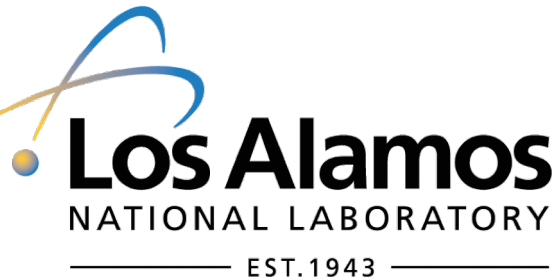


Measurements of the exotic tetraquark candidate $X(3872)$ in pp and pPb

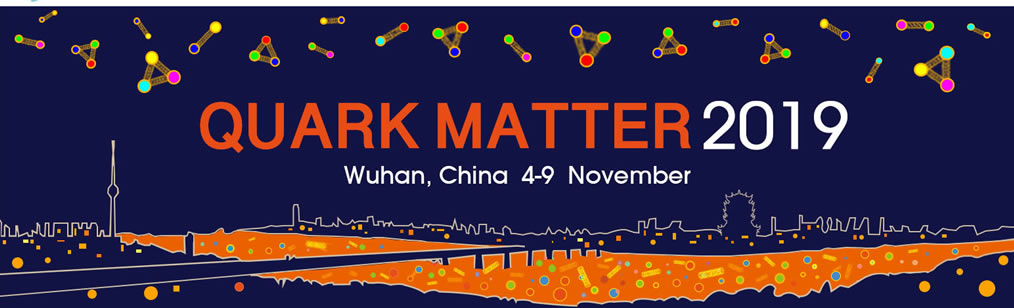
Matt Durham

for the LHCb Collaboration

durham@lanl.gov



THE 28TH INTERNATIONAL CONFERENCE ON ULTRARELATIVISTIC NUCLEUS-NUCLEUS COLLISIONS



Quark Model of Hadrons

Volume 8, number 3

PHYSICS LETTERS

1 February 1964

A SCHEMATIC MODEL OF BARYONS AND MESONS *

M. GELL-MANN

California Institute of Technology, Pasadena, California

Received 4 January 1964

A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon b if we assign to the triplet t the following properties: spin $\frac{1}{2}$, $z = -\frac{1}{3}$, and baryon number $\frac{1}{3}$. We then refer to the members $u^{\frac{2}{3}}$, $d^{-\frac{1}{3}}$, and $s^{-\frac{1}{3}}$ of the triplet as "quarks" q and the members of the anti-triplet as anti-quarks \bar{q} . Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(qq\bar{q})$, etc. It is assuming that the lowest baryon configuration (qqq) gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just **1** and **8**.

AN SU_3 MODEL FOR STRONG INTERACTION SYMMETRY AND ITS BREAKING

G. Zweig *)

CERN - Geneva

8182/TH.401

17 January 1964

In general, we would expect that baryons are built not only from the product of three aces, AAA , but also from $\bar{A}AAAA$, $\bar{A}AAAAAA$, etc., where \bar{A} denotes an anti-ace. Similarly, mesons could be formed from $\bar{A}A$, $\bar{A}AAA$ etc. For the low mass mesons and baryons we will assume the simplest possibilities, $\bar{A}A$ and AAA , that is, "deuces and treys".

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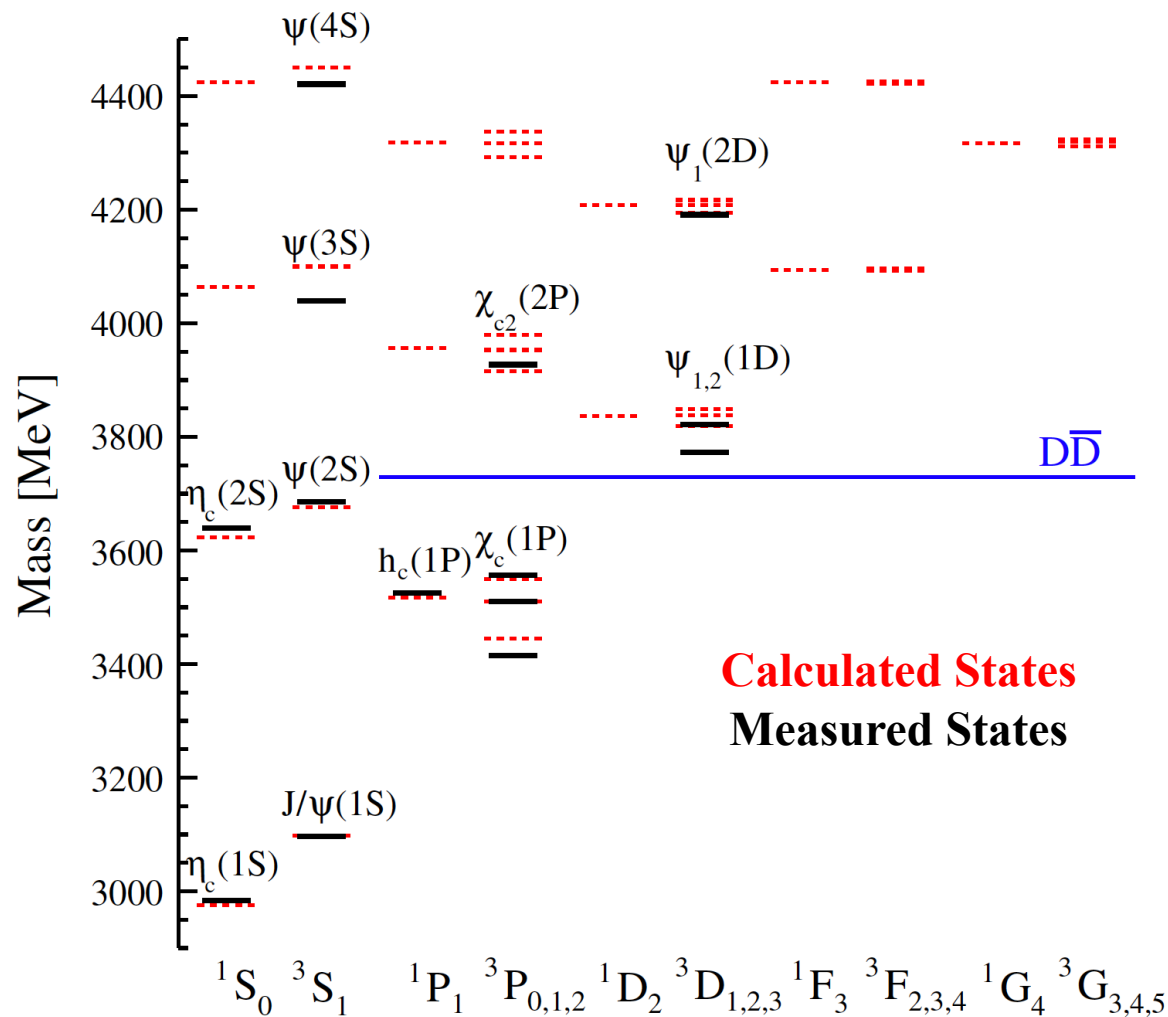
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States with >3 quarks have been expected since the beginning of the quark model

Conventional $c\bar{c}$ States



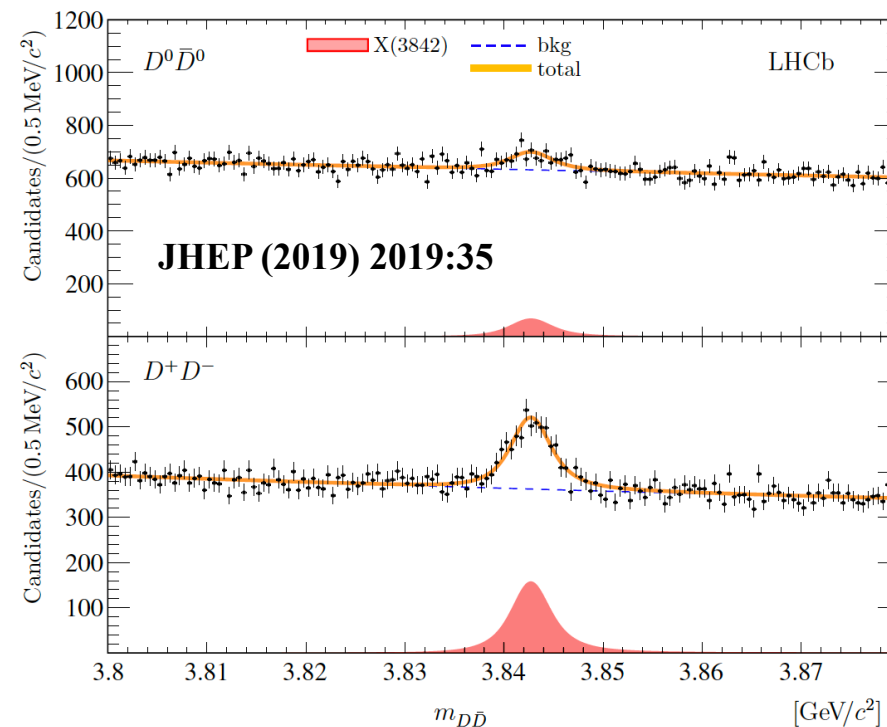
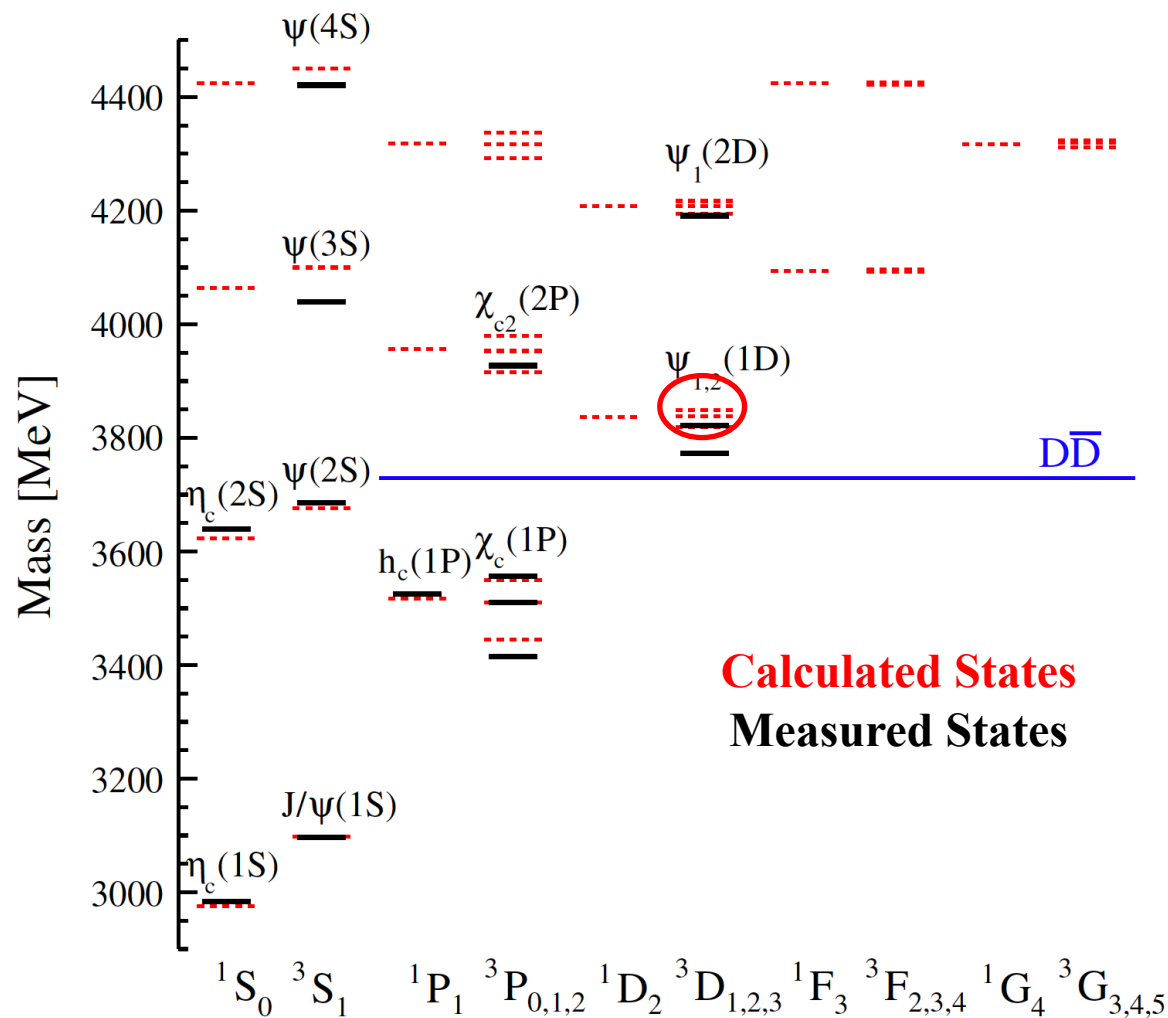
Nonrelativistic potential model: solve Schrodinger equation with the potential

$$V_0^{(c\bar{c})}(r) = -\frac{4}{3} \frac{\alpha_s}{r} + br + \frac{32\pi\alpha_s}{9m_c^2} \tilde{\delta}_\sigma(r) \vec{S}_c \cdot \vec{S}_{\bar{c}}$$

Barnes, Godfrey, Swanson,
Phys. Rev. D 72, 054026 (2005)

Rev. Mod. Phys. 90, 015003 (2018)

