

# Multi-Boson interactions 2019

## Summary Talk

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# Outline

## 1 Theory

- Standard Model
- Understanding Cross Sections
  - Diboson Production
  - VBF and VBS
  - Polarization
- Beyond the SM
  - EFTs
  - Other BSM strategies?

## 2 Experimental Searches

- Diboson Production
- VBF and VBS
- Polarization

## 3 Constraints on aTGCs

## 4 Constraints on aQGCs

## 5 Future sensitivity

## 6 Epilogue

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## 1 Theory

- Standard Model
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## 3 Constraints on aTGCs

## 4 Constraints on aQGCs

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## 6 Epilogue





1. *Journal of the American Medical Association*, 1997; 277: 1039-1043.

$$\begin{aligned} \frac{\mathcal{L}_{WWV}}{g_{WWV}} &= ig_1^V \left( W_{\mu\nu}^\dagger W^\mu V^\nu - W_\mu^\dagger V_\nu W^{\mu\nu} \right) + i\kappa_V W_\mu^\dagger W_\nu V^{\mu\nu} \\ &+ \frac{i\lambda_V}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu V^{\nu\rho} - g_4^V W_\mu^\dagger W_\nu (\partial^\mu V^\nu + \partial^\nu V^\mu) \\ &+ g_5^V \epsilon^{\mu\nu\rho\sigma} (W_\mu^\dagger \overleftrightarrow{\partial}_\rho W_\nu) V_\sigma + i\tilde{\kappa}_V W_\mu^\dagger W_\nu \tilde{V}^{\mu\nu} \\ &+ \frac{i\tilde{\lambda}_V}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu \tilde{V}^{\nu\rho} \end{aligned}$$

<sup>1</sup>K. Hagiwara, R. D. Peccei, D. Zeppenfeld and K. Hikasa, Nucl. Phys. B **282** (1987) 253.

# Anomalous Triple Gauge Couplings (aTGCs)

Historically, aTGCs parametrized by a general Lorentz invariant Lagrangian

$$\begin{aligned}
 \frac{\mathcal{L}_{WWV}}{g_{WWV}} &= ig_1^V \left( W_{\mu\nu}^\dagger W^\mu V^\nu - W_\mu^\dagger V_\nu W^{\mu\nu} \right) + i\kappa_V W_\mu^\dagger W_\nu V^{\mu\nu} \\
 &+ \frac{i\lambda_V}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu V^{\nu\rho} - g_4^V W_\mu^\dagger W_\nu (\partial^\mu V^\nu + \partial^\nu V^\mu) \\
 &+ g_5^V \epsilon^{\mu\nu\rho\sigma} (W_\mu^\dagger \overleftrightarrow{\partial}_\rho W_\nu) V_\sigma + i\tilde{\kappa}_V W_\mu^\dagger W_\nu \tilde{V}^{\mu\nu} \\
 &+ \frac{i\tilde{\lambda}_V}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu \tilde{V}^{\nu\rho}
 \end{aligned}$$

where

$V = V^\dagger = Z, \gamma$ ,  $V_{\mu\nu} = \partial_\mu V_\nu - \partial_\nu V_\mu$ ,  $\tilde{V}_{\mu\nu} = \frac{1}{2}\epsilon_{\mu\nu\rho\sigma} V^{\rho\sigma}$ .  $W^\mu \equiv W^{\mu-}$  and  $W_{\mu\nu} = \partial_\mu W_\nu - \partial_\nu W_\mu$   
 $\lambda_V, \tilde{\lambda}_V$  are dimension-6 operator couplings

$2 \times 7 = 14$  parameters

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$$\begin{aligned} \frac{\mathcal{L}_{WWV}}{g_{WWV}} &= ig_1^V \left( W_{\mu\nu}^\dagger W^\mu V^\nu - W_\mu^\dagger V_\nu W^{\mu\nu} \right) + i\kappa_V W_\mu^\dagger W_\nu V^{\mu\nu} \\ &+ \frac{i\lambda_V}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu V^{\nu\rho} - g_4^V W_\mu^\dagger W_\nu (\partial^\mu V^\nu + \partial^\nu V^\mu) \\ &+ g_5^V \epsilon^{\mu\nu\rho\sigma} (W_\mu^\dagger \overleftrightarrow{\partial}_\rho W_\nu) V_\sigma + i\tilde{\kappa}_V W_\mu^\dagger W_\nu \tilde{V}^{\mu\nu} \\ &+ \frac{i\tilde{\lambda}_V}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu \tilde{V}^{\nu\rho} \end{aligned}$$

Some terms in this parametrization explicitly break gauge invariance that we know it is valid from LEP.



