

New Challenges for open HF physic in URHICs

Yifeng Sun

INFN-LNS



Istituto Nazionale di Fisica Nucleare

In collaboration with:

S. Plumari, S. K. Das, V. Minissale, G. Coci, M. Ruggieri, V. Greco

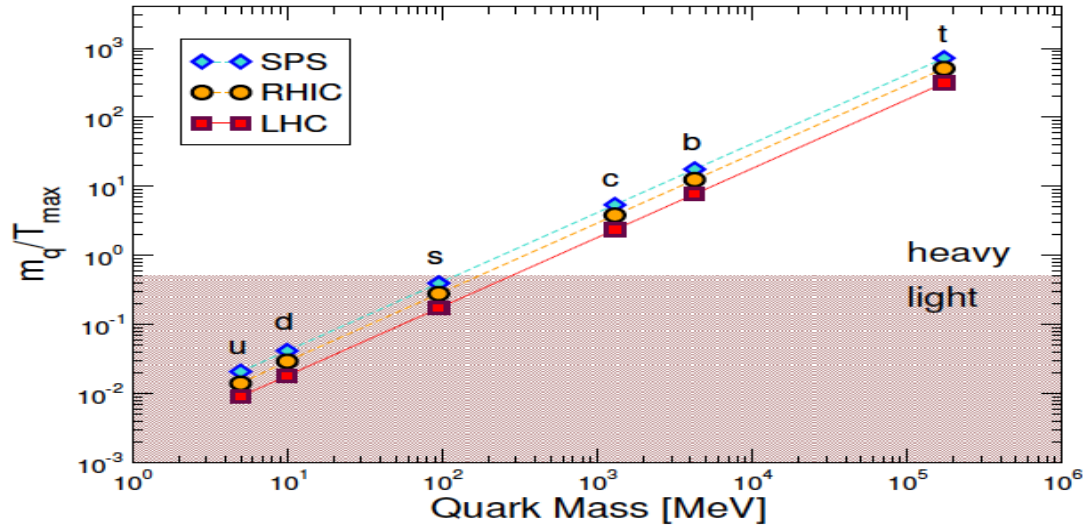


Outline

- ✧ HF dynamics in QGP: estimates of T dependence of D_s
- ✧ Hadronization mechanism on Λ_c/D ratios
- ✧ New observables: $v_n(\text{light}) - v_n(\text{heavy})$, σ_{v_n}/v_n
- ✧ Impact of initial stages on HF dynamics: EM field, vorticity and Glasma

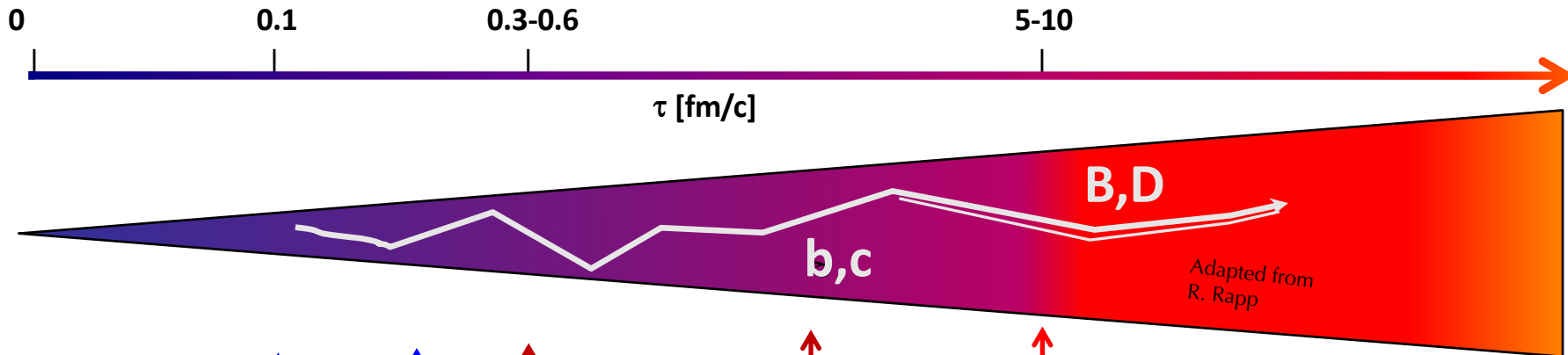
HF dynamics in QGP and estimates of T dependence of D_s

