



Physics prospects of future hadron colliders

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pp @ 14 TeV, 3ab⁻¹

**✓ Approved
2026-37**



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[link to CDR](#)

100km tunnel

- **e⁺e⁻ @ 91, 160, 240, 365 GeV**
- **pp @ 100 TeV**
- **e_{60GeV} p_{50TeV} @ 3.5 TeV**

LHC tunnel: HE-LHC

- **pp @ 27 TeV, 15ab⁻¹**

Lecture 2 by B. Ferrando



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100km tunnel

- e^+e^- @ 91, 240 GeV (but possibly 160 & 350)
- Future possible pp @ ~ 70 TeV and $e_{60\text{GeV}} p_{35\text{TeV}}$

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Readiness to address both scenarios is the best hedge for the field:

- *precision*
- *sensitivity (to elusive signatures)*
- *extended energy/mass reach*

Remark

the discussion of the **future** in HEP must start from the understanding that there is no experiment/facility, proposed or conceivable, in the lab or in space, accelerator or non-accelerator driven, which can **guarantee discoveries** beyond the SM, and **answers** to the big questions of the field

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- Exploration potential:
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 - *E.g. match the mass scales for new physics that could be exposed via indirect precision measurements in the EW and Higgs sector*

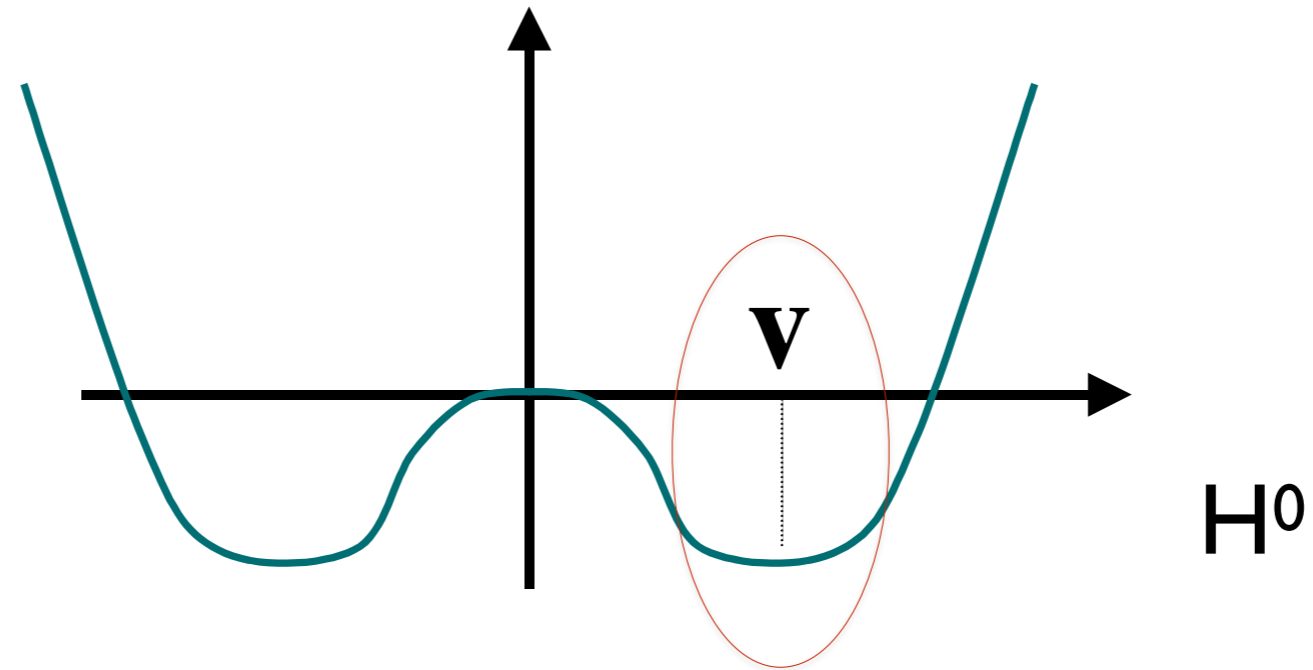
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 - *E.g. match the mass scales for new physics that could be exposed via indirect precision measurements in the EW and Higgs sector*
- Provide firm Yes/No answers to questions like:
 - is there a TeV-scale solution to the hierarchy problem?
 - is DM a thermal WIMP?
 - could the cosmological EW phase transition have been 1st order?
 - could baryogenesis have taken place during the EW phase transition?
 - could neutrino masses have their origin at the TeV scale?
 - ...

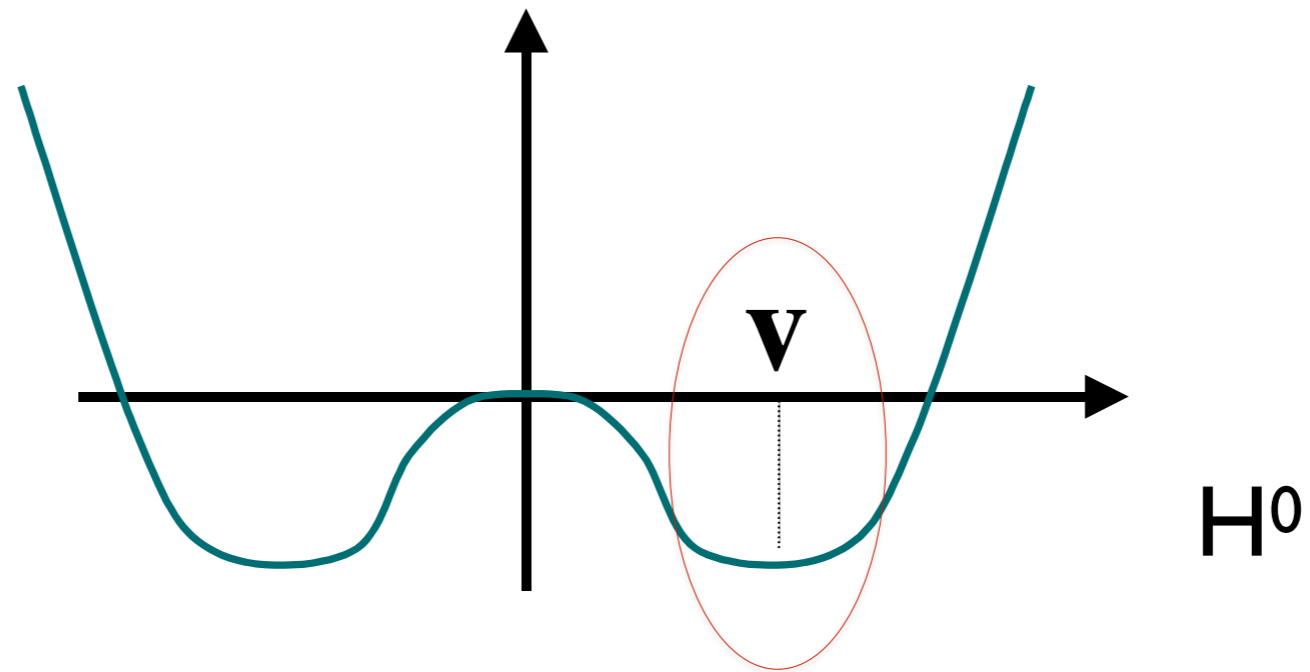
Guaranteed deliverables:

what more will we need to know about the Higgs after the HL-LHC? will it not get “boring” to keep studying the Higgs and the top?



$$V(H) = -\mu^2 |H|^2 + \lambda |H|^4$$

Who ordered that ?

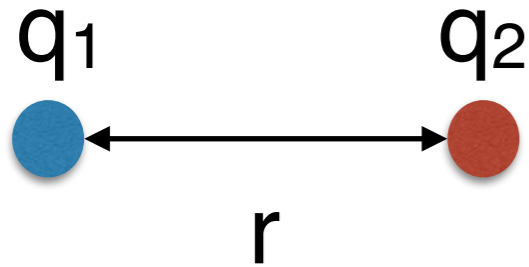


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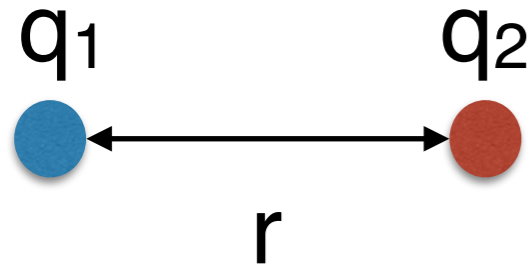
We must learn to appreciate the depth and the value of this question, which is set to define the future of collider physics

