

Axion-like dark matter search using ferromagnetic toroids

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BOSTON
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Northeastern University
Boston, MA

Ultralight axion dark matter

- Dark matter (DM) accounts for $\approx 1/4$ of the total energy density of the Universe
- The local DM energy density is

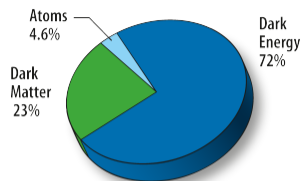
$$\rho_{\text{DM}} \approx 0.4 \text{ GeV}/\text{cm}^3$$

- QCD axions and axion-like particles are excellent DM candidates
- The number density of axions per de Broglie volume is large: $n_a/\lambda^3 \gg 1$
- As a result, axions form an oscillating classical field

$$a(t) = a_0 \sin(m_a t), \quad a_0 \approx \sqrt{2\rho_{\text{DM}}}/m_a \quad (\text{in natural units, i.e., } \hbar = c = 1)$$

- Coherence time of these oscillations is limited by the virial velocity $v_{\text{vir}} \approx 10^{-3}c$:

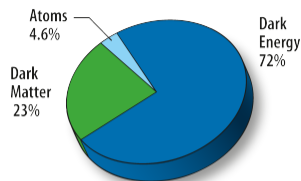
$$\frac{\Delta\omega_a}{m_a} \approx v_{\text{vir}}^2 \approx 10^{-6}, \quad \tau_c \approx 10^6 \frac{2\pi}{m_a}$$



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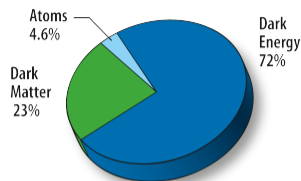
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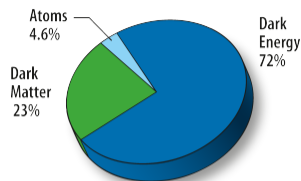
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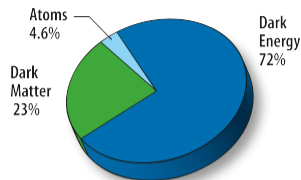
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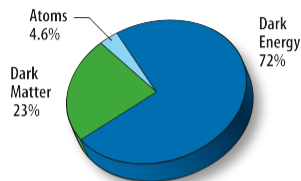
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- In the presence of a background axion field

$$a(t) = a_0 \sin(m_a t),$$

inhomogeneous Maxwell's equations take the form (in natural units, i.e., $\epsilon_0 = \mu_0 = 1$)

$$\nabla \cdot \mathbf{E} = \rho - g_{a\gamma\gamma} \nabla a \cdot \mathbf{B} \quad (\text{Gauss's law}),$$

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} + \mathbf{J} + g_{a\gamma\gamma} \left(\frac{\partial a}{\partial t} \mathbf{B} + \nabla a \times \mathbf{E} \right) \quad (\text{Ampère's law}).$$

- Under static magnetic field \mathbf{B}_0 , DM axions source an effective current density

$$\mathbf{J}_{\text{eff}} = g_{a\gamma\gamma} \frac{\partial a}{\partial t} \mathbf{B}_0 = g_{a\gamma\gamma} \sqrt{2\rho_{\text{DM}}} \mathbf{B}_0 \cos(m_a t).$$

P. Sikivie, PRL **51**, 1415 (1983)

P. Sikivie, N. Sullivan, D. B. Tanner, PRL **112**, 131301 (2014) — LC Circuit proposal

Axion electrodynamics

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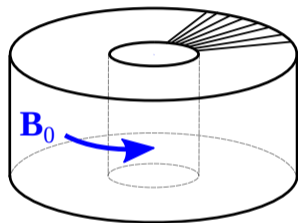
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- \mathbf{J}_{eff} generates an axial oscillating magnetic flux Φ_a
- Φ_a can be detected by a SQUID coupled to a pickup coil
- Similar to ABRACADABRA, but we use ferromagnetic core material to enhance \mathbf{B}_0 : $\mathbf{B}_0 = \mathbf{H}_0 + \mathbf{M}$ (natural units)
- We tried two core materials: gadolinium-iron garnet (GdIG) and Fe-Ni alloy powder