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Better formulated as deviations from SM TGCs

**SM tree level:**  $g_1^V = \kappa_V = 1$ ,  $\lambda_V = \tilde{\lambda}_V = g_4^V = g_5^V = \tilde{\kappa}_V = 0$   
 $g_{WW\gamma} = e$ ,  $g_{WWZ} = e \cot \theta_W$

# GBET and Unitarity

## Goldstone Boson Equivalence Theorem (GBET)<sup>1</sup>

At High Energies (HE) relative to the  $W$ -mass, massive gauge bosons  $W^\pm, Z$  can be replaced by the corresponding Goldstone Bosons  $G^\pm, G^0$  in scattering processes.

$$S[W_L^\pm, \text{physical}] = i^n \times S[G^\pm, \text{physical}]$$

Dynamics of  $V_L$ s is directly linked to GBET and  $SU(2)_L \times U(1)_Y$  invariance restoration at HE.

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<sup>1</sup>M. S. Chanowitz and M. K. Gaillard, Nucl. Phys. B **261** (1985) 379;  
G. J. Gounaris, R. Kogerler and H. Neufeld, Phys. Rev. D **34** (1986) 3257.

# GBET and Unitarity

## Tree Level Unitarity<sup>1</sup>

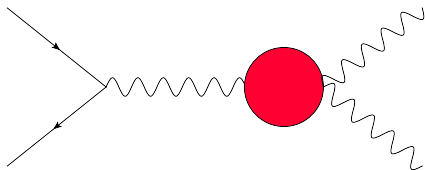
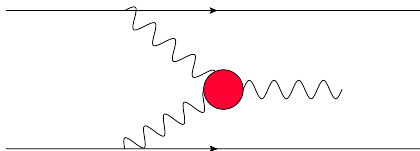
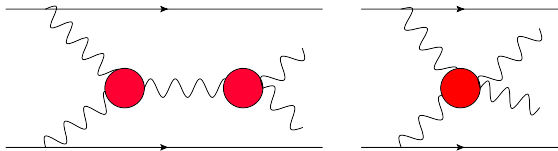
“ the N-particle S-matrix elements in the tree approximation must grow no more rapidly than  $E^{4-N}$  in the limit of HE, at fixed non-zero angles ”

In the SM all multi Goldstone Boson interactions do not contain momenta and therefore leading  $s$ -behaviour cancels against  $s, t$  and  $u$  exchanges of vector and Higgs-boson mediated amplitudes.

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<sup>1</sup>J. M. Cornwall, D. N. Levin and G. Tiktopoulos, Phys. Rev. D **10**, 1145 (1974);  
 B. W. Lee, C. Quigg and H. B. Thacker, Phys. Rev. D **16**, 1519 (1977);  
 C. E. Vayonakis, Lett. Nuovo Cim. **17**, 383 (1976).

# Processes

*Diboson Production**Vector Boson Fusion (VBF)**Vector Boson Scattering (VBS)*

# Diboson Production

- Experimental uncertainties have been surpassed the 10% uncertainty so we have to go to few percent theoretically.
- All NNLO QCD corrections have been completed! They are included in a code (MATRIX+OneLoops)<sup>2</sup>
- Excellent agreement between NNLO and data
- NNLO QCD + NLO EW corrections<sup>3</sup> giant K-factors in observables for  $ZZ, WW, WZ$
- At HE, NLO EW/LO = -40% - 50%.
- One must set jet-veto at high- $s$  since the process is driven away from aTGC searches

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<sup>2</sup>Review talk by Marius Wiesemann

<sup>3</sup>Review talk by Jonas Lindert

# VBF and VBS

- $pp \rightarrow e^+ \nu_e \mu^+ \mu^- jj$  at NLO EW/QCD corrections for  $W^+ Z$  scattering at the LHC: corrections of order  $O(\alpha_s \alpha^6)$  and  $O(\alpha^7)$  the latter being large  $\sim -17.5\%$  especially at high  $p_{Tj}$  (Sudakov enhancement)<sup>4</sup>
- Progress in QCD effects in borderline between perturbative and non-perturbative regime dictated by the factorization theorem discussed<sup>5</sup>

