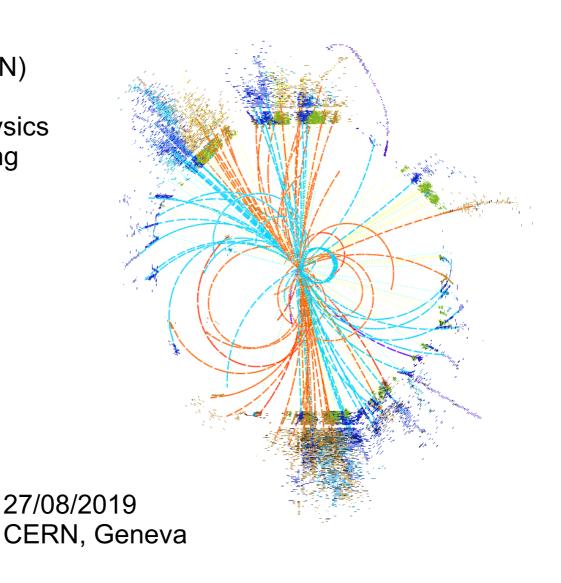
Beyond the Standard Model at future colliders



Philipp Roloff (CERN)

CLIC Detector and Physics Collaboration Meeting





Introduction and caveats

- This talk: mainly highlights from the BSM sessions at the Open Symposium in Granada (and some updates afterwards in view of the briefing book)
- Impossible to cover every model or signature, several representative cases for each thematic area
- Projections have very different levels of sophistication (theory-level studies → DELPHES → full detector simulations)
- Sensitivity to new physics (usually 95% CL. limits), no characterisation of potential discoveries discussed

Topics in BSM

Conveners: Gian Giudice, Paris Sphicas

1.) Electroweak symmetry breaking and new resonances

Andrea Wulzer (th.), Juan Alcaraz (exp.)

Composite Higgs, top partners, particles associated with EW symmetry breaking, heavy Z' and W'

2.) Supersymmetry

Andreas Weiler (th.), Monica D'Onofrio (exp.)

Collider searches, motivations for supersymmetry after the LHC, unexplored corners, new models

3.) Extended Higgs sectors and high-energy flavour dynamics

Veronica Sanz (th.), Ph. R. (exp.)

Two Higgs doublets, singlets, new particles accompanying the Higgs, leptoquarks, particles related to flavour dynamics at the EW scale, rare top decays

4.) Dark Matter

Matthew McCullough (th.), Caterina Doglioni (exp.)

Collider searches, simplified models, comparisons with direct/indirect searches

5.) Feebly-interacting particles

Gilad Perez (th.), Gaia Lanfranchi (exp.)

Long-lived particles, high-handed neutrinos at the EW scale, dark photons at colliders, dark scalar / relaxion, ALPs at colliders

The big questions

- 1.) To what extend can we tell whether the Higgs boson is fundamental or composite?
- → EWSB/new resonances, SUSY
- 2.) Are there new interactions or new particles around or above the electroweak scale?
- → EWSB/new resonances, SUSY, Ext-H/FlavourDyn, DM, FIPs
- 3.) What cases of thermal-relic particles are still unprobed and can be fully covered by future collider searches?
- → DM, SUSY, FIPs
- 4.) To what extend can current or future accelerators probe feebly-interacting sectors?
- → FIPs, SUSY

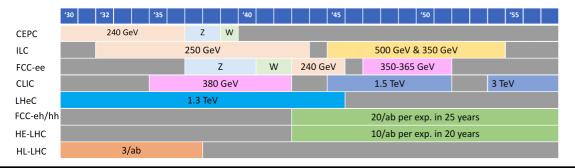
Collider parameters

Collider	Туре	\sqrt{s}	P [%]	N(Det.)	$\mathscr{L}_{ ext{inst}}$	$\mathscr{L}_{_{1}}$	Time
			$[e^{-}/e^{+}]$		$[10^{34}] \text{ cm}^{-2} \text{s}^{-1}$	$[ab^{-1}]$	[years]
HL-LHC	pp	14 TeV	-	2	5	6.0	12
HE-LHC	pp	27 TeV	-	2	16	15.0	20
FCC-hh	pр	100 TeV	-	2	30	30.0	25
FCC-ee	ee	M_Z	0/0	2	100/200	150	4
		$2M_W$	0/0	2	25	10	1-2
		240 GeV	0/0	2	7	5	3
		$2m_{top}$	0/0	2	0.8/1.4	1.5	5
							(+1)
ILC	ee	250 GeV	±80/±30	1	1.35/2.7	2.0	11.5
		350 GeV	$\pm 80/\pm 30$	1	1.6	0.2	1
		500 GeV	$\pm 80/\pm 30$	1	1.8/3.6	4.0	8.5
							(+1)
CEPC	ee	M_Z	0/0	2	17/32	16	2
		$2M_W$	0/0	2	10	2.6	1
		240 GeV	0/0	2	3	5.6	7
CLIC	ee	380 GeV	±80/0	1	1.5	1.0	8
		1.5 TeV	$\pm 80/0$	1	3.7	2.5	7
		3.0 TeV	$\pm 80/0$	1	6.0	5.0	8
							(+4)
LHeC	еp	1.3 TeV	-	1	0.8	1.0	15
HE-LHeC	ep	2.6 TeV	-	1	1.5	2.0	20
FCC-eh	ер	3.5 TeV	-	1	1.5	2.0	25

