

A Sub-GeV DM model

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Based on: Dutta, Ghosh, Kumar, 1905.02692

Introduction

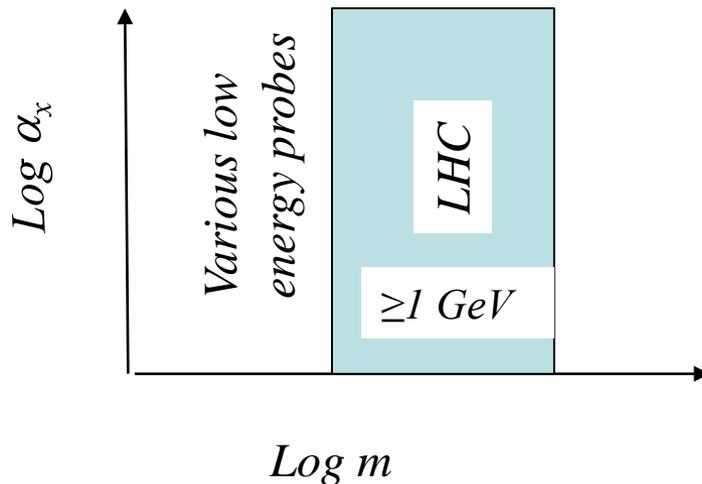
- ν sector of SM ($SU(3) \times SU(2) \times U(1)$) requires new physics for understanding the experimental results on masses and mixing angles

$$\text{tiny } \nu \text{ mass: } \frac{m_{\nu D}^2}{m_{Maj}}$$

- Similarly, DM explanation (with M_{DM} anywhere between 1 KeV to 100 TeV) requires new physics