Electromagnetic vs Higgs dynamics



$$V(r) = + \frac{q_1 \times q_2}{r^1}$$

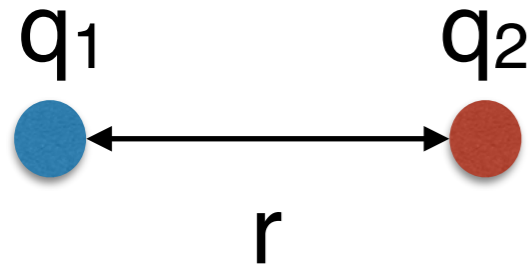
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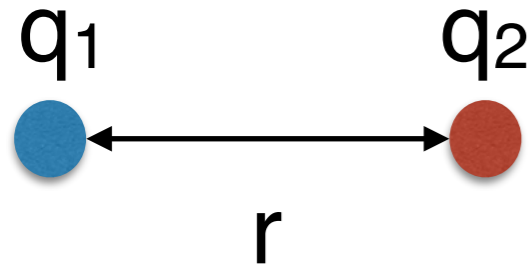
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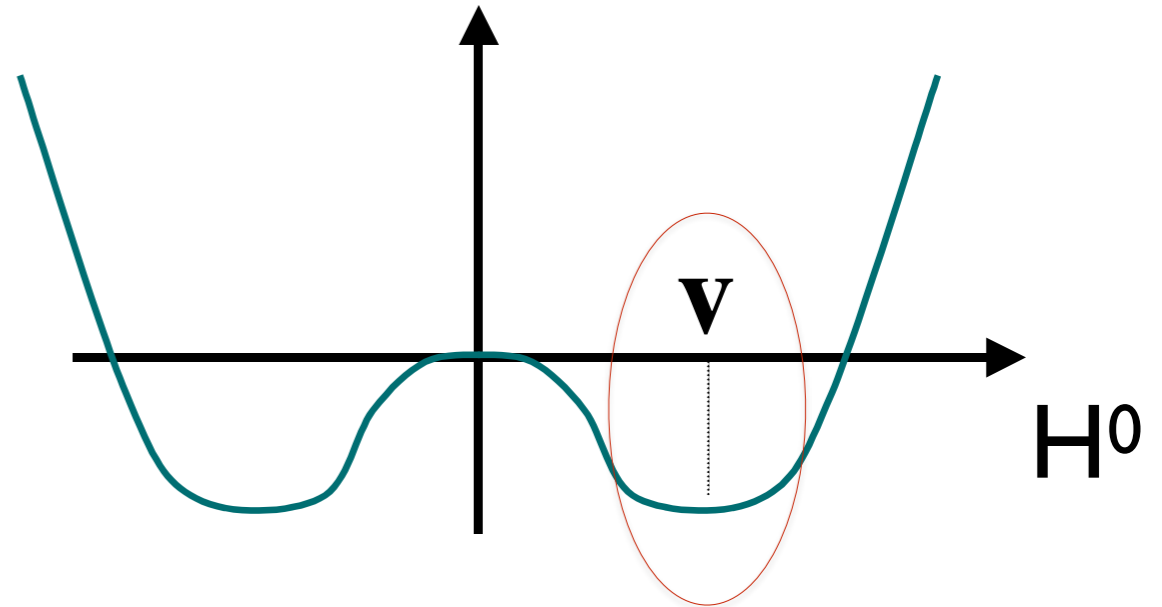


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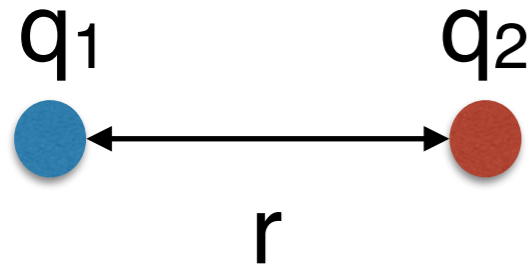
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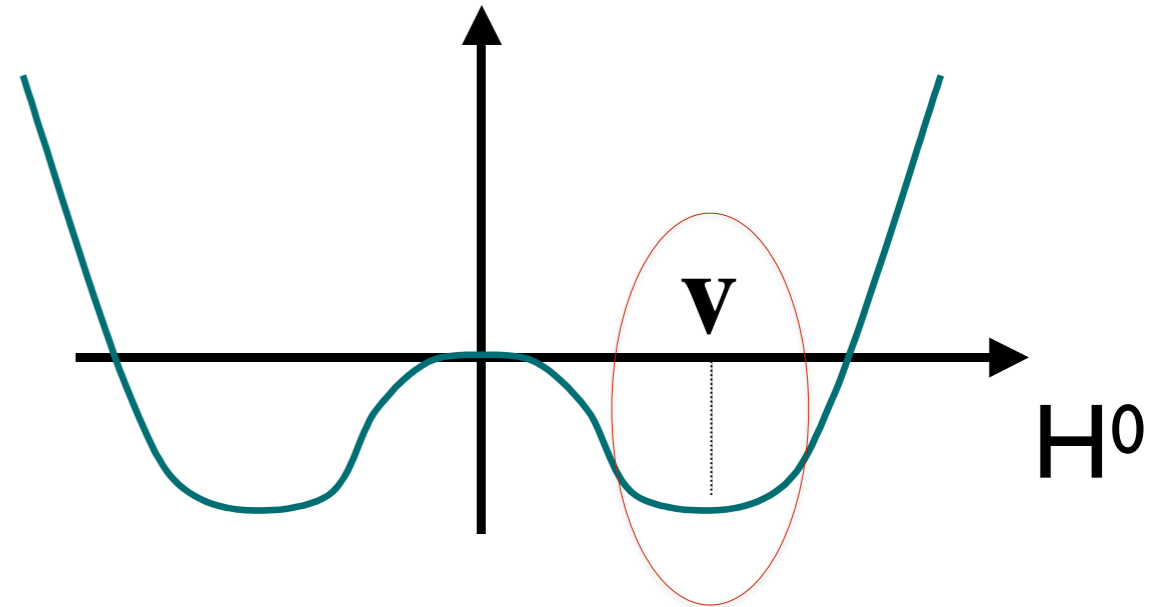


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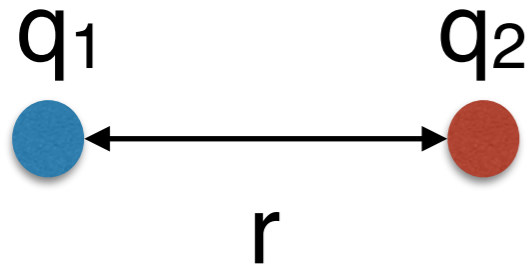


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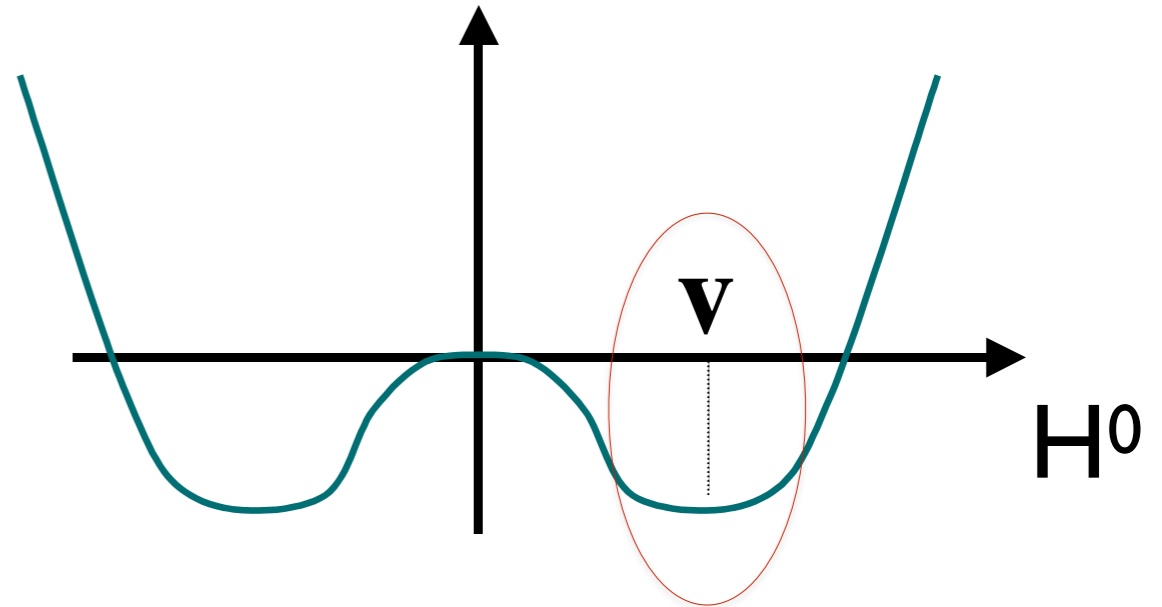


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any function of $|H|^2$ would be
ok wrt known symmetries

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- The relation between the Higgs phenomenon and the SM is similar to the relation between superconductivity and the Landau-Ginzburg theory of phase transitions: a quartic potential for a bosonic order parameter, with negative quadratic term, and the ensuing symmetry breaking. If superconductivity had been discovered after Landau-Ginzburg, we would be in a similar situations as we are in today: an experimentally proven phenomenological model. But we would still lack a deep understanding of the relevant dynamics.

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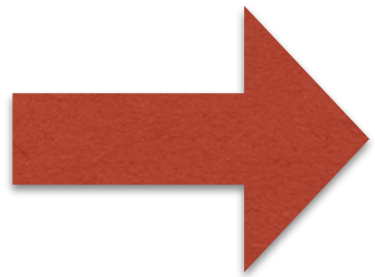
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- For superconductivity, this came later, with the identification of e^-e^- Cooper pairs as the underlying order parameter, and BCS theory. In particle physics, we still don't know whether the Higgs is built out of some sort of Cooper pairs (composite Higgs) or whether it is elementary, and in both cases we have no clue as to what is the dynamics that generates the Higgs potential. With Cooper pairs it turned out to be just EM and phonon interactions. With the Higgs, none of the SM interactions can do this, and **we must look beyond.**

examples of possible scenarios

- **BCS-like**: the Higgs is a composite object
- **Supersymmetry**: the Higgs is a fundamental field and
 - $\lambda^2 \sim g^2 + g'^2$, it is not arbitrary (MSSM, w/out susy breaking, has one parameter less than SM!)
 - potential is fixed by susy & gauge symmetry
 - EW symmetry breaking (and thus m_H and λ) determined by the parameters of SUSY breaking
- ...

furthermore ...

Hierarchy problem and naturalness !!



