Multi-stage evolution of heavy quarks in the quark-gluon plasma



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for the JETSCAPE Collaboration

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Natural Sciences and Engineering Research Council of Canada Conseil de recherches en sciences naturelles et en génie du Canada





UNIVERSITY

JETSCAPE simulates hard & soft sectors of heavy ion collisions



- Soft hadrons: Ongoing JETSCAPE Bayesian Analysis (see talk by Jean-François Paquet) constraining the QGP transport properties.
- Hard Partons: Ongoing multi-stage energy loss of hard partons with soft QGP (see also Amit Kumar, Wed) and jet-energy deposition into QGP (see talk by Yasuki Tachibana, Tue).

 JETSCAPE Physics Working Group focused on describing interactions of hard probes with the QGP: 38 members in total (see <u>http://jetscape.org/phys/</u>)



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Different physics mechanisms for in-medium energy loss in different kinematic regimes ⇒ a multi-stage approach is needed for accurate description

- MATTER (The Modular All Twist Transverse-scattering Elastic-drag and Radiation) valid for High E, High Q
 - Virtuality-ordered shower with splittings above $Q \gg Q_{\min}$
 - Splittings happen via the Sudakov form factor [Adv.Ser.Direct.HEP, 573 (1989); NPA 696, 788 (2001)] valid for all partonic processes and can include in-medium corrections

$$\Delta(Q_{\max}, Q \ge Q_{\min}) = \exp\left[-\frac{\alpha_s}{2\pi} \int_{Q^2}^{Q_{\max}^2} \frac{d(Q^2)}{Q^2} \int_{y_{\min}}^{y_{\max}} dy \,\mathcal{P}(y)\right]$$

• The limits and of the integral and the splitting(s) \mathcal{P} depend on the incoming and outgoing partons species and possible in-medium contributions.

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- Q_{\max} is set by hard cross section momentum exchange
- Emission Kinematics dictates that:

$$Q_{\min}^{2} = \frac{Q_{0}^{2}}{2} \left[1 + \sqrt{1 + \frac{4M^{2}}{Q_{0}^{2}}} \right] \qquad \qquad y_{\min} = \frac{Q_{0}^{2}}{2Q^{2}} + \frac{M^{2}}{M^{2} + Q^{2}} \qquad \qquad y_{\max} = 1 - \frac{Q_{0}^{2}}{2Q^{2}}$$

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• $\mathcal{P}(y)$: vacuum/medium splitting derived via SCET [**PRC 94, 054902 (2016)**] at LO in $\left(\alpha_{s}, \frac{M^{2}}{Q^{2}}\right)$ $\mathcal{P}(y) = P(y) + \frac{P(y)\left[\left\{\left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^{2}\right\}\left\{\int_{\tau_{i}}^{\tau_{f}} dt \,\hat{q}(t)\left[4\sin^{2}\left[\frac{t - \tau_{i}}{2\tau_{f}}\right]\right]\right\} + \hat{e}\xi...\} + \hat{e}_{2}\{\ldots\}\right]}{y(1 - y)Q^{2}(1 + \chi)^{2}}$ $\chi = \frac{y^{2}M^{2}}{y(1 - y)Q^{2} - y^{2}M^{2}}$ $\hat{q} \propto \alpha_{s}^{2}T^{3} \ln\left[\frac{cE}{\alpha_{s}T}\right]$

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• $\mathcal{P}(y)$ not yet derived using SCET for heavy flavor, so the light flavor prescription at LO in $\left(\alpha_{s}, \frac{1}{\alpha^{2}}\right)$ is used as ansatz $\mathcal{P}(\mathbf{y}) = P(\mathbf{y}) + \frac{P(\mathbf{y}) \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) 4\sin^2\left[\frac{t-\tau_i}{2\tau_f}\right]}{y(1-y)Q^2}$ $\hat{q} \propto \alpha_s^2 T^3 \ln\left[\frac{cE}{\alpha_s T}\right]$