pp colliders

e⁺e⁻ colliders

ep colliders



+ LE-FCC: pp, 15 ab⁻¹ at \sqrt{s} = 37.5 TeV

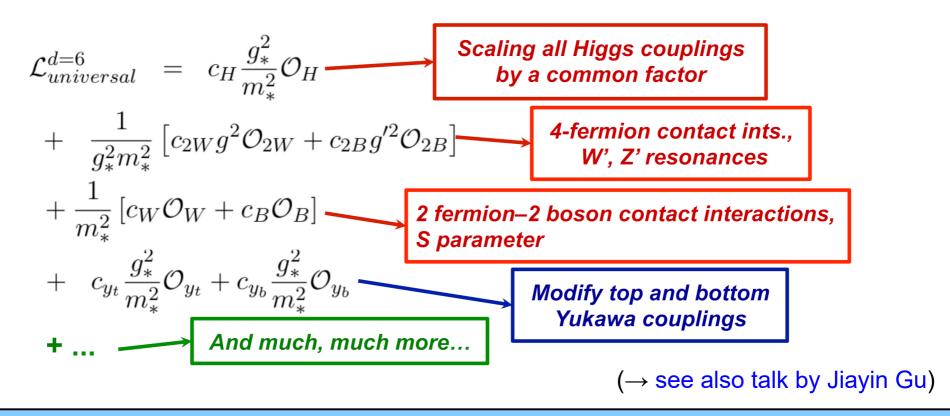
arXiv:1905.03764

Direct vs. indirect

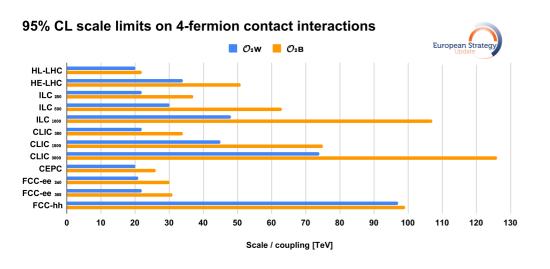
Direct searches: tuned for specific models (or classes of models)

Precision measurements: important complementary information

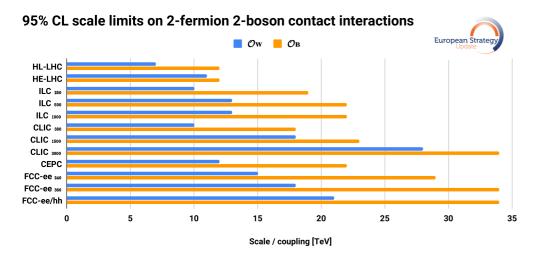
Standard Model Effective Field Theory: contributions from new physics expressed by dimension-6 operators



Contact interactions



- Projected limits from di-fermion final states (e⁺e⁻ → ff̄, Drell-Yan with neutral and charged currents)
- Sensitivity increases significantly with \sqrt{s}



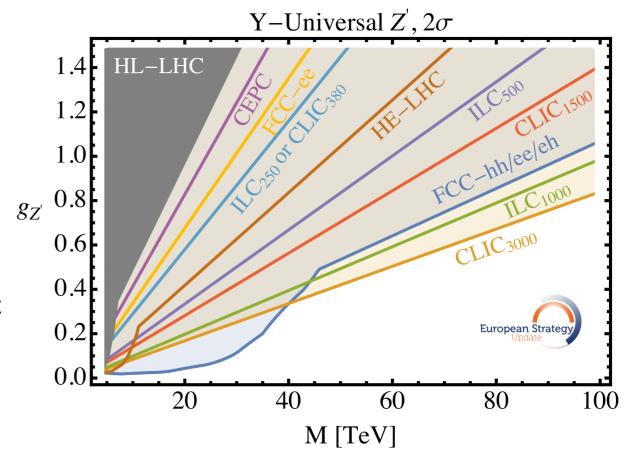
- New physics effects in the interaction between gauge and Higgs sectors
- O_W dominated by $e^+e^- \to ZH$ at CLIC (\to see talk by Matthias Weber), pp $\to WZ$ at FCC-hh
- Largest sensitivity in e^+e^- collisions at lower \sqrt{s} (and on O_B in general) from oblique parameter S

Y-Universal Z'

New neutral gauge boson Z'
with mass M and charges to
SM particles equal to hypercharge

$$\frac{c_{2B}}{\Lambda^2} = \frac{g_{Z'}^2}{g'^4 M_{Z'}^2}$$

- Direct reach inferior to the indirect one for high g_{7}
- NB: $g_{Z'} > 1.5 \rightarrow \text{width}$ exceeds 0.3 M

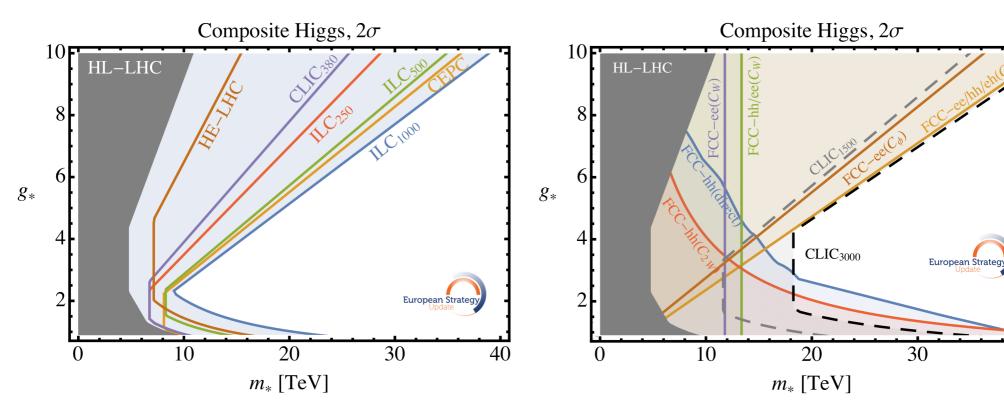


Composite Higgs (1)

m.: mass scale

g.: coupling

$$\frac{c_{\phi}}{\Lambda^2} \sim \frac{g_*^2}{m_*^2} \,, \qquad \frac{c_W}{\Lambda^2} \sim \frac{1}{m_*^2} \,, \qquad \frac{c_{2W}}{\Lambda^2} \sim \frac{1}{g_*^2 m_*^2} \,$$



Higgs factories (including 380 GeV CLIC) significantly better than HL-LHC

FCC-all and 3 TeV CLIC similar

Composite Higgs (2)

