

Continuity from neutron matter to color-superconducting quark matter with 3P_2 superfluidity

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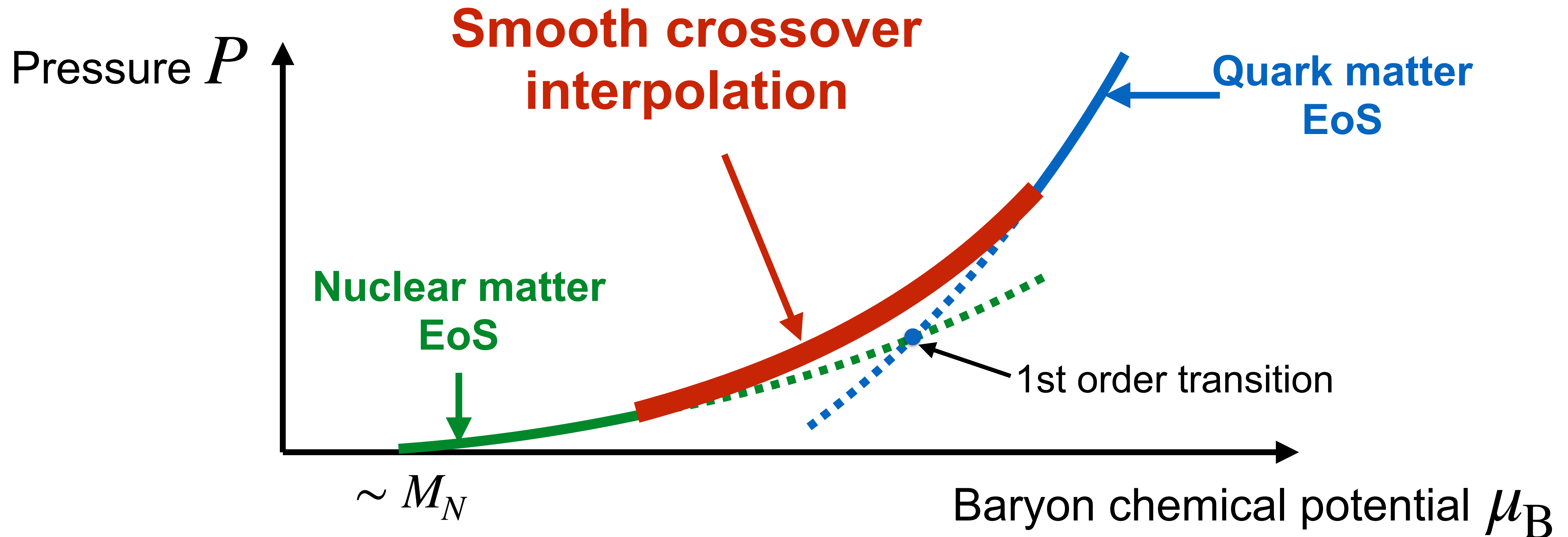
based on: **YF**, K. Fukushima, W. Weise, arXiv:1908.09360 [hep-ph]

6 Nov. 2019, Quark Matter 2019 @ Wuhan

Prelude: Dense Matter Equation of State (EoS)

Masuda, Hatsuda, Takatsuka (2012); Kojo, Powell, Song, Baym (2014)

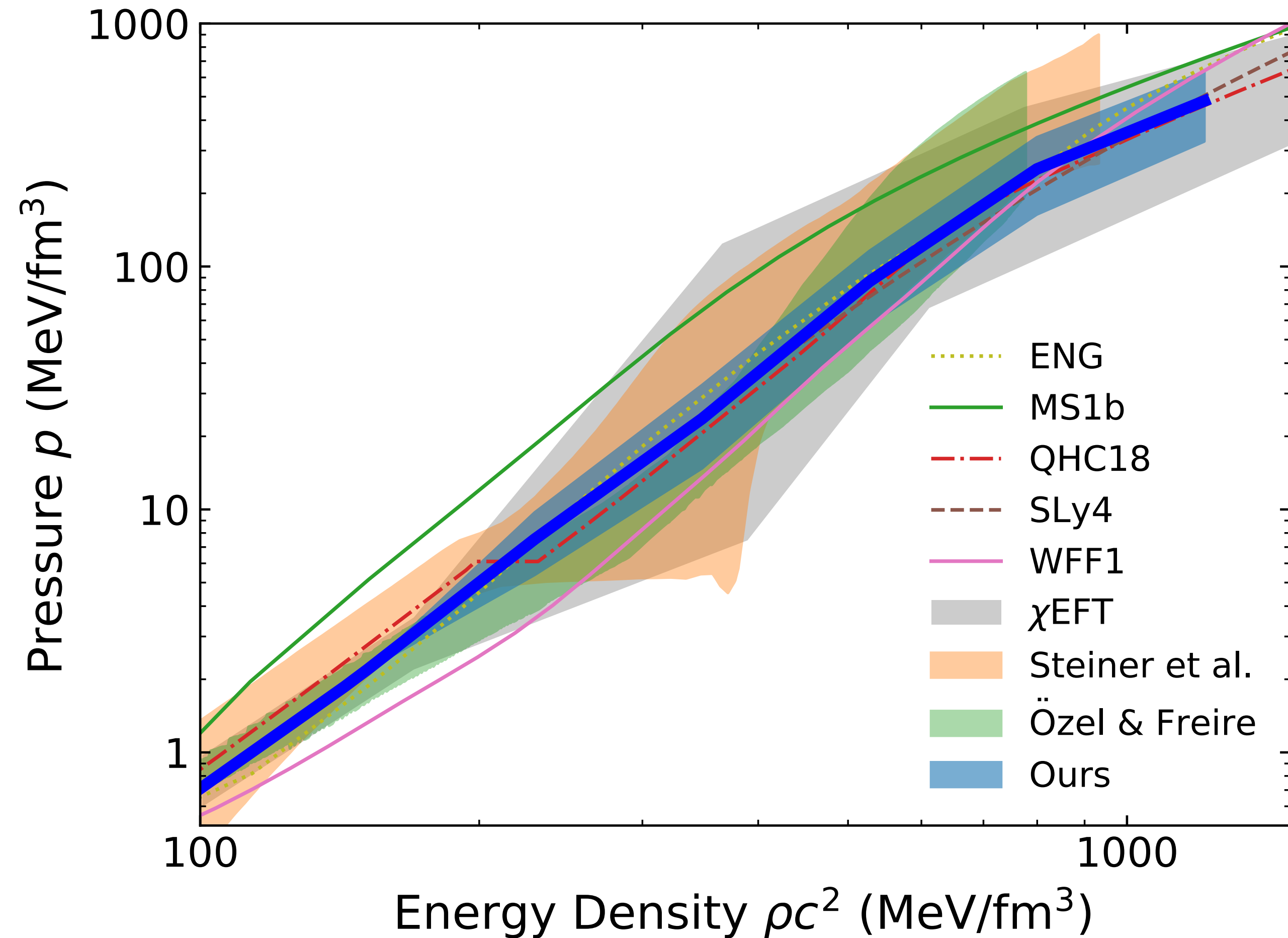
Neutron star phenomenology implies...



