

Chirality transfer & Chiral turbulence in gauge theories

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Based on:

Mace, Mueller, SS, Sharma [arXiv:1910:01654](https://arxiv.org/abs/1910.01654)

Mace, Mueller, SS, Sharma, PRD95 (2017) no.3, 036023;

Mueller, SS, Sharma PRL117 (2016) no.14, 142301;

Mace, SS, Venugopalan PRD 93 (2016) no.7, 074036;



 **UNIVERSITÄT
BIELEFELD**

 Faculty of Physics

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**CRC-TR 211**
Strong-interaction matter
under extreme conditions

Outline

- 1 Introduction & Motivation
- 2 Chirality transfer & Chiral turbulence in strongly coupled QED plasmas
- 3 Conclusions & Outlook

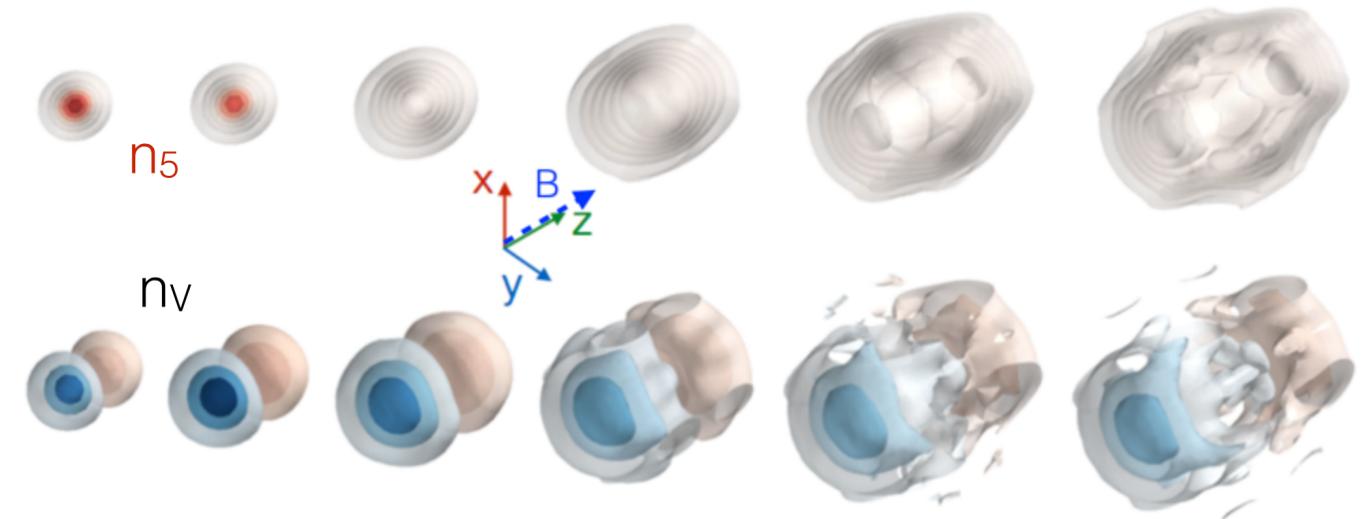
Chiral magnetic effect & anomalous transport

Discovery of new conductivities for systems with chiral fermions and (net-) chirality charge density

Fukushima, Kharzeev, Warringa PRD 78 (2008) 074033

$$\vec{j}_V \propto n_5 \vec{B}$$

j_V : vector current
 n_5 : chiral charge density
 B : magnetic field



Mueller, SS, Sharma PRL 117 (2016) no.14, 142301

Expected manifestations from Heavy-Ion collisions to Dirac/Weyl semi-metals

Difference to ordinary transport phenomena:

Energy-momentum tensor $T^{\mu\nu}$
 (Net-)baryon number n_B
 conserved

(Net-)chiral charge density n_5
 not conserved

-> Chiral transport phenomena are intrinsically non-equilibrium effects

Chiral anomaly in QCD and QED

Chiral charge density of fermions (n_5) not conserved due to the axial anomaly

$$\partial_\mu j_{5,f}^\mu = 2m_f \bar{q} \gamma_5 q - 2 \partial_\mu K^\mu \qquad \partial_\mu K^\mu = \frac{g^2}{32\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu}$$

quark mass Chern-Simons current (non)-abelian field-strength fluctuations

Gauge field contribution \propto E.B represents change of chiral charge density of gauge fields

$$N_{CS}(t) = \int d^3x K^0(t, x)$$

QED: Magnetic Helicity

QCD: Chern-Simons number

Chiral limit ($m \rightarrow 0$): Net chirality of the system $N_5 + 2N_{CS}$ conserved and thus shared & transferred between fermions and gauge fields

Extremely important to understand transfer/balance of chiral charge between fermions and gauge fields for macroscopic manifestations of anomaly

Chirality transfer in QED and QCD

Chirality transfer from gauge fields to fermions

QED: External E,B fields

QED & QCD: Space-time dependent fluctuations of (chromo-) electro-magnetic fields