Conventional $c\bar{c}$ States



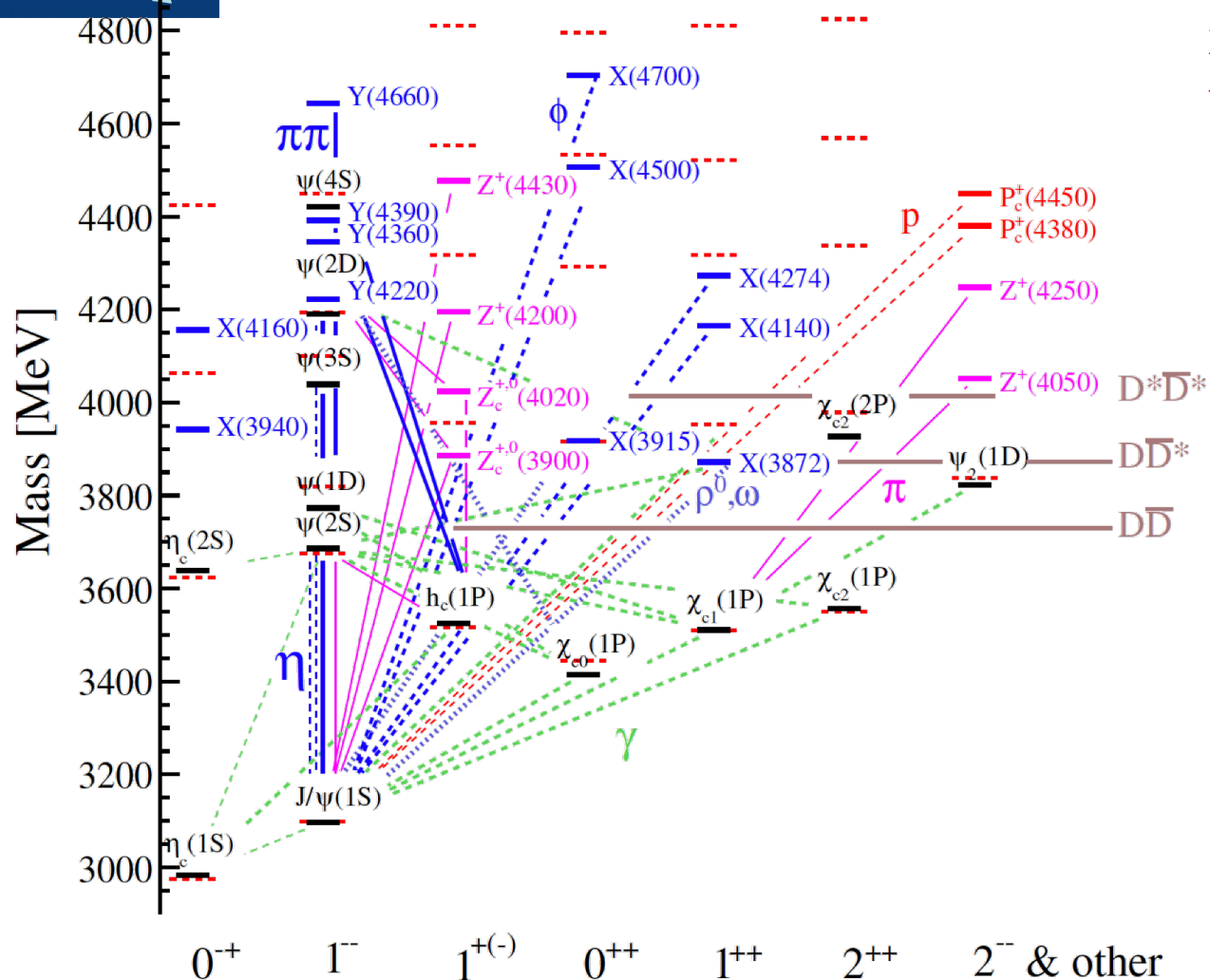
New charmonium states still being found: LHCb observed state consistent with $\psi_3(1^3D_3)$ found in $D\bar{D}$ and D^+D^- mass spectra in 2019

Crucial to account for conventional states when searching for exotics

Rev. Mod. Phys. 90, 015003 (2018)

Exotic $c\bar{c}$ States

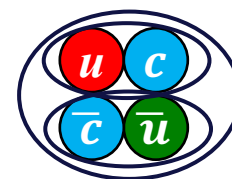
20+ states containing $c\bar{c}$ have been discovered since 2003
that do not fit in the picture of typical charmonium:
Collectively known as “XYZ” particles



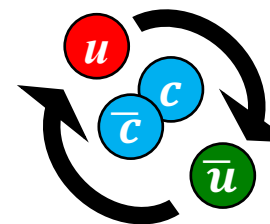
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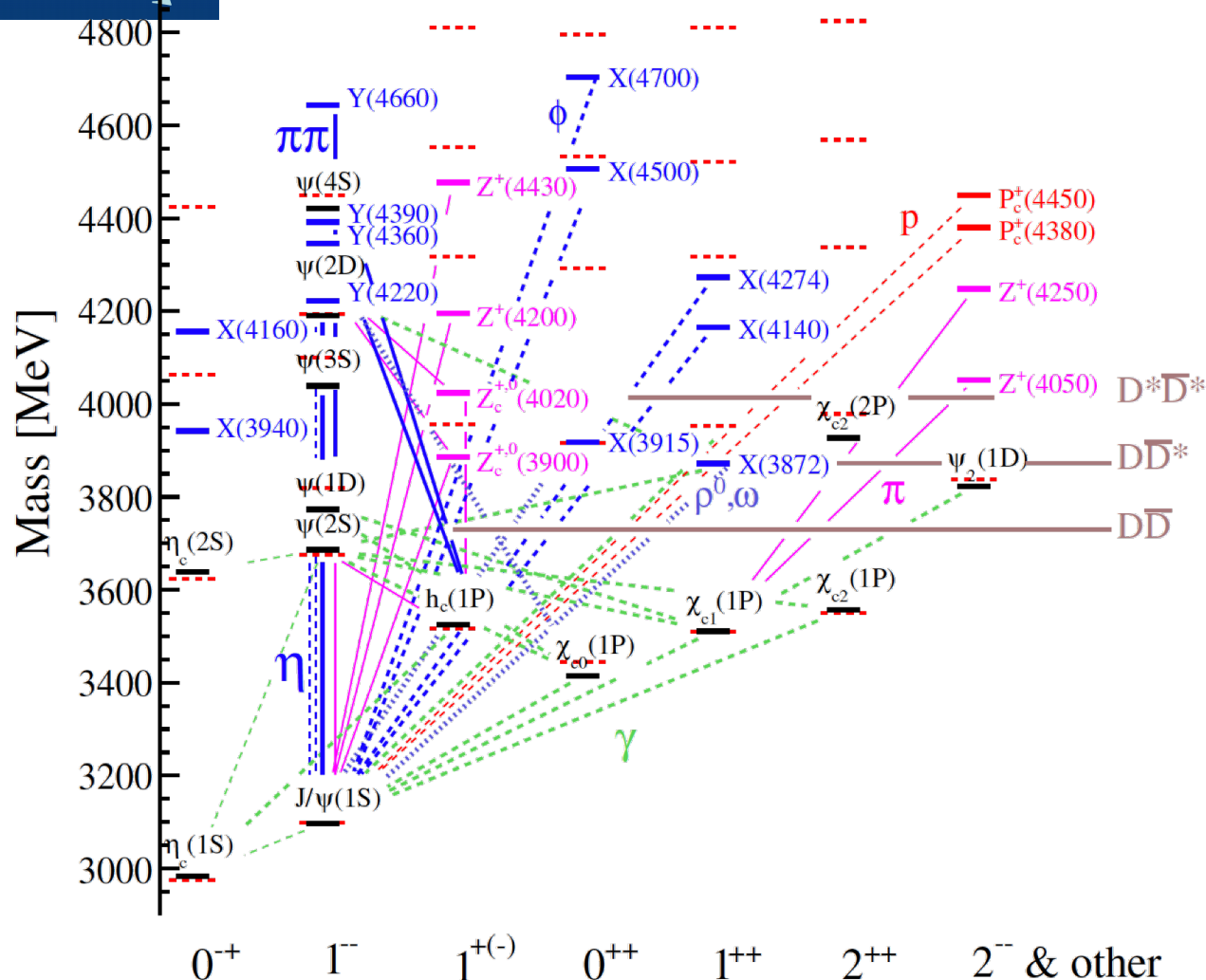
Compact tetraquark/pentaquark



Diquark-diquark
PRD 71, 014028 (2005)
PLB 662 424 (2008)



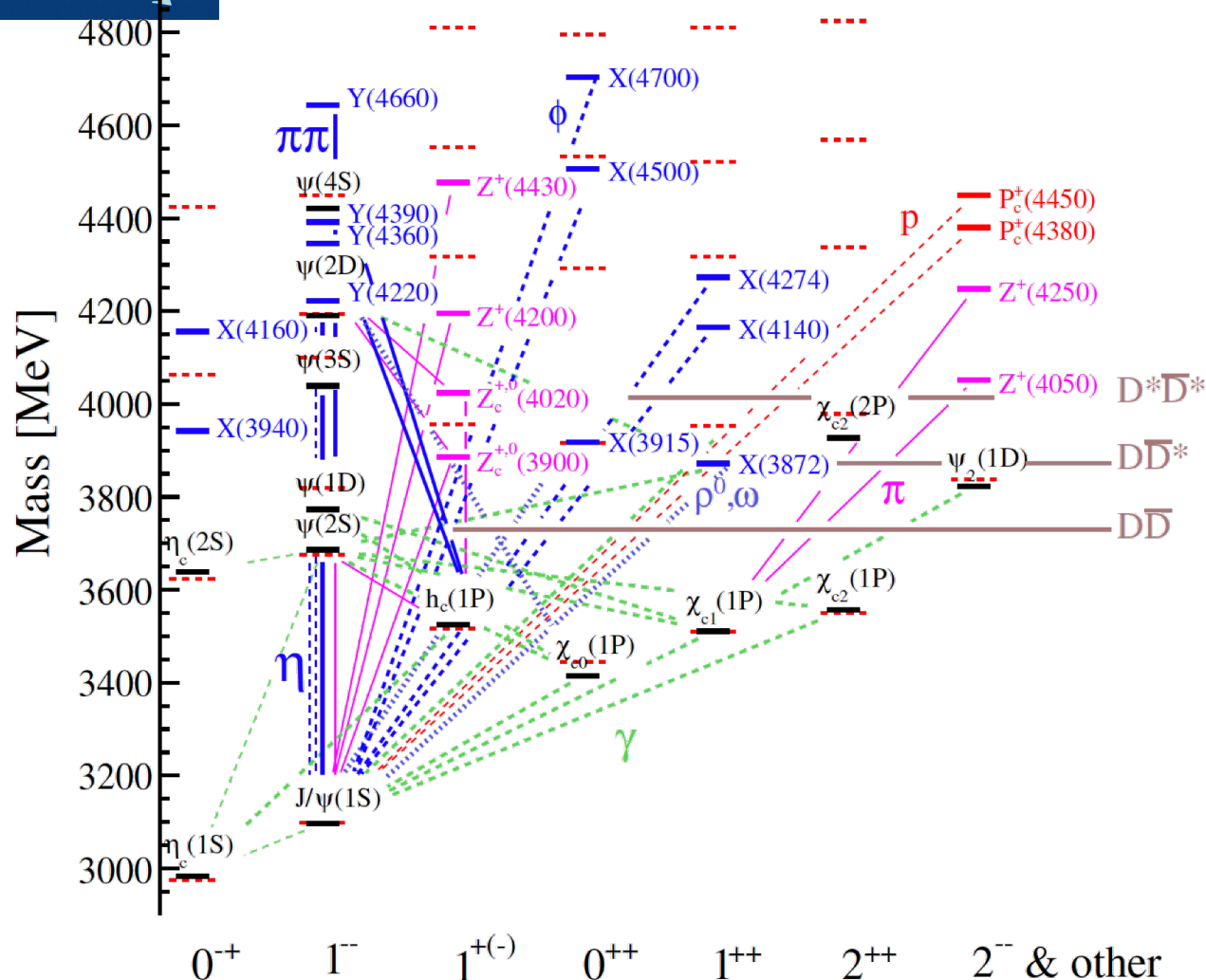
**Hadrocharmonium/
adjoint charmonium**
PLB 666 344 (2008)
PLB 671 82 (2009)



Rev. Mod. Phys. 90, 015003 (2018)

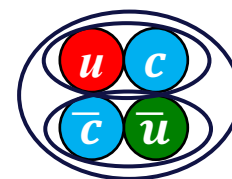
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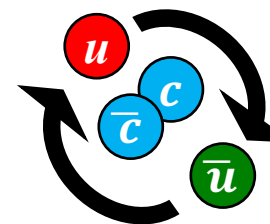


Multiple explanations explored in literature:

Compact
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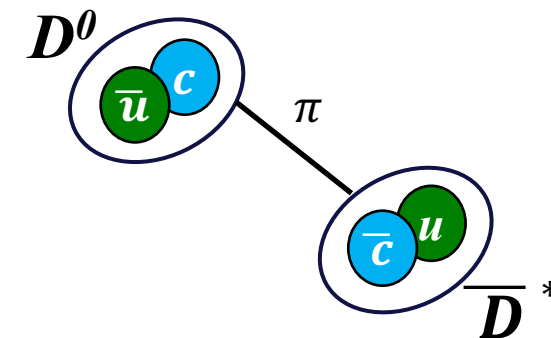
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Hadronic Molecules

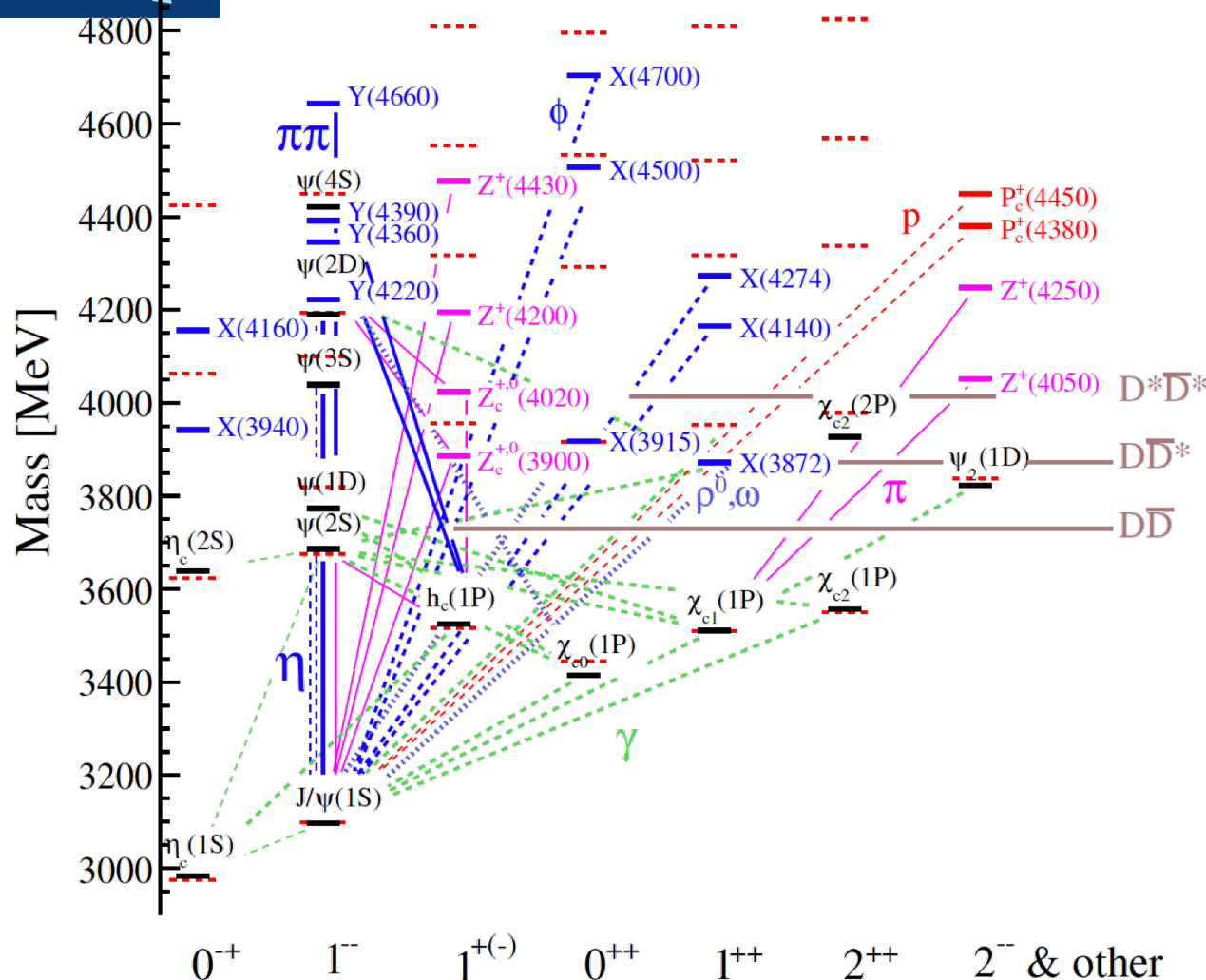
PLB 590 209 (2004)
PRD 77 014029 (2008)
PRD 100 0115029(R) (2019)