Prelude: Dense Matter Equation of State (EoS)

Fujimoto, Fukushima, Murase, arXiv:1903.03400

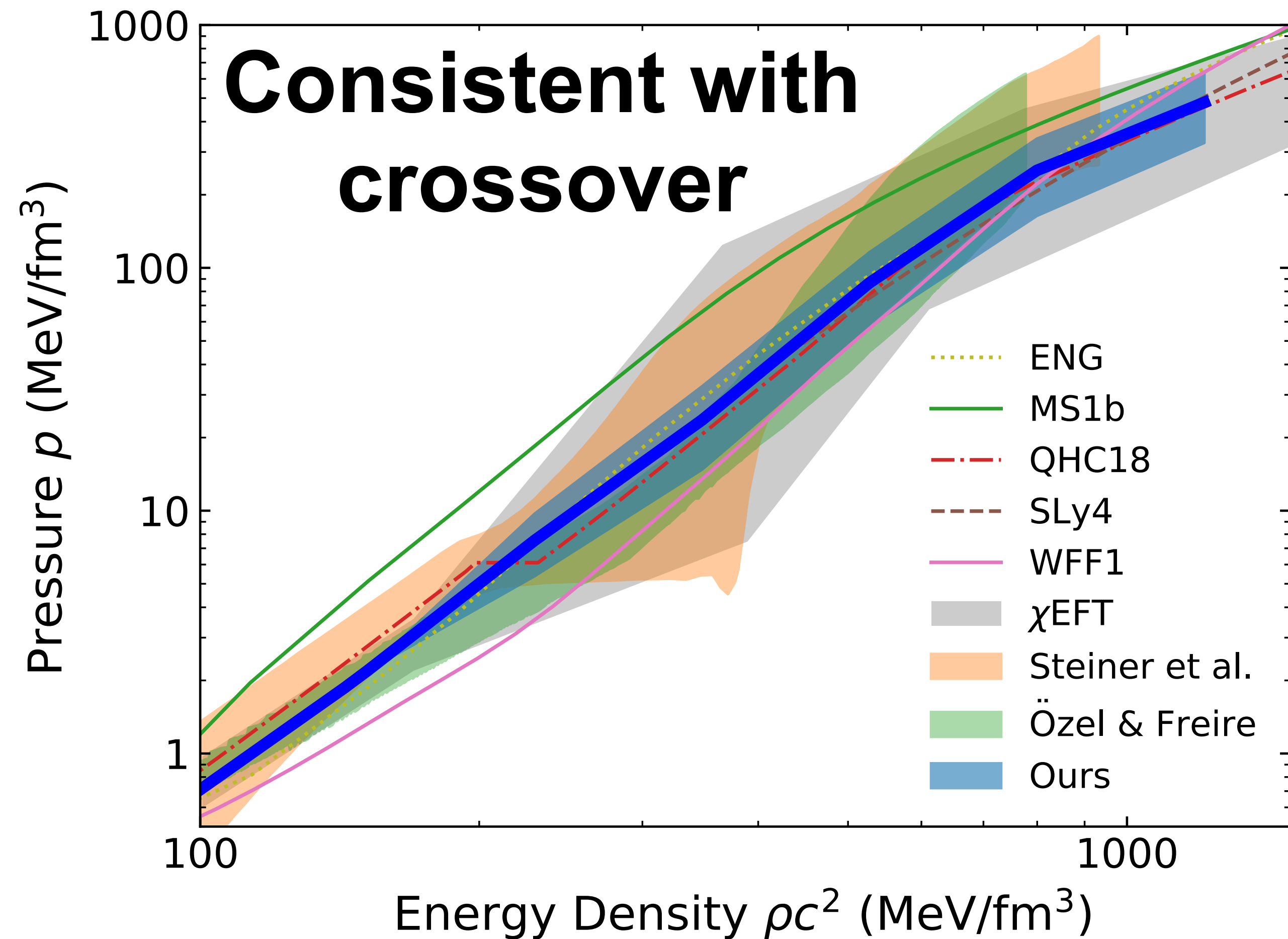
Deep learning: NS observations \rightarrow EoS



Prelude: Dense Matter Equation of State (EoS)

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Deep learning: NS observations \rightarrow EoS



Theoretical Background of Crossover EoS

Schafer, Wilczek (1998)

Alford, Berges, Rajagopal (1999)

Quark-Hadron Continuity (3-flavor)

= Both phases are indistinguishable

i.e. **(1) symmetries are the same;**

(2) no order parameter

... Assumes **3-flavor symmetric** matter $m_u = m_d = m_s$

- In reality, nuclear matter is **2-flavor symmetric** rather than three:

$$m_s \gg m_u \simeq m_d$$

- Most nuclear EoS are for 2-flavor matter \rightarrow **2-flavor matter continuity?**

No Quark-Hadron Continuity in Usual 2-flavor

Schafer, Wilczek (1998)

Alford, Berges, Rajagopal (1999)

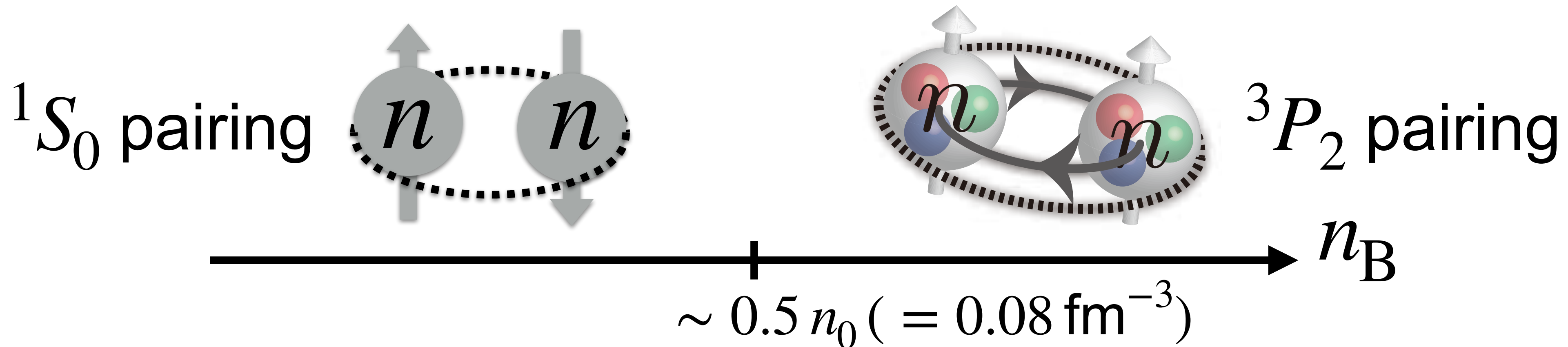
Quark-Hadron Continuity?? (2-flavor)

Conventional wisdom tells us...