Lecture 2 by M. Mc Cullough

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=> colliders are the only facilities that make this possible

Other important open issues on the Higgs sector

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- Is the Higgs the only (fundamental?) scalar field, or are there other Higgs-like states (e.g. $H^\pm, A^0, H^{\pm\pm}, \dots$, EW-singlets,) ?
 - Do all SM families get their mass from the **same** Higgs field?
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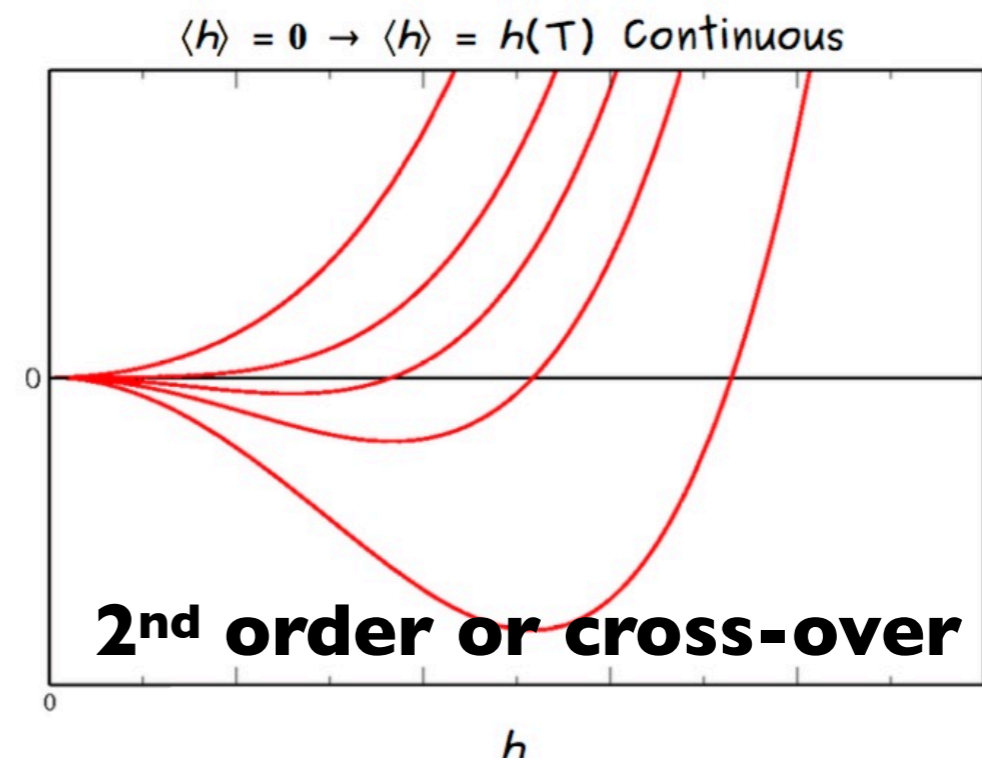
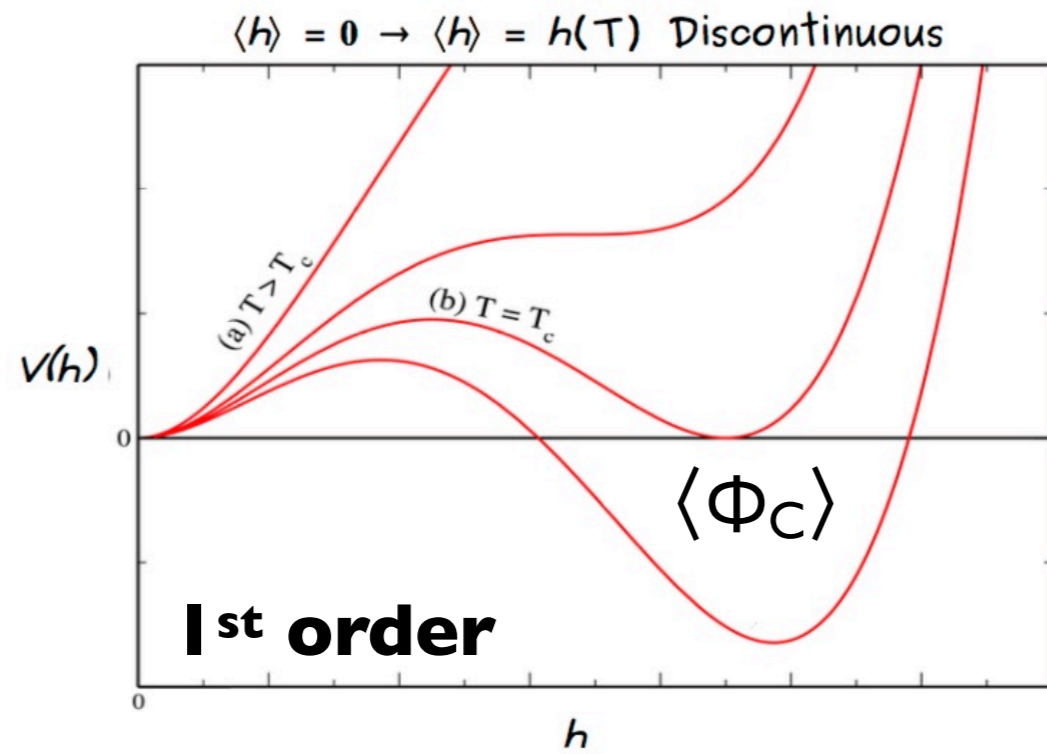
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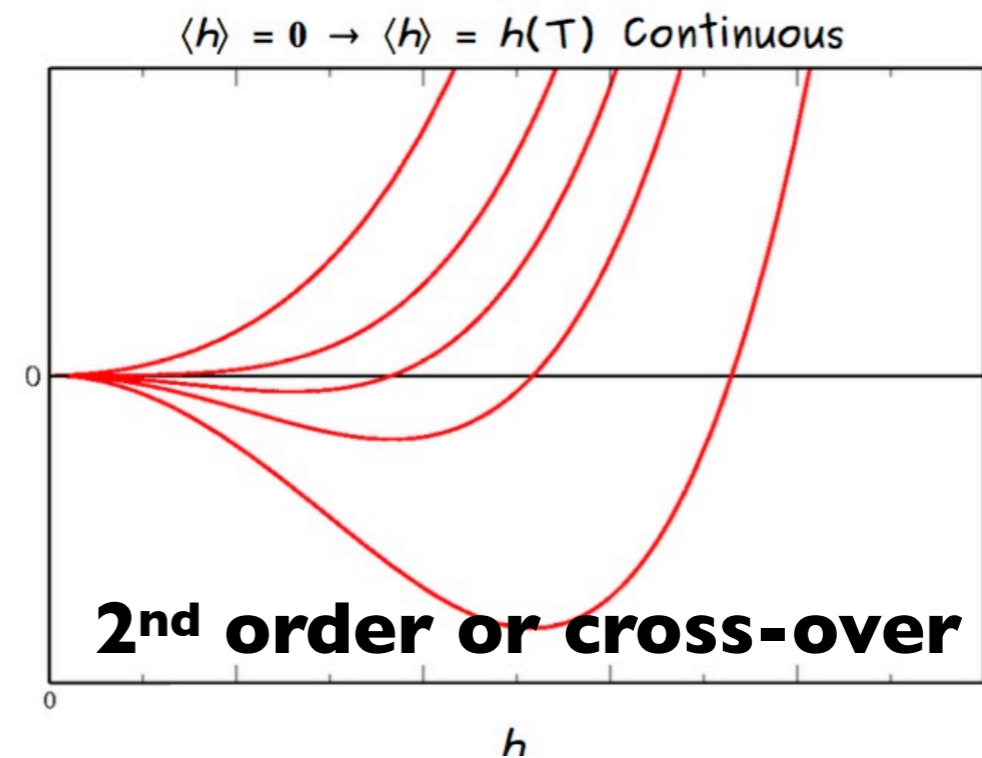
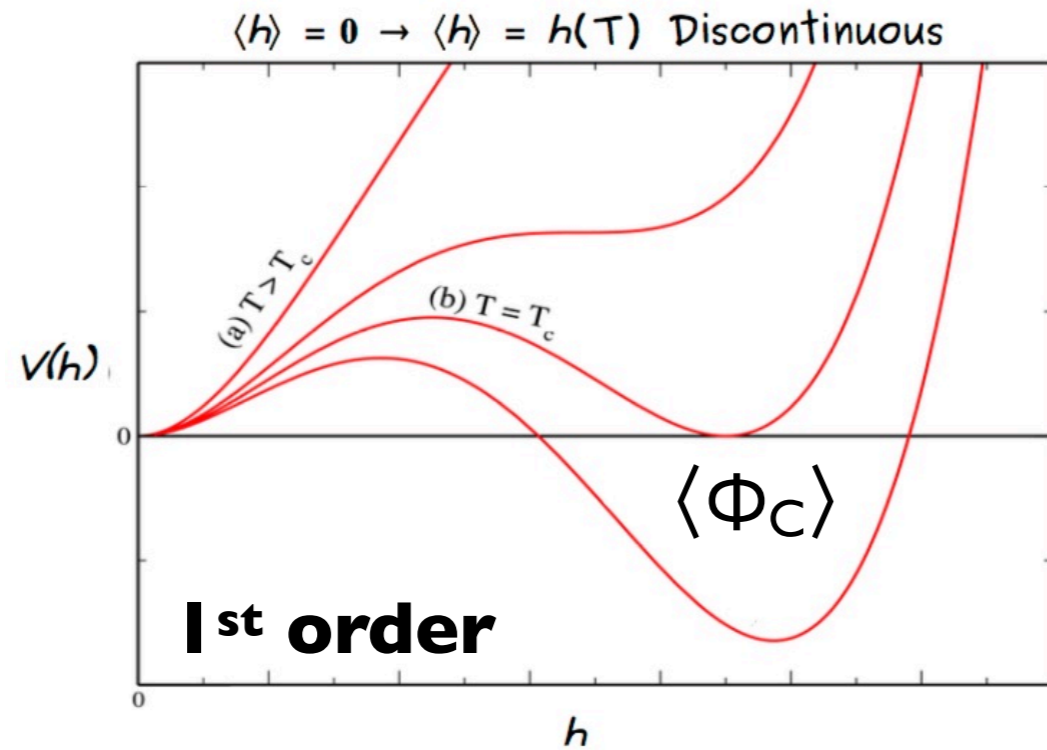
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- What happens at the EW phase transition (PT) during the Big Bang?
 - what's the order of the phase transition?
 - are the conditions realized to allow EW baryogenesis?