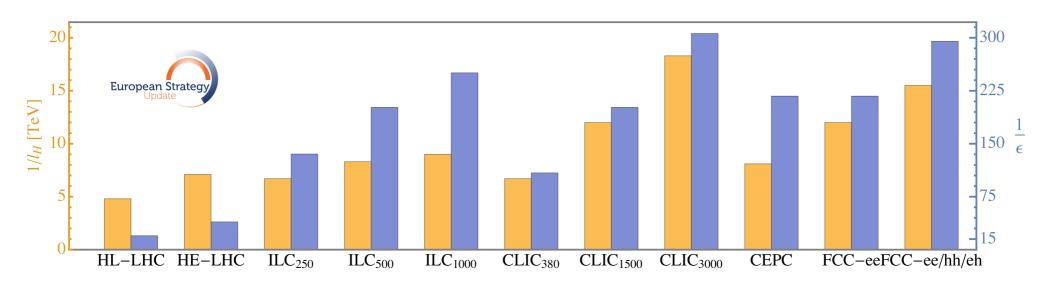
m_{*}: mass scale

g.: coupling

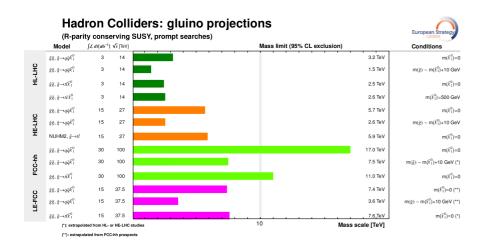
$$\frac{c_{\phi}}{\Lambda^2} \sim \frac{g_*^2}{m_*^2}, \quad \frac{c_W}{\Lambda^2} \sim \frac{1}{m_*^2}, \quad \frac{c_{2W}}{\Lambda^2} \sim \frac{1}{g_*^2 m_*^2}$$

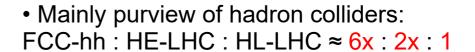
Inverse Higgs length: 1 / $\ell_H = m_*$

 \rightarrow 1 / $\ell_{\rm H}$ = 10 - 20 TeV is 4 orders of magnitude below the size of the proton!



Strongly-interacting SUSY





- CLIC at 3 TeV exceeds HL-LHC for top squarks for heavier LSP
- Powerful probes on the role of naturalness in the Higgs sector

Fine-tuning parameter: $1/\epsilon = \Delta m_H^2/m_H^2$

	ϵ	High-scale mediation	Low-scale mediation
•	stop	$5 \times 10^{-5} \left(\frac{10 \text{ TeV}}{m_{\tilde{t}}} \right)^2$	$2 \times 10^{-3} \left(\frac{10 \text{ TeV}}{m_{\tilde{t}}}\right)^2$
	gluino	$7 \times 10^{-6} \left(\frac{17 \text{ TeV}}{m_{\tilde{g}}}\right)^2$	$6 \times 10^{-3} \left(\frac{17 \text{ TeV}}{m_{\tilde{g}}}\right)^2$

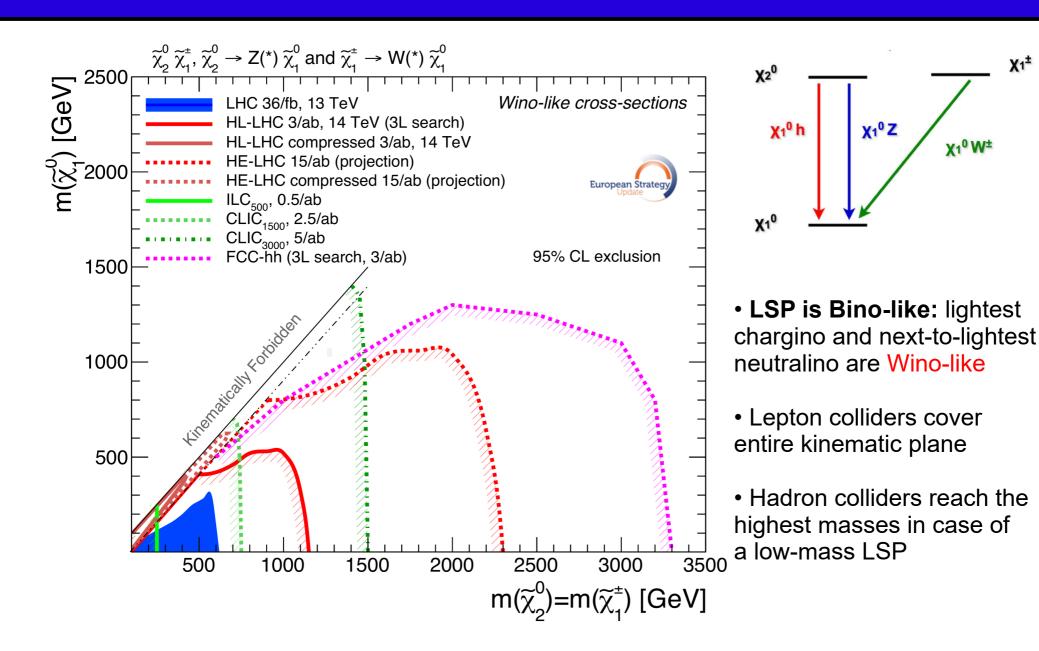
All Colliders: Top squark projections

(R-parity conserving SUSY, prompt searches)



