

# Collider signatures of dark CP violation

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In collaboration with

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1 Introduction and motivation

2 CPV DM in 3HDMs

3 CPV DM in colliders

4 Summary

# The SM and beyond

## SM predictions confirmed:

- the W & Z (1983)
- the top quark (1995)
- the tau neutrino (2000)
- “a” Higgs boson (2012)

## What is missing:

- Fermion mass hierarchy explanation
- Dark Matter candidate
- Baryon asymmetry

⇒ Non-minimal Higgs frameworks



# BSMs to the rescue

Scalar extensions with and without a  $Z_2$  symmetry:

- Higgs portal models: SM + scalar singlet

- $\phi_{SM}, S \Rightarrow$  CPV, DM

- $\phi_{SM}, S \Rightarrow$  DM, CPV

- 2HDM: SM + scalar doublet

- Types I, II, III, IV:  $\phi_1, \phi_2 \Rightarrow$  CPV, DM

- IDM - I(1+1)HDM:  $\phi_1, \phi_2 \Rightarrow$  DM, CPV

- 3HDM: SM + 2 scalar doublets

- Weinberg model:  $\phi_1, \phi_2, \phi_3 \Rightarrow$  CPV, DM

- I(1+2)HDM:  $\phi_1, \phi_2, \phi_3 \Rightarrow$  DM, CPV

- I(2+1)HDM:  $\phi_1, \phi_2, \phi_3 \Rightarrow$  CPV, DM

# CP-violating DM in 3HDMs

$$\phi_1, \phi_2, \phi_3$$

$$g_{Z_2} = \text{diag}(-1, -1, +1)$$

$$VEV = (0, 0, v)$$

# The scalar potential with explicit CPV

$$V_{3HDM} = V_0 + V_{Z_2}$$

$$V_0 = \sum_i^3 \left[ -\mu_i^2 (\phi_i^\dagger \phi_i) + \lambda_{ii} (\phi_i^\dagger \phi_i)^2 \right]$$

$$+ \sum_{i,j}^3 \left[ \lambda_{ij} (\phi_i^\dagger \phi_i) (\phi_j^\dagger \phi_j) + \lambda'_{ij} (\phi_i^\dagger \phi_j) (\phi_j^\dagger \phi_i) \right]$$

$$V_{Z_2} = -\mu_{12}^2 (\phi_1^\dagger \phi_2) + \lambda_1 (\phi_1^\dagger \phi_2)^2 + \lambda_2 (\phi_2^\dagger \phi_3)^2 + \lambda_3 (\phi_3^\dagger \phi_1)^2 + h.c.$$

$$+ \lambda_4 (\phi_3^\dagger \phi_1) (\phi_2^\dagger \phi_3) + \lambda_5 (\phi_1^\dagger \phi_2) (\phi_3^\dagger \phi_3) + \lambda_6 (\phi_1^\dagger \phi_2) (\phi_1^\dagger \phi_1)$$

$$+ \lambda_7 (\phi_1^\dagger \phi_2) (\phi_2^\dagger \phi_2) + \lambda_8 (\phi_3^\dagger \phi_1) (\phi_3^\dagger \phi_2) + h.c.$$

The  $Z_2$  symmetry

$$\phi_1 \rightarrow -\phi_1, \quad \phi_2 \rightarrow -\phi_2, \quad \phi_3 \rightarrow \phi_3, \quad \text{SM fields} \rightarrow \text{SM fields}$$

# The CP-mixed mass eigenstates

The doublet compositions

$$\phi_1 = \begin{pmatrix} H_1^+ \\ \frac{H_1^0 + iA_1^0}{\sqrt{2}} \end{pmatrix}, \quad \phi_2 = \begin{pmatrix} H_2^+ \\ \frac{H_2^0 + iA_2^0}{\sqrt{2}} \end{pmatrix}, \quad \phi_3 = \begin{pmatrix} G^+ \\ \frac{v+h+iG^0}{\sqrt{2}} \end{pmatrix}$$

