

Recent results on event-by-event fluctuations in ALICE

Mesut Arslandok

Physikalisches Institut, Heidelberg University
on behalf of the ALICE Collaboration

Quark Matter 2019

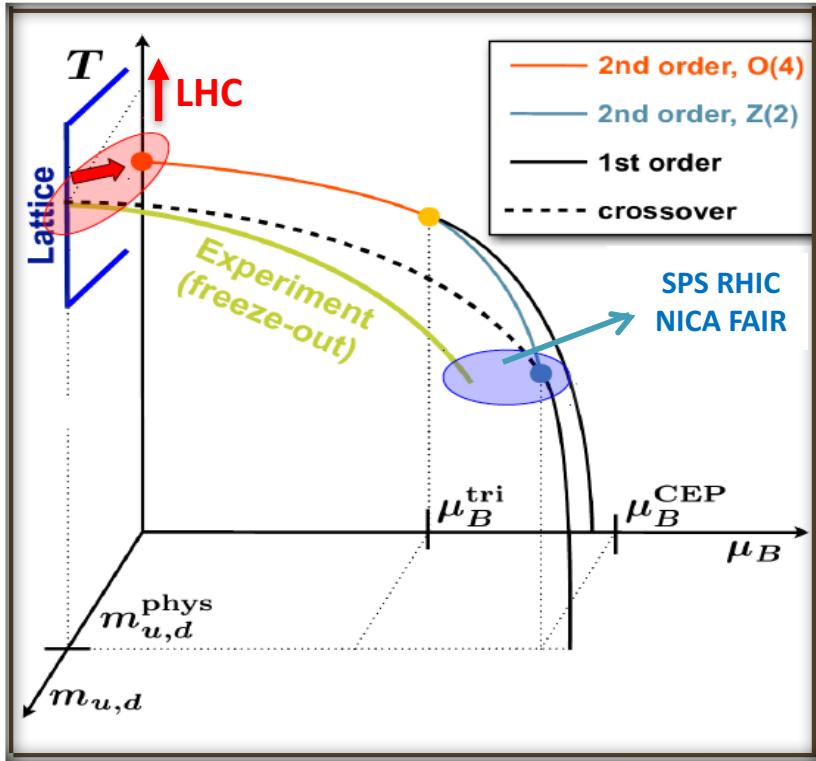
The 28th International Conference on Ultra-relativistic Nucleus-Nucleus Collisions
3-9 November 2019, Wuhan, China



UNIVERSITÄT
HEIDELBERG
ZUKUNFT
SEIT 1386



Nature of chiral phase transition



F. Karsch, Schleching 2016

small u, d quark masses
↔
vicinity to 2nd order O(4) criticality

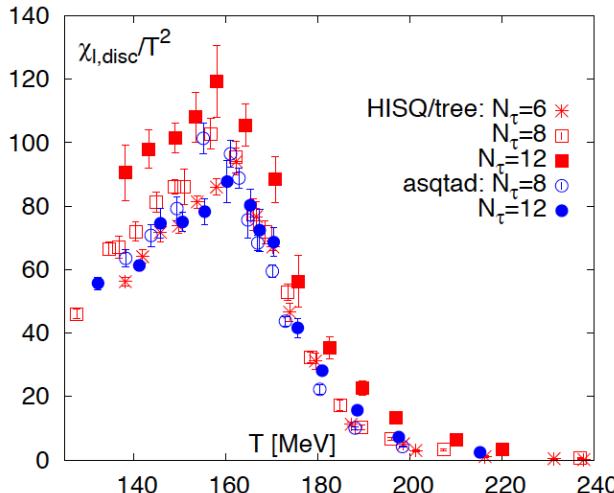
↓

pseudocritical features possible

Criticality at Crossover

HotQCD Collaboration

Phys.Rev. D85 (2012) 054503, Phys.Lett. B795 (2019) 15

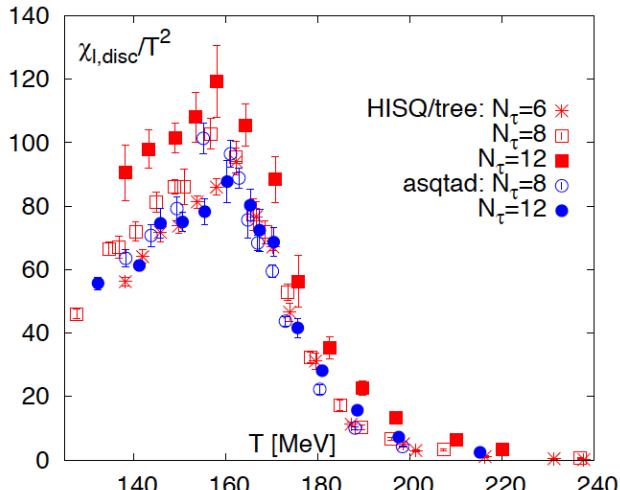


$$T_{pc} = 156.5 \pm 1.5 \text{ MeV}$$

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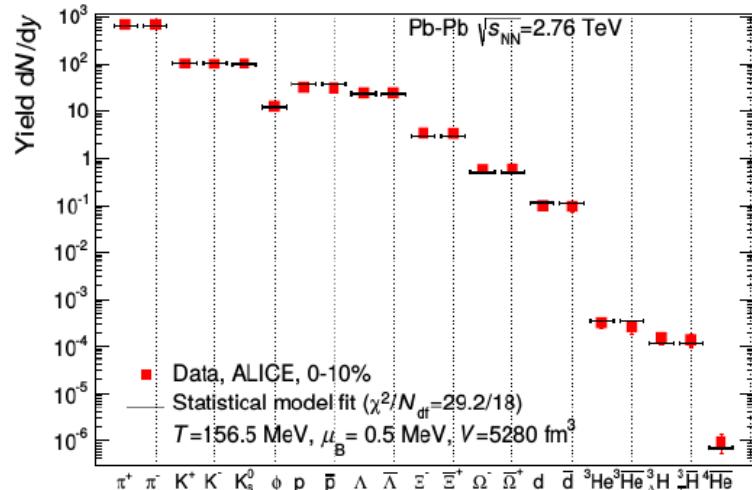
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A. Andronic, P. Braun-Munzinger, J. Stachel and K. Redlich

Nature 561, 321–330 (2018)

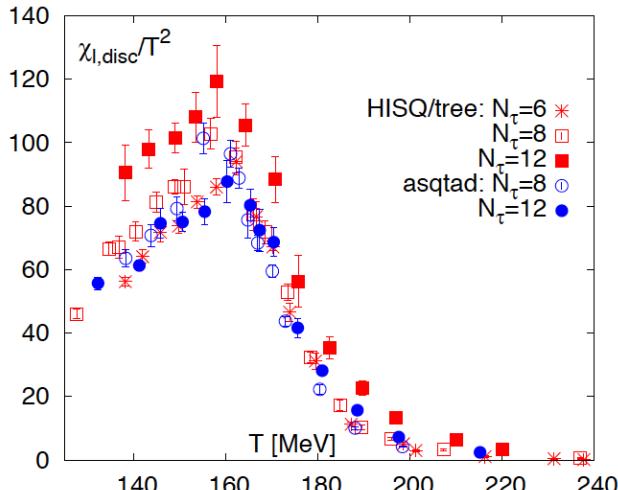


$$T_{\text{fo}}^{\text{ALICE}} = 156.5 \pm 3 \text{ MeV}$$

Criticality at Crossover

HotQCD Collaboration

Phys.Rev. D85 (2012) 054503, Phys.Lett. B795 (2019) 15



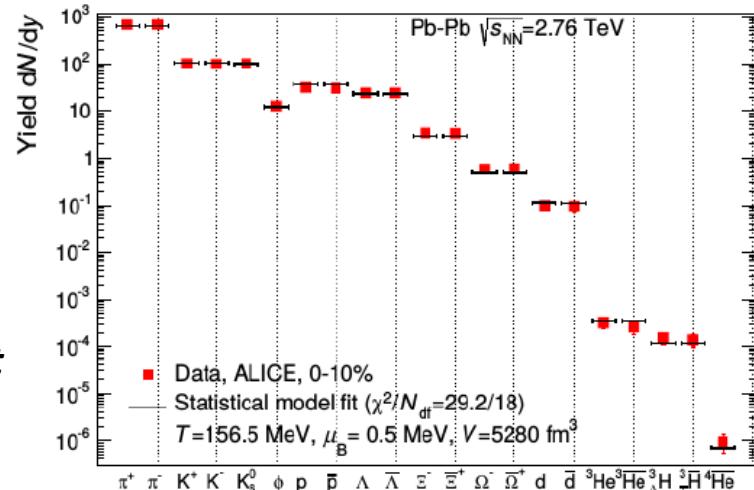
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*Chemical freeze-out
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phase boundary!*



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$$T_{fo}^{\text{ALICE}} = 156.5 \pm 3 \text{ MeV}$$

Chemical freeze-out near T_{pc} → motivation to look for higher order moments

Link to LQCD: Fluctuations of conserved charges

For a thermal system within the **Grand Canonical Ensemble**

$$\frac{P}{T^4} = \frac{1}{VT^3} \ln Z(V, T, \mu_{B,Q,S}) \rightarrow \hat{\chi}_n^{N=B,S,Q} = \frac{\partial^n P/T^4}{\partial (\mu_N/T)^n}$$

Susceptibilities

$$\rightarrow \hat{\chi}_2^B = \frac{\kappa_2(\Delta N_B)}{VT^3}$$

Cumulants

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Higher orders

P. Braun-Munzinger, A. Rustamov, J. Stachel
Nucl. Phys. A960 (2017) 114

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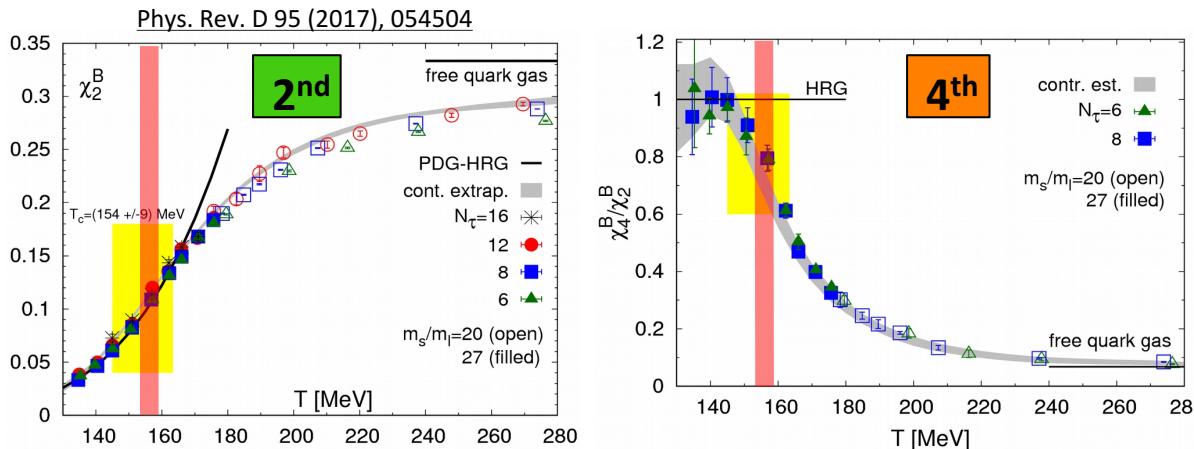
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- At 4th order LQCD shows a deviation ($\sim 30\%$ from unity) from Hadron Resonance Gas (HRG)

Baseline: Skellam distribution

$$X = N_B - N_{\bar{B}}$$

- **rth central moment:**

$$\mu_r \equiv \langle (X - \langle X \rangle)^r \rangle = \sum_X (X - \langle X \rangle)^r P(X)$$

- **First four cumulants**

$$\kappa_1 = \langle X \rangle, \quad \kappa_2 = \mu_2,$$

$$\kappa_3 = \mu_3, \quad \kappa_4 = \mu_4 - 3\mu_2^2$$

- **Uncorrelated Poisson limit:**

$$\langle N_B N_{\bar{B}} \rangle = \langle N_B \rangle \langle N_{\bar{B}} \rangle$$

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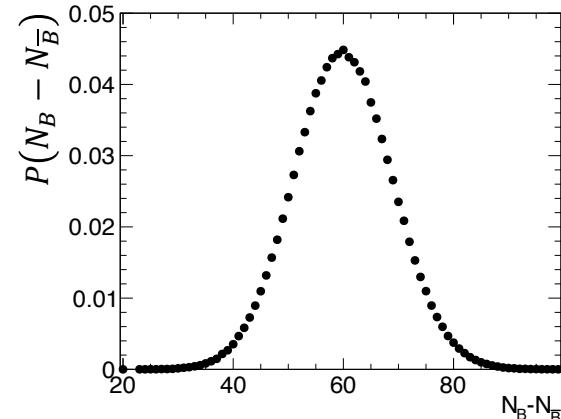
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Difference between two independent Poissonian distributions

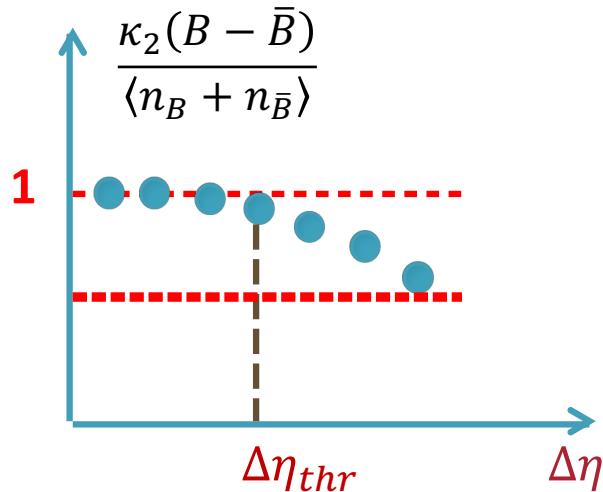
$$\kappa_n = \langle N_B \rangle + (-1)^n \langle N_{\bar{B}} \rangle$$



$$\frac{\kappa_{2n+1}}{\kappa_{2k}} = \frac{\langle n_B \rangle - \langle n_{\bar{B}} \rangle}{\langle n_B \rangle + \langle n_{\bar{B}} \rangle}$$