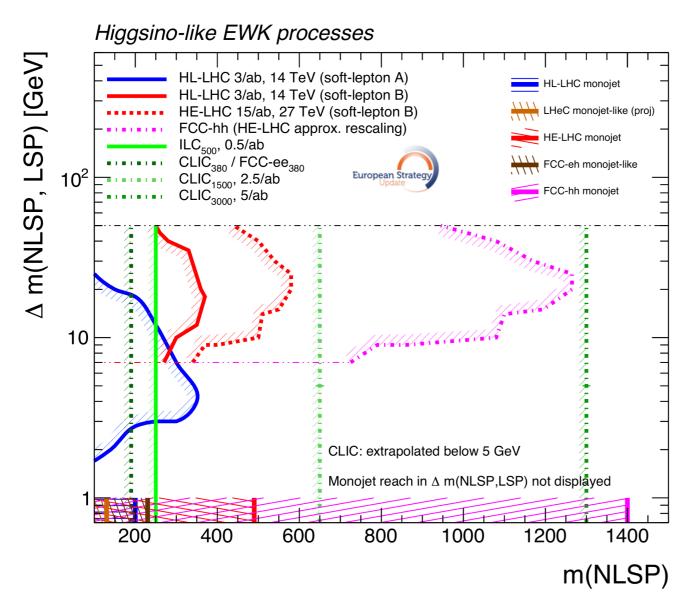
	Model	∫ <i>L dt</i> [ab	¹] √s [TeV]	Mass limit (95% CL exclusion)	Conditions
o	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{X}_1^0$	3	14	1.7 TeV	$m(\tilde{\chi}_1^0)=0$
HL-LHC	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{X}_1^0/3 \text{ bod}$, 3	14	0.85 TeV	$\Delta m(\tilde{r}_1, \tilde{X}_1^0) \sim m(t)$
Ξ	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}^0_1/4 \text{ bod}$	у 3	14	0.95 TeV	$\Delta m(\tilde{t}_1, \tilde{X}_1^0) \sim 5 \text{ GeV, monojet (*)}$
0	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}^{\pm}/t\tilde{\chi}_1^0, \tilde{\lambda}$	0 15	27	3.65 TeV	$m(\tilde{\chi}_1^0)=0$
H-LHC	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0/3\text{-bod}$	y 15	27	1.8 TeV	$\Delta m(\tilde{t}_1, \tilde{\chi}_1^0) \sim m(t)$ (*)
Ξ	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{X}_1^0/4-boo$	y 15	27	2.0 TeV	$\Delta m(\tilde{t}_1, \tilde{X}_1^0) \sim 5$ GeV, monojet (*)
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{X}_1^0$	15	37.5	4.6 TeV	$m(\bar{X}_1^0)=0$ (**)
LE-FCC	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0/3\text{-bod}$	y 15	37.5	4.1 TeV	$m(\tilde{\chi}_1^0)$ up to 3.5 TeV (**)
Ÿ	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}^0_1/4-boo$	y 15	37.5	2.2 TeV	$\Delta m(\tilde{t}_1, \tilde{\mathcal{X}}_1^0) \sim 5 \text{ GeV, monojet (**)}$
8	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}^{\pm}/t\tilde{\chi}_1^0$	2.5	1.5	0.75 TeV	$m(\tilde{\chi}_1^0)=0$
CLIC ₁₅₀₀	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}^{\pm}/t\tilde{\chi}^0_1$	2.5	1.5	0.75 TeV	$\Delta m(\tilde{r}_1, \tilde{X}_1^0) \sim m(t)$
Ö	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}^{\pm}/t\tilde{\chi}_1^0$	2.5	1.5	(0.75 - e) TeV	$\Delta m(\tilde{t}_1, \tilde{\chi}_1^0) \sim 50 \text{ GeV}$
90	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}^{\pm}/t\tilde{\chi}_1^0$	5	3.0	1.5 TeV	m($\tilde{\chi}_{1}^{0}$)~350 GeV
CLIC ₃₀₀₀	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}^{\pm}/t\tilde{\chi}_1^0$	5	3.0	1.5 TeV	$\Delta m(\tilde{r}_1, \tilde{X}_1^0) \sim m(t)$
0	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}^{\pm}/t\tilde{\chi}_1^0$	5	3.0	(1.5 - e) TeV	$\Delta m(\tilde{t}_1, \tilde{\chi}_1^0) \sim 50 \text{ GeV}$
£	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{X}_1^0$	30	100	10.8 TeV	$m(\tilde{\chi}_1^0)=0$
FCC-hh	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0/3\text{-bod}$	y 30	100	10.0 TeV	$m(\bar{\ell}_1^0)$ up to 4 TeV
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0/4$ -boo	у 30	100	5.0 TeV	$\Delta m(\tilde{t}_1, \tilde{X}_1^0) \sim 5$ GeV, monojet (*)
				10 ⁻¹ 1 Mass scale [TeV]	
	(*) indicates project (**) extrapolated fr			perimental searches	ne un to kinamatic limit 🎜 📶

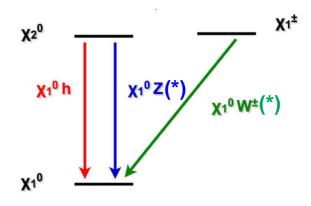
EW SUSY: Wino



X1[±]

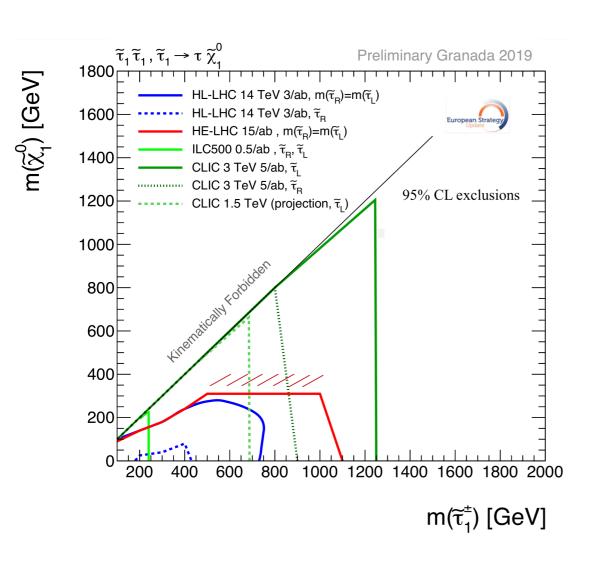
EW SUSY: Higgsino

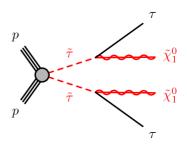




- LSP Higgsino-like: ΔM(NLSP, LSP) ~ O(GeV)
- LSP pure Higgsino: ΔM(NLSP, LSP) ~ 160 MeV → targeted by disappearing track analysis (see later)