The mass eigenstates

$$S_1 = \frac{\alpha H_1^0 + \alpha H_2^0 - A_1^0 + A_2^0}{\sqrt{2\alpha^2 + 2}}, \quad S_2 = \frac{-H_1^0 - H_2^0 - \alpha A_1^0 + \alpha A_2^0}{\sqrt{2\alpha^2 + 2}}$$

$$S_3 = \frac{\beta H_1^0 - \beta H_2^0 + A_1^0 + A_2^0}{\sqrt{2\beta^2 + 2}}, \quad S_4 = \frac{-H_1^0 + H_2^0 + \beta A_1^0 + \beta A_2^0}{\sqrt{2\beta^2 + 2}}$$

$$S_1^\pm = \frac{e^{\mp i\theta_{12}/2}}{\sqrt{2}} (H_2^\pm + H_1^\pm), \quad S_2^\pm = \frac{e^{\mp i\theta_{12}/2}}{\sqrt{2}} (H_2^\pm - H_1^\pm)$$

$S_1$  is assumed to be the DM candidate

# Input parameters and constraints

DM mass  $m_{S_1}$ , Mass splittings  $\delta_{S_2-S_1}$ ,  $\delta_{S_1^\pm-S_1}$ ,  $\delta_{S_2^\pm-S_1^\pm}$ ,  
Higgs-DM coupling  $g_{S_1 S_1 h}$ , CPV phases  $\theta_2$ ,  $\theta_{12}$

## Constraints taken into account include:

- Stability of the potential
- Positive-definiteness of the Hessian
- Limits from gauge bosons width:
- Limits on charged scalar mass and lifetime:
- Null DM collider searches excluding simultaneously:
- S,T,U parameters



# Relevant DM scenarios

In the low mass region ( $m_{S_1} < m_Z$ ):

- **Scenario A:** no coannihilation

$$m_{S_1} \ll m_{S_2}, m_{S_3}, m_{S_4}, m_{S_1^\pm}, m_{S_2^\pm}$$

- **Scenario B:** coannihilation with  $S_3$

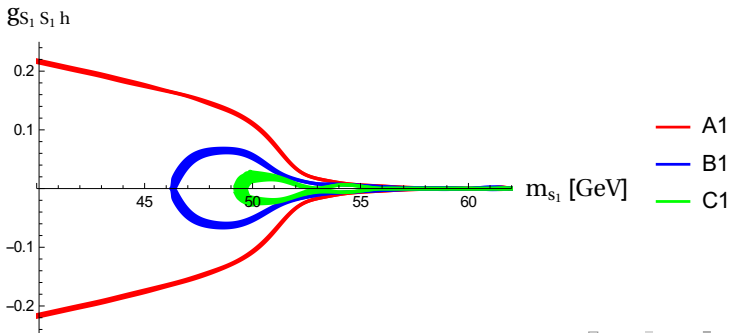
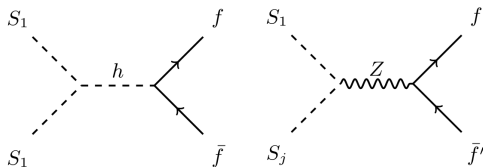
$$m_{S_1} \sim m_{S_3} \ll m_{S_2}, m_{S_4}, m_{S_1^\pm}, m_{S_2^\pm}$$

- **Scenario C:** coannihilation with all neutral particles

$$m_{S_1} \sim m_{S_3} \sim m_{S_2} \sim m_{S_4} \ll m_{S_1^\pm}, m_{S_2^\pm}$$

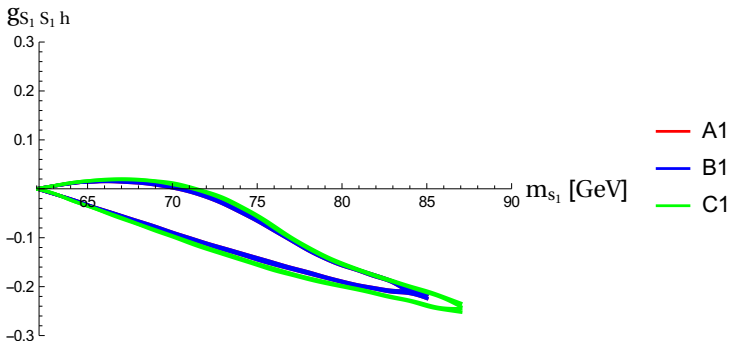
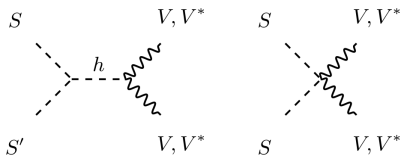
# Low DM mass region

