TMD evolution from SIDIS data

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Outlook

I will discuss some of the (TMD) physics in COMPASS measurements

Collins, Sivers, Unpolarized functions (crucial)

Challenges remain. I review some.

TMD evolution hard to see in data (Sivers and Collins effects). Accurate determination of unpolarized functions Issues with normalization Matching between small and large qT

Kinematics of applicability of factorization theorems

How to extract TMDs?

Ingredients:

Data

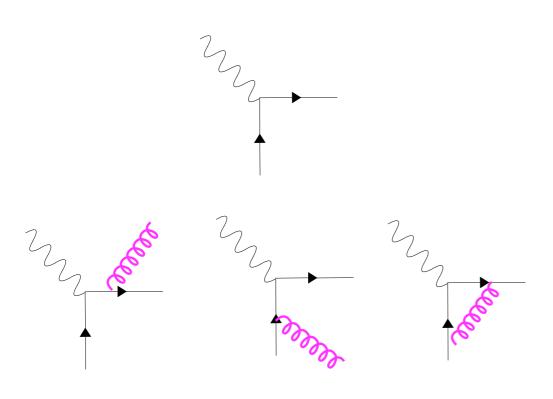
+

Theoretical framework

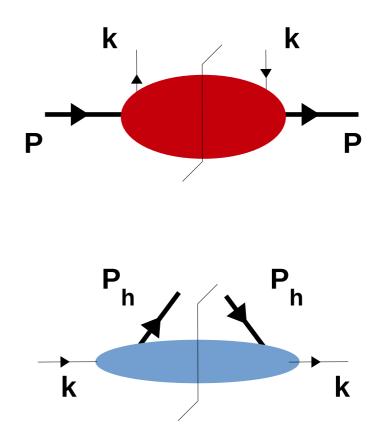
Recipe is tricky though, several challenges.

Theoretical Framework: Factorization theorems

Short distance effects.



Long distance physics



pQCD

Non-perturbative content

Theoretical Framework: Factorization theorems

W (TMD region)

$$\sum_{q} \mathcal{H}_{q} \text{ F.T.} \left\{ \tilde{D}_{h/q}(z, z \boldsymbol{b}_{\perp}; Q) \ \tilde{f}_{q/P}(x, \boldsymbol{b}_{\perp}; Q) \right\}$$

Fourier Transform of:

$$\begin{split} \tilde{F}_{j}(x,b_{T},Q,\zeta_{F}) &= \left(\frac{\sqrt{\zeta_{F}}}{\mu_{b}}\right)^{\tilde{K}(b_{*},\mu_{b})} \sum_{j} \int_{x}^{1} \frac{d\hat{x}}{\hat{x}} \underbrace{\tilde{C}_{ji}^{in}(x/\hat{x},b_{*},\mu_{b},\mu_{b}^{2}) f_{i}(\hat{x},\mu_{b})}_{\times \exp\left\{\int_{\mu_{b}}^{Q} \frac{d\mu}{\mu} \left(\underline{\gamma_{F}(\mu;1)} - \ln\left(\frac{\sqrt{\zeta_{F}}}{\mu}\right) \underline{\gamma_{K}(\mu)}\right)\right\}} \\ &\times \exp\left\{-g_{P}(x,b_{T}) - g_{K}(b_{T})\ln\left(\frac{\sqrt{\zeta_{F}}}{\sqrt{\zeta_{F0}}}\right)\right\}, \end{split}$$

pQCD

Input (extraction from collinear cross section)

Non-perturbative functions to extract from data.

General Strategy: (Possible recipe)

(1) Map kinematical dependencies of different experiments. Simple models, "snap shots" of TMDs.

(2) Study how much information can be inferred on certain effects, e.g. TMD evolution, strictly from data.

(3) Use information from (1), (2) to build full TMD picture (CSS, SCET, other TMD factorization schemes)

- i) Test importance on input information (collinear PDFs & FFs, TMD models).
- ii) Errors of factorization (optimal kinematical regime?).
- iii) Balance between constraints from theory and information obtained from statistical analyses, model comparison.

General Strategy: (Possible recipe)

(1) Map kinematical dependencies of different experiments. Simple models, "snap shots" of TMDs.

(2) Study how much information can be inferred on certain effects, e.g. TMD evolution, strictly from data.

Gaussian model for TMDs in momentum space

Transversity and Collins function.

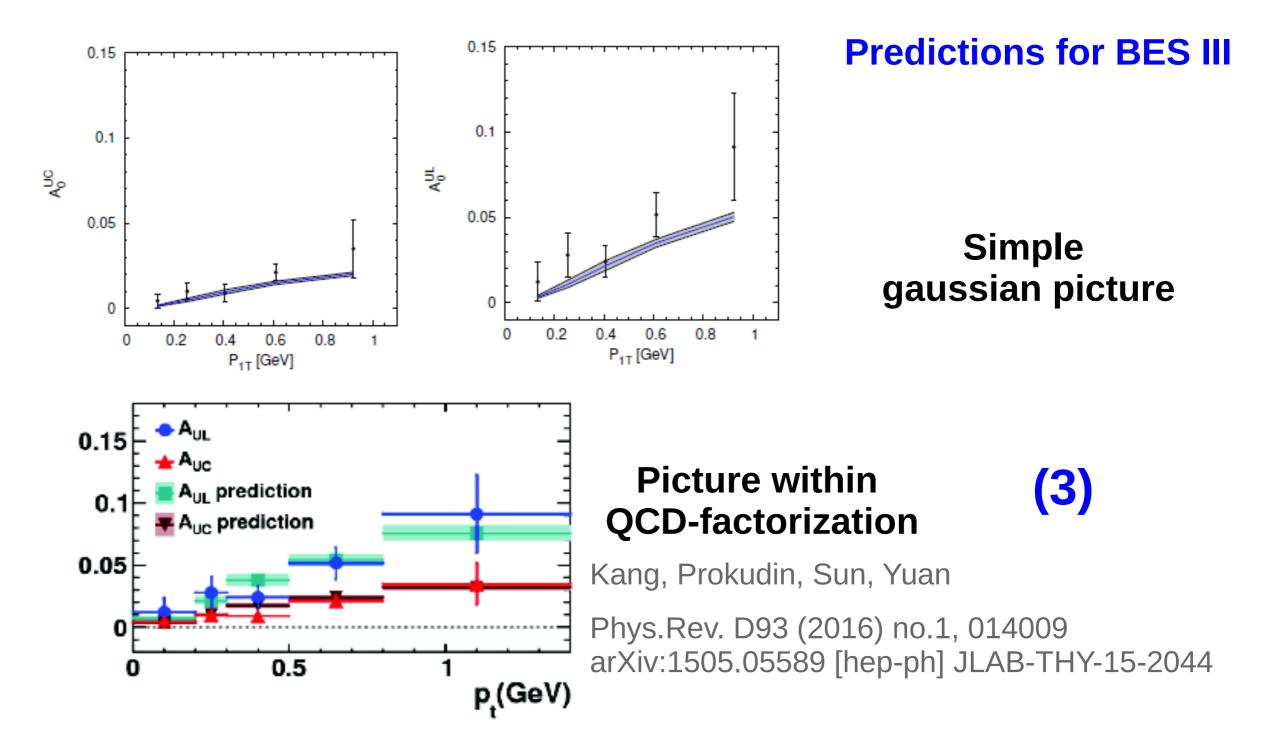
Transversity

(1) & (2)

Collins function

Scale dependence?