- The new physics: new mass scales and new couplings

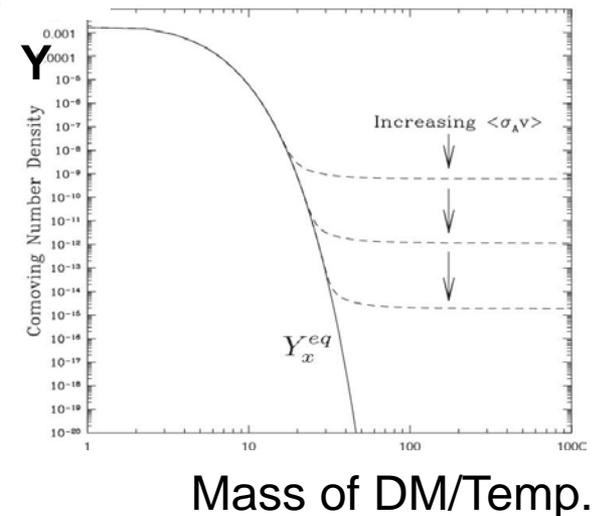
$$\alpha_x = \frac{g_x^2}{4\pi}$$



Thermal DM

$$H = n \langle \sigma v \rangle \quad \sigma \sim 1 pb$$

Y=DM abundance

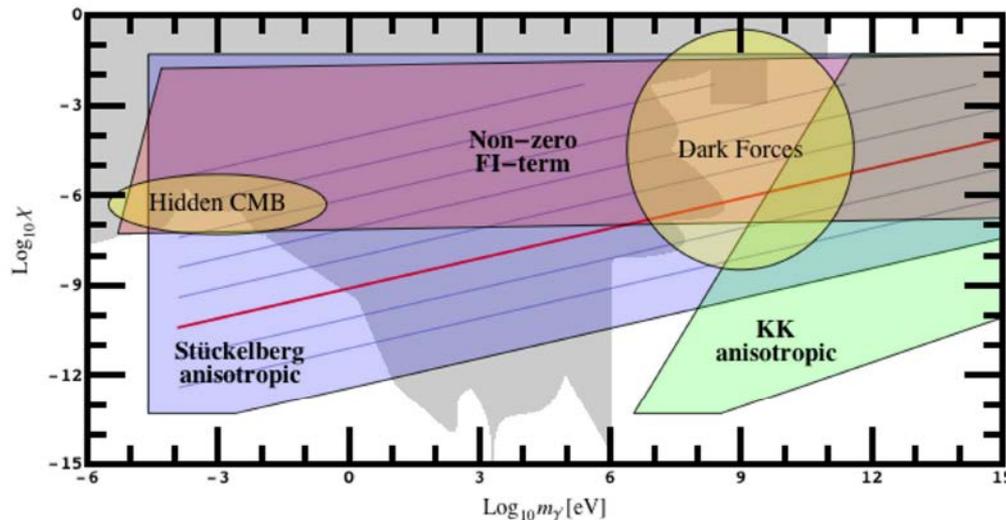


Low scales

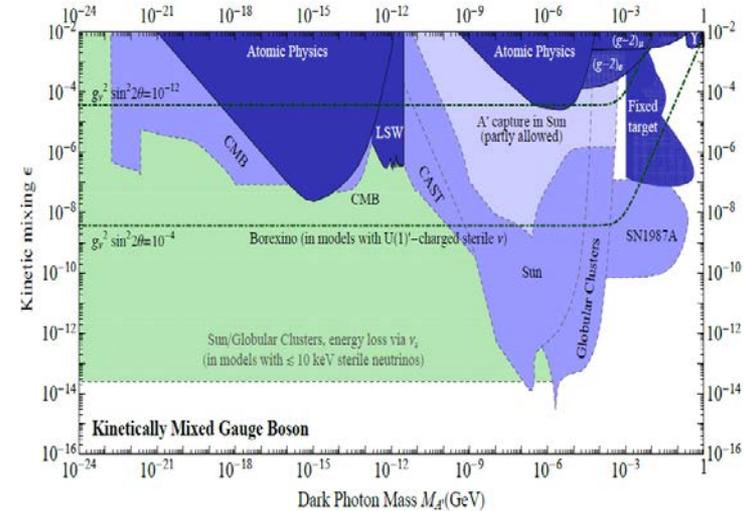
New physics: new symmetry breaking scale

Existence of new scales above or below the SM in many theories
Intermediate scale, GUT scale etc.

String theory: Many U(1) symmetry with the symmetry breaking scale can be anywhere



Cicoli, Goodsell, Jaeckel, Ringwald, 2011



Harnik, Kopp, Machado, 2012

Light mediator: Models

Various light mediators scenarios proposed:

- **Dark Matter scenarios based on hidden sectors:**
e.g., models of asymmetric DM,
Sommerfeld enhancements motivated by SIMP,
Decay of the observable sector DM into hidden sector
Essig et al, 2013 Review
- **$g-2$ of electron: 2.4σ discrepancy (recent)**
Marciano, Davoudiasl, 2018
- **Neutrino sector physics.**
New Neutrino interactions to satisfy MiniBooNE excess
Bertuzzo, Jana, Machado, Funchal, 2018
- **Solutions of Yukawa couplings hierarchies problem**
Dutta, Ghosh, Kumar, 2019

$$SU(2)_L \times U(1)_Y \times U(1)_{T3R}$$

Model for a sub GeV DM

*E.g., there may be a new symmetry breaking scale around GeV
 → 2nd and 1st generation fermion masses (~MeV to few GeV)*

Anomaly free

$$SU(2)_L \times U(1)_Y \times U(1)_{T3R}$$

field	q_{T3R}
q_R^u	-2
q_R^d	2
ℓ_R	2
ν_R	-2
η_L	1
η_R	-1
ϕ	-2

$U(1)_{T3R}$ is broken at 1-10 GeV down to Z_2

Low mass dark matter, gauge Boson, scalar

Predictions are testable at various low energy experiments

Dutta, Ghosh, Kumar, 2019

Similar model for with 3rd generation: Dutta, Kumar, 2011

U(1)_{T3R}

$$\mathcal{L}_{Yuk} = -\frac{\lambda_u}{\Lambda} \tilde{H} \phi^* \bar{Q}_L q_R^u - \frac{\lambda_d}{\Lambda} H \phi \bar{Q}_L q_R^d - \frac{\lambda_\nu}{\Lambda} \tilde{H} \phi^* \bar{L}_L \nu_R - \frac{\lambda_l}{\Lambda} H \phi \bar{L}_L \ell_R \\ - \lambda \phi \bar{\eta}_R \eta_L - \frac{1}{2} \lambda_L \phi \bar{\eta}_L^c \eta_L - \frac{1}{2} \lambda_R \phi^* \bar{\eta}_R^c \eta_R - \mu_\phi^2 \phi^* \phi - \lambda_\phi (\phi^* \phi)^2 + H.c.,$$

- **Scalar ϕ vev** $V = (-\mu_\phi^2/2\lambda_\phi)^{1/2}$ **breaks U(1)_{T3R} to Z₂**, vev is around **1-10 GeV** with $m_{\phi'} = 2\lambda_\phi^{1/2} V$.