## Higgs-mediated and Z-mediated (co)annihilation



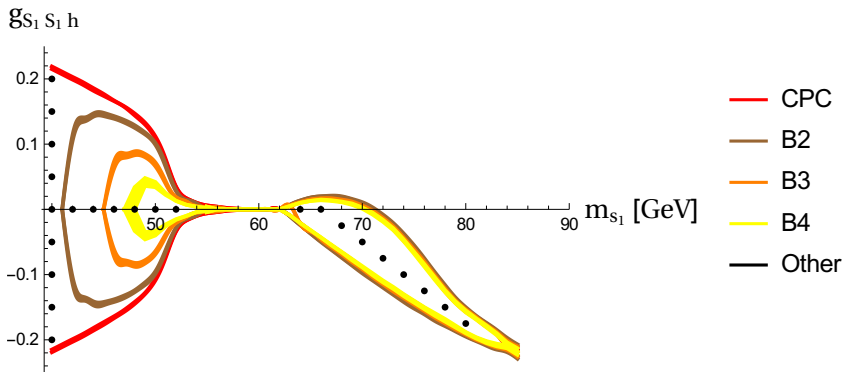
# Medium DM mass region

## Higgs-mediated and quartic (co)annihilation



# Filling the plot

C-type scenarios are the winners!



# Relevant DM scenarios

In the heavy mass region ( $m_{S_1} > 400$  GeV):

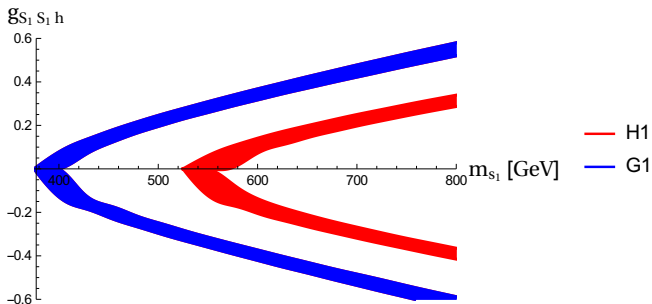
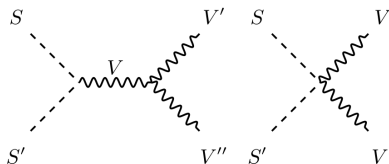
- **Scenario G:** coannihilation within the "family"

$$m_{S_1} \sim m_{S_3} \sim m_{S_1^\pm} \ll m_{S_2} \sim m_{S_4} \sim m_{S_2^\pm}$$

- **Scenario H:** coannihilation with all inert particles

$$m_{S_1} \sim m_{S_3} \sim m_{S_2} \sim m_{S_4} \sim m_{S_1^\pm} \sim m_{S_2^\pm}$$

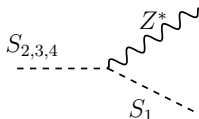
# Heavy DM mass region



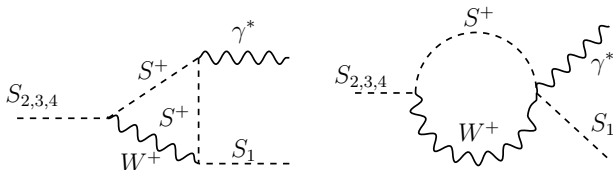
# Collider signatures of CPV DM

# Inert cascade decays at the LHC

When there is a **large mass splitting** between DM and other inert particles:



When there is a **small mass splitting** between DM and other inert particles (winning scenarios):





# $E_{miss}^T + e^+e^-$ signals at the LHC

- Higgs-strahlung at **tree level**:

$$q\bar{q} \rightarrow Z \rightarrow S_1 S_{2,3,4} \rightarrow S_1 S_1 Z^* \rightarrow S_1 S_1 e^+ e^-$$

- Higgs-strahlung at **loop level**:

$$q\bar{q} \rightarrow Z \rightarrow S_1 S_{2,3,4} \rightarrow S_1 S_1 \gamma^* \rightarrow S_1 S_1 e^+ e^-$$

- Gluon-fusion at **tree level**:

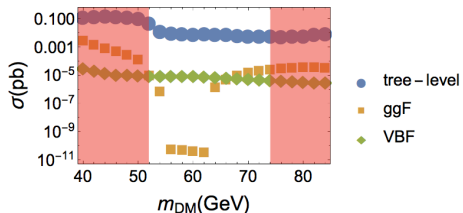
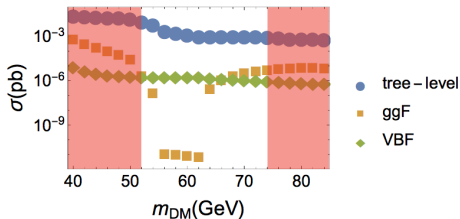
$$pp \rightarrow h \rightarrow S_1 S_{2,3,4} \rightarrow S_1 S_1 Z^* \rightarrow S_1 S_1 e^+ e^-$$

- Vector boson fusion at **loop level**:

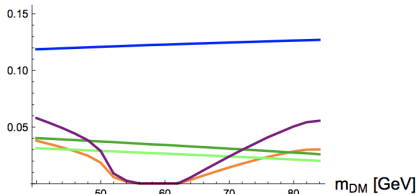
$$q_i q_j \rightarrow q_k q_l S_1 S_{2,3,4} \rightarrow q_k q_l S_1 S_1 \gamma^* \rightarrow q_k q_l S_1 S_1 e^+ e^-$$

Benchmark	$m_{H_2} - m_{H_1}$	$m_{A_1} - m_{H_1}$	$m_{A_2} - m_{H_1}$	$m_{H_1^\pm} - m_{H_1}$	$m_{H_2^\pm} - m_{H_1}$
A50	50	75	125	75	125
I5	5	10	15	90	95
I10	10	20	30	90	100

## Scenario A50

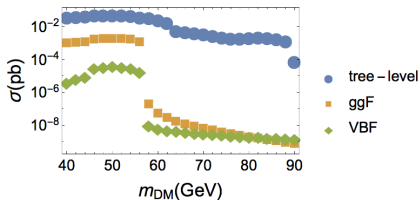
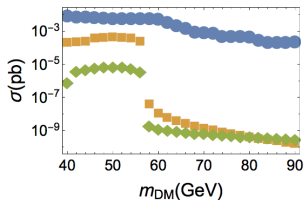


Scenario A50

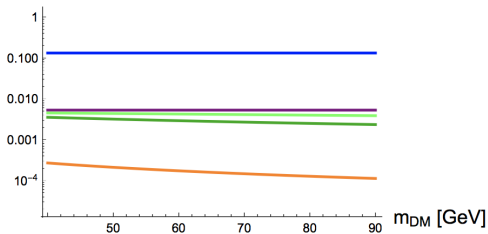


- $|g_{\text{ZH1A1}}^2|$  (tree level)
- $|g_{\text{hH1H2}}|$  (ggF)
- $|g_{\text{ZH1A1}} * g_{\text{ZH2A1}}|$  (neutral VBF)
- $|g_{\text{W}^+ \text{H1}^- \text{H1}} * g_{\text{W}^- \text{H1}^+ \text{H2}}|$  (charged VBF)
- $|g_{\text{hDM}} * g_{\text{hzz}}|$

## Scenario I10

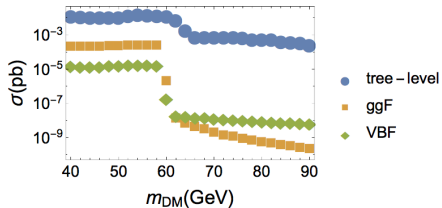
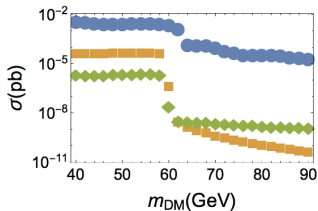


## Scenario I10

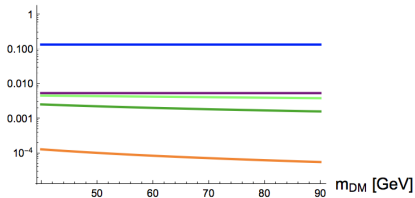


- $|g_{\text{ZH1A1}}^2|$  (tree level)
- $|g_{\text{hH1H2}}|$  (ggF)
- $|g_{\text{ZH1A1}} * g_{\text{ZH2A1}}|$  (neutral VBF)
- $|g_{\text{W}^+ \text{H1}^- \text{H1}} * g_{\text{W}^- \text{H1}^+ \text{H2}}|$  (charged VBF)
- $|g_{\text{hDM}} * g_{\text{hZZ}}|$

