

# Non-thermal behavior in the Lifshitz regime

Lifshitz regime ~ Overhauser-Migdal condensate (OM = chiral spirals):

RDP, V. Skokov & A. Tsvelik, 1801.08156

RDP, F. Rennecke, V. Skokov, S. Valgushev, & A. Tsvelik, 1912.....

Critical region *tiny*, Lifshitz regime is *large*:

Wei-Jie Fu, J. M. Pawłowski, & Fabian Rennecke, 1909.02991

Here: *cartoons* of theory, concentrate on a possible experimental signal

Fluctuations at *non-zero* momentum can be large *before* OM condensates emerge

At (relatively) high  $T$ , low  $\mu$ !

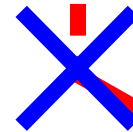
## Usual mean field theory

$$\mathcal{L} = (\partial_\mu \phi)^2 - h\phi + m^2 \phi^2 + \lambda \phi^4 + \kappa \phi^6$$

$m^2 \approx 0, \lambda > 0$ : crossover

$\lambda \uparrow$   
 $\langle \phi \rangle \neq 0$     $\langle \phi \rangle \approx 0$

$m^2 \approx \lambda \approx 0$ : critical endpoint



$m^2 \rightarrow$

$m^2 > 0, \lambda < 0$ : 1<sup>st</sup> order

$\langle \phi \rangle \neq 0$    1<sup>st</sup>  $\uparrow$

Plot out phase diagram by varying  $m^2$  &  $\lambda$ . Ok in mean field theory

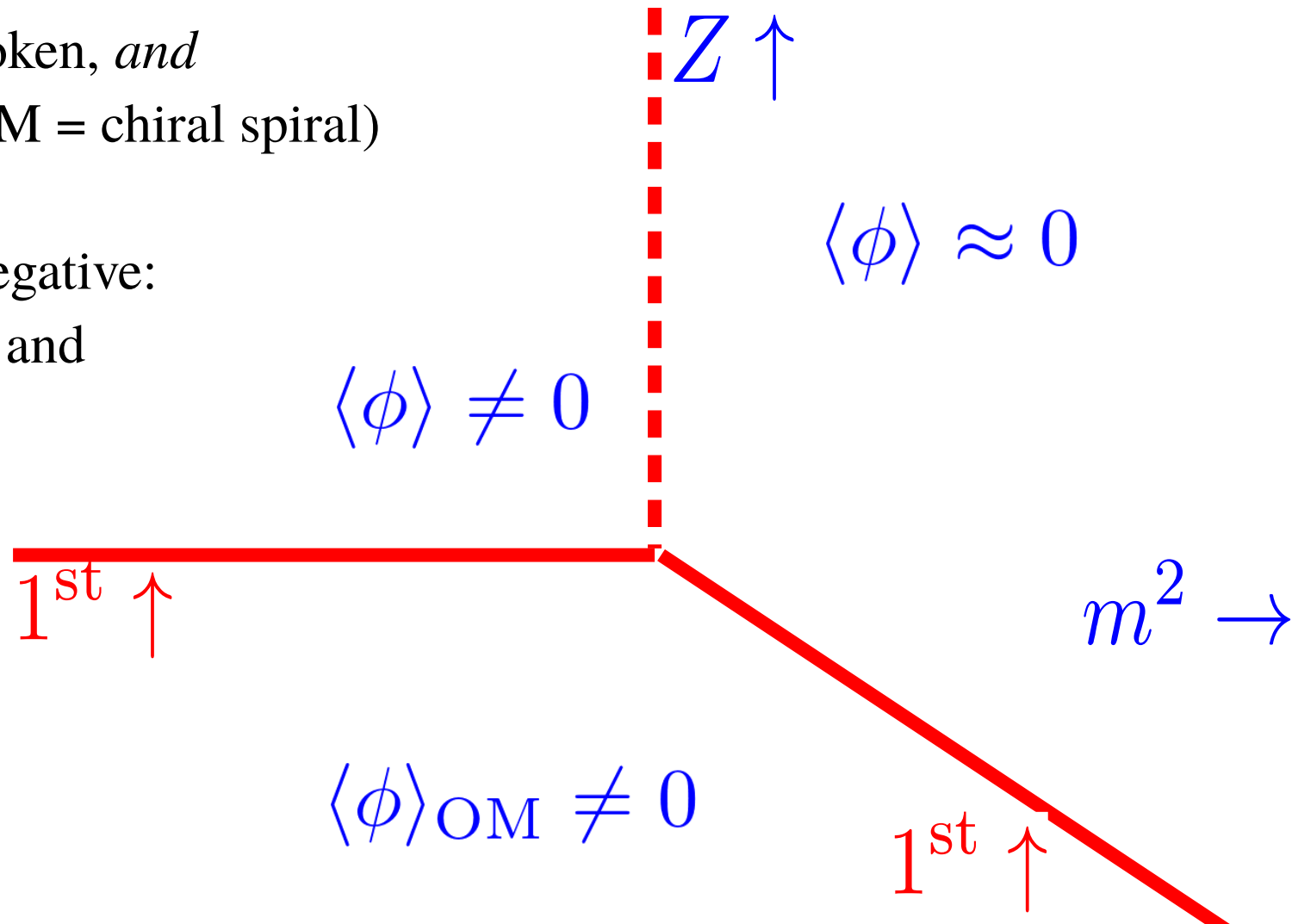
# Lifshitz mean field theory

Time derivatives must be positive (causality), spatial derivatives *not*:

$$\mathcal{L} = (\partial_0 \phi)^2 + Z(\partial_i \phi)^2 + \frac{1}{M^2}(\partial_i^2 \phi)^2 - h\phi + m^2 \phi^2 + \lambda \phi^4 + \kappa \phi^6$$