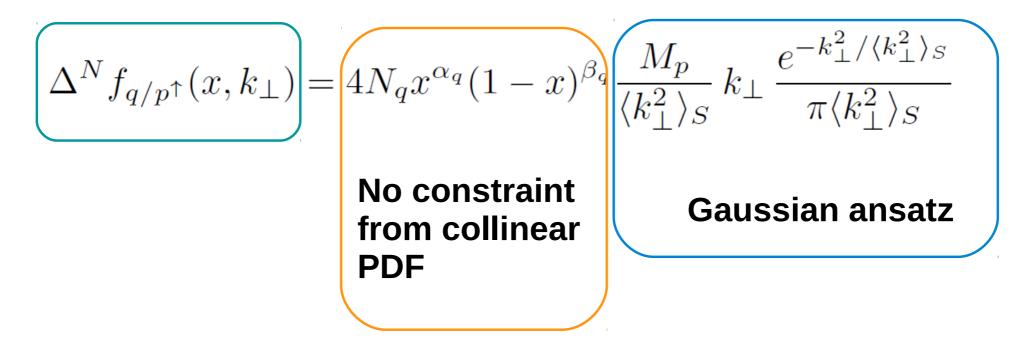
Q^2 = 13 GeV^2



Sivers asymmetry in SIDIS

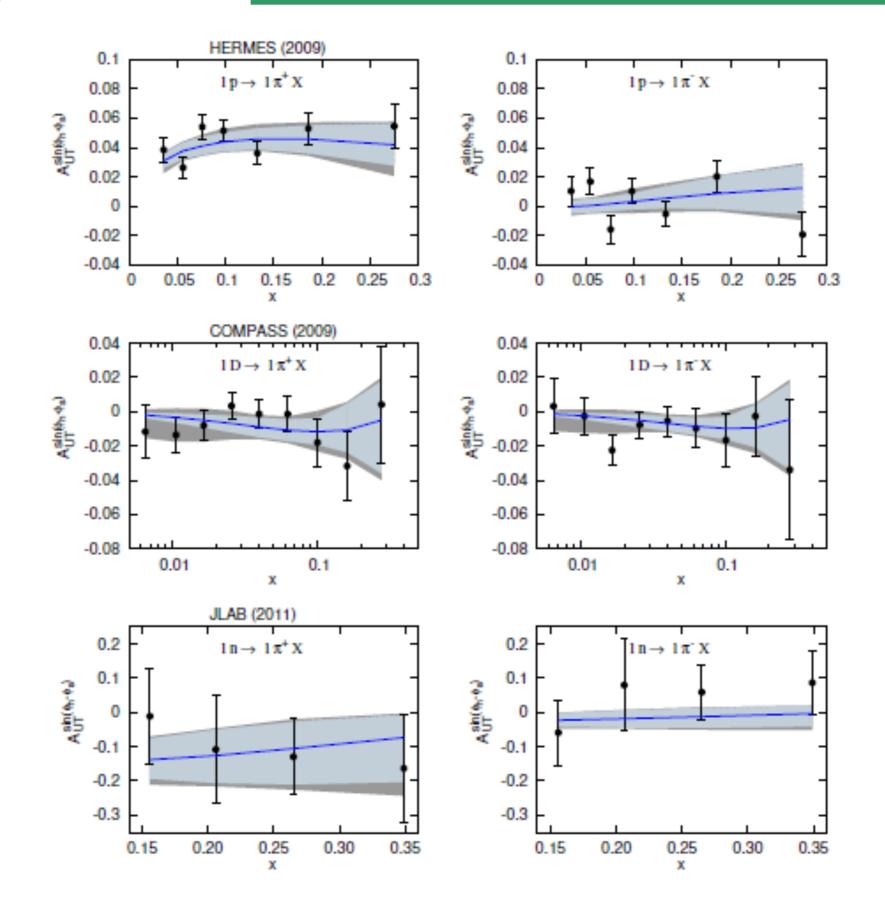
$$A_{UT}^{\sin(\phi_h - \phi_S)} = 2 \frac{\int d\phi_S d\phi_h \left[d\sigma^{\uparrow} - d\sigma^{\downarrow} \right] \sin(\phi_h - \phi_S)}{\int d\phi_S d\phi_h \left[d\sigma^{\uparrow} + d\sigma^{\downarrow} \right]} = \frac{F_{UT}^{\sin(\phi_h - \phi_S)}}{F_{UU}}$$

Generalized Parton Picture (no evolution)



• M. Boglione, U. D'Alesio, C. Flore, JOGH , JHEP 1807 (2018) 148

Sivers asymmetry in SIDIS



Sivers asymmetry in SIDIS

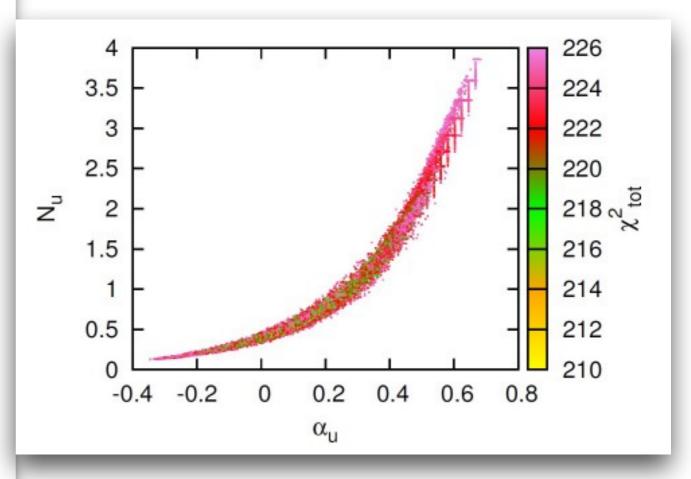
	s = 220
fits (3)	parameters)
$\chi^2_{ m tot}$	$\chi^2_{ m dof}$
408	1.88
914	4.21
• fits $(5$	parameters)
$\chi^2_{ m tot}$	$\chi^2_{ m dof}$
266	1.24
228	1.06
213	0.99
	408 914 • fits (5 χ^2_{tot} 266 228

Is it Reasonable to increase number of parameters?

Sivers asymmetry in SIDIS

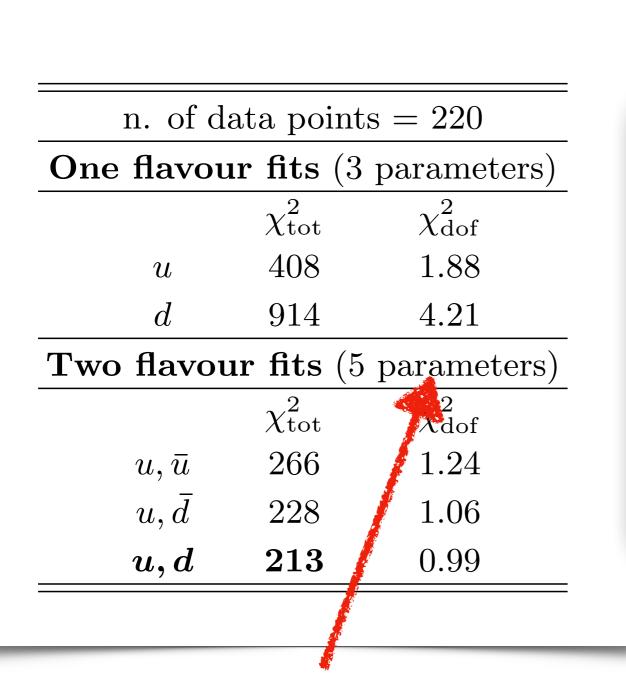
n. of data points $= 220$			
One flavour fits (3 parameters)			
$\chi^2_{ m tot}$	$\chi^2_{ m dof}$		
408	1.88		
914	4.21		
Two flavour fits (5 parameters)			
$\chi^2_{ m tot}$	$\chi^2_{ m dof}$		
266	1.24		
228	1.06		
213	0.99		
	r fits (3 χ^2_{tot} 408 914 r fits (5 χ^2_{tot} 266 228		

One more parameter (per flavor)