The nature of the EW phase transition



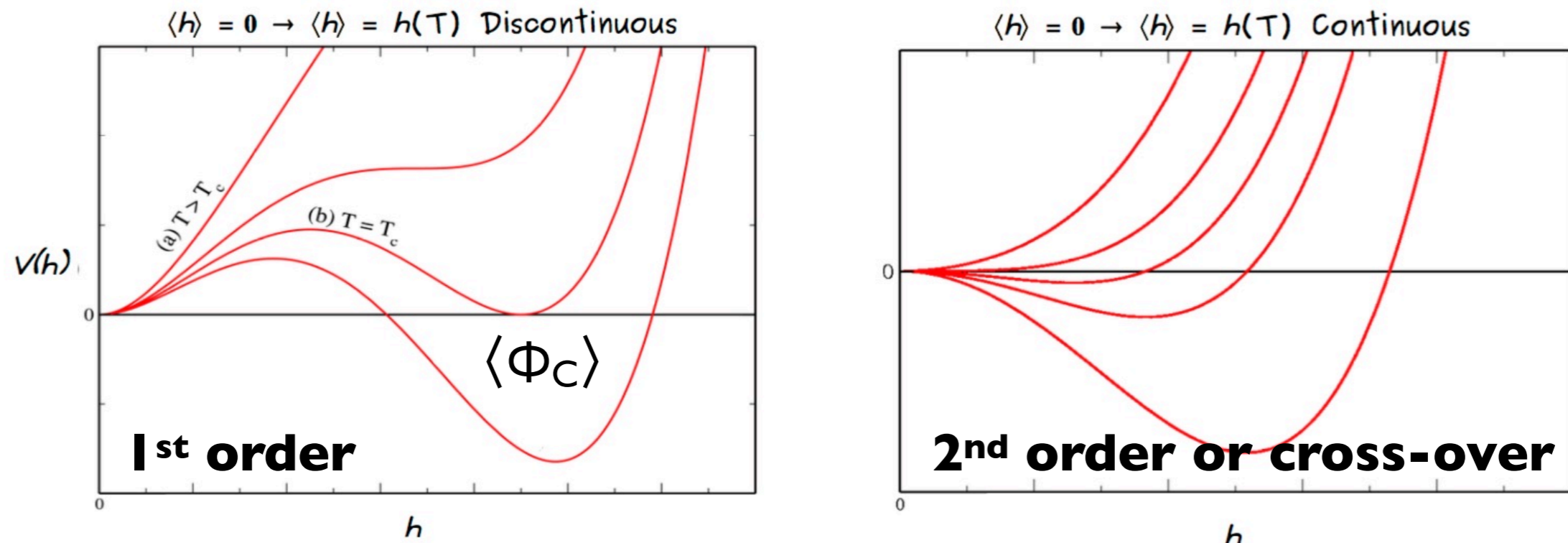
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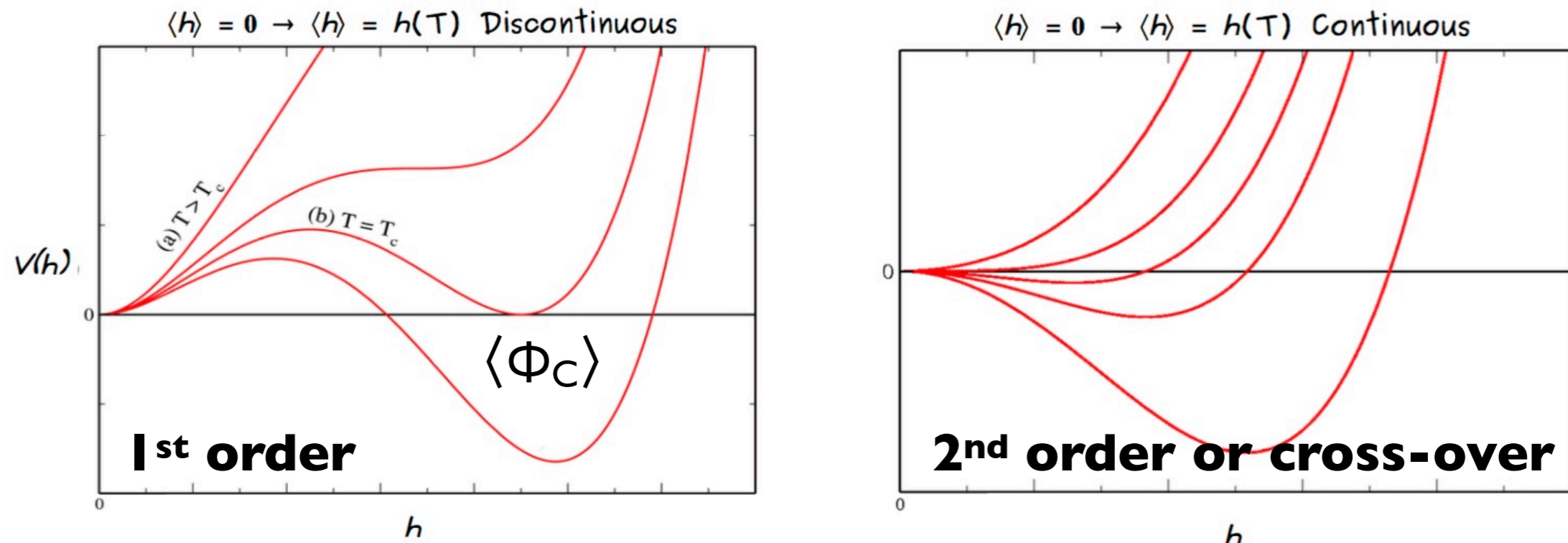
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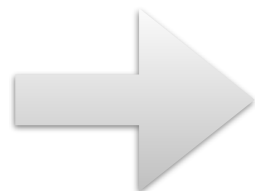


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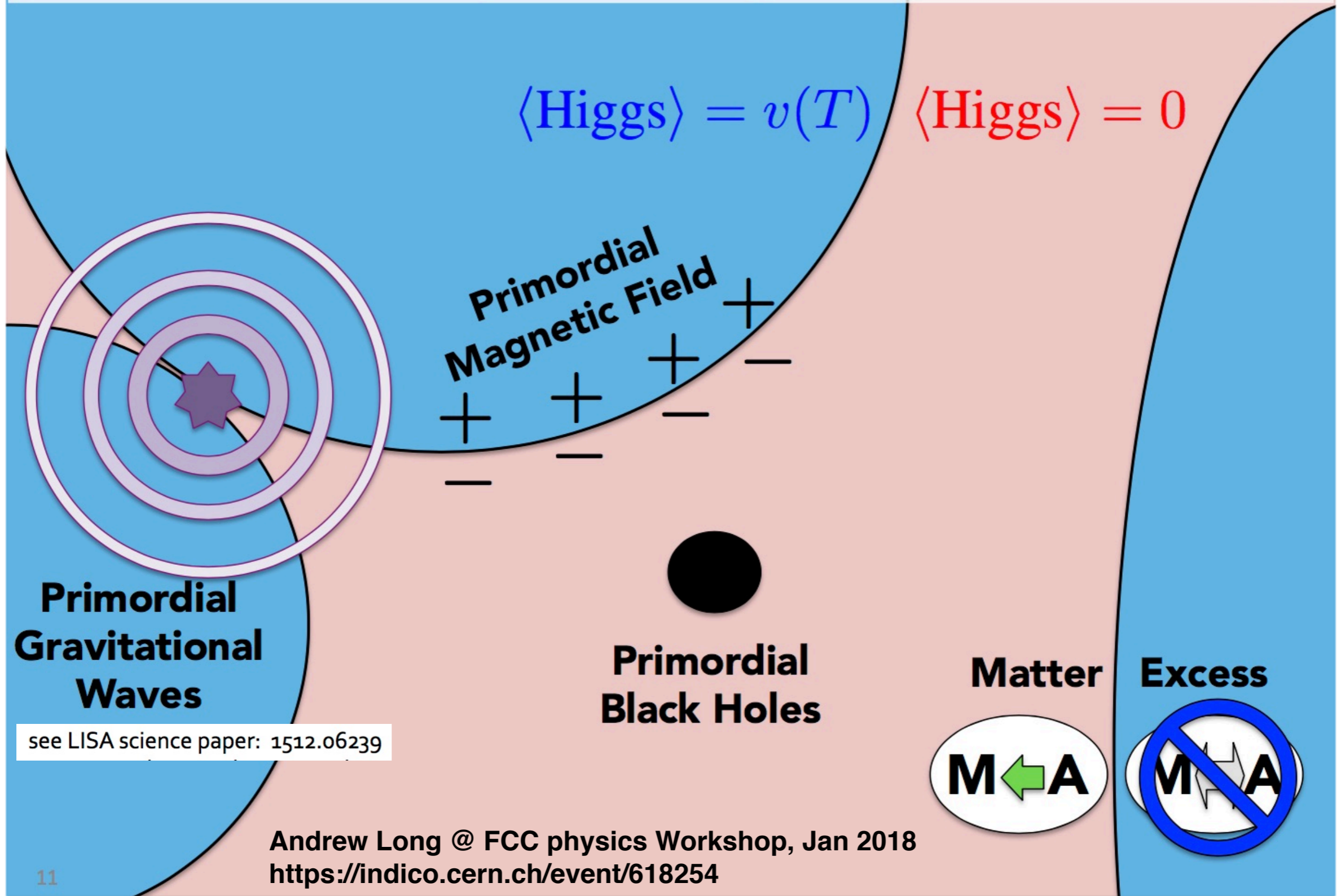
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- Probe higher-order terms of the Higgs potential (selfcouplings)
- Probe the existence of other particles coupled to the Higgs

1st Order EWPT has profound implications for cosmology



Andrew Long @ FCC physics Workshop, Jan 2018
<https://indico.cern.ch/event/618254>

Higgs couplings, beyond the HL-LHC: the e^+e^- phase

Collider	HL-LHC update	ILC ₂₅₀	CLIC ₃₈₀	LEP3 ₂₄₀	CEPC ₂₅₀	FCC-ee ₂₄₀₊₃₆₅		
Lumi (ab^{-1})	3	2	0.5	3	5	5 ₂₄₀	+1.5 ₃₆₅	+ HL-LHC
Years	25	15	7	6	7	3	+4	
$\delta\Gamma_{\text{H}}/\Gamma_{\text{H}}$ (%)	50	3.6	6.3	3.6	2.6	2.7	1.3	1.1
$\delta g_{\text{HZZ}}/g_{\text{HZZ}}$ (%)	1.5	0.3	0.40	0.32	0.25	0.20	0.17	0.16
$\delta g_{\text{HWW}}/g_{\text{HWW}}$ (%)	1.7	1.7	0.8	1.7	1.2	1.3	0.43	0.40
$\delta g_{\text{Hbb}}/g_{\text{Hbb}}$ (%)	3.7	1.7	1.3	1.8	1.3	1.3	0.61	0.56
$\delta g_{\text{Hcc}}/g_{\text{Hcc}}$ (%)	SM	2.3	4.1	2.3	1.8	1.7	1.21	1.18
$\delta g_{\text{Hgg}}/g_{\text{Hgg}}$ (%)	2.5	2.2	2.1	2.1	1.4	1.6	1.01	0.90
$\delta g_{\text{H}\tau\tau}/g_{\text{H}\tau\tau}$ (%)	1.9	1.9	2.7	1.9	1.4	1.4	0.74	0.67
$\delta g_{\text{H}\mu\mu}/g_{\text{H}\mu\mu}$ (%)	4.3	14.1	n.a.	12	6.2	10.1	9.0	3.8
$\delta g_{\text{H}\gamma\gamma}/g_{\text{H}\gamma\gamma}$ (%)	1.8	6.4	n.a.	6.1	4.7	4.8	3.9	1.3
$\delta g_{\text{H}tt}/g_{\text{H}tt}$ (%)	3.4	–	–	–	–	–	–	3.1
BR _{EXO} (%)	SM	< 1.7	< 3.0	< 1.6	< 1.2	< 1.2	< 1.0	< 1.0