Phases: symmetric, broken, *and*  
Overhauser-Migdal (OM = chiral spiral)

Both  $Z$  and  $\lambda$  can go negative:  
Lifshitz regime,  $Z < 0$ , and  
critical endpoint

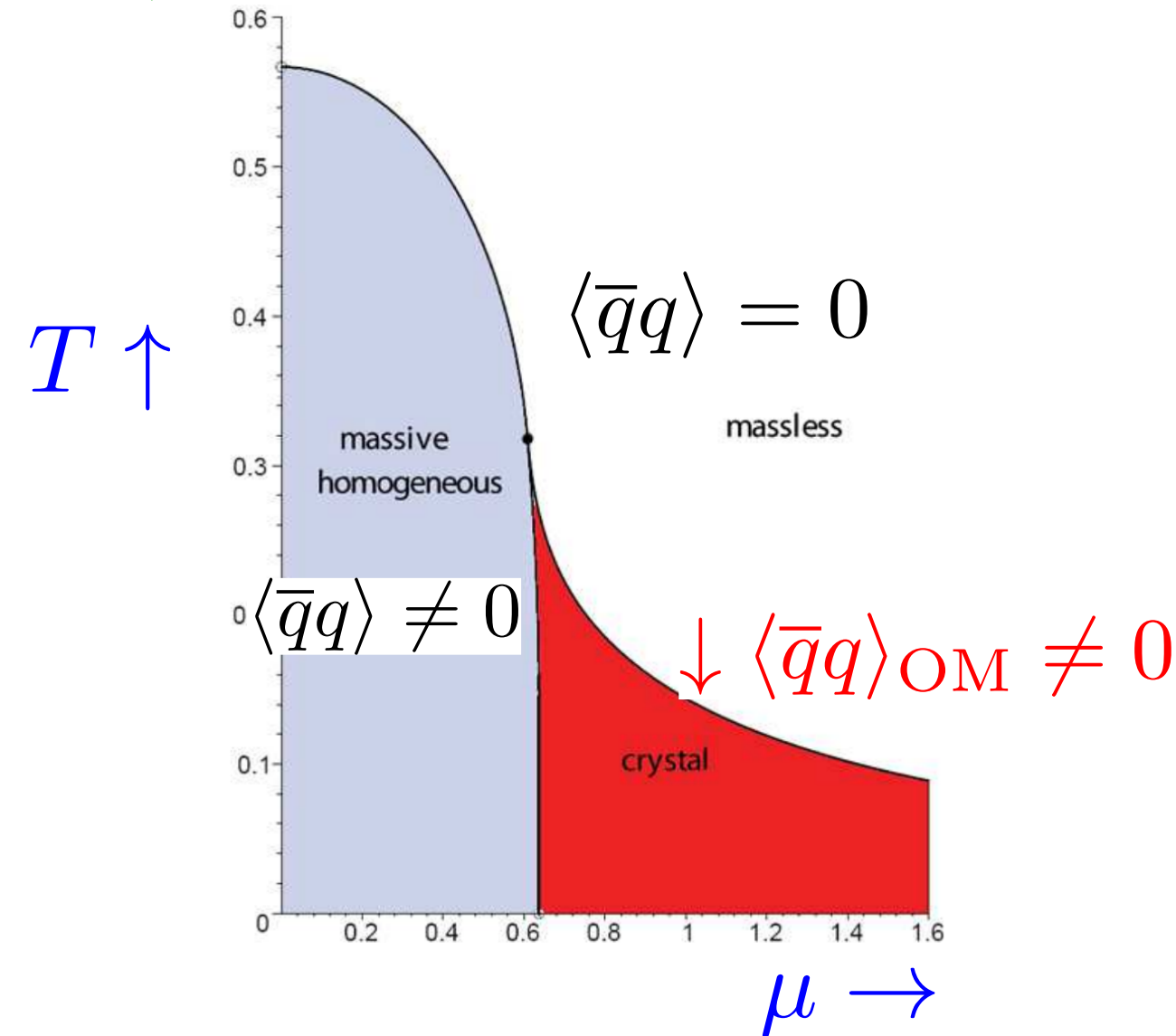


# OM condensates are *everywhere* in 1+1 dim.'s

In 1+1 dim's,  $\mu \neq 0$  can turn constant condensate into OM.