Vector-like quarks with mass 1 TeV or higher can generate this Lagrangian

→ *Larger Yukawa couplings for first two generations are possible*

$$\mathcal{L}_{Yuk} = -m_u \bar{q}_L^u q_R^u - m_d \bar{q}_L^d q_R^d - m_\nu \bar{\nu}_L \nu_R - m_\ell \bar{\ell}_L \ell_R - \frac{1}{2} m_1 \bar{\eta}_1 \eta_1 - \frac{1}{2} m_2 \bar{\eta}_2 \eta_2 \\ - \frac{m_u}{V} \bar{q}_L^u q_R^u \phi' - \frac{m_d}{V} \bar{q}_L^d q_R^d \phi' - \frac{m_\nu D}{V} \bar{\nu}_L \nu_R \phi' - \frac{m_\ell}{V} \bar{\ell}_L \ell_R \phi' - \frac{1}{2} \frac{m_1}{V} \bar{\eta}_1 \eta_1 \phi' - \frac{1}{2} \frac{m_2}{V} \bar{\eta}_2 \eta_2 \phi' + \dots$$

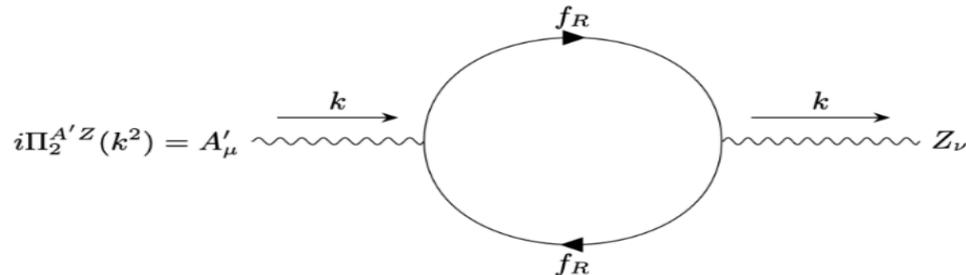
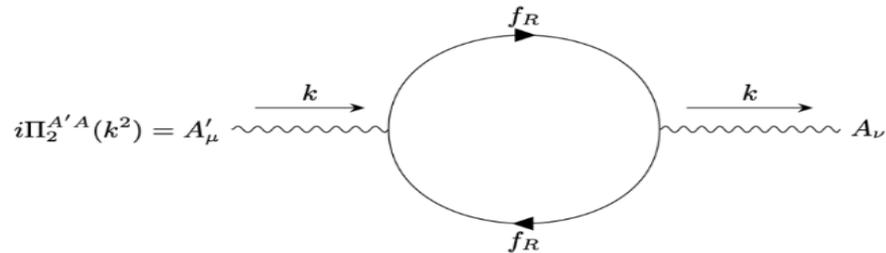
$$\eta_1 = -\frac{i}{\sqrt{2}} \begin{pmatrix} \eta_L - \eta_R^c \\ -\eta_L^c + \eta_R \end{pmatrix} \quad \eta_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} \eta_L + \eta_R^c \\ \eta_L^c + \eta_R \end{pmatrix}$$

Dark Matter (parity odd): $\eta_{1,2}$

U(1)_{T3R}

$$\mathcal{L}_{gauge} = \frac{i}{4} g_{T3R} A'_\mu (\bar{\eta}_1 \gamma^\mu \eta_2 - \bar{\eta}_2 \gamma^\mu \eta_1) + \frac{m_{A'}^2}{V} \phi' A'_\mu A'^\mu + i g_{T3R} A'_\mu (\phi' \partial^\mu \phi'^* - \phi'^* \partial^\mu \phi) - \frac{1}{2} g_{T3R} j_{A'}^\mu A'_\mu,$$

$$j_{A'}^\mu = \sum_f Q_{T3R}^f \bar{f} \gamma^\mu f. \quad m_{A'}^2 = 2g_{T3R}^2 V^2$$