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$\delta g_{\text{H}tt}/g_{\text{H}tt}$ (%)	3.4	–	–	–	–	–	–	3.1
BR _{EXO} (%)	SM	< 1.7	< 3.0	< 1.6	< 1.2	< 1.2	< 1.0	< 1.0

Table 1: Relative statistical uncertainty on the Higgs boson couplings and total decay width, as expected from the FCC-ee data, and compared to those from HL-LHC and other e^+e^- colliders exploring the 240-to-380 GeV centre-of-mass energy range. All numbers indicate 68% CL intervals, except for the last line which gives the 95% CL sensitivity on the "exotic" branching fraction, accounting for final states that cannot be tagged as SM decays. The FCC-ee accuracies are subdivided in three categories: the first sub-column give the results of the model-independent fit expected with 5 ab^{-1} at 240 GeV, the second sub-column in bold – directly comparable to the other collider fits – includes the additional 1.5 ab^{-1} at $\sqrt{s} = 365 \text{ GeV}$, and the last sub-column shows the result of the combined fit with HL-LHC. The fit to the HL-LHC projections alone (first column) requires two additional assumptions to be made: here, the branching ratios into $c\bar{c}$ and into exotic particles are set to their SM values.

* M. Cepeda, S. Gori, P. J. Ilten, M. Kado, and F. Riva, (conveners), et al, *Higgs Physics at the HL-LHC and HE-LHC*, CERN-LPCC-2018-04, <https://cds.cern.ch/record/2650162>.

1. To significantly improve the expected HL-LHC results, future facilities must push Higgs couplings' precision to the sub-% level
2. Event rates higher than what ee colliders can provide are needed to reach sub-% measurements of couplings such as $H\gamma\gamma$, $H\mu\mu$, $HZ\gamma$, Htt

The unique contributions of a 100 TeV pp collider to Higgs physics

- Huge Higgs production rates:
 - access (very) rare decay modes
 - push to %-level Higgs self-coupling measurement
 - new opportunities to reduce syst uncertainties (TH & EXP) and push precision
- Large dynamic range for H production (in p_T^H , $m(H+X)$, ...):
 - new opportunities for reduction of syst uncertainties (TH and EXP)
 - different hierarchy of production processes
 - develop indirect sensitivity to BSM effects at large Q^2 , complementary to that emerging from precision studies (eg *decay BRs*) at $Q \sim m_H$
- High energy reach
 - direct probes of BSM extensions of Higgs sector
 - SUSY Higgses
 - Higgs decays of heavy resonances
 - Higgs probes of the nature of EW phase transition
 - ...

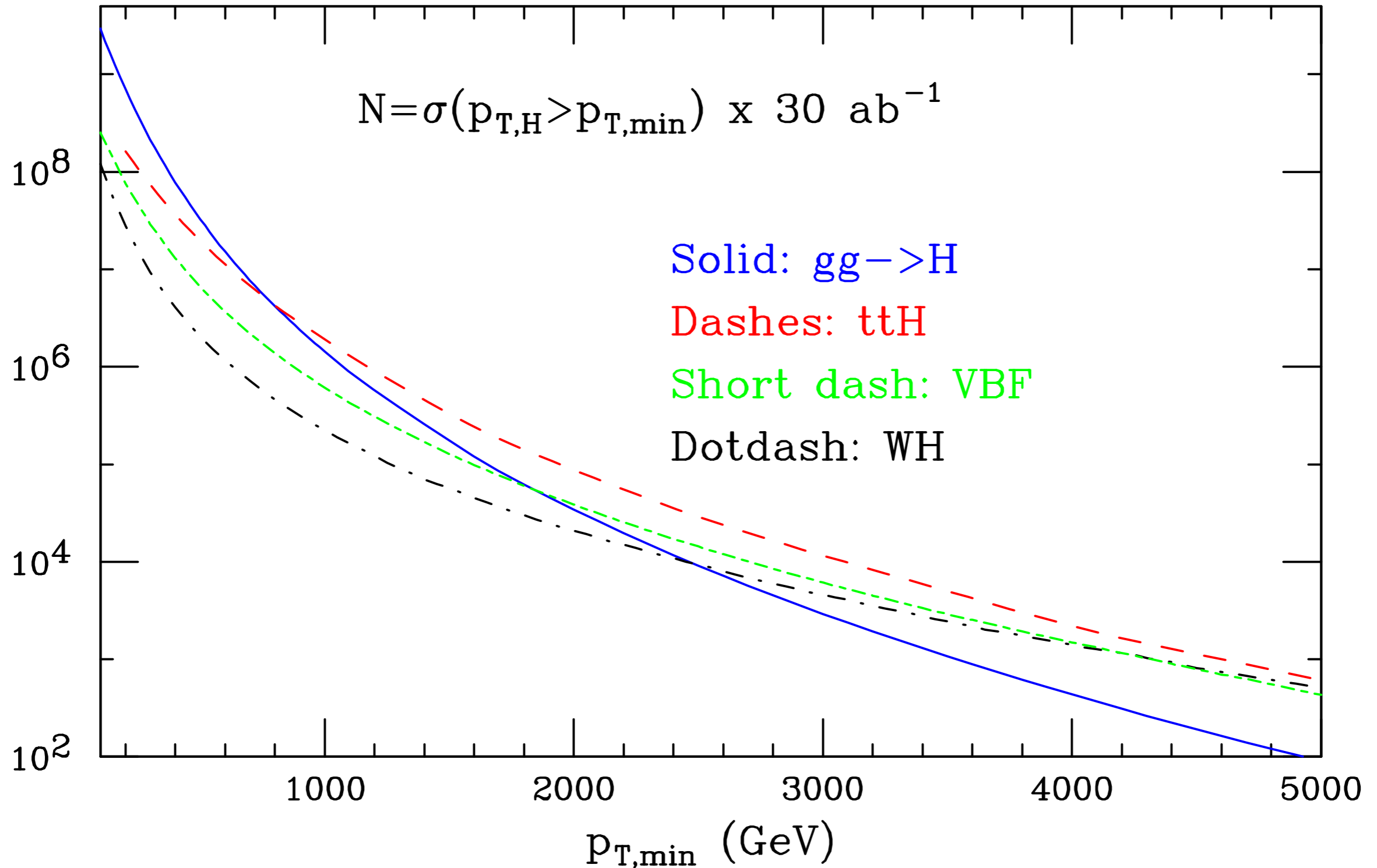
SM Higgs: event rates in pp@100 TeV

	gg→H	VBF	WH	ZH	ttH	HH
N_{100}	24 x 10 ⁹	2.1 x 10 ⁹	4.6 x 10 ⁸	3.3 x 10 ⁸	9.6 x 10 ⁸	3.6 x 10 ⁷
N_{100}/N_{14}	180	170	100	110	530	390