Gross-Neveu type models soluble for large number of flavors,  $N_f$ :

Basar, Dunne & Thies. 0903.1868; Dunne & Thies 1309.2443+ ...

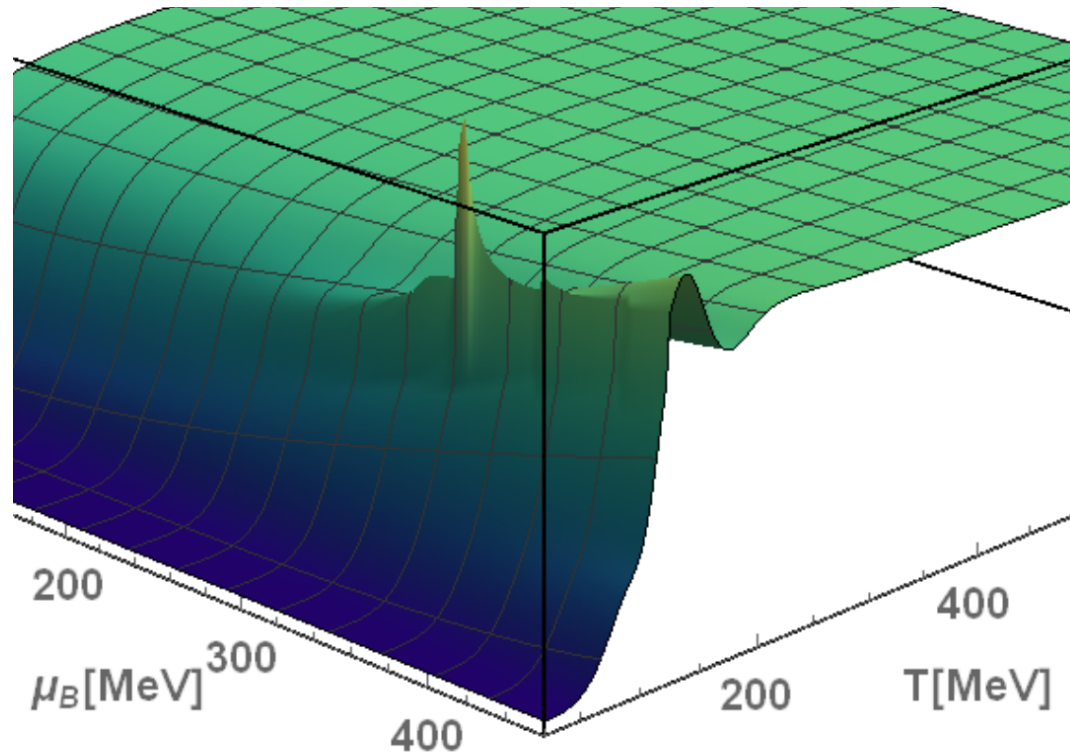


exact solutions,  
 $N_f = \infty$ :



## Width of critical region?

Parotto, Bluhm, Mroczek, Nahrgang, Noronha-Hostler, Rajagopal,  
Ratti, Schaefer, Stephanov, 1805.05249



width of critical region:

$w=1$  in Eqs. 3.1 & 3.2

*$w=1$ ?*

Wide critical region consistent  
with lattice about  $\mu=0$ .

*But is it wide?*

Vovchenko, Steinheimer, Philipsen, Stoecker, 1711.01261:

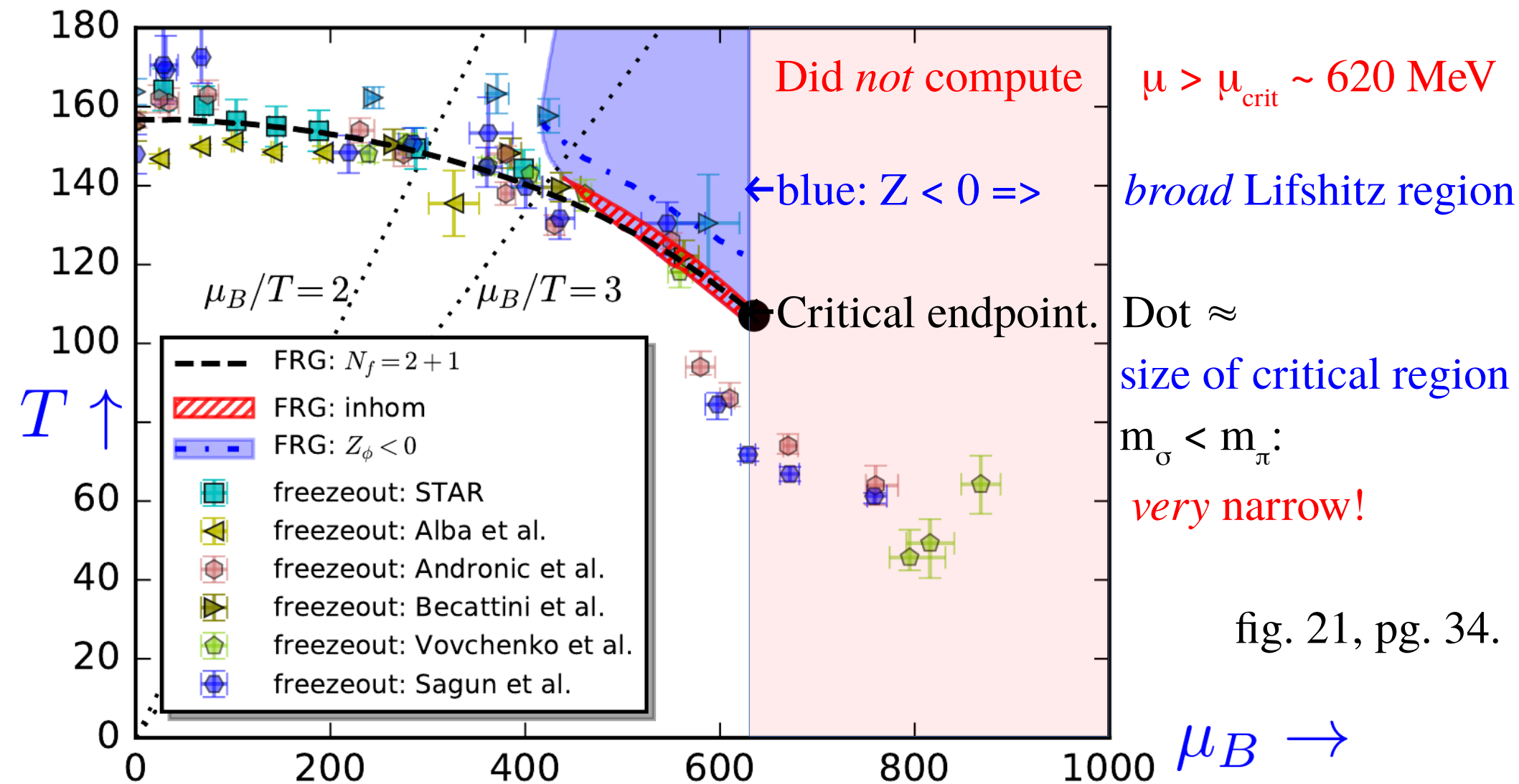
No critical endpoint up to  $\mu_B > \pi T$ .

# Phase diagram from FRG

Fu, Pawłowski, & Rennecke, 1909.02991

Functional Renormalization Group applied to QCD @  $T$  &  $\mu_B$ .

Lifshitz regime broad, critical region tiny.