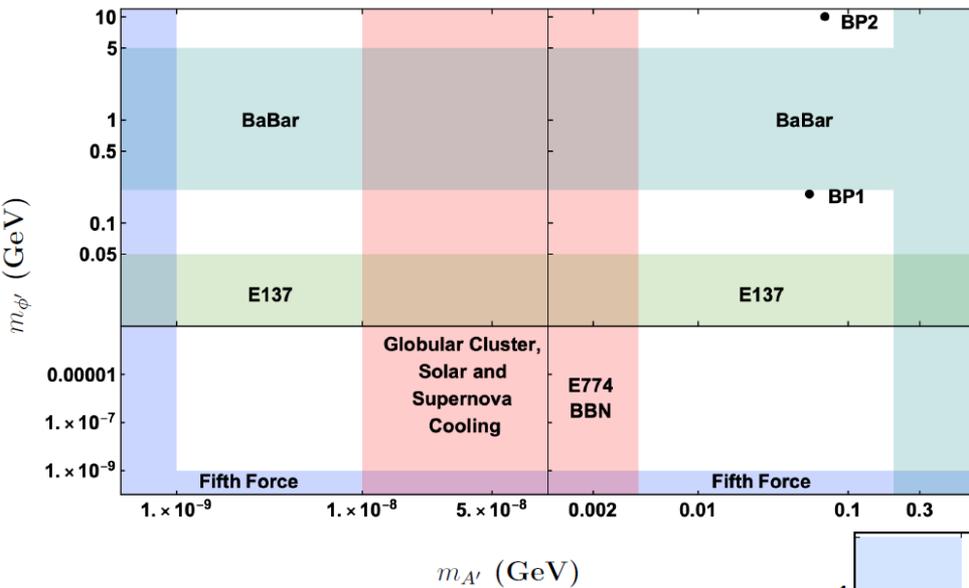


$U(1)_{T3R}$

- $\phi' : \phi' \rightarrow \bar{l}l, \nu_s \nu_A, \pi\pi, A' A' : \text{dominate, if kinematically allowed. Otherwise, } \phi' \rightarrow \gamma\gamma \text{ (one loop diagram) dominates}$
- $A' : A' \rightarrow \bar{l}l, \nu_s \nu_s, \pi\pi, \phi' \phi' : \text{dominate, if kinematically allowed. Otherwise, } A' \rightarrow \nu_L \nu_L \text{ (one loop diagram) dominates}$
- $\nu_s : \nu_s \rightarrow \nu_A \gamma\gamma : \text{mediated by an offshell } \phi' \text{ dominate}$

Parameter Space

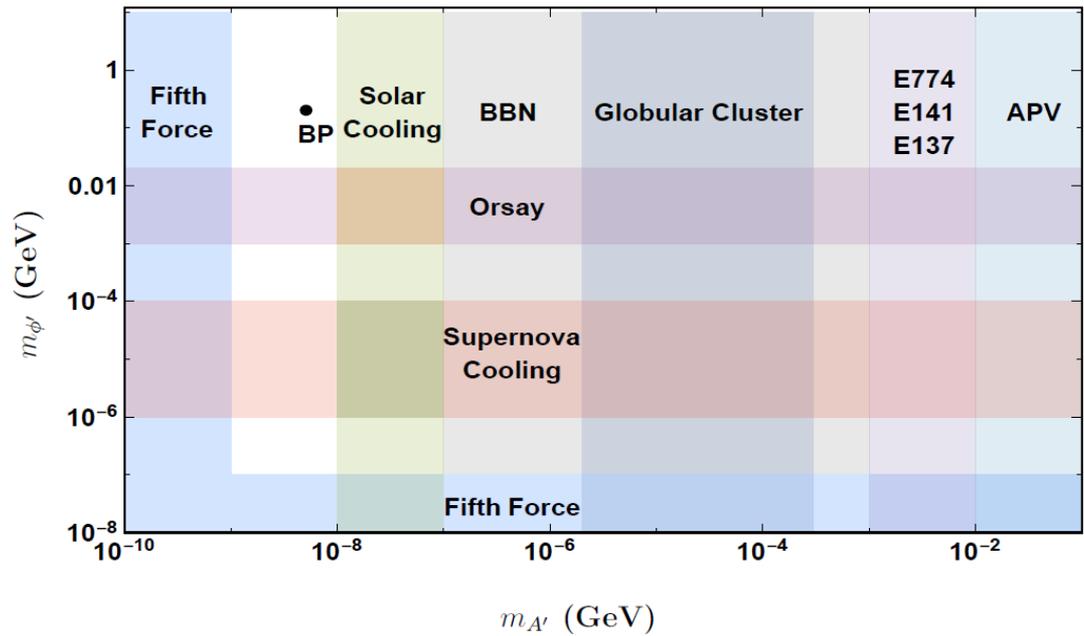
Various scenarios: Gauge boson (A')-scalar (ϕ') mediators parameter space



- *Muon model: u_R, d_R, ν_R, μ_R*

- *Electron model:
Muon model: u_R, d_R, ν_R, e_R*

- *Similarly, models with second generation quarks*



Direct Detection: Light DM&Med

Various ways of probing Sub-GeV DM:

Cosmic ray scattered

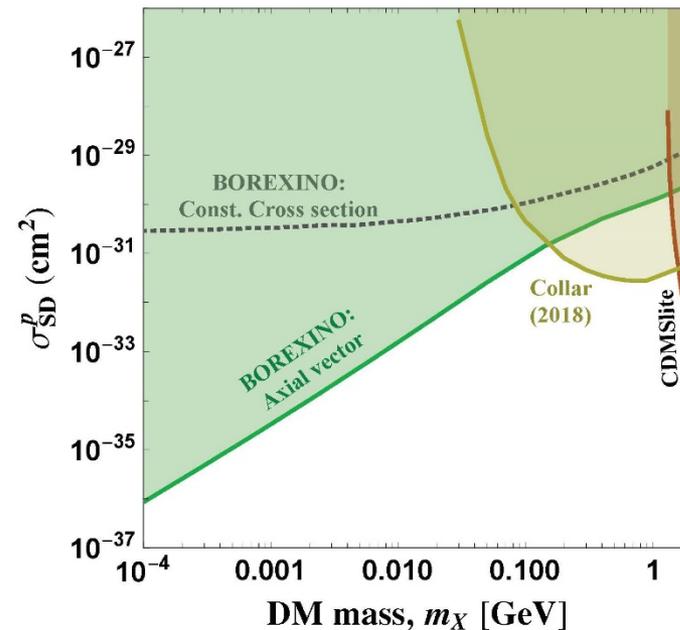
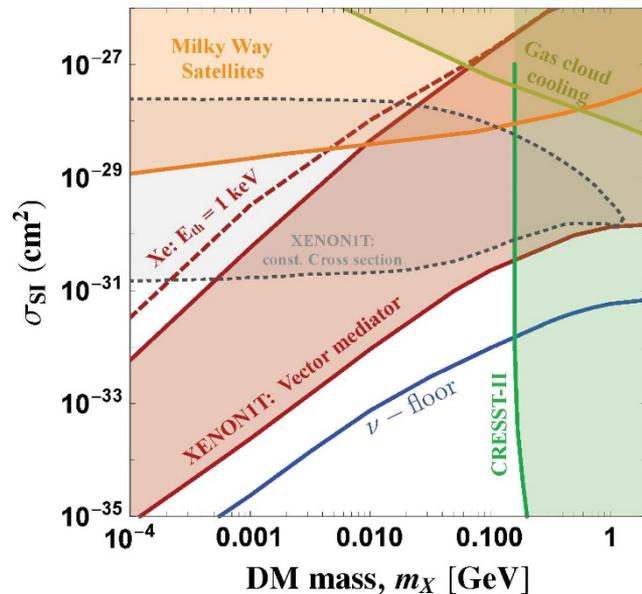


Bringmann, Pospelov, 2018

Ema, Sala, Sato, 2018

Dent, Dutta, Newstead, Shoemaker, 2019

*Low mass DM (up to 10 GeV)
becomes energetic → detection becomes easier*

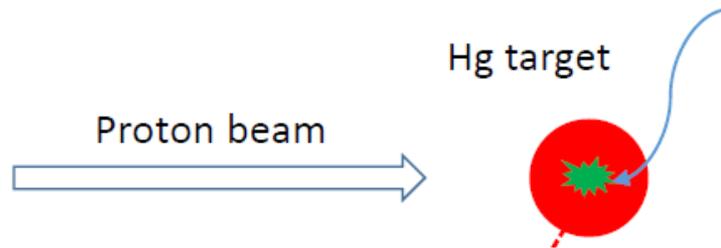


Dent, Dutta, Newstead, Shoemaker, 2019

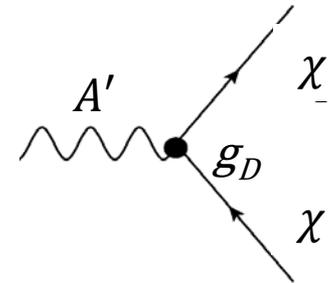
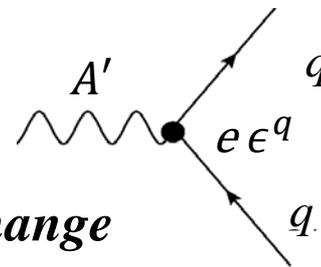
DM in ν experiments