## Scenario I5



Scenario I5

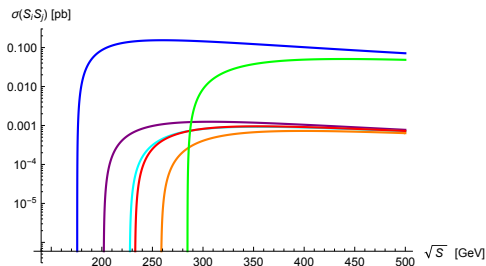


- $|g_{\text{ZH1A1}^2}|$  (tree level)
- $|g_{\text{hH1H2}}|$  (ggF)
- $|g_{\text{ZH1A1}^*g_{\text{ZH2A1}}|$  (neutral VBF)
- $|g_{\text{W}^+ \text{H1}^- \text{H1}^* g_{\text{W}^- \text{H1}^+ \text{H2}}|$  (charged VBF)
- $|g_{\text{hDM}}*g_{\text{hzz}}|$

# Production thresholds of the $Z$ boson

- The process  $e^+e^- \rightarrow Z^* \rightarrow S_i S_j$  ( $i, j = 1, \dots, 4$ )
- A simple collider energy scan combined with a trivial counting experiment in the detectors
- Three benchmark points with the cross section as large as a few picobarns at  $\sqrt{s}$  values accessible by future  $e^+e^-$  colliders
- The proximity (or otherwise) of these thresholds would serve as characteristic signatures of different BPs with different DM properties.

## Type A

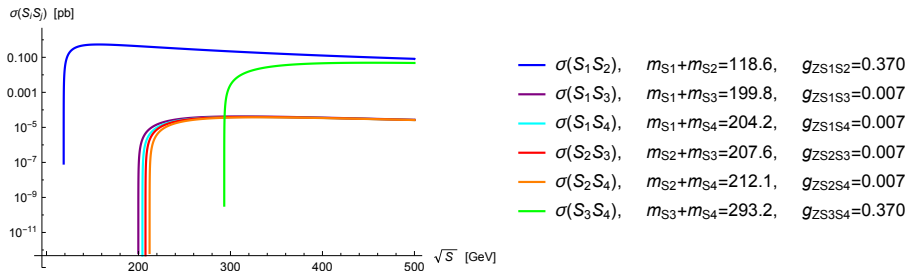


— $\sigma(S_1 S_2)$ ,	$m_{S_1} + m_{S_2} = 175.6$ ,	$g_{ZS_1 S_2} = 0.366$
— $\sigma(S_1 S_3)$ ,	$m_{S_1} + m_{S_3} = 201.8$ ,	$g_{ZS_1 S_3} = 0.0397$
— $\sigma(S_1 S_4)$ ,	$m_{S_1} + m_{S_4} = 227.5$ ,	$g_{ZS_1 S_4} = 0.0401$
— $\sigma(S_2 S_3)$ ,	$m_{S_2} + m_{S_3} = 232.8$ ,	$g_{ZS_2 S_3} = 0.04006$
— $\sigma(S_2 S_4)$ ,	$m_{S_2} + m_{S_4} = 258.5$ ,	$g_{ZS_2 S_4} = 0.0397$
— $\sigma(S_3 S_4)$ ,	$m_{S_3} + m_{S_4} = 284.6$ ,	$g_{ZS_3 S_4} = 0.366$

$$m_{S_1} = 72.331 \text{ GeV}, \quad m_{S_2} = 103.313 \text{ GeV}, \quad m_{S_1^\pm} = 106.235 \text{ GeV},$$

$$m_{S_3} = 129.467 \text{ GeV}, \quad m_{S_4} = 155.178 \text{ GeV}, \quad m_{S_2^\pm} = 157.588 \text{ GeV}.$$

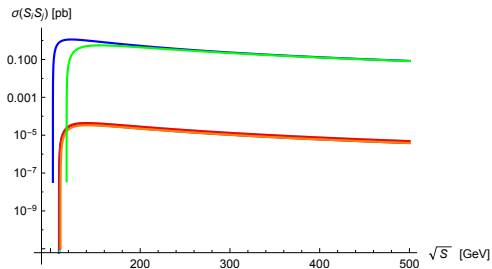
## Type B



$$m_{S_1} = 55.441 \text{ GeV}, \quad m_{S_2} = 63.219 \text{ GeV}, \quad m_{S_1^\pm} = 79.184 \text{ GeV},$$

$$m_{S_3} = 144.377 \text{ GeV}, \quad m_{S_4} = 148.842 \text{ GeV}, \quad m_{S_2^\pm} = 159.203 \text{ GeV}.$$

