

# Meson masses in external magnetic fields from Lattice QCD

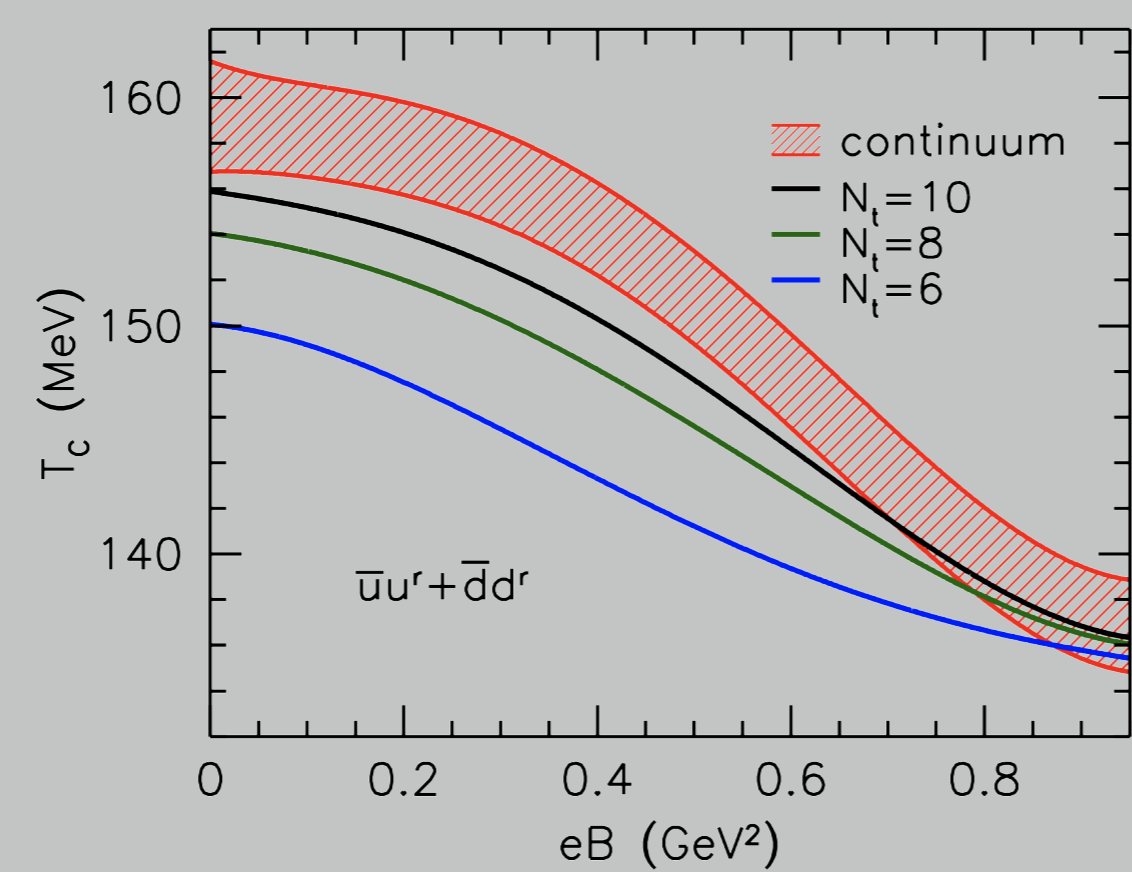
