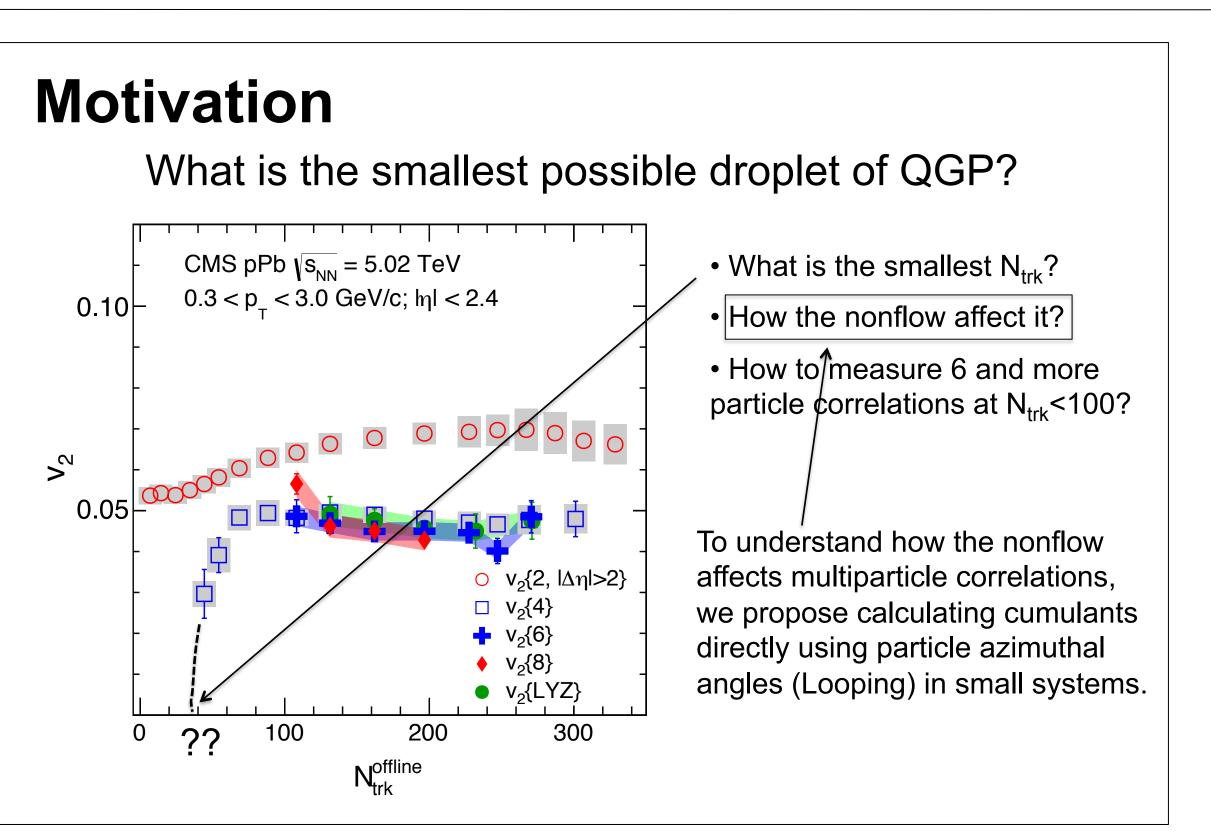
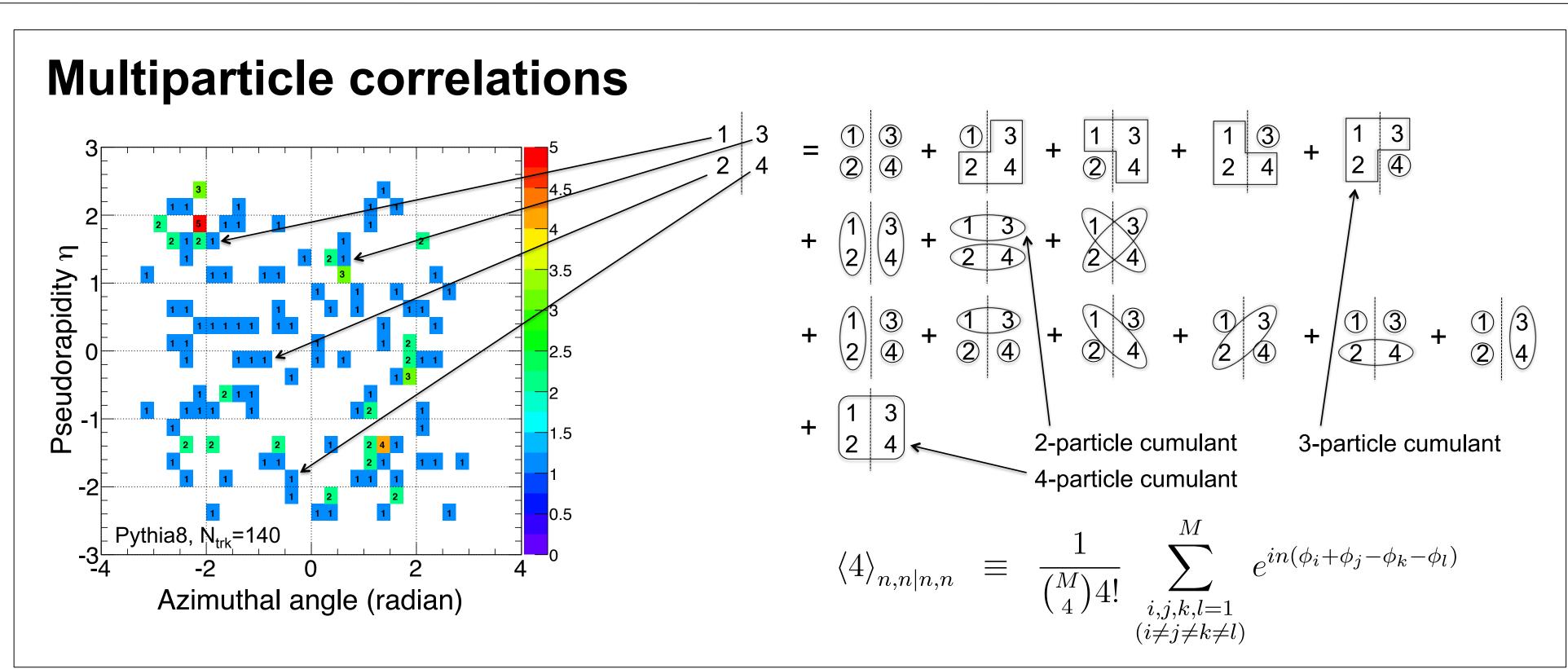
## Multiparticle correlations from the direct calculation of cumulants using particle azimuthal angles