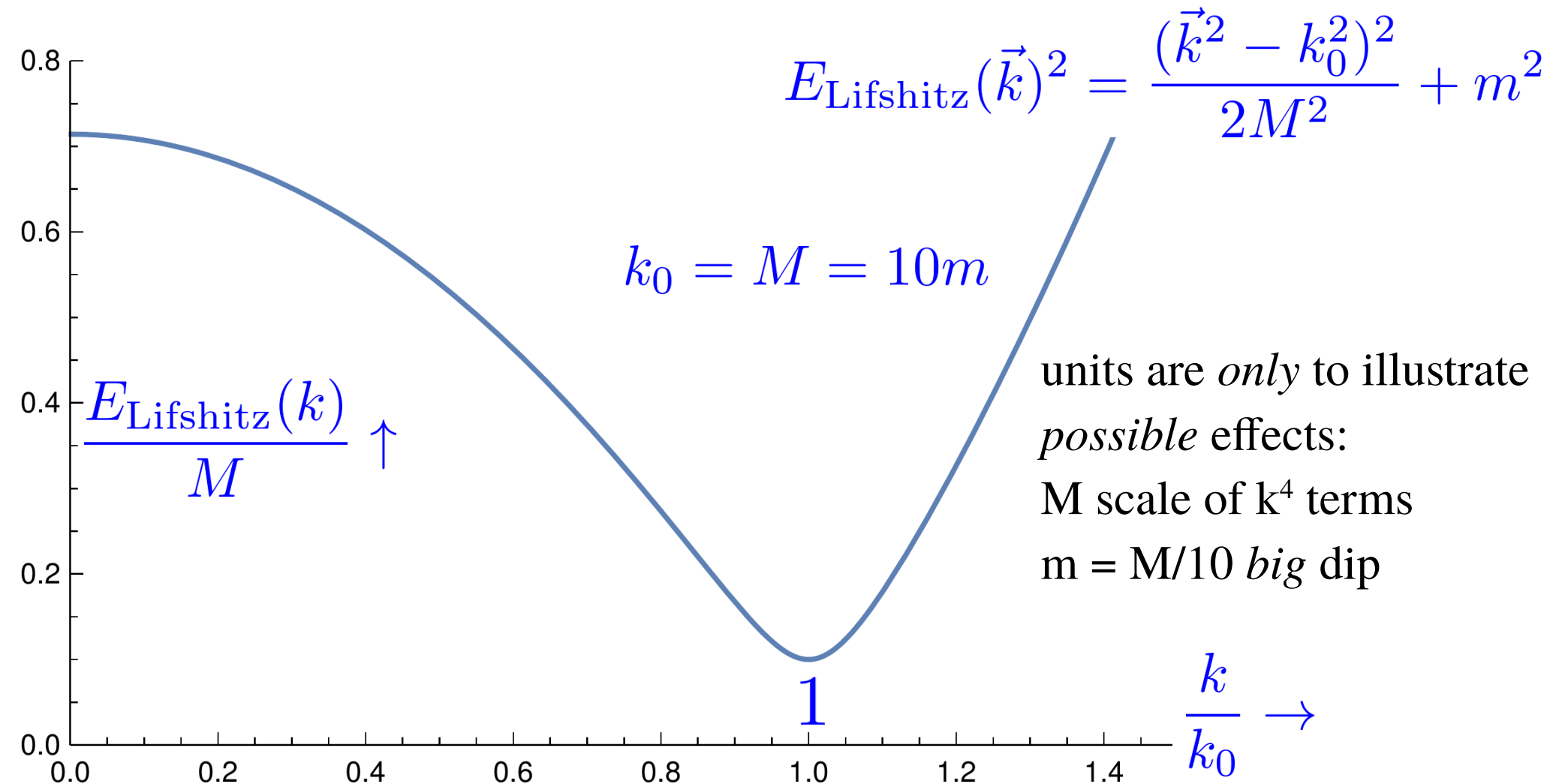


# Cartoon of dispersion relation in the Lifshitz regime

Lifshitz regime:  $Z < 0 \Rightarrow$  minimum of energy at *non-zero* momentum,  $k_0$ .

Overhauser-Migdal condensate when curve hits zero,  $E_{\text{Lifshitz}}(k_0) = 0$ .

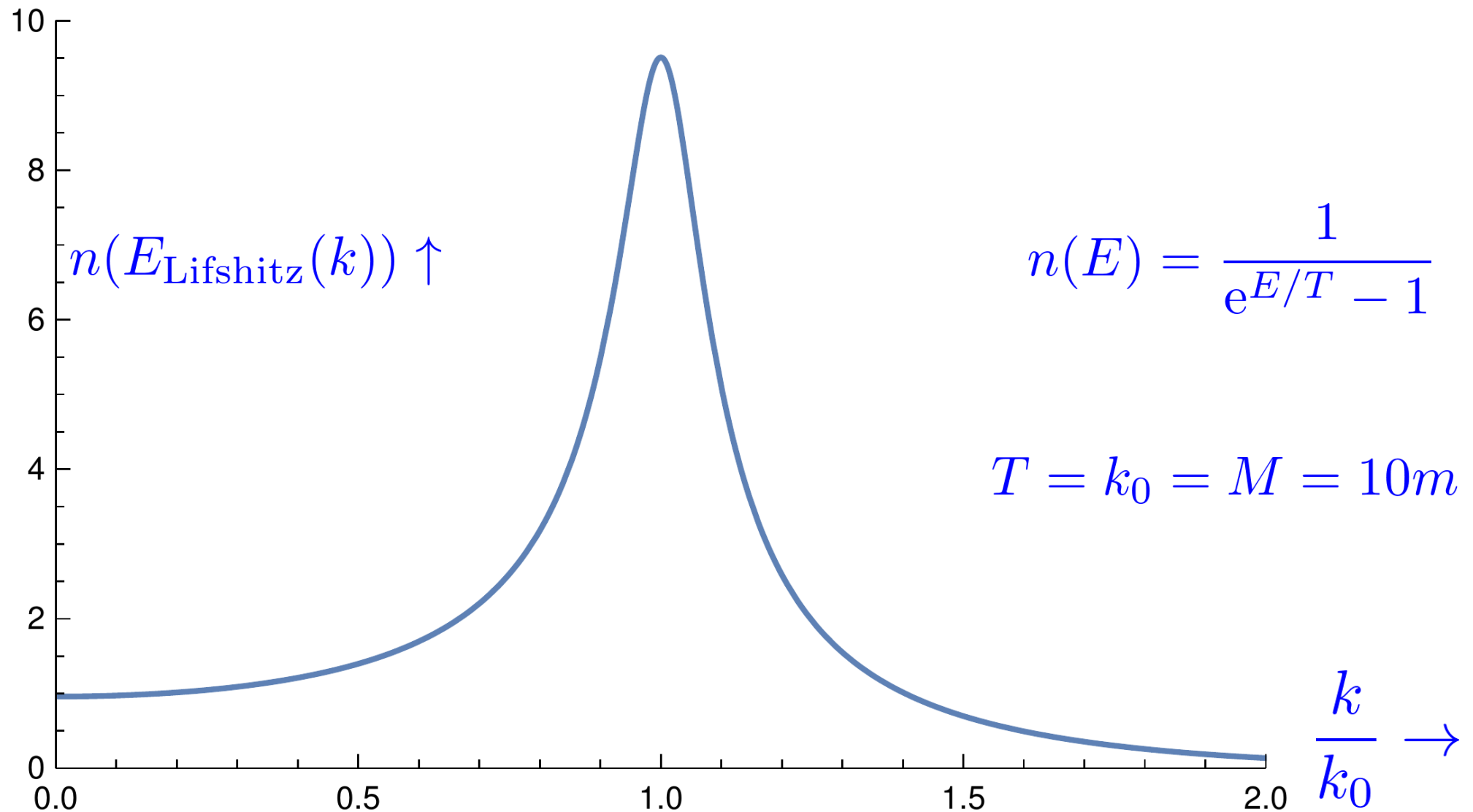
But interesting even when  $E_{\text{Lifshitz}}(k_0) > 0$ !