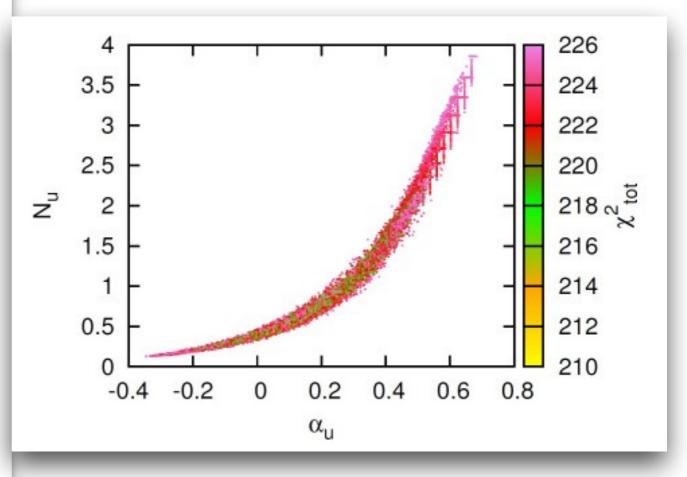


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Sivers asymmetry in SIDIS

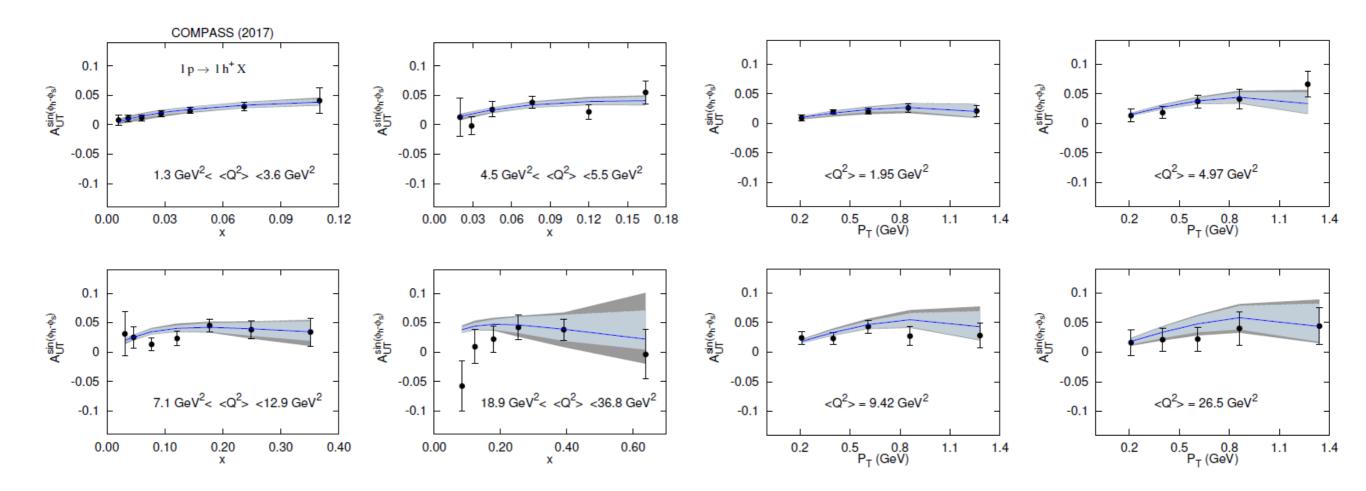


One more parameter (per flavor)



Rough limit on number of parameters (benchmark)

Sivers asymmetry in SIDIS

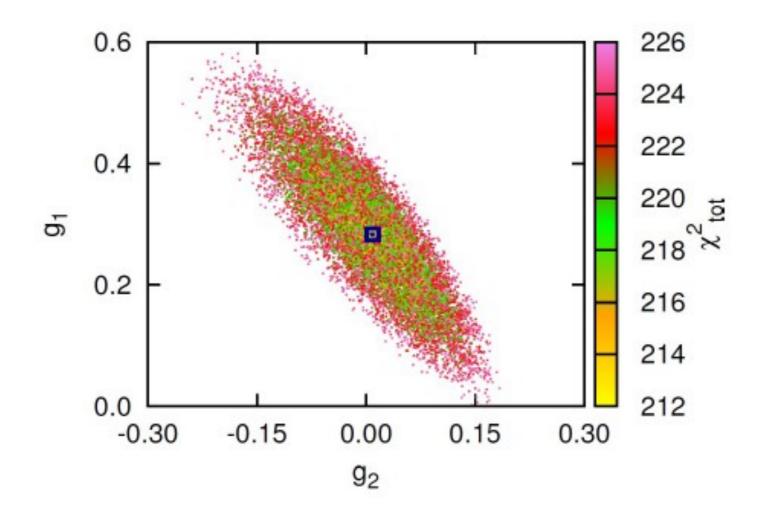


Sivers asymmetry in SIDIS

Signals of scale dependence

$$\langle k_{\perp}^2 \rangle_S = g_1 + g_2 \ln \frac{Q^2}{Q_0^2}$$

g2 here to "mimic" TMD evolution



No "visible" sign TMD evolution, expected: It washes out in the ratio of the asymmetry

Recapitulating so far:

Signals of TMD evolution are not so "visible" in asymmetries (ratios)

Important to look at correlations of parameters.

One may get a rough idea of a reasonable number of parameters appropriate for an analysis by comparing to some 'benchmark' (simple model)

Note that parameter number may increase if adding more constraints (whether correct or incorrect).

TMD, QCD definition (CSS2 scheme) To many moving parts ...

W (TMD region)

$$\sum_{q} \mathcal{H}_{q} \text{ F.T.} \left\{ \tilde{D}_{h/q}(z, z \boldsymbol{b}_{\perp}; Q) \ \tilde{f}_{q/P}(x, \boldsymbol{b}_{\perp}; Q) \right\}$$

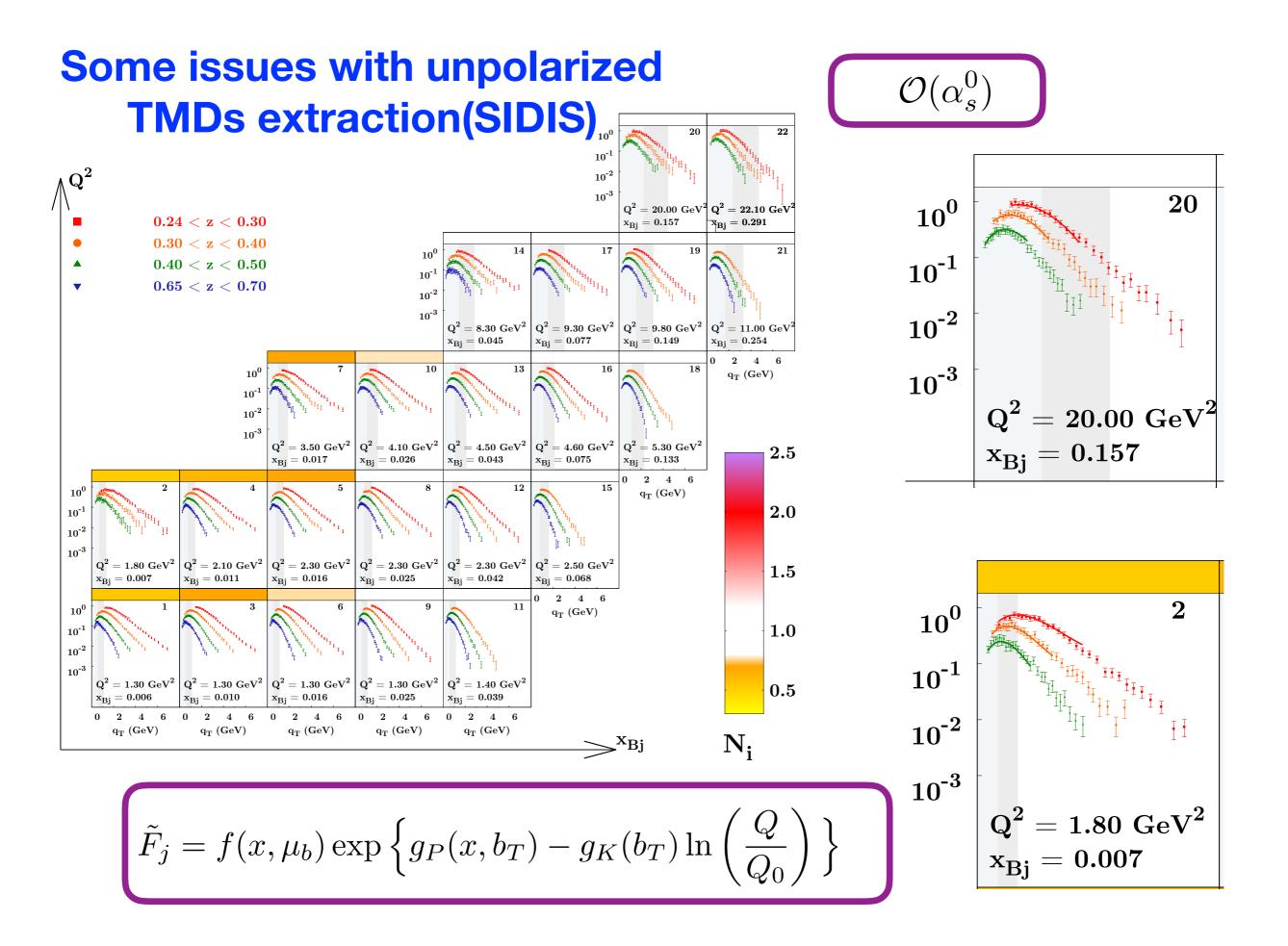
Fourier Transform of:

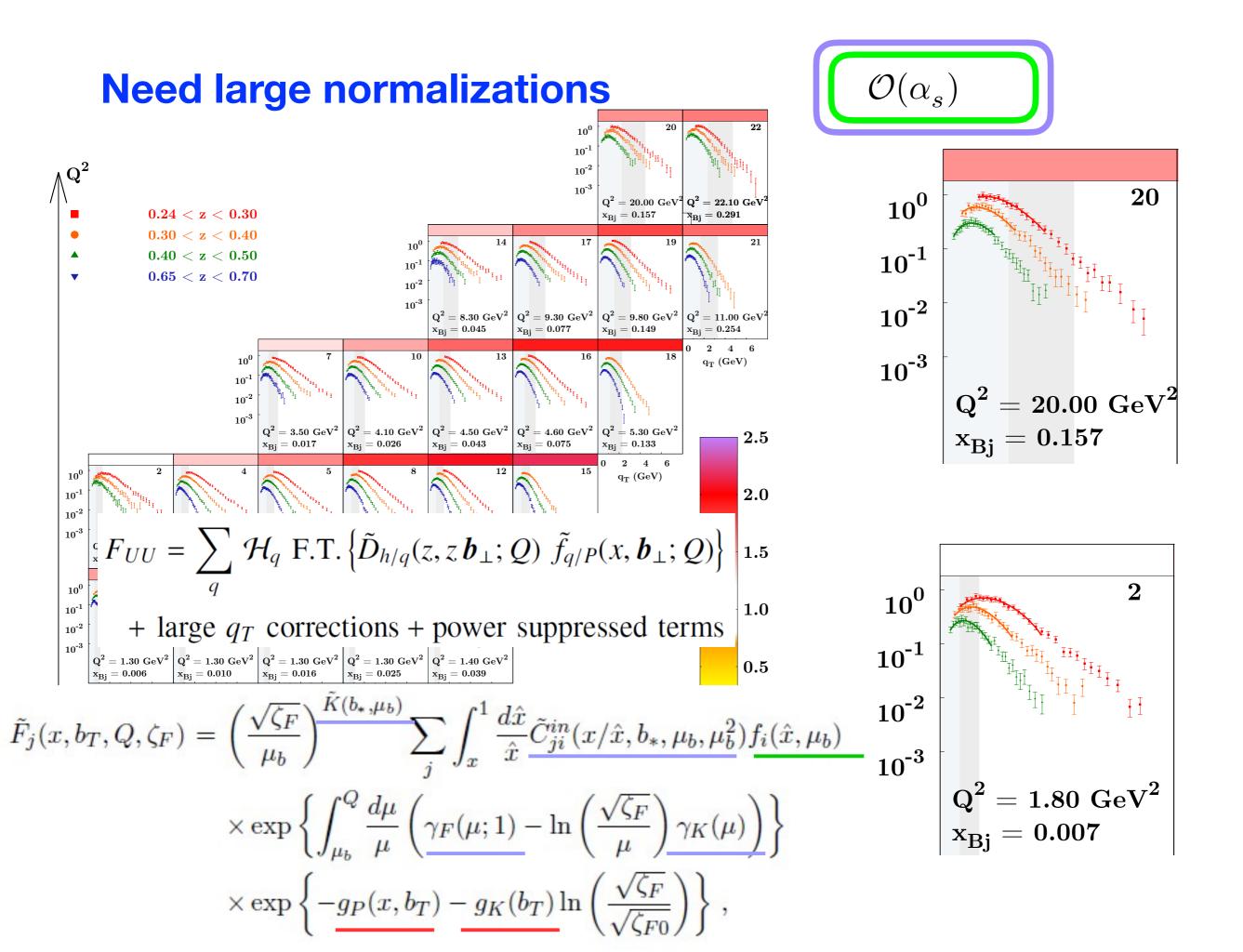
$$\begin{split} \tilde{F}_{j}(x,b_{T},Q,\zeta_{F}) &= \left(\frac{\sqrt{\zeta_{F}}}{\mu_{b}}\right)^{\tilde{K}(b_{*},\mu_{b})} \sum_{j} \int_{x}^{1} \frac{d\hat{x}}{\hat{x}} \underbrace{\tilde{C}_{ji}^{in}(x/\hat{x},b_{*},\mu_{b},\mu_{b}^{2}) f_{i}(\hat{x},\mu_{b})} \\ &\times \exp\left\{\int_{\mu_{b}}^{Q} \frac{d\mu}{\mu} \left(\underline{\gamma_{F}(\mu;1)} - \ln\left(\frac{\sqrt{\zeta_{F}}}{\mu}\right) \underline{\gamma_{K}(\mu)}\right)\right\} \\ &\times \exp\left\{-g_{P}(x,b_{T}) - g_{K}(b_{T})\ln\left(\frac{\sqrt{\zeta_{F}}}{\sqrt{\zeta_{F0}}}\right)\right\}, \end{split}$$

pQCD

Input (extraction from collinear cross section)

- Non-perturbative functions to extract from data.





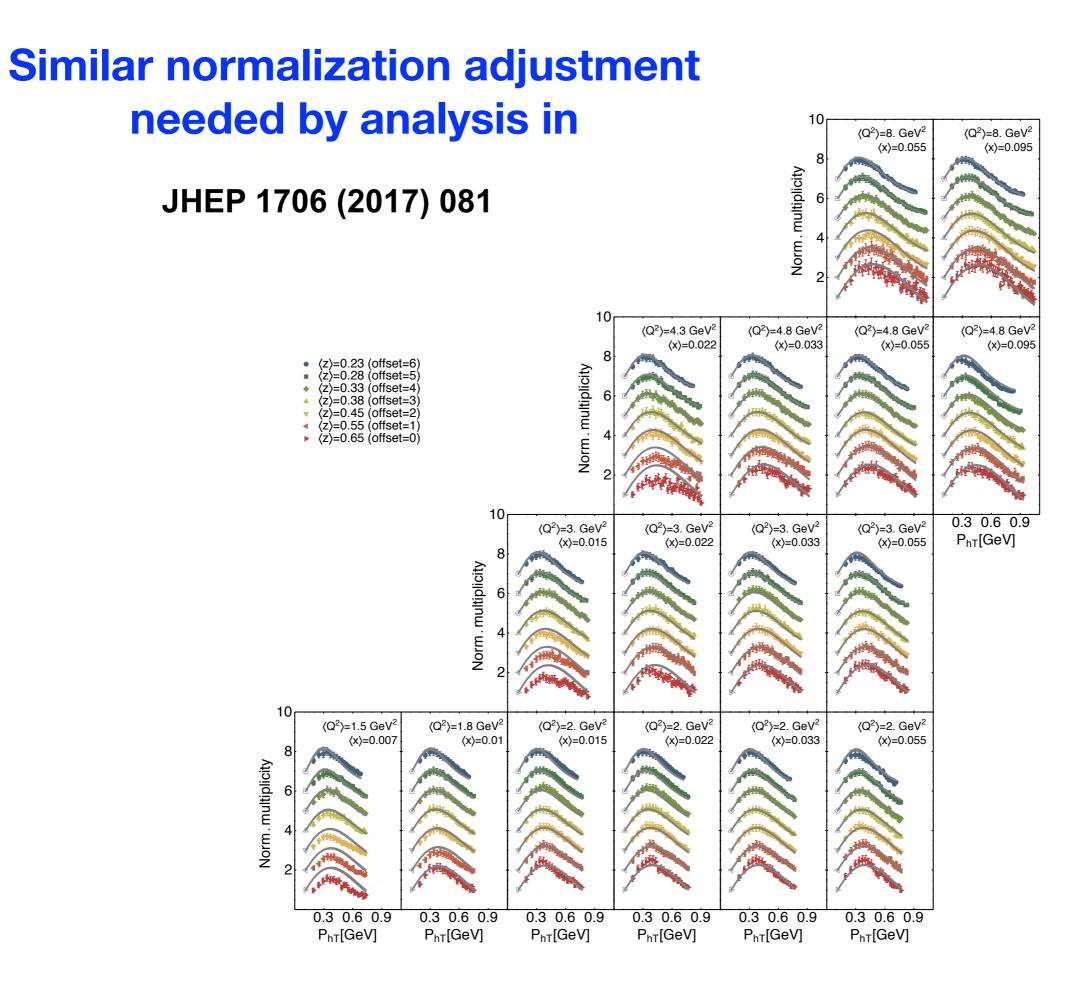
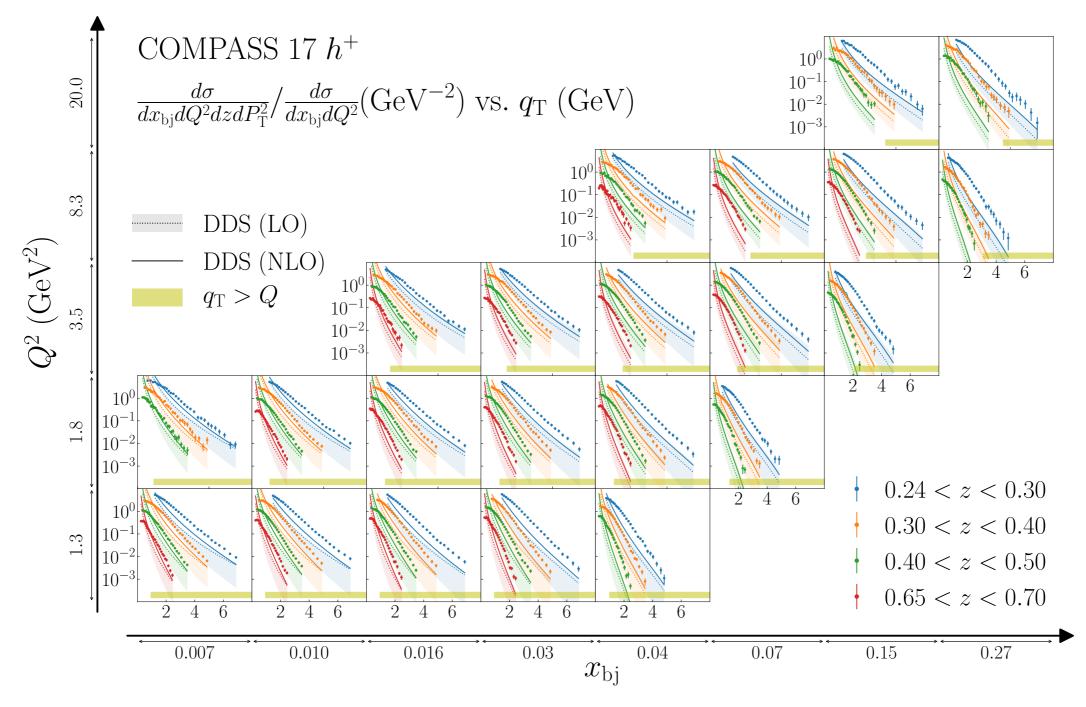