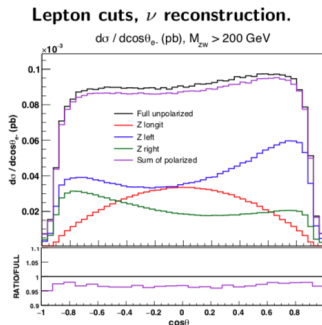
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<sup>4</sup>Talk by Christopher Schwan

<sup>5</sup>Talk by Simon Plaetzer

# Polarization

- Important to provide accurate theory predictions for polarized VBS for LHC analyses
- A nice formula but assumes no lepton cuts so interference effects vanish<sup>6</sup>



<sup>6</sup>Talk by Giovanni Pelliccioli

# EFTs

Systematic deviations from the SM can be studied effectively within EFT as long as the scale of New Physics,  $\Lambda$ , is  $\Lambda^2 \gg M_W^2$  **and**  $s = (p_1 + p_2)^2$ .

A particularly interesting EFT scenario is SMEFT.

**SMEFT:** Assuming there is nothing but the SM below, say  $\Lambda \sim 1$  TeV, and the Higgs field belongs to the SU(2)-doublet, SM is augmented with high dimensional gauge invariant operators

“Warsaw basis”<sup>7</sup> is a non-redundant basis. The complete set of Feynman Rules have been completed for  $d \leq 6$  operators in Unitary and  $R_\xi$ -gauges.<sup>8</sup>

**SmeftFR code**<sup>9</sup> generates FRs and interfaces them to UFO and FeynArts for various event generator and symbolic calculations.

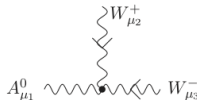
<sup>7</sup>B. Grzadkowski, M. Iskrzynski, M. Misiak and J. Rosiek, arXiv:1008.4884

<sup>8</sup>AD, W. Materkowska, M. Paraskevas, J. Rosiek and K. Suxho, arXiv:1704.03888

<sup>9</sup>AD, M. Paraskevas, J. Rosiek, K. Suxho and L. Trifyllis, arXiv:1904.03204.



# $\gamma W^+ W^-$ in SMEFT with $d \leq 6$ operators



$$\begin{aligned}
 & + \frac{i\bar{g}\bar{g}'}{\sqrt{\bar{g}^2 + \bar{g}'^2}} (\eta_{\mu_1\mu_2} (p_1 - p_2)^{\mu_3} + \eta_{\mu_2\mu_3} (p_2 - p_3)^{\mu_1} + \eta_{\mu_3\mu_1} (p_3 - p_1)^{\mu_2}) \\
 & - \frac{6i\bar{g}'}{\sqrt{\bar{g}^2 + \bar{g}'^2}} C^W (p_3^{\mu_1} p_1^{\mu_2} p_2^{\mu_3} - p_2^{\mu_1} p_3^{\mu_2} p_1^{\mu_3} + \eta_{\mu_1\mu_2} (p_1^{\mu_3} p_2 \cdot p_3 - p_2^{\mu_3} p_1 \cdot p_3) \\
 & + \eta_{\mu_2\mu_3} (p_2^{\mu_1} p_1 \cdot p_3 - p_3^{\mu_1} p_1 \cdot p_2) + \eta_{\mu_3\mu_1} (p_3^{\mu_2} p_1 \cdot p_2 - p_1^{\mu_2} p_2 \cdot p_3)) \\
 & + \frac{i\bar{g}^2 v^2}{(\bar{g}^2 + \bar{g}'^2)^{3/2}} C^{\varphi WB} \left( \eta_{\mu_1\mu_2} (\bar{g}^2 p_1^{\mu_3} + \bar{g}'^2 p_2^{\mu_3}) + \eta_{\mu_2\mu_3} (\bar{g}'^2 p_3^{\mu_1} - \bar{g}^2 p_2^{\mu_1}) \right. \\
 & \left. + \eta_{\mu_3\mu_1} (-\bar{g}'^2 p_3^{\mu_2} - \bar{g}^2 p_1^{\mu_2}) \right) \\
 & - \frac{2i\bar{g}'}{\sqrt{\bar{g}^2 + \bar{g}'^2}} C^{\widetilde{W}} \left( \epsilon_{\mu_1\mu_2\mu_3\alpha_1} (p_1^{\alpha_1} p_2 \cdot p_3 + p_2^{\alpha_1} p_1 \cdot p_3 + p_3^{\alpha_1} p_1 \cdot p_2) \right. \\
 & + \epsilon_{\mu_1\mu_2\alpha_1\beta_1} (p_1 - p_2)^{\mu_3} p_1^{\alpha_1} p_2^{\beta_1} + \epsilon_{\mu_2\mu_3\alpha_1\beta_1} (p_2 - p_3)^{\mu_1} p_2^{\alpha_1} p_3^{\beta_1} \\
 & \left. + \epsilon_{\mu_3\mu_1\alpha_1\beta_1} (p_3 - p_1)^{\mu_2} p_3^{\alpha_1} p_1^{\beta_1} \right) \\
 & + \frac{i\bar{g}^2 v^2}{\sqrt{\bar{g}^2 + \bar{g}'^2}} C^{\varphi \widetilde{W} B} \epsilon_{\mu_1\mu_2\mu_3\alpha_1} p_1^{\alpha_1}
 \end{aligned}$$