# Cartoon of Lifshitz statistical distribution function

Lifshitz regime:  $Z < 0 \Rightarrow$  minimum of energy at *non-zero* momentum,  $k_0$ .

Hence peak of  $n(E_{\text{Lifshitz}}(k))$  at  $k_0$ : *possibly, non-thermal behavior can be large.*



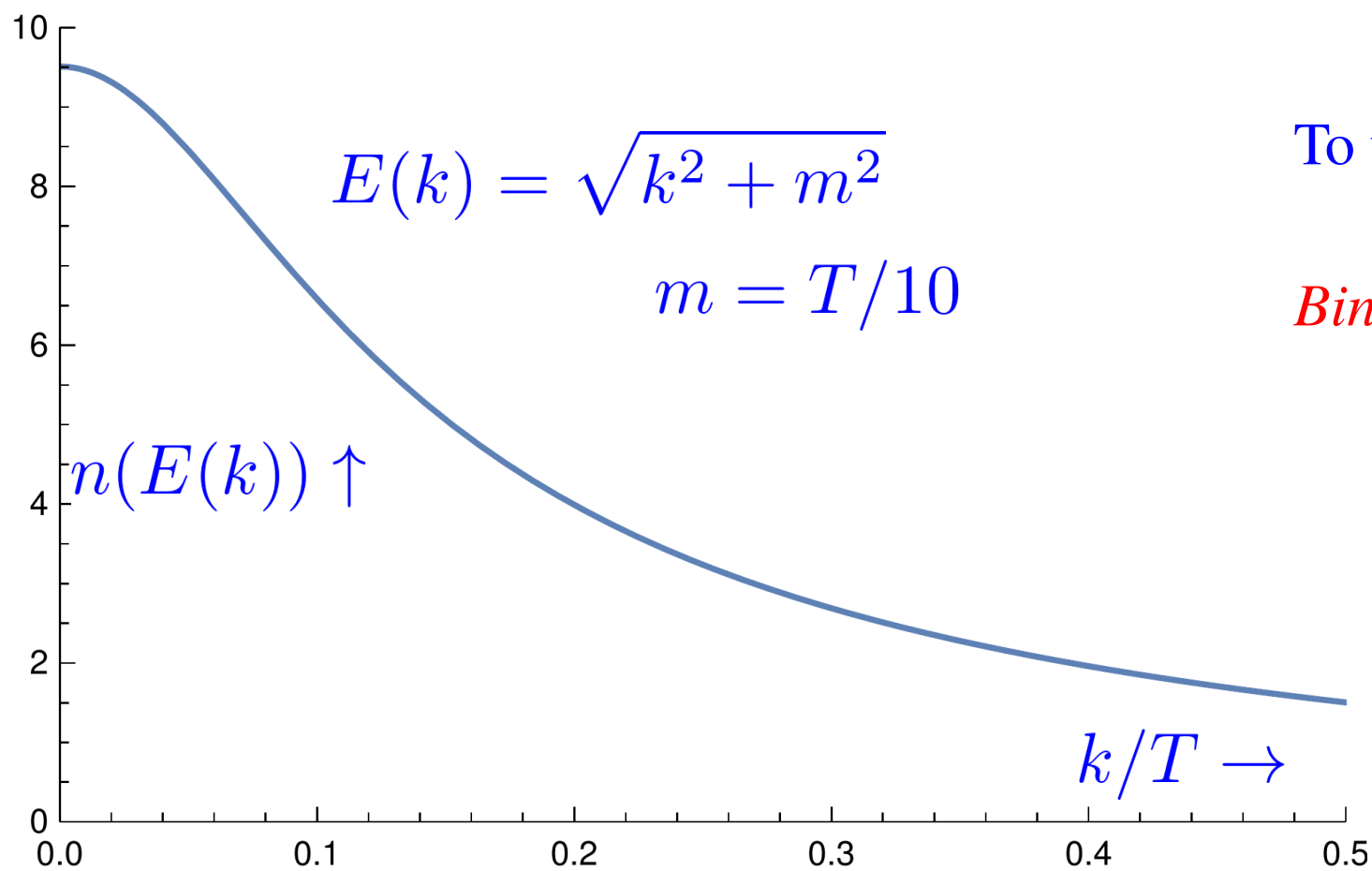


# Cartoon of Lifshitz statistical distribution function

Fluctuations in Lifshitz regime greatest at *non-zero* momentum.