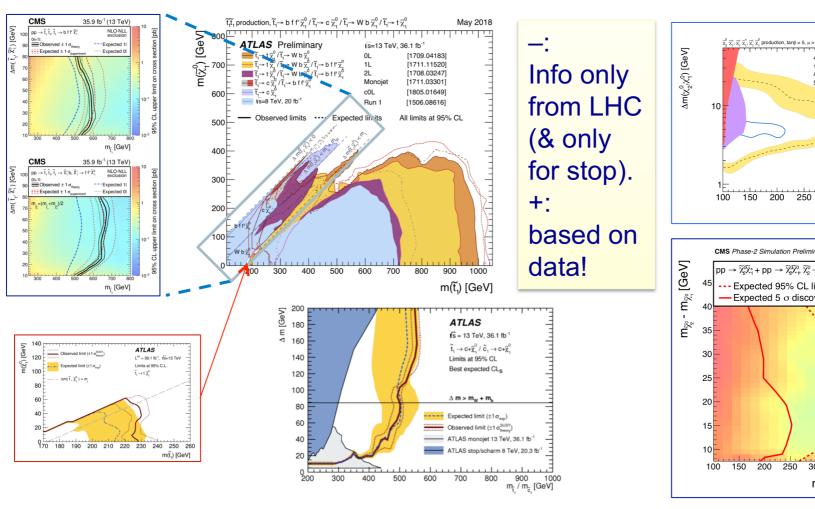
EW SUSY: staus

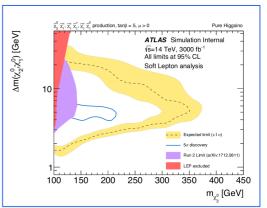


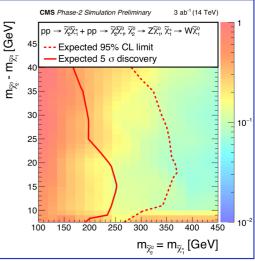


- Mixing of RH and LH component
- → large impact on cross section
- FCC-hh might push limit to 3 4 TeV for low LSP mass (boosed tau reconstruction to be studied)
- Lepton colliders cover large part of kinematic plane for LH stau
- → see talk by Ulrike Schnoor

SUSY: unexplored corners?



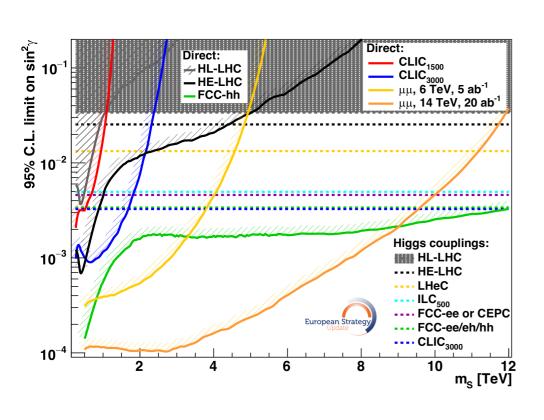




Indeed, after LHC, there will be holes [in low mass regions]; closing or looking at how to close them at HL-LHC; for EWKinos, some regions will remain difficult @ pp.

Extended Higgs sectors: SM + real scalar singlet

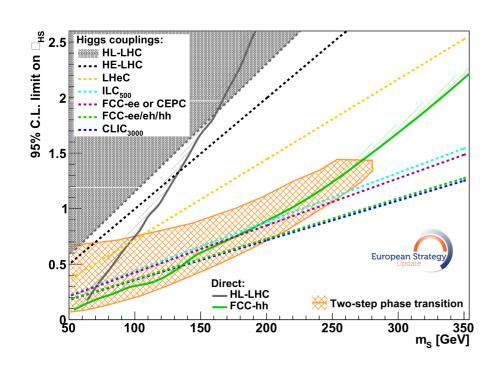
$$|V_0 = -\mu^2 |H|^2 + \lambda |H|^4 - \frac{1}{2} \mu_S^2 S^2 + \frac{1}{4} \lambda_S S^4 + \lambda_{HS} |H|^2 S^2$$



- H & S can mix: sin²γ
- Direct searches: S → ZZ in pp
 and S → HH in e⁺e⁻ give the best sensitivity
- Indirect searches: sensitivity from c_φ
 (overall scaling of the Higgs couplings)
- Direct & indirect proves provide complementary information
- Direct reach at FCC-hh better than precision Higgs couplings for m_s < 12 TeV

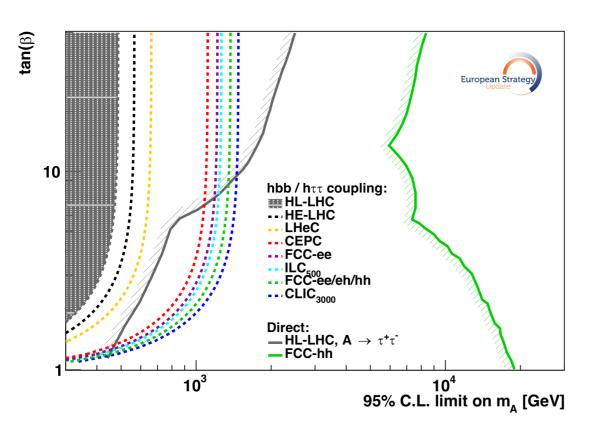
No mixing limit

$$|V_0 = -\mu^2 |H|^2 + \lambda |H|^4 - \frac{1}{2} \mu_S^2 S^2 + \frac{1}{4} \lambda_S S^4 + \lambda_{HS} |H|^2 S^2$$



- No mixing \rightarrow limit on portal coupling λ_{HS}
- Direct searches: S escapes undetected, event selection using VBF jets: pp → SSjj
- Indirect searches: sensitivity from C (overall scaling of the Higgs couplings)
- FCC-all or 3 TeV CLIC would cover most of the region compatible with two-step phase transition (where the singlet supports the Higgs in delivering a strong first-order phase transition)
- Strong first-order phase transitions could lead to sizeable gravitational wave signals at future experiments like LISA