FIG. 4: Cross section as a function of p_T , data and cuts as in Figure 3.

Also issues in non-TMD region (large nowever the difference between LO and NLO decreases p_T increases.

The uncertainty due to the choice of a fragmentation functions set is also quite noticeable, this fact driven by the different graon content of the two sets considered here. Low Q^2 time seem to prefer KKP set, which have a larger gluon-fragmentation content, whereas for larger Q^2 both sets agree with the data within errors. LO estimates show a much smaller sensitivity on the choice of fragmentation functions, since gluon fragmentation does not contribute significantly to the cross section at this order.

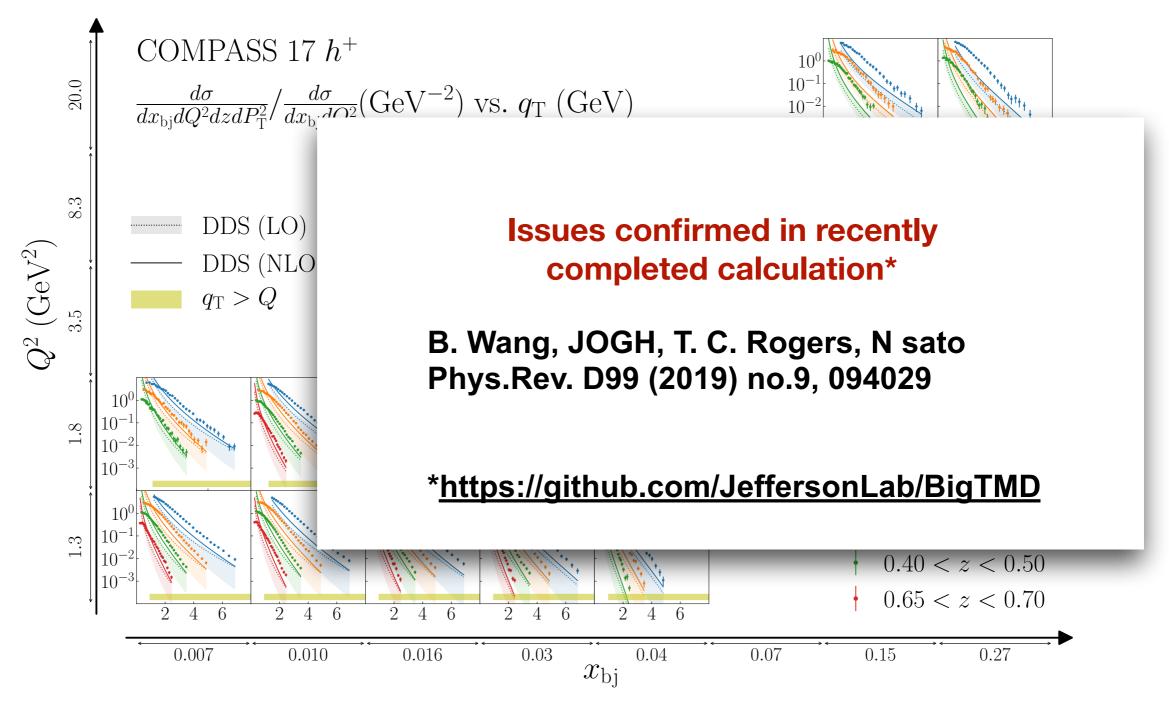


JOGH, Rogers, Sato, Wang Phys.Rev. D98 (2018) no.11, 114005

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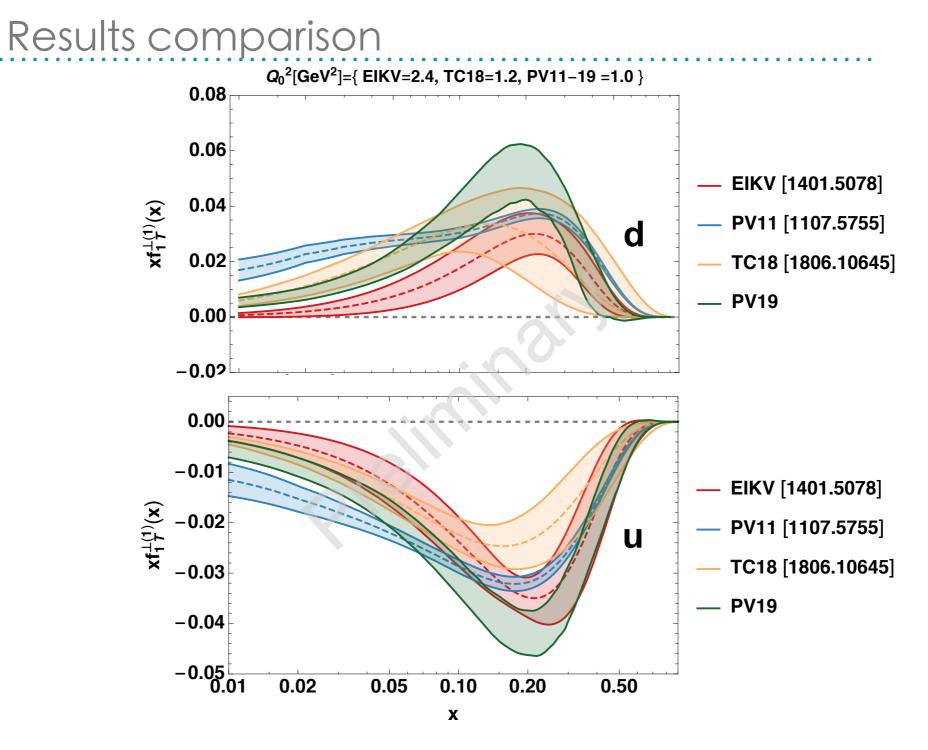