QCD: Sphaleron transitions

Chirality transfer from fermions to gauge fields

QCD: Bias of sphaleron transition rate

McLerran, Mottola, Shaposhnikov PRD43 (1991) 2027-2035

QED & QCD: Chiral plasma instabilities

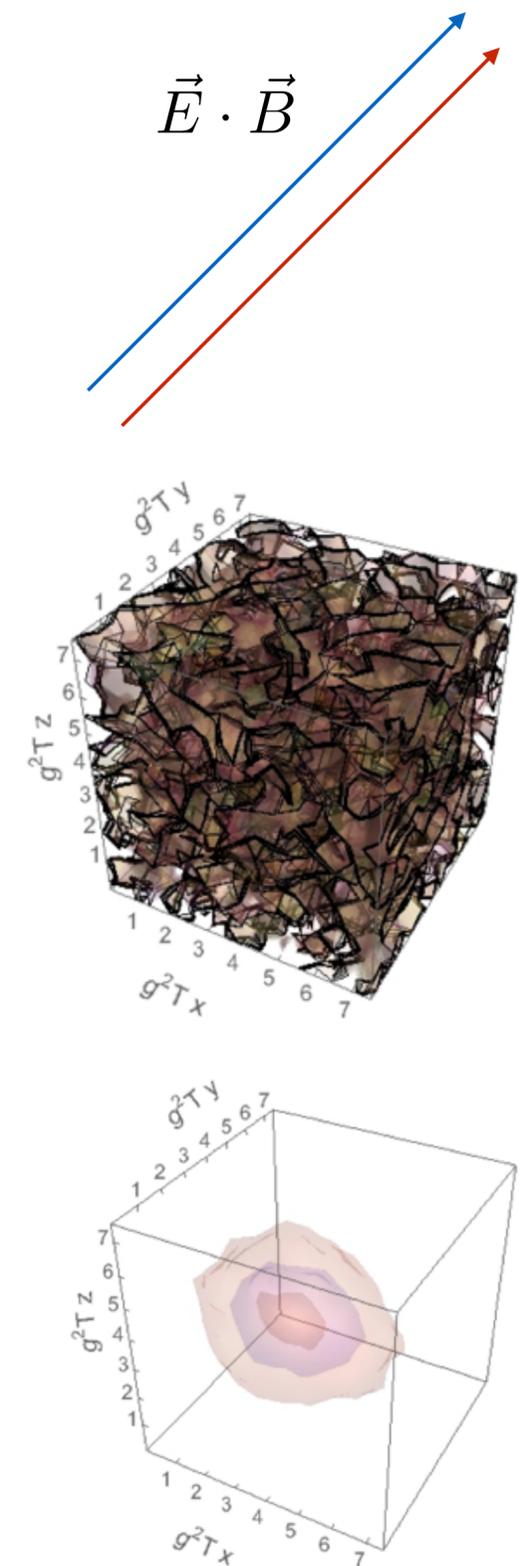
Akamatsu, Yamamoto PRL111 (2013) 052002

Crucial for understanding of chiral transport in HIC

Interesting applications in Astrophysics & Cosmology

=> Chirality imbalance through weak interaction processes

Masada et al. PRD98 (2018) no.8, 083018; Brandenburg et al. Astrophys.J. 845 (2017) no.2, L21



Mace, SS, Venugopalan
PRD93 (2016) no.7, 074036

Microscopic description for QED

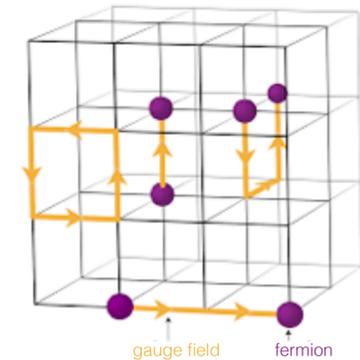
Classical-statistical lattice simulations of dynamics of underlying QFT

Solve operator Dirac equation
with QED gauge fields

$$i\gamma^0 \partial_t \hat{\psi} = (-i\not{D}_W^s + m)\hat{\psi}$$

Solve Maxwell's equations
with fermion currents

$$D_\mu F^{\mu\nu} = \langle \hat{\psi}(x) \gamma^\mu \hat{\psi}(x) \rangle$$

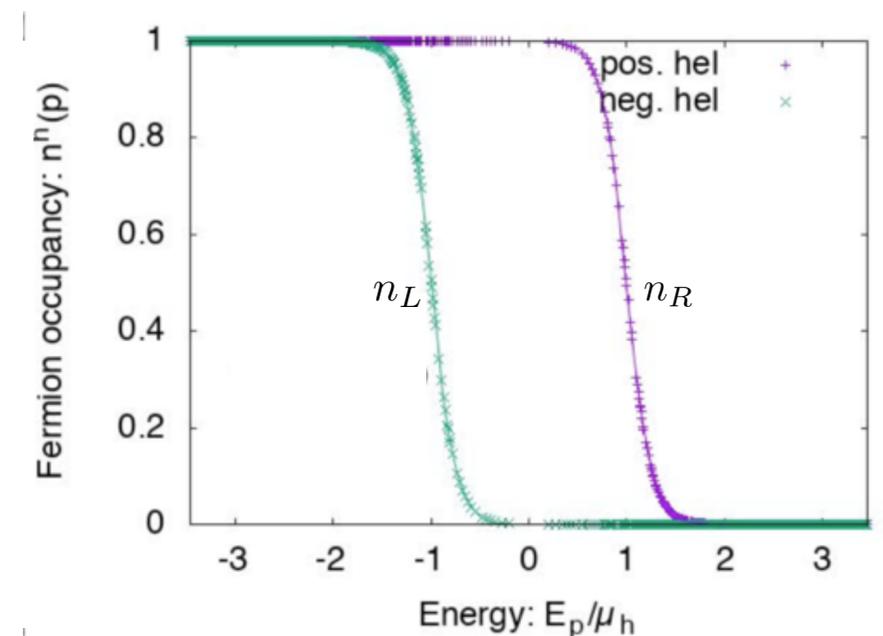


Start with spin polarization of the fermion sector, characterized by helicity
chemical potential $\mu_h \gg T$

=> energy density & chiral charge
initially concentrated in fermion sector

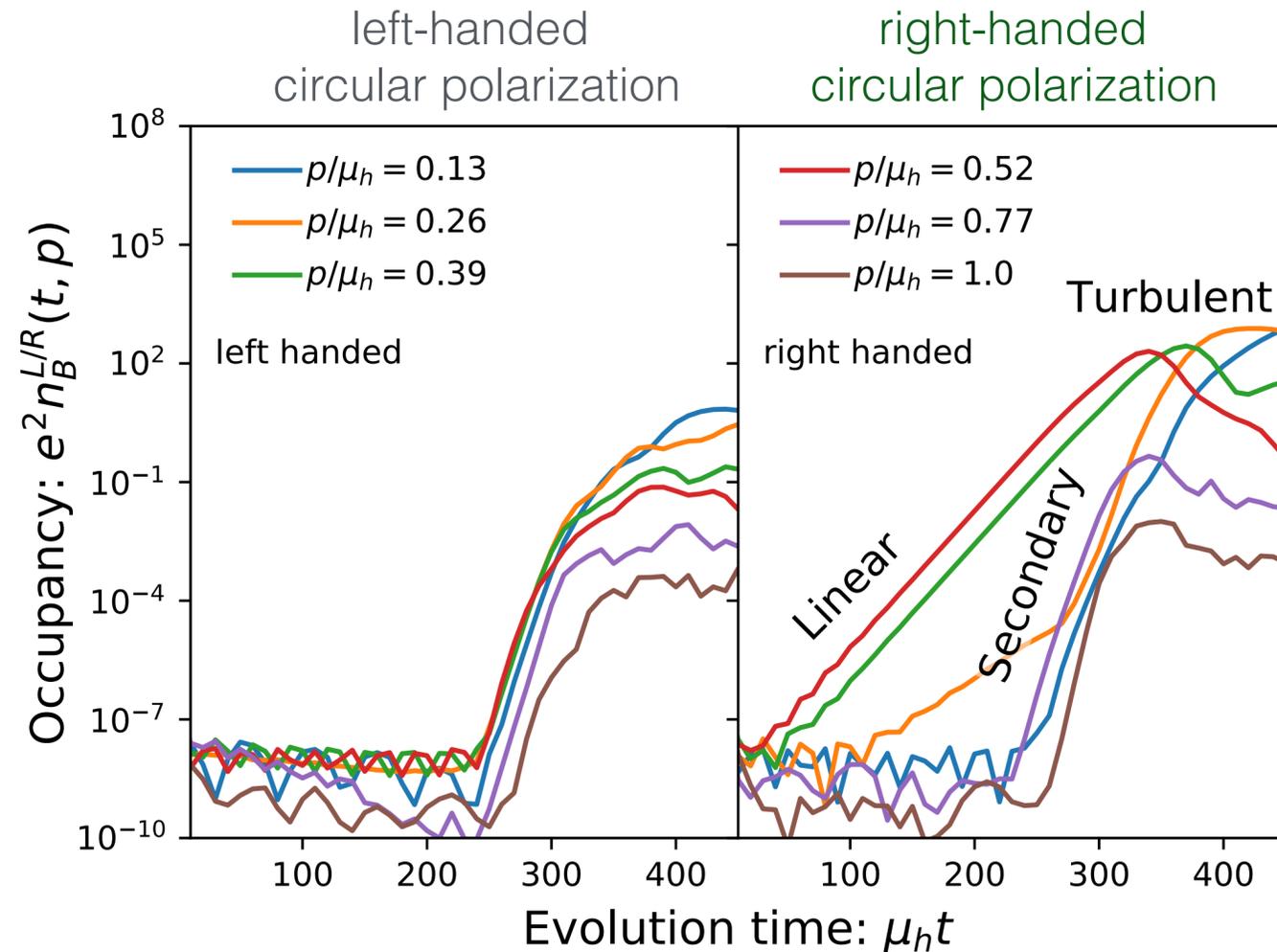
Simulate dynamics in **strongly coupled QED** $e^2 N_f = 64$ to properly resolve all
dynamical scales on the lattice