Specific of Heavy Quarks



- $m_{c,b} \gg \Lambda_{\text{QCD}}$ production by pQCD hard processes
- $m_{c,b} \gg T_{\text{RHIC,LHC}}$ negligible thermal production
- $m_{c,b} \gg q^2 \sim gT_{\text{RHIC,LHC}}$ soft scatterings \rightarrow Brownian motion
- $\tau_0 \approx 1/2 m_Q \ll \tau_{\text{QGP}}$ witness of all the QGP evolution
- $\tau_{\text{th}} \approx \tau_{\text{QGP}} \gg \tau_{q,g}$ carry more information of their evolution

HF in all stages of URHICs



HQs prod.

HQs in:
- Glasma
- EM field

Formation of QGP
with large vorticity

• Dynamics in QGP
- Transp. coeff. of QGP
- Thermalization time

• Hadronization mechanism
- coalescence and/or fragm.

Relativistic Boltzmann at finite η/s

$$p^\mu \partial_\mu f_q(x, p) + m(x) \partial_\mu^x m(x) \partial_p^\mu f_q(x, p) = C[f_q, f_g]$$

$$p^\mu \partial_\mu f_g(x, p) + m(x) \partial_\mu^x m(x) \partial_p^\mu f_g(x, p) = C[f_q, f_g]$$

free streaming

Interaction $\varepsilon - 3p \neq 0$

collision term

Heavy Quarks evolution

$$p^\mu \partial_\mu f_Q(x, p) = C[f_q, f_g, f_Q](x, p)$$

M scattering matrix by QPM model
fit to IQCD thermodynamics

gauged to some $\eta/s \neq 0$

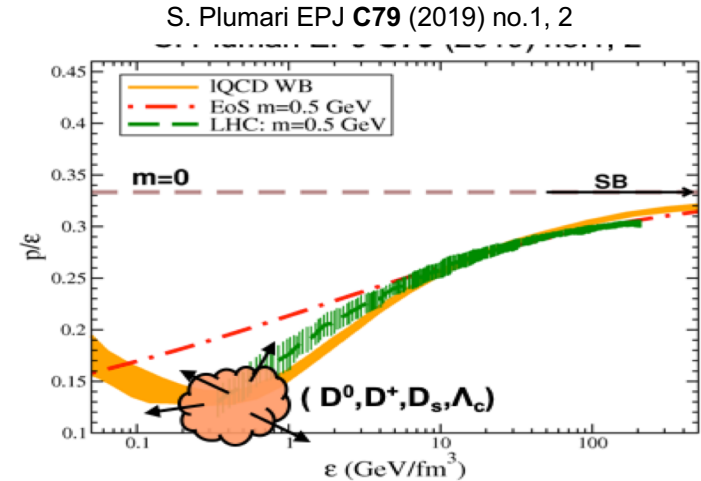
Equivalent to viscous hydro

Hadronization by coalescence plus fragmentation

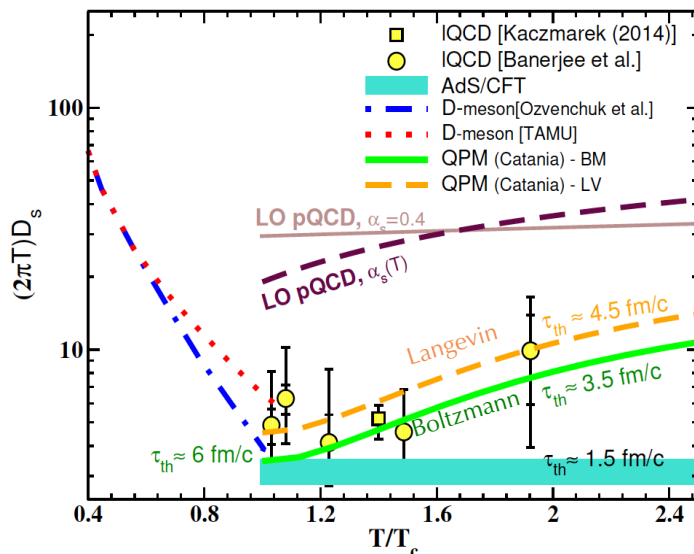
S. Plumari et al., EPJC78 (2018) no.4, 348