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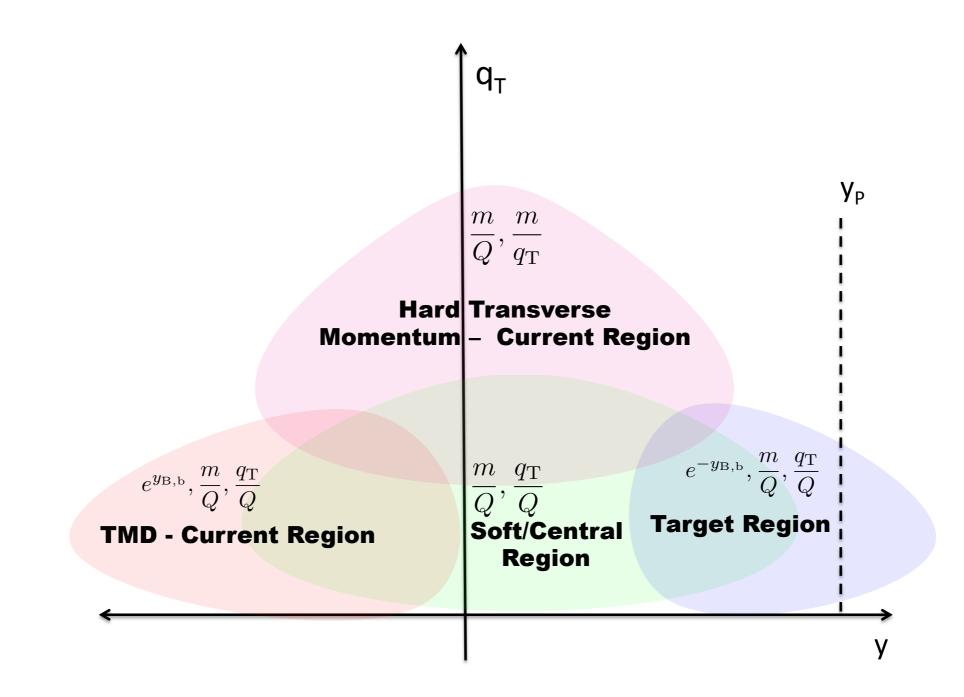
Some challenges: **One can still infer information about** the evolution W (TMD region) $\sum_{q} \mathcal{H}_{q} \text{ F.T.} \left\{ \tilde{D}_{h/q}(z, z \boldsymbol{b}_{\perp}; Q) \ \tilde{f}_{q/P}(x, \boldsymbol{b}_{\perp}; Q) \right\}$ Fourier Transform of: $$\begin{split} \tilde{F}_{j}(x,b_{T},Q,\zeta_{F}) &= \left(\frac{\sqrt{\zeta_{F}}}{\mu_{b}}\right)^{\tilde{K}(b_{*},\mu_{b})} \sum_{j} \int_{x}^{1} \frac{d\hat{x}}{\hat{x}} \underbrace{\tilde{C}_{ji}^{in}(x/\hat{x},b_{*},\mu_{b},\mu_{b}^{2})f_{i}(\hat{x},\mu_{b})}_{\times \exp\left\{\int_{\mu_{b}}^{Q} \frac{d\mu}{\mu} \left(\gamma_{F}(\mu;1) - \ln\left(\frac{\sqrt{\zeta_{F}}}{\mu}\right)\gamma_{K}(\mu)\right)\right\}} \\ &\times \exp\left\{-g_{P}(x,b_{T}) - g_{K}(b_{T})\ln\left(\frac{\sqrt{\zeta_{F}}}{\sqrt{\zeta_{F0}}}\right)\right\}, \end{split}$$ pQCD Input (extraction from collinear cross section)

Non-perturbative functions to extract from data.

Issues on normalization likely propagate here



Slide by Filippo Delcarro (JLAB) -QCD evolution 2019-

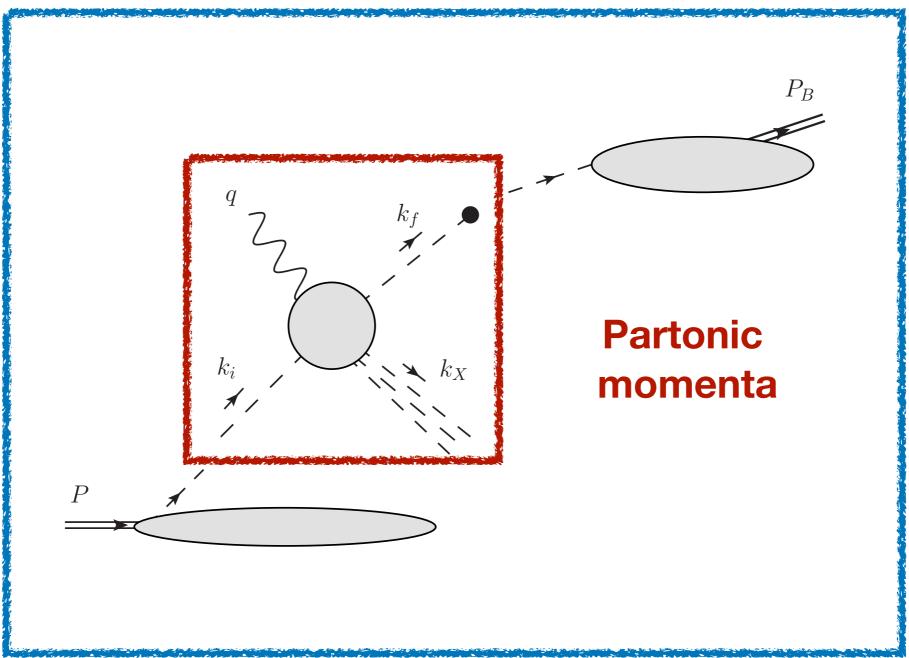


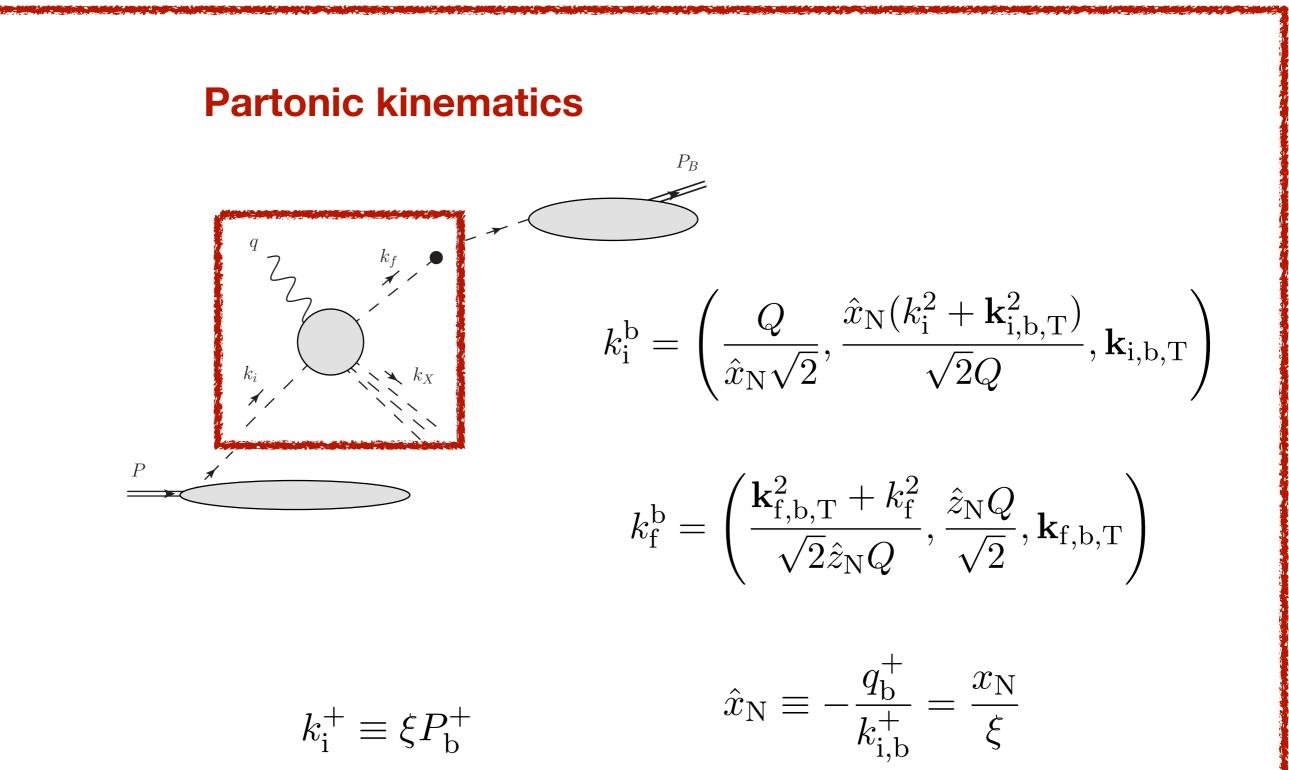
Based on:

M.Boglione, A. Dotson, L. Gamberg, S. Gordon, JOGH, A. Prokudin, T. C. Rogers, N. Sato Submitted to J.High Energy Phys. JLAB-THY-19-2920 e-Print: <u>arXiv:1904.12882</u>