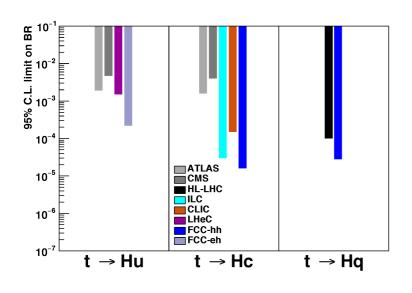
Heavy neutral scalars in minimal SUSY

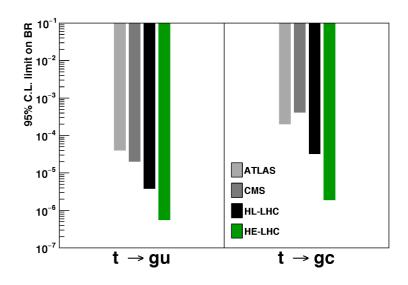


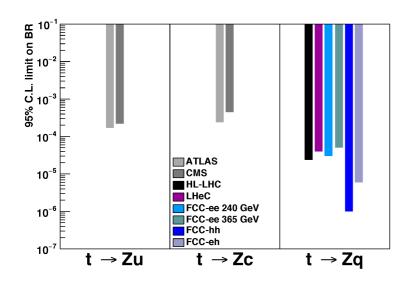
Example: CP-odd scalar A

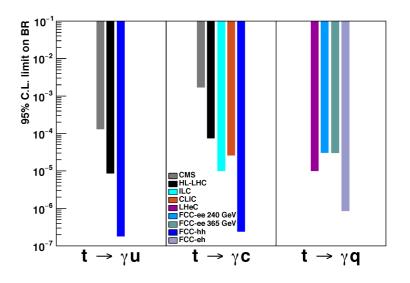
- Mass reach generally close to \sqrt{s} / 2 at lepton colliders
- HL-LHC sensitive to heavy neutral scalars up to 2.5 TeV for tan β > 50 using $\tau^+\tau^-$
- Exclusion limits > 5 TeV at FCC-hh (20 TeV at low tan β)
- Indirect sensitivity from g_{hbb} or $g_{h\tau\tau}$ (further improvement from global fit)

Flavour dynamics: FCNC





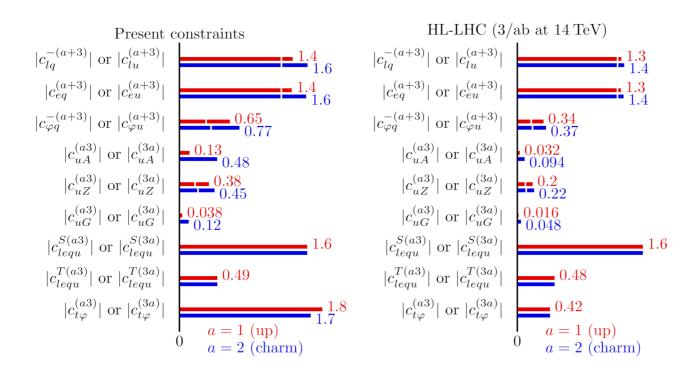




Top-quark FCNC: EFT for HL-LHC

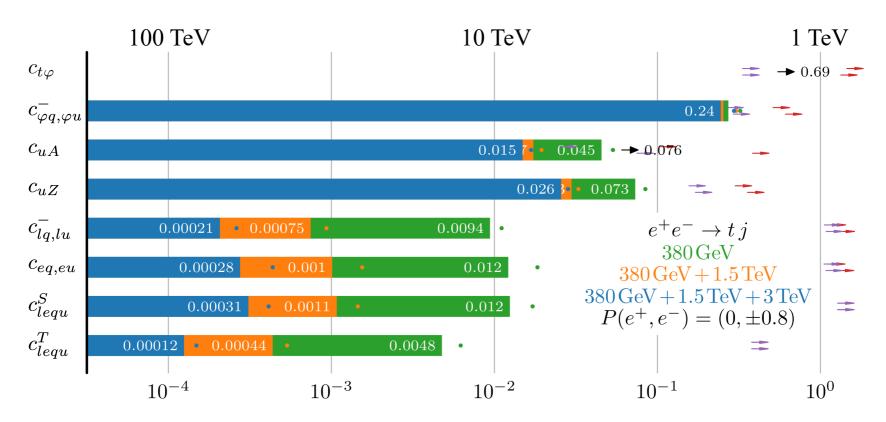
Sensitivity to top-quark FCNC effects can be studied using EFT

Input: limits on FCNC branching ratios, limits on e⁺e⁻ → tj from LEP II



White marks: individual limits

Top-quark FCNC: e⁺e⁻ → tj at CLIC



95% C.L. limits on top-quark FCNC operator coefficients

Black arrows: decays at CLIC (see slide X)

Red arrows: current LHC

Magenta arrows: HL-LHC projections Dots: CLIC without beam polarisation

- The high-energy runs significantly improve the sensitivity for "four-fermion" operators
- e⁺e⁻ → tj much more powerful than the decays at high-energy lepton colliders

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Dark Matter

Cosmology (thermal freeze-out mechanism) provides a strong motivation for direct, indirect and collider searches

- → DM masses from multi-keV to 100 TeV
- → Couplings to SM particles comparable or weaker than EW interactions

Focus on GeV - TeV mass range, two approaches:

- 1.) "Classic" electroweak WIMP candidates (SUSY inspired):
 - Winos and Higgsinos
 - 2.) Simplified models with mediator particles:
 - Axial-vector simplified models
 - Scalar simplified models

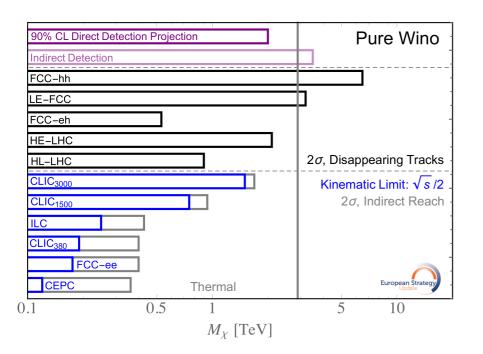
WIMPs: Wino and Higgsino

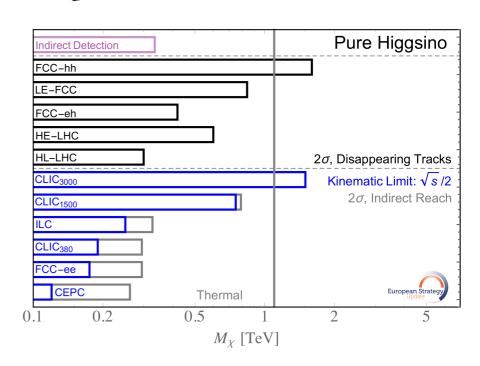
In SUSY terminology: "pure Wino" and "pure Higgsino" (spin ½ particles transforming as

doublets or triplets under SU(2) symmetry)