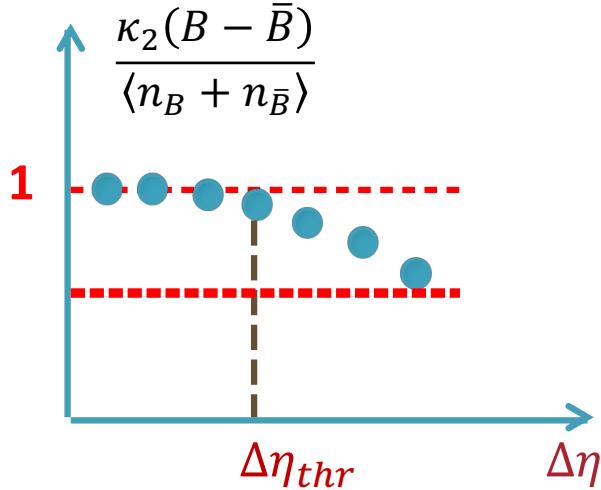
Importance of acceptance and baryon number conservation

- Fluctuations of conserved charges appear **only inside finite acceptance**
- **In the limit of very small acceptance**
→ only Poissonian fluctuations



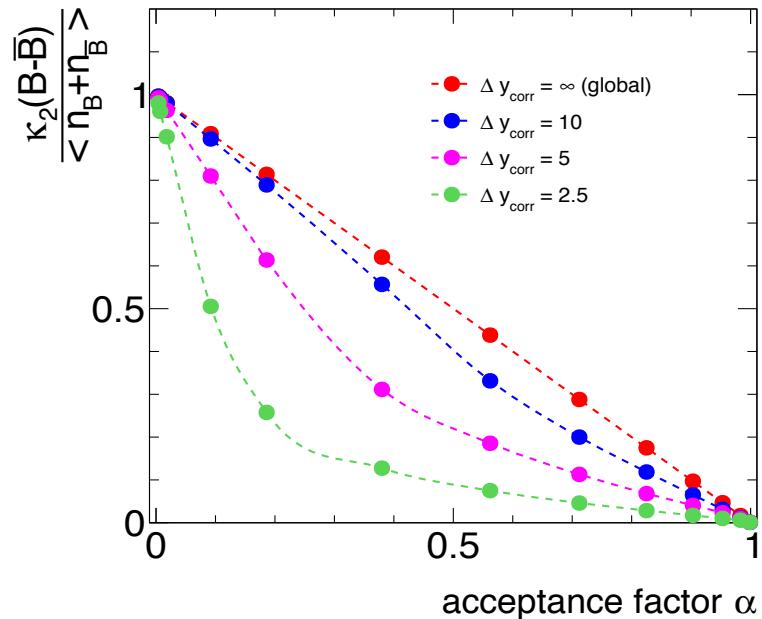
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- **Baryon number conservation** imposes subtle correlations

$$\alpha = \frac{\langle N_B^{acc} \rangle}{\langle N_B^{4\pi} \rangle} \quad |y_{\bar{B}} - y_B| < \frac{\Delta y_{corr}}{2}$$



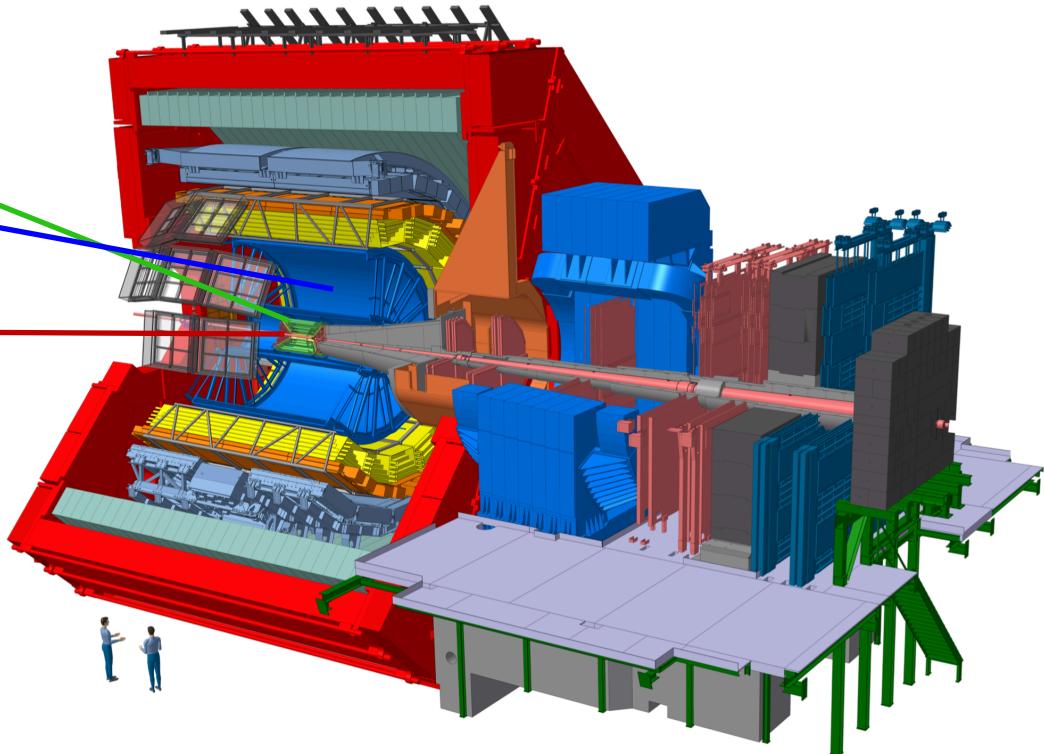
P. Braun-Munzinger, A. Rustamov, J. Stachel, arXiv:1907.03032

RESULTS

A Large Ion Collider Experiment

Main detectors used:

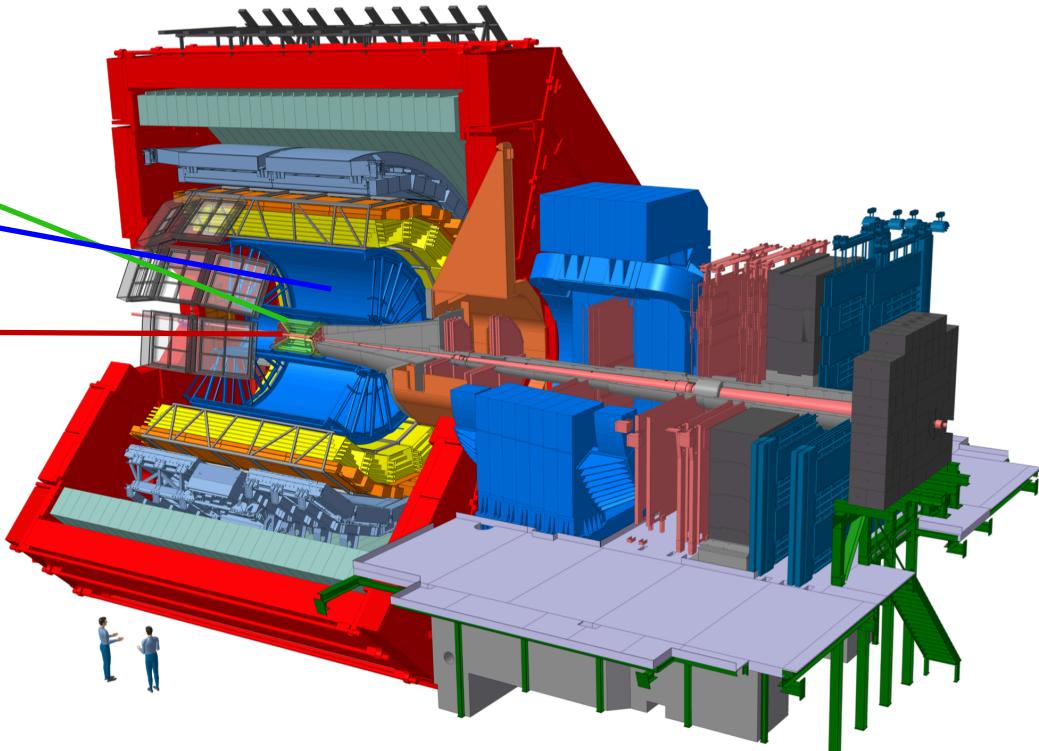
- Inner Tracking System (**ITS**)
→ Tracking and vertexing
- Time Projection Chamber (**TPC**)
→ Tracking and
Particle Identification (PID)
- Vertex 0 (**V0**)
→ Centrality determination



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Data Set:

- $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$, $\sim 78 \text{ M events}$
- $\sqrt{s_{\text{NN}}} = 2.76 \text{ TeV}$, $\sim 12 \text{ M events}$

Kinematic acceptance:

- $0.6 < p < [1.5, 2] \text{ GeV}/c$
- $|\eta| < 0.2, 0.4, \dots, 0.8$

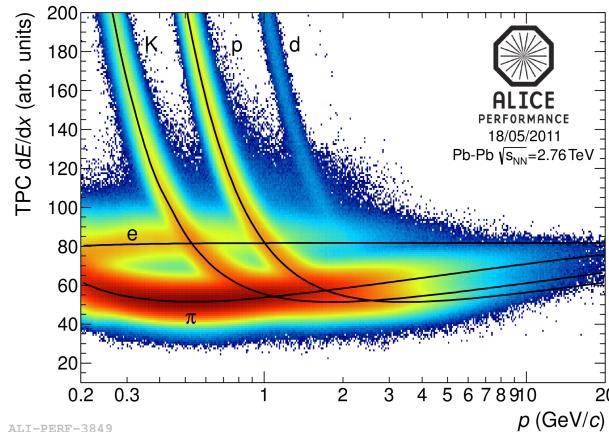
Identity Method

Cut-based approach:

- count tracks of a given particle type

Identity method:

- count probabilities to be of a given particle type



A. Rustamov, M. Gazdzicki, M. I. Gorenstein, PRC 86, 044906 (2012),
M. Arslandok, A. Rustamov, NIM A 946 (2019) 162622

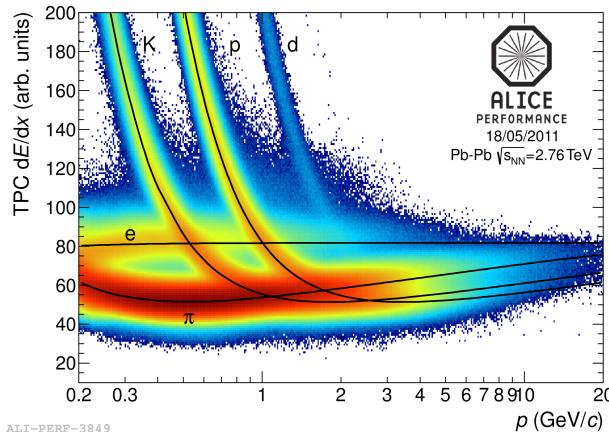
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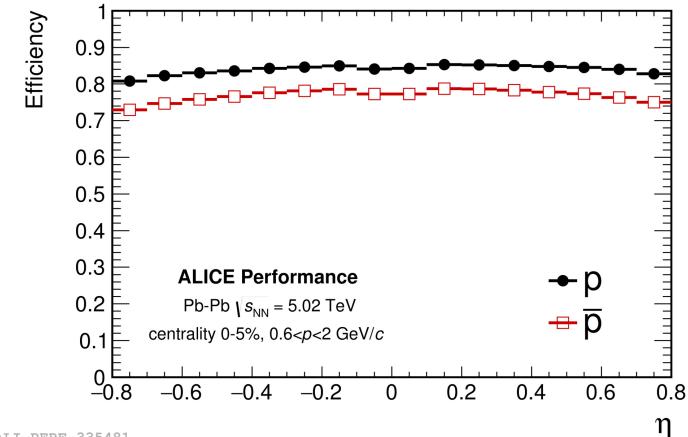
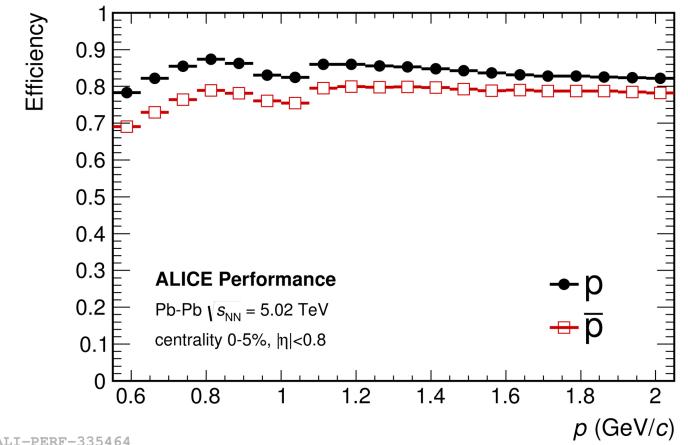
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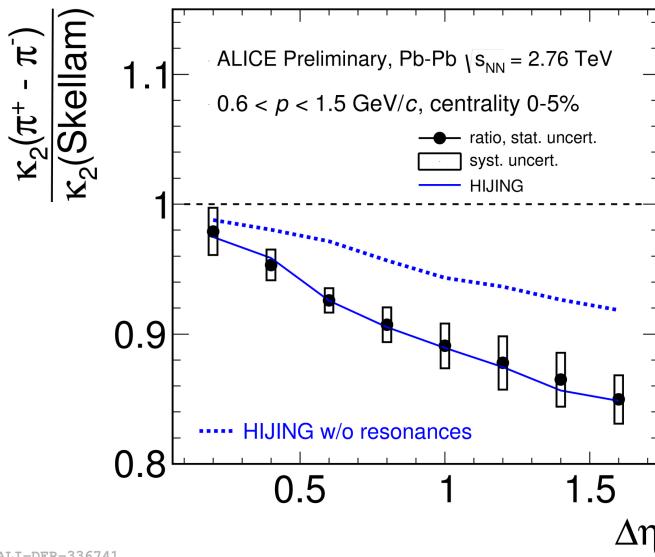
Advantages:

- Gives folded multiplicity distribution
- Allows for **larger efficiencies** → smaller correction

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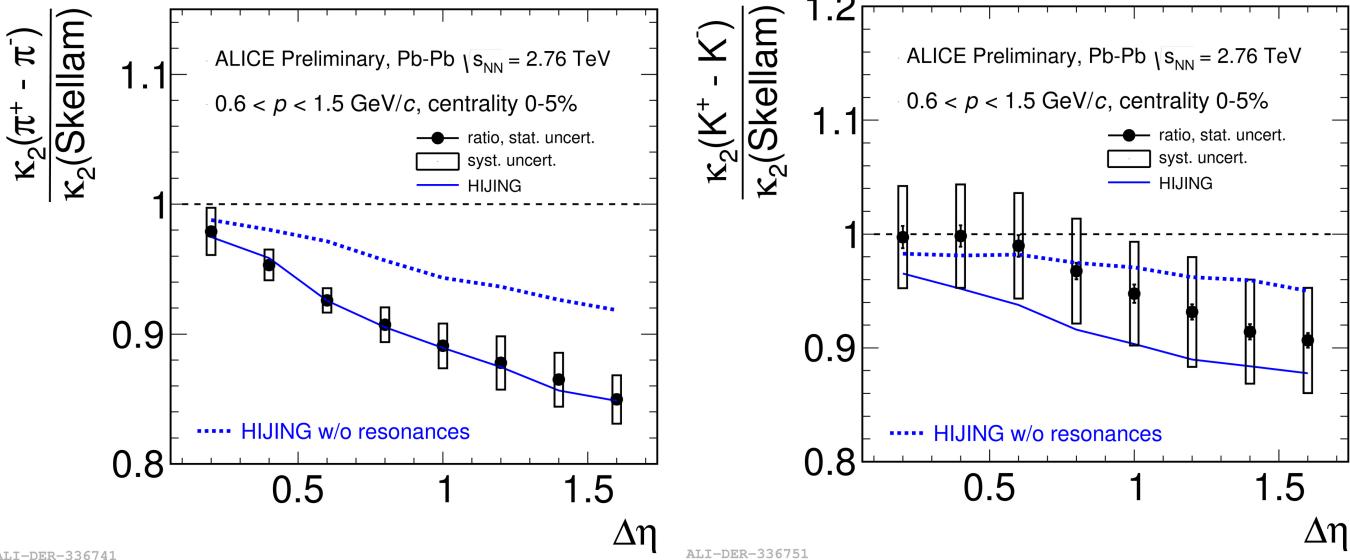


Net-(global)charge fluctuations



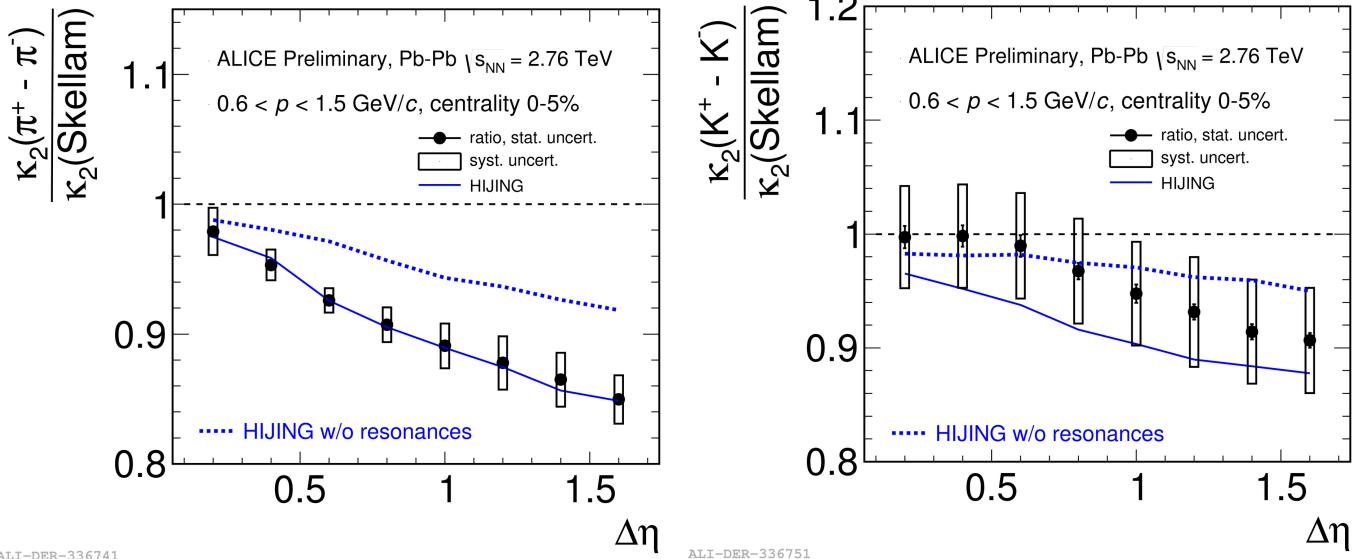
- **Net-electric-charge:** → Strongly dominated by **resonance contributions** (Poster: Shaista Khan)

Net-(global)charge fluctuations



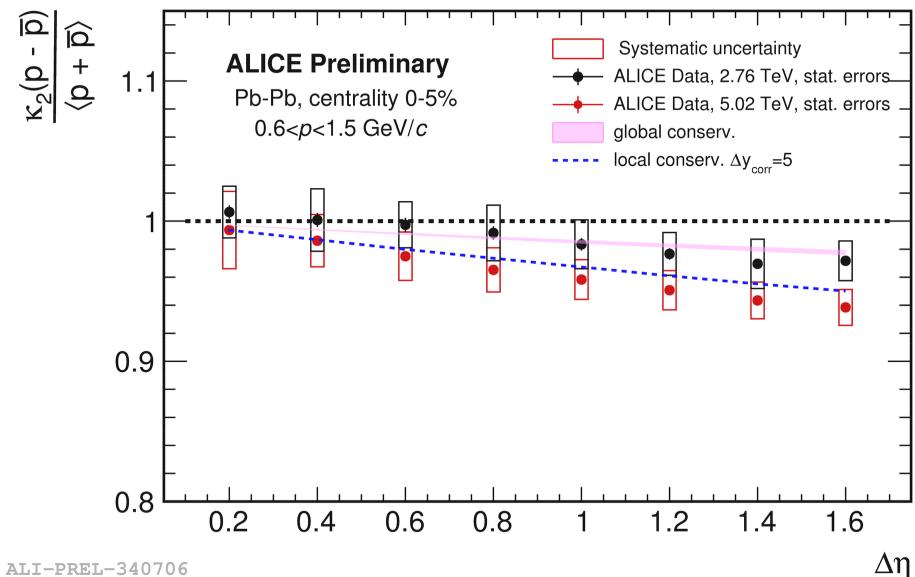
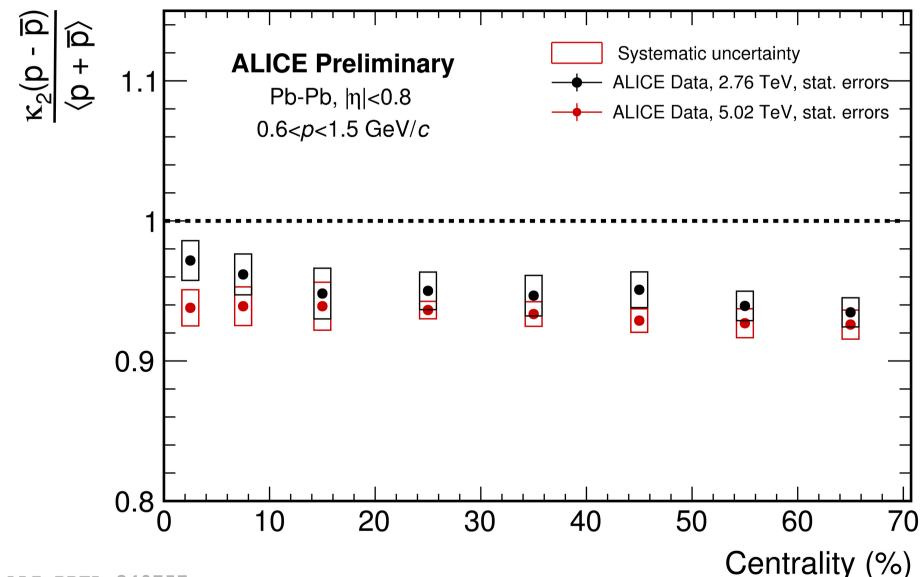
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Net-(global)charge fluctuations



- **Net-electric-charge:** → Strongly dominated by **resonance contributions** (Poster: Shaista Khan)
- **Net-strangeness:** → Kaons are dominated by Φ -decay
- **Net-baryon:**
 - Due to **isospin randomization**, at $\sqrt{s_{\text{NN}}} > 10 \text{ GeV}$ **net-baryon** fluctuations can be obtained from corresponding **net-proton** measurements ([M. Kitazawa, and M. Asakawa, Phys. Rev. C 86, 024904 \(2012\)](#))
 - No resonance feeding $p + \bar{p}$
 - **Best candidate for measuring charge susceptibilities**

2nd order cumulants of net-p: Energy dependence

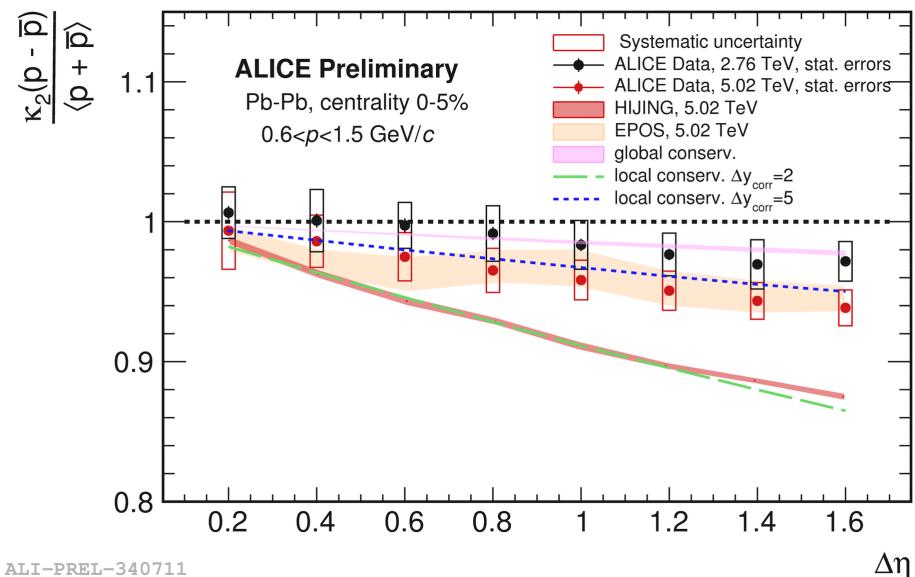
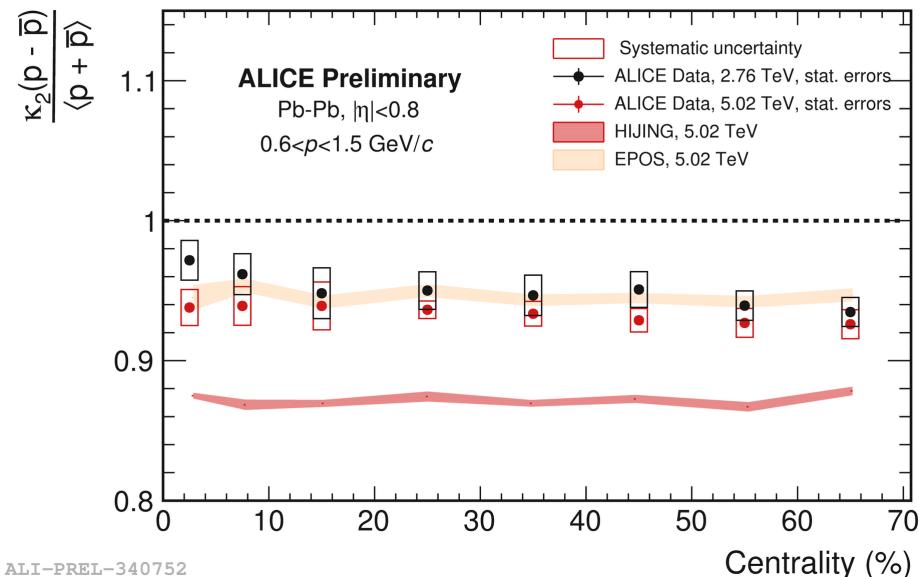