Rev. Mod. Phys. 90, 015003 (2018)

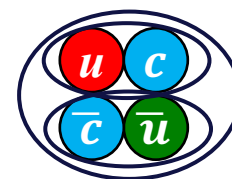
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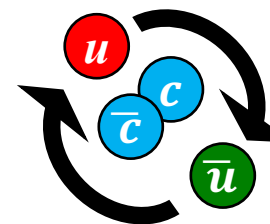


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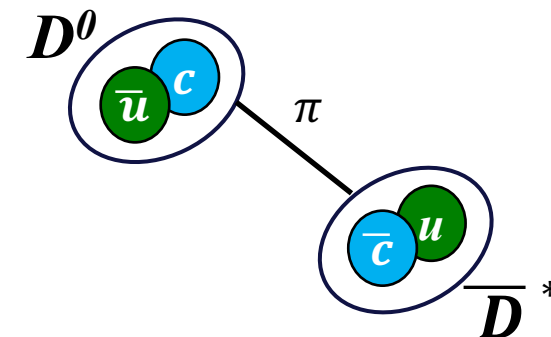
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PRD 100 0115029(R) (2019)



Mixtures of exotic + conventional states

$$X = a |c\bar{c}\rangle + b |c\bar{c}q\bar{q}\rangle$$

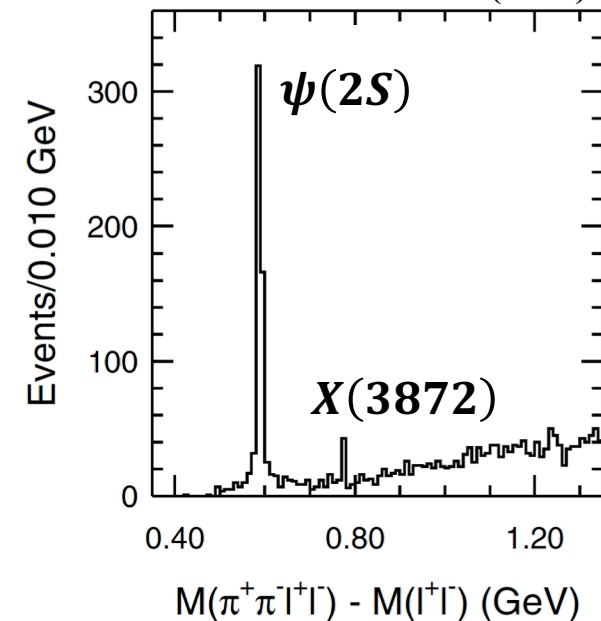
PLB 578 365 (2004)
PRD 96 074014 (2017)

Rev. Mod. Phys. 90, 015003 (2018)

X(3872) - a puzzle

Recently renamed
 $\chi_{c1}(3872)$ by PDG

Belle Collaboration
PRL 91 262001 (2003)

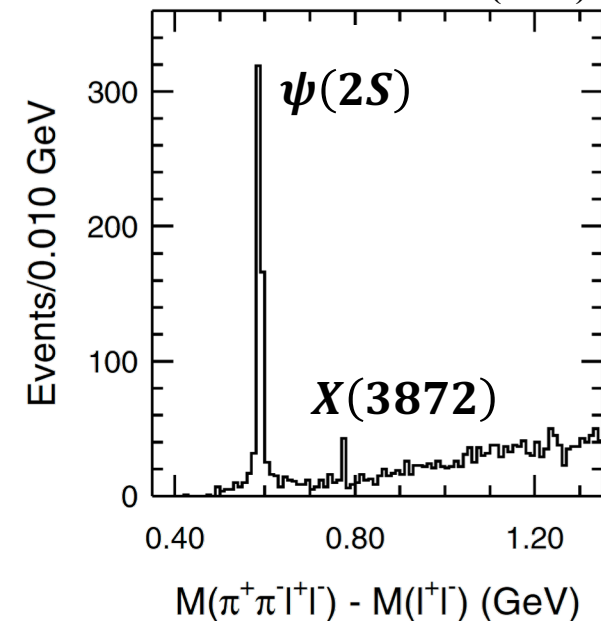


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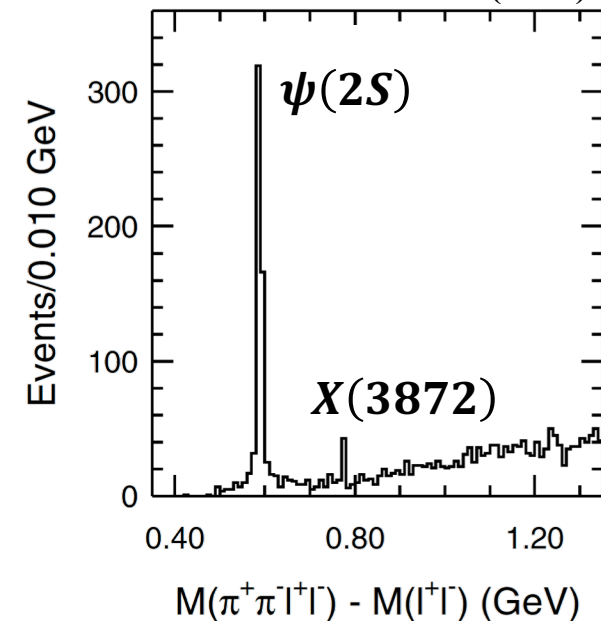
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$$M_{\chi_{c1}(3872)} - (M_{D^0} + M_{\bar{D}^{*0}}) = 0.01 \pm 0.27 \text{ MeV}$$

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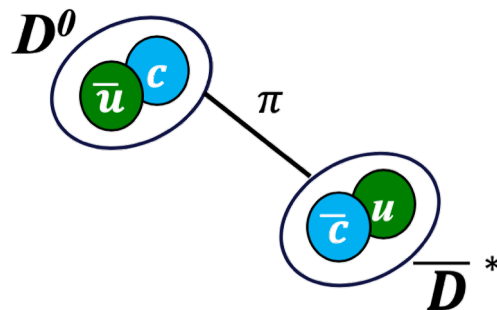
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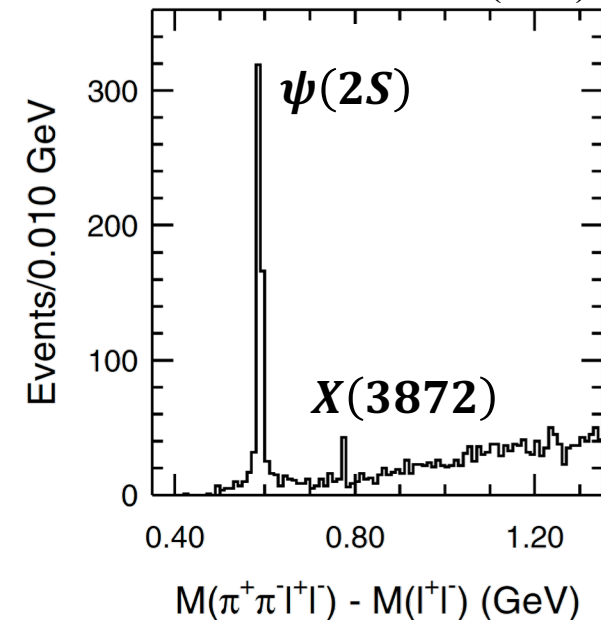


VERY small binding energy
VERY large radius, ~ 7 fm

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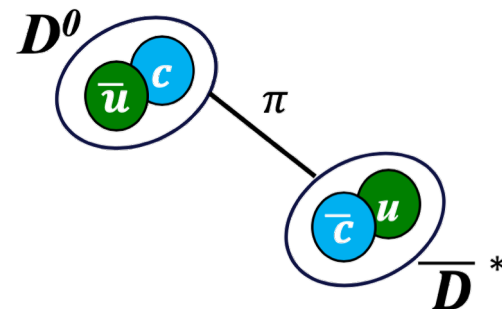
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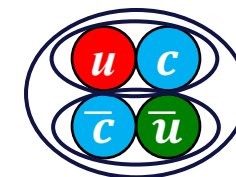
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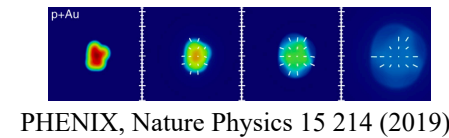
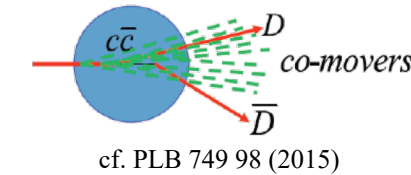
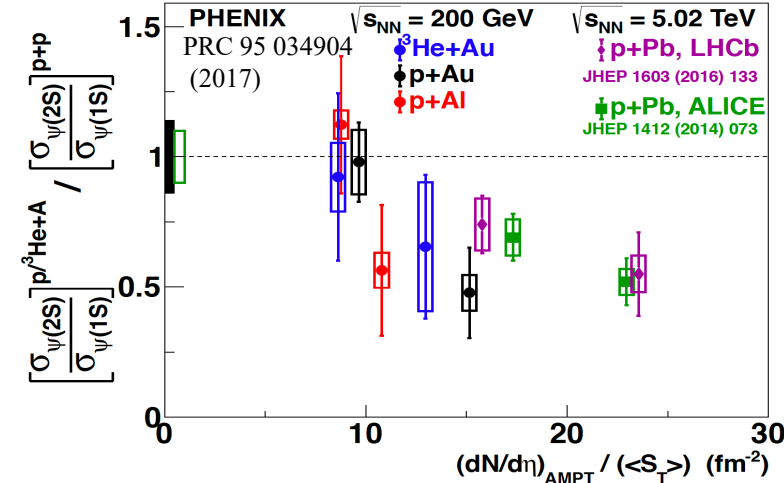
Compact tetraquark



Tightly bound via color exchange between diquarks
Small radius, ~1 fm

Effects of Binding Energy

- Suppression of weakly-bound quarkonia states has been studied for decades in pA collisions
 - Ratios of $\psi(2S)/J/\psi$ and $\Upsilon(2S,3S)/\Upsilon(1S)$

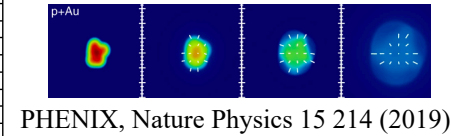
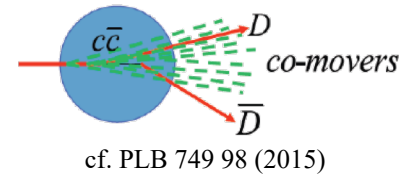
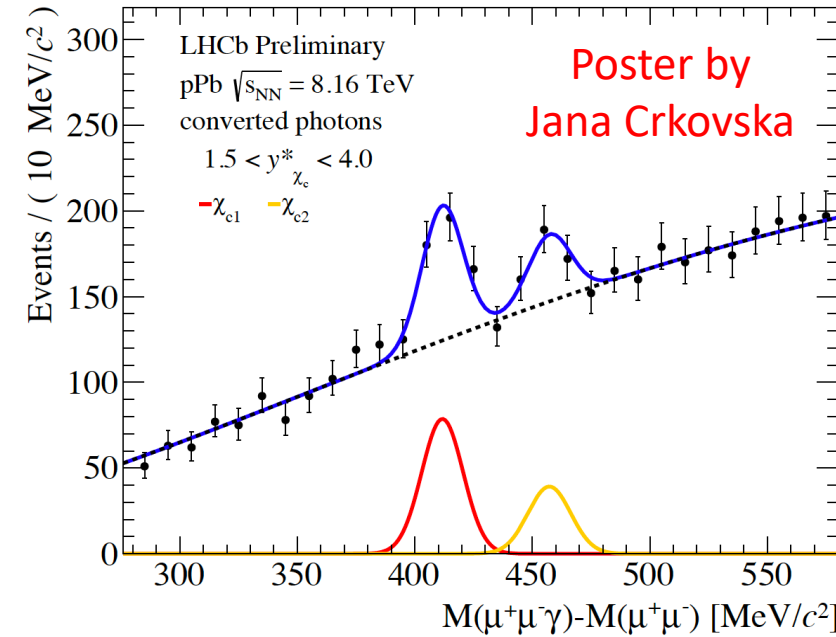


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mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69
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Satz, J. Phys. G 32 (3) 2006