- In quark phase, 2-flavor color-superconductivity (**2SC**) has: $\langle ud \rangle_{L/R}$

→ **NO $U(1)_B$ breaking** (i.e. no superfluidity)

- In hadron phase, neutron superfluid → **$U(1)_B$ breaking**



No Quark-Hadron Continuity in Usual 2-flavor

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Quark-Hadron Continuity?? (2-flavor)

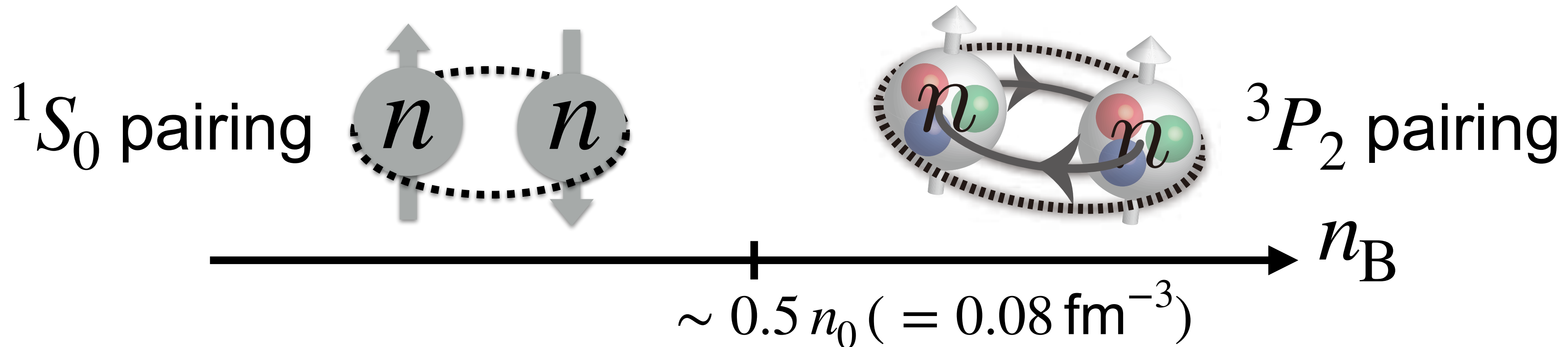
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No continuity?

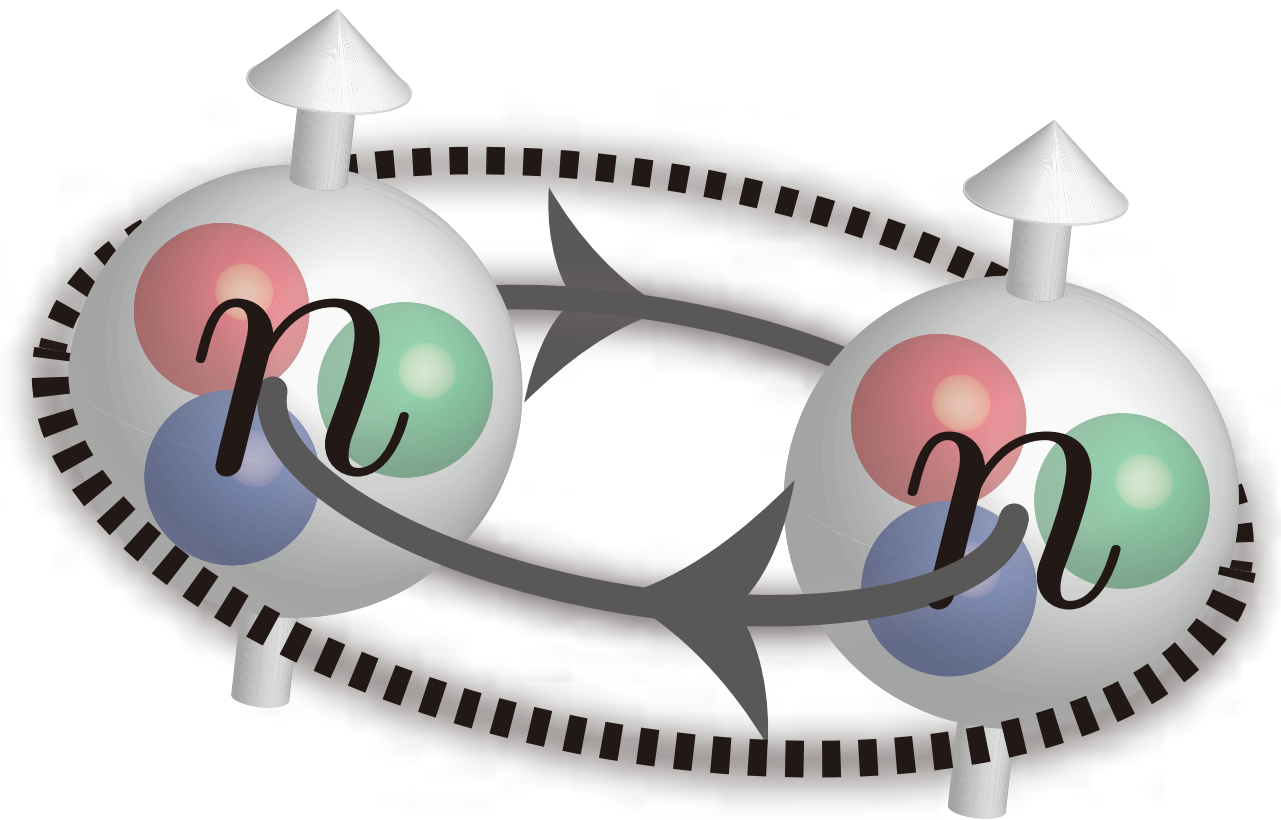
- In **hadron phase**, neutron superfluid → **$U(1)_B$ breaking**