ALI-PREL-340757

ALI-PREL-340706

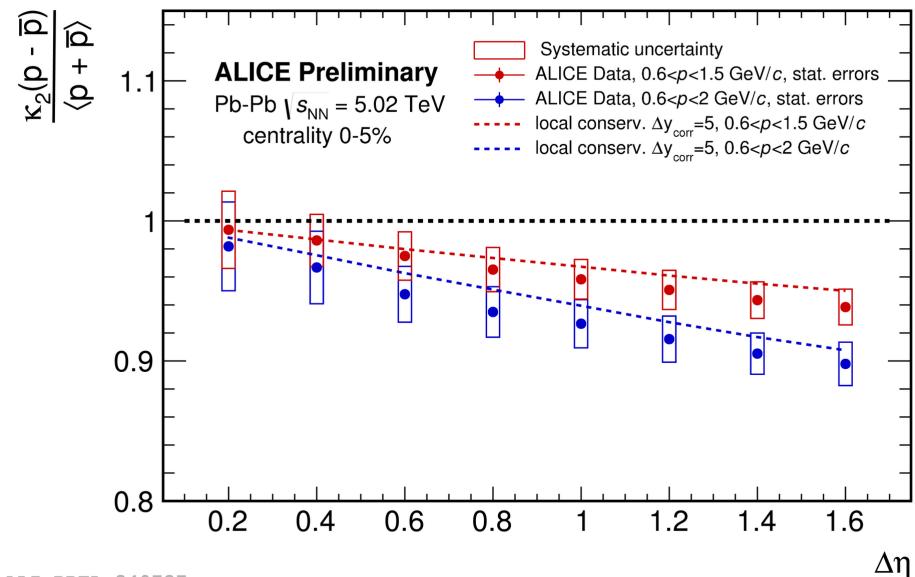
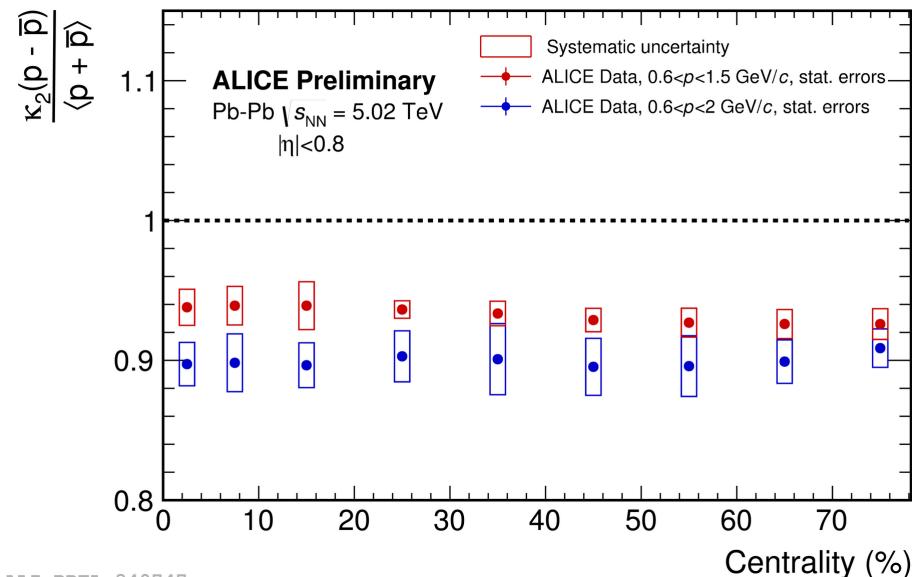
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- ALICE data suggest **long range correlations**, $\Delta y = \pm 2.5$ unit or longer

2nd order cumulants of net-p: Energy dependence



- Deviation from Skellam baseline is due to **baryon number conservation**
- ALICE data suggest **long range correlations**, $\Delta y = \pm 2.5$ unit **or longer**
- EPOS agrees with ALICE data but HIJING deviates significantly
 - Event generators based on string fragmentation (HIJING) conserve baryon number over $\Delta y = \pm 1$ unit

2nd order cumulants of net-p: Acceptance dependence

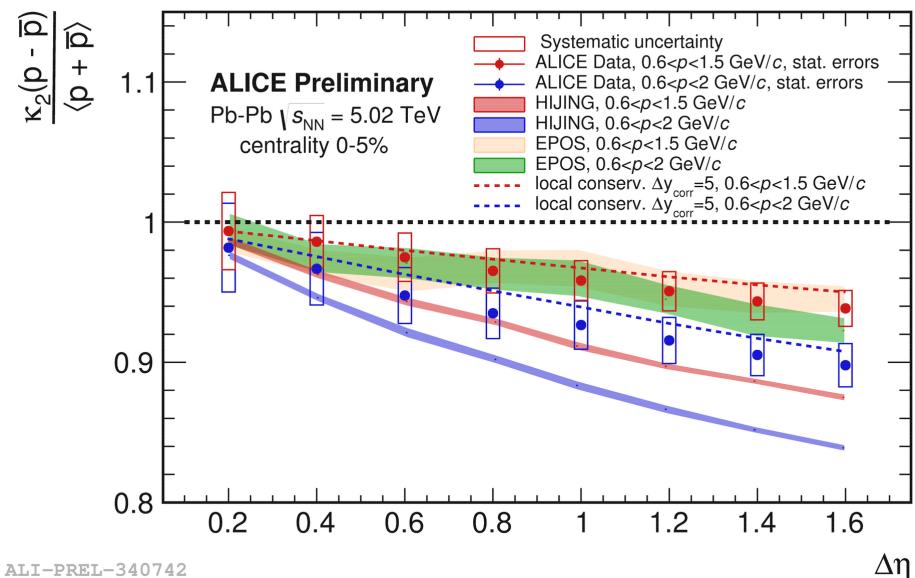
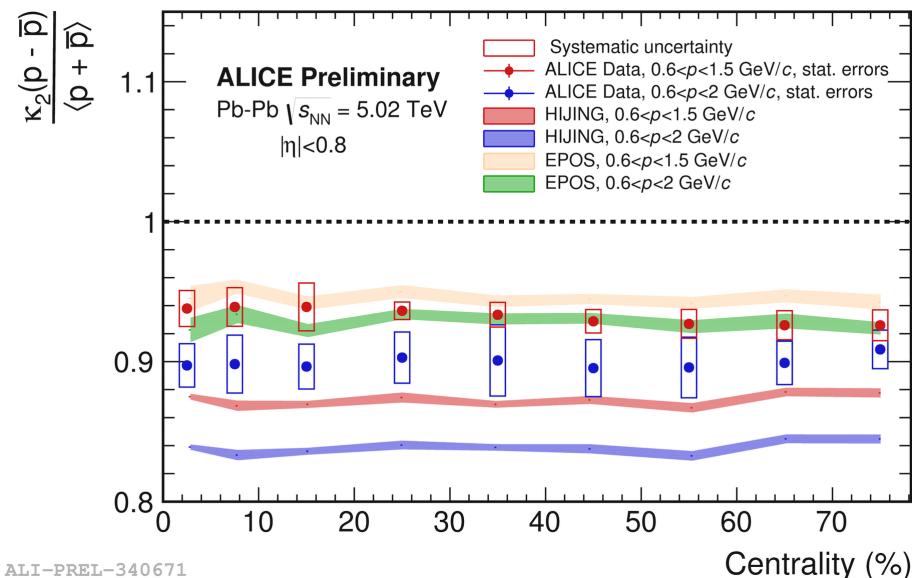


ALI-PREL-340747

ALI-PREL-340737

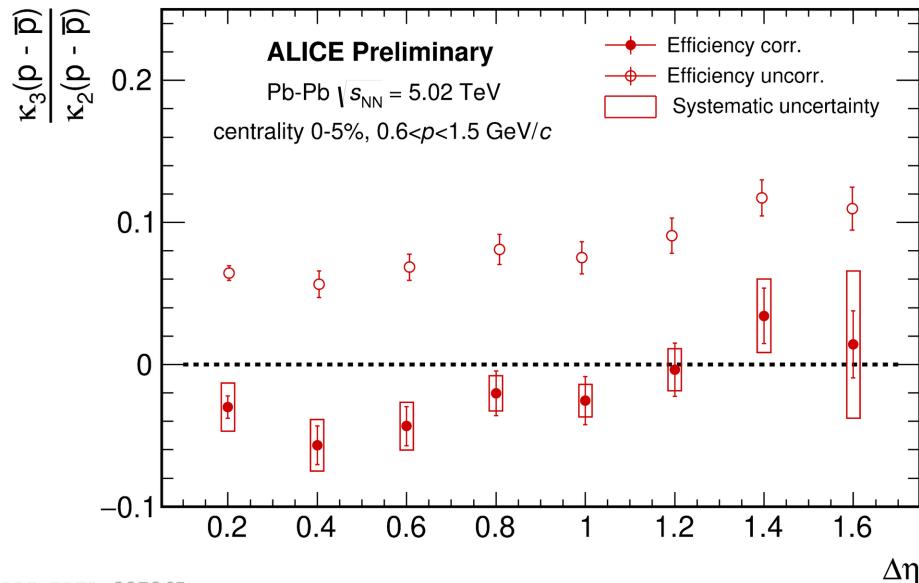
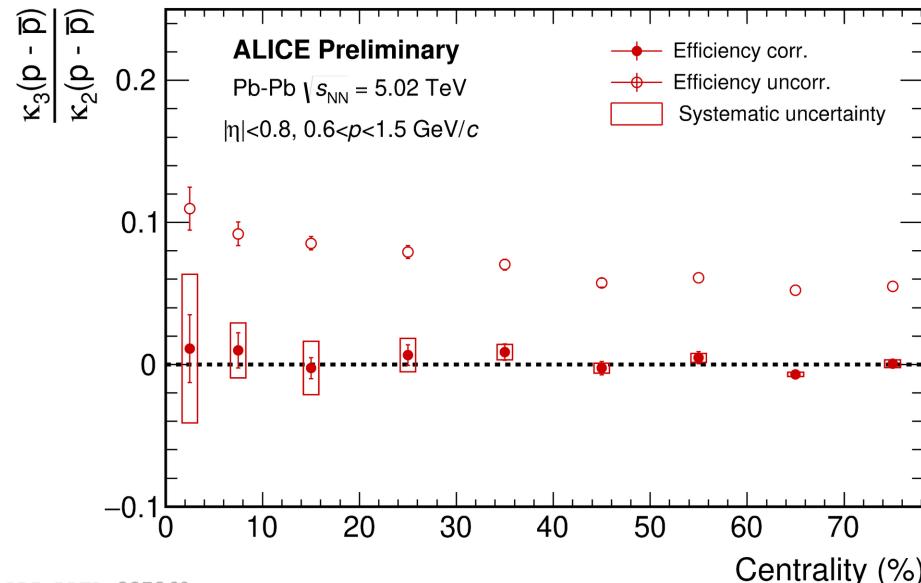
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 - Increase in fraction of accepted $p, \bar{p} \rightarrow$ stronger constraint of fluctuations due to baryon number conservation

2nd order cumulants of net-p: Acceptance dependence



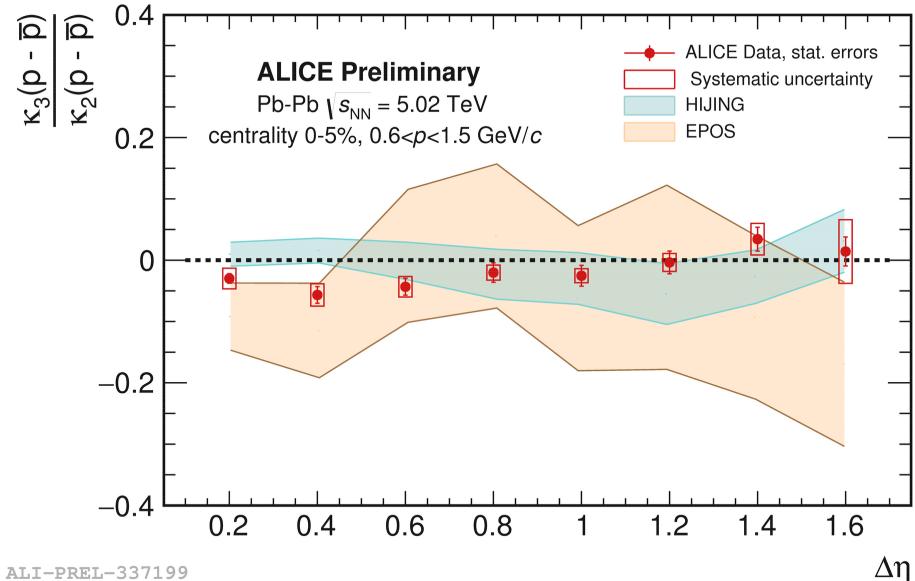
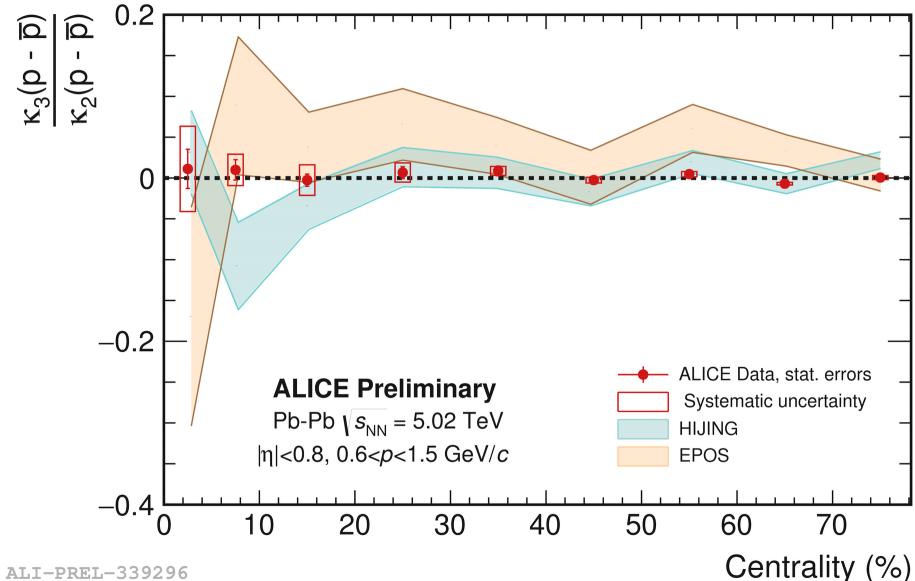
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- EPOS & HIJING show this drop qualitatively

3rd order cumulants of net-p: Centrality and pseudorapidity dependence



- Data agree with Skellam baseline “0” as a function of centrality and pseudorapidity
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- **EPOS and HIJING in agreement with data**
 - Both models conserve global charge → net-p within acceptance is ~ 0