“Triboson” parameters in EFT with  $d \leq 6$  operators:

CP-invariant parameters : 5 new - 2 constraints = 3

CP-violating parameters = 2

One should be very careful with EFTs !

- $s, M_W^2 \ll \Lambda^2$  – otherwise using EFT is nonsense!
- Bounds on Wilson coefficients are given in particular **operator basis** (e.g. Warsaw, SILH, etc)
- Bounds on Wilson coefficients should be given in a particular **input scheme** e.g. one has to trade  $\bar{g}, \bar{g}', v$  with measurable quantities like for example<sup>10</sup>
  - $\{\alpha_{em}, m_Z, G_F, m_h, m_t, \dots\}$ -scheme
  - $\{m_W, m_Z, G_F, m_h, m_t, \dots\}$ -scheme
  - ....

<sup>10</sup>I. Brivio and M. Trott, arXiv:1701.06424

# SMEFT@NLO

- The S-matrix must be renormalization scale independent (up to a fixed order in loop and EFT expansion).
- in SMEFT one has to include running from all operators so that the result is gauge invariant
- Overview of NLO calculations in SMEFT<sup>11</sup>. Several observable studies fully at 1-loop in SMEFT.
- A first step towards SMEFT@NLO QCD has been done. NLO EW fully automated?

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<sup>11</sup>Review talk by Cen Zhang

# Non-linear EFT

- Non Linear EFT<sup>12</sup> : the Higgs field **does not** reside in the  $SU(2)_L$  doublet
- EFT is now more involved.  $d = 8$  operators are promoted to  $d = 6$  operators that reach easily unitarity bound,  $|a_J(s)| = 1.$ , because for example VBS amplitudes grow with  $A \sim s^2$
- There must be a cut-off at scales much before the typical  $\Lambda^2$ : there are unitarization suggestions

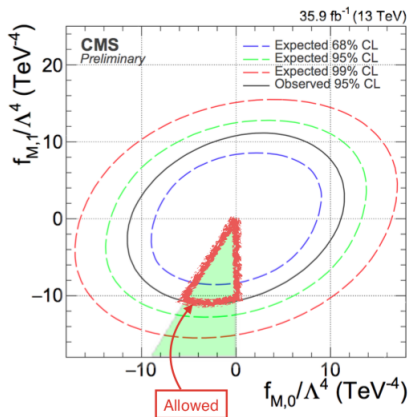
**My opinion: a prototype model (if it exists) could be used as a benchmark to be used in experimental studies**

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<sup>12</sup>Talk by Rafael Lopez Delgado

**Positivity constraints:** In every QFT based on analyticity, unitarity, Lorentz invariance there are certain bounds on certain linear combinations of  $d = 8$  operators in SMEFT<sup>13</sup>

$$\sum_i C_i^{(8)} x_i \geq 0$$



<sup>13</sup>Talk by Cen Zhang

## BSM Models and EFTs<sup>14</sup>

- SMEFT at HE: there are 4 parameters that grow with  $\sigma \sim s$  in diboson processes. Observables are WW, WZ, HW, HZ. Possible to probe "weak coupling regime" where  $g_*^2/M^2$  with  $g^* \sim g \lesssim 0.1\%$ .
- Currently we are probing the strong coupling regime  $g_* \sim (4\pi)$  in diboson searches.
- A composite model related to TGCs presented.
- One way to compete with LEP precision is by going to HE and study ZH production. Due to GBET we have a contact term that dominates the amplitude far from resonances.<sup>15</sup>