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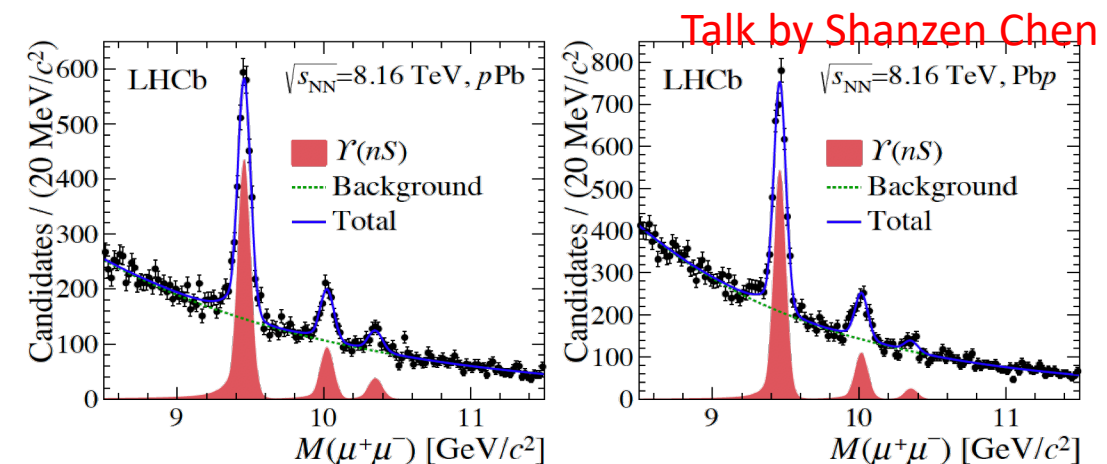
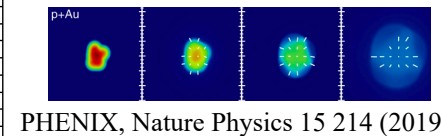
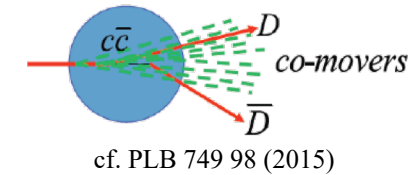
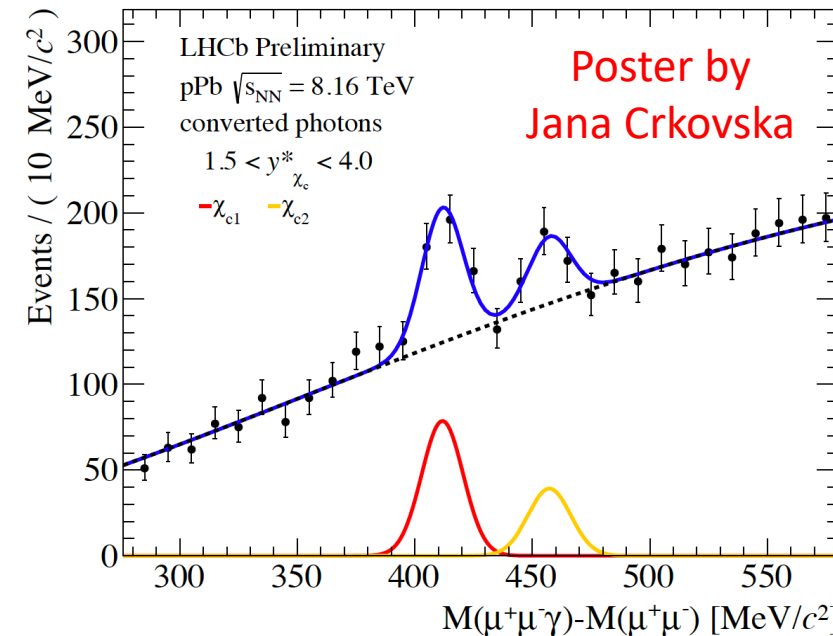
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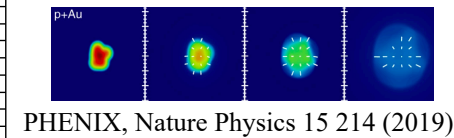
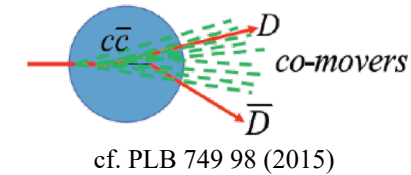
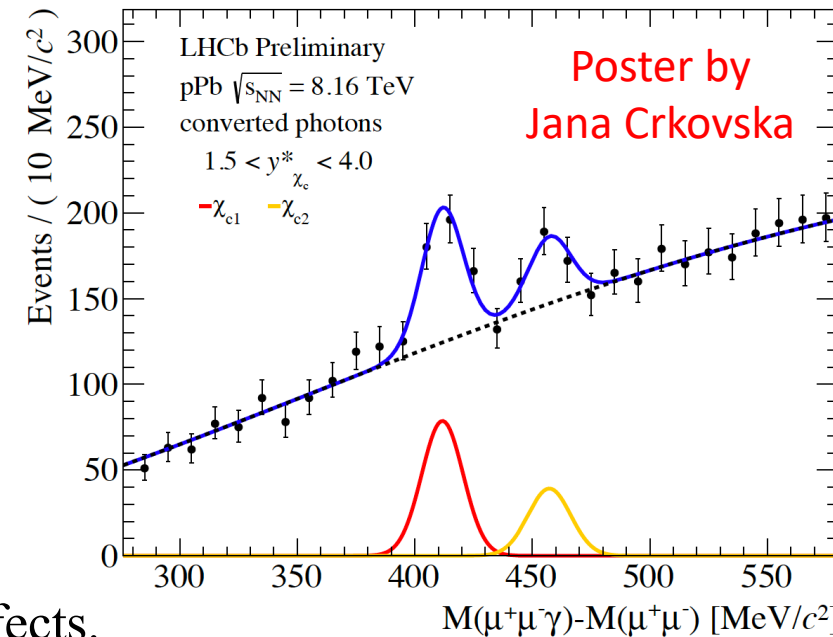
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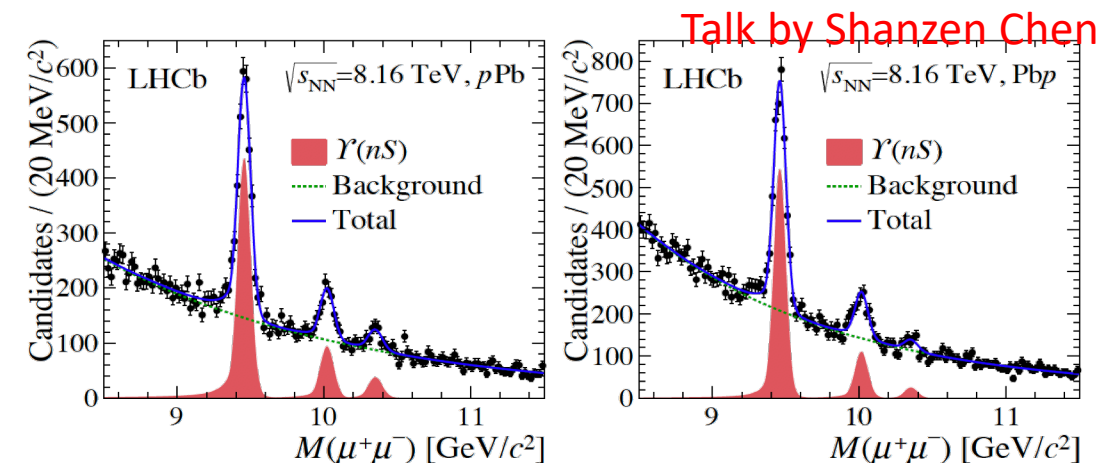


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- Prevalent in regions with high particle multiplicity
- Weakly bound hadronic molecules may show similar effects.



state	η_c	J/ψ	χ_{c0}	χ_{c1}	χ_{c2}	ψ'	$D\bar{D}^*$ Molecule
mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69	X(3872)
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Satz, J. Phys. G 32 (3) 2006

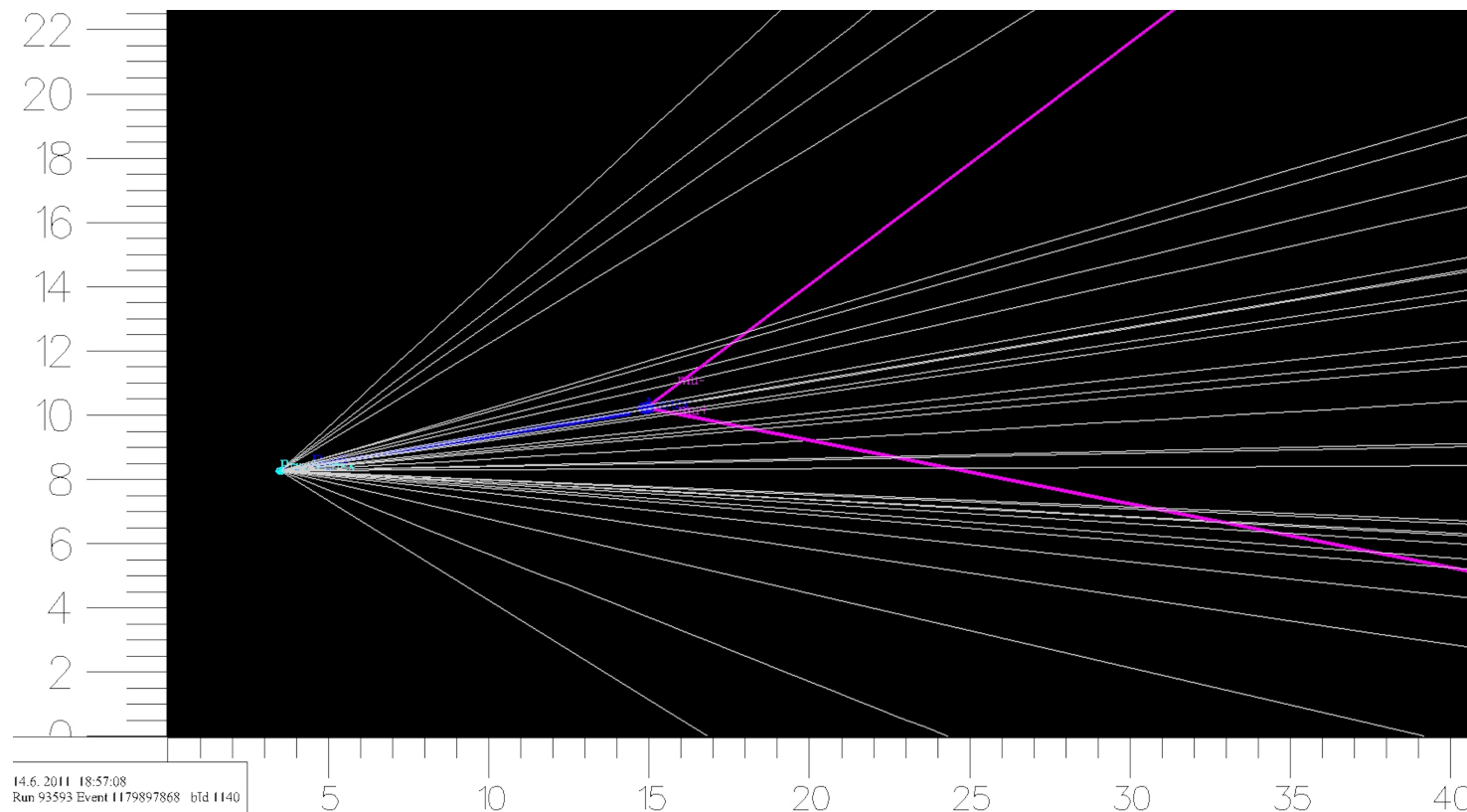


Probing X(3872) structure via interactions with the underlying event

Prompt production:

- X(3872) produced at collision vertex can be subject to further interactions with co-moving particles (medium?) produced in the event
- Potentially subject to breakup effects

Event display of $B_s^0 \rightarrow \mu^+ \mu^-$ candidate, PRL 118 191801 (2017)



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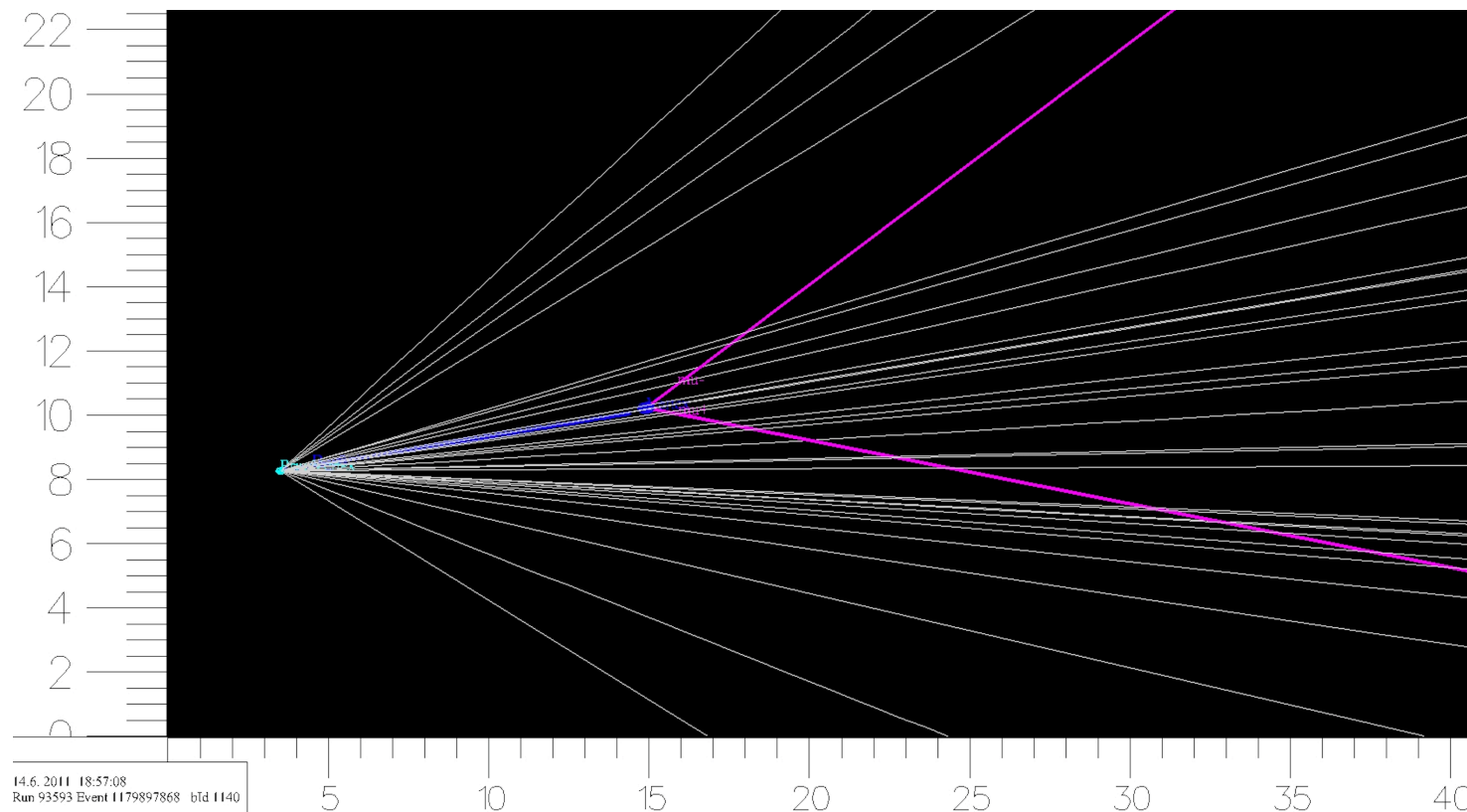
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Production in b -decays:

- Hadrons containing b travel down the beampipe and decay away from the primary vertex and decay in vacuum
- X(3872) from decays not subject to further interactions
- Control sample

Event display of $B_s^0 \rightarrow \mu^+ \mu^-$ candidate, PRL 118 191801 (2017)



$$X(3872) \rightarrow J/\psi \pi^+ \pi^-$$

Vertex detector (VELO):