Production of DM(χ_1) at COHERENT^T

Deniverville, Pospelov, Ritz, '15



- $\pi^- + p \rightarrow n + A', A' \rightarrow \chi\chi$



There is also another process: Charge exchange

- $\pi^- + p \rightarrow \pi^0 + n$

$$\pi^0 \rightarrow \gamma + A' \quad [\text{JSNS}^2 \text{ TDR}]$$

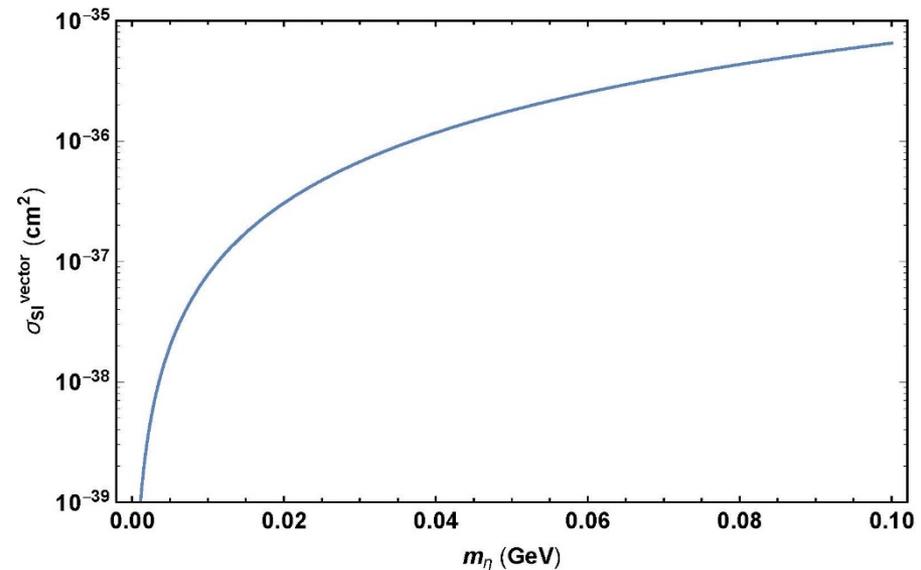
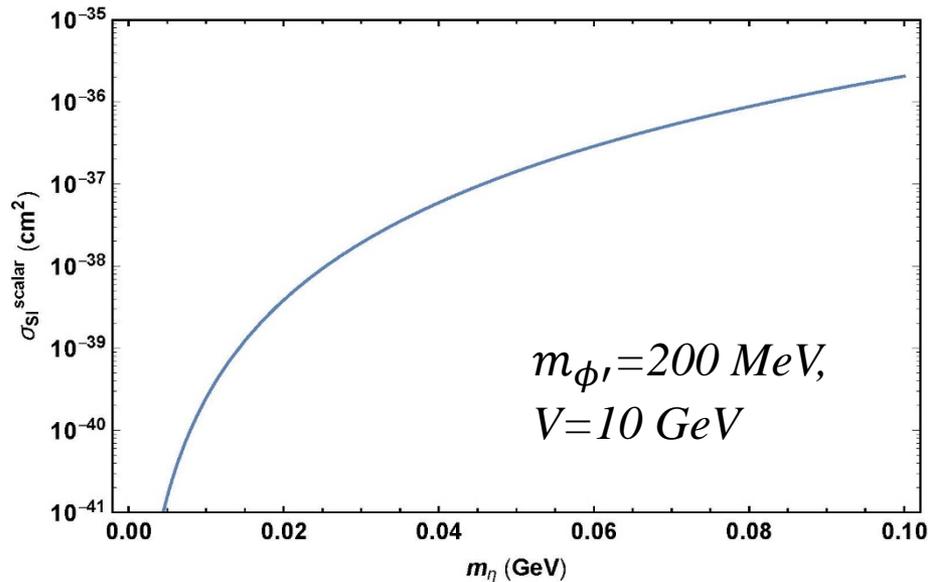
Probe Dark Photon (A') decay into DM utilizing: timing measurements

- *DM mass needs to be less than 70 MeV*

Dutta, Kim, Liao, Park, Shin, Strigari: 1906.10745

direct detection

Constraints from CRESST, Xenon1T (cosmic ray scattered,
Dent Dutta, Newstead, Shoemaker, 2019)



Elastic scattering: $\bar{\eta}\eta\bar{q}_Lq_R$

ϕ' mediated

$$\sigma_{\text{SI}}^{\text{scalar}(p,n)} = \frac{\mu_{\eta N}^2 m_\eta^2}{\pi V^4 m_{\phi'}^4} f_{p,n}^2$$

