



# Unearthing kinematic information in WH production

Based on `arXiv:1908.XXXX`

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In collaboration with J. Brehmer, S. Dawson, F. Kling, and T. Plehn

**APS DPF Meeting, August 1, 2019**  
**Northeastern University**

# The Higgs Legacy of the LHC

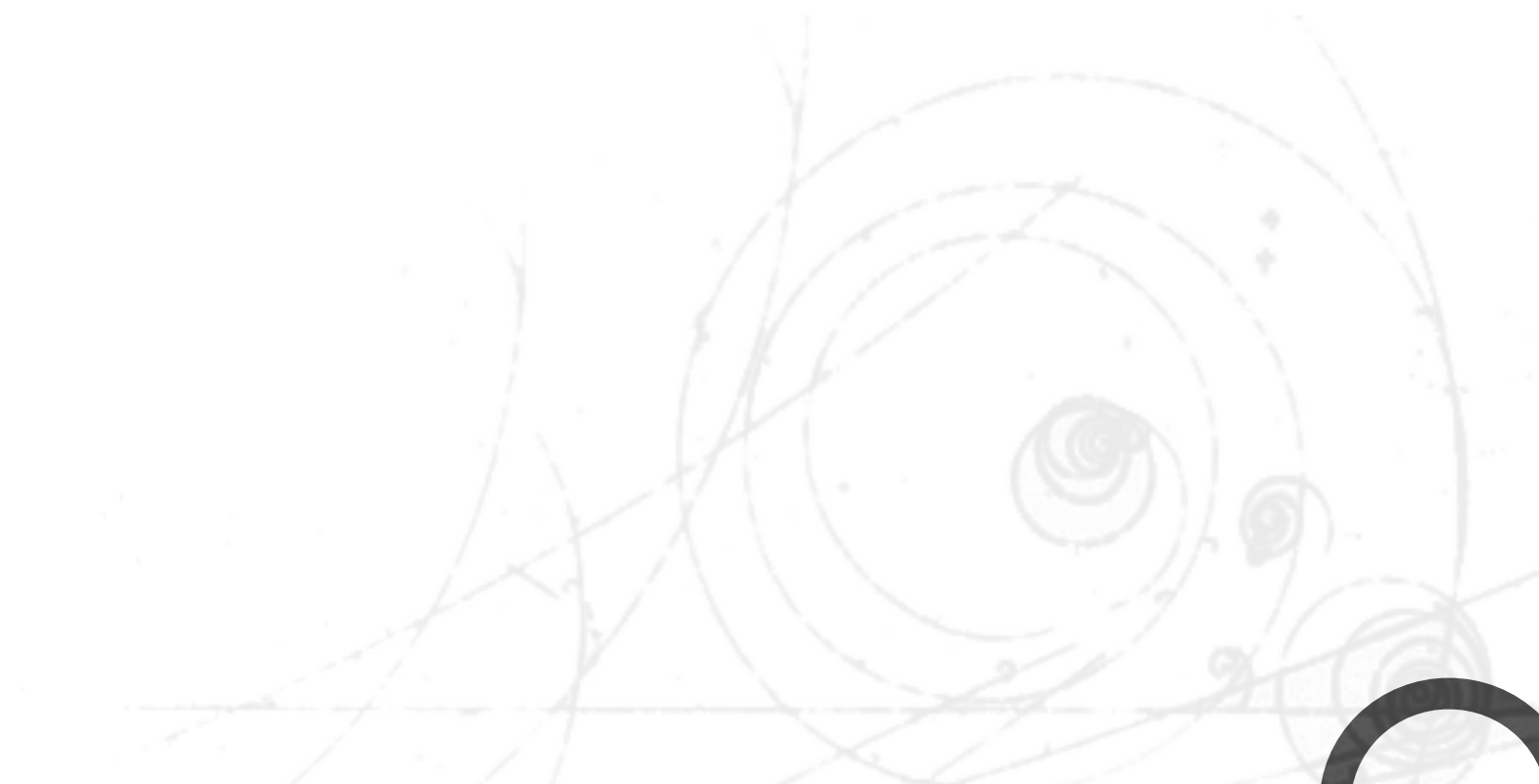
With the Standard Model complete, the next steps are to measure the theory as precisely as possible

SMEFT: Parameterizes BSM effects in terms of higher dimensional operators

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_{d,k} \frac{C_k^d}{\Lambda^{d-4}} \mathcal{O}_k^d$$

In this work: truncate at  $d = 6$ , and consider effects only up to  $\mathcal{O}(1/\Lambda^2)$

**Goal: Understand what the legacy measurements of the LHC will tell us about BSM physics**



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**Goal: Understand what the legacy measurements of the LHC will tell us about BSM physics**

**Where do global analyses of Higgs-Gauge Sector get their information?**

A. Butter, O. Éboli, J. Gonzalez-Fraile, M. C. Gonzalez-Garcia, T. Plehn, M. Rauch, 1604.03105

J. Ellis, C. Murphy, V. Sanz, T. You, 1803.03252

A. Biekötter, D. Gonçalves, T. Plehn, M. Takeuchi, D. Zerwas, 1811.08401

A. Biekötter, T. Corbett, T. Plehn, 1812.07587

Here: consider example of  $WH$  production, in the  $\ell\nu b\bar{b}$  channel

Compare traditional analysis methods with modern inference techniques

Use Information Geometry to make these questions quantitative

# WH Production in the SMEFT

In the Warsaw Basis:

There are four relevant operators\*

$$\mathcal{O}_{HD} = |H^\dagger D^\mu H|^2 \quad \left. \vphantom{\mathcal{O}_{HD}} \right\} \begin{array}{l} \text{Finite Higgs} \\ \text{wave-function} \\ \text{renormalization} \end{array}$$

$$\mathcal{O}_{H\Box} = (H^\dagger H)\Box(H^\dagger H)$$

$$\mathcal{O}_{HW} = (H^\dagger H) W_{\mu\nu}^a W^{a\mu\nu}$$

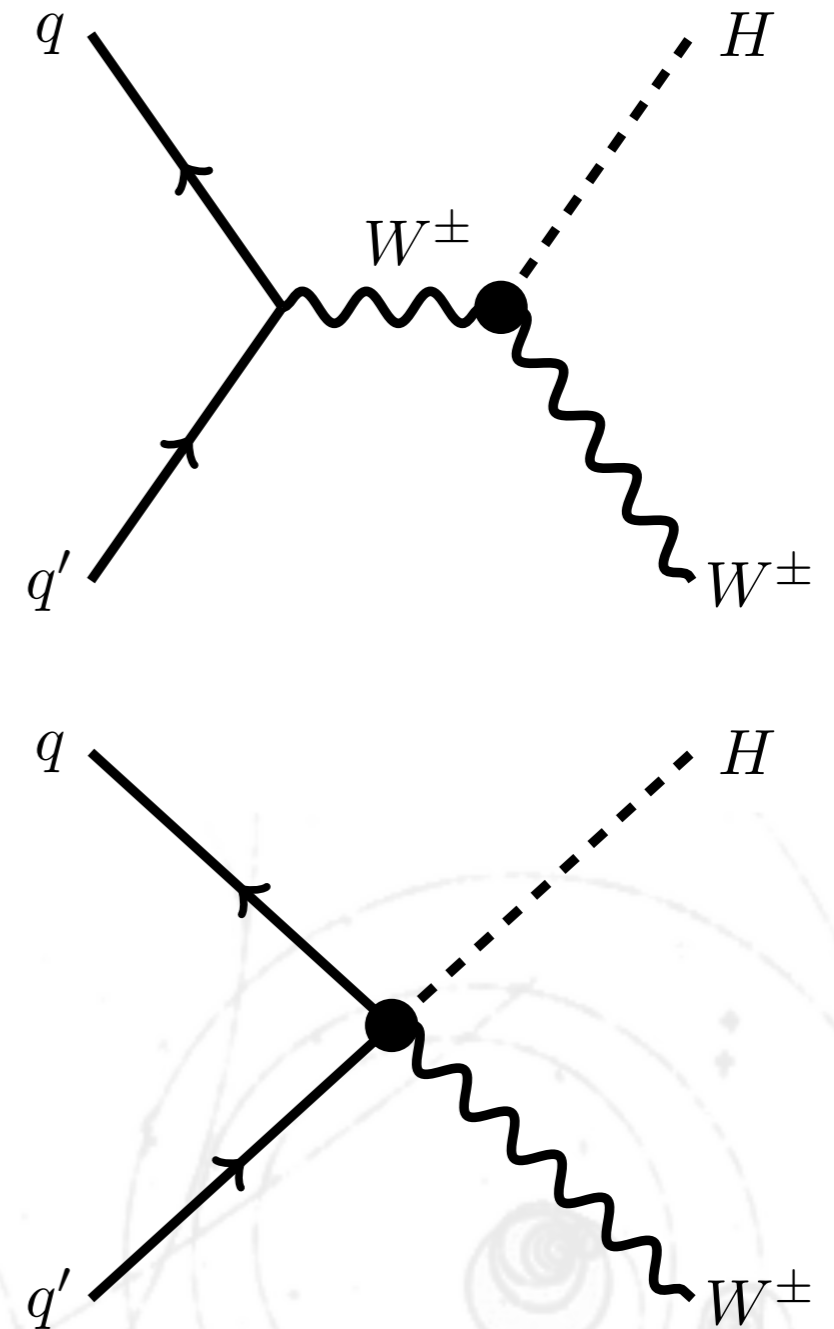
$$\mathcal{O}_{Hq}^{(3)} = (H^\dagger i\overleftrightarrow{D}_\mu^a H)(Q_L \sigma^a \gamma^\mu Q_L)$$

$\mathcal{O}_{HD}, \mathcal{O}_{H\Box}$  always enter in the combination

$$\frac{\tilde{C}_{HD}}{\Lambda^2} \tilde{\mathcal{O}}_{HD} \equiv \frac{\tilde{C}_{HD}}{\Lambda^2} \left( \mathcal{O}_{H\Box} - \frac{1}{4} \mathcal{O}_{HD} \right)$$

So we're left with 3 theory parameters:

$$\left\{ \tilde{C}_{HD}, C_{HW}, C_{Hq}^{(3)} \right\}$$



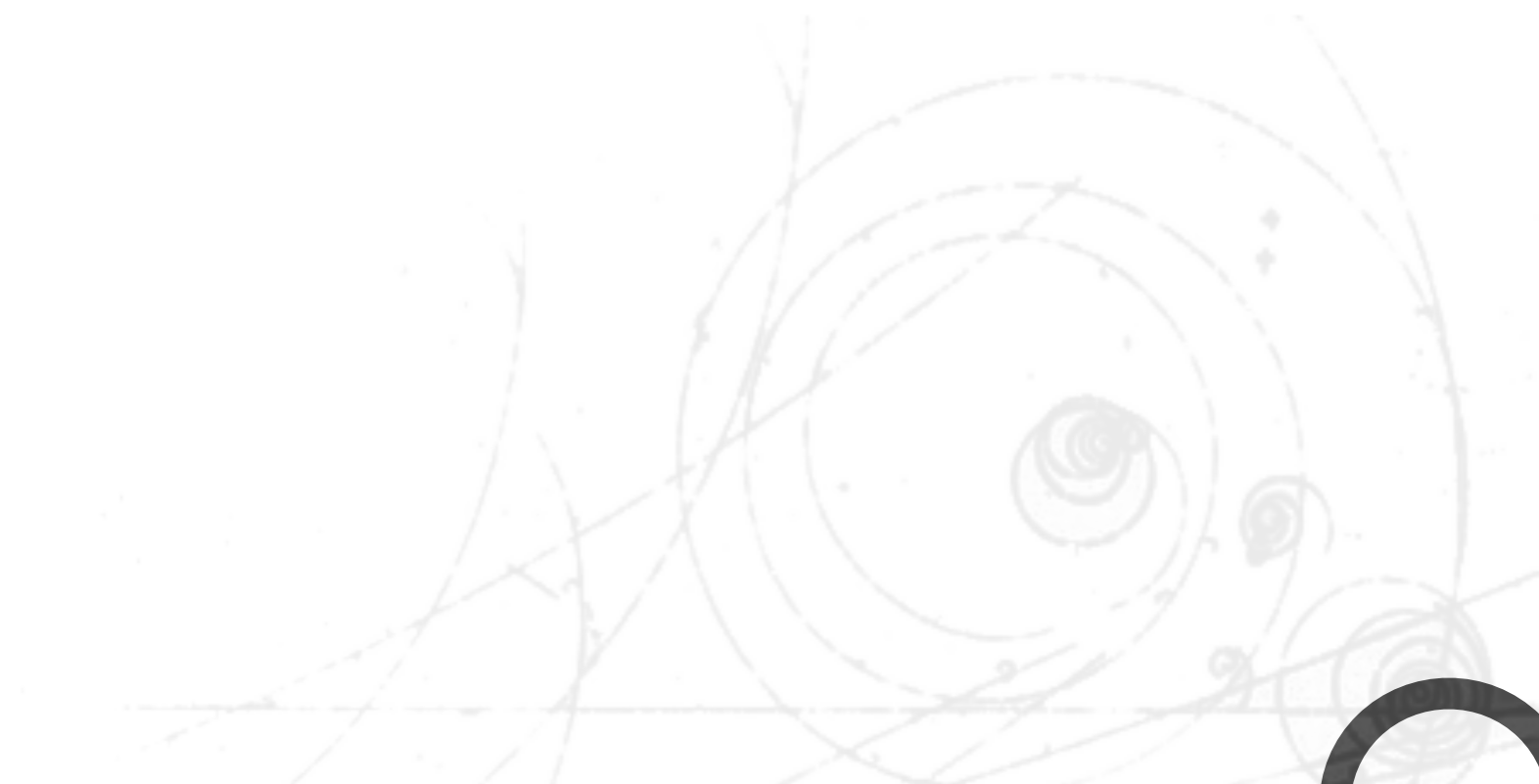
\*neglecting operators modifying the Higgs decay, which are not well measured here

# How do we estimate theory parameters?

The object linking observed events to theory parameters is the likelihood function:

$$p(x|\theta) = \int dz_d \int dz_s \int dz_p \underbrace{p(x|z_d)}_{\text{detector smearing}} \underbrace{p(z_d|z_s)}_{\text{shower-splitting}} \underbrace{p(z_s|z_p)}_{\text{parton-level}} p(z_p|\theta)$$

Doing this integral explicitly is intractable  $(\sim |\mathcal{M}(z_p|\theta)|^2)$



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But we can make progress by choosing simple observables

Rate:  $x = N_{\text{events}}$

Histogram:  $x = \{N_{\text{bin}}\}$

(e.g., Simplified Template Cross Sections)

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(e.g., Simplified Template Cross Sections, STXS)

Close to the Standard Model, we can use *optimal observables*

$$\mathcal{O}_n(x) \sim \left. \frac{\partial}{\partial \theta_n} \log p(x|\theta) \right|_{\theta=0} \equiv t_n(x|\theta=0)$$

D. Atwood, A. Soni PRD 45, 7 (1992)

M. Diehl, O. Nachtmann, Z. Phys. C 62 (1994)

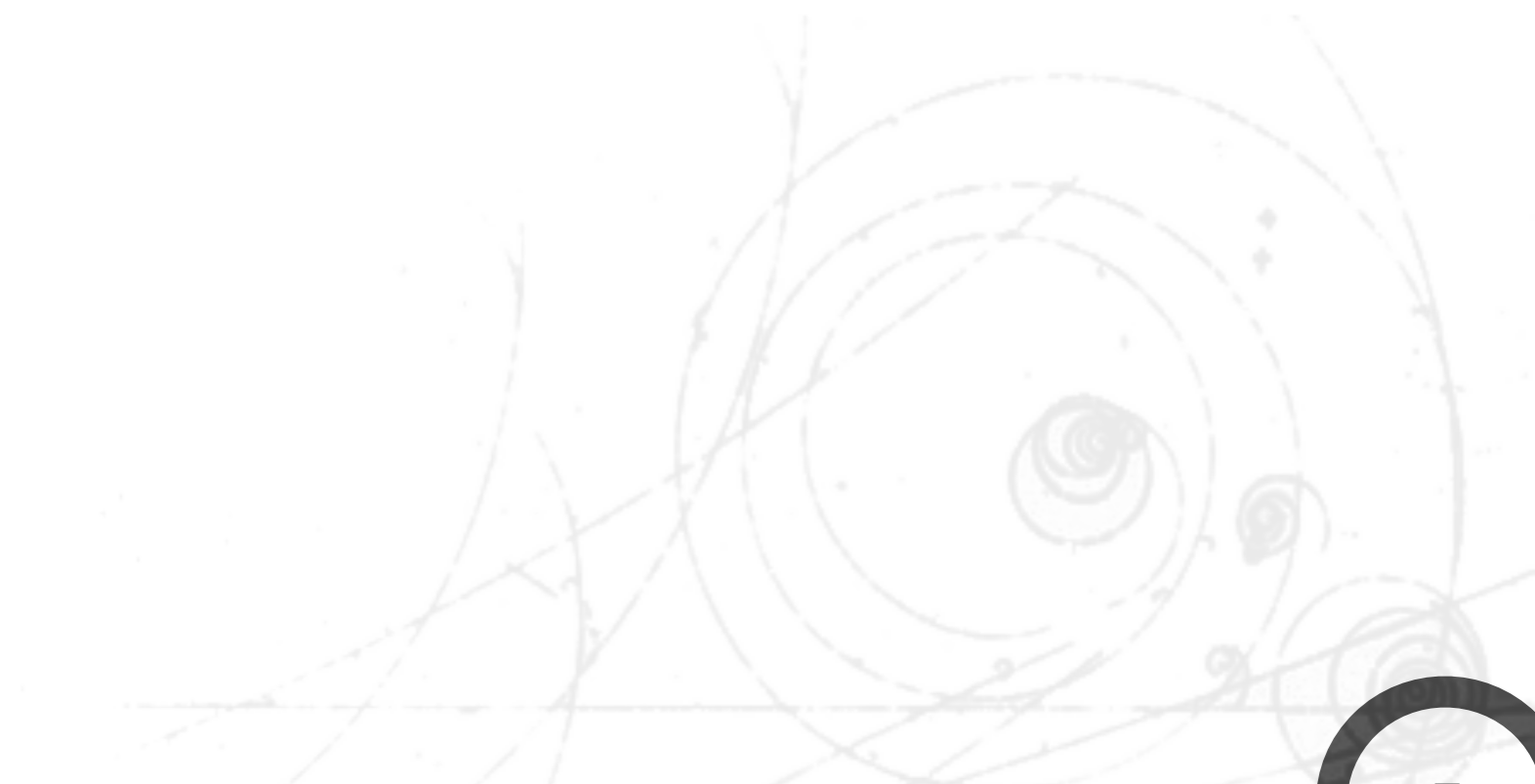
# Information Geometry

J. Brehmer, K. Cranmer, F. Kling, T. Plehn,  
1612.05261

J. Brehmer, F. Kling, T. Plehn, T. Tait,  
1712.02350

Expected limits can be nicely summarized by the *Fisher Information*:

$$I_{ij}(\theta_0) \equiv - \int dx p(x | \theta_0) \frac{\partial \log p(x | \theta)}{\partial \theta_i} \frac{\partial \log p(x | \theta)}{\partial \theta_j} = \langle t_i t_j | \theta_0 \rangle$$





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1612.05261

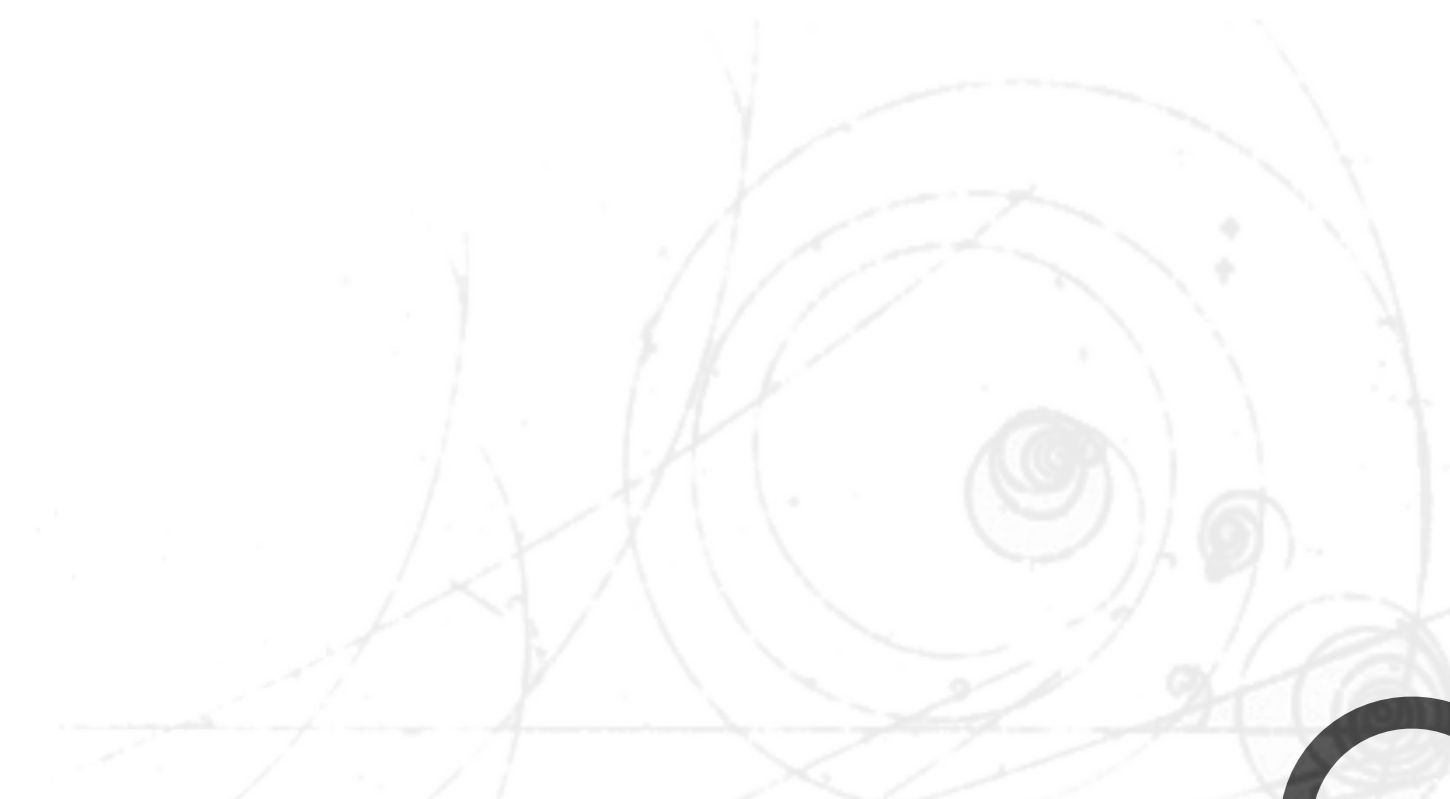
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which has a number of useful properties:

- Expected limits bounded by the Inverse Fisher Information (Cramèr-Rao)
- Transforms covariantly under parameter transformations (EFT bases)
- Additive for different processes & phase space regions
- Can include systematic uncertainties as nuisance parameters



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$$I_{ij}^{\text{Rate}} = \begin{pmatrix} C_{H\Box} & C_{HD} & C_{HW} & C_{Hq}^{(3)} \\ 28.0 & -7.0 & 101 & 272 \\ -7.0 & 1.75 & -25.2 & -68 \\ 101 & -25.2 & 360.1 & 979 \\ 272 & -68 & 979 & 2660 \end{pmatrix}$$