$$\left\{ \begin{aligned} \frac{dN_{Hadron}}{d^2 p_T} &= g_H \int \prod_{i=1}^n p_i d\sigma_i \frac{d^3 p_i}{(2\pi)^3} f_q(x_i, p_i) f_W(x_1, \dots, x_n; p_1, \dots, p_n) \delta(p_T - \sum_i p_{iT}) \\ \frac{dN_h}{d^2 p_h} &= \sum_f \int dz \frac{dN_f}{d^2 p_f} D_{f \rightarrow h}(z) \end{aligned} \right.$$

Good description of spectra $D^0, D^+, D_s, \Lambda_c, B, \Lambda_b$



What is the underlying D_s ?



Scardina et al., PRC96(2017) 044905

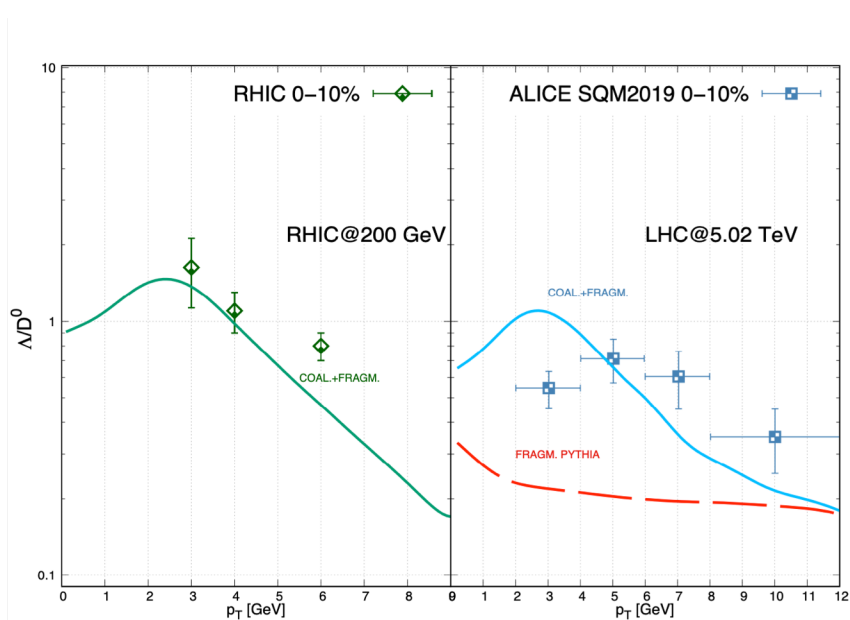
Not a model fit to
IQCD data!
The result is gotten
from the simultaneous
description of $R_{AA}(p_T)$ &
 $v_2(p_T)$

- ❖ Get close AdS/CFT at $T \approx T_c$
- ❖ We have a probe with $\tau_{\text{therm}} \approx \tau_{\text{QGP}}$

$$\tau_{th} = \frac{M}{2\pi T^2} (2\pi T D_s) \cong 1.8 \frac{2\pi T D_s}{(T/T_c)^2} \text{ fm/c}$$