Main techniques:

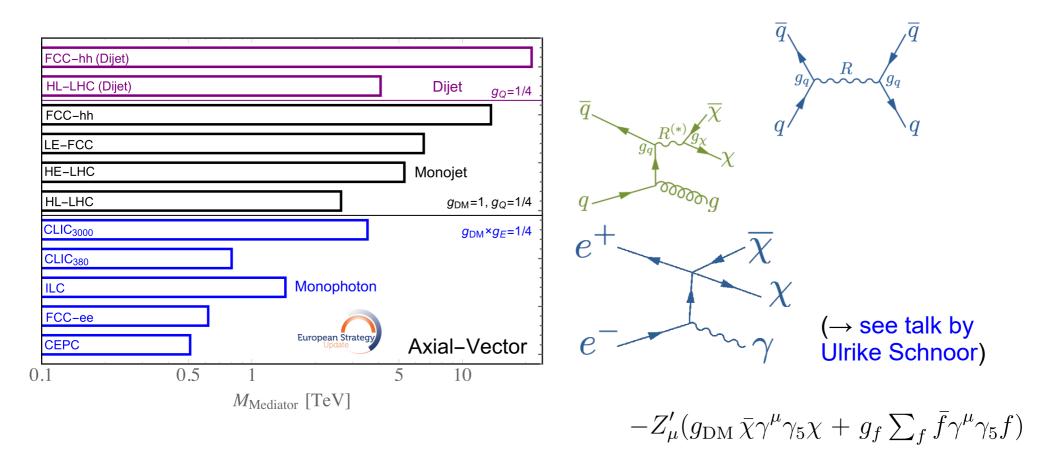
- Dark Matter in loops
- Disappearing tracks (→ see talk by Cecilia Ferrari)





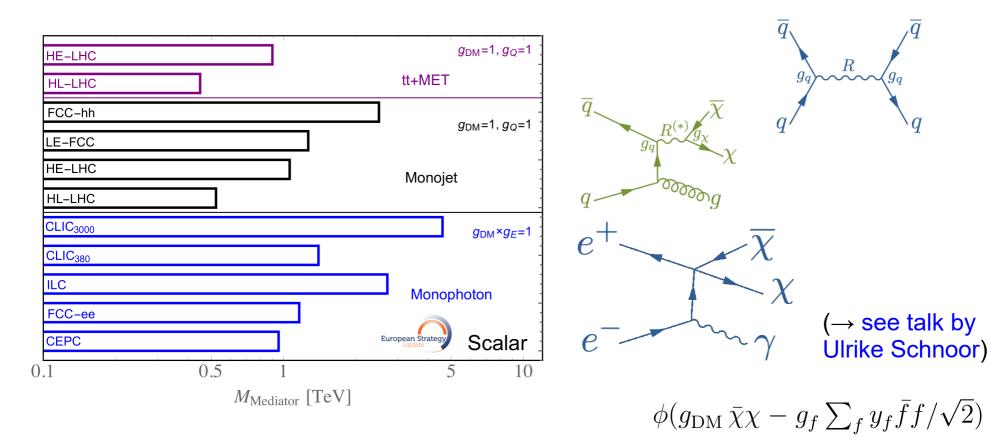
→ Thermal Wino within reach of FCC-hh and LE-FCC, thermal Higgsino within reach of FCC-hh and 3 TeV CLIC

Simplified models: axial vector



- Mediator is spin-1 particle (Z') coupled to an axial-vector current (reach of direct DM searches limited → interesting for colliders)
- pp colliders assume couplings to quarks only, e⁺e⁻ colliders assume couplings to leptons only → projections not directly comparable

Simplified models: scalar



y_f: Yukawa couplings

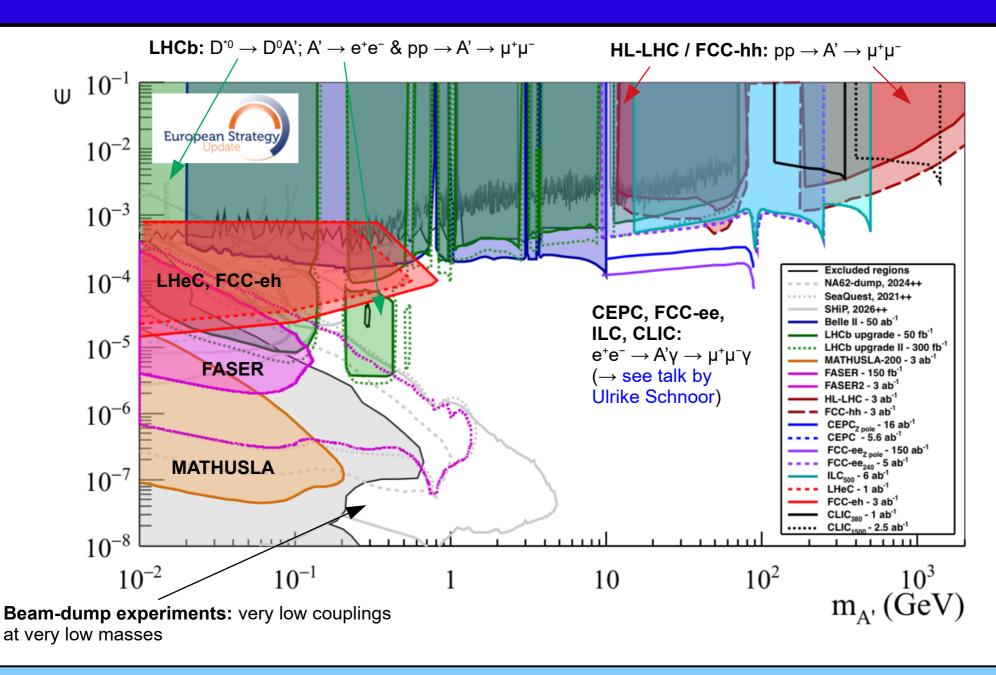
- Mediator is spin-0 particle (φ)
 (reach of direct DM searches limited → interesting for colliders)
- pp colliders assume couplings to quarks only, e⁺e⁻ colliders assume couplings to leptons only → projections not directly comparable