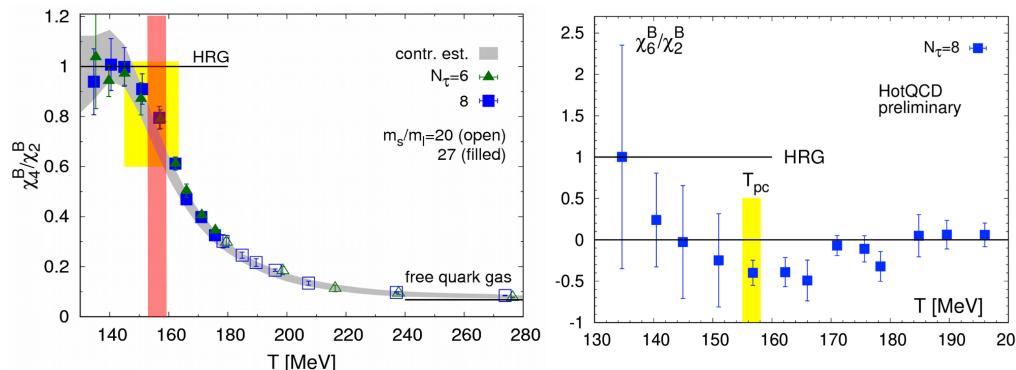
Summary & Outlook

- **Net-electric-charge fluctuations:** Challenge are the dominant **resonance contributions**
 - **Net-proton fluctuations:**
 - ✓ **1st order:** $T_{fo}^{ALICE} \sim T_{pc}^{LQCD}$
 - ✓ **2nd order:** Deviation from Skellam baseline is due to baryon number conservation
 - ALICE data suggests **long range correlations**
 - ✓ **3rd order:** Agrees with Skellam baseline “0” as a function of centrality and pseudorapidity
 - Achieved precision of **better than 5%** for the κ_3/κ_2 results is promising for the higher order cumulants
 - **Up to 3rd order** ALICE data agree with the LQCD expectations
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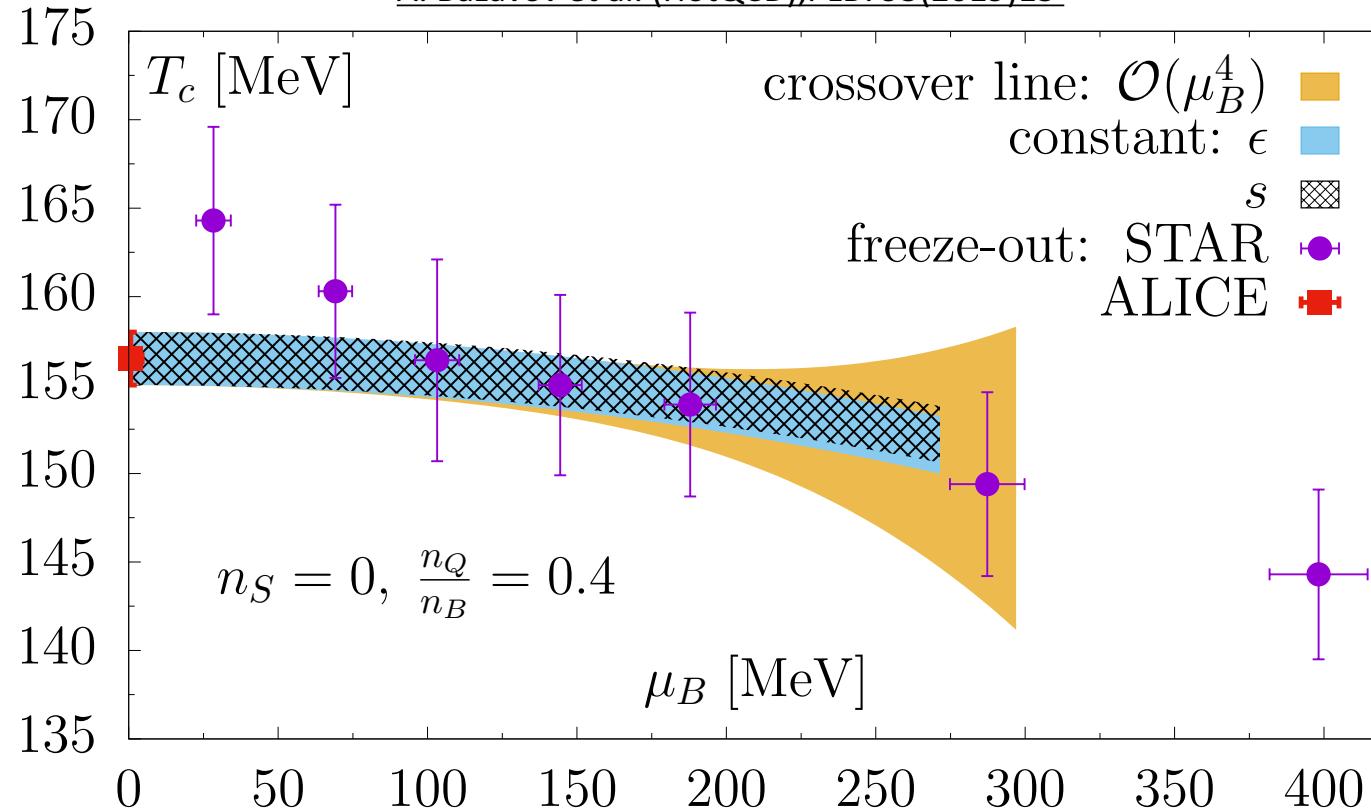
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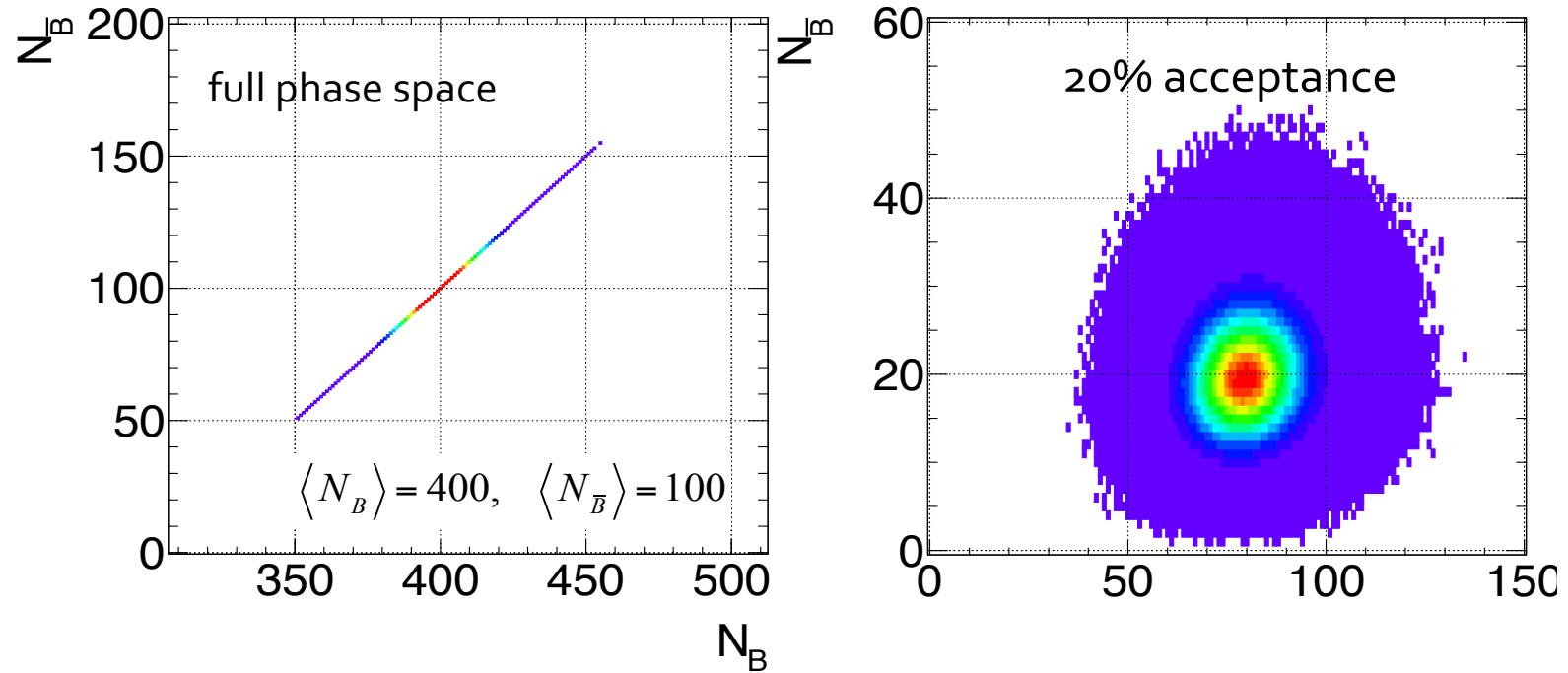
- **Holy grail: see critical behavior in 6th and higher order cumulants**

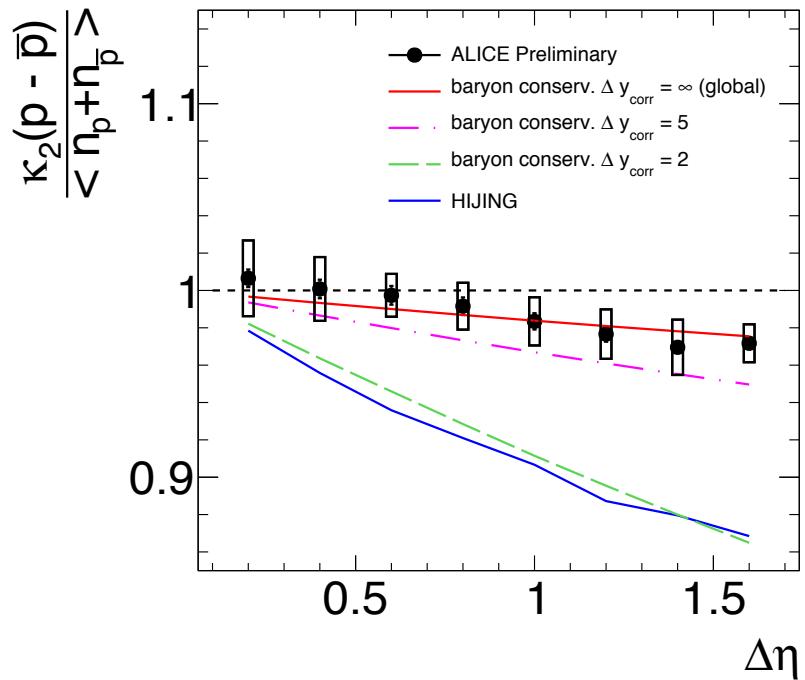
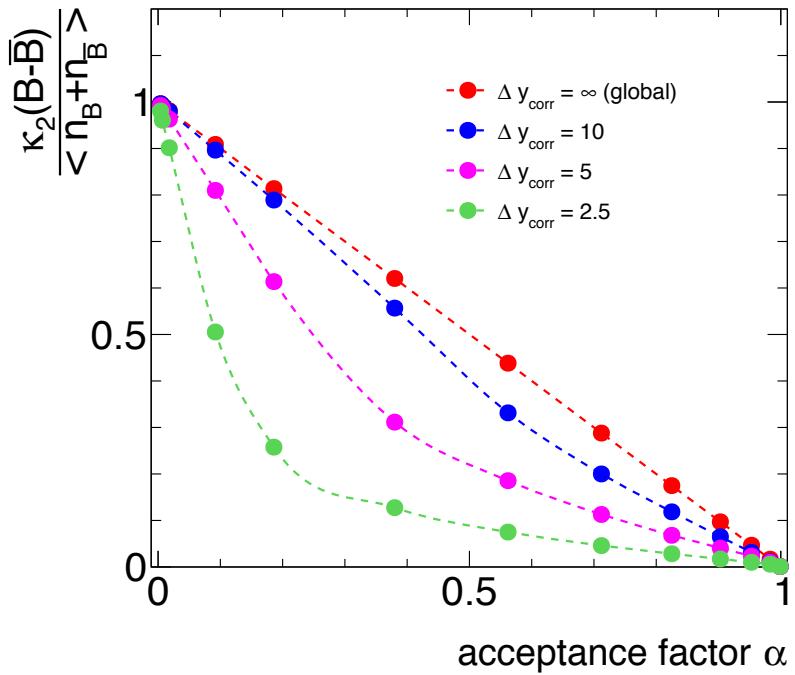


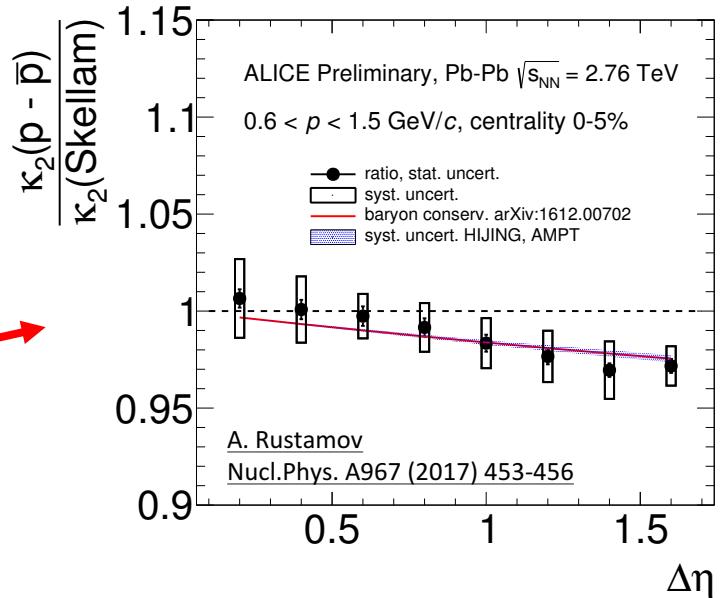
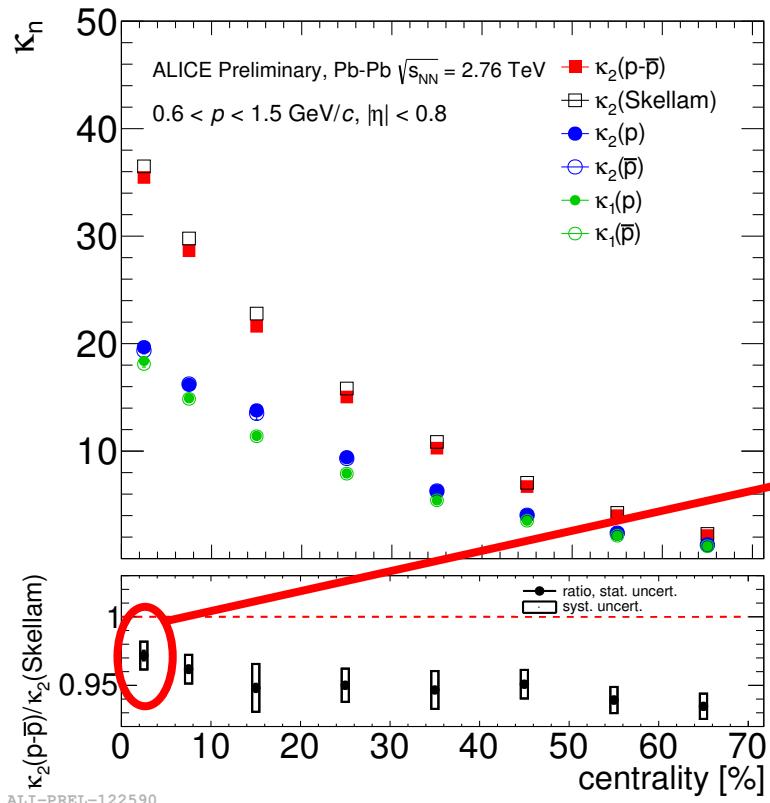
RUN1: 2nd order (~13M min. bias events)
RUN2: 4th order (~150M central events)
RUN3: 6th ... (>1000M central events)