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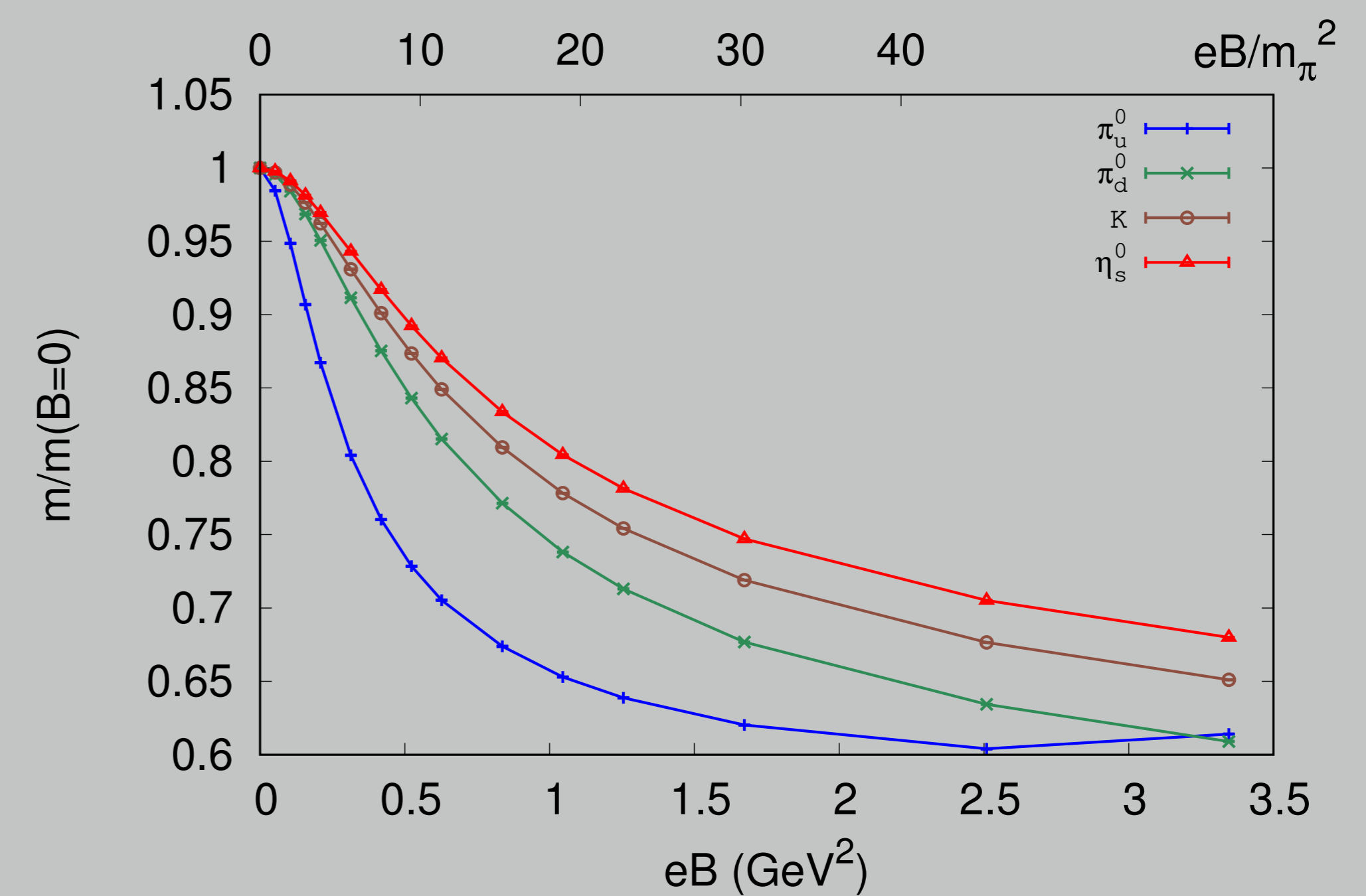