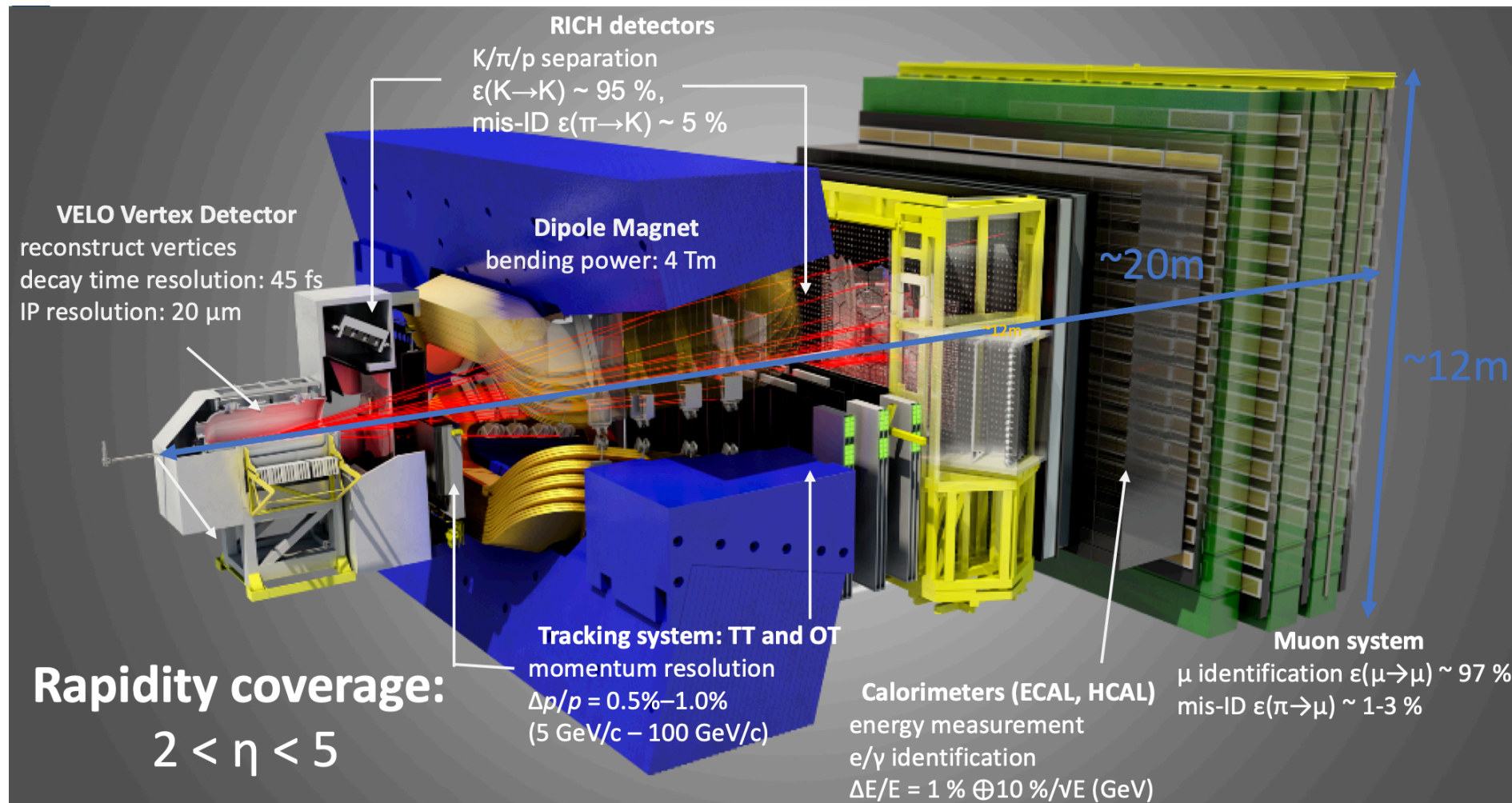
- Separation of prompt and b -decay production
- Number of VELO tracks gives measure of event activity

Two RICH detectors:

- Pion identification

Muon System:

- Layers of absorber/tracking
- Muon hardware trigger



X(3872) selection

LHCb-CONF-2019-005

Reconstruct the $\mu^+ \mu^- \pi^+ \pi^-$ final state from the decays:

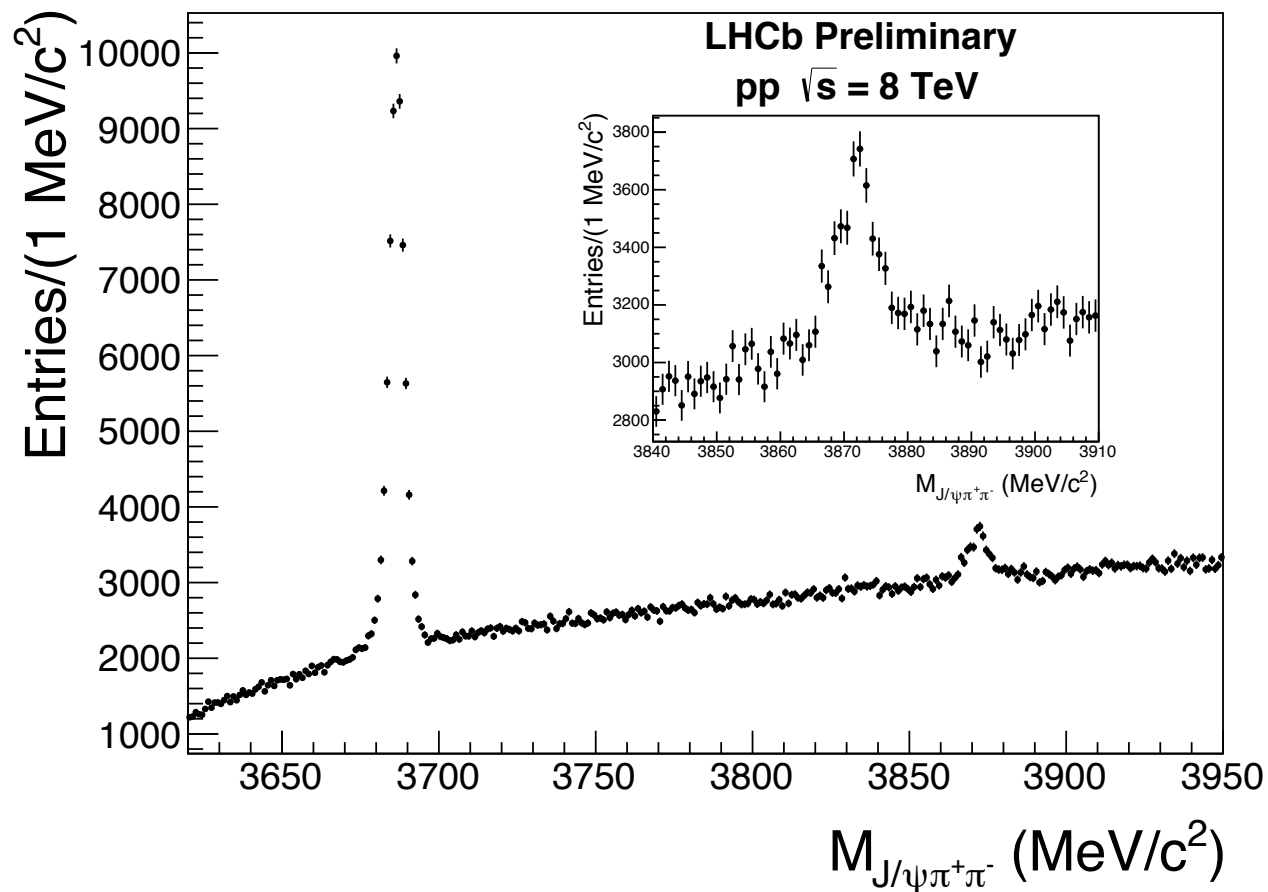
$$X(3872) \rightarrow J/\psi(\rightarrow \mu^+ \mu^-) \rho(\rightarrow \pi^+ \pi^-)$$

$$\psi(2S) \rightarrow J/\psi(\rightarrow \mu^+ \mu^-) \pi^+ \pi^-$$

Select J/ψ from dimuons, combine with two identified pions. Perform kinematic refit, constraining J/ψ mass to known value and all four tracks to identical vertex.

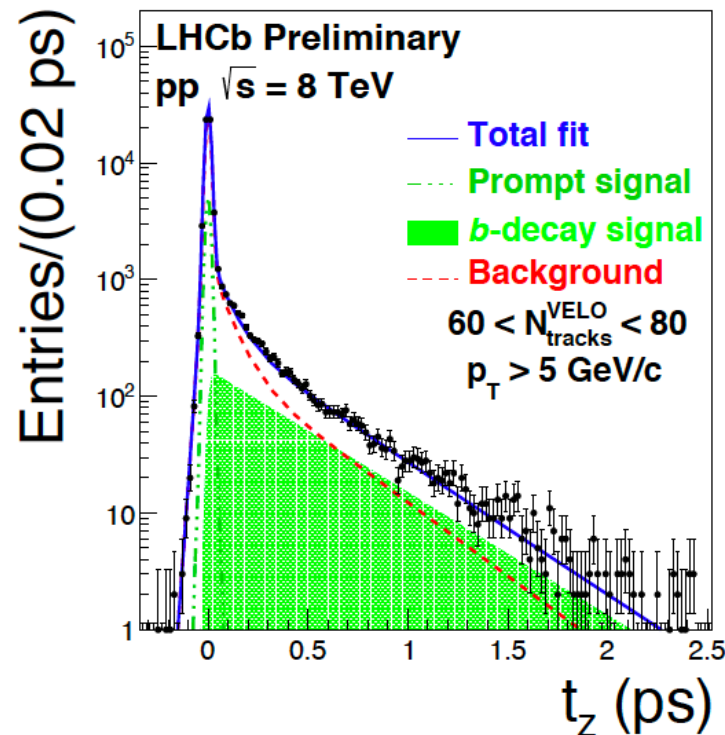
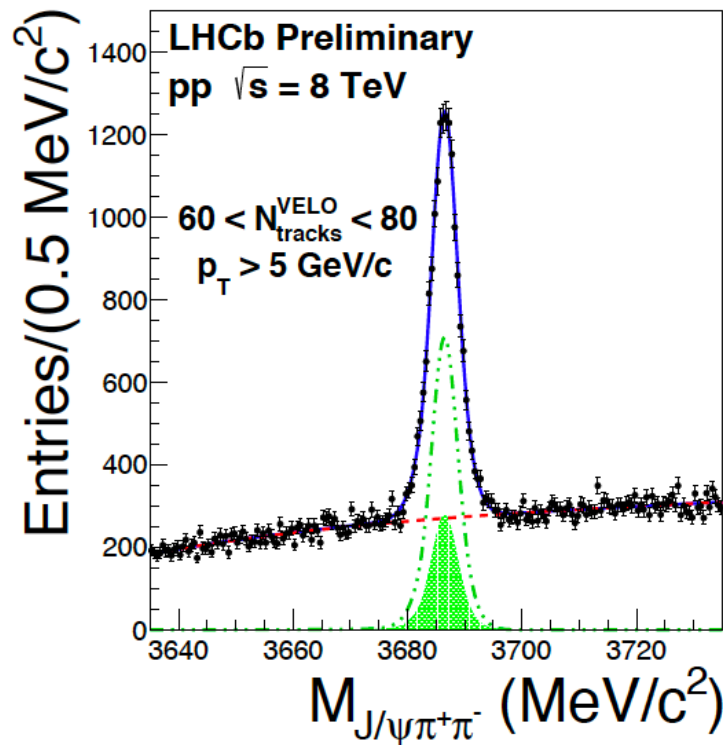
Direct comparison between conventional charmonium $\psi(2S)$ and exotic $X(3872)$ via ratio of cross sections:

$$\frac{\sigma_{\chi_{c1}(3872)}}{\sigma_{\psi(2S)}} \times \frac{\mathcal{B}[\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-]}{\mathcal{B}[\psi(2S) \rightarrow J/\psi \pi^+ \pi^-]}$$



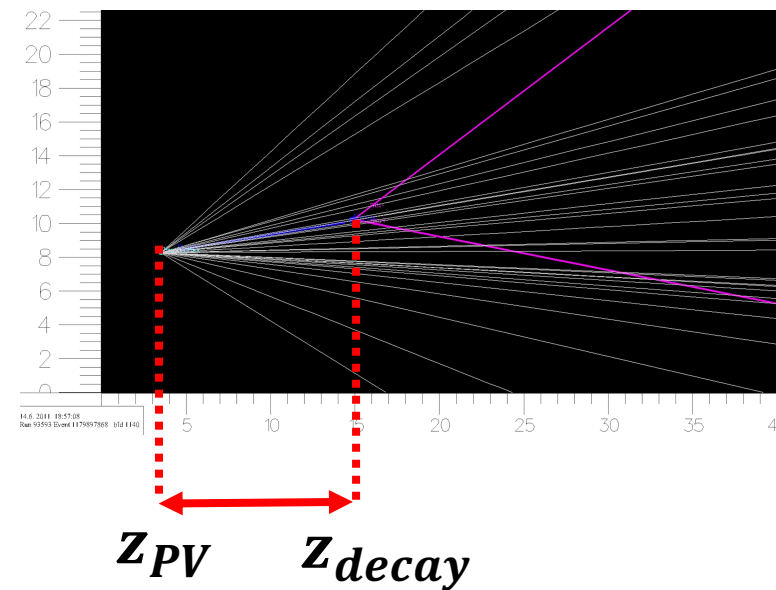
Prompt / b -decay separation

LHCb-CONF-2019-005



Simultaneous fit to invariant mass and pseudo proper time spectrum:

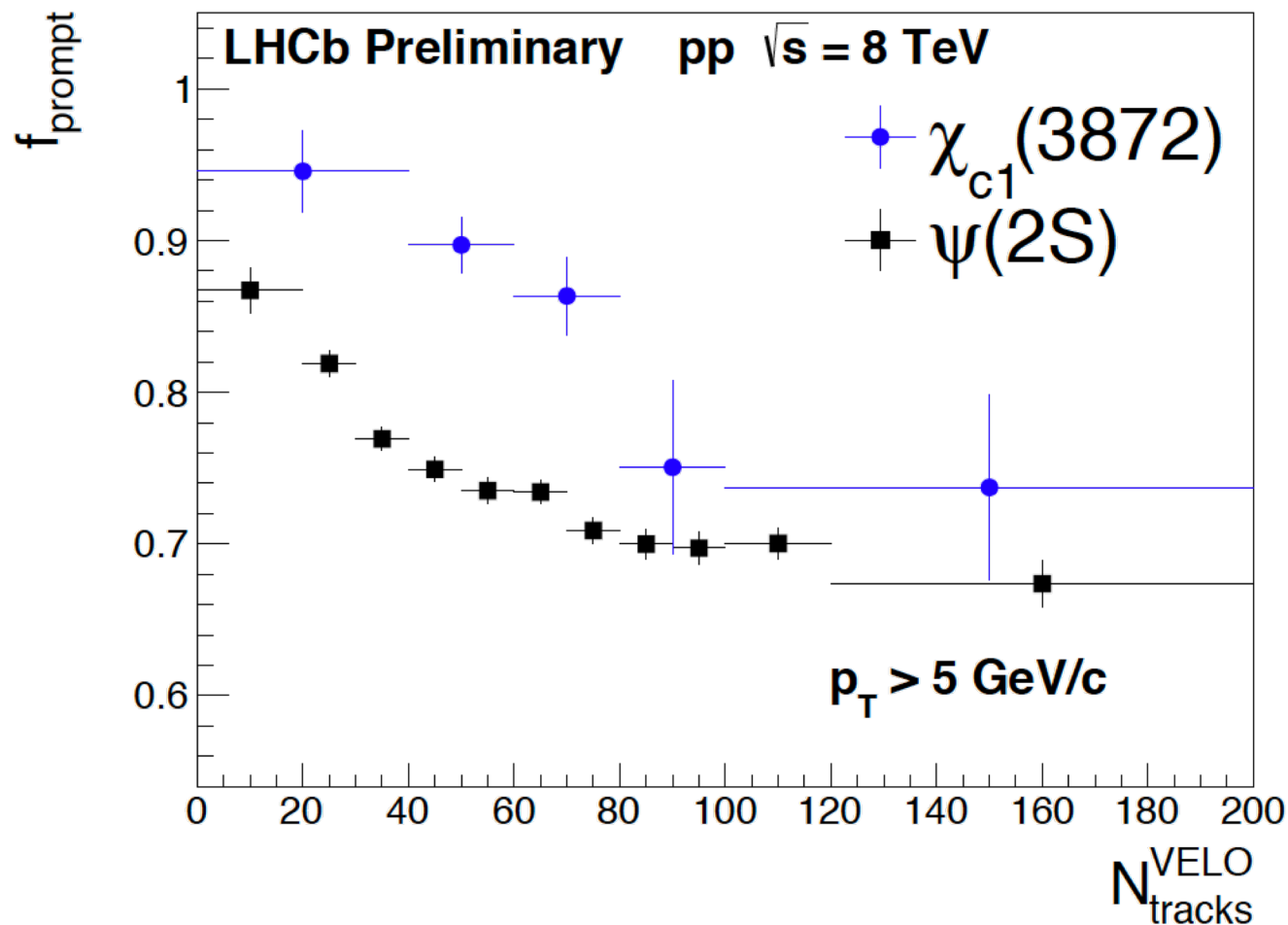
$$t_z = \frac{z_{\text{decay}} - z_{\text{PV}}}{p_z} M$$



Fit to mass constrains S/B while fit to t_z constrains prompt fraction

Prompt fraction

LHCb-CONF-2019-005

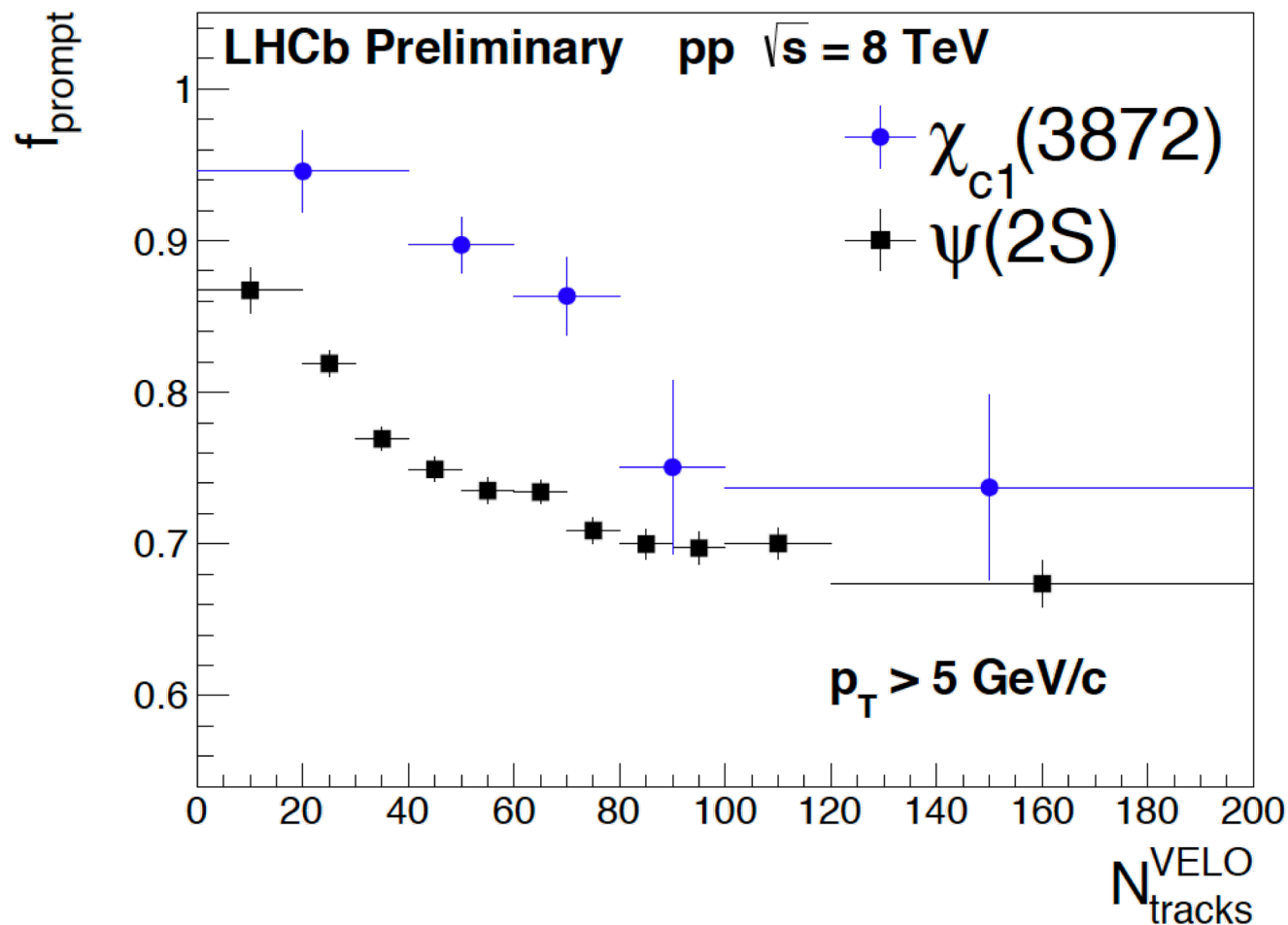


$$f_{\text{prompt}} = \frac{N_{\text{prompt}}}{N_{\text{prompt}} + N_{b\text{-decay}}}$$

- Significant decrease in prompt fraction of both $X(3872)$ and $\psi(2S)$ as event activity increases:

Prompt fraction

LHCb-CONF-2019-005

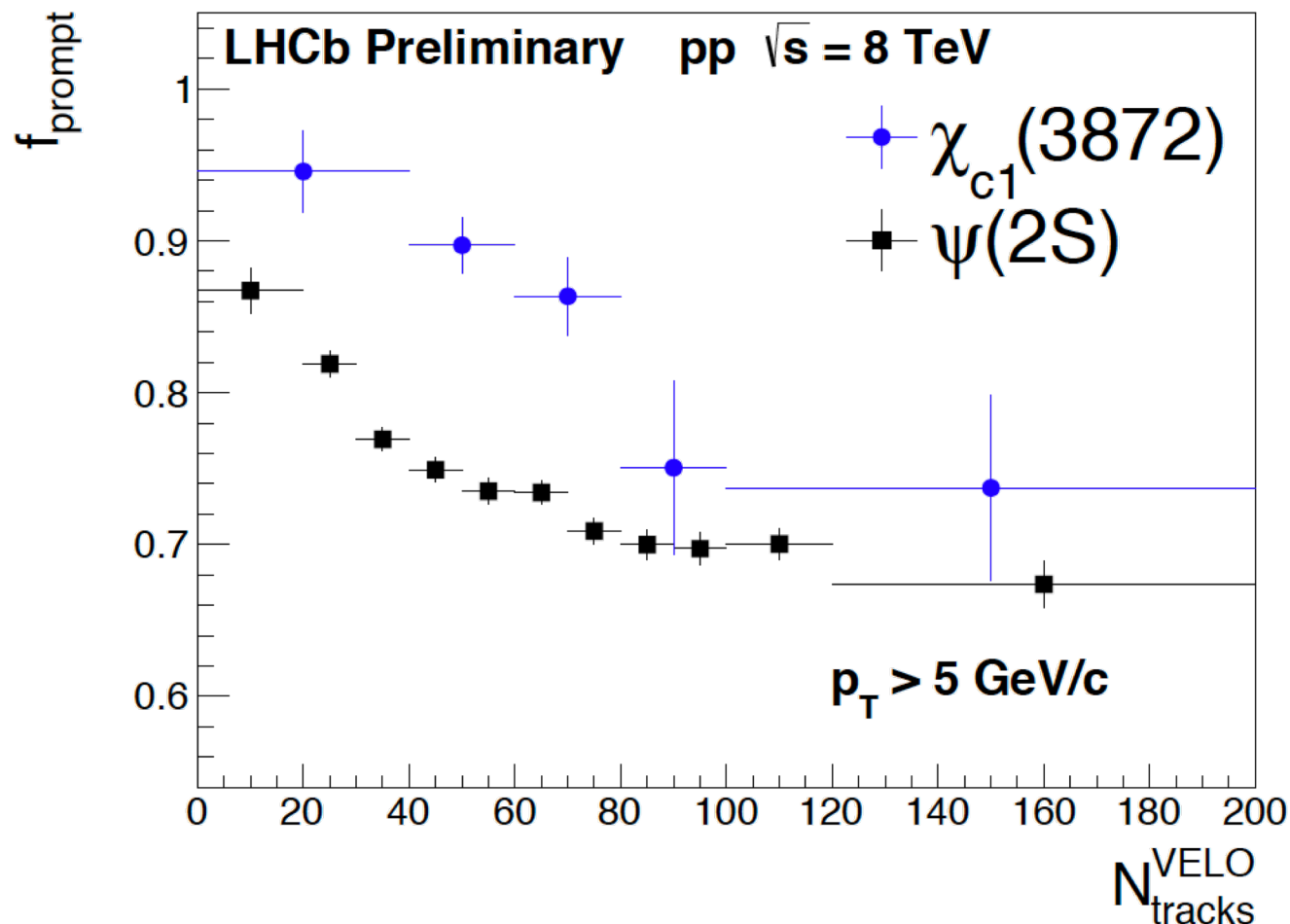


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- Significant decrease in prompt fraction of both $X(3872)$ and $\psi(2S)$ as event activity increases:
- Events with $b\bar{b}$ production naturally have higher multiplicity, due to fragmentation and decays
 - OPAL, PLB 550 33 (2002)

Prompt fraction

LHCb-CONF-2019-005



$$f_{\text{prompt}} = \frac{N_{\text{prompt}}}{N_{\text{prompt}} + N_{b\text{-decay}}}$$

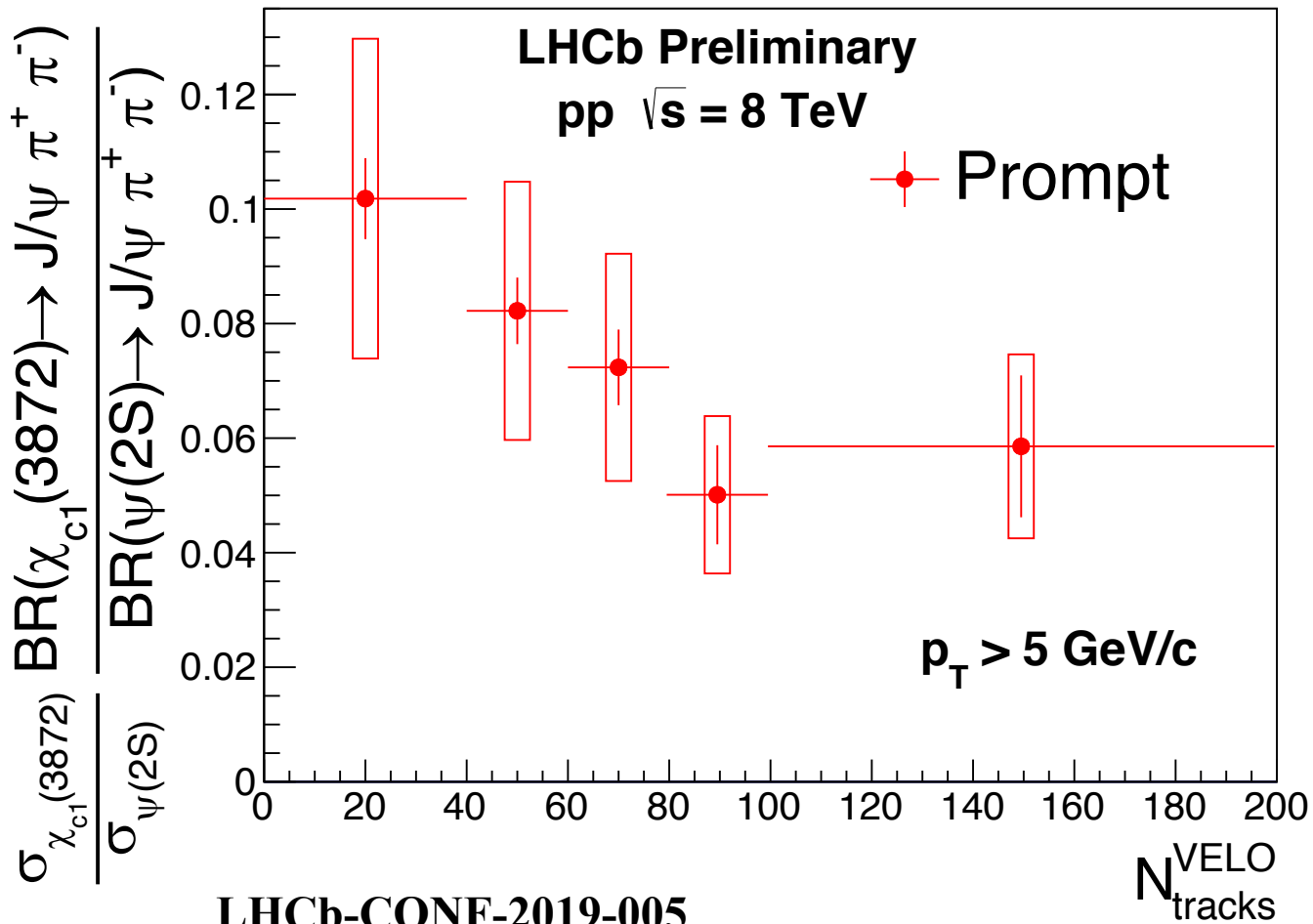
- Significant decrease in prompt fraction of both $X(3872)$ and $\psi(2S)$ as event activity increases:
- Events with $b\bar{b}$ production naturally have higher multiplicity, due to fragmentation and decays
 - OPAL, PLB 550 33 (2002)
- Formation of prompt $X(3872)$ and $\psi(2S)$ may be disrupted at the vertex, which cannot affect production via b decays in vacuum.