$$N_{100} = \sigma_{100\text{TeV}} \times 30 \text{ ab}^{-1}$$

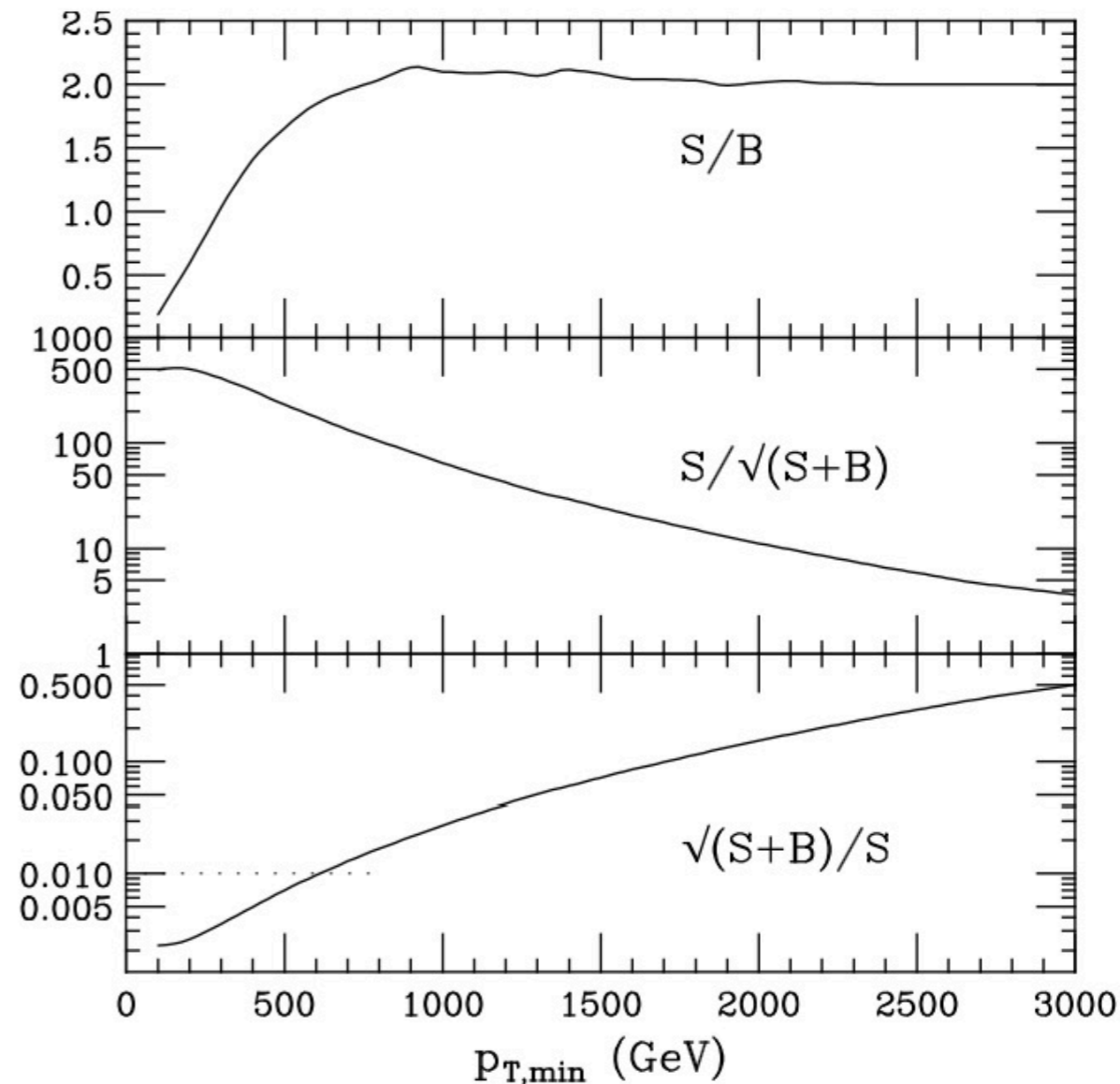
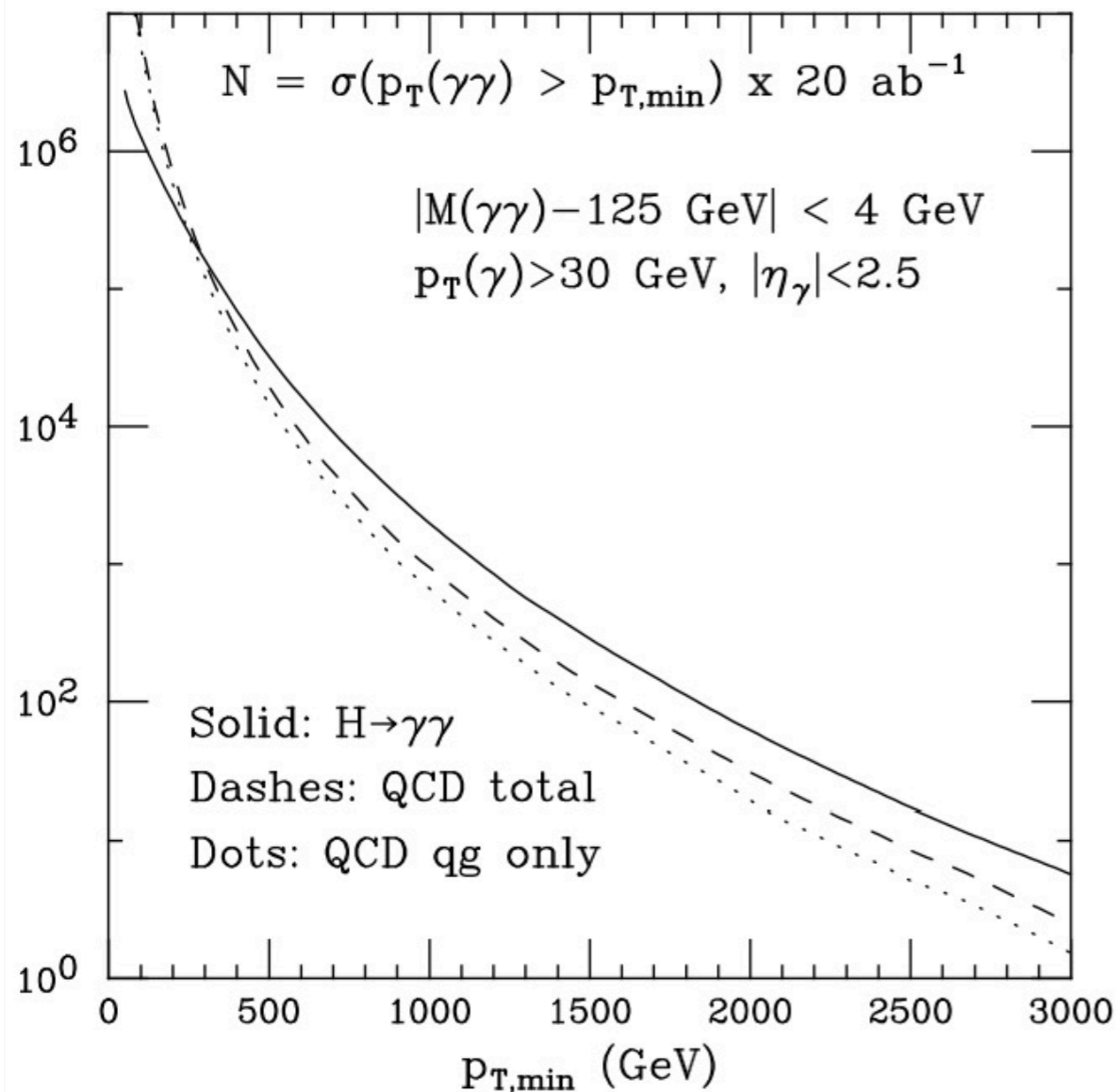
$$N_{14} = \sigma_{14\text{TeV}} \times 3 \text{ ab}^{-1}$$

H at large p_T



- Hierarchy of production channels changes at large $p_T(H)$:
 - $\sigma(ttH) > \sigma(gg \rightarrow H)$ above 800 GeV
 - $\sigma(\text{VBF}) > \sigma(gg \rightarrow H)$ above 1800 GeV

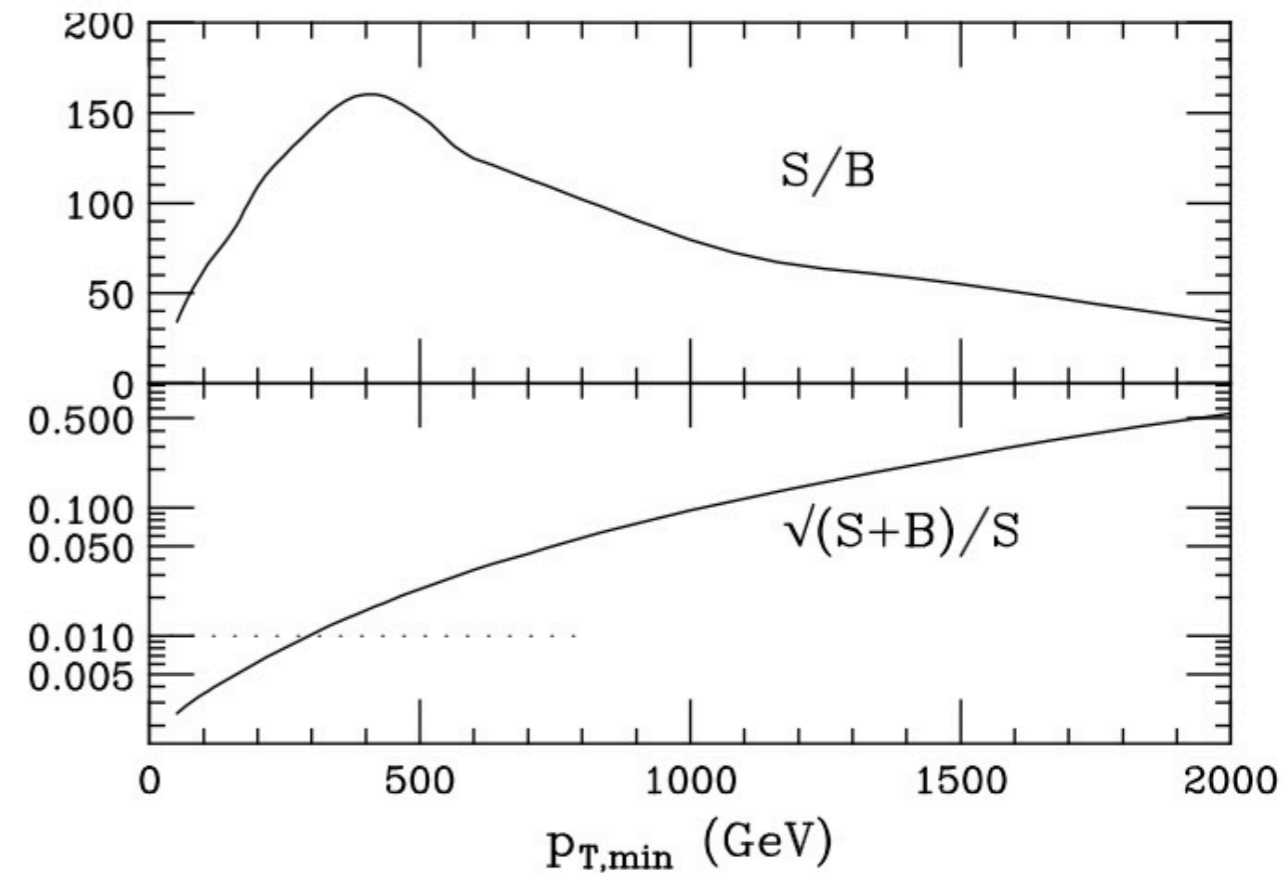
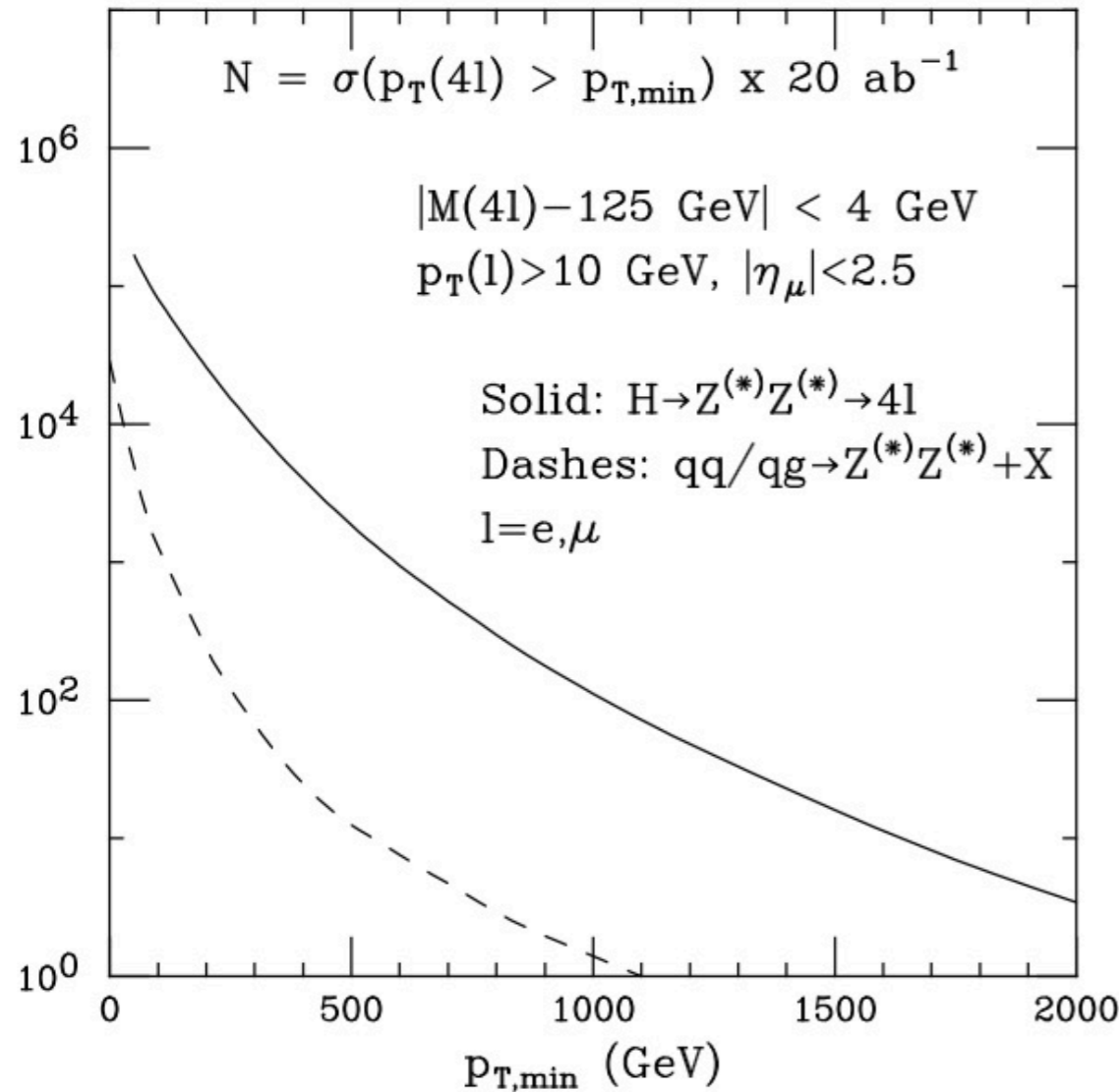
$gg \rightarrow H \rightarrow \gamma\gamma$ at large p_T