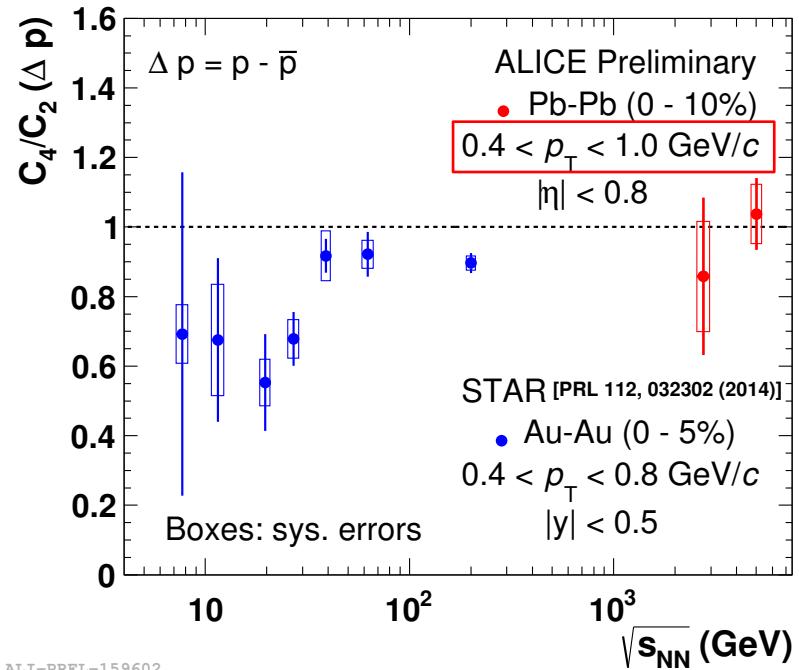
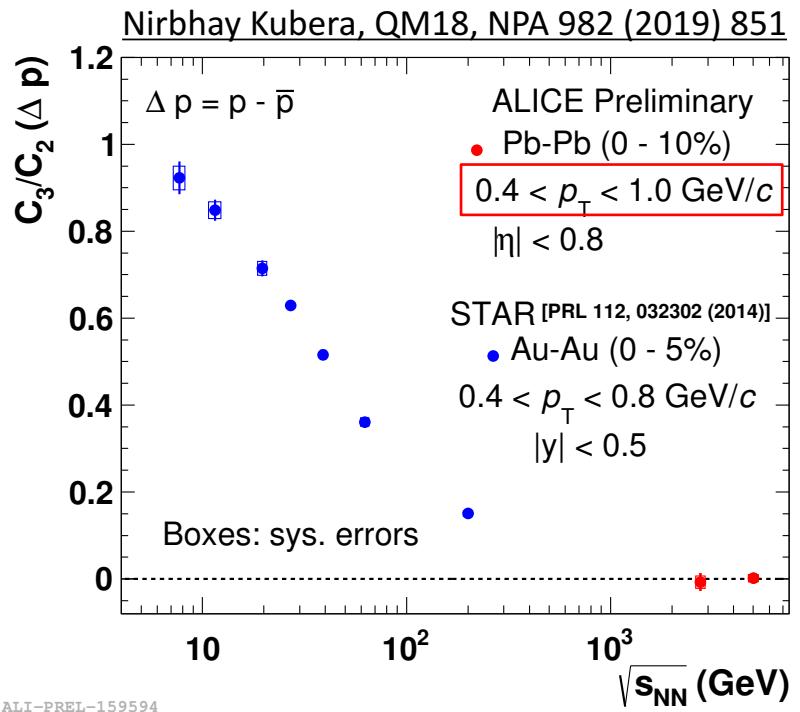
BACKUP

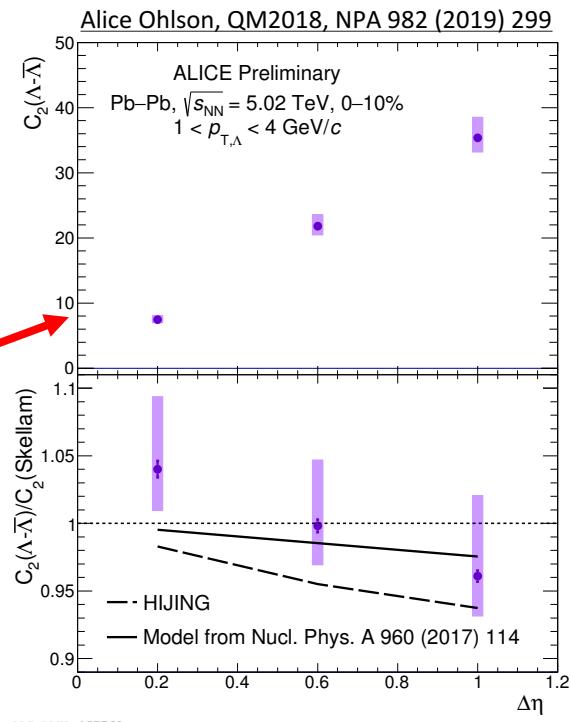
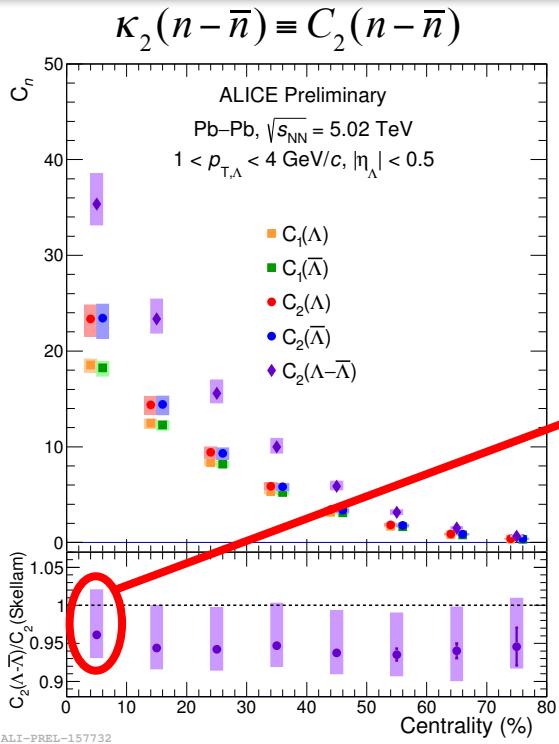








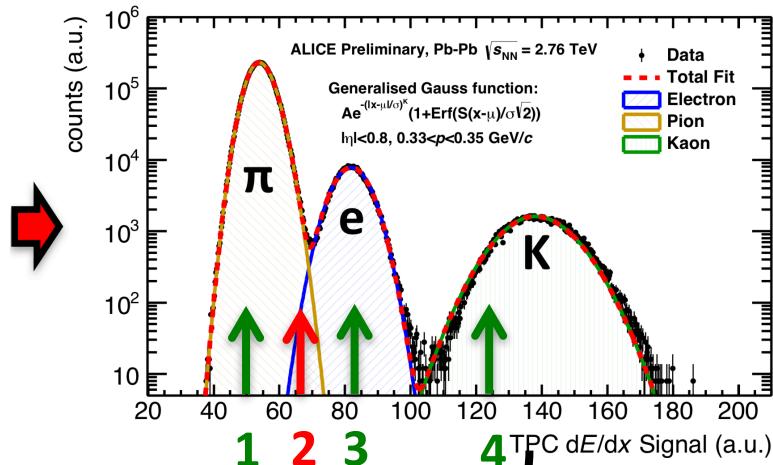
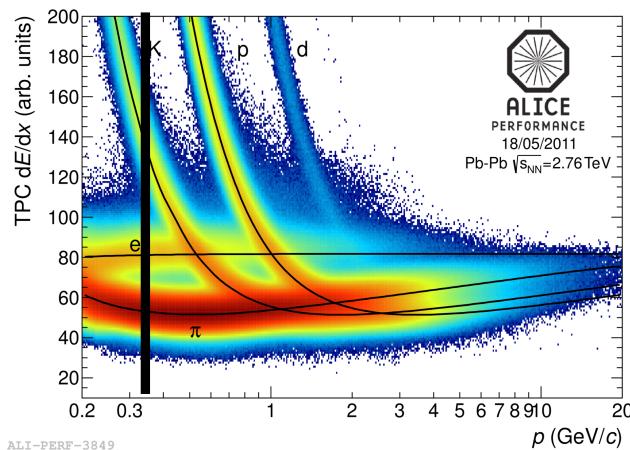




- Similar trend as for **net-p**
- Better precision is needed to disentangle global vs local conservation laws

Cut-based approach: count tracks of a given particle type

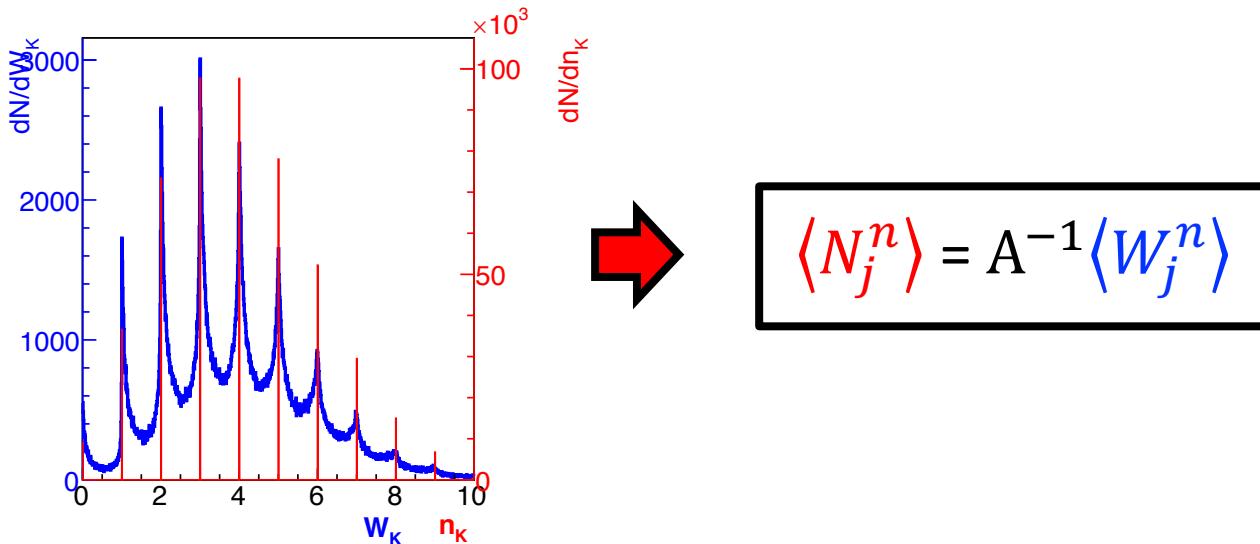
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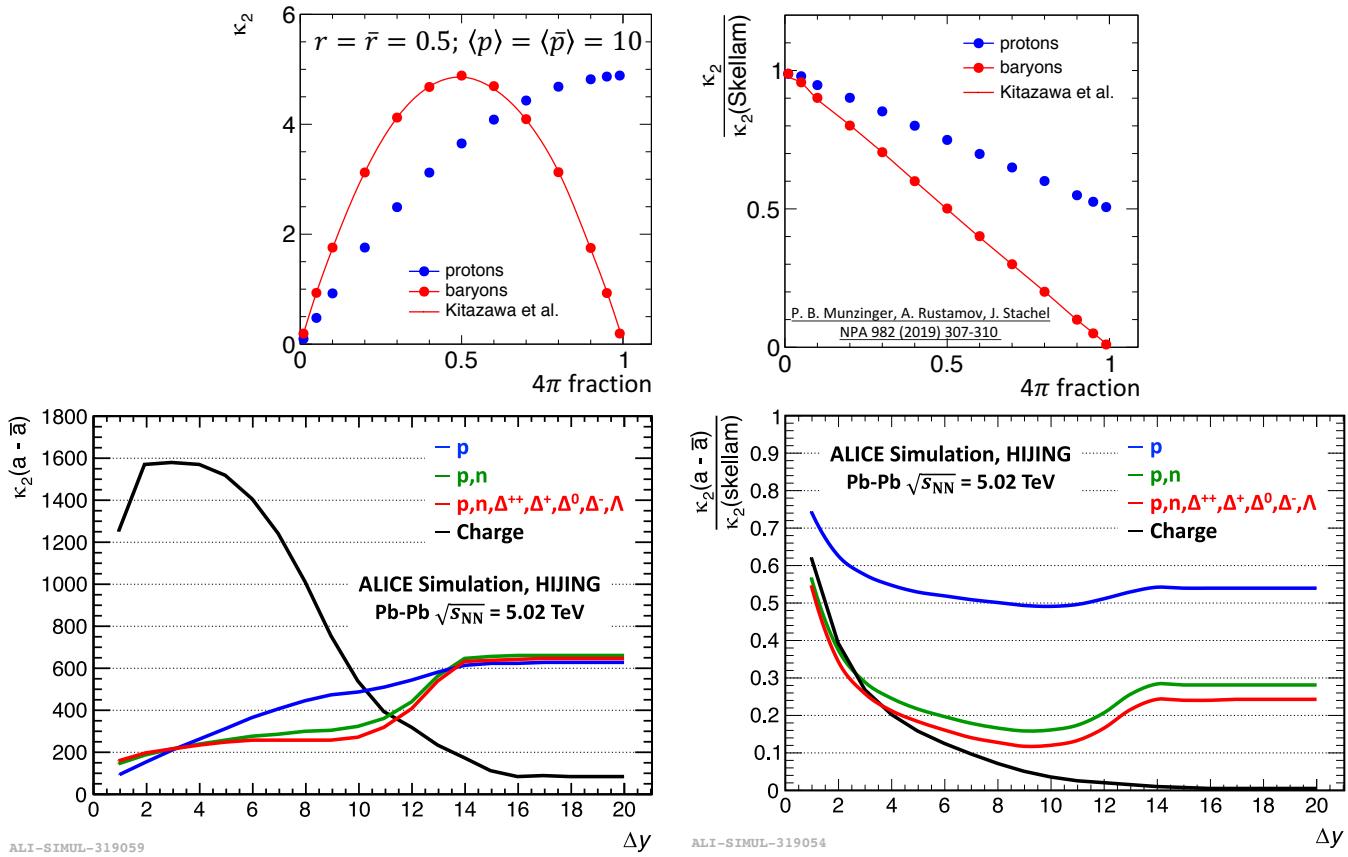
$$\omega_\pi^{(1)} = 1, \quad \omega_\pi^{(2)} \approx 0.6, \quad \omega_\pi^{(3)} = 0, \quad \omega_\pi^{(4)} = 0 \quad \Rightarrow \boxed{W_\pi = 1.6 \neq N_\pi}$$