**Degeneracy between rate only operators is apparent**

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Moreover, the Fisher Information can be reliably estimated *in the presence of detector effects* using Machine Learning techniques

J. Brehmer, K. Cranmer, G. Louppe, J. Pavez  
1805.00013, 1805.00020, 1805.12244

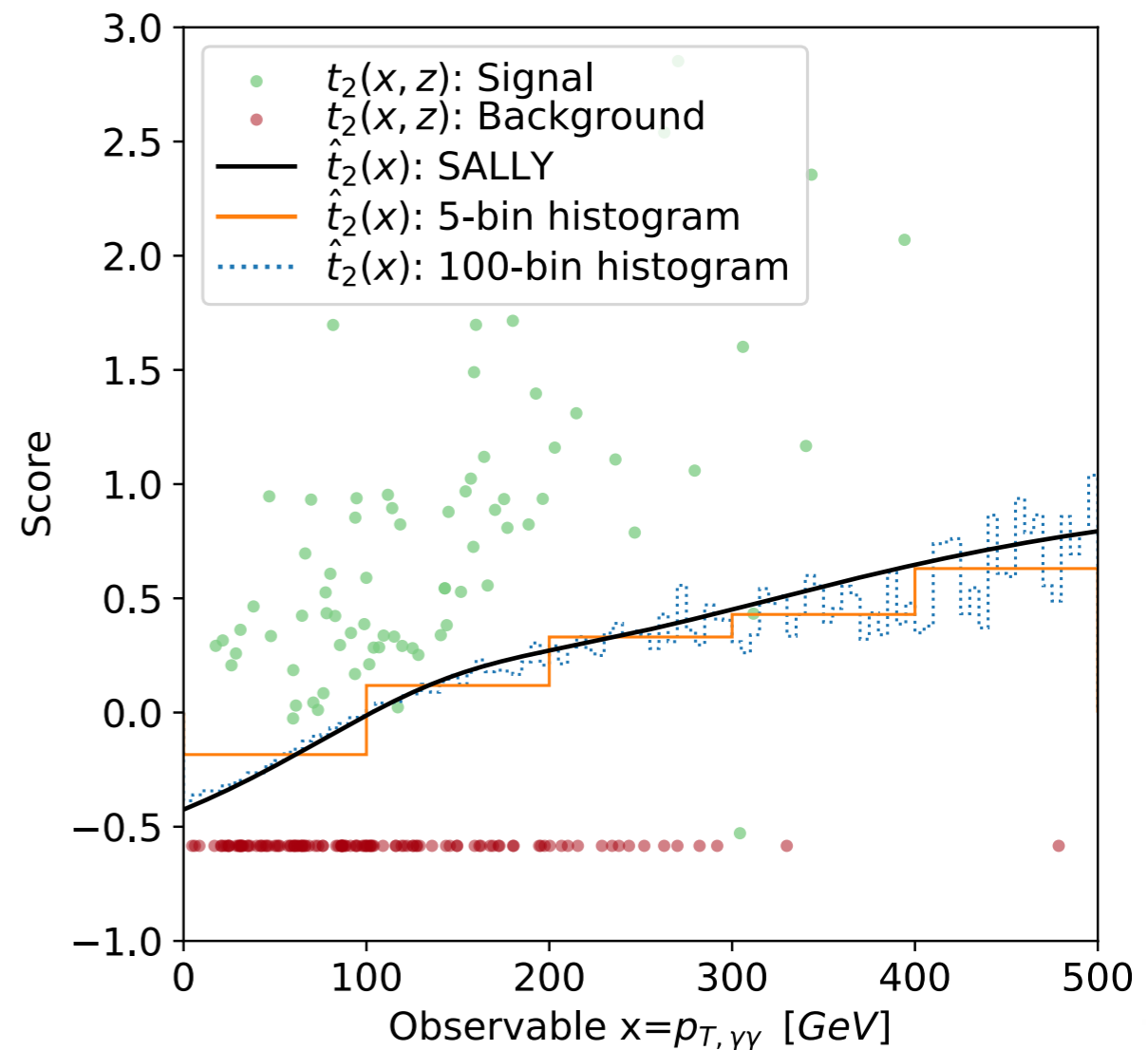
J. Brehmer, F. Kling, I. Espejo, K. Cranmer  
1907.10621

# MadMiner

An Inference Toolkit for Particle Physics

J. Brehmer, F. Kling, I. Espejo, K. Cranmer, 1907.10621

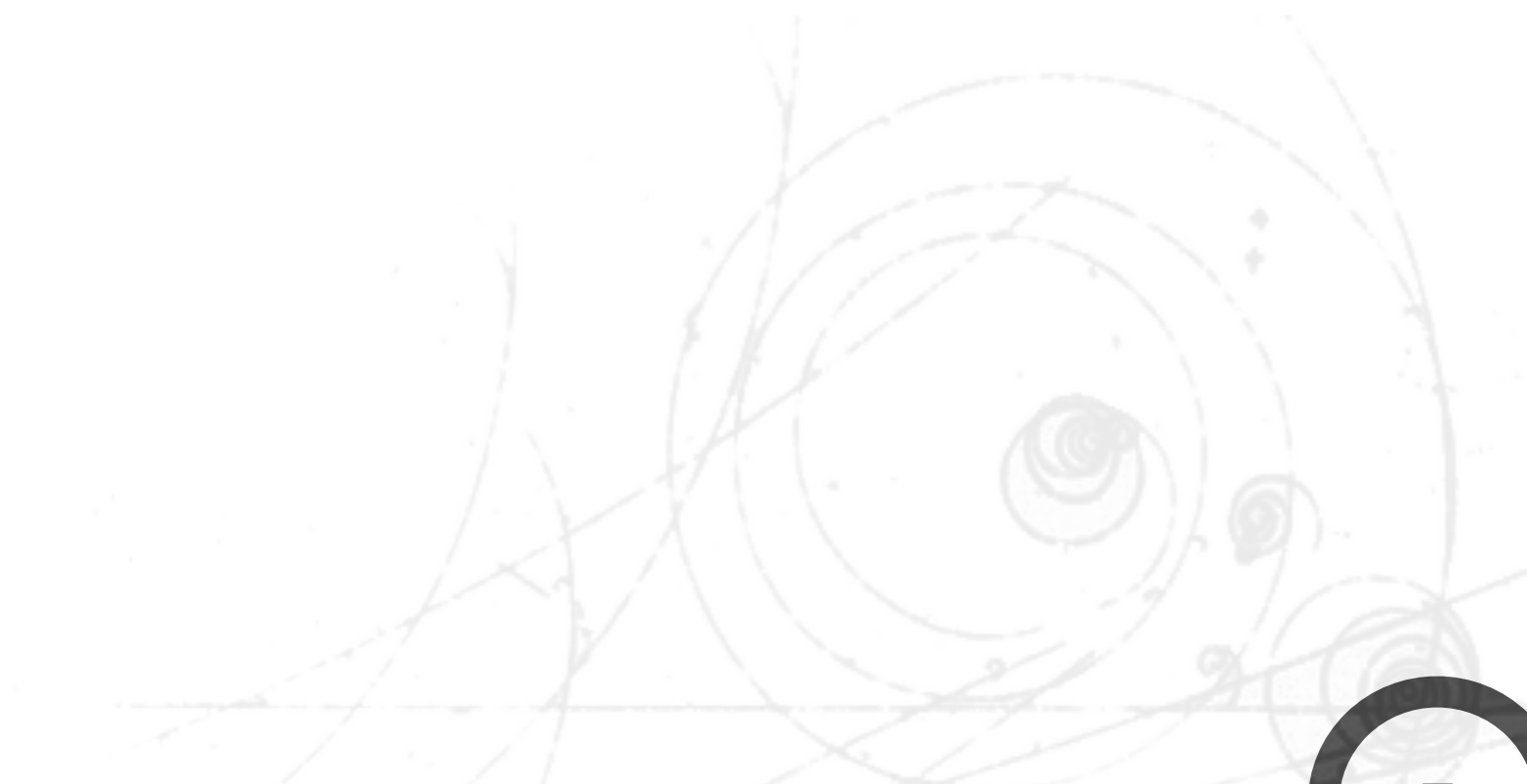
- Combines the power of ML inference methods with intuition of the Matrix Element Method
- Automizes score & likelihood ratio estimation techniques, fully interfaced with simulation tools
- Out of the box: pheno-level analysis
  - ▶ MadGraph, Pythia, Delphes
  - ▶ Backgrounds
  - ▶ PDF/scale uncertainties
  - ▶ ML uncertainties
- Scalable to state-of-the-art experimental tools
- Python package
  - ▶ Modular interface
  - ▶ Extensive documentation
  - ▶ On GitHub: [github.com/diana-hep/madminer](https://github.com/diana-hep/madminer)
  - ▶ Easy to install: `pip install madminer`



# MadMiner Workflow

Setup Morphing Basis  
(choose theory parameters)