What We Found Here: 2-flavor Continuity

Fujimoto, Fukushima, Weise, arXiv:1908.09360

Nuclear matter

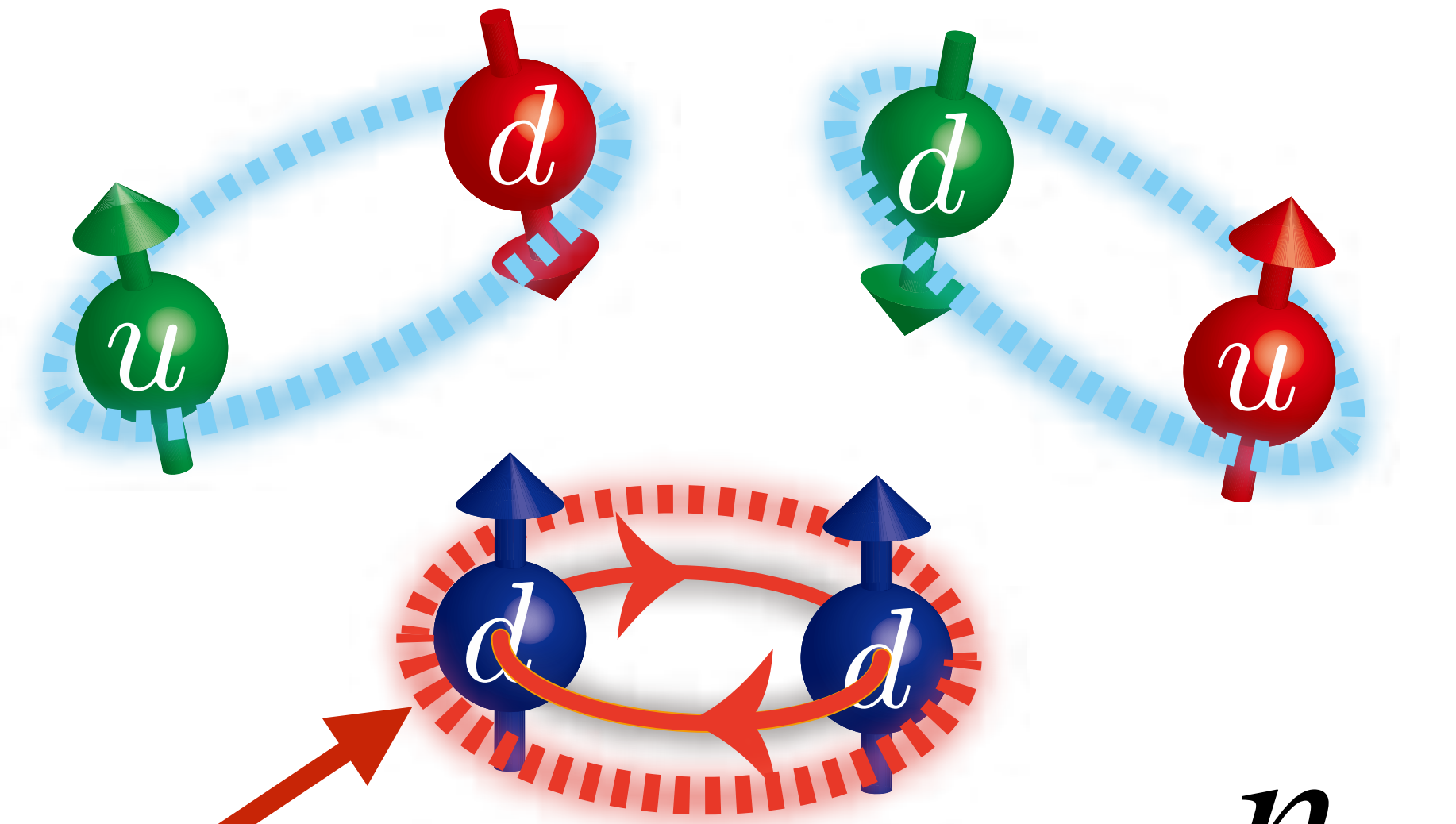


$\sim n_0 (= 0.16 \text{ fm}^{-3})$

Continuity



Quark matter




$\sim 10 n_0$

n_B

**New element here; ensures continuity
“2SC + <dd>”**

[Review] Quark-Hadron Continuity: 3-flavor

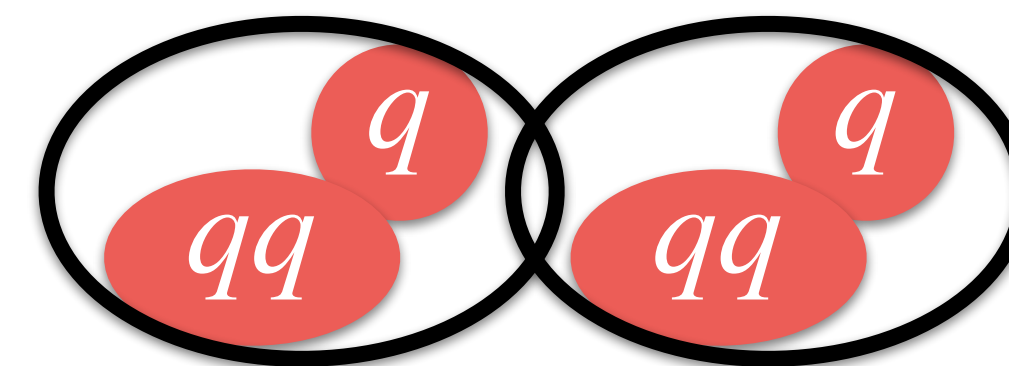
Alford, Baym, Fukushima, Hatsuda, Tachibana (2018)

- Here, define CFL diquark operator:  (gauge variant; color $\bar{3}$)
- **Gauge invariant** order parameters with the same quark constituent:

in the CFL phase



in the hadronic phase



- Compare the expectation value (in the mean field approx.):

$$H \left\langle \left(\text{CFL diquarks} \right) \right\rangle \approx \left\langle \left(\text{hadronic baryons} \right) \right\rangle \Lambda \Lambda \quad \begin{array}{l} \text{3 diquarks} \\ \sim 2 \text{ baryons} \end{array}$$

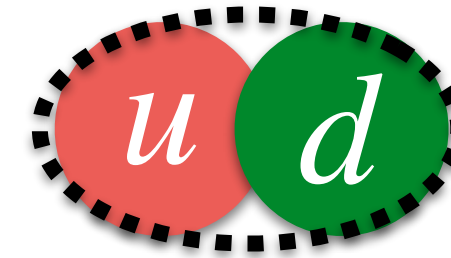
... **No gauge-invariant order parameter** to distinguish both phases

Quark-Hadron Continuity: 2-flavor Case

Fujimoto, Fukushima, Weise, arXiv:1908.09360

- Doing just the same as in 3-flavor case...

- We already know 2SC diquark operator:

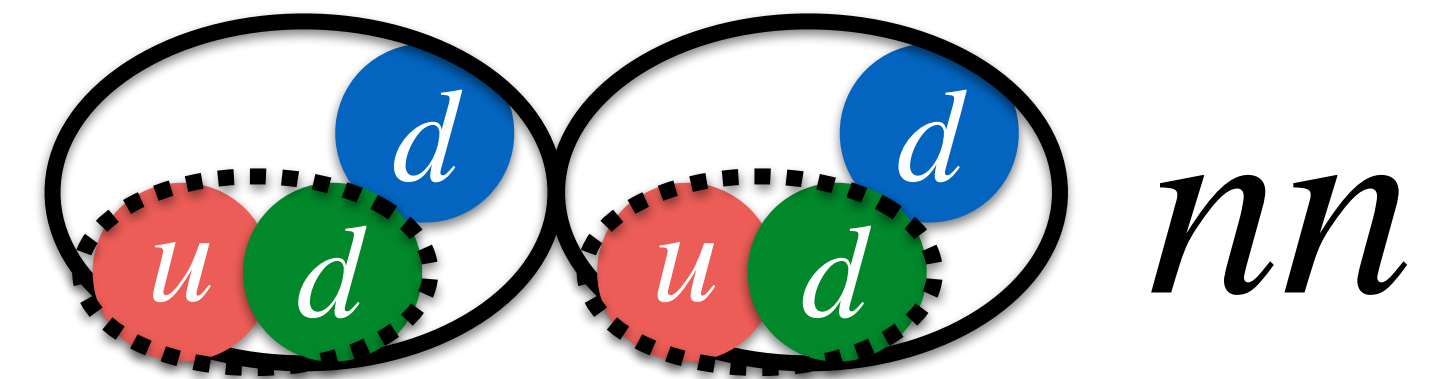


- **Gauge invariant** order parameters with the same quark constituent:

in the quark phase

?

in the hadronic phase



- Find the quark phase counterpart by rearrangement (in the mean field approx.):

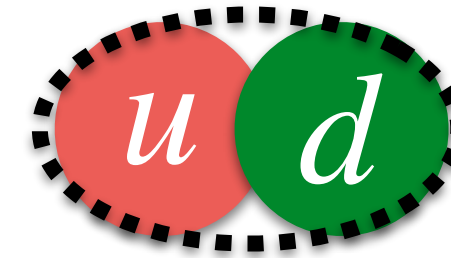
$$nn \left\langle \left(\text{diagram of two nucleons} \right) \right\rangle \approx \left\langle \text{diagram of two diquarks} \right\rangle \left\langle \text{diagram of two diquarks} \right\rangle \left\langle \text{diagram of two diquarks} \right\rangle 2SC + \langle dd \rangle$$

