# A collection of fun numerical methods that [might] lead us to Lattice QCD 

yes i think they're fun

Based on: Lattice QCD: a practical guide; SUPA lecture series by Christine Davies,
and G. P. Lepage, Lattice QCD for Novices, hep-lat/0506036
Vithyaban Anjelo Narendran @ Graduate Symposium 2024

## Quick Intro

## Immediately starts assuming things

- Assume you know of QCD


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- Assume you know of QCD
- Problems when doing perturbation theory
- Power series expansion in $\alpha_{s}$, the QCD coupling constant.
- Blows up if $\alpha_{s}$ is large.
- Calculate numerically the properties we want to know about in QCD - masses and properties of the bound states.


## Muon g-2

## Houston, we have a result!



## Quantum Mechanics

$$
\hbar=c=1
$$

## 1 particle, 1 dimension



$$
[\hat{x}, \hat{p}]=i
$$

Position and Momentum Operators don't commute

## Quantum Mechanics

## 1 particle, 1 dimension

$$
\hat{H}=\frac{p^{2}}{2 m}+V(x) \quad[\hat{x}, \hat{p}]=i
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- Solve Schrödinger's equation.


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- Find Eigenfunctions, Eigenvalues


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- Solve Schrödinger's equation.
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- 

Party!

## Path Integral Formulation

Transition Amplitude and Action

$$
K\left(x, t ; x_{i}, t_{i}\right)=\left\langle x(t) \mid x_{i}\left(t_{i}\right)\right\rangle=\int \mathscr{D} x(t) e^{i S[x]}
$$

## Path Integral Formulation

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Weighted by exponential of classical action

## Path Integral Formulation

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$$

Weighted by exponential

$$
S[x] \equiv \int_{t_{i}}^{t} d t L(x, \dot{x}) \equiv \int d t\left[\frac{m \dot{x}(t)^{2}}{2}-V(x(t))\right]
$$

Lagrangian

## Classical Paths and Quantum Fluctuations



Classical Paths and Quantum Fluctuations


Classical Paths and Quantum Fluctuations


## Connecting to the Eigenstates of the Hamiltonian

 Inserting Hamiltonian via time evolution$$
\langle x| \exp (-i H(T))\left|x_{i}\right\rangle=\sum_{n}\langle x \mid n\rangle \exp \left(-i E_{n} T\right)\langle n|\left|x_{i}\right\rangle
$$

## Connecting to the Eigenstates of the Hamiltonian

 Inserting Hamiltonian via time evolution$$
\begin{aligned}
& \sum_{n}\langle n||n\rangle=1 \\
& \langle x| \exp (-i H(T))\left|x_{i}\right\rangle=\sum_{n}\langle x \mid n\rangle \exp \left(-i E_{n} T\right)\langle n|\left|x_{i}\right\rangle
\end{aligned}
$$

## Connecting to the Eigenstates of the Hamiltonian

 Inserting Hamiltonian via time evolution$$
\begin{array}{cc}
\sum_{n}\langle n||n\rangle=1 & \text { Overlaps of Eigenstate with Position Eigenstate } \\
\langle x| \exp (-i H(T))\left|x_{i}\right\rangle=\sum_{n}\langle x \mid n\rangle \exp \left(-i E_{n} T\right)\langle n|\left|x_{i}\right\rangle
\end{array}
$$

## Connecting to the Eigenstates of the Hamiltonian

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\begin{gathered}
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\langle x| \operatorname{lexp}(-i H(T))\left|x_{i}\right\rangle=\sum_{n}\langle x \mid n\rangle \exp \left(-i E_{n} T\right)\langle n|\left|x_{i}\right\rangle \\
=\sum_{n} \psi_{n}^{*}\left(x_{i}\right) \psi_{n}(x) \exp \left(-i E_{n} T\right)
\end{gathered}
$$

## Connecting to the Eigenstates of the Hamiltonian

 Inserting Hamiltonian via time evolution$$
\begin{gathered}
\sum_{n}\langle n||n\rangle=1 \quad\left|\begin{array}{c}
\mid \text { Insert } \\
\left\langle x^{2}\right| \exp (-i H(T))
\end{array} x_{i}\right\rangle=\sum_{n}\langle x \mid n\rangle \exp \left(-i E_{n} T\right)\langle n|\left|x_{i}\right\rangle \\
=\sum_{n} \psi_{n}^{*}\left(x_{i}\right) \psi_{n}(x) \exp \left(-i E_{n} T\right)
\end{gathered}
$$

$\square$

## We Rotate Time <br> $$
x_{i}=x_{f}=x
$$

Minkowski -> Euclidean space-time ( $t \rightarrow-i t$ )

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New Transition Amplitude

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Minkowski -> Euclidean space-time $(t \rightarrow-i t)$

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\begin{gathered}
\text { New Transition Amplitude } \\
\langle x| e^{-H(T)}|x\rangle=\int \mathscr{D} x e^{-S[x]} \quad S[x] \equiv \int_{0}^{T} d t L(x, \dot{x}) \equiv \int d t\left[\frac{m \dot{x}(t)^{2}}{2}+V(x(t))\right]
\end{gathered}
$$

Weighting is $\Re$ valued

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& \text { Inserting Eigenstates as before: }
\end{aligned}
$$

$$
\langle x| e^{-H(T)}|x\rangle=\sum_{n} \psi_{n}^{*}(x) \psi_{n}(x) e^{-E_{n} T}
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$$