$$\left\{ \tilde{C}_{HD}, C_{HW}, C_{Hq}^{(3)} \right\}$$



# MadMiner Workflow

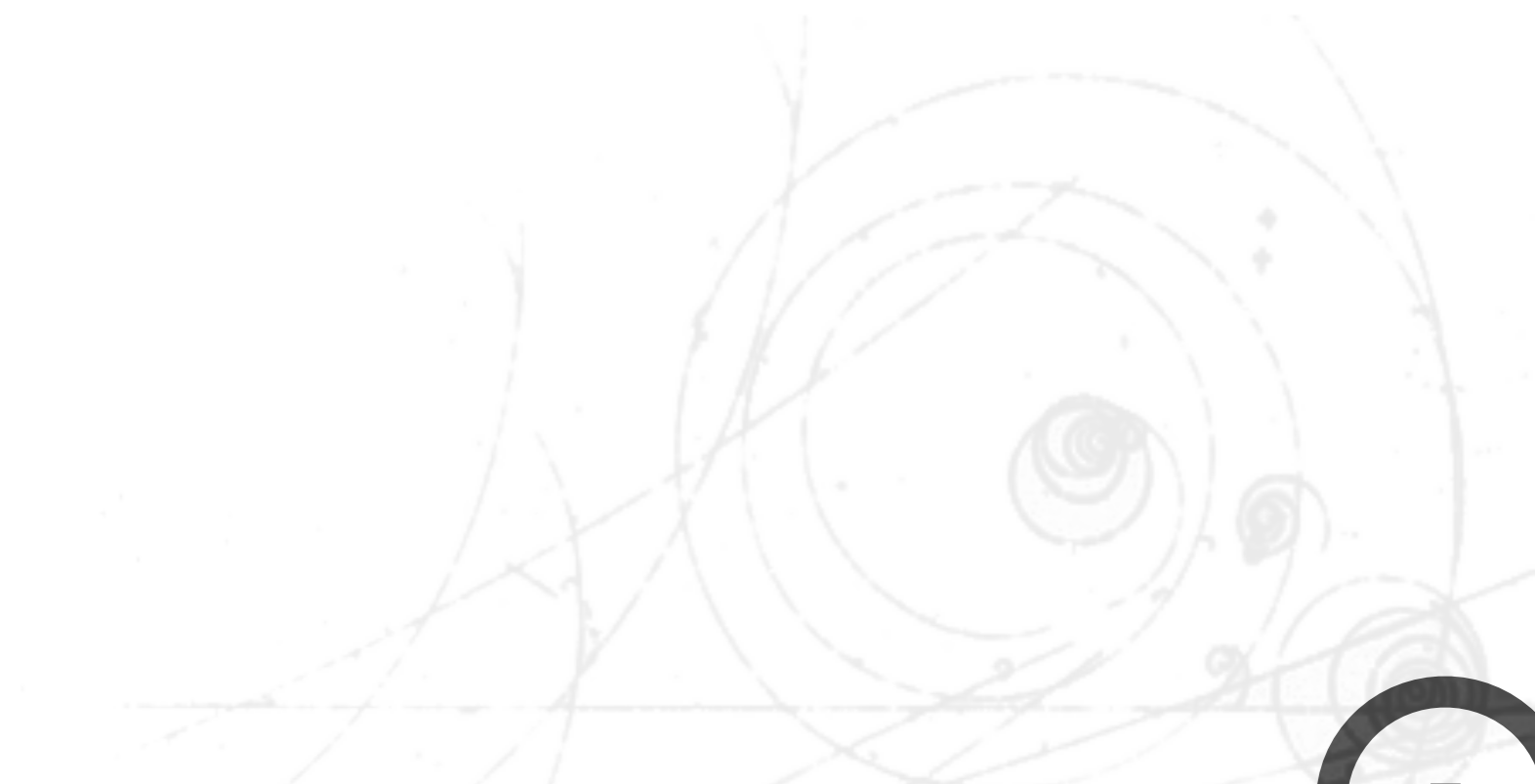
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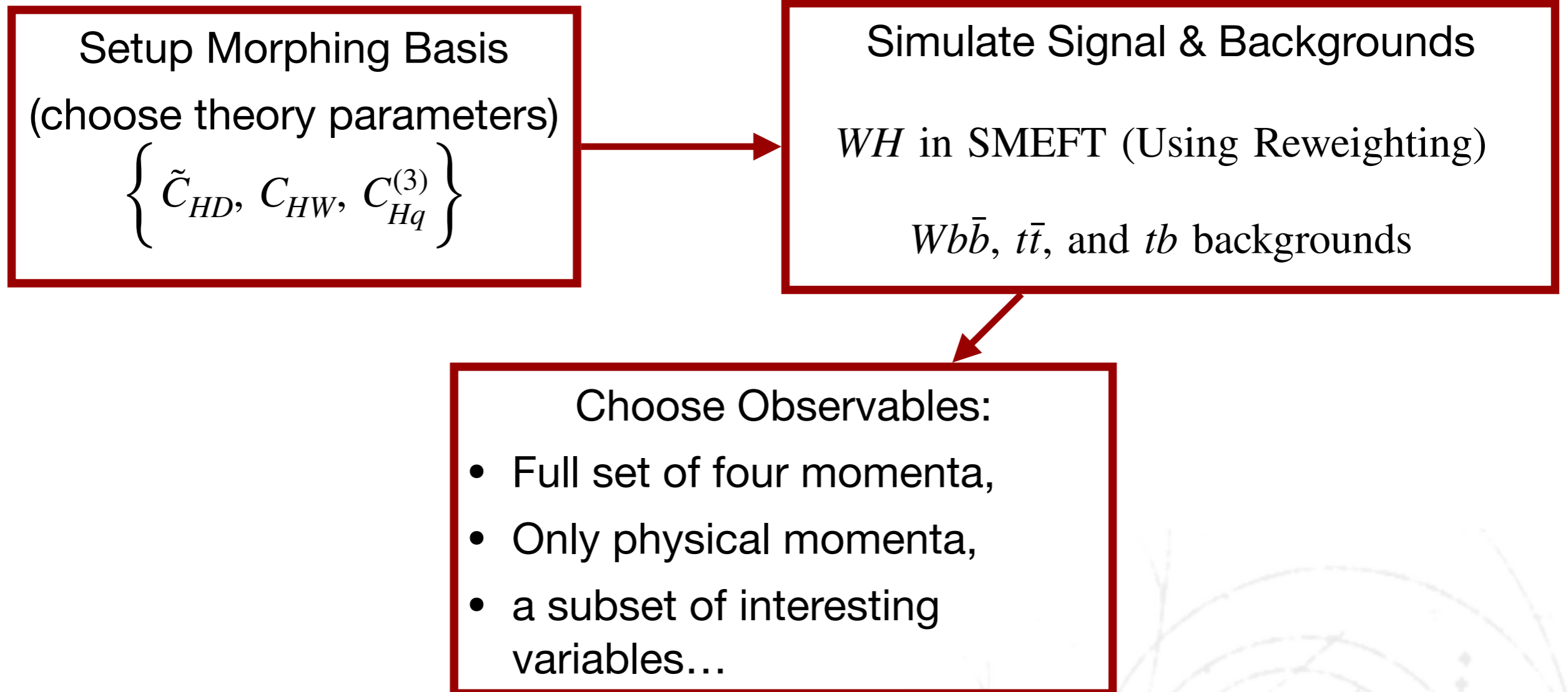
Simulate Signal & Backgrounds

*WH* in SMEFT (Using Reweighting)

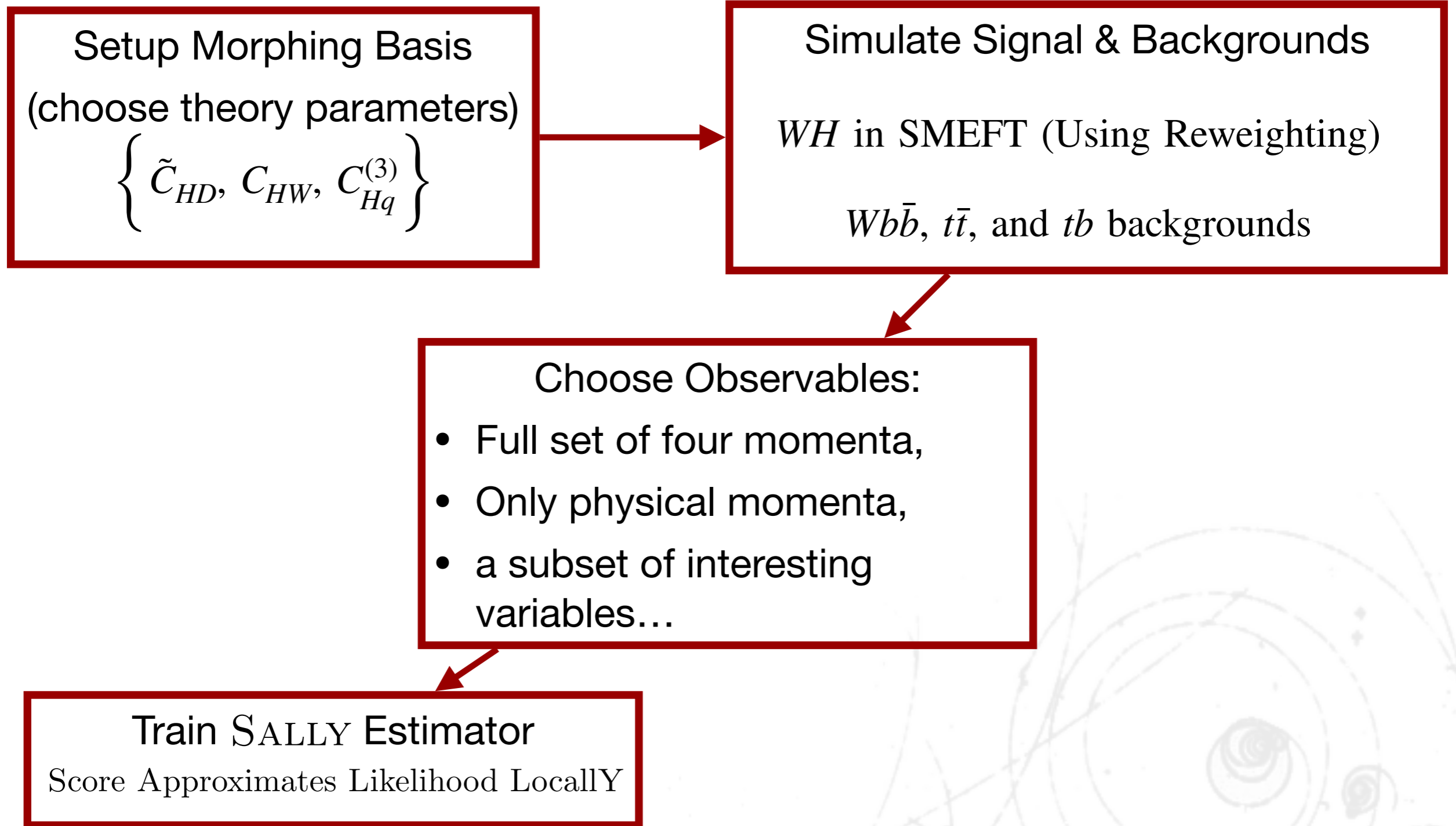
*Wb* $\bar{b}$ , *t* $\bar{t}$ , and *tb* backgrounds