Simulations performed **close to chiral limit** $m \ll \mu_h$ where net-chirality is effectively conserved
trivial to extend to finite quark mass but not considered so far

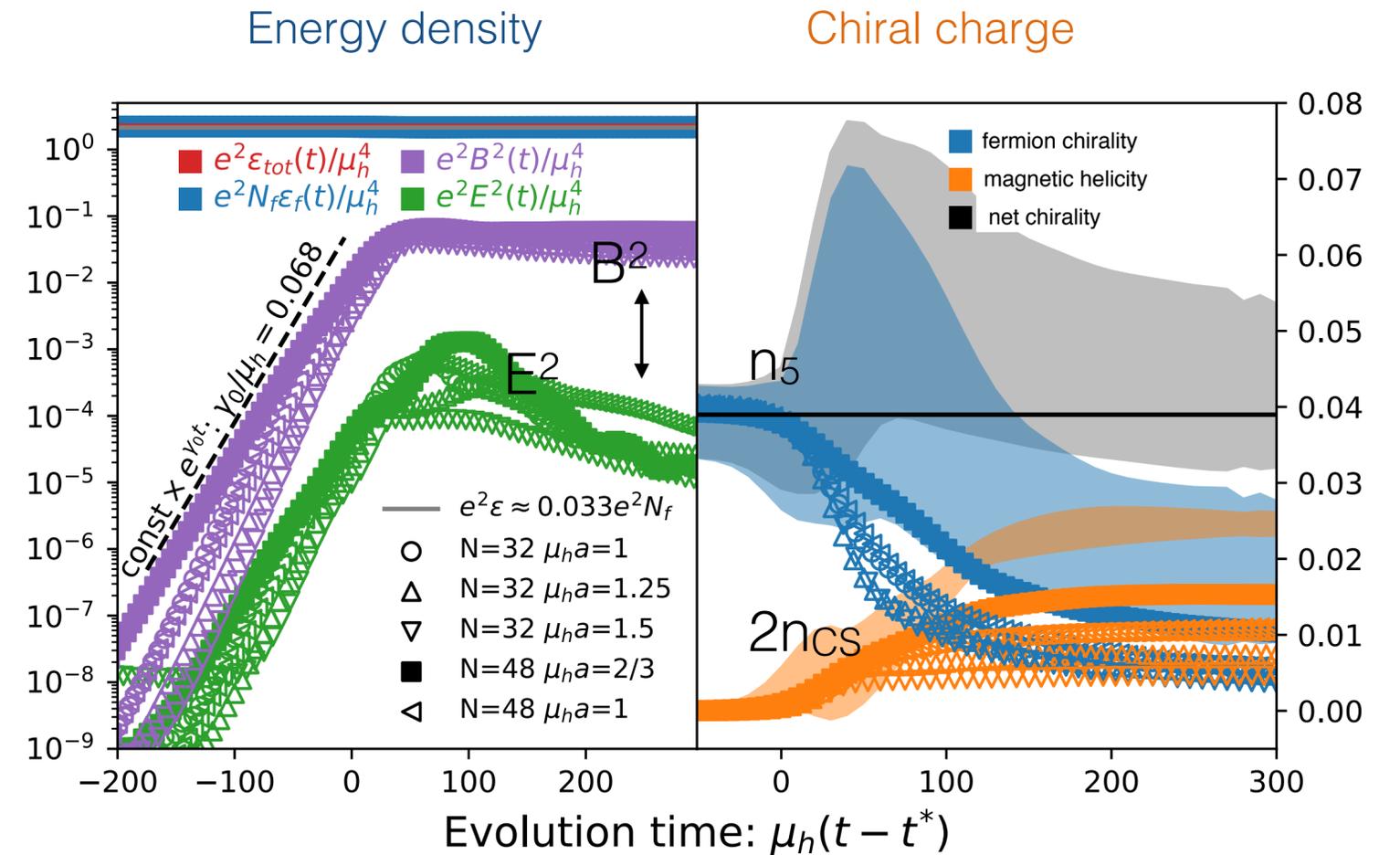


