

# Top Quark Interactions in Simplified Dark Matter Models

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- Top Quark Physics
  - Top Quark Production Mechanisms
- Simplified Dark Matter Models
  - Spin-0 Mediators
  - Spin-1 Mediators
- Conclusions and Outlook

# Top Quark Physics

# Top Quark in the SM

- First observed in 1995: **Top quark pair production**

F. Abe et al., Phys. Rev. Lett. 74 (1995) 2626;

S. Abachi et al., Phys. Rev. Lett., 74 (1995) 2632.

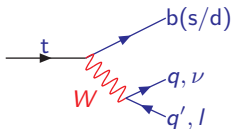
- Observed again in 2009: **Single top quark production**

S. Abachi et al., Phys. Rev. Lett., 103 (2009) 092001;

T. A. Aaltonen et al., Phys. Rev. Lett. 103 (2009) 092002.

- It is the heaviest elementary particle in the SM ( $m_t \approx 173$  GeV).
- Production ( $1/m_{top}$ ) < Lifetime ( $\tau_t \simeq 1/\Gamma_t \approx 5 \times 10^{-25}$  s) < Hadronization ( $\tau_{had} \simeq 1/\Lambda_{QCD} \approx 3 \times 10^{-24}$  s) < Spin decorrelation ( $m_{top}/\Lambda^2$ ).
- As a consequence, it is possible to measure t quark polarisation, spin correlations and  $W^\pm$  boson helicity states by studying angular distributions of the decay products.

- $t \rightarrow Wb$  in the SM.



W decay	BR
$W \rightarrow l\nu$	0.32
$W \rightarrow qq'$	0.68

- It plays a very important role in the determination of the EWSB mechanism ( $\lambda_t \sim 1$ ) and also in NP connected to the EWSB.

Some Top Quark Reviews:

W. Bernreuther, J. Phys. G35 (2008) 083001;

V. del Duca and E. Laenen, Int. J. Mod. Phys. A30 (2015) no. 35, 1530063;

U. Husemann, Prog. Part. Nucl. Phys. 95 (2017) 48-97;

M. Cristinziani and M. Mulders, J. Phys. G44 (2017) no. 6, 063001.

# Top Quark Production Mechanisms

# Top Quark Pair Production (at LO QCD)

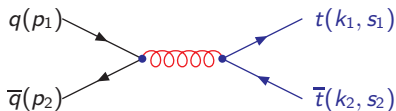
At **leading order (LO)** the **partonic cross section** for  $t\bar{t}$  production is of order  $\mathcal{O}(\alpha_s^2)$ . The subprocesses that contribute to the cross section at this level are

M. Glück, J. F. Owens and E. Reya, Phys. Rev. D17 (1978) 2324;

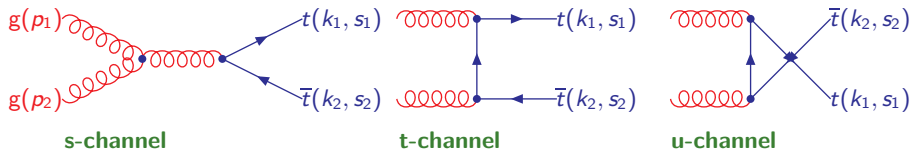
J. Babcock, D. Silvers and S. Wolfram, Phys. Rev. D18 (1978) 162;

H. Georgi et al., Ann. Phys. 114 (1978) 273.

## $q\bar{q}$ Annihilation



## Gluon Fusion



The differential cross section for the two particle scattering process can be written as

$$\frac{d\sigma}{dz} = \frac{\beta_t}{32\pi s} \sum |\mathcal{M}(\hat{s}, m_t, z)|^2$$

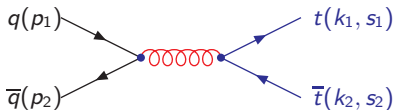
with the spin and color averaged square matrix element, and where  $z = \frac{\hat{u}-\hat{t}}{\hat{s}} = \cos\theta$ , with  $\theta$  being the scattering angle, and  $\beta_t$  is the top quark velocity defined by

$$\beta_t = \sqrt{1 - \frac{4m_t^2}{\hat{s}}}$$

W. Benakker, A. Denner, W. Hollik, T. Mertig, R. Sack and D. Wackerth, Nucl. Phys. B (1994) 343



### $q\bar{q}$ Annihilation



$$\overline{\sum} |\mathcal{M}_B^{q\bar{q}}|^2 = \frac{8}{9} \alpha_s^2 \pi^2 (2 - \beta_t^2 (1 - z^2))$$