Quark-Hadron Continuity: 2-flavor Case

Fujimoto, Fukushima, Weise, arXiv:1908.09360

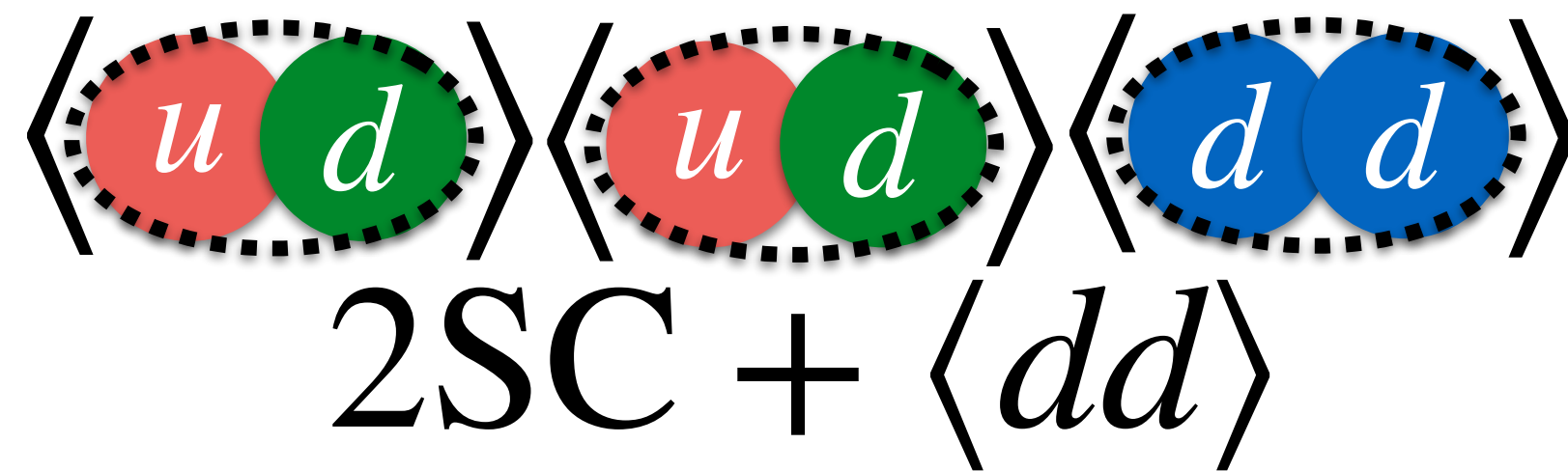
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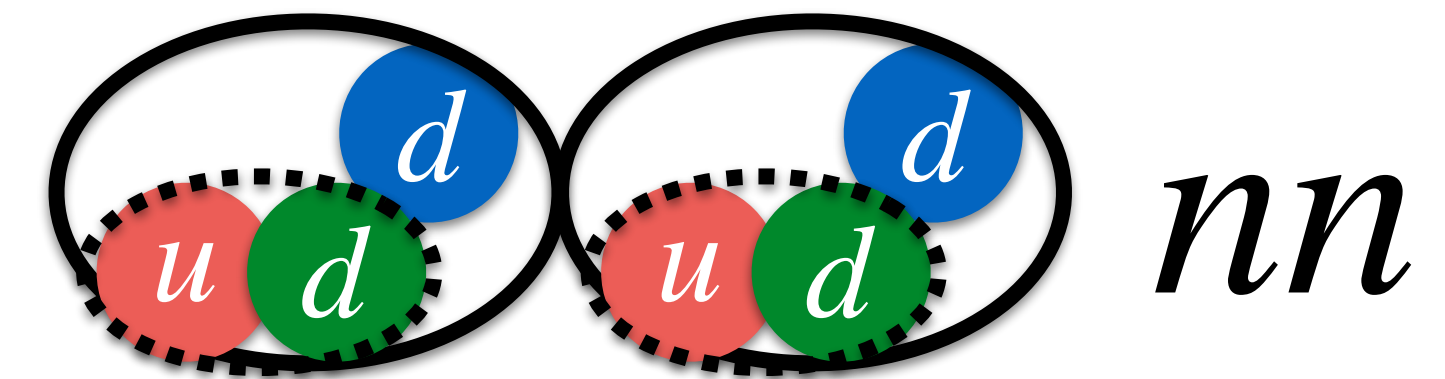


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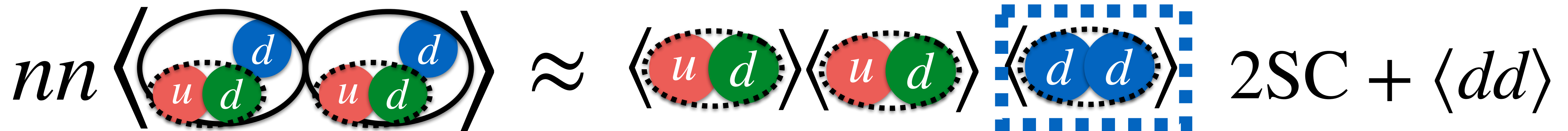
in the quark phase



in the hadronic phase



- Find the quark phase counterpart by rearrangement (in the mean field approx.):



Quark-Hadron Continuity: 2-flavor Case

Fujimoto, Fukushima, Weise, arXiv:1908.09360

- We have **2SC+<dd>**:

$$\langle \mathbf{ud} \rangle_{L/R} + \langle \hat{\mathbf{d}}^T C \gamma^i \nabla^j \hat{\mathbf{d}} \rangle$$

... <dd> breaks $U(1)_B$

→ matter is superfluid, the same as hadronic phase → **Continuity!**

- $\langle \hat{d}_\alpha^T C \gamma^i \nabla^j \hat{d}_\beta \rangle \rightarrow$ Color should be **symmetric** (in sextet channel **6**)

One-gluon exchange (i.e. short-range) between dd **repulsive!?**

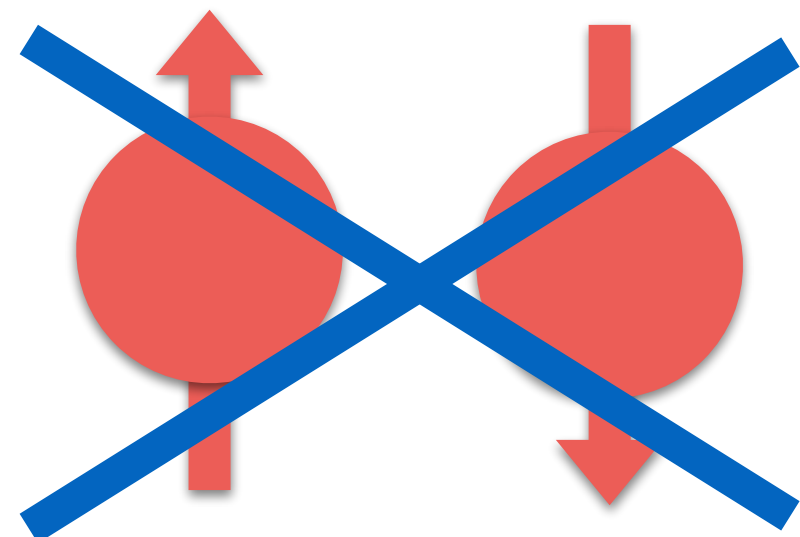
→ **No Cooper pair formation?**

Repulsive interaction between dd ?

Fujimoto, Fukushima, Weise, arXiv:1908.09360

- Repulsive nature is not harmful, rather **useful**...