# MadMiner Workflow

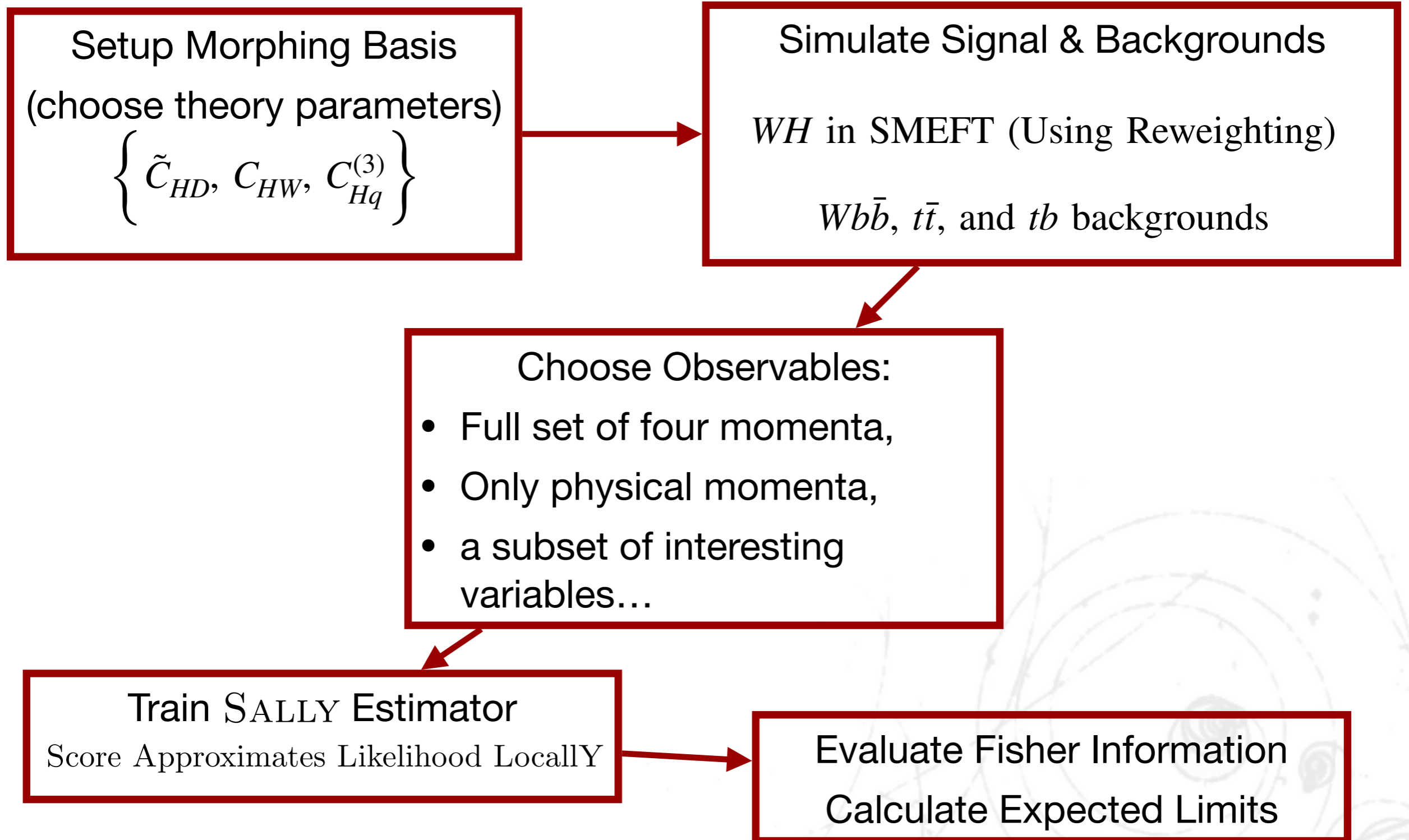


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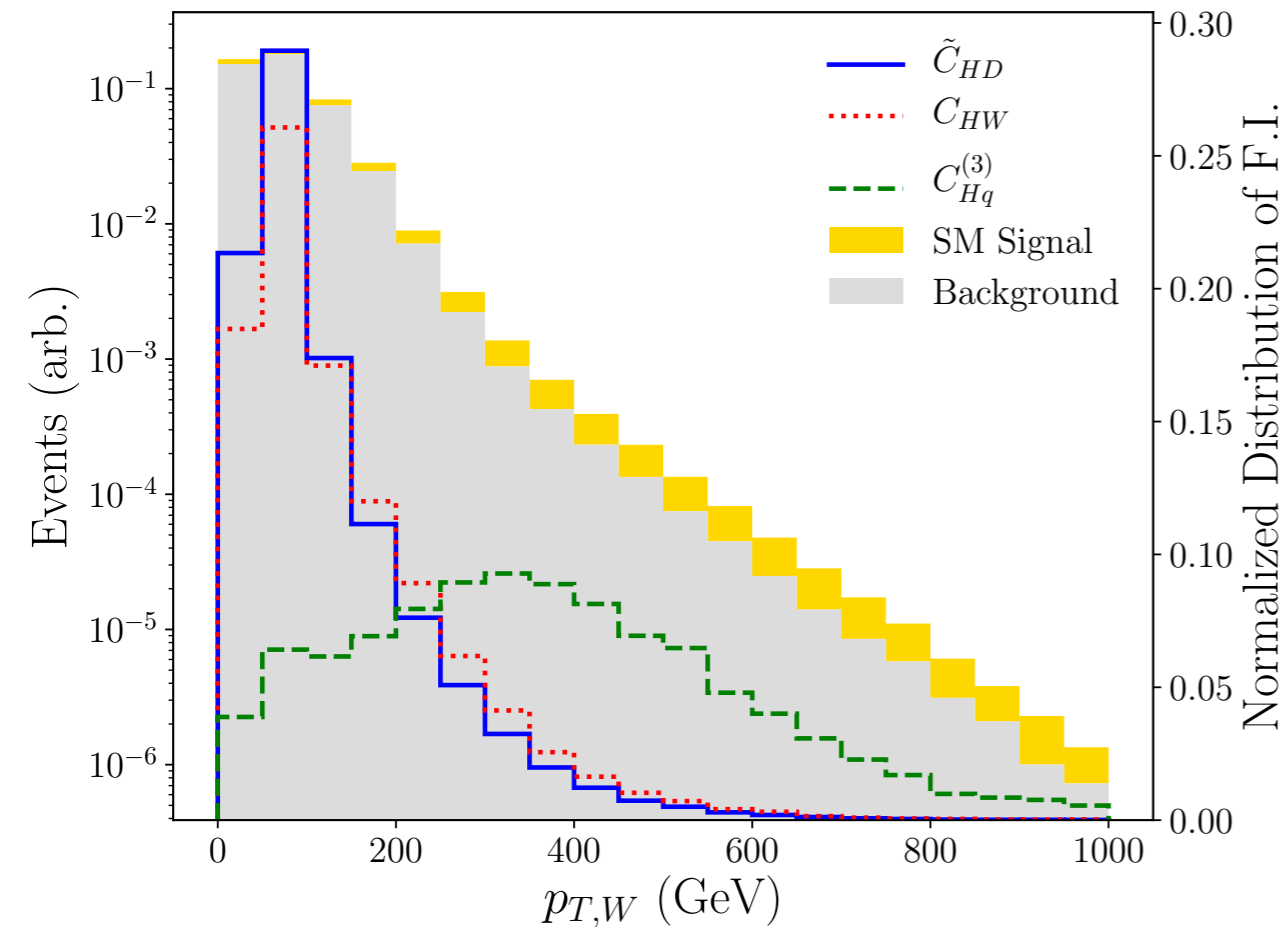