## Type C



—	$\sigma(S_1 S_2)$ ,	$m_{S_1} + m_{S_2} = 102.7$ ,	$g_{ZS_1 S_2} = 0.370$
—	$\sigma(S_1 S_3)$ ,	$m_{S_1} + m_{S_3} = 109.4$ ,	$g_{ZS_1 S_3} = 0.0025$
—	$\sigma(S_1 S_4)$ ,	$m_{S_1} + m_{S_4} = 110.4$ ,	$g_{ZS_1 S_4} = 0.0028$
—	$\sigma(S_2 S_3)$ ,	$m_{S_2} + m_{S_3} = 110.3$ ,	$g_{ZS_2 S_3} = 0.0028$
—	$\sigma(S_2 S_4)$ ,	$m_{S_2} + m_{S_4} = 111.2$ ,	$g_{ZS_2 S_4} = 0.0025$
—	$\sigma(S_3 S_4)$ ,	$m_{S_3} + m_{S_4} = 118.0$ ,	$g_{ZS_3 S_4} = 0.370$

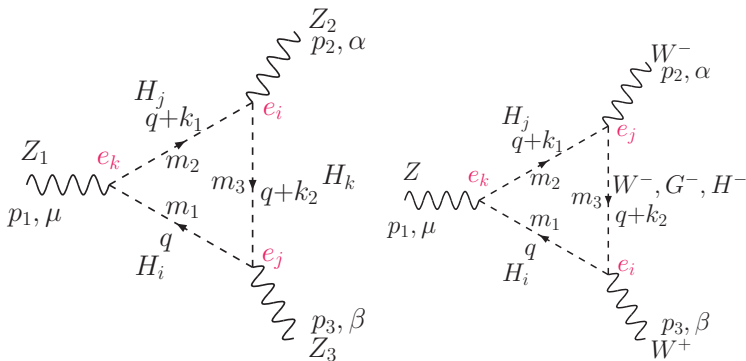
$$m_{S_1} = 50.925 \text{ GeV}, \quad m_{S_2} = 51.793 \text{ GeV}, \quad m_{S_1^\pm} = 99.176 \text{ GeV},$$

$$m_{S_3} = 58.555 \text{ GeV}, \quad m_{S_4} = 59.459 \text{ GeV}, \quad m_{S_2^\pm} = 111.136 \text{ GeV}.$$



# Other CPV observables

## ZZZ and ZWW vertices



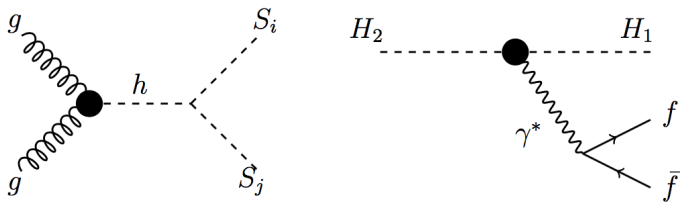
JHEP1605,025(2016)

# Summary

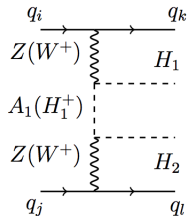
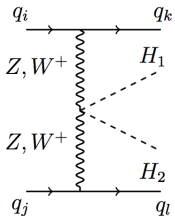
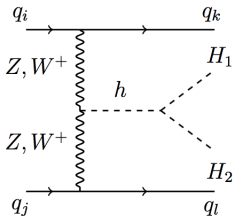
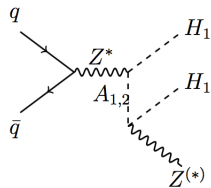
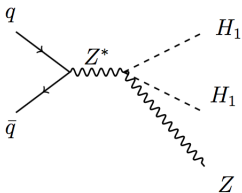
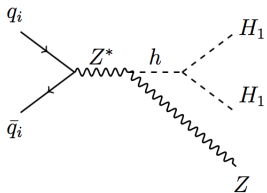
- CP-Violation in **I(1+2)HDM**
  - IDM-like inert sector: **CPC DM**
  - CPV in the active sector:  $\tilde{H}_1, \tilde{H}_2, \tilde{H}_3$
  - Interesting LHC phenomenology, however, very limited CPV
- CP-Violation in **I(2+1)HDM**
  - SM-like active sector:  $H_3 \equiv h^{SM}$
  - Unbounded CPV in the inert sector:  $H_{1,2}, A_{1,2} \rightarrow S_{1,2,3,4}$  **CPV DM**
  - opens up new regions of parameter space
  - New observables at the LHC:  $S_i S_j Z$  vertices