Different types of approximations

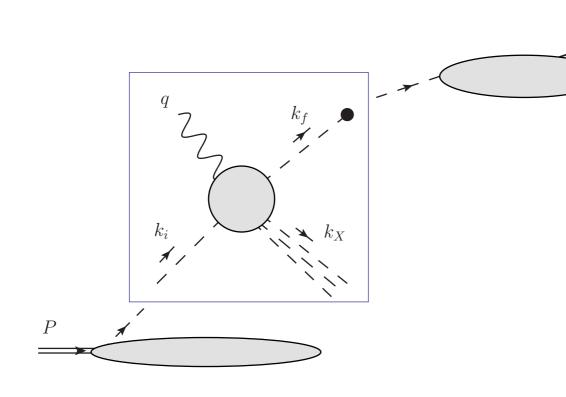
External momenta kinematics

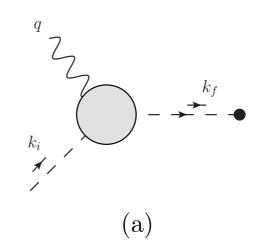


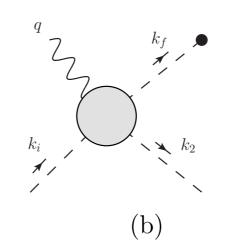


 $P_{\rm B,b}^{-} \equiv \zeta k_{\rm f}^{-}$

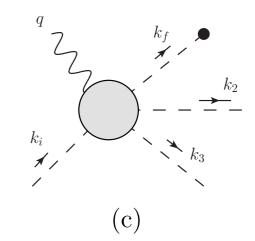
 $\hat{z}_{\mathrm{N}} \equiv \frac{k_{\mathrm{f,b}}^{-}}{q_{\mathrm{b}}^{-}} = \frac{z_{\mathrm{N}}}{\zeta}$

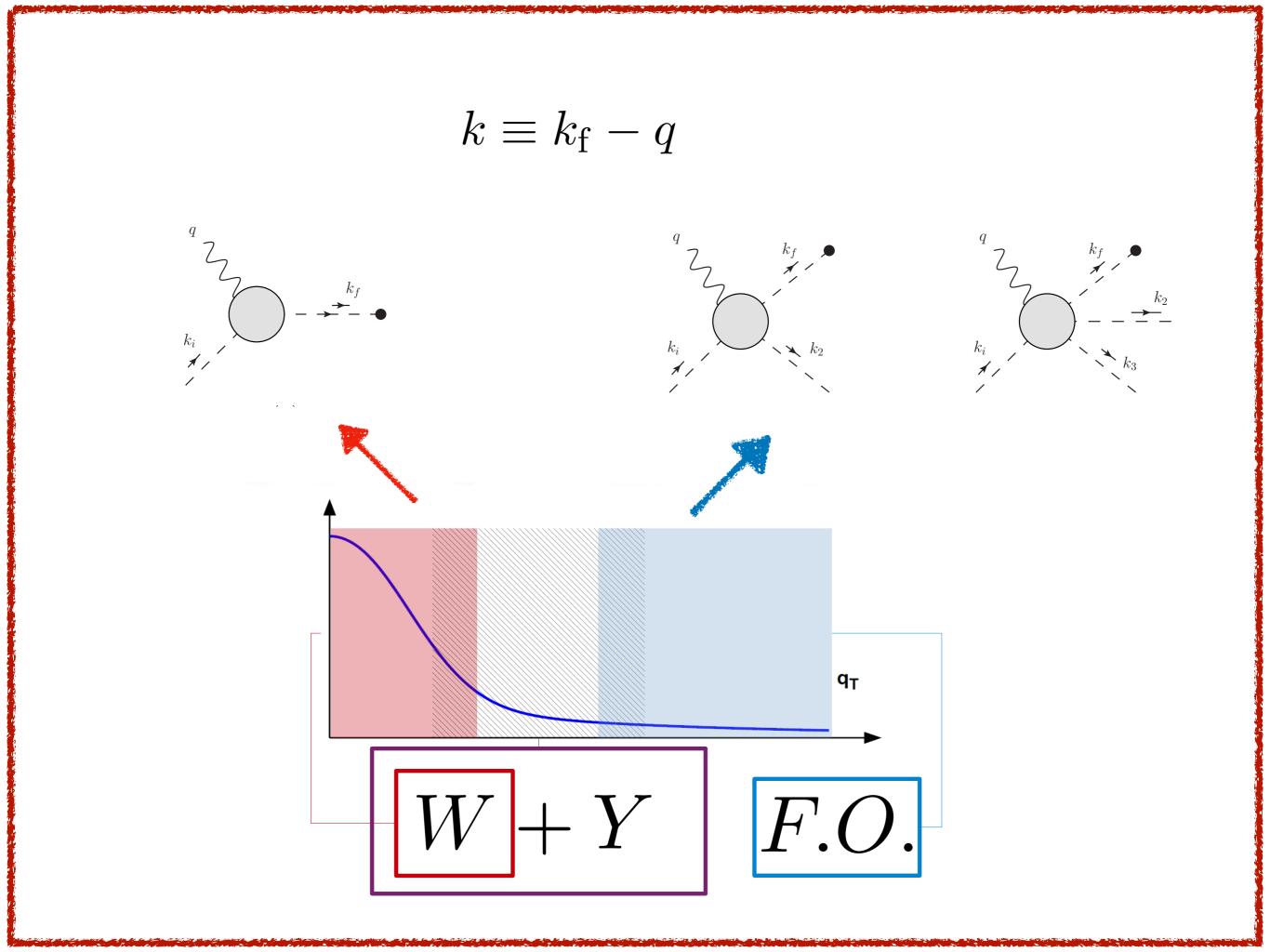






 P_B





$$k^{2} = (k_{f} - q)^{2}$$

$$k^{2} = 0$$

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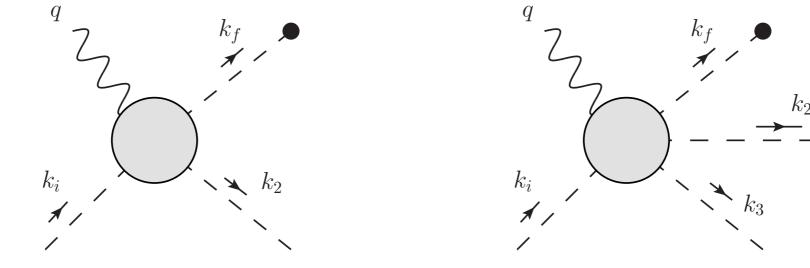
$$k^{2} = 0$$

$$k^{2} = 0$$

$$k^{2} \neq 0$$

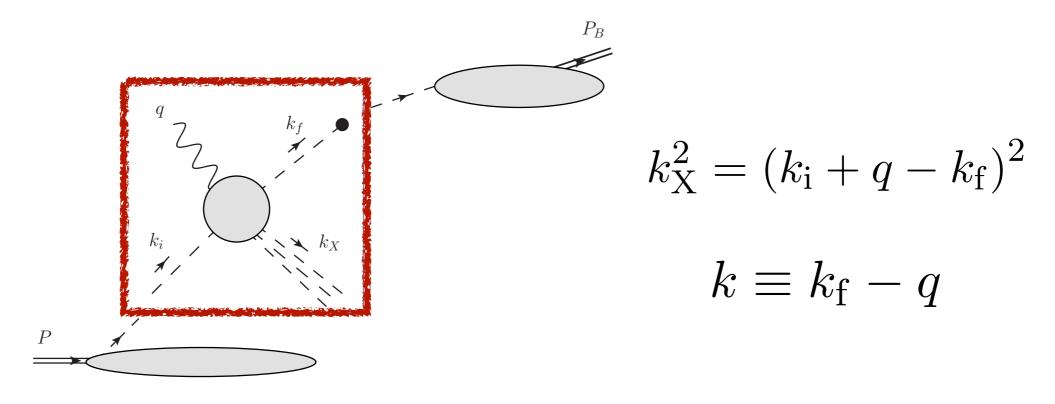
Allows to distinguish handbag from real emission kinematics

 $k_{\rm X}^2 = (k_{\rm i} + q - k_{\rm f})^2$



 $k_X^2 = 0$ $k_X^2 \neq 0$ Higher order pQCD corrections in large qT cross section associated to larger values of

cross section associated to larger values of virtuality spectator



Collinearity =
$$R_1 \equiv \frac{P_{\rm B} \cdot k_{\rm f}}{P_{\rm B} \cdot k_{\rm i}}$$