- At LHC, S/B in the $H \rightarrow \gamma\gamma$ channel is $O(\text{few } \%)$
- At FCC, for $p_T(H) > 300 \text{ GeV}$, $S/B \sim 1$
- Potentially accurate probe of the H p_T spectrum up to large p_T

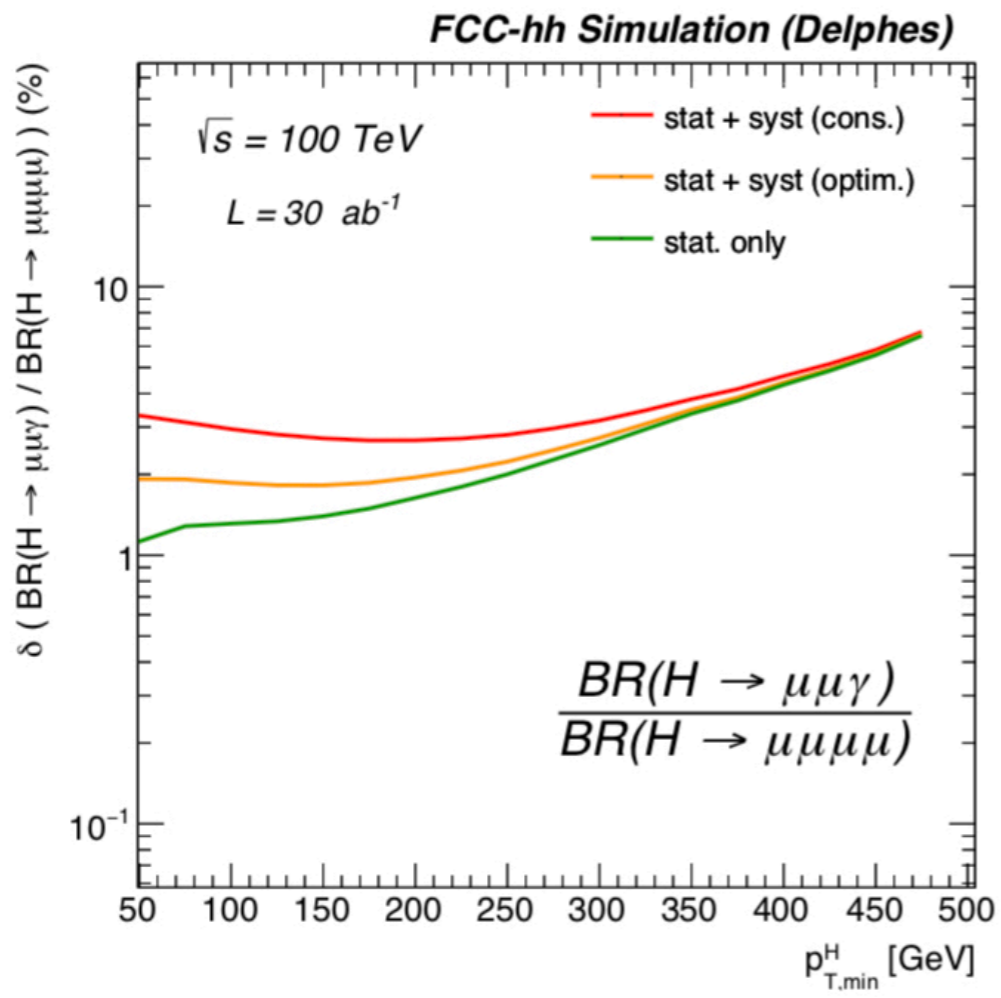
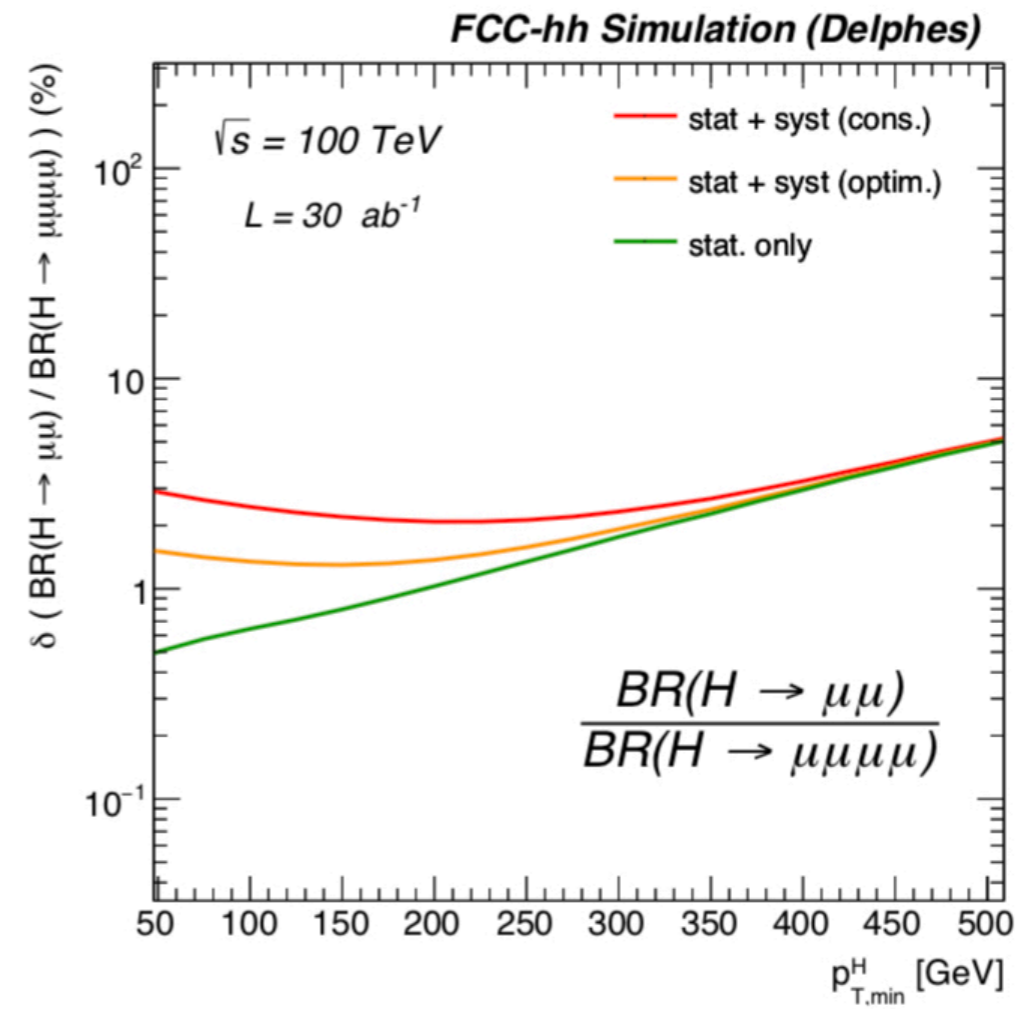
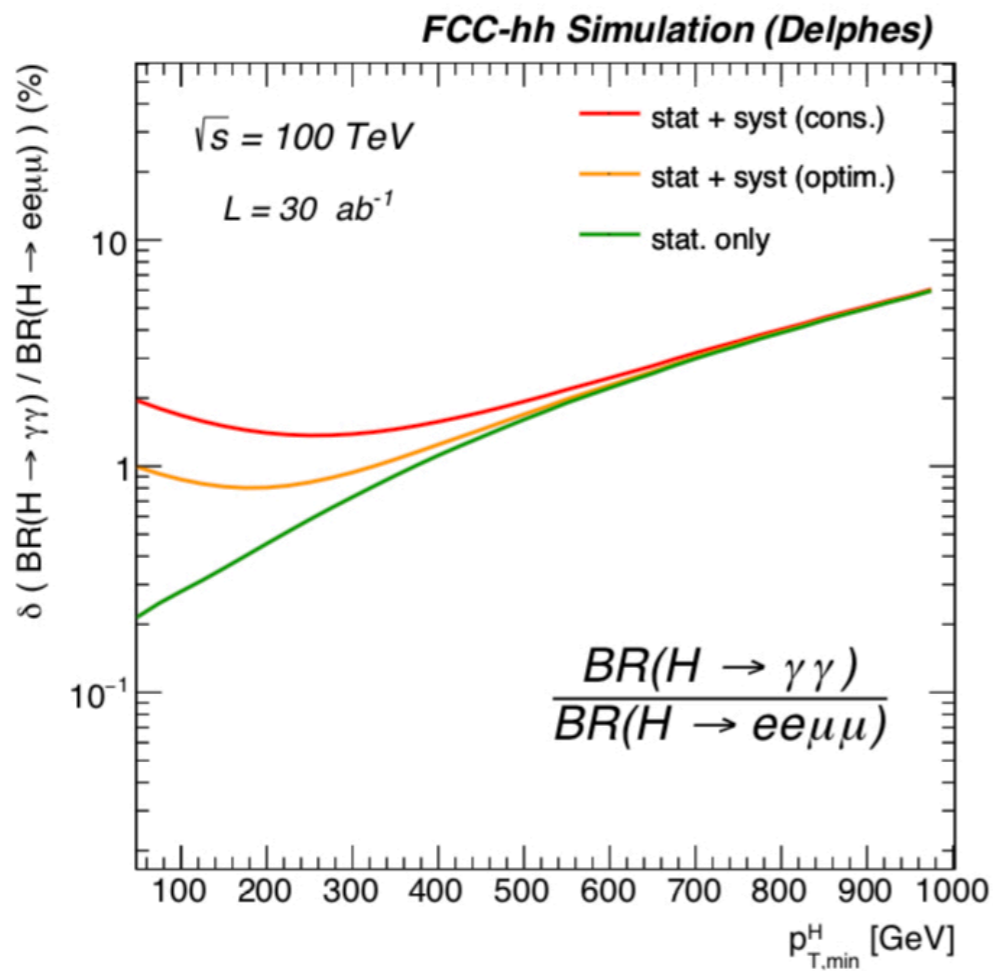
$p_{T,\min}$ (GeV)	δ_{stat}
100	0.2%
400	0.5%
600	1%
1600	10%