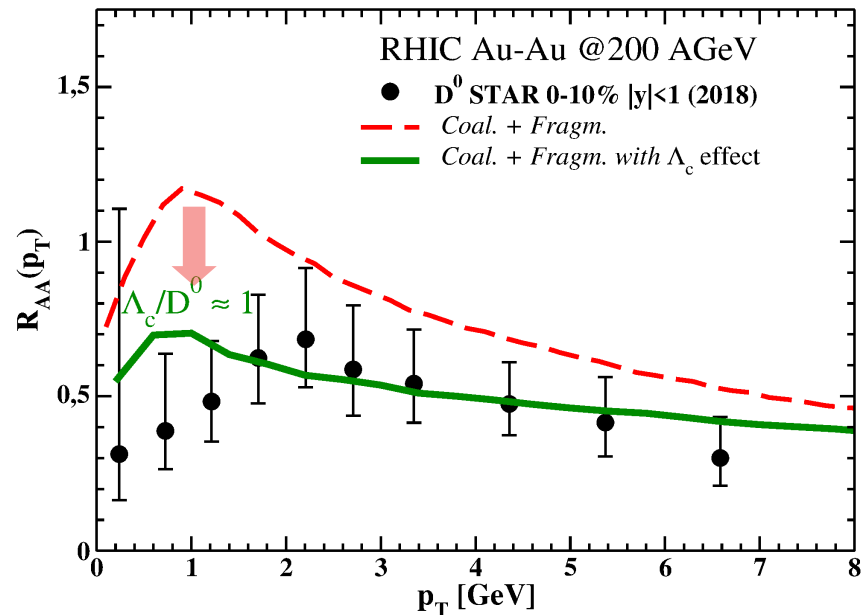
Hadronization mechanism on Λ_c/D ratio

Λ_c/D from hadronization mechanism



S. Plumari et al., EPJC78 (2018) no.4, 348

- With the same coalescence plus fragmentation model we describe the Λ_c/D^0 for RHIC and LHC

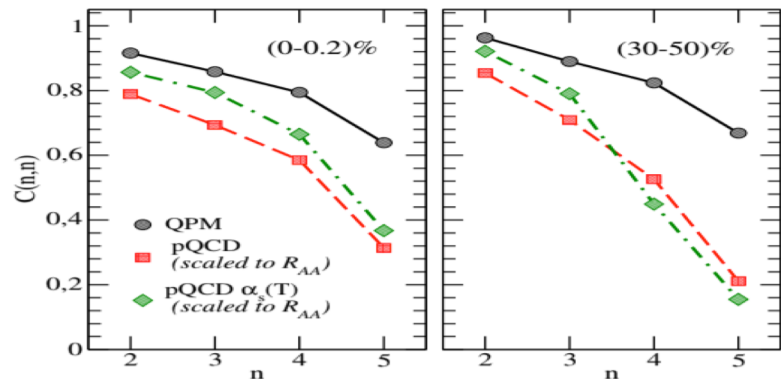


STAR Collaboration PRC99 (2019) no.3, 034908

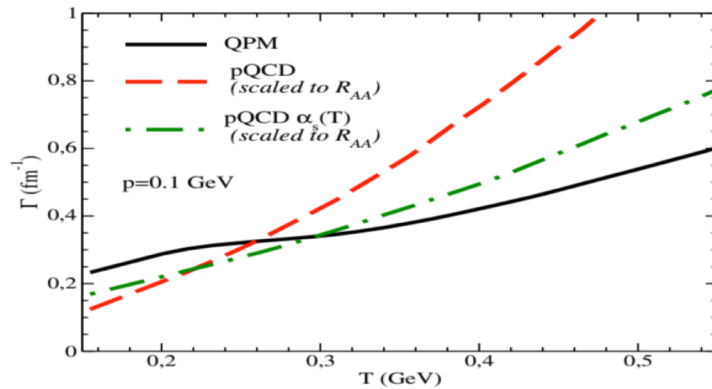
- This opens a new paradigm in studying HF
- The effect will be even larger for because larger Λ_b/B