Fluctuations near critical endpoint greatest at *zero* momentum.

Usual  $\chi_2$ ,  $\chi_4$ ,  $\chi_6 \dots$  integral over momenta.

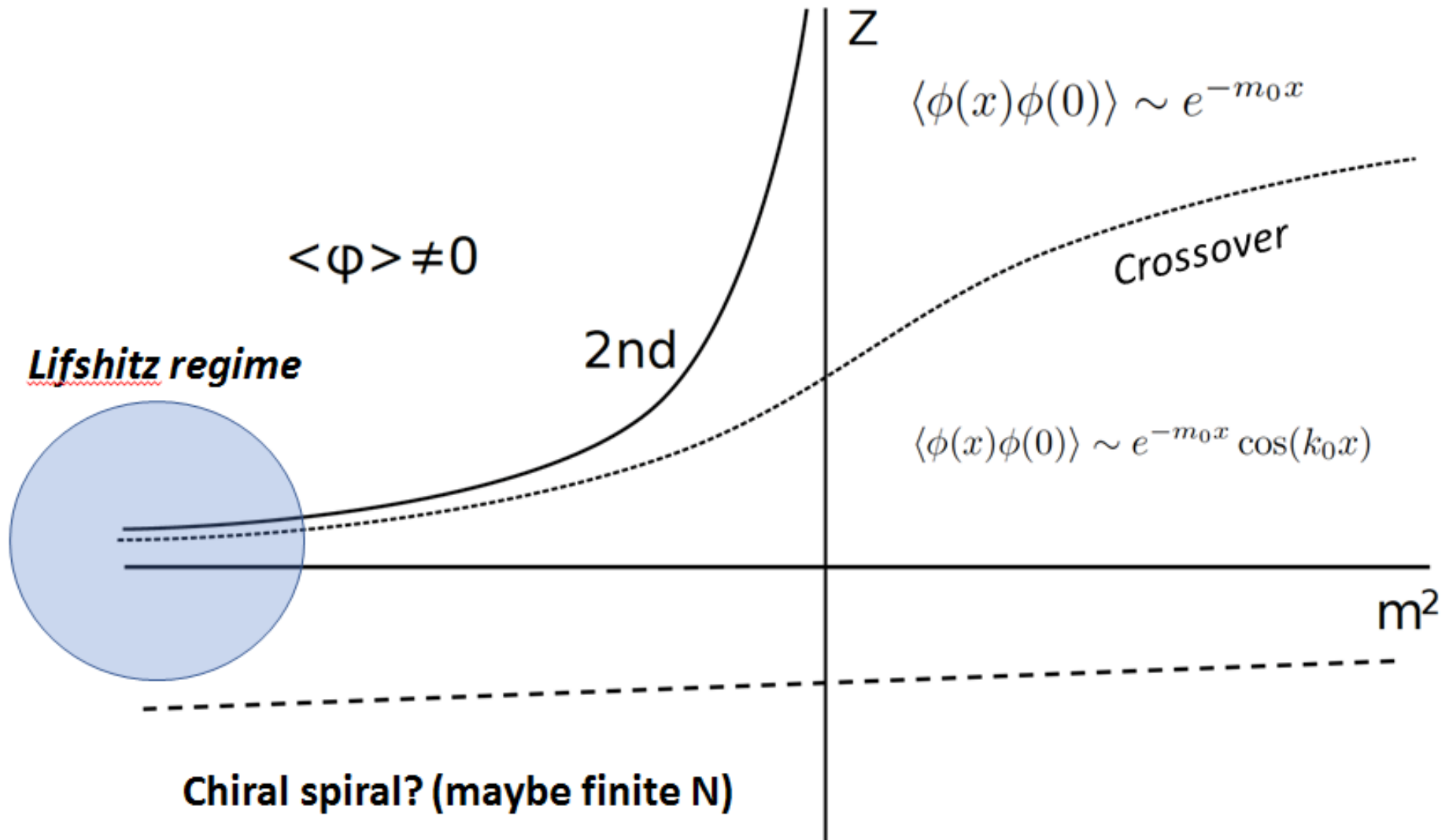


To test Lifshitz regime:

*Bin* in momenta!