## Motivation

- QCD in strong magnetic field has interesting effects like **Inverse magnetic catalysis**<sup>[1],[2]</sup> (transition critical temperature decreases as magnetic field increases), corresponding lightest Goldstone boson particle mass ( $m_\pi$ ) is thus interesting to be investigated.
- In the framework of NJL model, QCD vacuum transforms into a **superconducting** state in sufficient strong magnetic field signaled by  $\rho^\pm$  meson condensation<sup>[3],[4]</sup> i.e.  $m_{\rho^\pm} = 0$ . A mass relation:  $m_{\rho^\pm} \geq (m_{\pi_u^0} + m_{\pi_d^0})/2$  given by QCD inequality<sup>[5]</sup> can be used to check the condensation of  $\rho^\pm$ .
- Results obtained from lattice simulation can be compared with the ground state energy given by the **Lowest Landau Level (LLL) approximation**  $E^2 = p_z^2 + 2n|qB| + m^2$ .



## Neutral Pseudo-scalar particle masses $\pi_u^0, \pi_d^0, K, \eta_s^0$ v.s. $eB$



- Neutral PS particles' masses decrease as  $eB$  grows and saturate at large  $eB$ .
- Lighter hadrons are more affected by magnetic field.
- Neutral PS particles' masses reduce about **30 ~ 40%**, which is different from effective theory predictions.
- The decreasing of  $m_\pi$  explains the decreasing of transition temperature<sup>[2]</sup>.

## Recent results from Lattice

Mass v.s. $eB$ behavior	$\pi^0$	$\pi^\pm$	$\rho_0^0$	$\rho_+^0$	$\rho_+^+$	$\rho_-^+$	$K$	$\eta^0$
quenched Wilson <sup>[6]</sup> /Overlap <sup>[7]</sup>	↘	↗	↘	↗	↗	↘	↗	
dynamical stout <sup>[6]</sup>	↘							

## Lattice setup

- (2+1)-flavor dynamical HISQ fermion at  $T=0$
- Lattice size:  $32^3 \times 96$ ,  $a^{-1} = 1.6852$  GeV
- $m_\pi = 230$  MeV
- Magnetic field ranges from  $0 \sim 3.3$  GeV<sup>2</sup>
- Multipule wall sources have been used to improve the signal.

## Methodology

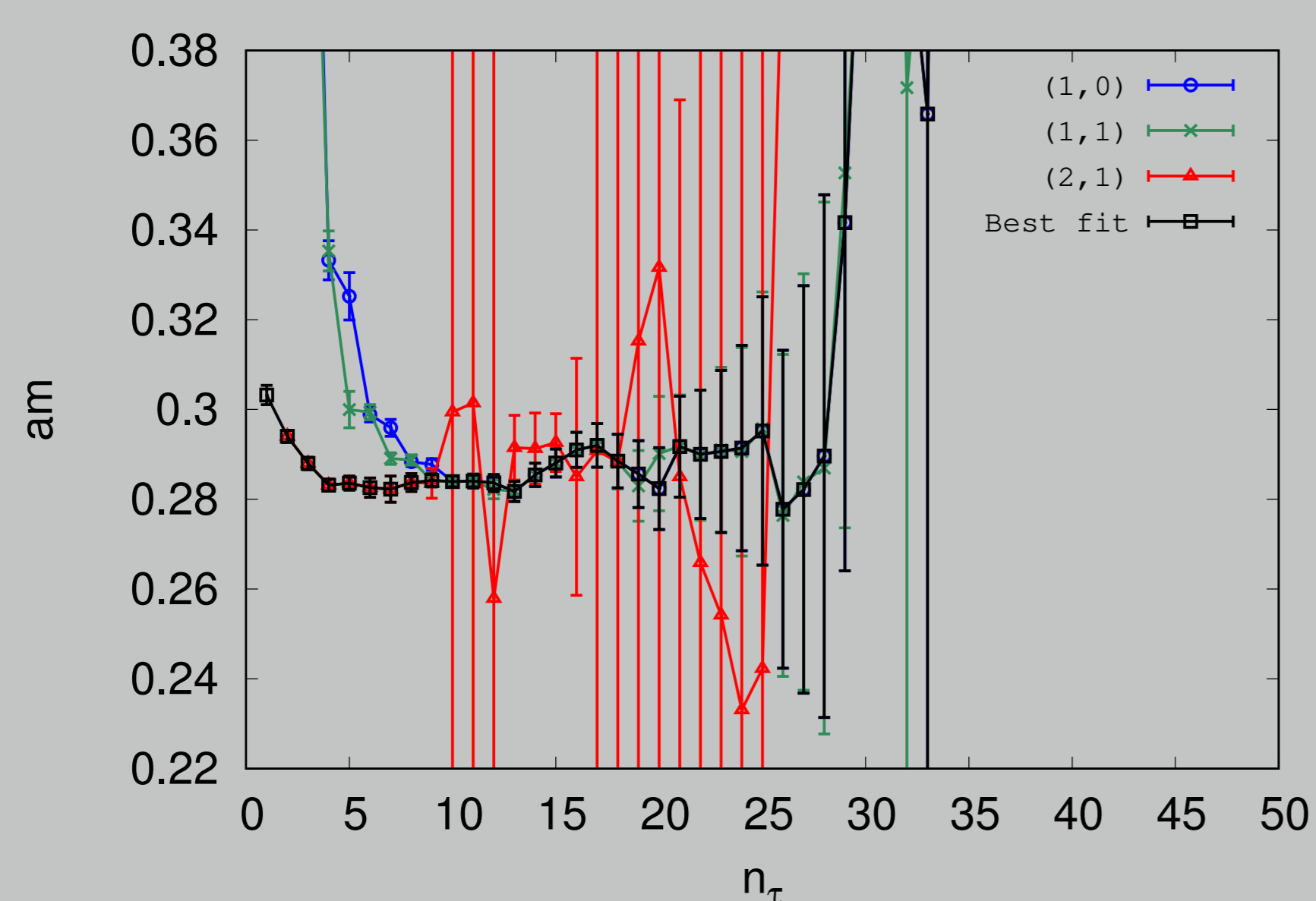
- Correlator formular for staggered fermion:

$$G(\tau) = - \sum_n \zeta(\vec{n}) \text{Tr} [M^{-1\dagger}(n, 0) M^{-1}(n, 0)] \quad (1)$$

- Fit ansatz:

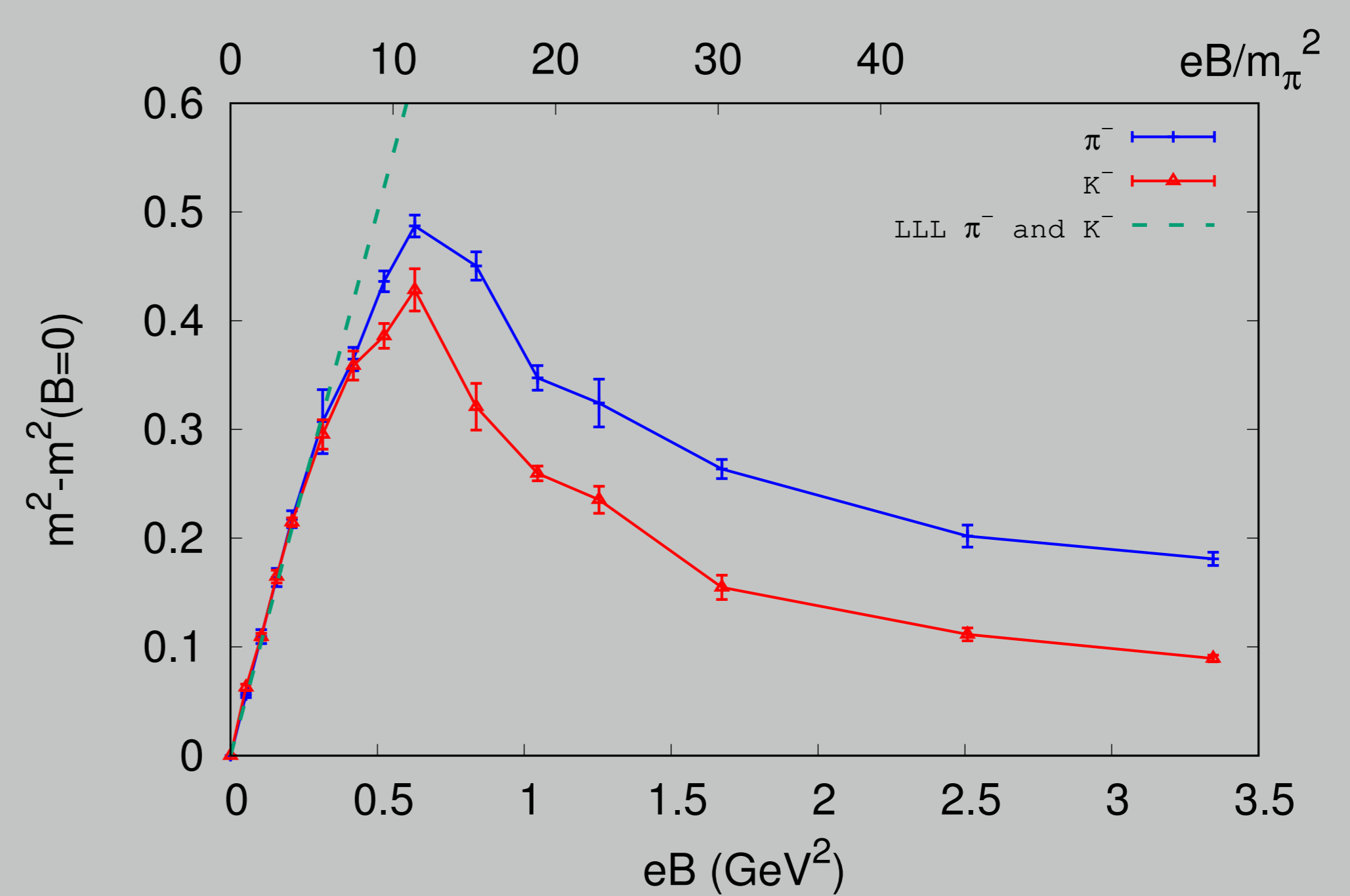
$$G_{\text{fit}}(n_\tau) = \sum_{i=1}^{N_s, no} A_{no,i} \exp(-m_{no,i} n_\tau) - (-1)^{n_\tau} \sum_{i=0}^{N_s, osc} A_{osc,i} \exp(-m_{osc,i} n_\tau) \quad (2)$$

- Selected best fit through AICc<sup>[8]</sup>:



$$\text{AICc} = 2k - \ln(\hat{L}) + \frac{2k^2 + 2k}{n - k - 1} \quad (3)$$

## Charged Pseudo-scalar particle masses $\pi^-, K^-$ v.s. $eB$



- Charged PS particles' masses have non-monotonous behavior, which is different from quenched lattice result.
- At  $0 < eB \lesssim 0.3$  GeV<sup>2</sup>, PS particles' masses can be well described by Lowest Landau level approximation.

## Conclusion

- Masses of neutral pseudo-scalar particles decrease as the magnetic field grows but saturate at large  $eB$ .
  - Neutral particles interact with magnetic field.
- Charged PS particle showed non-monotonous behavior as  $eB$  increases.
- According to the mass relation between  $\rho^\pm$  meson with  $\pi_u^0$  and  $\pi_d^0$ , there is no  $\rho^\pm$  condensation in current window of  $eB$ .

## References

- [1] Bali, Bruckmann et al., Phys.Rev. D86 (2012) 071502
- [2] Bali, Bruckmann et al., JHEP 1202 (2012) 044
- [3] M. N. Chernodub Phys. Rev. Lett. 106, 142003 (2011)
- [4] M. N. Chernodub, Phys. Rev. D 82, 085011 (2010)
- [5] Y. Hidaka and A.Yamamoto Phys. Rev. D 87, 094502 (2013)
- [6] Bali, Brandt et al., PHYS. REV. D 97, 034505 (2018)
- [7] E.V. Luschevskaya et al. Nuclear Physics B 898 (2015)
- [8] J. E. Cavanaugh Statistics & Probability Letters 33.2 (1997)

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