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S. Chaudhuri et al., Phys. Rev. D **92**, 075012 (2015) — DM Radio proposal

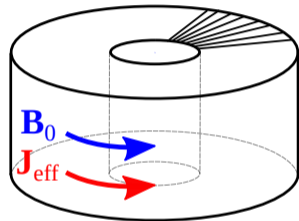
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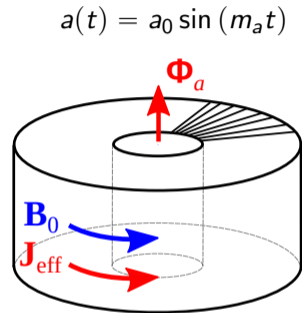
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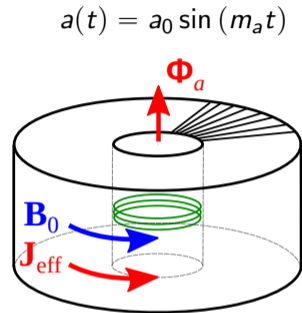
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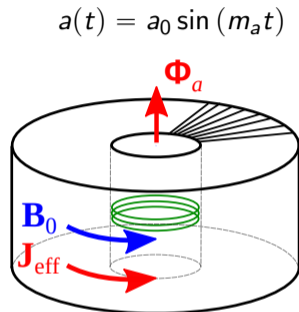
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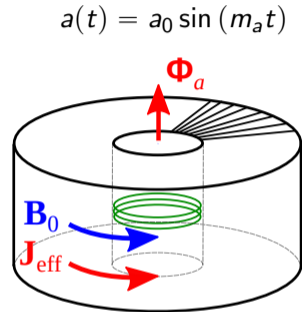
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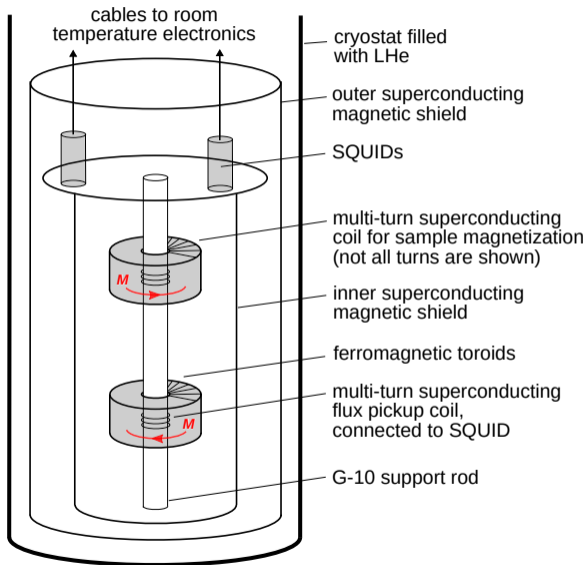
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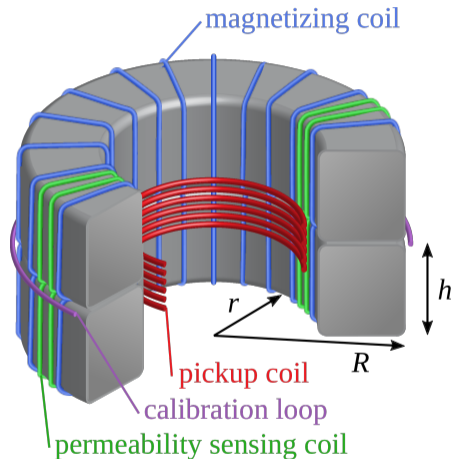
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Experimental apparatus



Fe-Ni alloy powder core toroids:



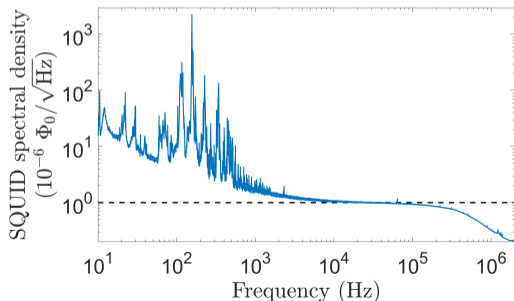
$$r = 24 \text{ mm}, R = 39 \text{ mm}, h = 16 \text{ mm}$$

SQUID magnetometers

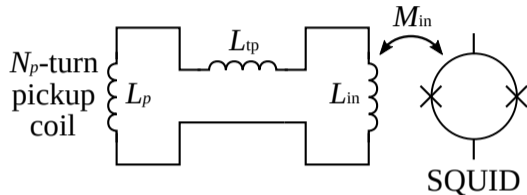
SQUIDS from Magnicon GmbH (Germany)



Flux noise at 4 K: $\approx 1 \mu\Phi_0/\sqrt{\text{Hz}}$



Broadband readout circuit

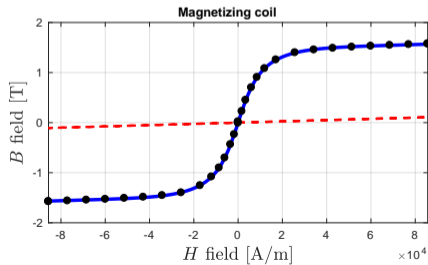
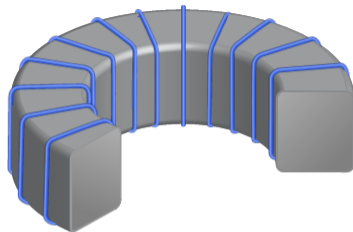
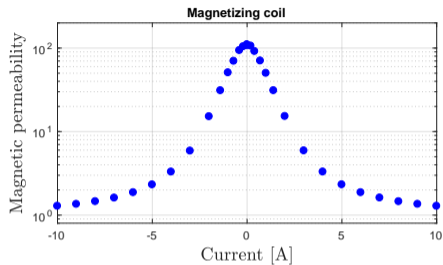


$$\Phi_{\text{SQUID}} = \frac{N_p M_{\text{in}}}{L_p + L_{\text{tp}} + L_{\text{in}}} \Phi_a$$

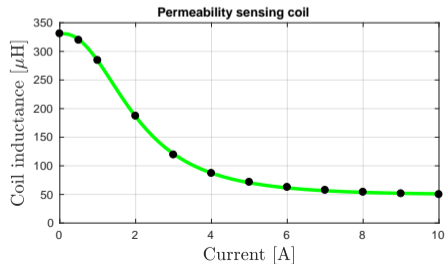
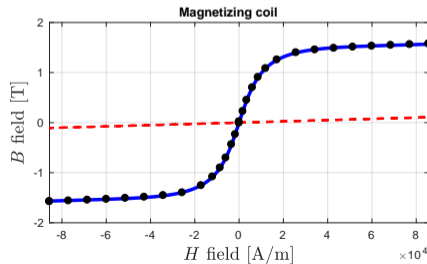
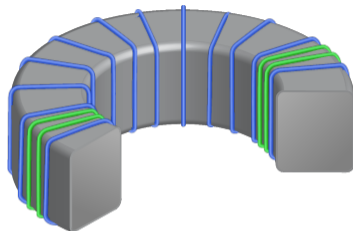
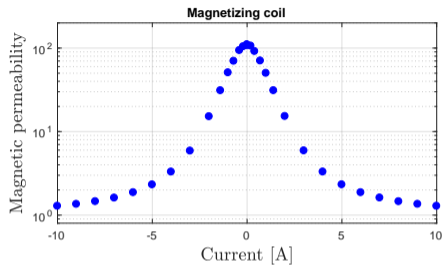
$$M_{\text{in}} = 9 \text{ nH}, L_{\text{in}} = 1.8 \mu\text{H}, L_p = 3 \mu\text{H}$$