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<sup>14</sup>Review talk by Marc Montull

<sup>15</sup>Talk by Sandeepan Gupta

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## 2 Experimental Searches

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- VBF and VBS
- Polarization

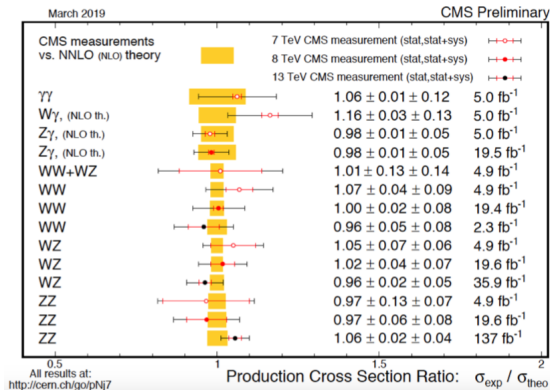
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CMS  $\mathcal{L} = 137 \text{ fb}^{-1}$  (partly Run-II data included) total and differential Xsections. Report <sup>16</sup> on  $WW, WZ, ZZ$  leptonic final states as well as on

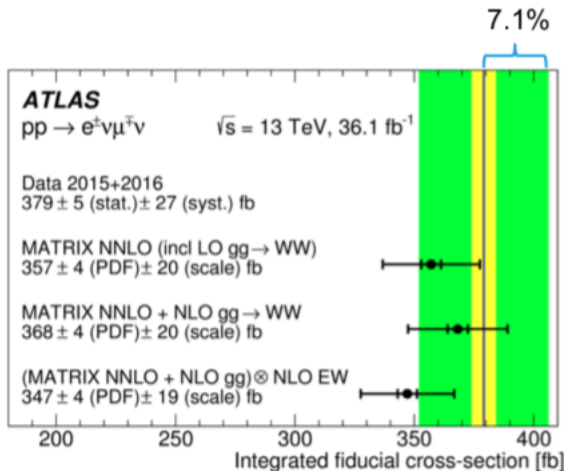


Precision of 5% reached.

<sup>16</sup>Talk by Alicia Calderon



# ATLAS fiducial and diff Xsections Reports on WW and ZZ<sup>16</sup>

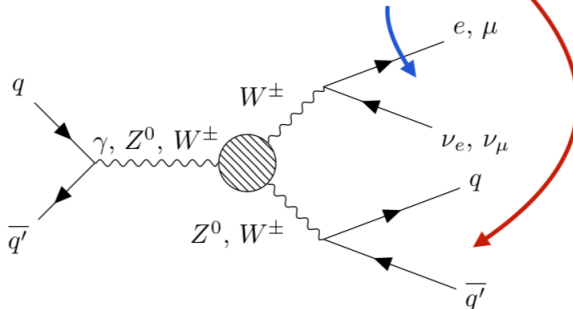


<sup>16</sup>Talk by Valerie Lang

# Measurement of Diboson production in semileptonic decay modes and anomalous Couplings<sup>17</sup>

## Semileptonic:

Diboson final state **with one V decay to quarks ( $\rightarrow$  jets)**,  
and **other V decay to final state with lepton ( $e, \mu$ )**



<sup>17</sup>Review talk of ATLAS and CMS results by Robin Cameron Aggleton

Measurement of Diboson production in **semileptonic** decay modes and anomalous Couplings<sup>17</sup>

Direct Searches for NP in diboson searches: limits on KK gravitons,  $V'$  and  $W'$  masses  $\geq 1.5$  TeV.<sup>18</sup>

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<sup>17</sup>Review talk of ATLAS and CMS results by Robin Cameron Aggleton

<sup>18</sup>Talk by Antonis Agapitos

- $ssWWjj$ ,  $WZjj$  and  $ZZjj$  at  $\sqrt{s} = 13$  TeV seem to have been observed at ATLAS
- Also photons in the final state<sup>19</sup>:  $Z\gamma jj$  (strong evidence for both ATLAS and CMS) or  $\gamma\gamma jj$  (observation in heavy ions by ATLAS)
- Higgs VBF and Z/W VBF presented<sup>20</sup>. The former agrees with the SM while the latter are dominated by systematics.
- Jet-veto technics based on event-by-event selection may help in VBF/VBS signal/background in SM and BSM searches<sup>21</sup>

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<sup>19</sup>Talk by Narei Lorenzo Martinez