S -wave Cooper pairing ($L = 0$)

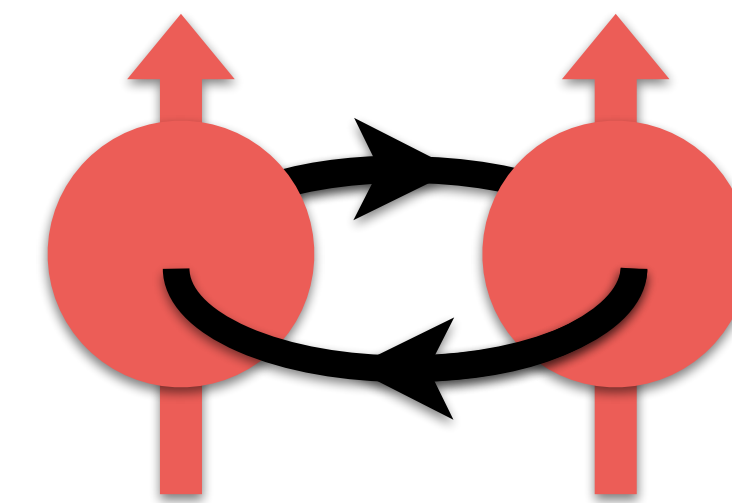


suppression



Short-range
repulsive force

P -wave pairing ($L = 1$)



Higher partial-wave state
favored!

- Due to $\langle dd \rangle$ structure: $\langle \mathbf{d}^T C \gamma^i \nabla^j \hat{\mathbf{d}} \rangle$ spin & angular momentum $\neq 0$

- Prototype example: **neutron 3P_2 pairing**

3P_2 Neutron Superfluidity

Hoffberg *et al.* (1970); Tamagaki (1970)

- Inside neutron star, 3P_2 **neutron superfluid** is presumed

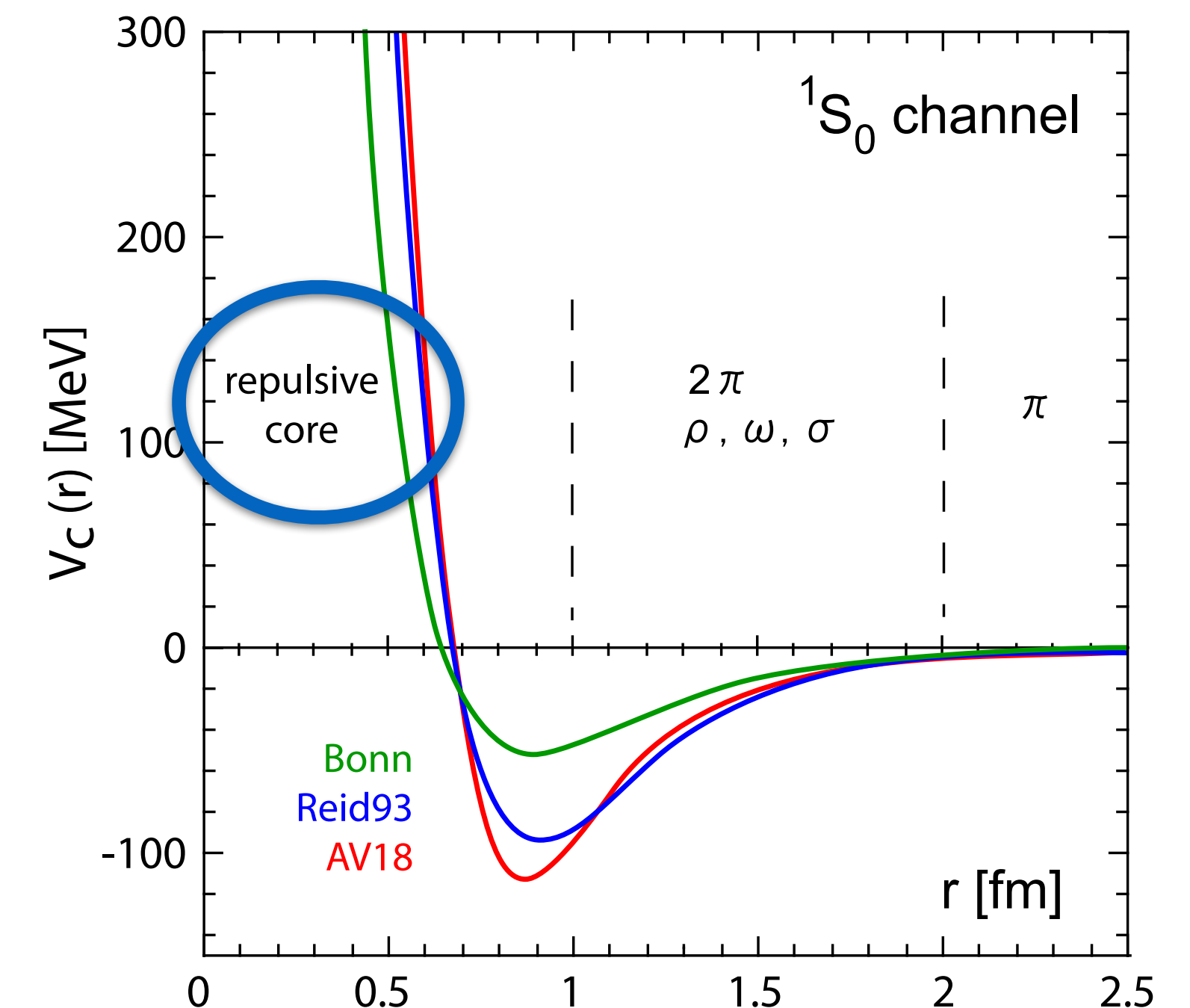
- Increasing densities, NN interaction becomes repulsive due to the **repulsive core**

→ - S -wave ($L = 0$) superfluidity suppressed

- P -wave ($L = 1$) takes over

- **Spin triplet** ($S = 1$): fixed by anti-symmetry

Aoki (2011)



3P_2 Neutron Superfluidity

- $J = 2$ is favored among $J = 0, 1, 2$

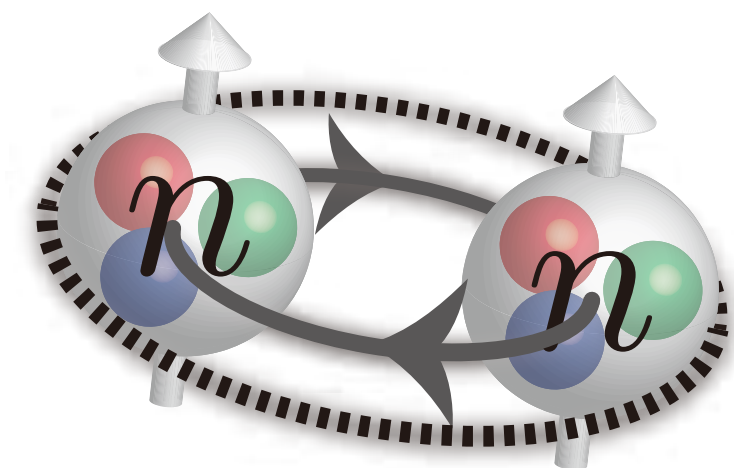
by attractive LS potential V_{LS} :

$$V_{LS} \propto -\mathbf{L} \cdot \mathbf{S} = \begin{cases} +2 & (J = 0) \text{ Repulsive} \\ +1 & (J = 1) \text{ Repulsive} \\ -1 & (J = 2) \text{ Attractive} \end{cases}$$

- Phase shift (effective NN interaction):