Integrating the differential cross section between the kinematical limits  $-1 < z < 1$  we get the total cross section for the  $q\bar{q}$  annihilation subprocess

$$\hat{\sigma}_B^{q\bar{q}}(\hat{s}) = \frac{4}{27} \frac{\alpha_s^2 \pi \beta_t}{s} (3 - \beta_t^2)$$

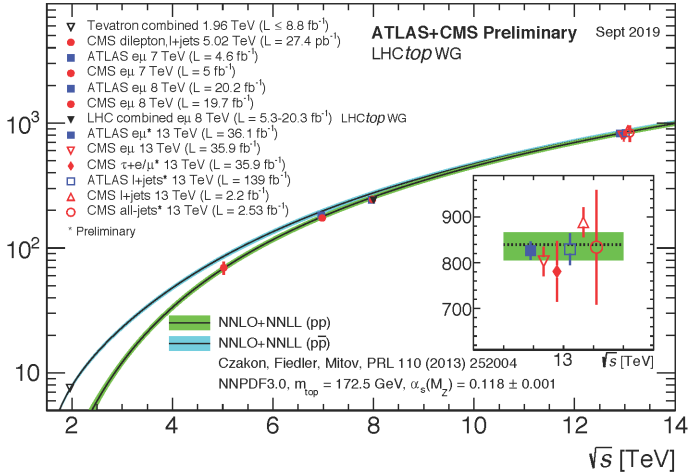


# Top Quark Production at the Tevatron and LHC

	$q\bar{q} \rightarrow t\bar{t}$	$gg \rightarrow t\bar{t}$
Tevatron ( $p\bar{p}$ at $\sqrt{s} = 1.96 \text{ TeV}$ )	85%	15%
LHC ( $pp$ at $\sqrt{s} = 14 \text{ TeV}$ ( $\sqrt{s} = 7 \text{ TeV}$ ))	10%	90% ( $\approx 80\%$ )

M. Tanabashi et al. (Particle Data Group), Phys. Rev. D 98 (2018) 030001

Inclusive  $t\bar{t}$  cross section [pb]



<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/LHCTopWGSummaryPlots>

# Single Top Quark Production

- The most important production process at hadron colliders is  $t\bar{t}$ , which is mediated by the **strong interaction**.
- Single top quarks (antiquarks) production is mediated by **electroweak interactions**.
- The single top quark signal is smaller than the  $t\bar{t}$  signal and it is difficult to separate from the background.
- The single top quark production cross section is within an order of magnitude of top quark pair production.

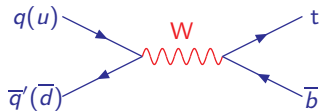
- $t \rightarrow Wb$  vertex in production and decay.
- Top is produced polarized, almost 100%.
- Cross sections are proportional to  $|V_{tb}|^2$  in all channels.
- BSM physics can appear in cross sections and properties.

E. Boos and L. Dudko, *Int. J. Mod. Phys. A*27 (2012) 1230026;

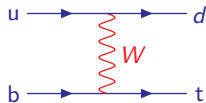
A. Giammanco and R. Schwienhorst, *Rev. Mod. Phys.* 90 (2018) no.3 035001;