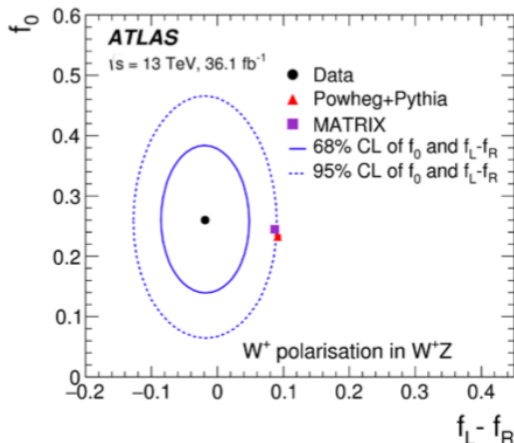
<sup>20</sup>Talk by Dag Gillberg

<sup>21</sup>Talk by Richard Ruiz

# Polarization

Probing GBET directly. ATLAS search for  $WZ$  channel<sup>22</sup>

- $f_0$  is measured different from 0 at more than 3 sigma and in agreement with predictions
- FL-FR at  $2\sigma$  from predictions in  $W^+$



<sup>22</sup>Talk by Corinne Goy

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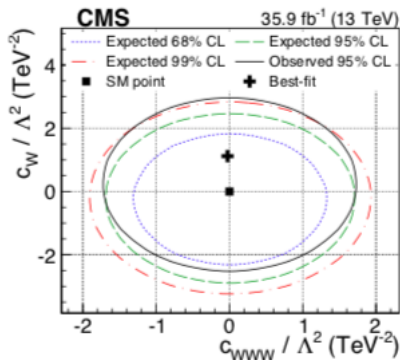
CMS at 13 TeV (1D and 2D limits). Semileptonic final state.<sup>23</sup>. Big improvements w.r.t 8 TeV results

Parametrization	aTGC	Expected limit	Observed limit	Observed best-fit	CMS 8 TeV observed limit
EFT	$c_{WWW}/\Lambda^2$ (TeV $^{-2}$ )	[-1.44, 1.47]	[-1.58, 1.59]	-0.26	[-2.7, 2.7]
	$c_W/\Lambda^2$ (TeV $^{-2}$ )	[-2.45, 2.08]	[-2.00, 2.65]	1.21	[-2.0, 5.7]
	$c_B/\Lambda^2$ (TeV $^{-2}$ )	[-8.38, 8.06]	[-8.78, 8.54]	1.07	[-14, 17]
LEP	$\lambda_Z$	[-0.0060, 0.0061]	[-0.0065, 0.0066]	-0.0010	[-0.011, 0.011]
	$\Delta g_1^Z$	[-0.0070, 0.0061]	[-0.0061, 0.0074]	0.0027	[-0.009, 0.024 ]
	$\Delta \kappa_Z$	[-0.0074, 0.0078]	[-0.0079, 0.0082]	-0.0010	[-0.018, 0.013 ]

<sup>23</sup>Talk by Robin Cameron Aggleton

# Dibosons

CMS at 13 TeV (1D and 2D limits). Semileptonic final state.<sup>23</sup>. Big improvements w.r.t 8 TeV results



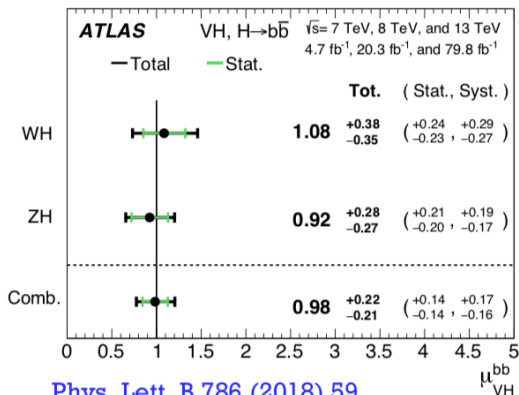
<sup>23</sup>Talk by Robin Cameron Aggleton



# VH and HH

$V(\ell\ell)H(\bar{b}b)$  has been observed<sup>24</sup>. HH is currently being searched for.