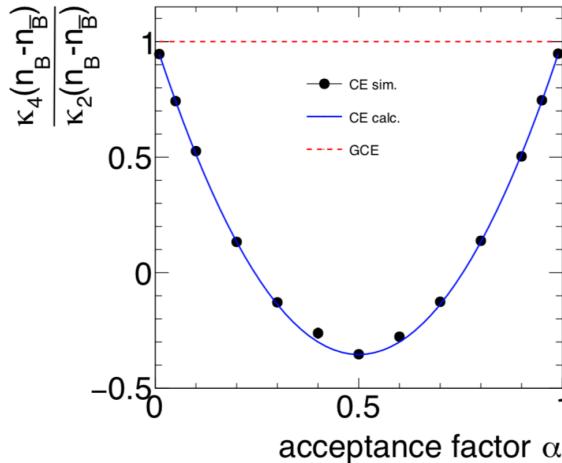
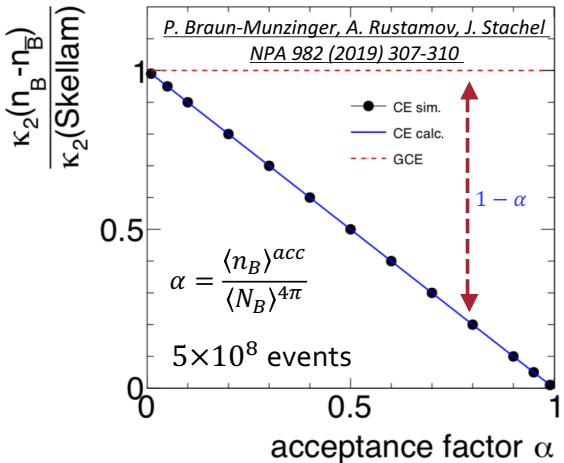
A. Rustamov, M. Gazdzicki, M. I. Gorenstein, PRC 86, 044906 (2012), PRC 84, 024902 (2011)

A. Rustamov, M. Arslanbek, arXiv:1807.06370, NIM in print

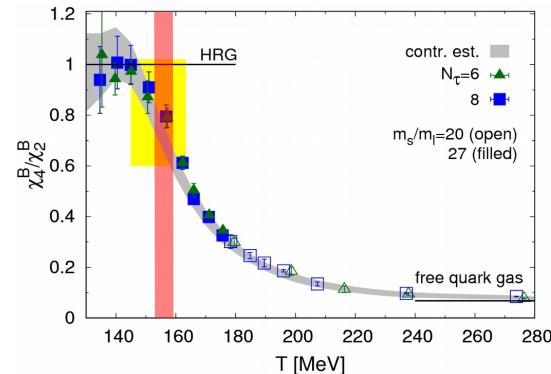


- **Cut-based approach**
 - Uses additional detector information or reject a given phase space bin
 - Challenge: efficiency correction and contamination
- **Identity Method**
 - Gives folded multiplicity distribution
 - Allows for larger efficiencies → smaller correction needed
 - Ideal approach for low momentum ($p < 2$ GeV/c)



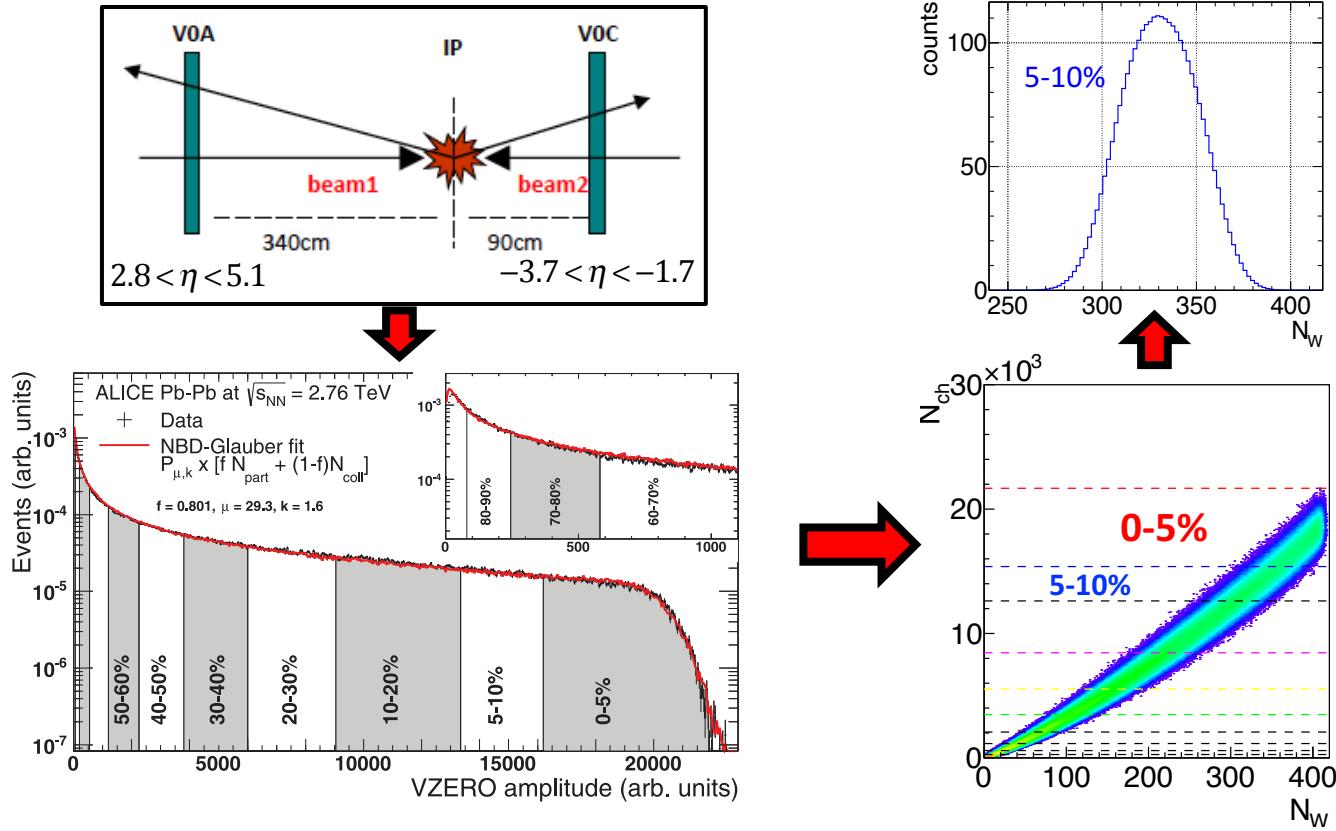


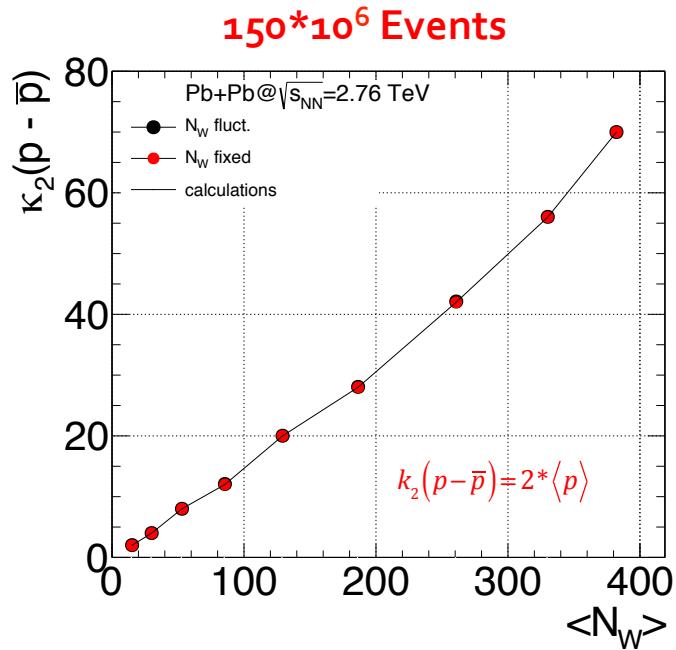
**Deviations from unity
are driven
by different mechanisms**



A. Bzdak, V. Koch, V. Skokov, PRC87 (2013) 014901

K. Redlich and L. Turko, Z. Phys. C5 (1980) 201

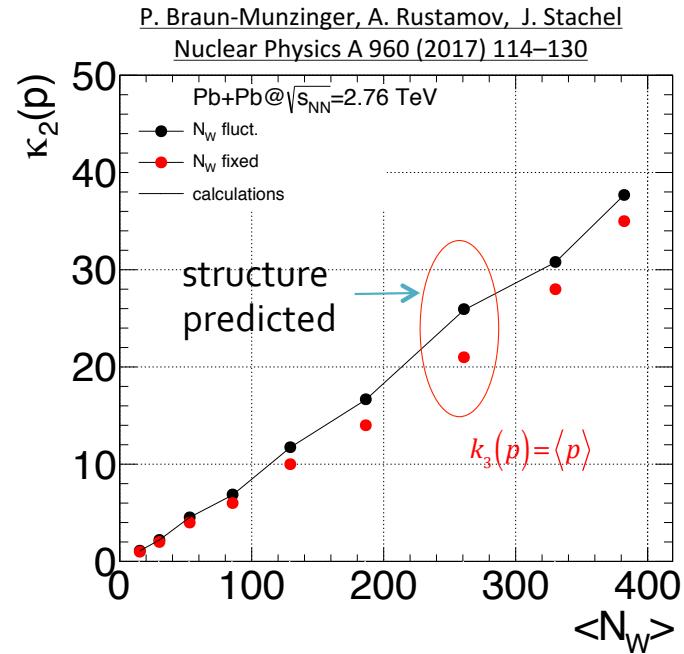




$$k_2(p - \bar{p}) = \langle N_w \rangle k_2(n - \bar{n}) + \langle n - \bar{n} \rangle^2 k_2(N_w)$$

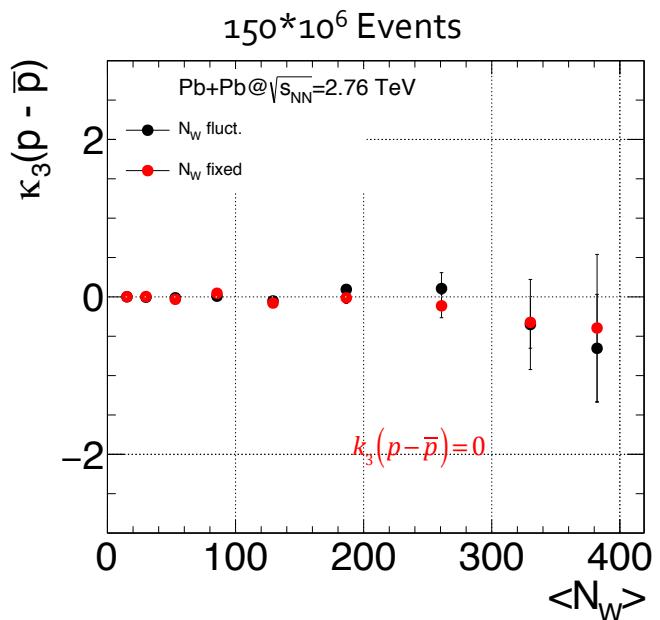
↓
vanishes for ALICE

n, \bar{n} from single wounded nucleon



$$k_2(p) = \langle N_w \rangle k_2(n) + \langle n \rangle^2 k_2(N_w)$$

↓
does not vanish

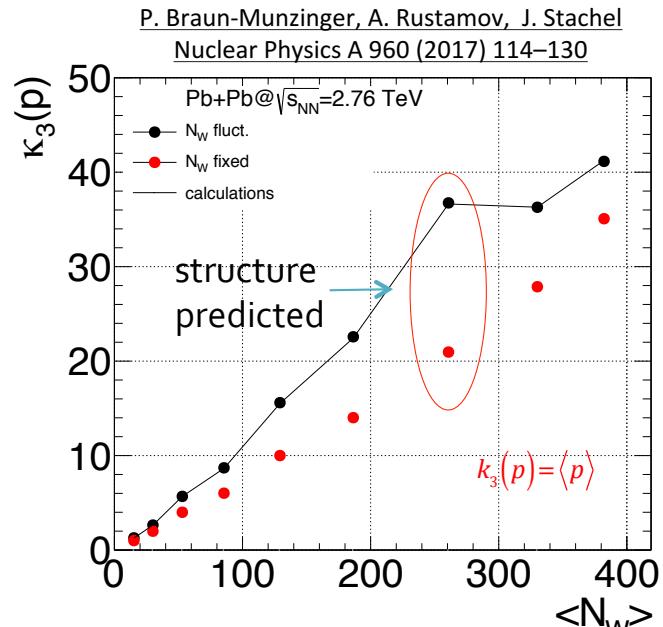


$$k_3(p - \bar{p}) = \langle N_w \rangle k_3(n - \bar{n}) + \langle n - \bar{n} \rangle (\dots)$$