Ratio of cross sections

$$\frac{\sigma_{\chi_{c1}(3872)}}{\sigma_{\psi(2S)}} \times \frac{\mathcal{B}[\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-]}{\mathcal{B}[\psi(2S) \rightarrow J/\psi \pi^+ \pi^-]} = \frac{N_{\chi_{c1}(3872)} f_{\text{prompt}}^{\chi_{c1}(3872)}}{N_{\psi(2S)} f_{\text{prompt}}^{\psi(2S)}} \times \frac{\varepsilon_{\psi(2S)}}{\varepsilon_{\chi_{c1}(3872)}}$$

Ratio of cross sections

$$\frac{\sigma_{\chi_{c1}(3872)}}{\sigma_{\psi(2S)}} \times \frac{\mathcal{B}[\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-]}{\mathcal{B}[\psi(2S) \rightarrow J/\psi \pi^+ \pi^-]} = \frac{N_{\chi_{c1}(3872)} f_{\text{prompt}}^{\chi_{c1}(3872)}}{N_{\psi(2S)} f_{\text{prompt}}^{\psi(2S)}} \times \frac{\varepsilon_{\psi(2S)}}{\varepsilon_{\chi_{c1}(3872)}}$$

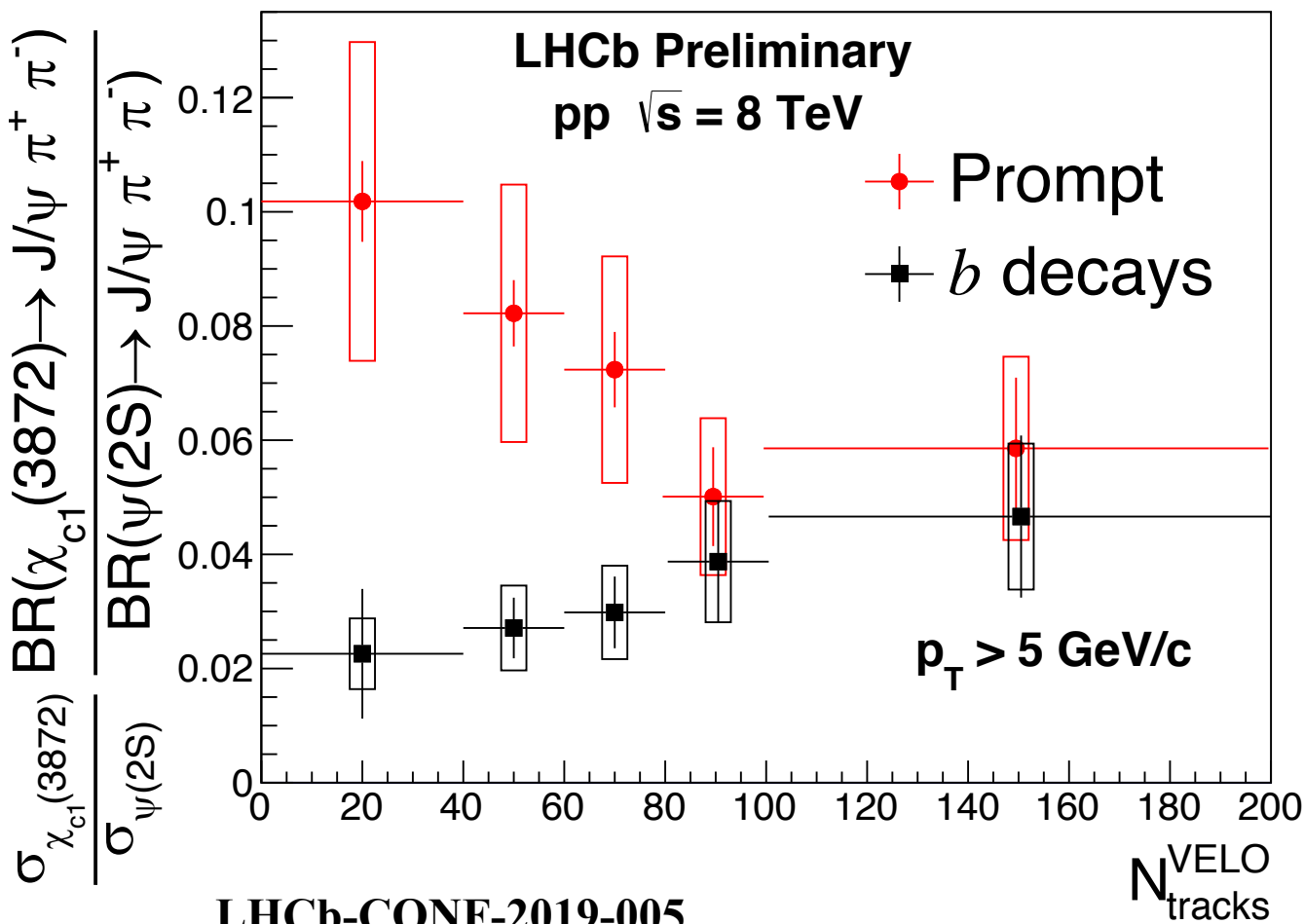


Prompt Component:
Increasing suppression of
 $\chi(3872)$ production relative to
 $\psi(2S)$ as event activity increases

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Ratio of cross sections

$$\frac{\sigma_{\chi_{c1}(3872)}}{\sigma_{\psi(2S)}} \times \frac{\mathcal{B}[\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-]}{\mathcal{B}[\psi(2S) \rightarrow J/\psi \pi^+ \pi^-]} = \frac{N_{\chi_{c1}(3872)} f_{\text{prompt}}^{\chi_{c1}(3872)}}{N_{\psi(2S)} f_{\text{prompt}}^{\psi(2S)}} \times \frac{\varepsilon_{\psi(2S)}}{\varepsilon_{\chi_{c1}(3872)}}$$



Prompt Component:
Increasing suppression of **$X(3872)$** production relative to **$\psi(2S)$** as event activity increases

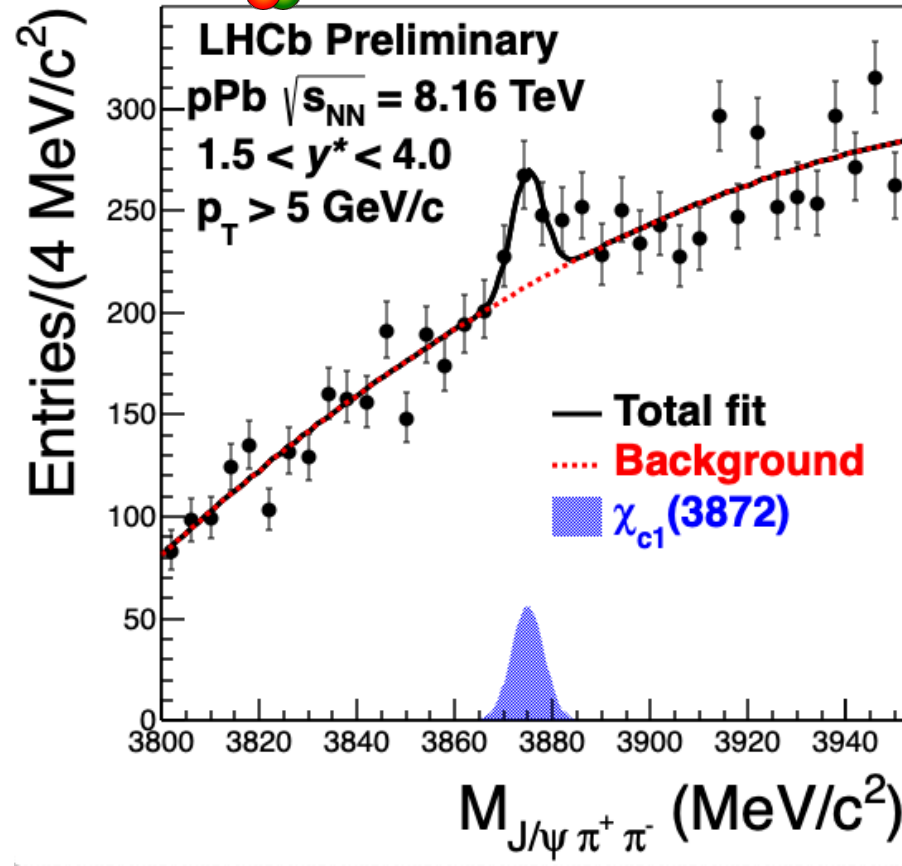
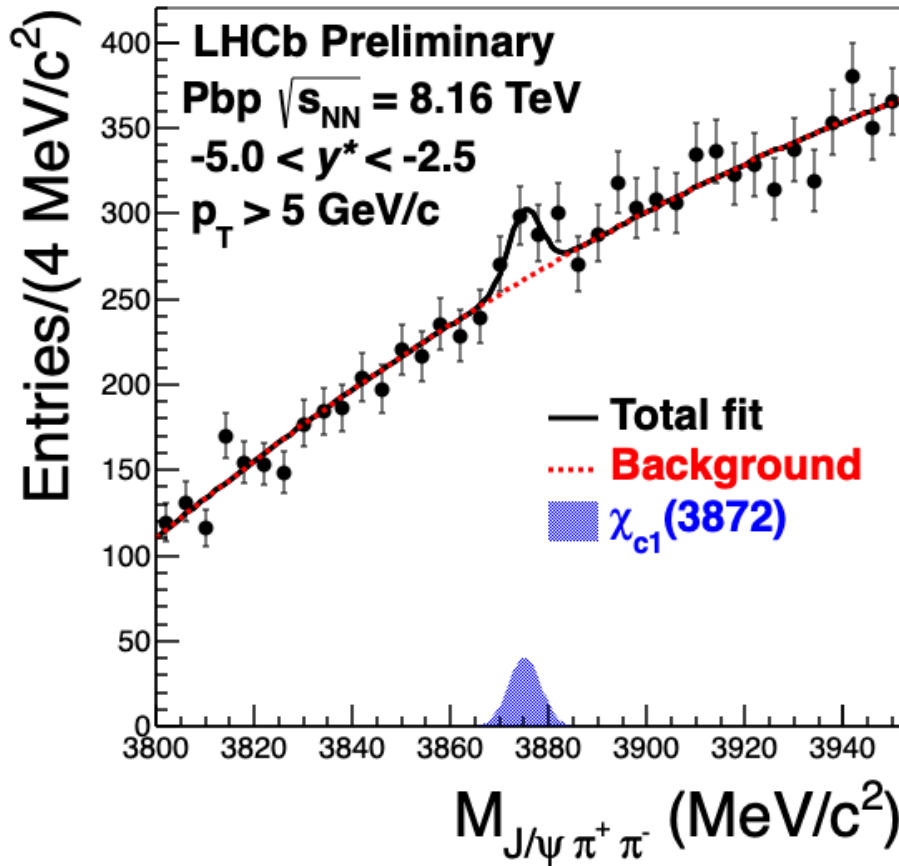
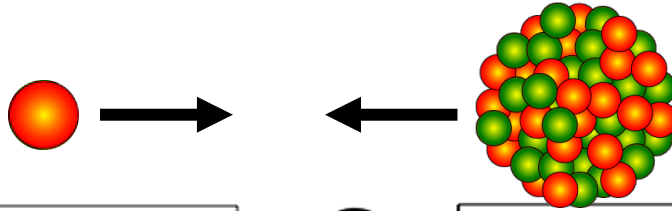
b -decay component:
No significant change in relative production, as expected for decays in vacuum. Ratio is set by b decay branching fractions.

Consistent with ATLAS measurement
 $R = 0.0395 \pm 0.0032 \pm 0.0008$ ($p_T > 10 \text{ GeV}/c$)

JHEP 2017:117 (2017)

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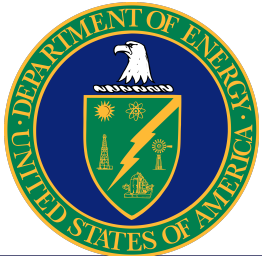
X(3872) in pPb collisions