Inserting Eigenstates as before:
Decaying Exponential of all E Eigenstates

$$
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Inserting Eigenstates as before:

$$
\langle x| e^{-H(T)}|x\rangle=\sum_{n} \psi_{n}^{*}(x) \psi_{n}(x) e^{-E_{1} T}
$$

If we have large $T$, smallest $E$ gets picked out -> Extract $E_{0}$

## Transition Amplitudes -> Eigenstates of the Hamiltonian <br> -> Physics

## Excited States

$$
\frac{\langle x(T)| x\left(t_{2}\right) x\left(t_{1}\right)|x(0)\rangle}{\langle x(T) \mid x(0)\rangle}=\frac{\int \mathscr{D} x x\left(t_{2}\right) x\left(t_{1}\right) e^{-S[x]}}{\int \mathscr{D} x e^{-S[x]}}
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\text { Average } x\left(t_{1}\right) x\left(t_{2}\right) \\
\text { on each path }
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Plug a complete set of Eigenstates in as before:

$$
\left.=\left|\left\langle E_{0}\right| x\right| E_{1}\right\rangle\left.\right|^{2} e^{-\left(E_{1}-E_{0}\right)\left(t_{2}-t_{1}\right)}, t_{2}-t_{1} \rightarrow \infty
$$

## Excited States

## Average $x\left(t_{1}\right) x\left(t_{2}\right)$ on each path

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Average over all paths generated
Plug a complete set of Eigenstates in as before:

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& \text { Matrix Element }
\end{aligned}
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Matrix Element

[^0]
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## Discretise Time

Divide up the line from $t_{i} \rightarrow t_{f}$ to a set of points

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\int \mathscr{D} x(t)=\int d x_{1} d x_{2} \ldots d x_{n}
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Divide up the line from $t_{i} \rightarrow t_{f}$ to a set of points

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$$

Integrate over all the possible values of $x$ for each $x_{j}$
n-dimensional integral

## Calculating the Path Integral Numerically

## Discretising the Action

Call $t_{j+1}-t_{j}=a$, the lattice spacing

$$
\int_{t_{j}}^{t_{j+1}} d t L \approx a\left[\frac{m}{2}\left(\frac{x_{j+1}-x_{j}}{a}\right)^{2}-\frac{1}{2}\left(V\left(x_{j+1}\right)+V\left(x_{j}\right)\right)\right]
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Call $t_{j+1}-t_{j}=a$, the lattice spacing

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& \\
& \\
& \hline
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\dot{x}(t)^{2} \text { between } t_{j+1} \text { and } t_{j} \\
\text { Euler Method }
\end{gathered}
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$$

## Moving to QCD <br> The Path Integral in QCD



Gluon Fields


Weighted by the exponential of the QCD action

## Moving to QCD

## Build a Lattice



## Moving to QCD

## Impose Periodic Boundary Conditions



## 4D Doughnut? <br> Rendering error: Try installing additional dimensions via: sudo apt-get xtra-dims

## Finishing Up

- Path Integrals for Eigenstates, and useful physics
- Calculating excited states (particles)
- Discretisation
- Make-up of the lattice
- There are so very many cool things in this field
- Big thanks to Christine Davies - Physicist in Lattice QCD


## More on the Lattice

Local Gauge Invariance of QFTs

$$
\begin{array}{ll}
\mathscr{L}_{Q E D}(x)=-\frac{1}{4} F_{\mu \nu}(x) F^{\mu \nu}(x)+\bar{\psi}(x)(\gamma \cdot D+m) \psi(x) \\
F_{\mu \nu}(x)=\partial_{\mu} A_{\nu}(x)-\partial_{\nu} A_{\mu}(x) \\
D_{\mu}=\partial_{\mu}+i e A_{\mu}(x) & \mathrm{U}(1) \text { rotations: } \Omega(x)=e^{i e \omega(x)}
\end{array}
$$



Photon field 'picks up' the difference of the gauge transformation - at the ends of a 'Link'

## More on the Lattice

## QCD Lagrangian

$$
\mathscr{L}_{Q C D}(x)=-\frac{1}{4} F_{\mu \nu}^{a}(x) F^{\mu \nu, a}(x)+\bar{\psi}(x)(\gamma \cdot D+m) \psi(x)
$$

## More on the Lattice

Local Gauge Invariance of QFTs

$$
\mathscr{L}_{Q E D}(x)=-\frac{1}{4} F_{\mu \nu}(x) F^{\mu \nu}(x)+\bar{\psi}(x)(\gamma \cdot D+m) \psi(x)
$$

$$
F_{\mu \nu}(x)=\partial_{\mu} A_{\nu}(x)-\partial_{\nu} A_{\mu}(x)
$$

$$
D_{\mu}=\partial_{\mu}+i e A_{\mu}(x)
$$

$$
\Omega(x)=e^{i e \omega(x)}
$$

$$
\psi^{(g)}=e^{i e \omega(x)} \psi, \bar{\psi}^{(g)}(x)=e^{-i e \omega())} \bar{\psi}(x)
$$

$$
A_{\mu}^{(g)}=A_{\mu}-\partial_{\mu} \omega(x)
$$


[^0]:    Excitation Energy between ground and first excited state