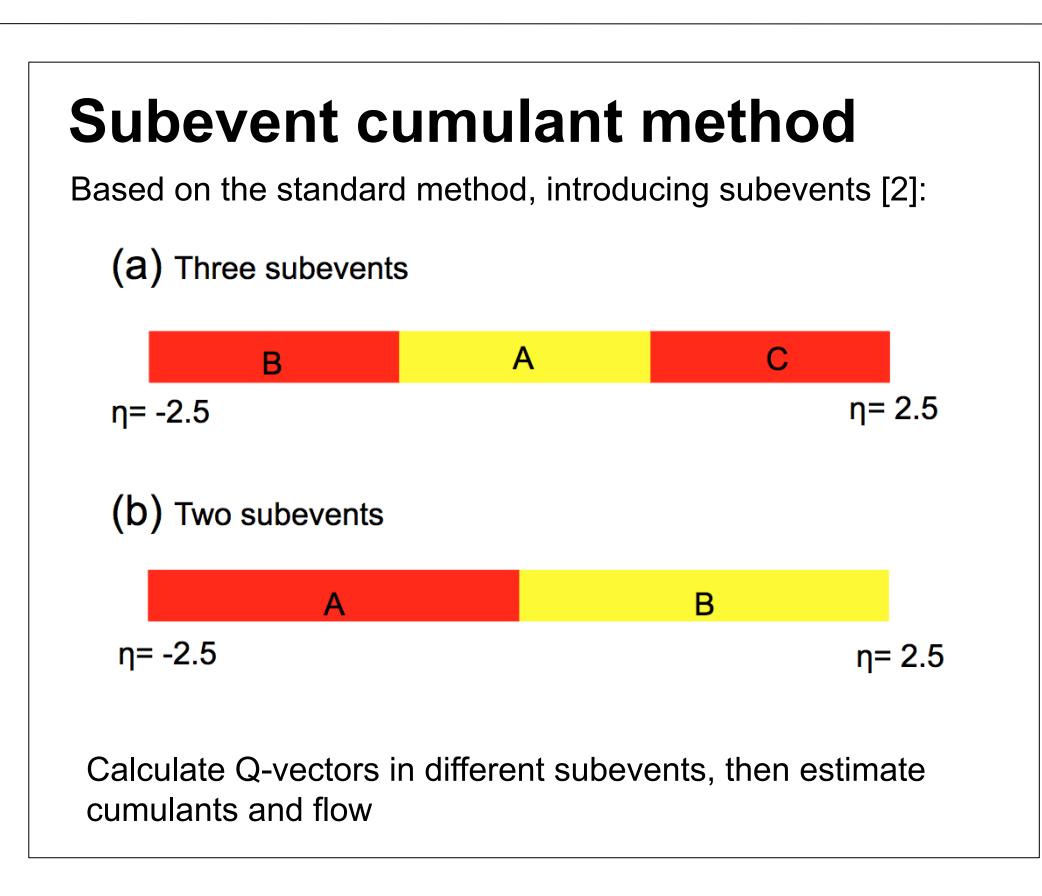
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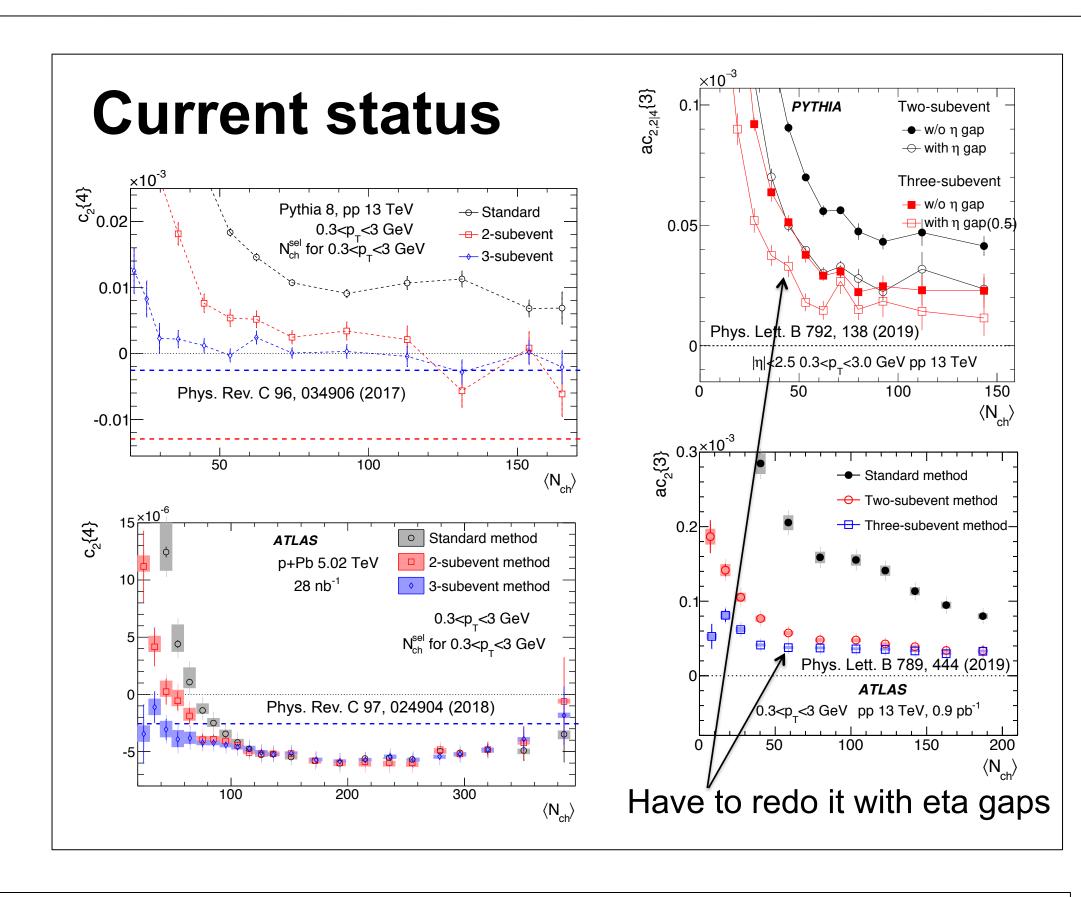