Optimal number of turns: $N_p = 6$

Magnetization measurements



Magnetization measurements



Sensitivity scaling

- Axion flux scales with the static magnetic field and the toroid effective volume as

$$\Phi_a = g_{a\gamma\gamma} \sqrt{2\rho_{\text{DM}}} H_{\text{min}} V$$

- For 6 A current in the magnetizing coil:

$$H_{\text{min}} = 41 \text{ kA/m}, \quad V = 396 \text{ cm}^3$$

- A factor of 30 enhancement compared to an air-core toroid ($V = 13.4 \text{ cm}^3$)
- Sensitivity scales with the integration time, t , as

$$\text{sensitivity} \propto \begin{cases} \sqrt{t}, & \text{if } t \ll \tau_c \text{ (coherent averaging),} \\ \sqrt[4]{\tau_c t}, & \text{if } t \gg \tau_c \text{ (incoherent averaging),} \end{cases}$$

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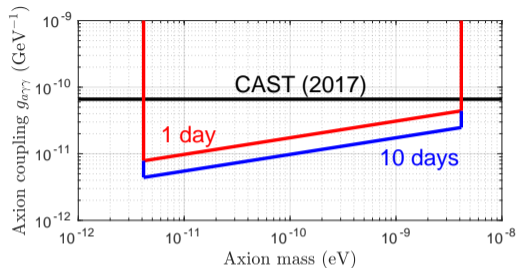
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Sensitivity reach of the experiment

$$\Phi_{\text{SQUID}} < \frac{1 \mu\Phi_0/\sqrt{\text{Hz}}}{\sqrt[4]{\tau_c t}}$$

$$g_{a\gamma\gamma} < \frac{1 \mu\Phi_0/\sqrt{\text{Hz}}}{\sqrt[4]{\tau_c t}} \times \frac{L_p + L_{\text{tp}} + L_{\text{in}}}{N_p M_{\text{in}}} \times \frac{1}{\sqrt{2\rho_{\text{DM}} H_{\text{min}} V}}$$



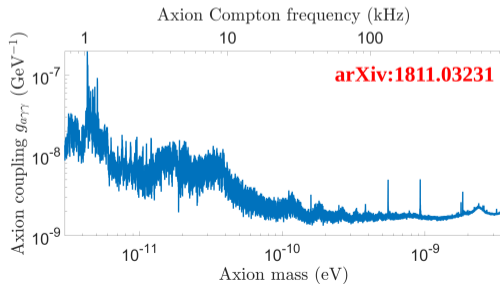
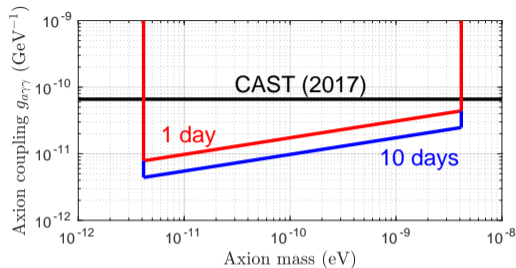
*V. Anastassopoulos et al. (CAST Collaboration),
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Summary and outlook

- Ultralight axion DM behaves like an oscillating classical field
- It generates an oscillating magnetic field in the presence of a static one
- Using SQUIDs we can search for axion DM in peV–neV mass range
- We enhance the static magnetic field with ferromagnetic toroidal cores
- Fe-Ni alloy powder cores provide a factor of 30 increase in sensitivity
- Projected sensitivity of the experiment surpasses the existing laboratory limits
- Work in progress: reduction of electromagnetic interference and data analysis

Summary and outlook

- Ultralight axion DM behaves like an oscillating classical field
- It generates an oscillating magnetic field in the presence of a static one
- Using SQUIDs we can search for axion DM in peV–neV mass range
- We enhance the static magnetic field with ferromagnetic toroidal cores
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Thank you for your attention!