At High  $p_T$  is an interesting probe for NP (see theory talks on EFTs)



[Phys. Lett. B 786 \(2018\) 59](#)

<sup>24</sup>Talk by Stephane Cooperstein

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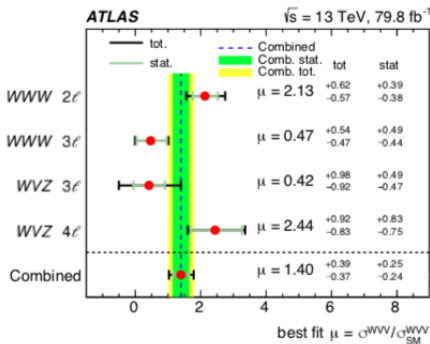
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<sup>24</sup>Talk by Stephane Cooperstein

# Triboson Production

There is evidence for triboson production at ATLAS<sup>25</sup>.

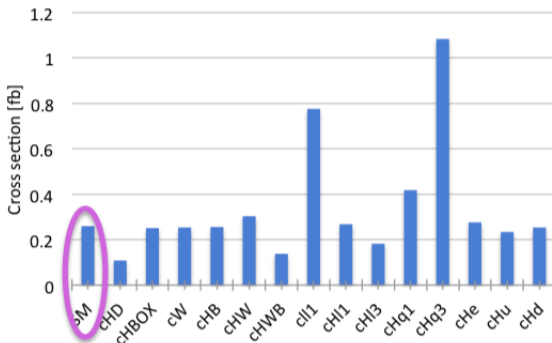


Triboson searches are under investigation also at CMS.<sup>26</sup>

<sup>25</sup>Talk by Andrea Sciandra

<sup>26</sup>Talk by Miaoyuan Liu

# Effects on VBS Xsection from SMEFT parameters (Warsaw basis) in WZ production<sup>27</sup>



<sup>27</sup>Talk by Despoina Sampsonidou

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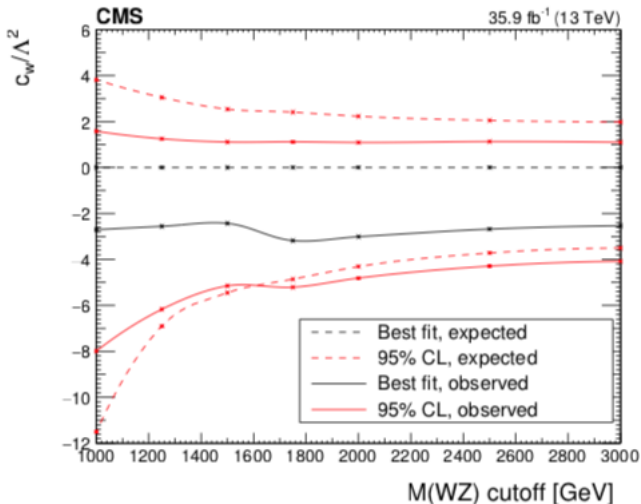
Limits on  $d = 8$  Wilson coefficients in VBS (CMS, EWK VV production, semileptonic mode)<sup>28</sup>

	Obs Low	Obs High	Exp Low	Exp High
$F_{S,0}$	-2.7	2.7	-4.2	4.2
$F_{S,1}$	-3.4	3.4	-5.2	5.2
$F_{M,0}$	-0.69	0.70	-1.0	1.0
$F_{M,1}$	-2.0	-2.1	-3.0	3.0
$F_{M,6}$	-1.3	1.3	-1.4	1.4
$F_{M,7}$	-3.4	3.4	-5.1	5.1
$F_{T,0}$	-0.12	0.11	-0.17	0.16
$F_{T,1}$	-0.12	0.13	-0.18	0.18
$F_{T,2}$	-0.28	0.28	-0.41	0.41

Bounds on  $F/\Lambda^4$  are in  $\text{TeV}^{-4}$  units. These are the best limits so far.

<sup>28</sup>Talk by Andrew Levin

The plot created most discussions!<sup>29</sup>



<sup>29</sup>Talk by Hannes Mildner

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- Polarization in VBS at HL-LHC : projective evidence for LL at few  $\sigma$ s per experiment but for the  $VV \rightarrow V_L V_L$
- But, using machine learning techniques ATLAS+CMS at  $3000 fb^{-1}$  each may reach  $5\sigma$  for  $ssWW$  and VBS ZZ scattering<sup>30</sup>
- HH prospects .....<sup>31</sup>

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<sup>30</sup>Talk by Meng Lu

<sup>31</sup>Talk by Stephane Cooperstein

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- 4 Constraints on aQGCs
- 5 Future sensitivity
- 6 Epilogue**

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