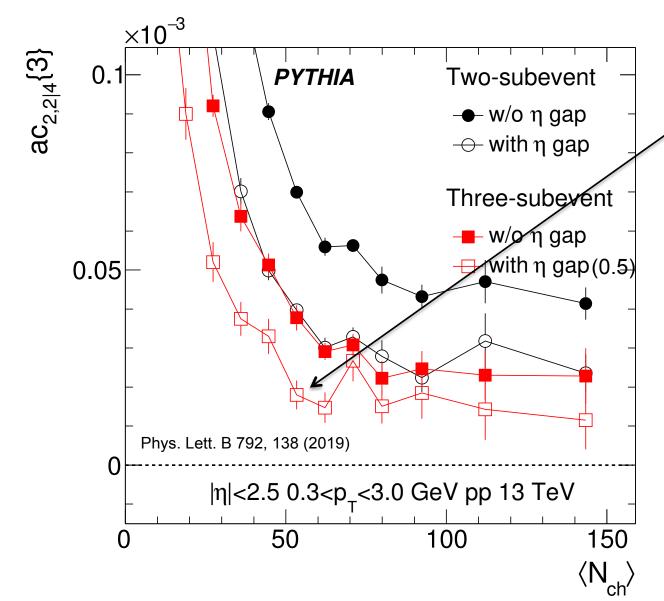


# Standard cumulant method To avoid using nested loops, introduce the Q-vector [1]: $Q_n \equiv \sum_{i=1}^M e^{in\phi_i} \qquad \langle 2 \rangle = \frac{|Q_n|^2 - M}{M(M-1)}$ $|Q_n|^4 = Q_n Q_n Q_n^* Q_n^* = \sum_{i,j,k,l=1}^M e^{in(\phi_i + \phi_j - \phi_k - \phi_l)}$ $\langle 4 \rangle = \frac{|Q_n|^4 + |Q_{2n}|^2 - 2 \cdot \Re \epsilon \left[Q_{2n} Q_n^* Q_n^*\right]}{M(M-1)(M-2)(M-3)}$ $-2 \frac{2(M-2) \cdot |Q_n|^2 - M(M-3)}{M(M-1)(M-2)(M-3)}.$ $c_n\{4\} = \langle \langle 4 \rangle \rangle - 2 \cdot \langle \langle 2 \rangle \rangle^2 \qquad v_n\{4\} = \sqrt[4]{-c_n\{4\}}$