## Soluble model at infinite N

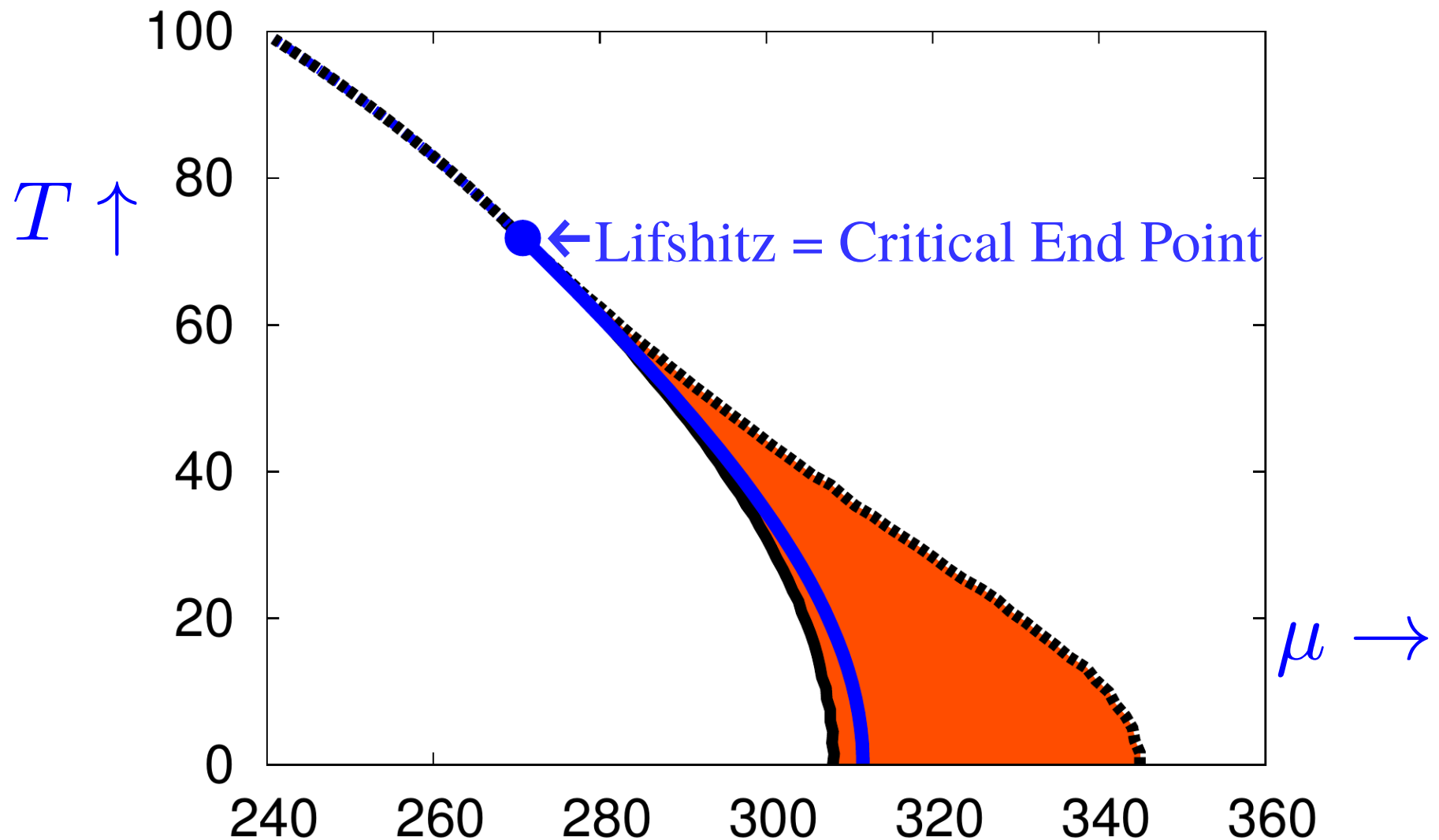
$O(N)$  scalar model soluble at infinite  $N$ . Interesting phase structure in  $Z$  &  $m^2$ .  
Numerical simulations at small  $N$  by Monte Carlo.  
Clear signal for  $Z < 0$ : oscillatory behavior in 2-point function.



# Chiral Spirals in 3+1 dimensions

In 3+1, *common* in NJL models: [Nickel, 0902.1778](#) + ... [Buballa & Carignano 1406.1367](#) + ...

In reduction to 1-dim,  $\Gamma_5^{1\text{-dim}} = \gamma_0 \gamma_z$ , so chiral spiral between  $\bar{q}q$  &  $\bar{q}\gamma_0\gamma_z\gamma_5q$



# Chiral spirals in 1+1 dimensions

In 1+1 dim., can eliminate  $\mu$  by chiral rotation:

$$q' = e^{i\mu z \Gamma_5} q, \quad \bar{q}(\not{D} + i\mu\Gamma_0)q = \bar{q}' \not{D} q', \quad \Gamma_5 \Gamma_z = \Gamma_0$$

Thus a constant chiral condensate automatically becomes a chiral spiral:

$$\bar{q}' q' = \cos(2\mu z) \bar{q} q + i \sin(2\mu z) \bar{q} \gamma_5 q$$

Argument is only suggestive.

N.B.: anomaly ok, gives quark number:  $\langle \bar{q} \Gamma_0 q \rangle = \mu/\pi$

Pairing is between quark & quark-hole, both at edge of Fermi sea.

Thus chiral condensate varies in  $z$  as  $\sim 2\mu$ .

# Fluctuations at 7 GeV

Beam Energy Scan, down to 7 GeV.

Fluctuations *MUCH* larger when up to 2 GeV than to 0.8 GeV

Trivial multiplicity scaling? ... or Chiral Spiral?

But fluctuations in nucleons, not pions.

X. Luo & N. Xu, 1701.02105, fig. 37; Jowazee, 1708.03364