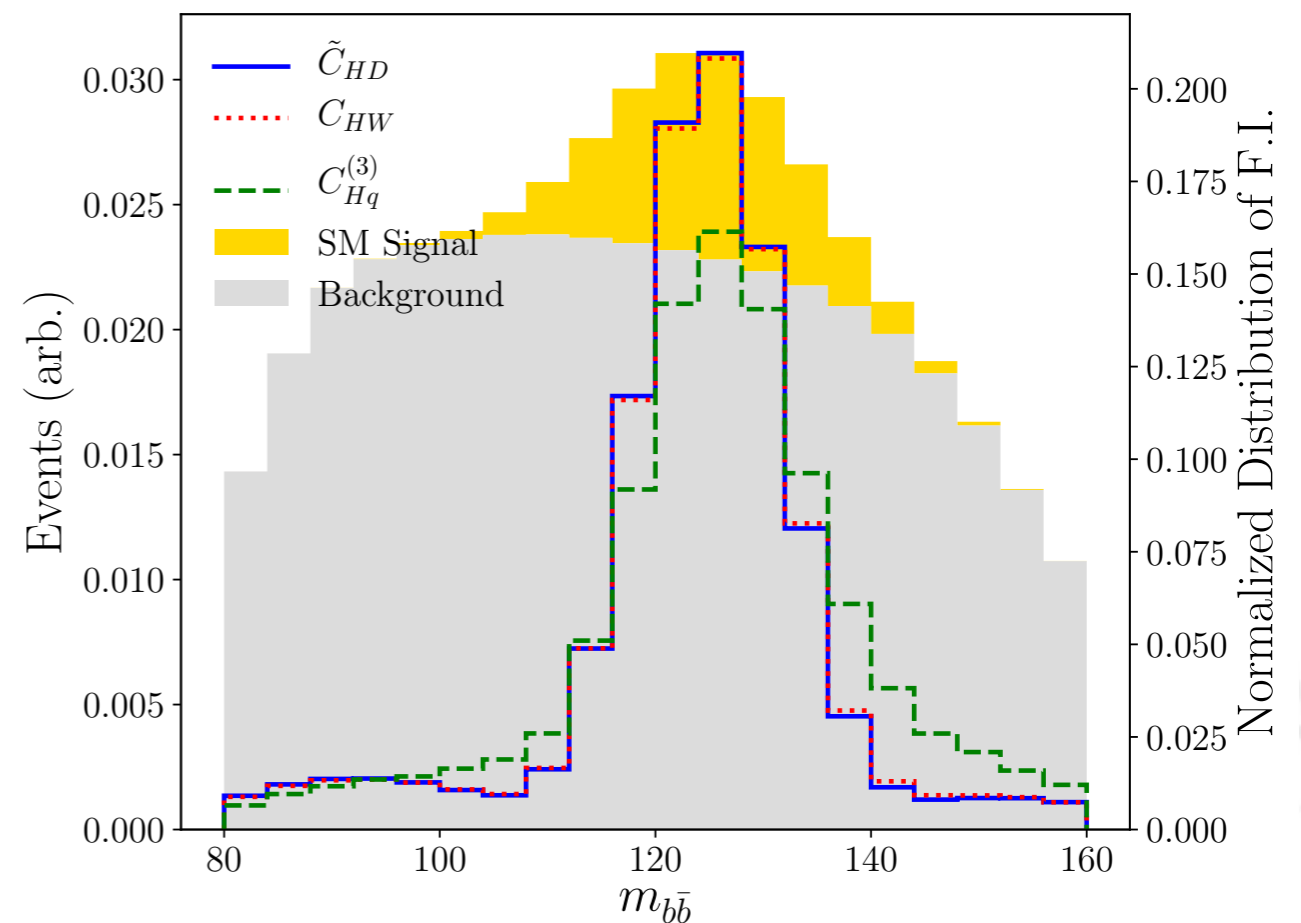
# MadMiner Workflow



# Where is the Information?



Momentum enhanced operators have Fisher Information peaked in high momentum bins



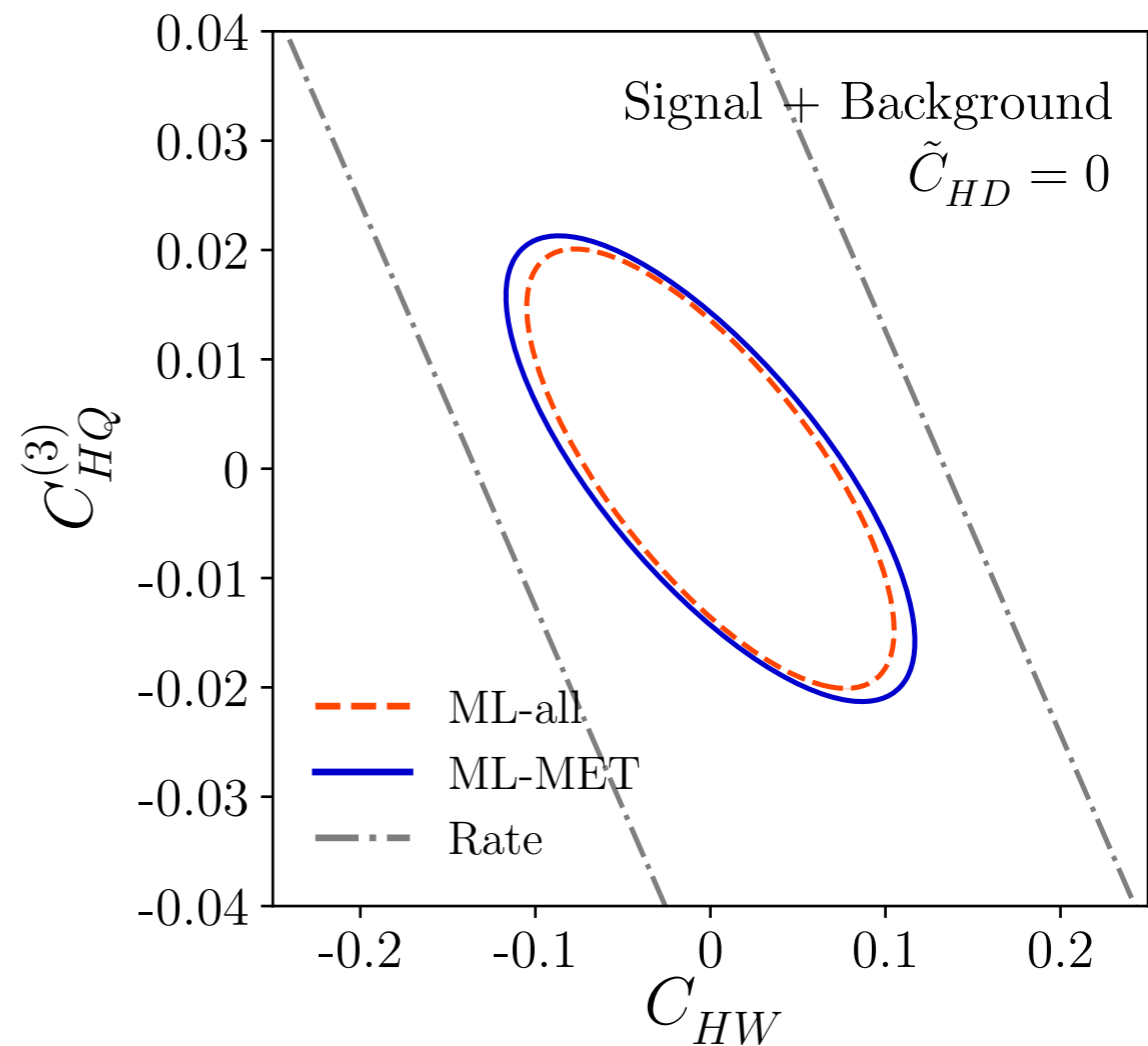
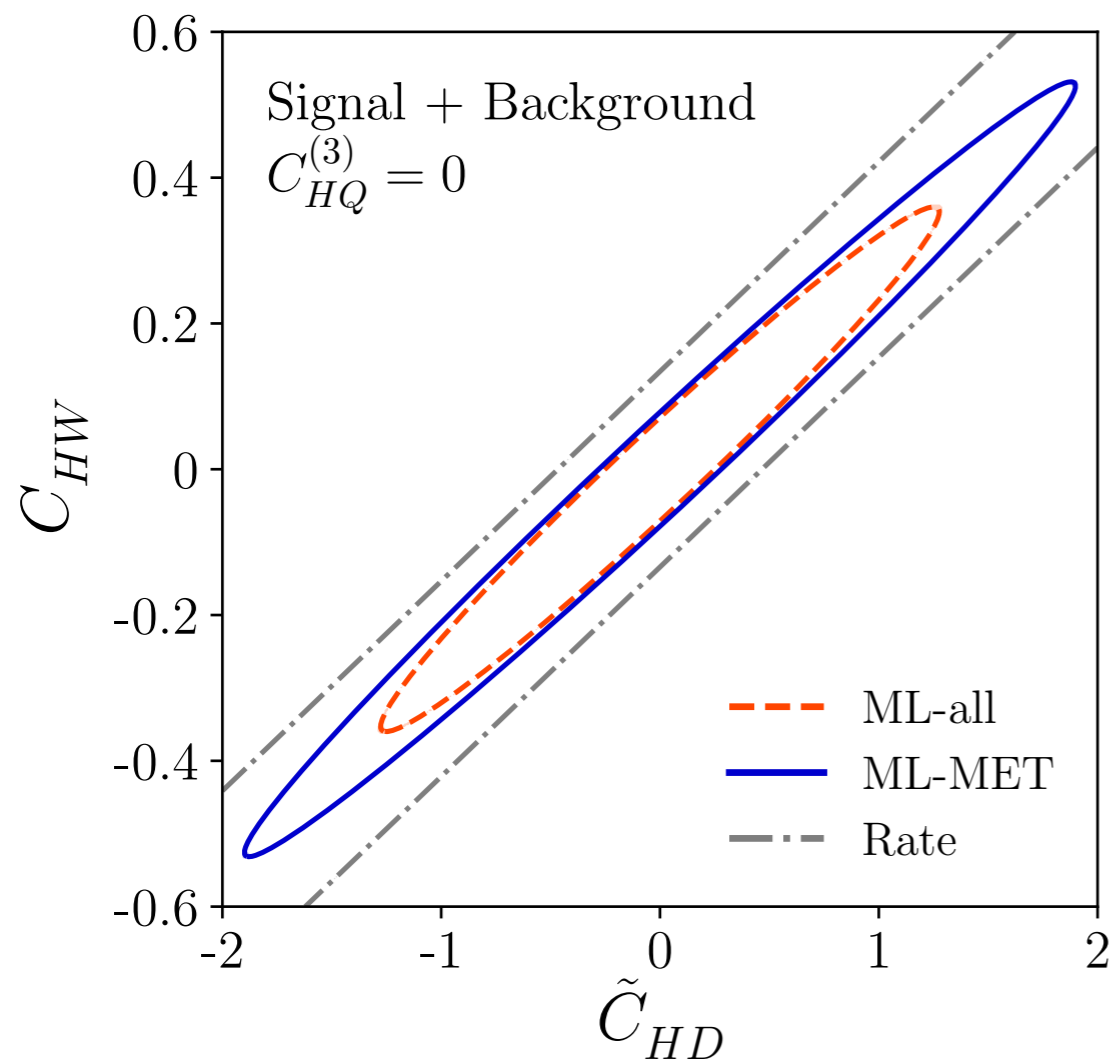
Including backgrounds can change naive distribution of information

$$\mathcal{M}^0 \sim 1 + \frac{sv^2}{M_W^2 \Lambda^2} C_{Hq}^{(3)}$$

$$\mathcal{M}^\pm \sim \frac{M_W}{\sqrt{s}} \left( 1 + \frac{sv^2}{M_W^2 \Lambda^2} (C_{HW} + C_{Hq}^{(3)}) \right)$$

# What is lost in the Neutrino Momentum?

We can train our estimator on different sets of observables:  
with and without the neutrino energy & longitudinal momentum



This allows us to quantify how much information is lost in missing momentum!

# How Many Bins are Necessary?

We can compute the Fisher Info for a histogram of  $p_{T,W}$  starting with bins at 0, 150, and 250 GeV

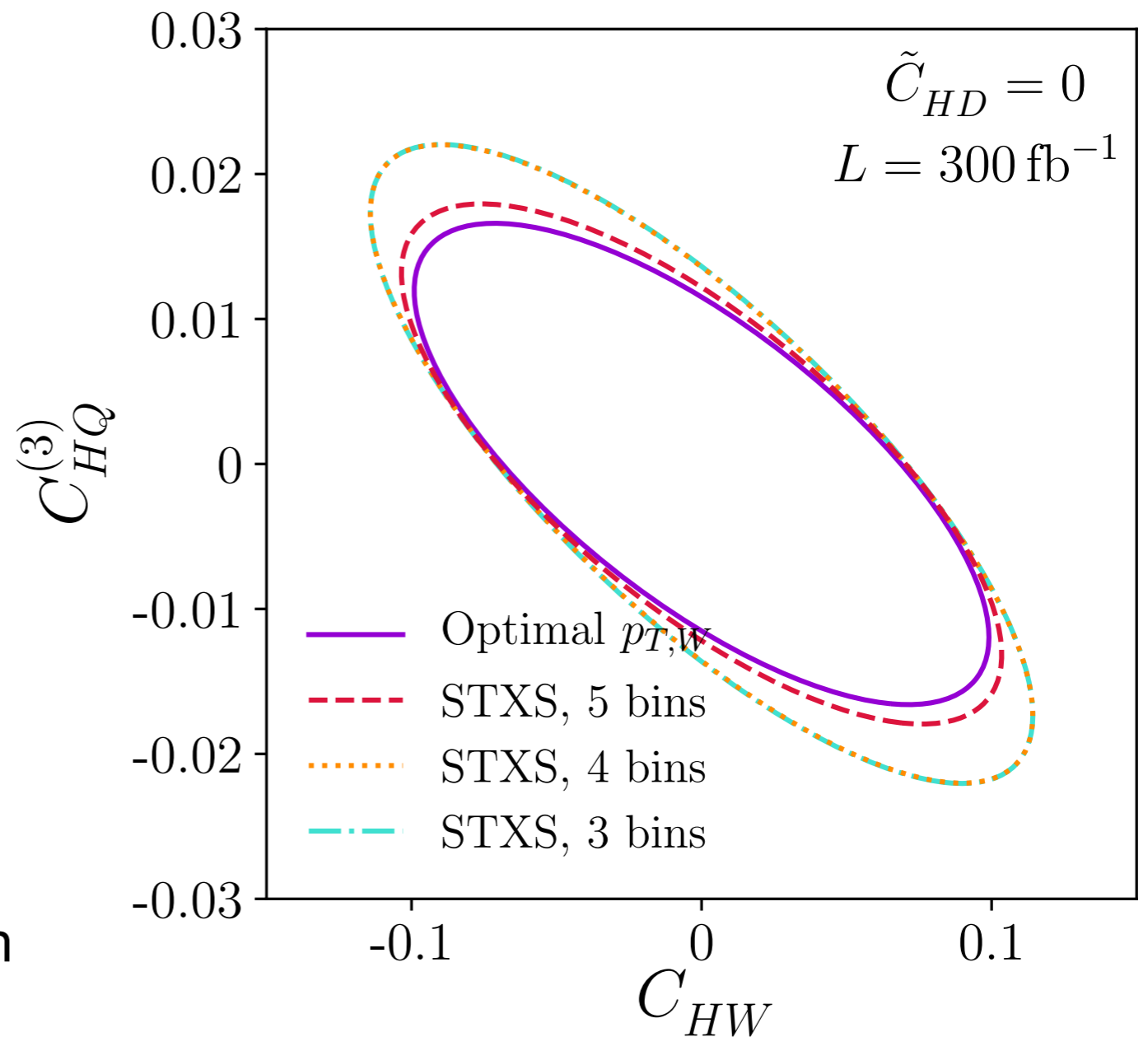
STXS Stage 1 [1610.07922]

Include the additional bins at 75 GeV (4 bins) and 400 GeV (5 bins)

STXS Stage 1.1 [1906.02754]