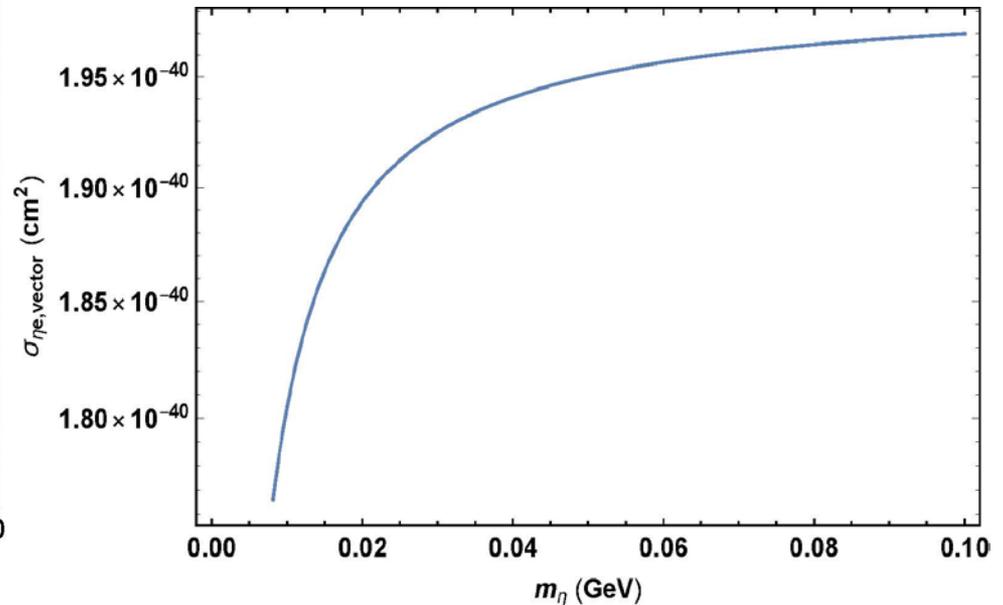
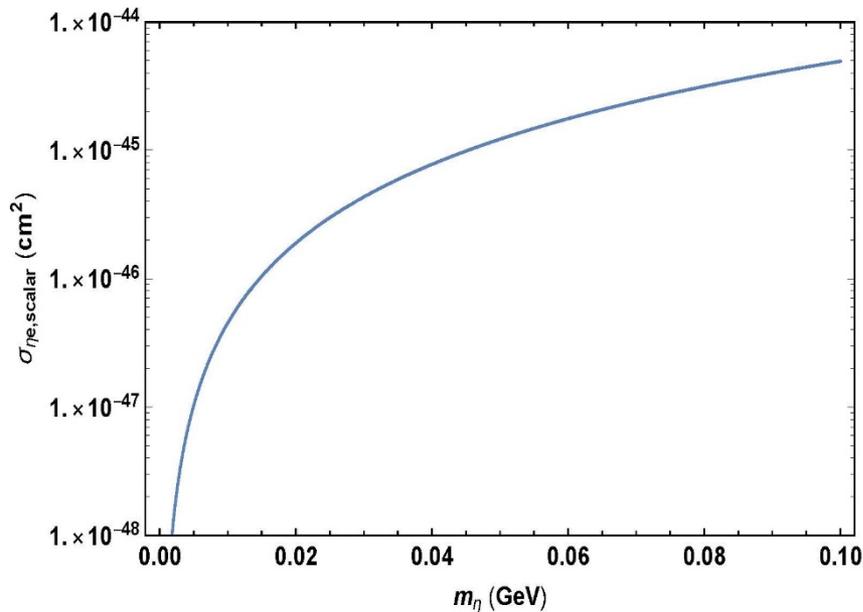
Inelastic scattering: $\bar{\eta}_1\gamma^\mu\eta_2\bar{q}_L\gamma^\mu q_L$