Chirality transfer in large N_f QED

Evolution of left/right-handed magnetic field modes



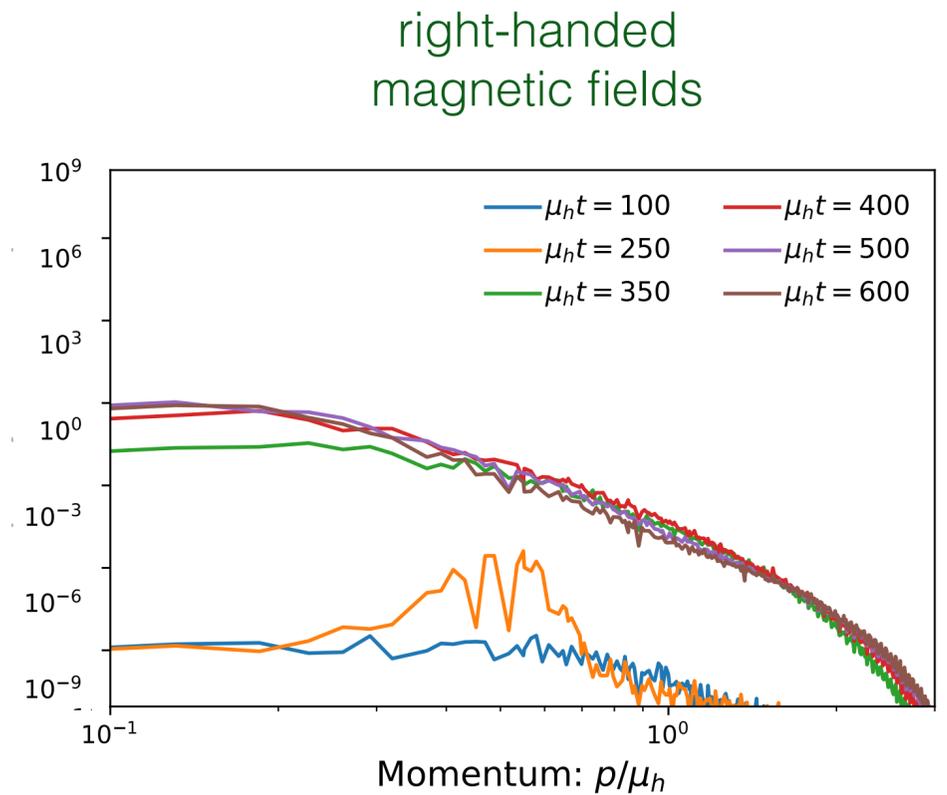
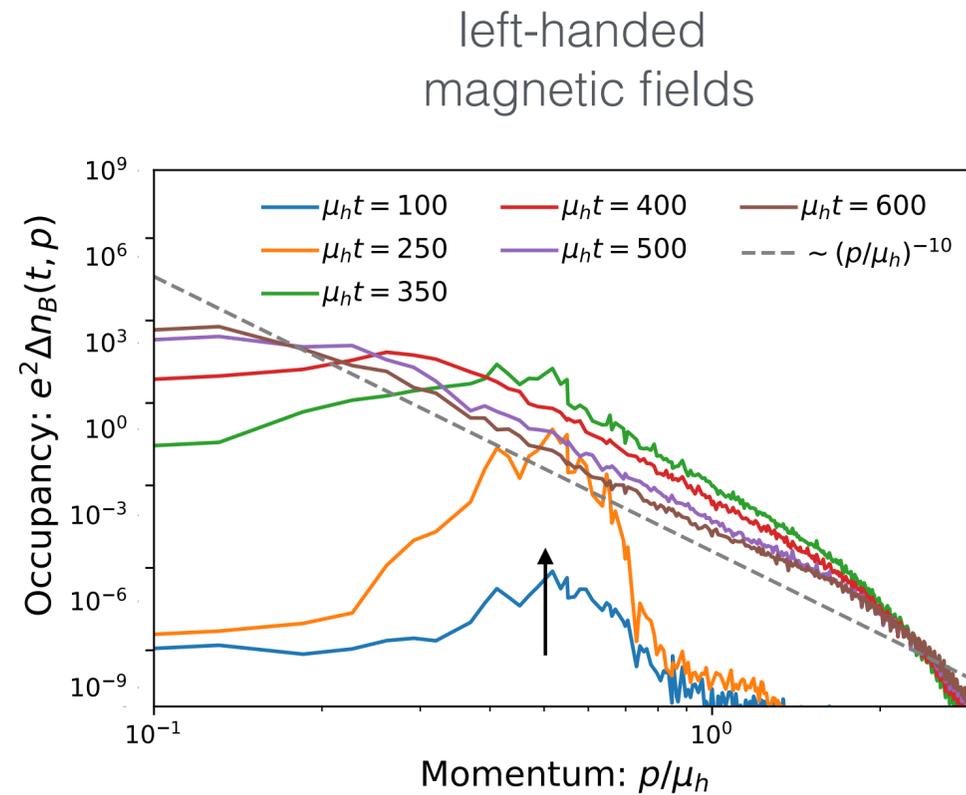
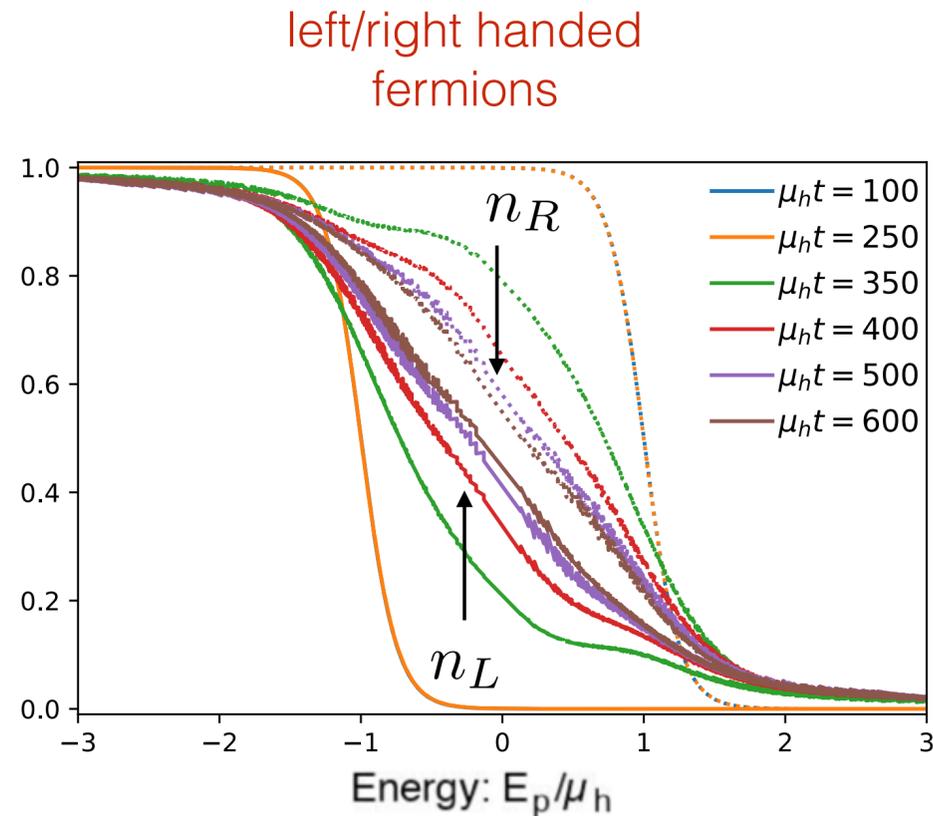
Evolution of conserved quantities