# BACKUP SLIDES

# ggF processes and the effective vertex



# HS and VBF processes



# LHC bounds on CPV DM

# Higgs invisible branching ratio and total decay

From ATLAS and CMS

$$\text{Br}(h \rightarrow \text{inv}) < 0.23 - 0.36$$

for  $m_{i,j} < m_h/2$  if long lived

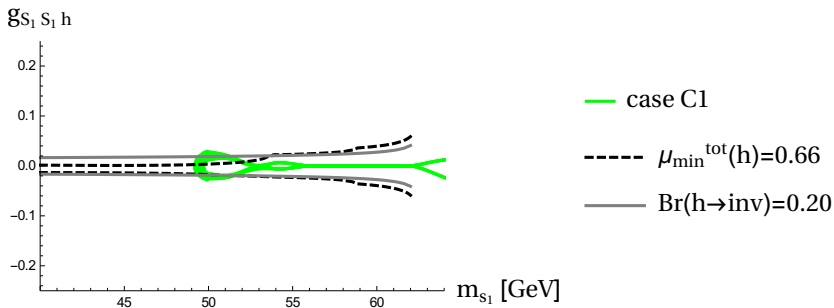
$$\text{BR}(h \rightarrow \text{inv}) = \frac{\sum_{i,j} \Gamma(h \rightarrow S_i S_j)}{\Gamma_h^{\text{SM}} + \sum_i \Gamma(h \rightarrow S_i S_j)}$$

The **total decay** signal strength

$$\mu_{\text{tot}} = \frac{\text{BR}(h \rightarrow \text{XX})}{\text{BR}(h_{\text{SM}} \rightarrow \text{XX})} = \frac{\Gamma_{\text{tot}}^{\text{SM}}(h)}{\Gamma_{\text{tot}}^{\text{SM}}(h) + \Gamma^{\text{inert}}(h)}$$

We use  $\mu_{\text{tot}} = 1.17 \pm 0.17$  at  $3\sigma$  level.

# Relic density vs. Higgs decay bounds





# $h \rightarrow \gamma\gamma$ signal strength bounds

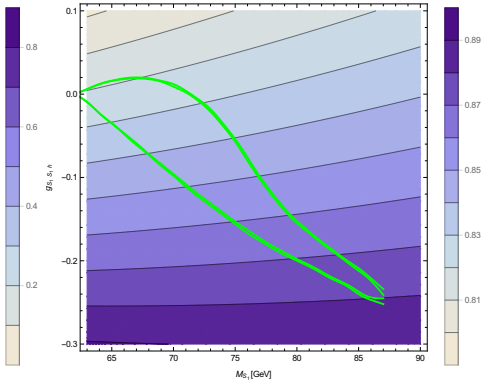
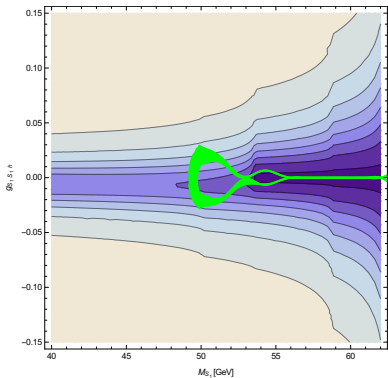
From ATLAS and CMS:  $\mu_{\gamma\gamma} = 1.16^{+0.20}_{-0.18}$

$$\mu_{\gamma\gamma} = \frac{\Gamma(h \rightarrow \gamma\gamma)^{3\text{HDM}} \Gamma(h)^{\text{SM}}}{\Gamma(h \rightarrow \gamma\gamma)^{\text{SM}} \Gamma(h)^{3\text{HDM}}}$$