A' mediated

$$\sigma_{\text{SI}}^{\text{vector}(p,n)} = \frac{\mu_{\eta N}^2}{16\pi V^4}$$

direct detection

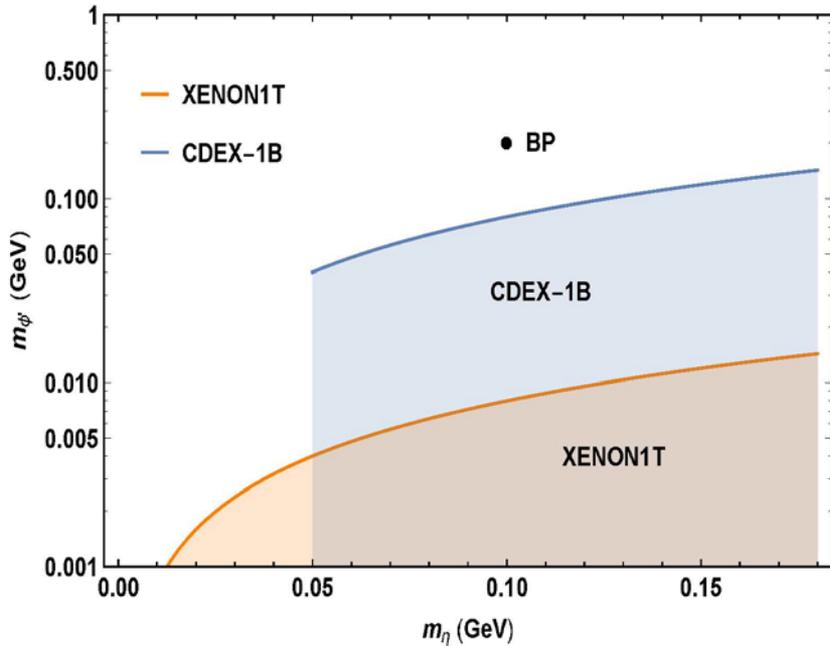
Dark Matter-electron Scattering



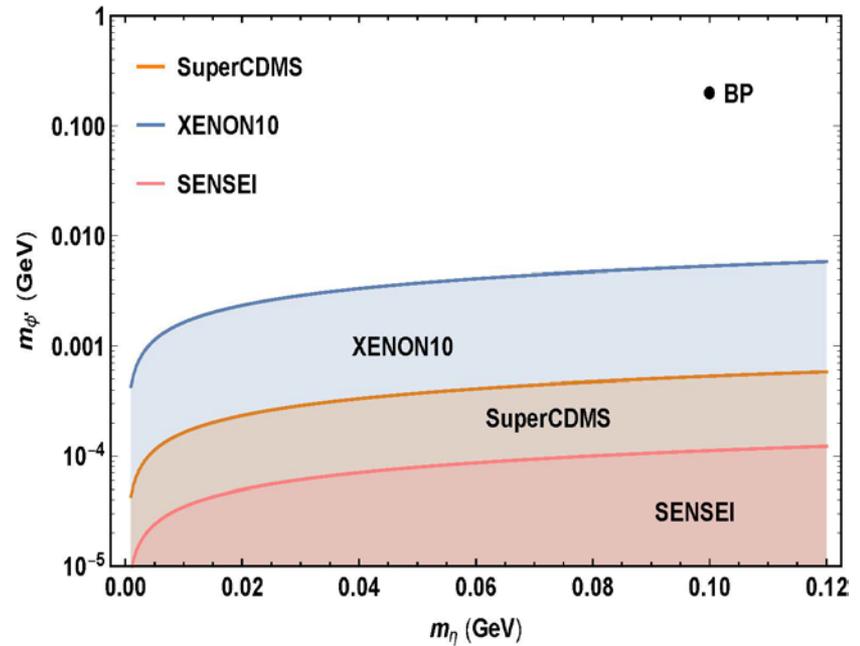
The allowed cross-section $\leq 10^{-40} \text{ cm}^2$

*Ema, Sala, Sato, 2018
Xenon 1T, 1906.04717*

direct detection



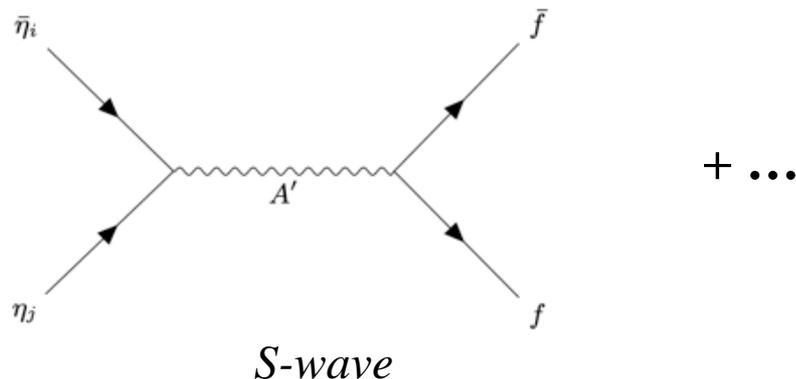
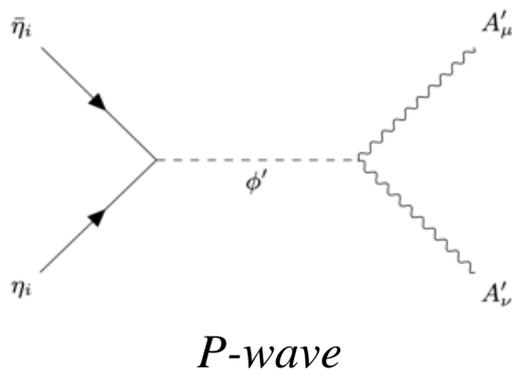
Dark Matter-nucleon



Dark Matter-electron

Thermal Relic Abundance

Dominant two body final states: $\bar{\ell}\ell, \bar{\nu}\nu, \pi\pi, \pi^0(\phi', A', \gamma)$
 +
 $A'A', \phi'\phi'$ and $\phi'A'$



Resonance/non-resonance:

	$m_{A'}$ (MeV)	$m_{\phi'}$ (MeV)	m_{η} (MeV)	m_{ν_s} (MeV)	m_{ν_D} (MeV)	$\langle\sigma v\rangle$ (cm ³ /sec)	$\sigma_{\text{SI}}^{\text{scalar}}$ (pb)	$\sigma_{\text{SI}}^{\text{vector}}$ (pb)
muon case	55	200	100	10	10 ⁻³	3×10 ⁻²⁶	2.05	6.50
	70	10 ⁴	50	10 ¹⁶	10 ⁴	3×10 ⁻²⁶	3.29×10 ⁻⁷	1.80
electron case	5 × 10 ⁻⁶	200	100	10	10 ⁻³	3×10 ⁻²⁶	2.05	6.50

Outlook

- **What is the scale of new physics?**
- **Models with light mediators are very interesting:
Dark Matter, $g-2$ of electron, neutrino masses,
Yukawa coupling hierarchy, MiniBooNE excess**
- **We constructed a model, where DM mass ~ 100 MeV,
two light mediators: scalar and gauge boson**
- **The model has sufficient allowed parameter space after satisfying
various experimental constraints**
- **DM relic abundance can occur thermally**