Chiral plasma instability of leads to a characteristic sequence of dynamics in unstable systems

Small fraction of energy but significant fraction of **chiral charge transferred from fermions to gauge fields**

Chirality transfer in large N_f QED



Non-linearities lead to **depletion of helicity imbalance** (stabilisation)

Strong **heating of the fermion sector** upon saturation of instability

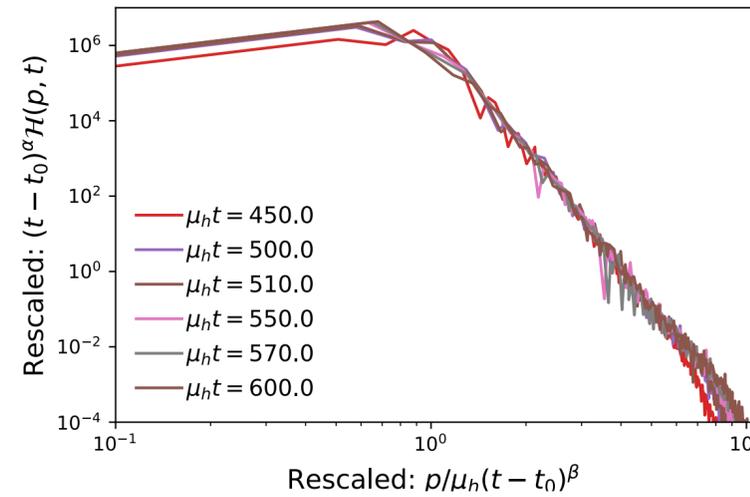
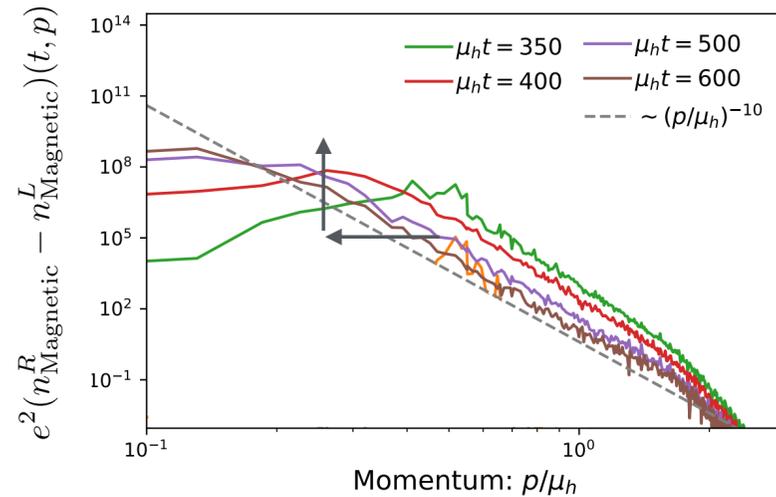
Non-perturbatively large occupations of left & right-handed magnetic fields

Strong IR excess of right vs. left-handed modes characterized by power law $\sim (p/\mu)^{-10}$ with large spectral exponent

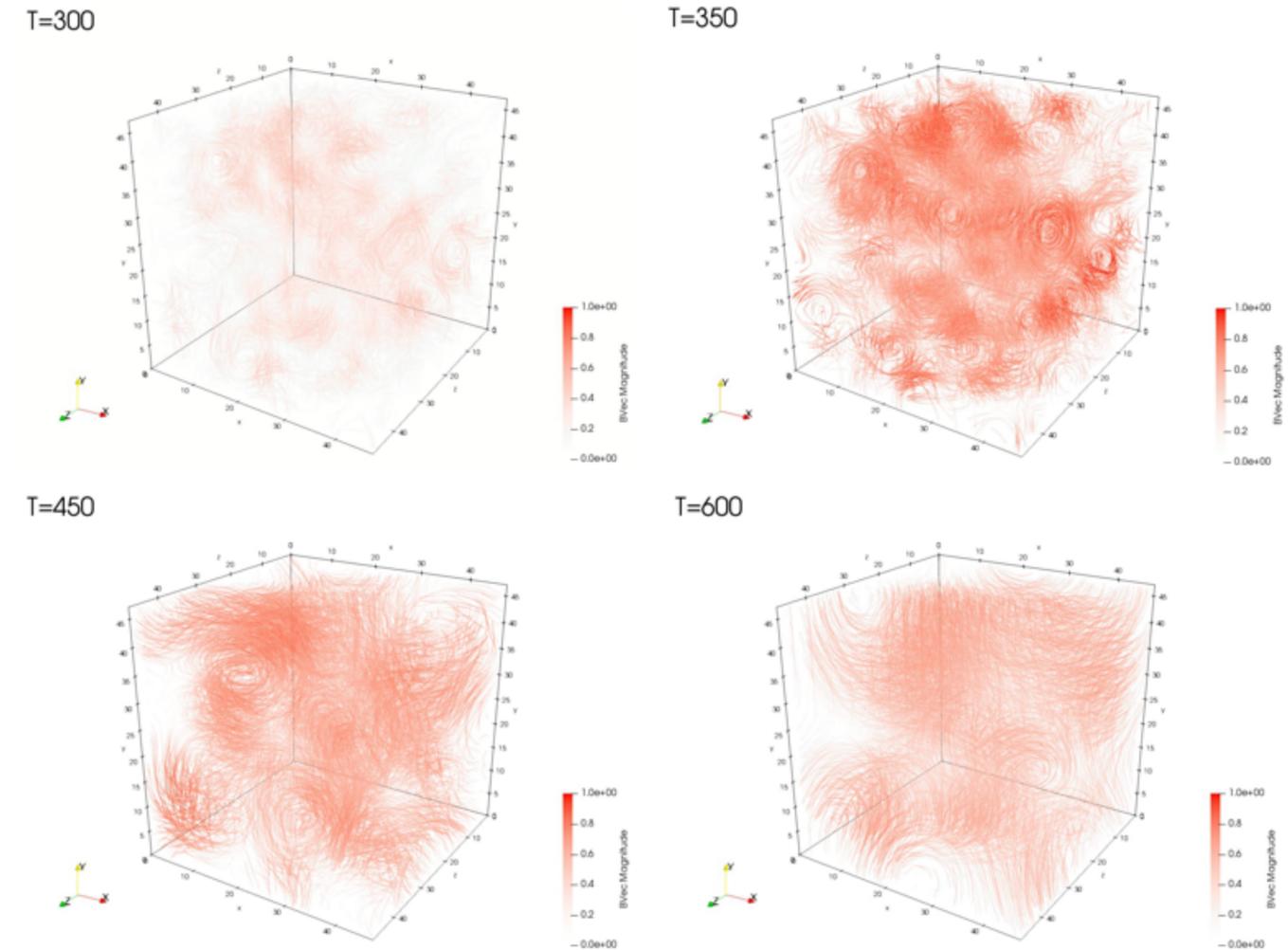
-> Chirality imbalance eventually absorbed by large scale helical magnetic fields