New observables: $v_n(\text{light}) - v_n(\text{heavy})$, σ_{v_n}/v_n

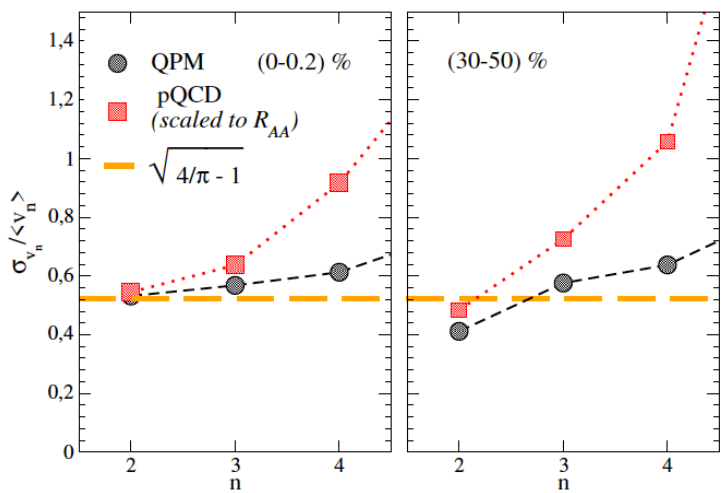
Sensitivity of $C(n,n)$ and v_n variances on $D_s(T)$



$$C(v_n^{light}, v_m^{heavy}) = \left(\frac{(v_n^{light} - \langle v_n^{light} \rangle)(v_m^{heavy} - \langle v_m^{heavy} \rangle)}{\sigma_{v_n^{light}} \sigma_{v_m^{heavy}}} \right)$$



- $C(n,n)$ and σ_{v_n}/v_n is sensitive to T dependence of D_s



Impact of initial stages on HF dynamics by EM field, vorticity and Glasma

HF dynamics in vorticity and EM field

We solve the relativistic Boltzmann eq coupled with the **external EM field**.

$$p^\mu \partial_\mu f_q(x, p) + m(x) \partial_\mu^x m(x) \partial_p^\mu f_q(x, p) + q F_{ext}^{\mu\nu} p_\mu \partial_\mu f_q(x, p) = C[f_q, f_g]$$

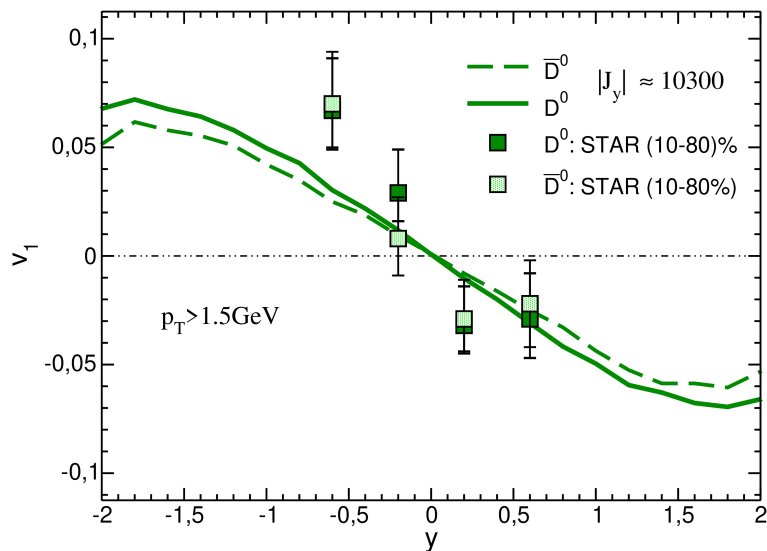
$$p^\mu \partial_\mu f_g(x, p) + m(x) \partial_\mu^x m(x) \partial_p^\mu f_g(x, p) = C[f_q, f_g]$$

Heavy quark evolution

$$p^\mu \partial_\mu f_Q(x, p) + q F_{ext}^{\mu\nu} p_\mu \partial_\mu f_Q(x, p) = C[f_q, f_g, f_Q]$$

- Charm transport coefficient constrained by experimental data on the $R_{AA}(p_T)$ and v_2 of D meson
- Vorticity employed for Λ polarization and v_1 of light particles
- EM field is calculated with fast moving protons in colliding nuclei assuming a constant conductivity of medium