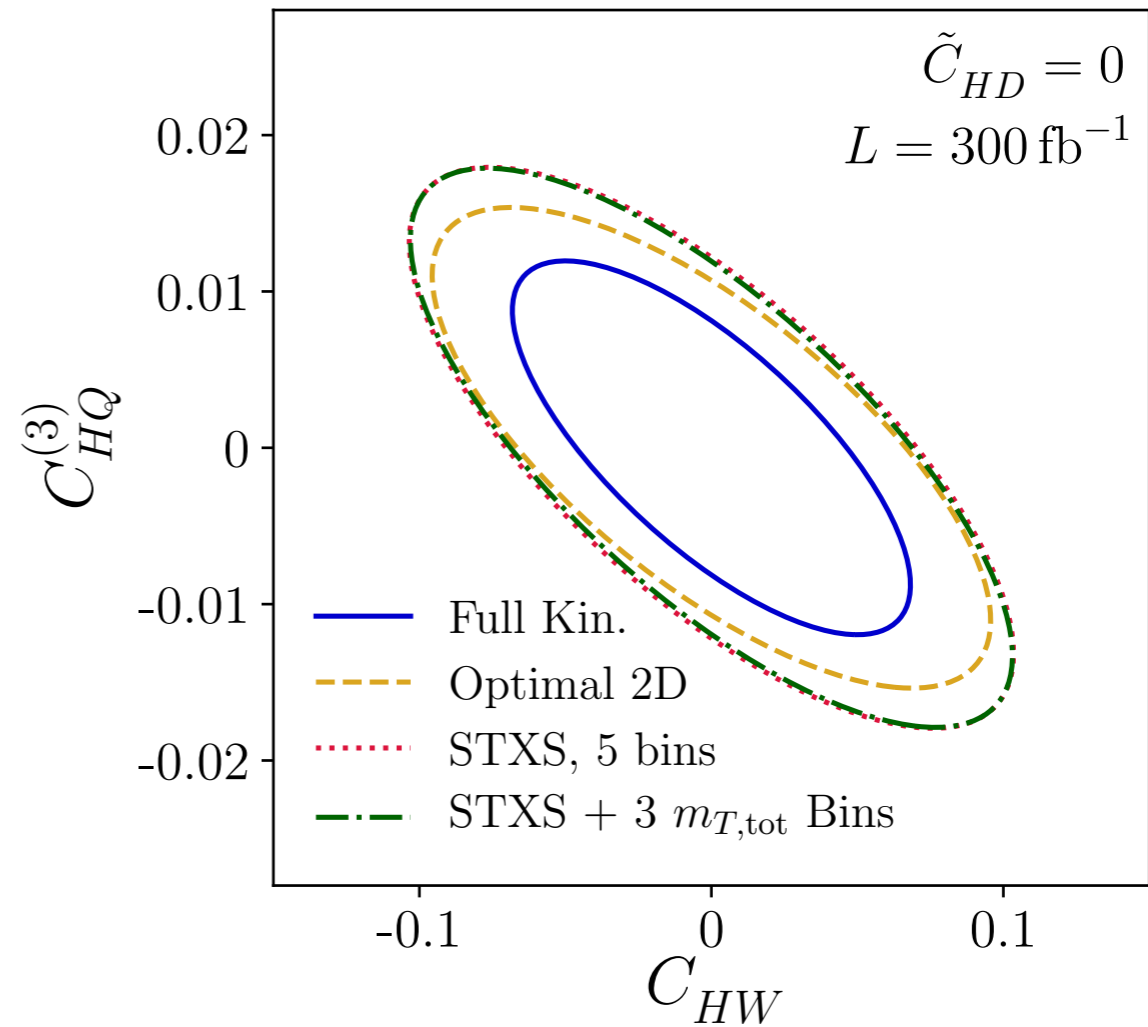
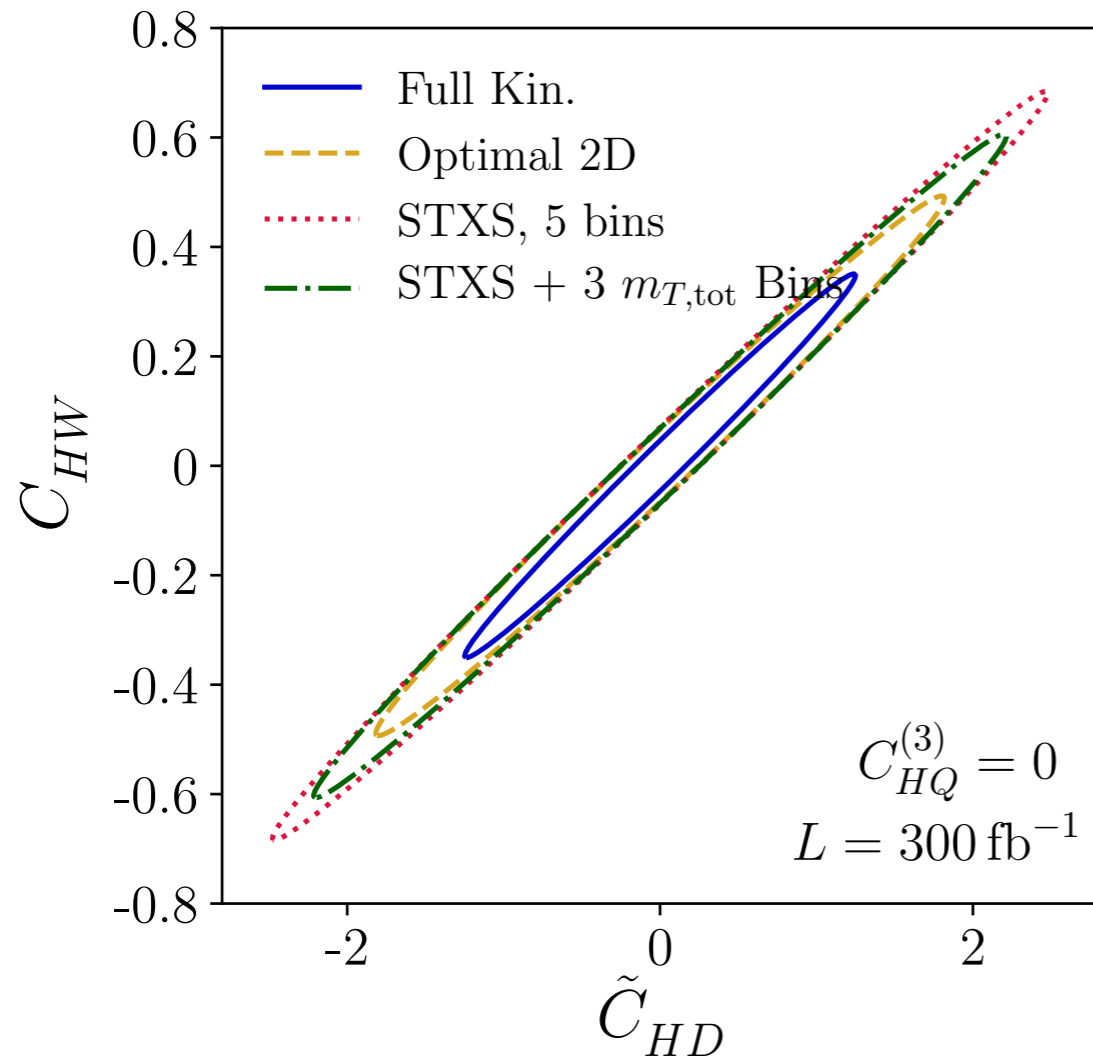
Compare to SALLY Estimator trained on only  $p_{T,W}$

➔ Additional high momentum bin essential for constraining  $C_{Hq}^{(3)}$



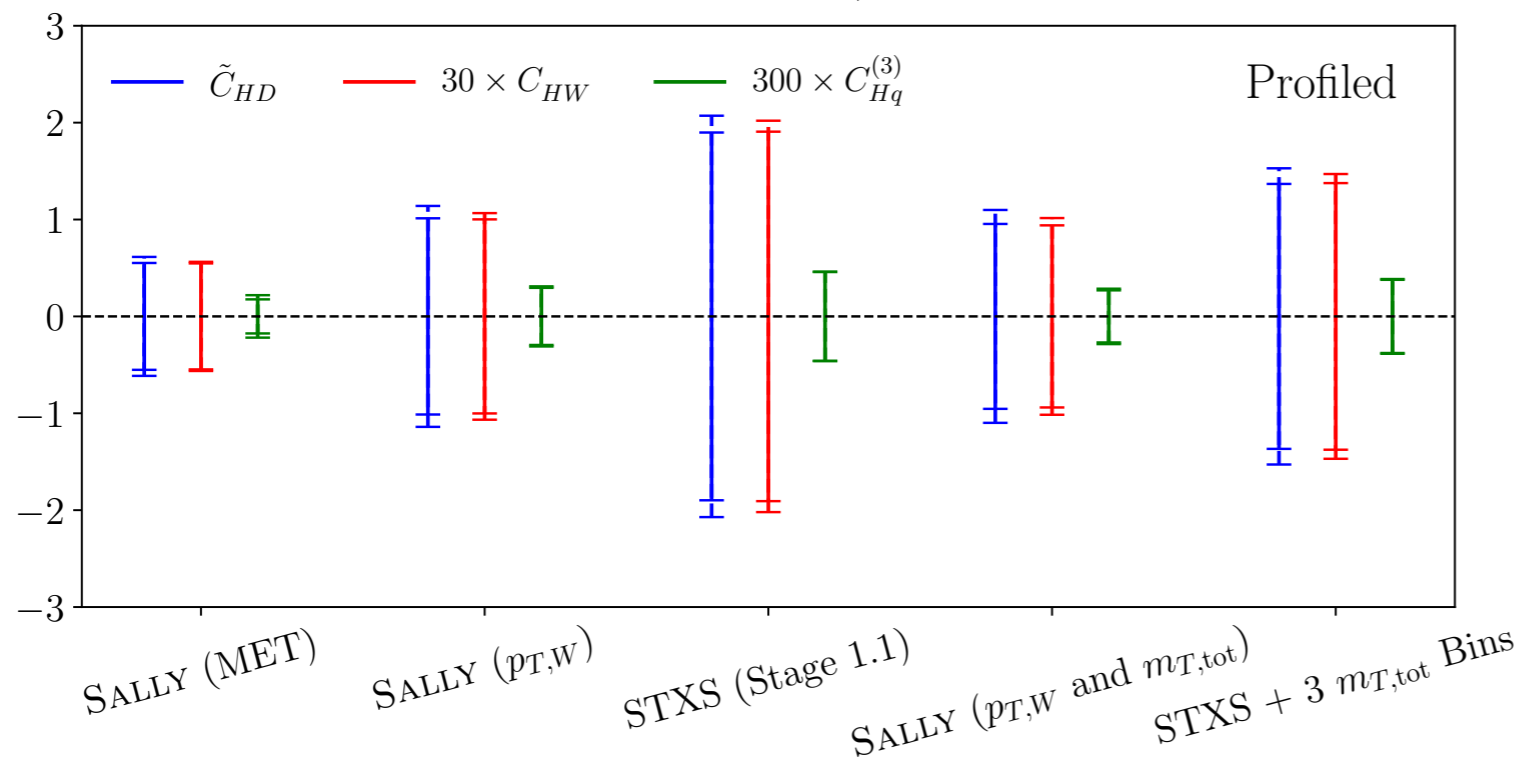
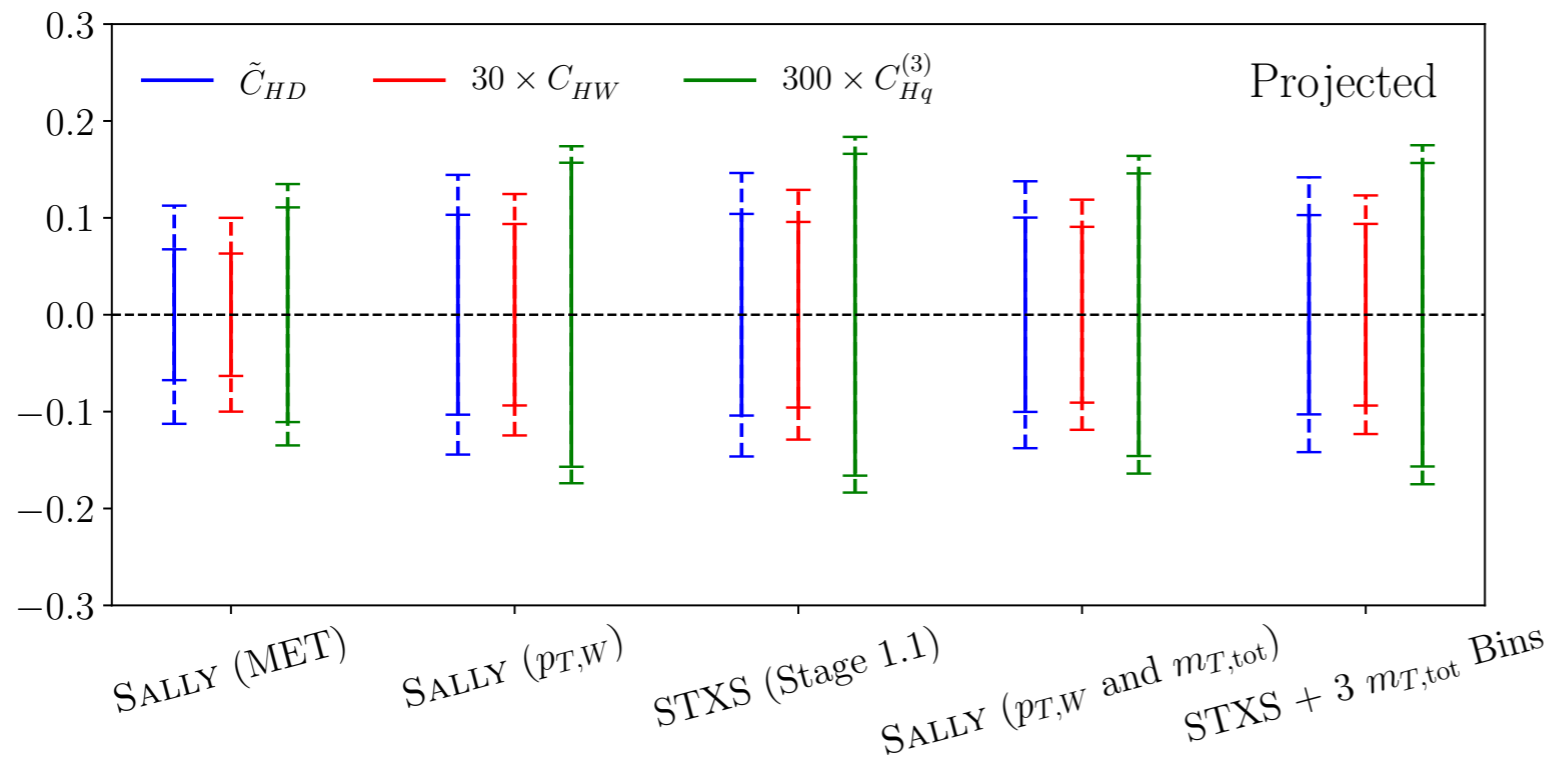
# Is One Observable Enough?