Chirality turbulence

Evolution of net-magnetic helicity spectra



Evolution of magnetic fields in coordinate space



Evolution in the turbulent regime proceeds via **self-similar evolution of net magnetic helicity**

Scaling exponents α, β determined from statistical scaling analysis

$$\Delta n_B(t, \mathbf{p}) = \tau^\alpha f_s(\tau^\beta |\mathbf{p}|)$$

$$\tau \equiv \mu_h(t - t^*)$$

$$\alpha = 1.14 \pm 0.50$$

$$\beta = 0.37 \pm 0.13$$

$$\alpha \approx 3\beta$$

Generation of large scale magnetic fields via self-similar inverse cascade of magnetic helicity

Conclusions & Outlook

Non-conservation of chiral charge complicates the macroscopic description

-> Important to understand chirality transfer between fermions and gauge fields

QED plasmas: Chiral plasma instability provides efficient mechanism to transfer chiral charge from fermions to gauge fields

-> Eventually chiral charge imbalance leads to generation of large scale helical magnetic fields via self-similar (inverse) cascade of magnetic helicity

Established for the first time from real-time simulations of underlying microscopic theory

QCD plasmas: Sphaleron transitions provide an additional mechanism to absorb chirality imbalance into topology of non-abelian gauge fields

Competition between effectively abelian and non-abelian effects requires further investigation

Need to take into account in macroscopic descriptions of anomalous transport phenomena

$$\partial_\mu j_a^\mu = -4\Gamma_{\text{sph}} \frac{\mu_A}{T}$$

SS, deBruin work in progress

Backup

Classical-statistical lattice simulations

Discretize theory on 3D spatial lattice using the Hamiltonian lattice formalism

$$i\gamma^0 \partial_t \hat{\psi} = (-i\mathcal{D}_W^s + m)\hat{\psi}$$

quantum fermions

$$D_\mu F^{\mu\nu} = \langle \hat{\psi}(x) \gamma^\mu \hat{\psi}(x) \rangle$$

classical gauge fields

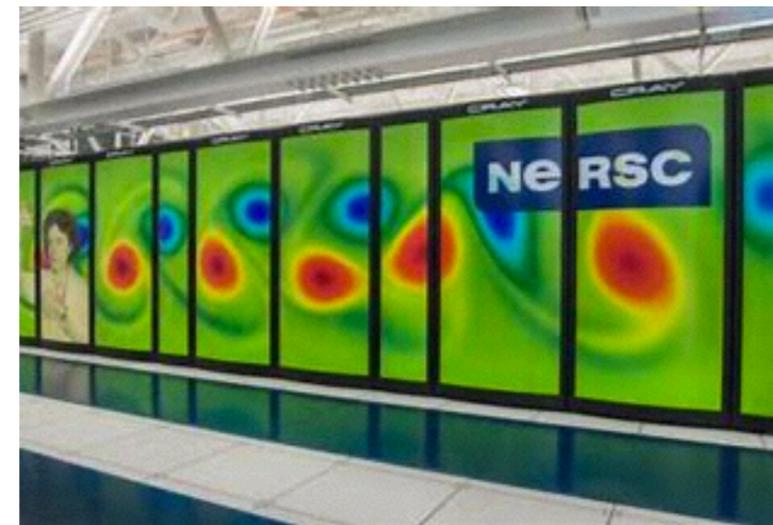
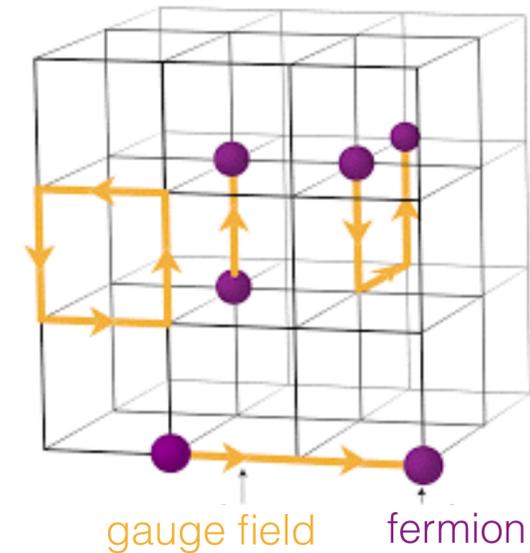
Expansion to leading order in e^2 but all orders in $e^2 N_f$ (exact at large N_f)

Solution to operator Dirac obtained by expanding the fermion field in operator basis at initial time