Transverse Hardness Ratio =
$$R_2 \equiv \frac{|k^2|}{Q^2}$$

Spectator Virtuality Ratio =
$$R_3 \equiv \frac{|k_X^2|}{Q^2}$$

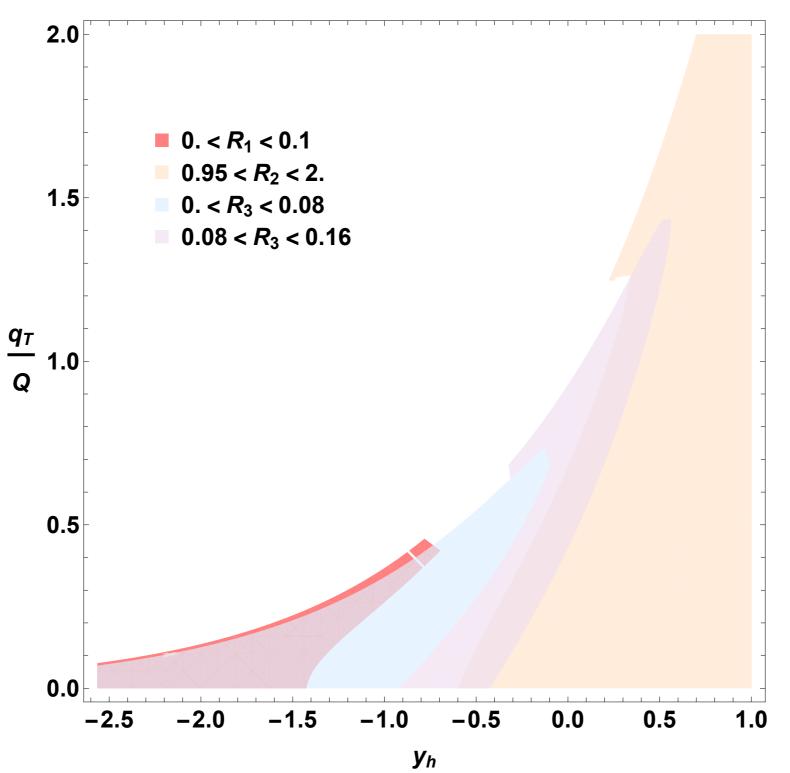
The size of these ratios determine partonic configurations (factorization theorem) and map to kinematical regions of the observables

Collinearity =
$$R_1 \equiv \frac{P_{\rm B} \cdot k_{\rm f}}{P_{\rm B} \cdot k_{\rm i}}$$

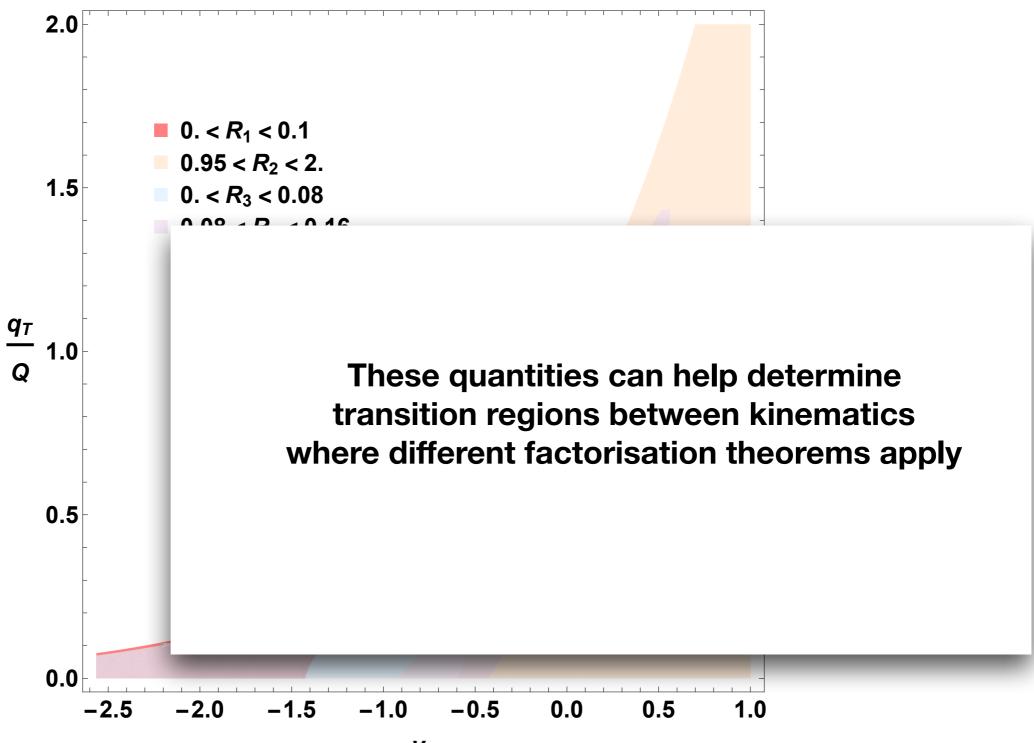
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(caveat: Parton momenta have to be estimated, so this is really just an example)



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Final remarks

Simple model like gaussian in momentum space are useful "snapshots" of TMDs at a given scale. They may also help assessing signals of evolution when used as a benchmark.

This is important since several issues remain in the extraction of TMDs.

Fits in CSS formalism (TMD evolution) undershoot the data.

One can probably study "shape" and evolution of TMDs separately, although ultimately, the normalisation issues should be resolved.

The issues propagate to other extractions.

Among other things, assessing the kinematics validity of different factorisation theorems seems important.

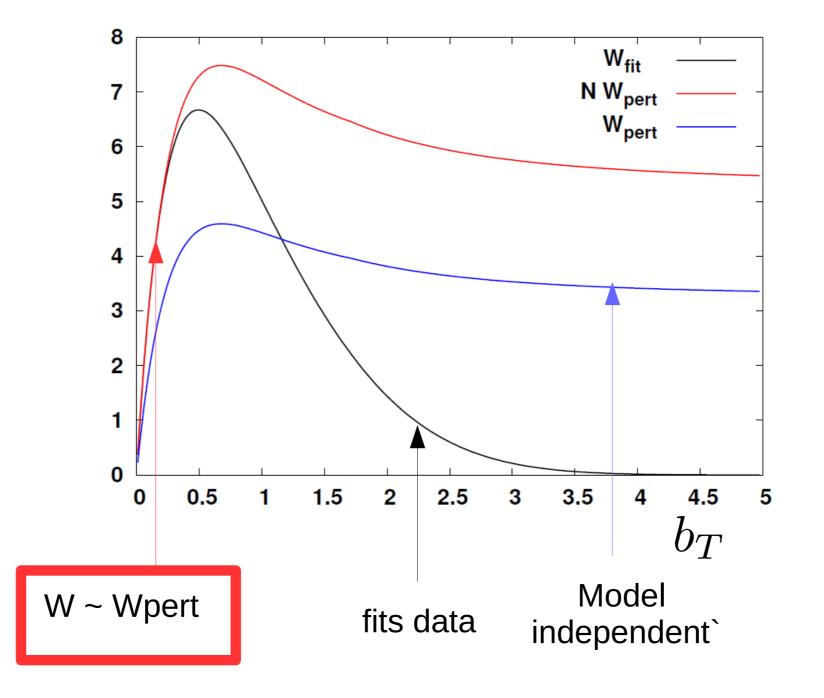
Quantities R1, R2, R3 presented here can serve as tool to do this.

Thanks.

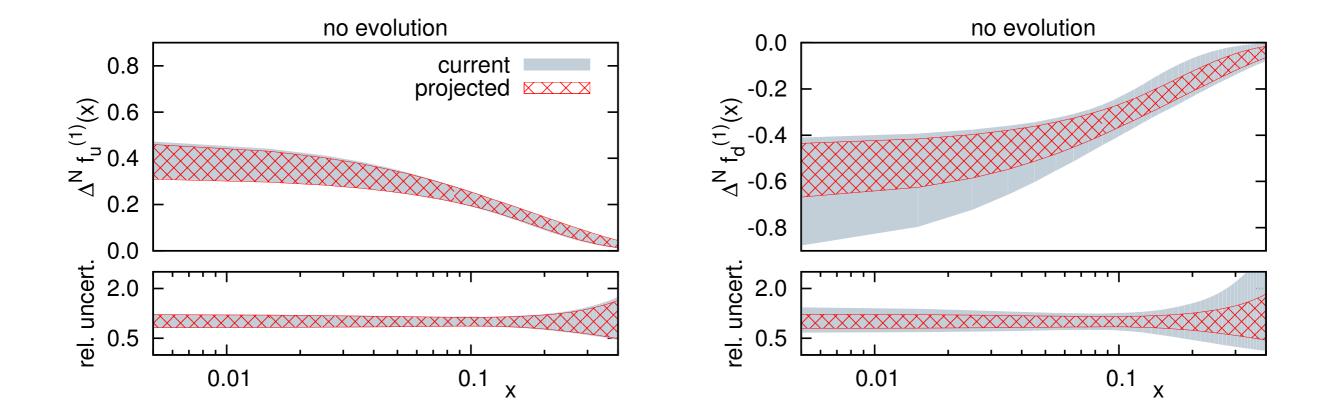
backup.

Some challenges:

These normalizations are hard to justify, but they do have and impact in the "shape" of the TMDs.



Sivers asymmetry in SIDIS



Uncertainty bands corresponding to projected errors for future COMPASS run on Deuteron target