Compare STXS to Estimator trained on 2 Observables & 2D Histogram



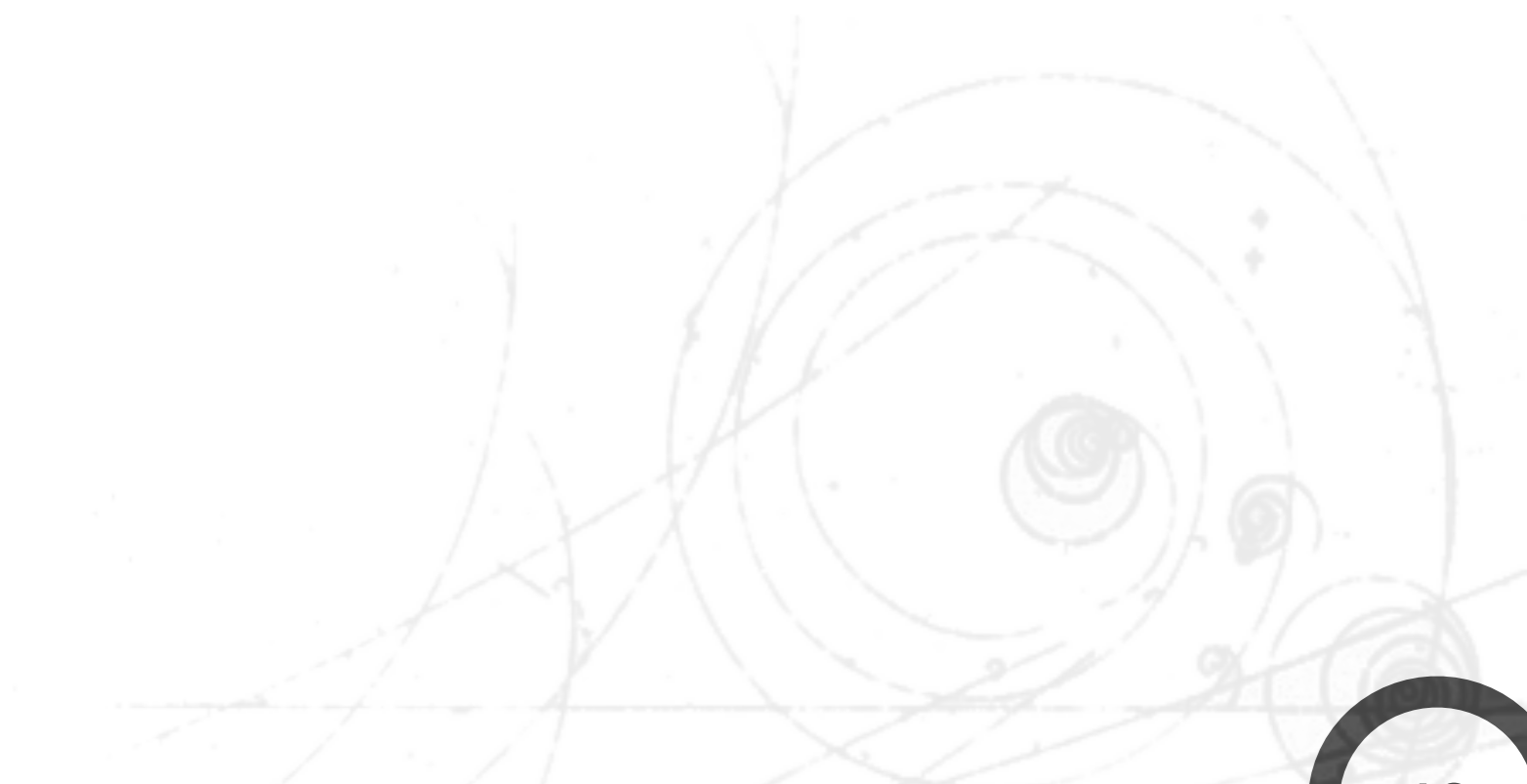
Additional observables help discriminate against background

# Conclusions



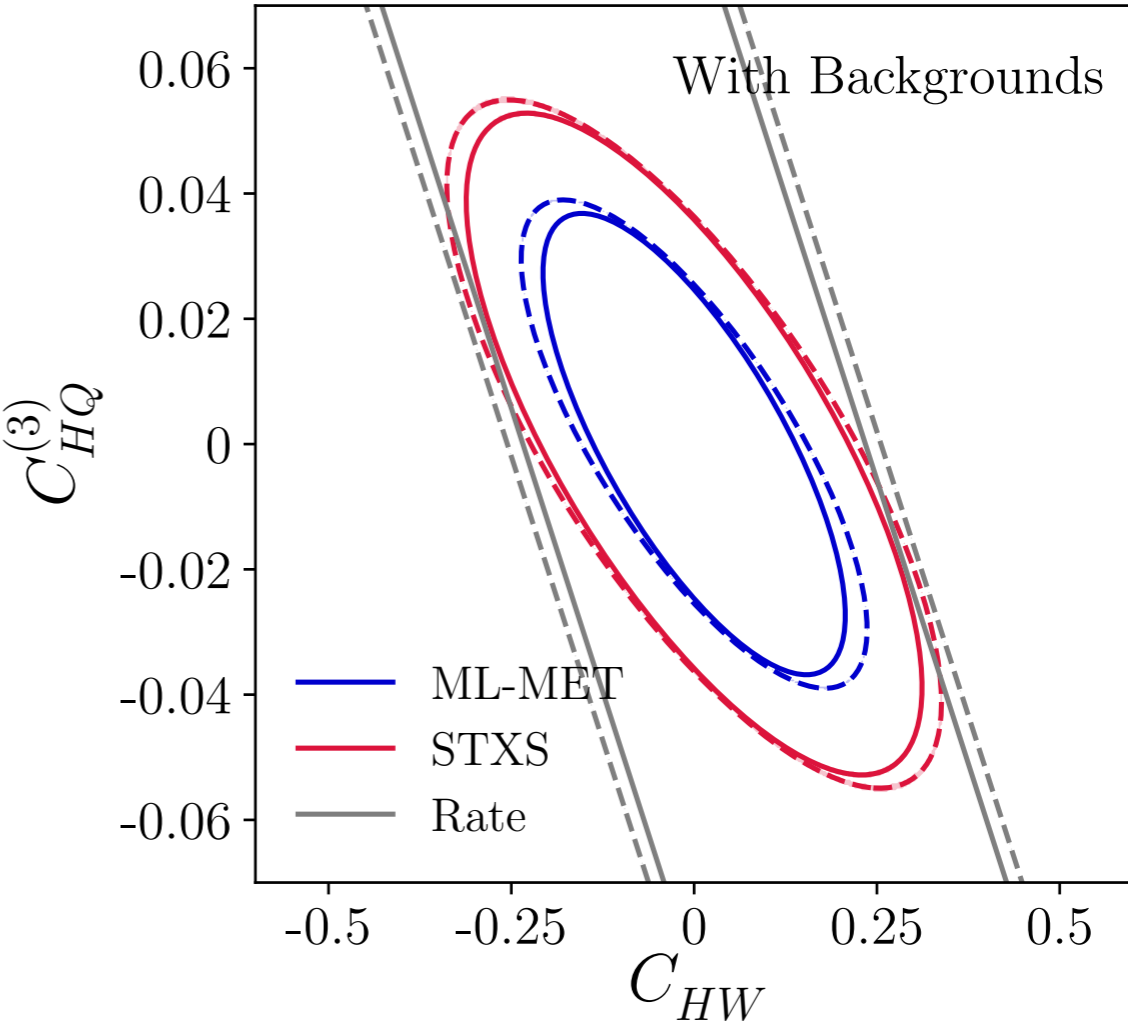
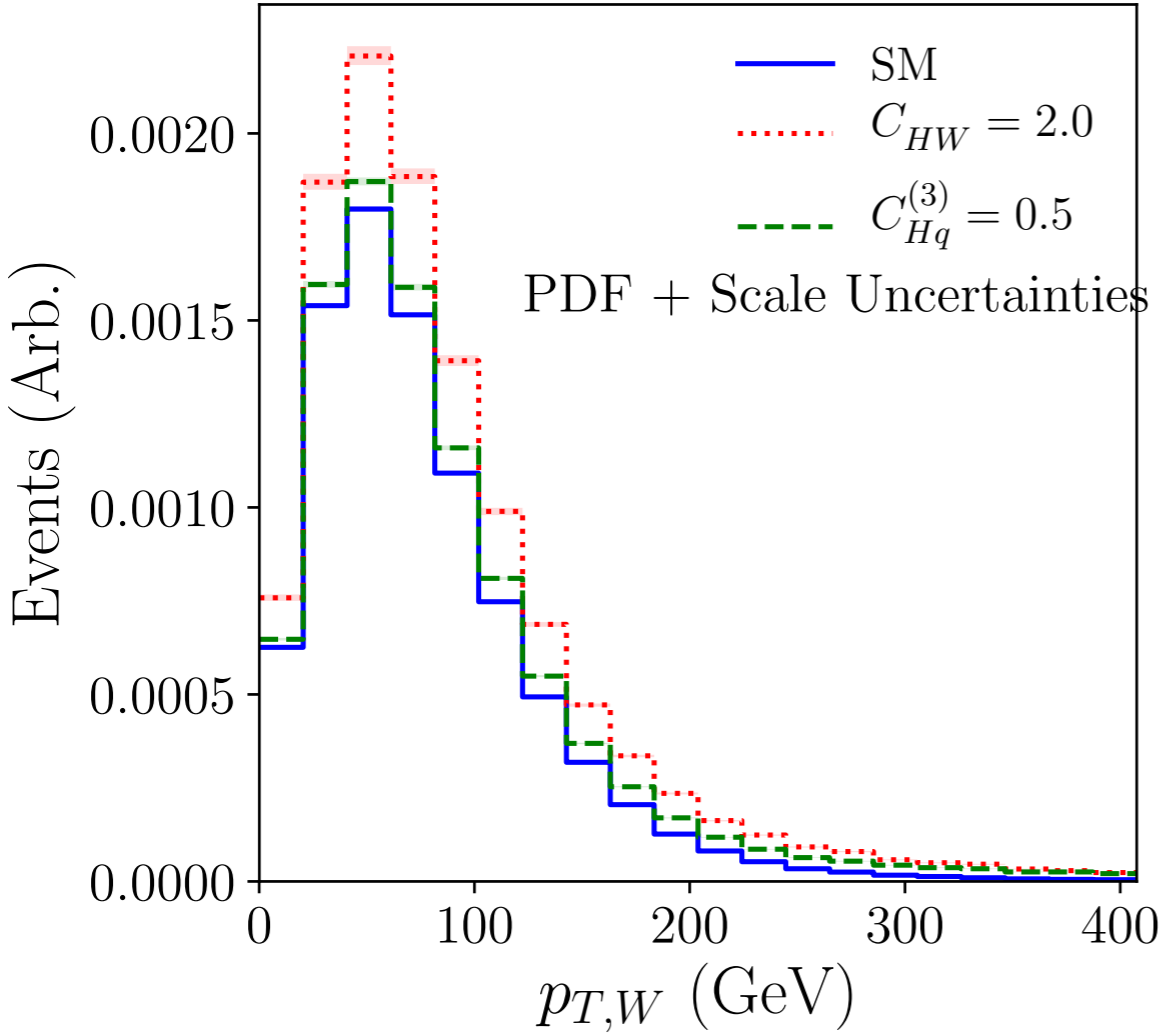
- Our analysis captures effects at order  $1/\Lambda^2$
- Momentum-enhanced 4-point interactions require high momentum bins to constrain
- Can be further constrained with additional observables
- These effects are well understood in the context of the Fisher Information with the help of MadMiner

# Backup



# Including Systematics

Scale & PDF Uncertainties can be treated as additional (nuisance) parameters



We include these for the WH Signal and marginalize over them in the Fisher Information.