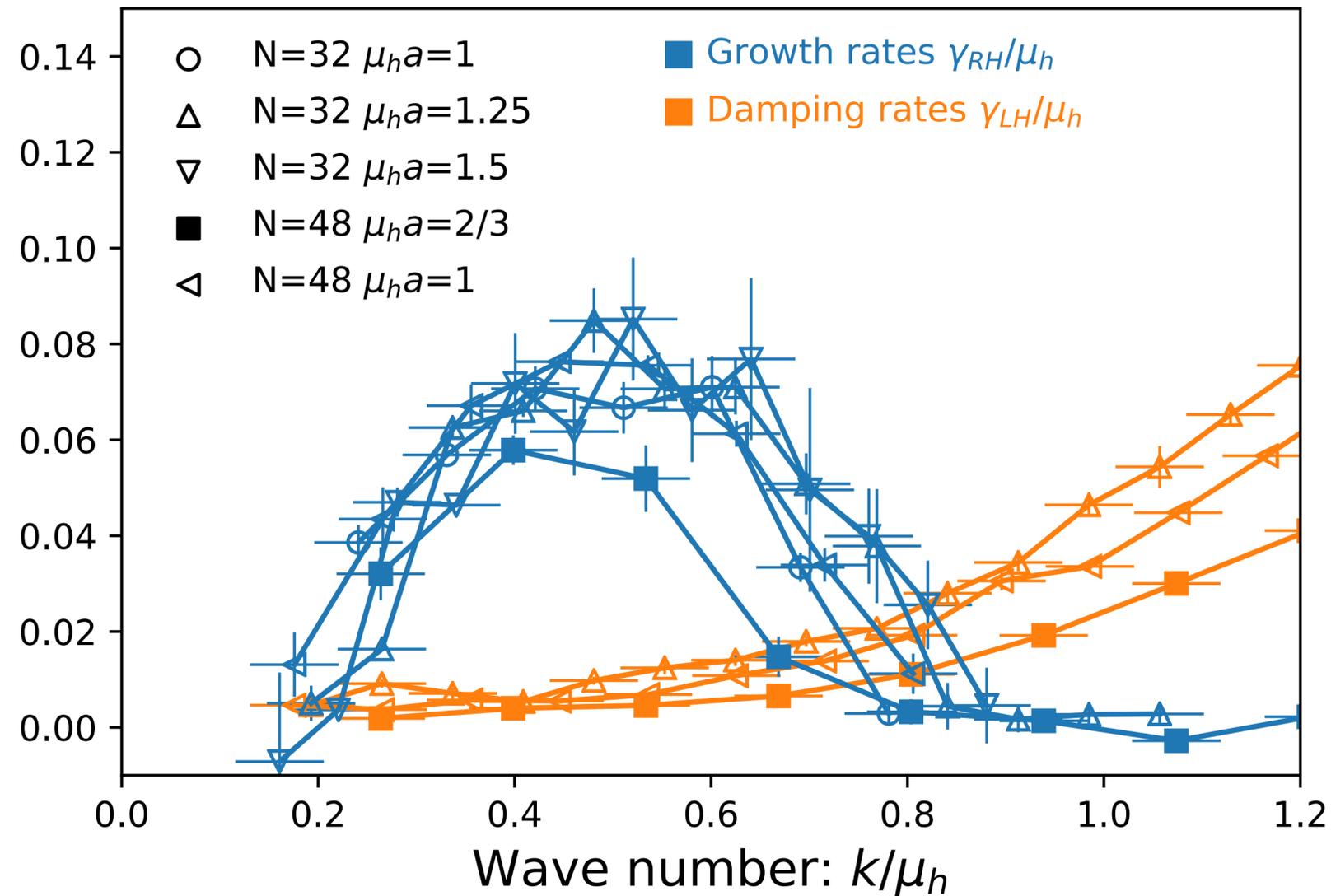
$$\hat{\psi}(x, t) = \sum_{p, \lambda} \hat{b}_{p, \lambda}(t=0) \phi_u^{p, \lambda}(x, t) + \hat{d}_{p, \lambda}^\dagger(t=0) \phi_v^{p, \lambda}(x, t)$$

and solving the Dirac equation for evolution of $4N_c N^3$ wave-functions

Computationally extremely demanding -> lattice sizes up to 48^3
(~5M CPU hours per simulation on Cori@NERSC)



Growth & Damping rates in large N_f QED

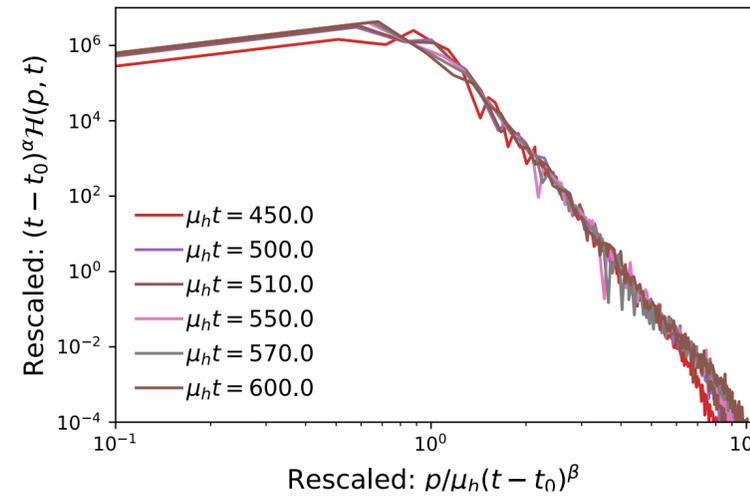
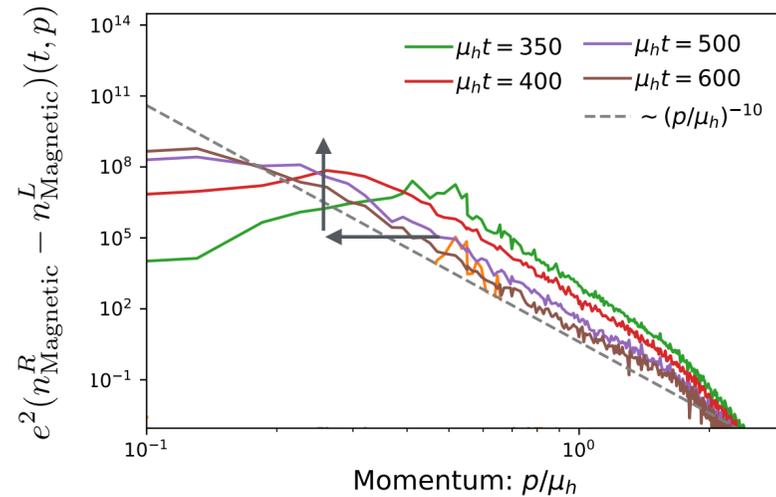


Growth and damping rates in qualitative agreement with expectation from one-loop HTL calculation

By use of improved operator definitions (Wilson NLO) errors due to finite lattice size and lattice spacing appear to be under control