Feebly-interacting particles

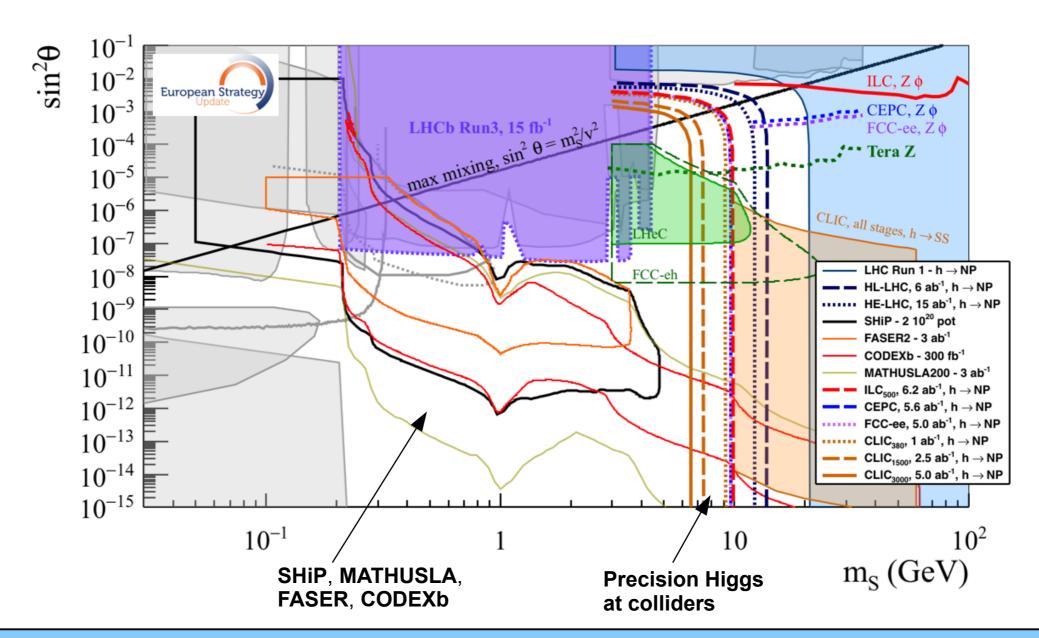
- Undetected particles feebly interacting with SM particles: part of Hidden or dark sector
- Masses and interactions of dark sector particles unknown, focus here on MeV to tens of GeV range → important motivation is Dark Matter
- Very wide range of theoretical models: simplified models (4 portals) to compare experiments from Physics Beyond Colliders study
- From portals: identify benchmarks to evaluate experimental sensitivity
- → common ground to compare machines / experiments

Portal	Coupling
Vector (Dark Photon, A_{μ})	$-\frac{\epsilon}{2\cos\theta_W}F'_{\mu\nu}B^{\mu\nu}$
Scalar (Dark Higgs, S)	$(\mu S + \lambda_{HS} S^2) H^{\dagger} H$
Fermion (Sterile Neutrino, N)	$y_N LHN$
Pseudo-scalar (Axion, a)	$\frac{a}{f_a}F_{\mu\nu}\tilde{F}^{\mu\nu}, \frac{a}{f_a}G_{i,\mu\nu}\tilde{G}_i^{\mu\nu}, \frac{\partial_{\mu}a}{f_a}\overline{\psi}\gamma^{\mu}\gamma^5\psi$

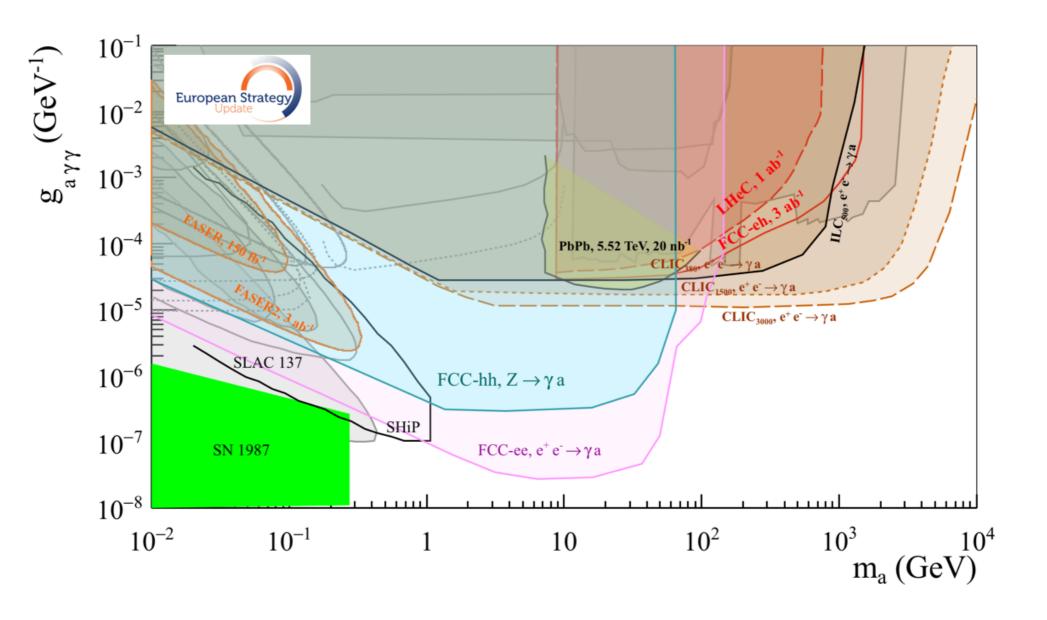
Dark Photons



Dark scalar mixing with the Higgs



ALPs coupled to photons



Summary

- Substantial improvement with respect to HL-LHC possible for all discussed physics topics
- Precision Higgs measurements are central for many BSM scenarios (e.g. composite Higgs models, extended Higgs sectors)
- Large amount of complementarity between direct and indirect searches (e.g. Higgs couplings, EW precision measurements at Z-pole, flavour sector, SM processes at highest energies) for new particles and interactions
- For feebly-interacting particles (including forms of Dark Matter), the reach of energy-frontier colliders is complemented by beam-dump and fixed target facilities at low masses

Thank you!

Backup slides

Heavy Neutral Lepton mixed with electron neutrino