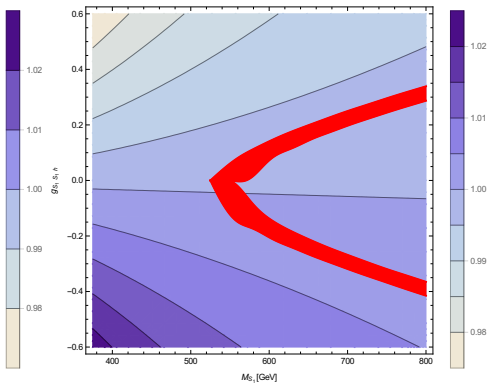
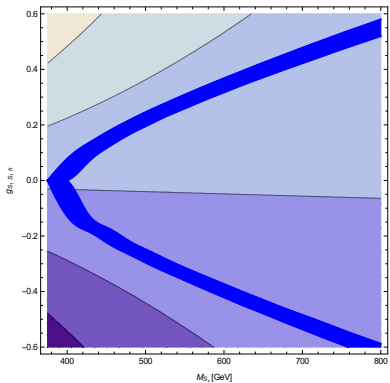
Modified by

- charged scalars contribution to  $\Gamma(h \rightarrow \gamma\gamma)^{3\text{HDM}}$
- light neutral scalars contribution to  $\Gamma(h)^{3\text{HDM}}$

# Relic density vs. $\mu_{\gamma\gamma}$ - scenario C



# Relic density vs. $\mu_{\gamma\gamma}$ - scenarios G & H



# Parameters of the model

- no new phenomenology from  $\lambda_4, \dots, \lambda_8$  terms  $\rightarrow \lambda_{4-8} = 0$
- “dark” parameters  $\lambda_1, \lambda_{11}, \lambda_{22}, \lambda_{12}, \lambda'_{12}$
- “dark democracy” limit  
 $\mu_1^2 = \mu_2^2, \quad \lambda_3 = \lambda_2, \quad \lambda_{31} = \lambda_{23}, \quad \lambda'_{31} = \lambda'_{23}$
- fixed by the Higgs mass  $\mu_3^2 = v^2 \lambda_{33} = m_h^2/2$

## 7 important parameters

- CPV and mass splittings  $\mu_{12}^2 = |\mu_{12}^2| e^{i\theta_{12}}, \quad \lambda_2 = |\lambda_2| e^{i\theta_2}$
- Higgs-DM coupling  $\lambda_2, \lambda_{23}, \lambda'_{23}$
- Mass scale of inert particles  $\mu_2^2$

In the CPC limit

$$\alpha = \frac{-|\mu_{12}^2| \cos \theta_{12} + v^2 |\lambda_2| \cos \theta_2 - \Lambda}{|\mu_{12}^2| \sin \theta_{12} + v^2 |\lambda_2| \sin \theta_2} \rightarrow \infty$$

$$\beta = \frac{|\mu_{12}^2| \cos \theta_{12} + v^2 |\lambda_2| \cos \theta_2 - \Lambda'}{|\mu_{12}^2| \sin \theta_{12} - v^2 |\lambda_2| \sin \theta_2} \rightarrow \infty$$

where

$$\Lambda = \sqrt{v^4 |\lambda_2|^2 + |\mu_{12}^2|^2 - 2v^2 |\lambda_2| |\mu_{12}^2| \cos(\theta_{12} + \theta_2)},$$

$$\Lambda' = \sqrt{v^4 |\lambda_2|^2 + |\mu_{12}^2|^2 + 2v^2 |\lambda_2| |\mu_{12}^2| \cos(\theta_{12} + \theta_2)}.$$

# Benchmark scenarios

$$A1 : \delta_{12} = 125 \text{ GeV}, \delta_{1c} = 50 \text{ GeV}, \delta_c = 50 \text{ GeV}, \theta_2 = \theta_{12} = 1.5$$

$$B1 : \delta_{12} = 125 \text{ GeV}, \delta_{1c} = 50 \text{ GeV}, \delta_c = 50 \text{ GeV}, \theta_2 = \theta_{12} = 0.82$$

$$C1 : \delta_{12} = 12 \text{ GeV}, \delta_{1c} = 100 \text{ GeV}, \delta_c = 1 \text{ GeV}, \theta_2 = \theta_{12} = 1.57$$

$$G1 : \delta_{12} = 2 \text{ GeV}, \delta_{1c} = 1 \text{ GeV}, \delta_c = 1 \text{ GeV}, \theta_2 = \theta_{12} = 0.82$$

$$H1 : \delta_{12} = 50 \text{ GeV}, \delta_{1c} = 1 \text{ GeV}, \delta_c = 50 \text{ GeV}, \theta_2 = \theta_{12} = 0.82$$