# Single Top Quark Production Channels

Electroweak single top quark production

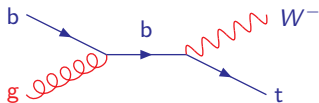


s-channel

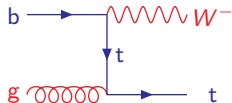


t-channel

Top - W boson

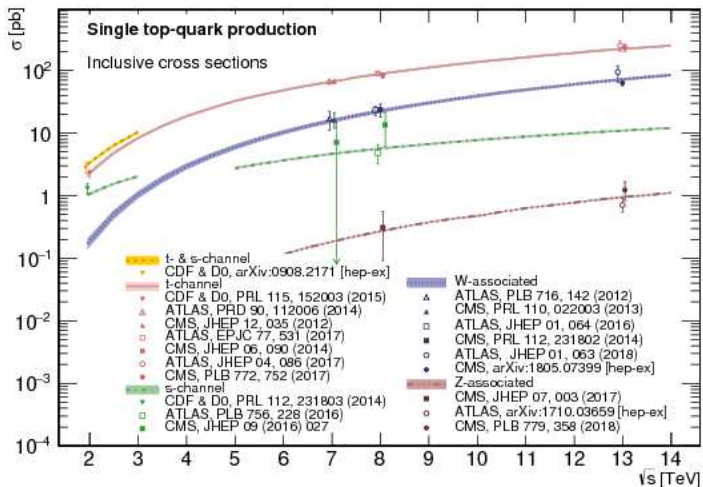


s-channel



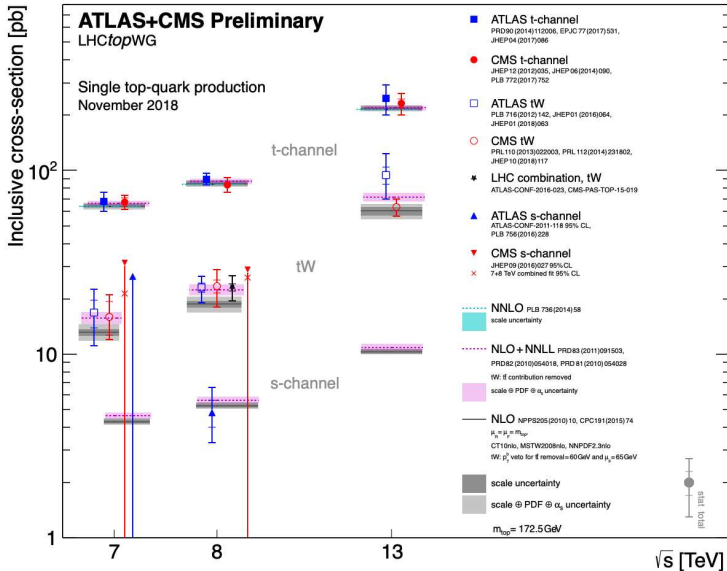
t-channel

# Single Top Quark Production: Inclusive Cross Sections



A. Giammanco and R. Schwienhorst, *Rev. Mod. Phys.* 90 (2018) 035001

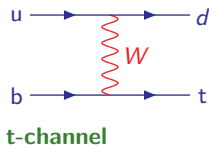
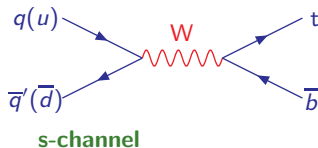




<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/LHCTopWGSummaryPlots>

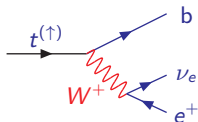
# s-channel and t-channel: Spin Correlations

**Spin correlations** may appear when the top quark is highly polarized in its production and decay

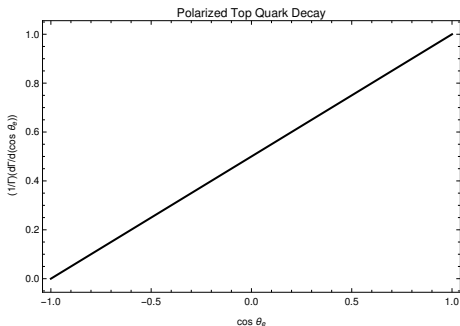


# Polarized Top Quark Decay

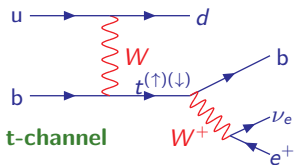
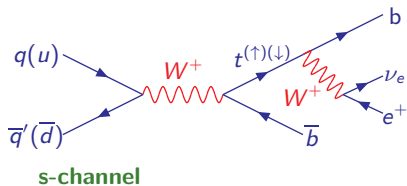
$t^{(\uparrow)} \rightarrow b\nu_e e^+$  in the SM.



$$\frac{1}{\Gamma_T} \frac{d\Gamma_{t^{(\uparrow)}}}{d(\cos \theta_{e^+})} = \frac{1}{2}(1 + \cos \theta_{e^+})$$



# Single Top Quark Production and Decay



We define the spin asymmetry factor  $A_{\uparrow\downarrow}$  as

$$A_{\uparrow\downarrow} = \frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}}$$

which defines the size of the observable angular correlations when there is a mixture of spin up and spin down top quarks.

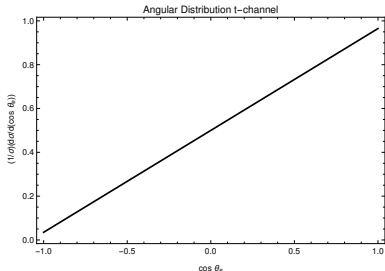
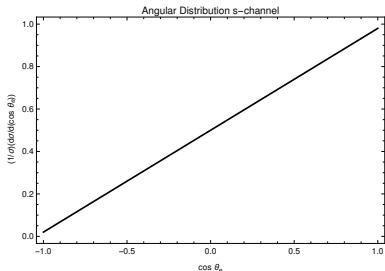
Angular distributions are linear in the cosine of the decay angles:

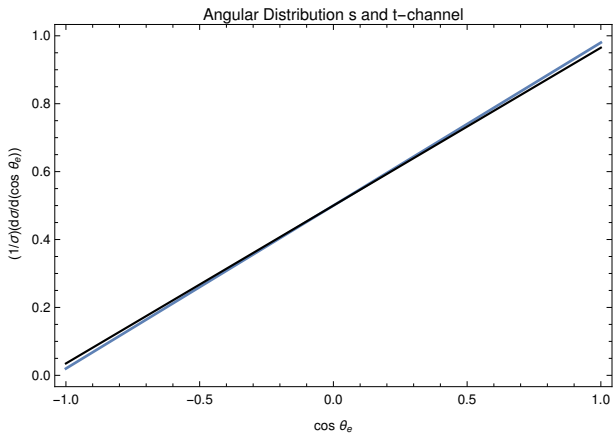
$$\frac{1}{\sigma_T} \frac{d\sigma_T}{d(\cos\theta_{e^+})} = \frac{1}{2}(1 + A_{\uparrow\downarrow} \cos\theta_{e^+})$$

G. Mahlon, arXiv:hep-ph/0011349v1

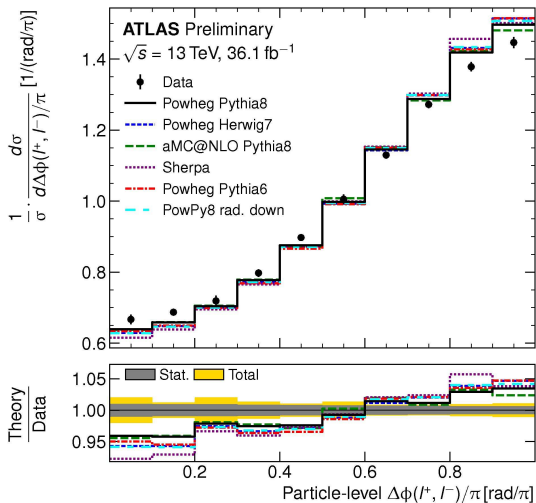
For the s-channel,  $\theta_{e^+}$  is the angle between the momenta of the outgoing positron and the incoming  $\bar{d}$  ( $A_{\uparrow\downarrow} = 0.96$ ).

For the t-channel,  $\theta_{e^+}$  is the angle between the momenta of the outgoing positron and the outgoing  $d$  quark ( $A_{\uparrow\downarrow} = 0.93$ ).





# Top Quarks as Probes of New Physics



ATLAS Collaboration. Jul 5, 2018. ATLAS-CONF-2018-027



# Simplified Dark Matter Models

# Dark Matter

- The nature of the **dark matter (DM)** is a fundamental open question in physics.
- Its discovery is an important goal of the **LHC**.
- **Potential interactions between DM and SM particles.**
- A viable **DM** candidate must be:
  - electrically neutral
  - stable
  - weakly interacting
  - decay lifetime larger than the age of the Universe

- **DM is invisible at the LHC**: Experimental signature of DM production at colliders is an event with a visible final state object recoiling against  $E_t^{miss}$  associated with DM.
- **Simplified Models** have few assumptions about DM and a minimal particle content.
  - D. Pinna et al., Phys. Rev. D96 (2017) 035031;
  - P. Pani and G. Polesello, Phys. Dark Univ. 21 (2018) 8;
  - CMS Collaboration. 2018. CMS-PAS-EXO-18-010.
- **Associated production of DM with top quarks may affect the spin correlations.**

# Simplified Dark Matter Models: Spin-0 Mediators

The **Lagrangian** with the interactions between **SM particles** and **DM** ( $\chi$ , Dirac fermions) mediated by a **massive electrically neutral scalar or pseudo-scalar**  $\varphi$ , is given by

$$\mathcal{L}_\varphi \supset g_\chi \varphi \bar{\chi} \chi + \frac{g_v \varphi}{\sqrt{2}} \sum_f (y_f \bar{f} f)$$
$$\mathcal{L}_A \supset i g_\chi A \bar{\chi} \gamma^5 \chi + \frac{i g_v A}{\sqrt{2}} \sum_f (y_f \bar{f} \gamma^5 f)$$

where  $y_f = \sqrt{2} m_f / v$  are the Yukawa couplings, with  $v = 246$  GeV,  $g_\chi$  is the **DM mediator coupling** and  $g_v$  is the **fermion mediator coupling**. **Minimal set of four free parameters** (with MFV):  $m_\chi$ ,  $m_\varphi$ ,  $g_\chi$  and  $g_v$ .