Impact of EM Field + Vorticity on D



V. Greco, QM2017, NPA (2017)

S. Das et al., PLB768 (2017)

- Large average v_1 of D meson due to vorticity, a lot larger than for charged particles
- A sizable v_1 splitting for D meson than for charged particles predicted $v_1 \approx 10^{-3}-10^{-4}$ [Gursoy et al., PRC89 (2014)] using the same EM field

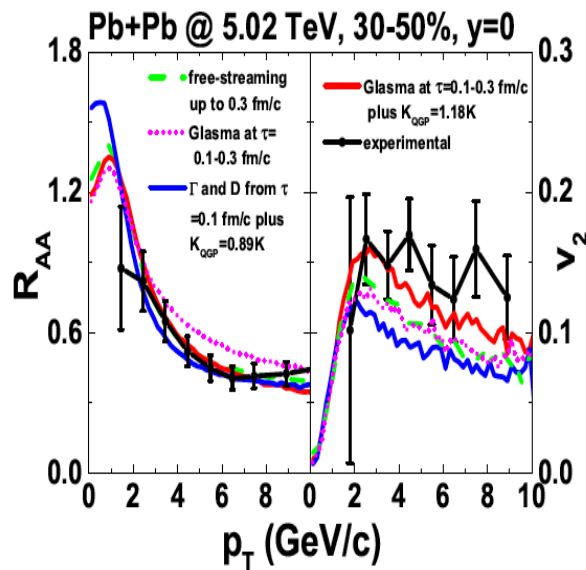
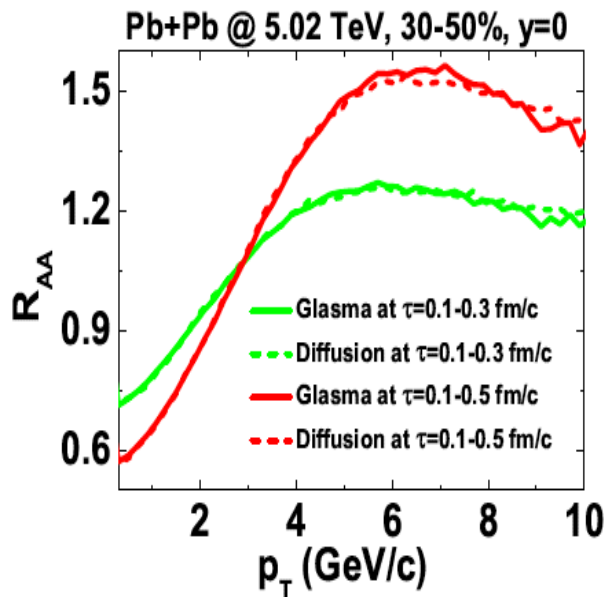
- At LHC energy we get much less than experimental data (???)

HQ in Glasma

$$\frac{dx_i}{dt} = \frac{p_i}{E},$$

$$E \frac{dp_i}{dt} = Q_a F_{iv}^a p^v,$$

$$E \frac{dQ_a}{dt} = -Q_c \varepsilon^{cba} A_b \cdot p,$$



Y. Sun et al., PLB **798** (2019) 134933

✧ **Enhancement of $R_{AA}(p_T)$**

✧ The effect is similar to Fokker-Planck with only diffusion and negligible drag

✧ **Gain in v_2** : larger interaction in the QGP to have same $R_{AA}(p_T)$

Summary & Outlook

- ❖ Prediction of both R_{AA}, v_2 gives us **$D_s(T)$ of QGP \approx IQCD**
- ❖ Λ_c/D from Coal.+Fragm. explained in AA (RHIC & LHC)
- ❖ **$v_n(\text{HQ})-v_n(\text{QGP})$ and σ_{v_n}/v_n** , new observables probing $D_s(T)$
- ❖ Charm directed flow can probe the initial strong EM field and vorticity
- ❖ Glasma effect on HF

There are more things to do and to learn than expected!