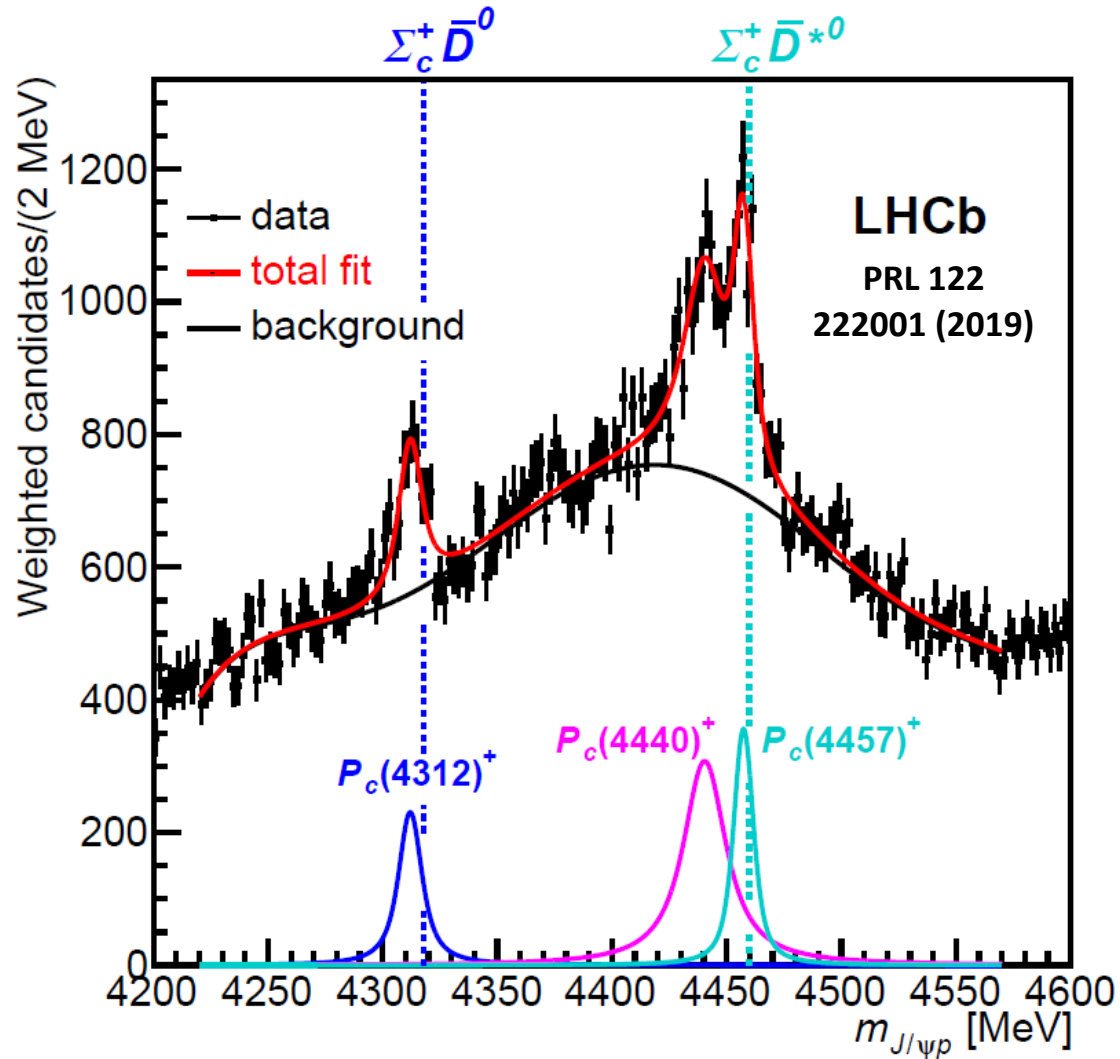
Theorists: predictions welcome

Summary

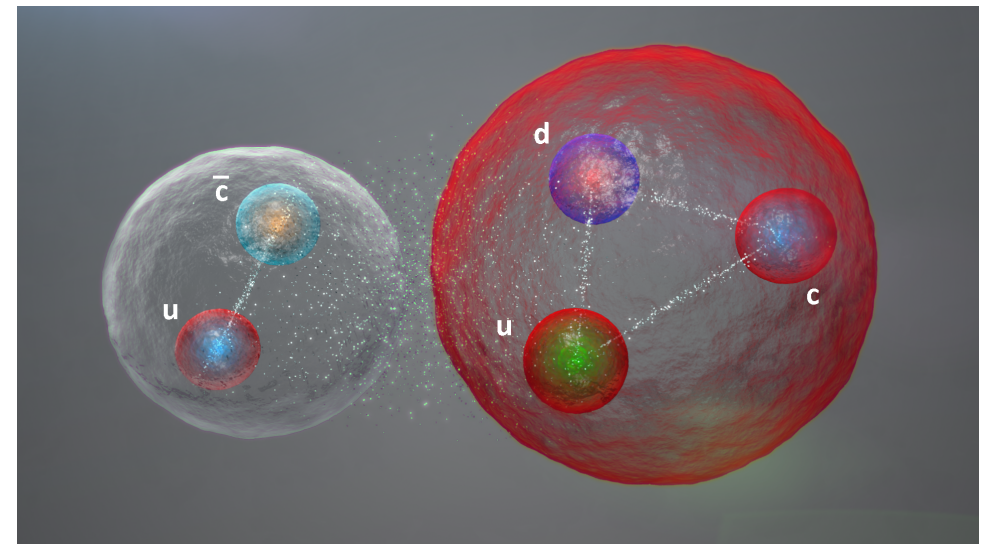
- The study of exotic hadrons is an active area of discovery in QCD
- Data and techniques from heavy ion physics give us a new window into dynamics of exotic states in a dense QCD environment:
 - Prompt fraction of $X(3872)$ and $\psi(2S)$ decreases with multiplicity in pp
 - Relative production in b -decays shows no significant change with multiplicity
 - Indications that prompt $X(3872)$ may be suppressed more than prompt $\psi(2S)$ as multiplicity increases
- Consistent with the interpretation of $X(3872)$ as a large, weakly bound state such as a hadronic molecule.



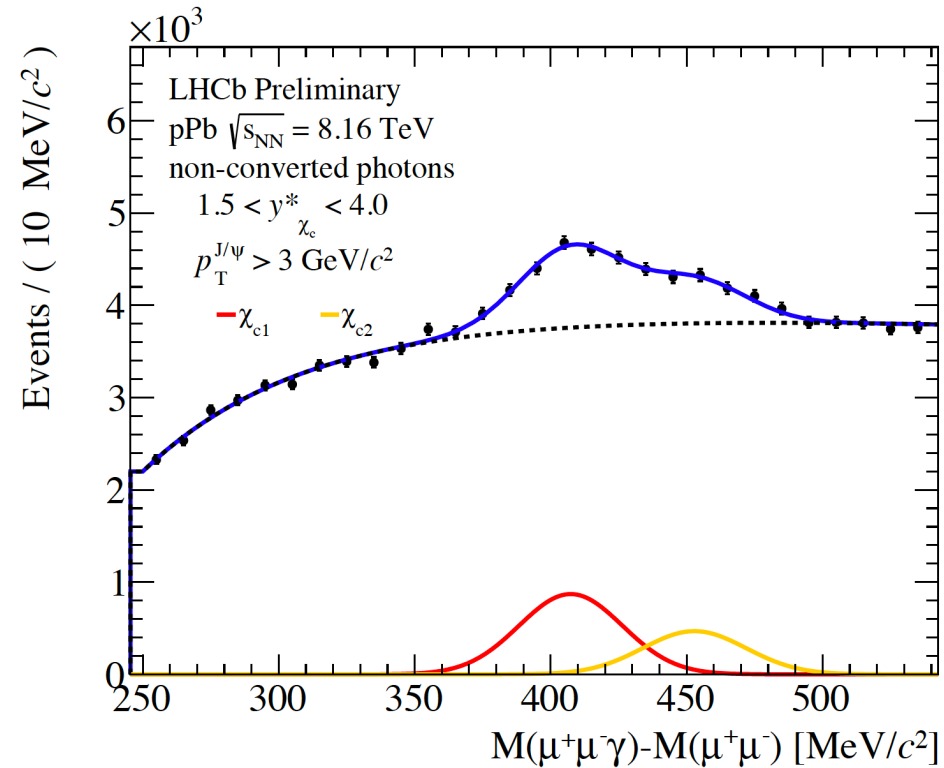
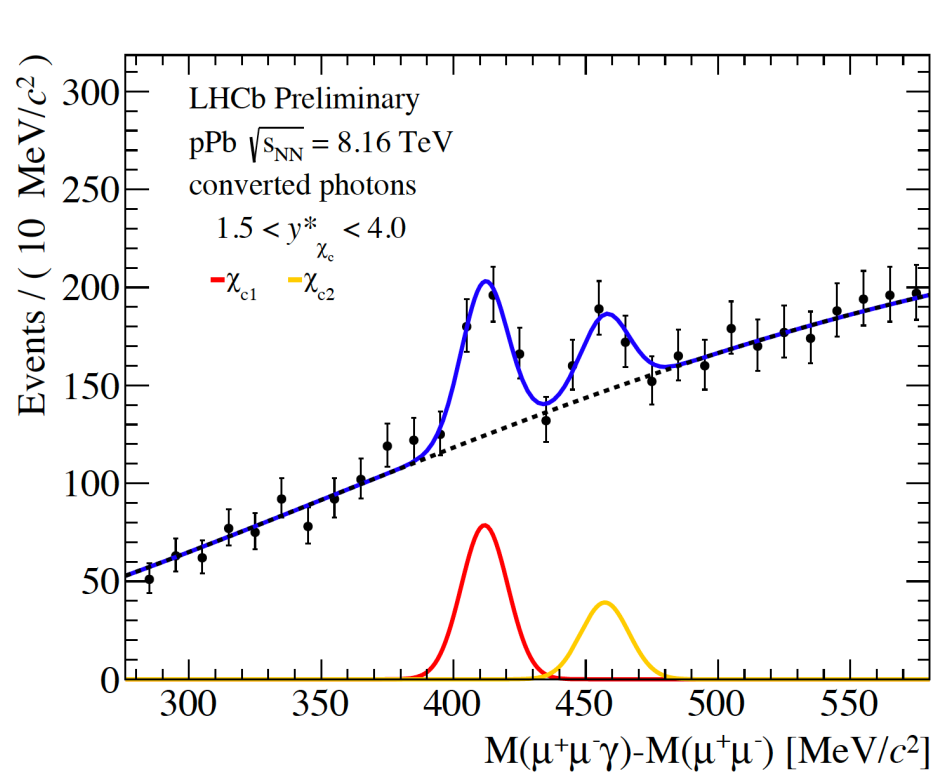
This work is supported by the US Dept. of Energy/Office of Science/Nuclear Physics Division



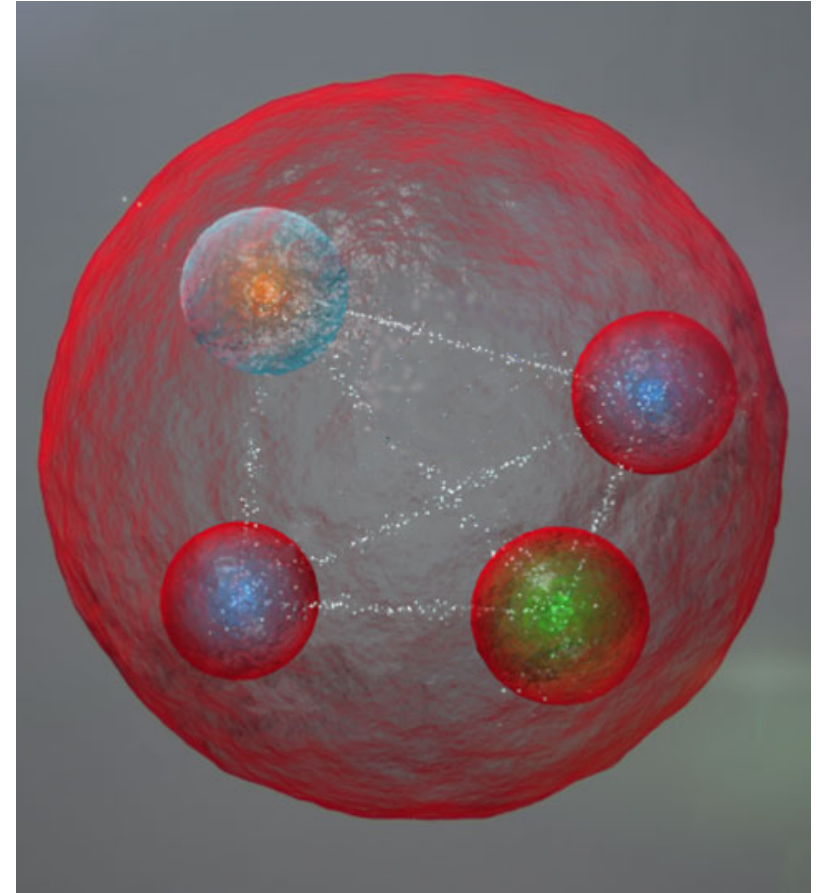
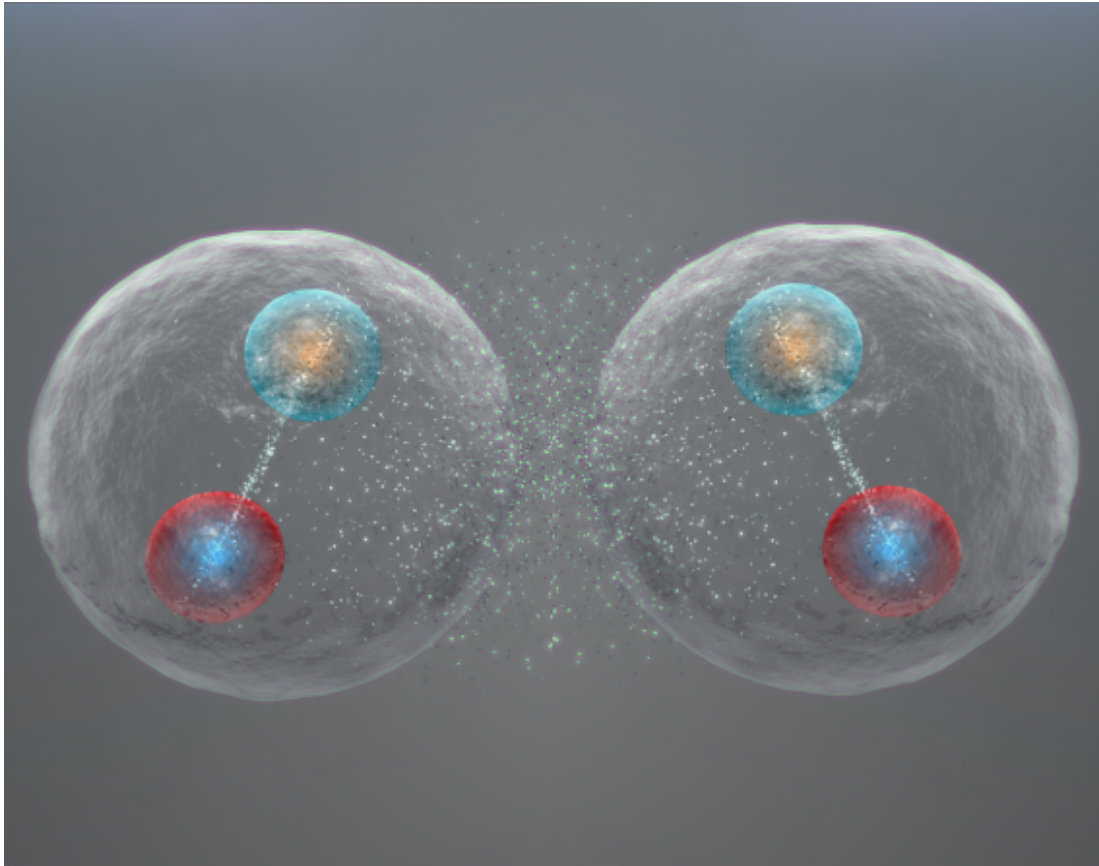
P_c pentaquark states recently discovered by LHCb are very close to mass thresholds for hadronic molecules.



χ_c States in pPb



See Poster by Jana Crkovska





The LHCb collaboration

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ABSTRACT: Using proton-proton collision data, collected with the LHCb detector and corresponding to 1.0, 2.0 and 1.9 fb⁻¹ of integrated luminosity at the centre-of-mass energies of 7, 8, and 13 TeV, respectively, the decay $\Lambda_b^0 \rightarrow \chi_{c1}(3872)pK^-$ with $\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-$ is observed for the first time. The significance of the observed signal is in excess of seven standard deviations. It is found that $(58 \pm 15)\%$ of the decays proceed via the two-body intermediate state $\chi_{c1}(3872)\Lambda(1520)$. The branching fraction with respect to that of the $\Lambda_b^0 \rightarrow \psi(2S)pK^-$ decay mode, where the $\psi(2S)$ meson is reconstructed in the $J/\psi \pi^+ \pi^-$ final state, is measured to be:

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \chi_{c1}(3872)pK^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \psi(2S)pK^-)} \times \frac{\mathcal{B}(\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-)}{\mathcal{B}(\psi(2S) \rightarrow J/\psi \pi^+ \pi^-)} = (5.4 \pm 1.1 \pm 0.2) \times 10^{-2},$$

where the first uncertainty is statistical and the second is systematic.

KEYWORDS: B physics, Branching fraction, Exotics, Hadron-Hadron scattering (experiments)

ARXIV EPRINT: [1907.00954](https://arxiv.org/abs/1907.00954)