Linear Boltzmann Transport for Heavy Quarks

- Valid for high E, assuming particles are on-shell
- Solves the time-ordered evolution for the phase space distribution function

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• The G_{inel} calculates medium-induced stimulated $1 \rightarrow 2$ emission at LO in $\left(\alpha_s, \frac{M^2}{Q^2}\right)$ [see **PRC 94, 054902 (2016)**]

$$\begin{aligned} \mathcal{G}_{inel} &= \int \frac{d(Q^2)}{Q^2} \frac{\alpha_s(Q^2)}{2\pi} \int dy \,\mathcal{P}(y) \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \left\{ \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) \left[4\sin^2 \left[\frac{t - \tau_i}{2\tau_f} \right] \right] \right\} + \hat{e}_{1} \left\{ \dots \right\} + \hat{e}_{2} \left\{ \dots \right\} \right] \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \left\{ \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) \left[4\sin^2 \left[\frac{t - \tau_i}{2\tau_f} \right] \right] \right\} + \hat{e}_{2} \left\{ \dots \right\} \right] \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \left\{ \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) \left[4\sin^2 \left[\frac{t - \tau_i}{2\tau_f} \right] \right] \right\} + \hat{e}_{2} \left\{ \dots \right\} \right\} \right] \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \left\{ \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) \left[4\sin^2 \left[\frac{t - \tau_i}{2\tau_f} \right] \right] \right\} + \hat{e}_{2} \left\{ \dots \right\} \right\} \right] \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \left\{ \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) \left[4\sin^2 \left[\frac{t - \tau_i}{2\tau_f} \right] \right] \right\} \right\} + \hat{e}_{2} \left\{ \dots \right\} \right\} \right] \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \left\{ \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) \left[4\sin^2 \left[\frac{t - \tau_i}{2\tau_f} \right] \right] \right\} \right\} \right\} \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \left\{ \int_{\tau_i}^{\tau_f} dt \,\hat{q}(t) \left[4\sin^2 \left[\frac{t - \tau_i}{2\tau_f} \right] \right] \right\} \right\} \right\} \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \right\} \right] \\ \mathcal{P}(y) &= P(y) + \frac{P(y) \left[\left\{ \left(1 - \frac{y}{2}\right) - \chi + \left(1 - \frac{y}{2}\right)\chi^2 \right\} \right\} \right\} \right\}$$

About the QGP medium simulations



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 - TRENTO initial conditions
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- The same underlying QGP simulation is used to study
 - Light flavor high- p_T hadrons and jets (see talk by **Amit Kumar, Wed**)
 - Jet-medium interactions (see talk by Yasuki Tachibana, Tue)
 - Heavy flavor interactions with the QGP (this talk)



- Jet and charged hadron R_{AA} used to tune parameters, i.e. $\alpha_s = 0.25$ and $Q_s = 2GeV$ (for more details see talk by **Amit Kumar**)
- No additional tuning was done for D^0 meson R_{AA}

D^0 mesons R_{AA} vs CMS data from PbPb at $\sqrt{s_{NN}} = 5.02 \ TeV$



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 - $\Rightarrow R_{AA}$ can go much closer to 0.1
- Using a multi-scale approach allows to balance the contributions from few vs multiple scattering, and ultimately gives an improved description of R_{AA}

Conclusion and outlook

- A multi-scale formalism, such as that present inside the JETSCAPE framework, allows for an improved description of D^0 energy loss inside QGP, because the effects of multiple vs few scatterings is included.
- Furthermore, JETSCAPE provides a unified framework where these formalisms can be combined, which allows for a simultaneous description of both light and heavy parton energy loss.
- Future studies will explore:
 - Open bottom hadron *R*_{AA}
 - the effects of additional transport coefficients relevant for heavy flavor production, i.e. longitudinal drag (\hat{e}) and diffusion (\hat{e}_2) coefficients.