3P_2 channel is dominantly attractive above ρ_0

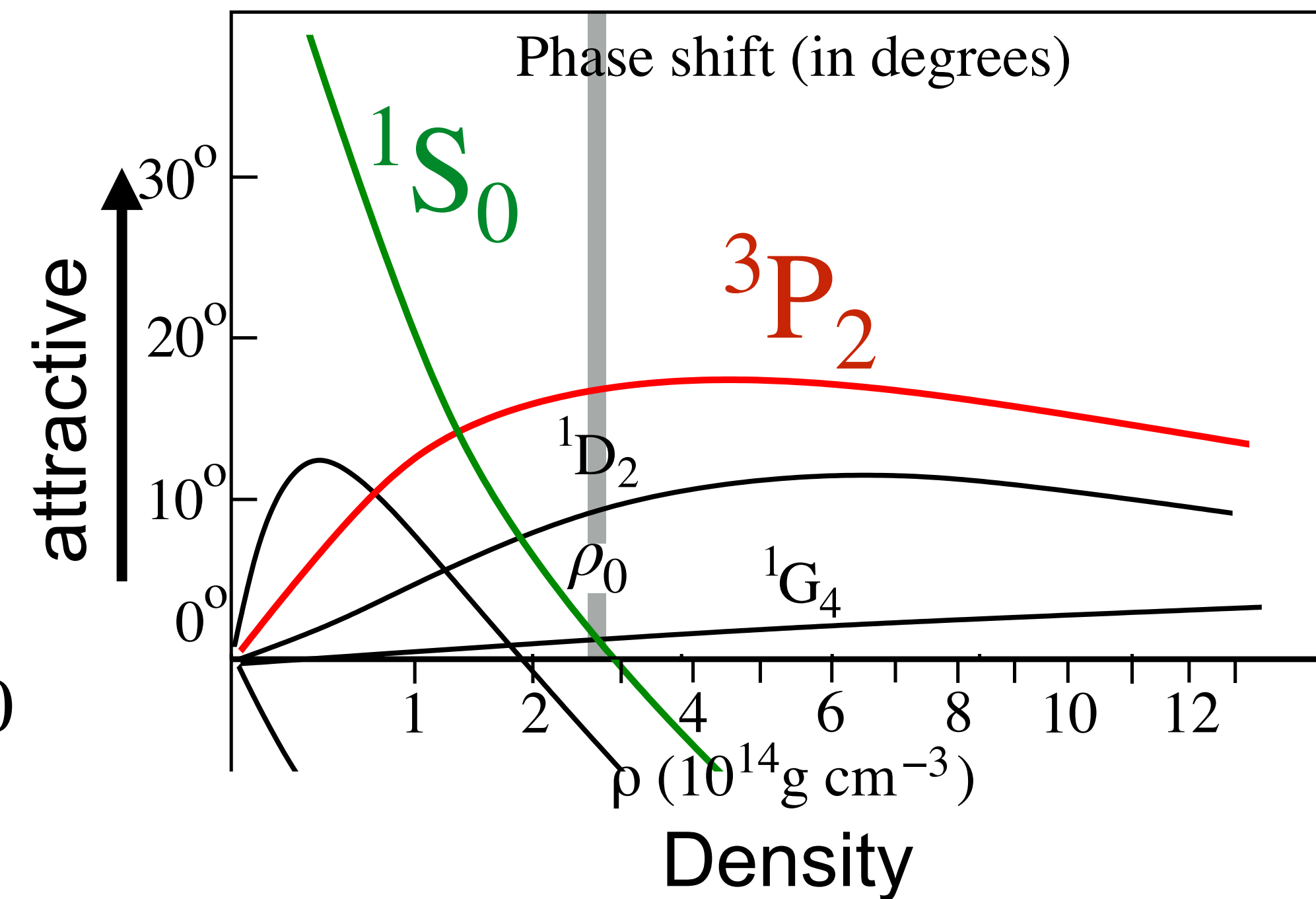
→ 3P_2 superfluidity



Hoffberg *et al.* (1970); Tamagaki (1970)

Page, Prakash, Lattimer, Steiner (2011)