#### Introducing the Looping method



- Need to study correlations vs. η gap
- Run out of statistics quickly
- Any solutions?
- Recall that the idea behind Q-cumulant method is to avoid using nested loops
- If we just loop over particle azimuthal angles, it is very easy to study the η gap dependence
- Also with much better statistics since it keeps all possible combinations

### $\langle 4 \rangle_{n,n|n,n} \equiv \frac{1}{\binom{M}{4} 4!} \sum_{\substack{i,j,k,l=1\\(i \neq j \neq k \neq l)}}^{M} e^{in(\phi_i + \phi_j - \phi_k - \phi_l)}$

#### Time complexity:

- $O(N_{trk}^{n})$ , for n particle cumulant with  $N_{trk}$  total number of particles per event
- The 8 particle cumulant in an event with 1000 particles will take ~1 billion years
- However, our interest is in the small system with  $N_{trk}$  less than 100
- It takes a few seconds to calculate the 4 particle cumulant with  $N_{trk}$ =100
- It could be much faster after applying η gaps between particles

#### Asymmetric Cumulants Symmetric Cumulants Results $\langle 3 \rangle_{n,m|n+m} = \langle e^{i(n\phi_1 + m\phi_2 - (n+m)\phi_3)} \rangle$ Cumulants $SC(n,m) = \langle \langle \{4\}_{n,m} \rangle - \langle \langle \{2\}_n \rangle \rangle \langle \langle \{2\}_m \rangle \rangle = \langle \langle e^{in(\phi_1 - \phi_2) + im(\phi_3 - \phi_4)} \rangle - \langle \langle e^{in(\phi_1 - \phi_2)} \rangle \rangle \langle \langle e^{im(\phi_1 - \phi_2)} \rangle \rangle$ $ac_{n,m|n+m}\{3\} = \langle \langle 3 \rangle \rangle_{n,m|n+m}$ $c_n\{4\} = \langle \langle 4 \rangle \rangle - 2 \cdot \langle \langle 2 \rangle \rangle^2$ $SC(n,m)_{flow} = \langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle$ $ac_{n,m|n+m}\{3\}_{\text{flow}} = \langle v_n v_m v_{n+m} \cos(n\Phi_n + m\Phi_m - (n+m)\Phi_{n+m}) \rangle$ Pythia 8, pp 13 TeV Pythia 8, pp 13 TeV Pythia 8, pp 13 TeV Cumulants 1.0 1.0 Cumulants 0.15 Standard SC $\{2,3\}$ Looping $|\Delta \eta| > 0.8$ - ac<sub>2.2|4</sub>{3} Looping |Δη|>0.8 2-sub |Δη|>0.0 Cumulant c<sub>2</sub>{4} SC $\{2,4\}$ Looping $|\Delta \eta| > 0.8$ \_\_ ac<sub>2.2|4</sub>{3} Looping l∆ηl>1.2 □ 2-sub l∆ηl>0.8 $ac_{2.3|5}$ {3} Looping $|\Delta \eta| > 0.8$ 0.1 $ac_{2,3|5}^{(3)}$ (3) Looping $|\Delta \eta| > 1.2$ Looping I∆ηI>0.8 0.05 Looping l∆ηl>1.2 Asymmetric of the state of the Symmetr -0.05 Introducing subevents makes a huge difference Looping method suppresses the nonflow for symmetric cumulants • Not enough for $ac_{2.2|4}\{3\}$ with an eta gap > 0.8 Looping method removes much more nonflow

#### Conclusions

- Turn-on of flow is important for understanding the smallest possible droplet
- The looping method is introduced to suppress nonflow in small systems
- The method could also be used for Symmetric/Asymmetric Cumulants

#### References

- [1] A. Bilandzic, R. Snellings, and S. Voloshin, Phys. Rev. C 83, 044913 (2011)
- [2] J. Jia, M. Zhou, and A. Trzupek, Phys. Rev. C 96, 034906 (2017)

