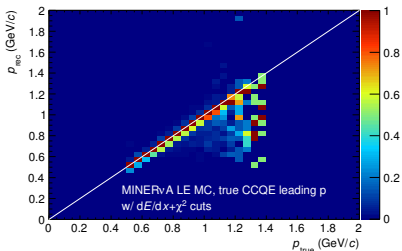
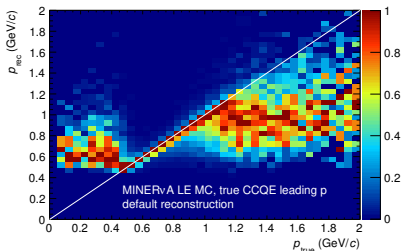
Figure:  $p_n$

Figure:  $\delta\alpha_T$

Figure:  $\delta\phi_T$

Figure:  $\delta\phi_T$

## Event Selection



- Low Energy Neutrino datasets,  $3.28 \times 10^{20}$  proton on target.
- Proton  $dE/dX$  and  $\chi^2$  cuts performed to select protons with well reconstructed momentum

## Correction to GENIE

GENIE:

$$\nu_{GENIE} = \sqrt{M_p^2 + (\mathbf{k} + \mathbf{q})^2} + \frac{\mathbf{k}^2}{2M_{A-1}} - M_p + S^p \quad (3)$$

The corrected:

$$\nu_{corr} = \sqrt{M_p^2 + (\mathbf{k} + \mathbf{q})^2} - |U_{opt}| + \frac{\mathbf{k}^2}{2M_{A-1}^*} - M_p + S^p + E_x \quad (4)$$

We have:

$$\nu_{corr} = \nu_{GENIE} - |U_{opt}| + E_x \quad (5)$$

$$E_{\mu}^{GENIE} = E_{\mu}^{corr} - |U_{opt}| + E_x \quad (6)$$

$$E_{\mu}^{corr} = E_{\mu}^{GENIE} + |U_{opt}| - E_x \quad (7)$$

## Correction to GENIE

Hadronic Side:

$$E_p^{GENIE} = E'_p{}^{GENIE} - \Delta(25) \quad (8)$$

$$E'_p{}^{GENIE} = \sqrt{M_p^2 + (\mathbf{k} + \mathbf{q})^2} \quad (9)$$

$$E_p^{corr} = E'_p{}^{GENIE} - |U_{opt}| = E_p^{GENIE} + \Delta - |U_{opt}| \quad (10)$$

In the case of GENIE we need to put  $\Delta = 25$  to the generated proton.