NN Scattering



3P_2 Neutron Superfluidity

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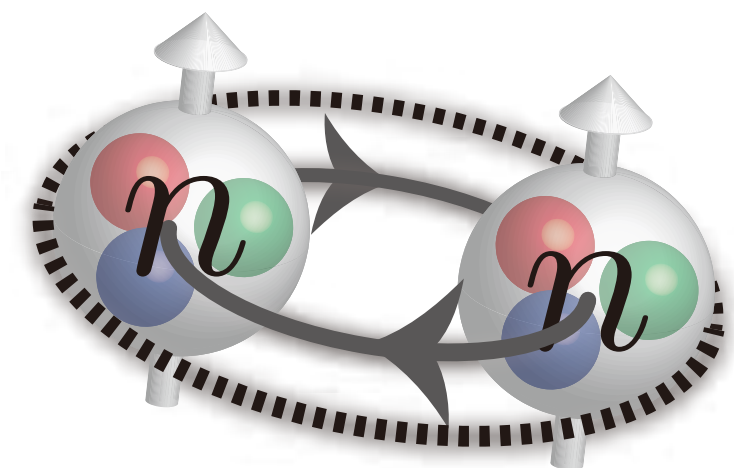
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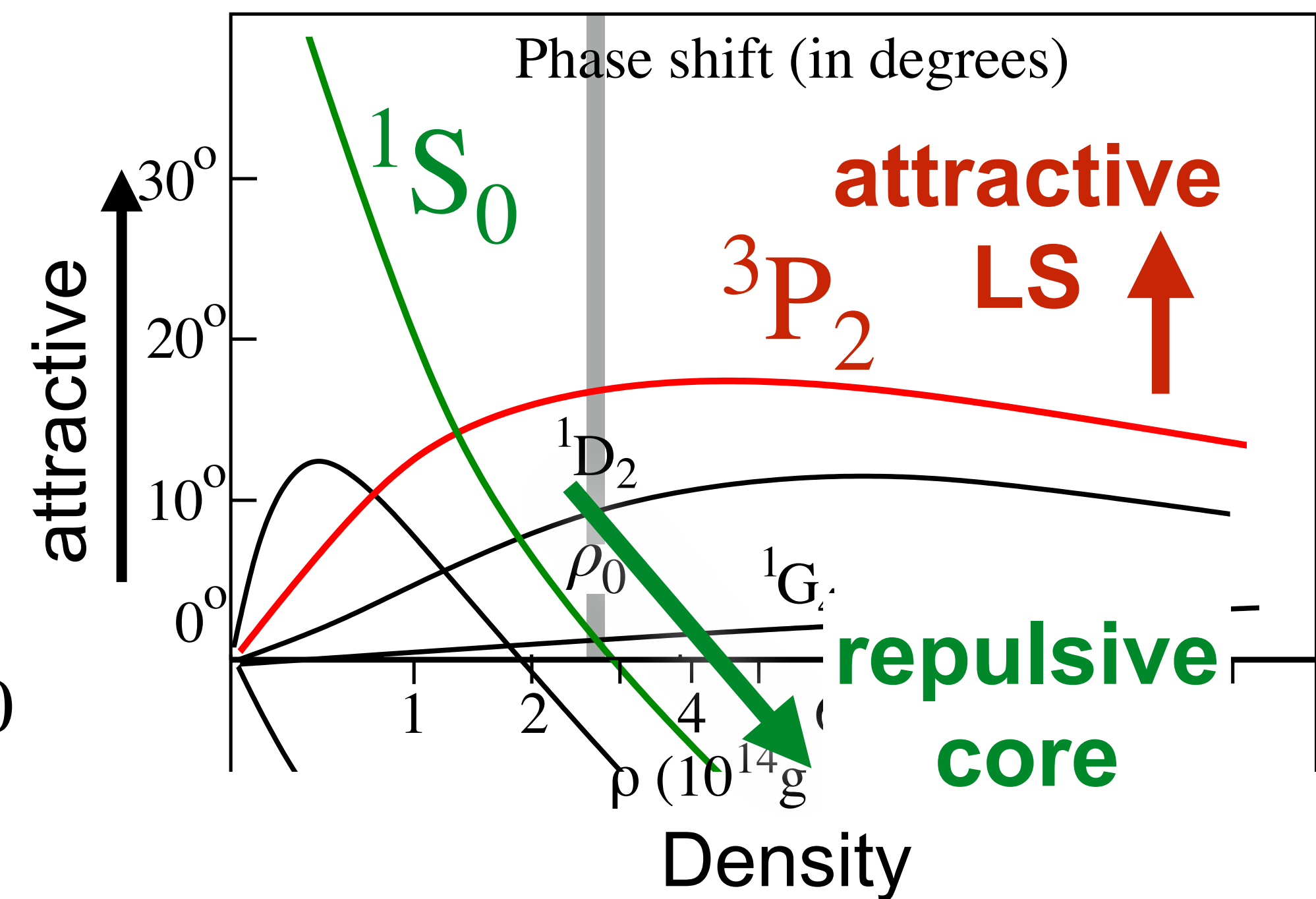
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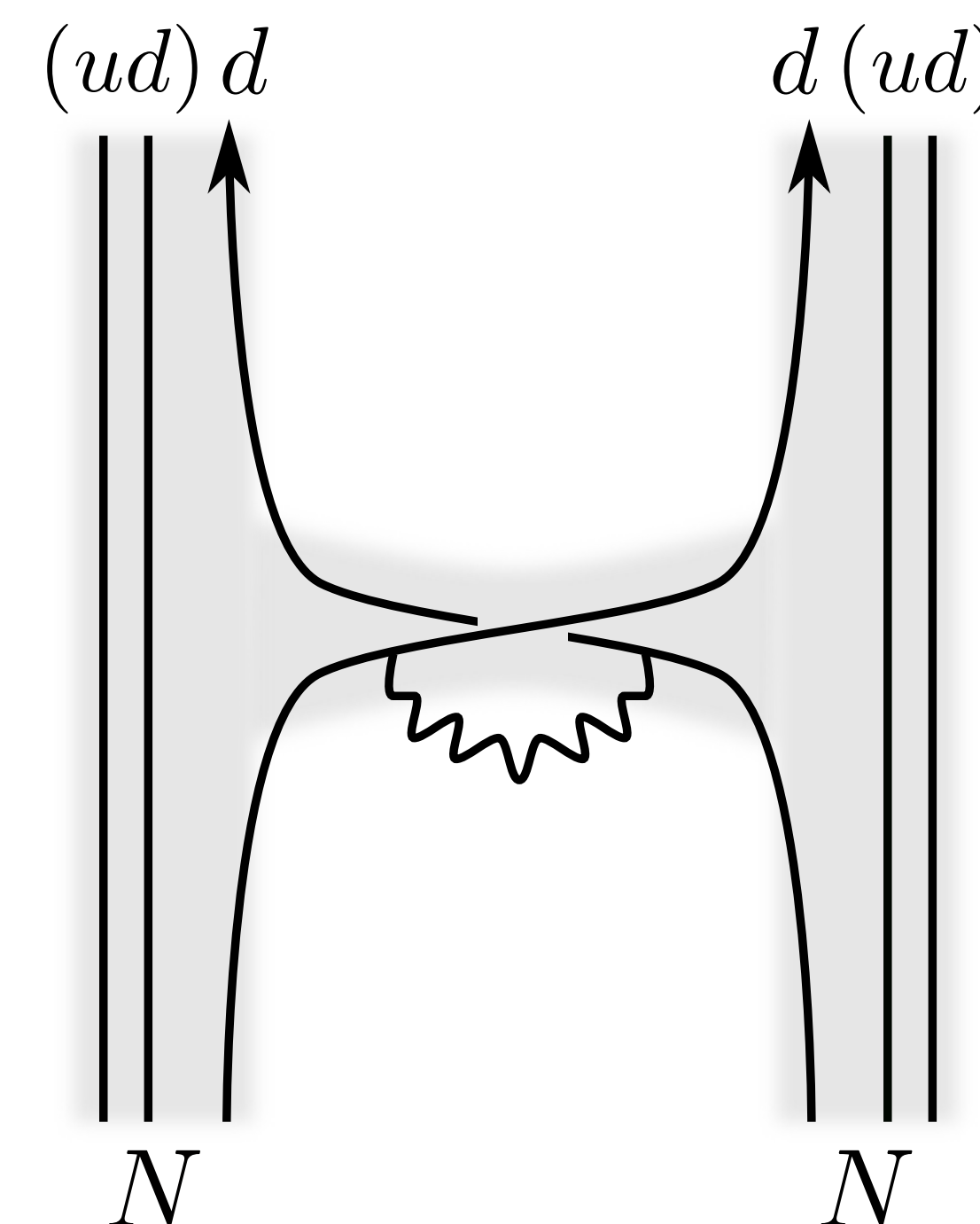


Quark-Hadron Continuity: 2-flavor Case

Fujimoto, Fukushima, Weise, arXiv:1908.09360

Physics mechanism that favors $2SC + \langle dd \rangle$

- Universal mechanism favoring 3P_2 superfluidity
- Short-range repulsion between dd disfavors S -wave:
derived from dd quark exchange
[cf. Repulsive core of NN force: Oka, Yazaki (1980)
derived from **quark exchange between NN s]**



- $J = 2$ state is favored:
confirmed this behavior employing NJL-type model with vector coupling

Coupling to Energy-Momentum Tensor (EMT)

Fujimoto, Fukushima, Weise, arXiv:1908.09360

- Gauge invariant order parameter for $\langle dd \rangle$:

$$\hat{\mathcal{F}} = (\bar{\psi}_d \gamma^i \nabla^j C \bar{\psi}_d^\top) (\psi_d^\top C \gamma_i \nabla_j \psi_d)$$



Fierz
transformation

EMT: $T^{\mu\nu} = \bar{\psi}_d i\gamma^\mu \partial^\nu \psi_d$
 In equilibrium: $T^{\mu\nu} = \text{diag}[\mathcal{E}, P, P, P]$

$\langle \hat{\mathcal{F}} \rangle \approx \frac{3}{4} P^2$ **Macroscopic value related to EoS**

Summary and Outlooks

- Established 2-flavor quark-hadron continuity with **2SC+ $\langle dd \rangle$** phase [Fujimoto,Fukushima,Weise, arXiv:1908.09360]
- Consistent with real data deduced from deep learning analysis [Fujimoto,Fukushima,Murase, arXiv:1903.03400]
- Implication to NS physics:
 - Theoretical background for smooth crossover EoS construction
 - NS cooling — 2SC+X scenario [Grigorian,Blaschke,Voskresensky (2005)]
- Outlooks:
 - Dynamical calculation of $\langle dd \rangle$ gap matrix
 - Vortex continuity of 2SC+ $\langle dd \rangle$ — relation with glitch phenomena? [Alford,Baym,Fukushima,Hatsuda,Tachibana (2018)]