D. Pinna et al., Phys. Rev. D96 (2017) 035031;

D. Abercrombie et al., arXiv:1507.00966;

M. R. Buckley, D. Feld and D. Goncalves, Phys. Rev. D91 (2015) 015017.

The decay width of the scalars is

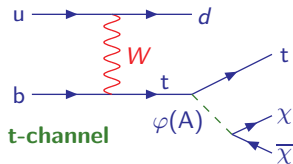
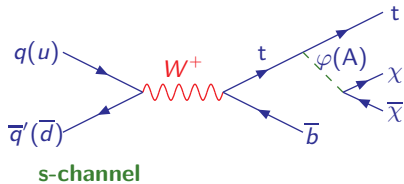
$$\Gamma_\varphi = \frac{g_\chi^2 m_\varphi}{8\pi} \left(1 - \frac{4m_\chi^2}{m_\varphi^2}\right)^{n/2} + \sum_f \frac{g_v^2 y_f^2 m_\varphi}{16\pi} \left(1 - \frac{4m_f^2}{m_\varphi^2}\right)^{n/2}$$

D. Pinna et al., Phys. Rev. D96 (2017) 035031;

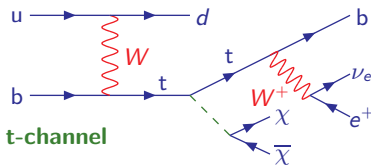
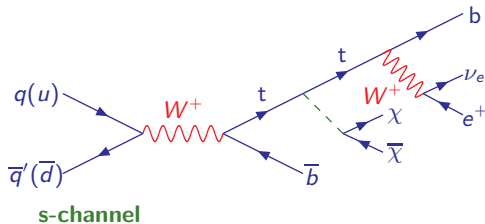
D. Abercrombie et al., arXiv:1507.00966;

M. R. Buckley, D. Feld and D. Goncalves, Phys. Rev. D91 (2015) 015017.

# Dark Matter and Single Top Quark Production

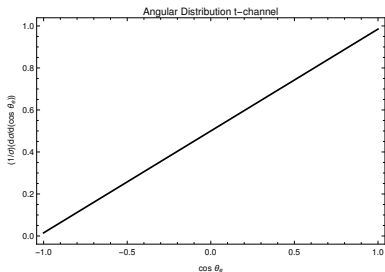
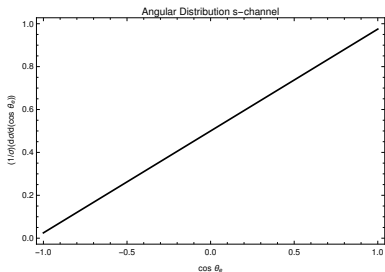


# Dark Matter and Single Top Quark Production and Decay



For the s-channel,  $\theta_{e^+}$  is the angle between the momenta of the outgoing positron and the incoming  $\bar{d}$ .

For the t-channel,  $\theta_{e^+}$  is the angle between the momenta of the outgoing positron and the outgoing  $d$  quark.





# Simplified Dark Matter Models: Spin-1 Mediators

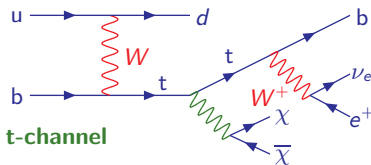
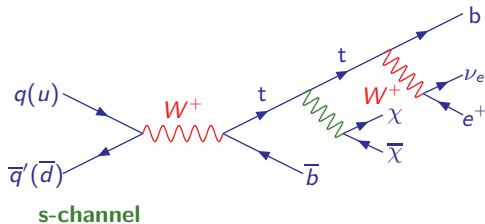
The **Lagrangian** with the interactions between **SM particles** and **DM** ( $\chi$ , Dirac fermions) mediated by a **spin-1 particle**, is given by

$$\mathcal{L}_V \supset V_\mu \bar{\chi} \gamma^\mu (g_\chi^V - g_\chi^A) \chi + \sum_{f=q,l,\nu} V_\mu \bar{f} \gamma^\mu (g_f^V - g_f^A \gamma_5) f,$$

Minimal set of six free parameters (with MFV):  $m_\chi$ ,  $M_V$ ,  $g_\chi^V$ ,  $g_u^V$ ,  $g_d^V$  and  $g_l^V$ .

J. Abdallah et al., Phys. Dark. Univ. 9-10 (2015) 8-23.

# Dark Matter and Single Top Quark Production and Decay



# Conclusions and Outlook

- A detailed MC study is needed to study in detail effects that could be observed.
- A detailed study of the **spin correlations in DM production in association with a single top quark** at the LHC is in progress.