↓

vanishes for ALICE

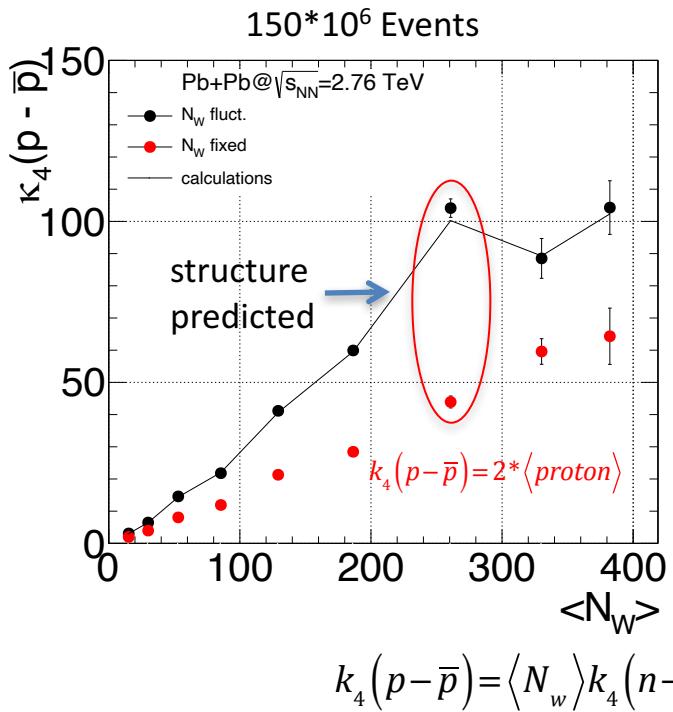
n, \bar{n} from single wounded nucleon



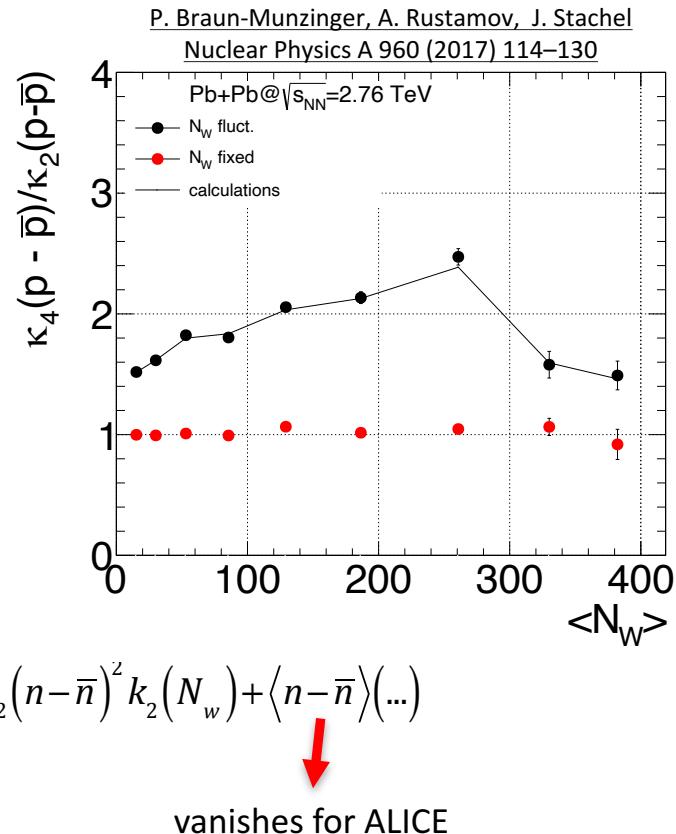
$$k_3(p) = \langle N_w \rangle k_3(n) + \langle n \rangle (\dots)$$

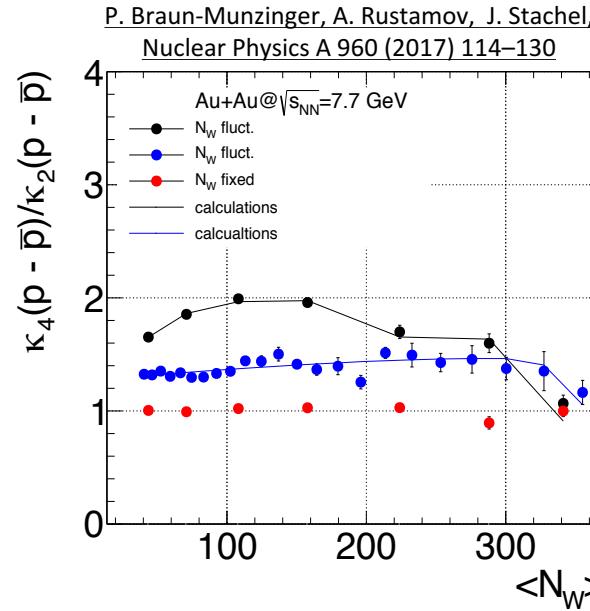
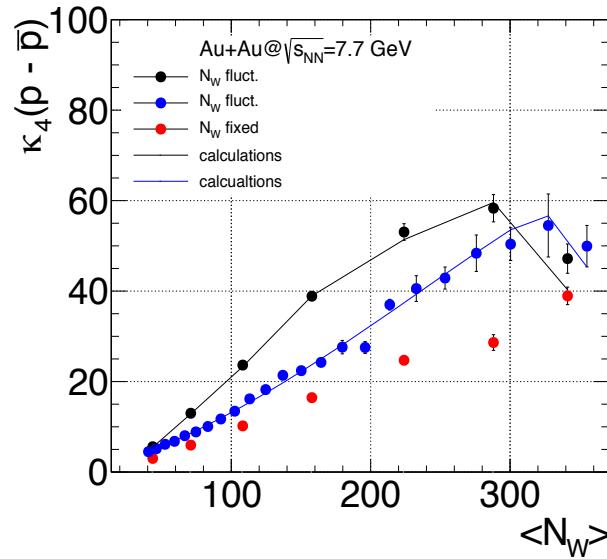
↓

does not vanish



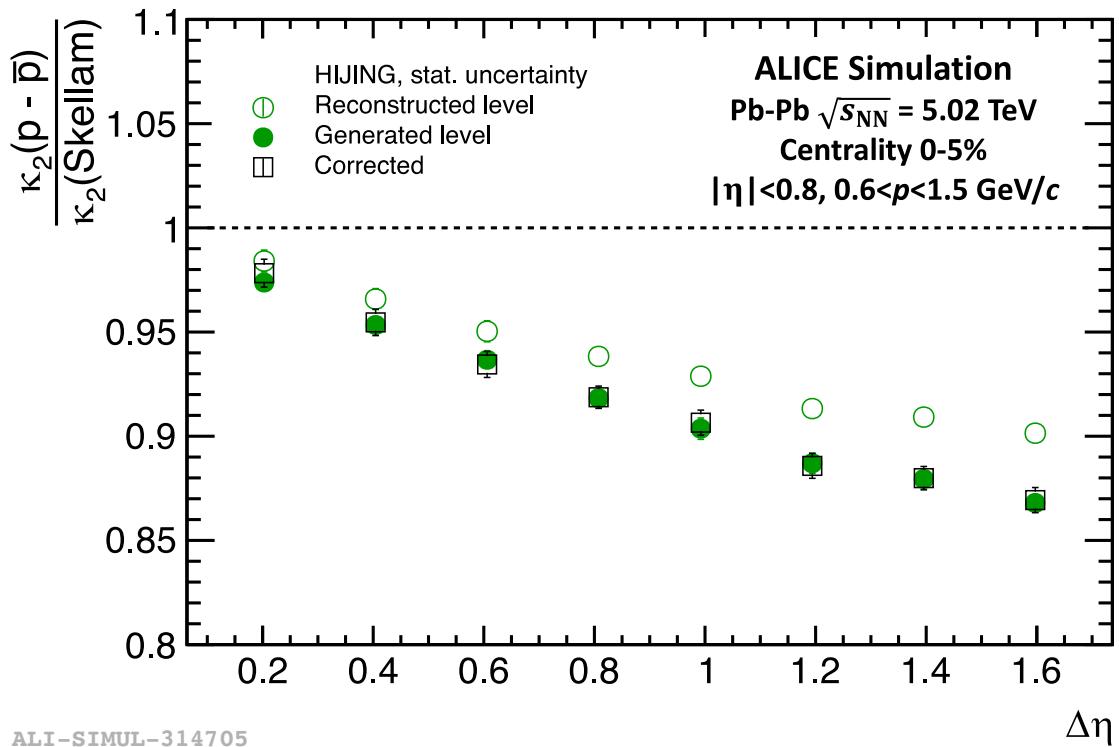
$n, \bar{n} \rightarrow$ from single wounded nucleon





- Subdividing a given centrality bin into smaller ones and then merging them together **incoherently**.
- Incoherent addition of data from intervals with very small centrality bin width will **eliminate true dynamical fluctuations**.

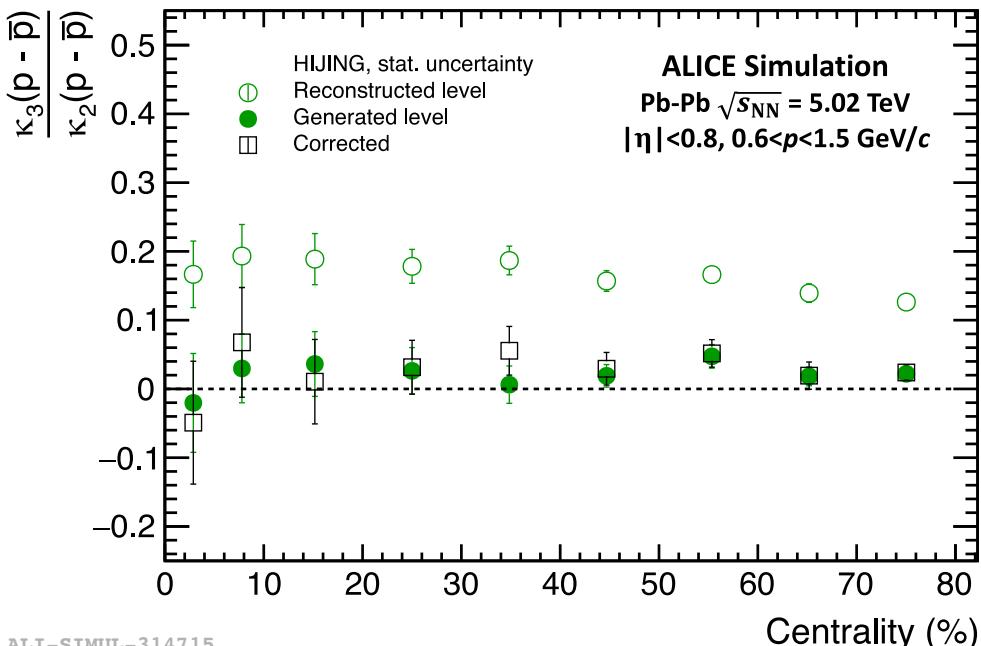
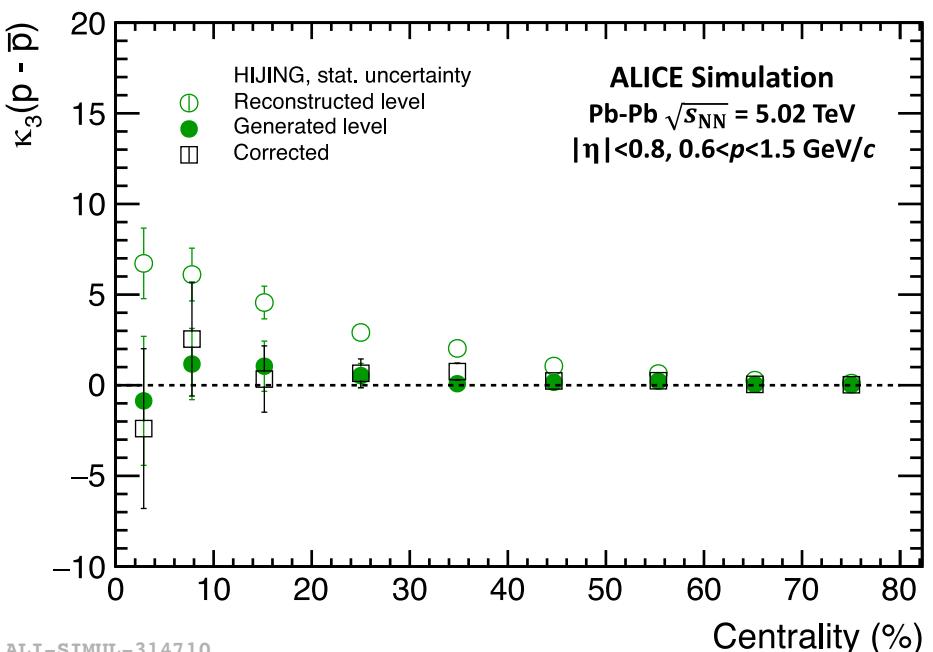
Better publish uncorrected results



Efficiency correction with binomial assumption:

T. Nonaka, M. Kitazawa, S. Esumi, Phys. Rev. C 95, 064912 (2017)

Adam Bzdak, Volker Koch, Phys. Rev. C 86, 044904 (2012)



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