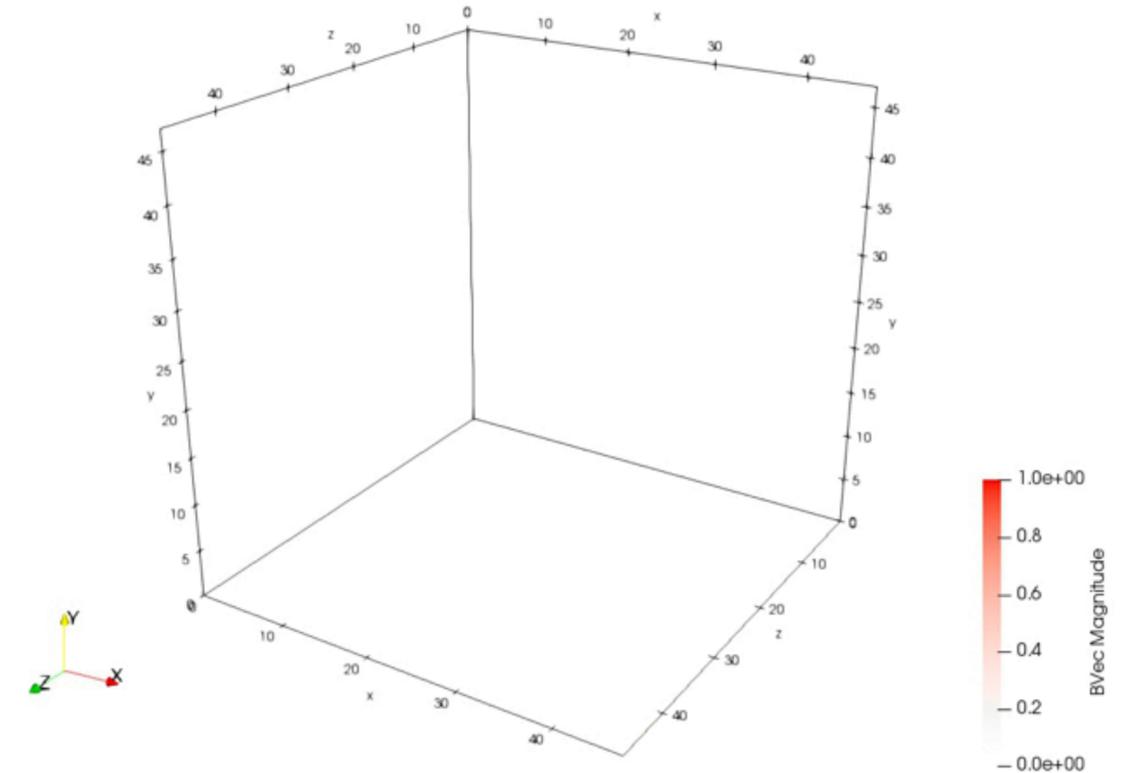
Chirality turbulence

Evolution of net-magnetic helicity spectra



Evolution of magnetic fields in coordinate space

T=200



Evolution in the turbulent regime proceeds via **self-similar evolution of net magnetic helicity**

Scaling exponents α, β determined from statistical scaling analysis

$$\Delta n_B(t, \mathbf{p}) = \tau^\alpha f_s(\tau^\beta |\mathbf{p}|)$$

$$\tau \equiv \mu_h(t - t^*)$$

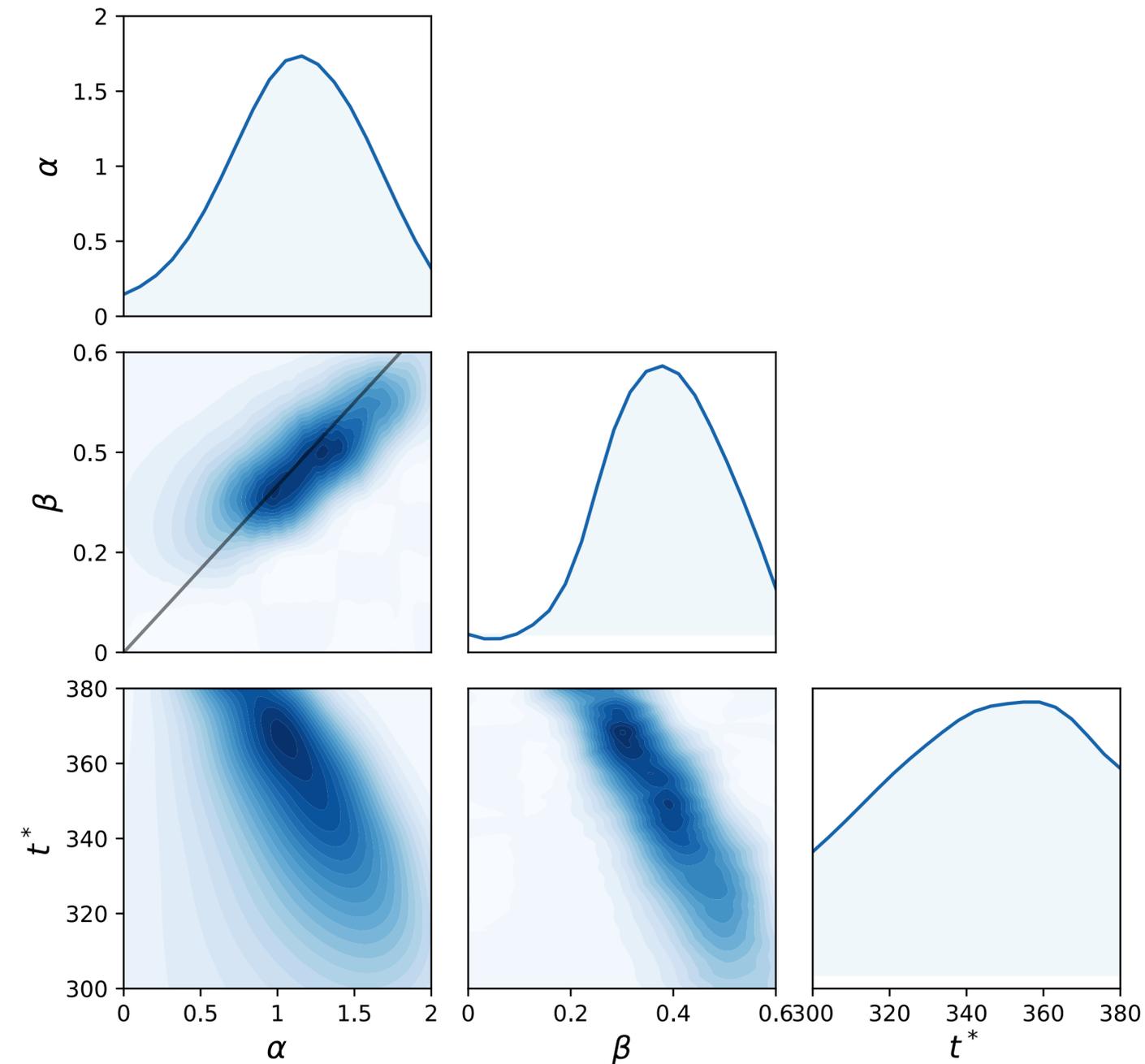
$$\alpha = 1.14 \pm 0.50$$

$$\beta = 0.37 \pm 0.13$$

$$\alpha \approx 3\beta$$

Generation of large scale magnetic fields via self-similar inverse cascade of magnetic helicity

Extraction of scaling exponents



Statistical scaling analysis

$$f_{\text{ref}}(t, \tilde{p} = \tau_{\text{ref}}^\beta p) = \log \tau_{\text{ref}}^{-\alpha} \Delta n_B(\tau_{\text{ref}}^\beta p),$$

$$\chi^2(\alpha, \beta, t^*) = \frac{1}{N_{\text{test}}} \sum_{\tau_{\text{test}}} \frac{\int \frac{d\tilde{p}}{\tilde{p}} (f_{\text{ref}}(\tilde{p}) - f_{\text{test}}(\tilde{p}))^2}{\int \frac{d\tilde{p}}{\tilde{p}} (f_{\text{ref}}(\tilde{p}))^2},$$

determines likelihood

$$W(\alpha, \beta, t^*) = \frac{1}{\mathcal{N}} \exp \left[-\frac{\chi^2(\alpha, \beta, t^*)}{\chi_{\text{min}}^2} \right]$$