$gg \rightarrow H \rightarrow ZZ^* \rightarrow 4l$ at large p_T



- $S/B \sim 1$ for inclusive production at LHC
- Practically bg-free at large p_T at 100 TeV, maintaining large rates

$p_{T,\min}$ (GeV)	δ_{stat}
100	0.3%
300	1%
1000	10%



Importance of standalone precise “ratios-of-BRs” measurements:

- independent of α_S , m_b , m_c , Γ_{inv} systematics
- sensitive to BSM effects that typically influence BRs in different ways. Eg

$$\mathbf{BR(H \rightarrow \gamma\gamma) / BR(H \rightarrow ZZ^*)}$$

loop-level

tree-level

$$\mathbf{BR(H \rightarrow \mu\mu) / BR(H \rightarrow ZZ^*)}$$

2nd gen'n Yukawa

gauge coupling

$$\mathbf{BR(H \rightarrow \gamma\gamma) / BR(H \rightarrow Z\gamma)}$$

different EW charges in the loops of the two procs

$$\mathbf{BR(H \rightarrow inv) / BR(H \rightarrow \gamma\gamma)}$$

tree-level neutral

loop-level charged

Higgs couplings after FCC-ee / hh

	HL-LHC	FCC-ee	FCC-hh
$\delta\Gamma_H / \Gamma_H$ (%)	SM	1.3	tbd
$\delta g_{HZZ} / g_{HZZ}$ (%)	1.5	0.17	tbd
$\delta g_{HWW} / g_{HWW}$ (%)	1.7	0.43	tbd
$\delta g_{Hbb} / g_{Hbb}$ (%)	3.7	0.61	tbd
$\delta g_{Hcc} / g_{Hcc}$ (%)	~70	1.21	tbd
$\delta g_{Hgg} / g_{Hgg}$ (%)	2.5 (gg->H)	1.01	tbd
$\delta g_{H\tau\tau} / g_{H\tau\tau}$ (%)	1.9	0.74	tbd
$\delta g_{H\mu\mu} / g_{H\mu\mu}$ (%)	4.3	9.0	0.65 (*)
$\delta g_{H\gamma\gamma} / g_{H\gamma\gamma}$ (%)	1.8	3.9	0.4 (*)
$\delta g_{Htt} / g_{Htt}$ (%)	3.4	~10 (indirect)	0.95 (**)
$\delta g_{HZ\gamma} / g_{HZ\gamma}$ (%)	9.8	–	0.9 (*)
$\delta g_{HHH} / g_{HHH}$ (%)	50	~44 (indirect)	6.5
BR_{exo} (95%CL)	$BR_{\text{inv}} < 2.5\%$	< 1%	$BR_{\text{inv}} < 0.025\%$

* From BR ratios wrt $B(H \rightarrow 4\text{lept})$ @ FCC-ee

** From $pp \rightarrow ttH$ / $pp \rightarrow ttZ$, using $B(H \rightarrow bb)$ and ttZ EW coupling @ FCC-ee

Higgs self-coupling, $gg \rightarrow HH$

From the detector performance studies:

Pheno-level studies:

	$bb\gamma\gamma$	$bbZZ[\rightarrow 4l]$	$bbWW[\rightarrow 2jlv]$	$4b+j$	$2b2\tau+j$
$\delta\kappa_\lambda$ (%)	6.5	14	40	30	8

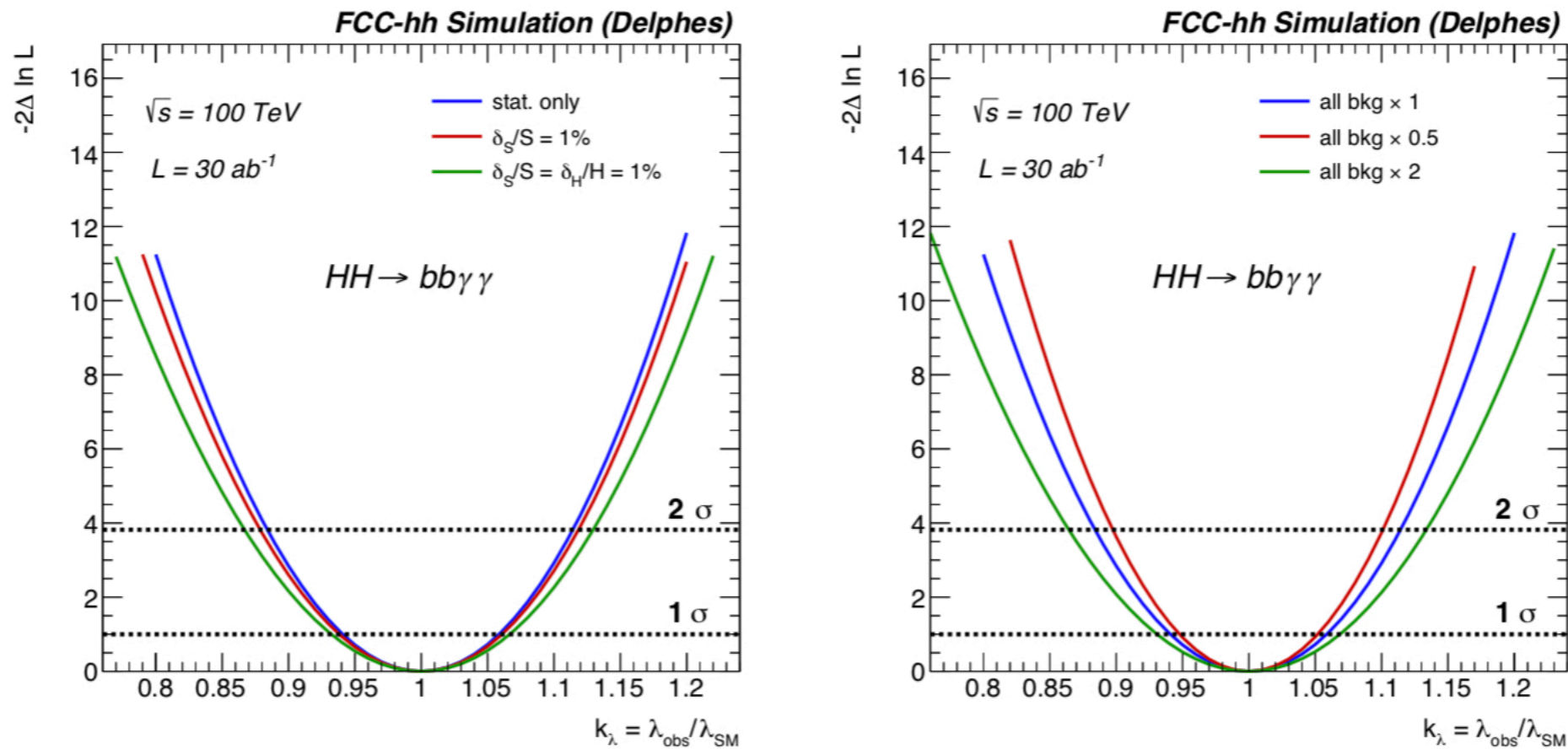


Figure 10.4: Expected precision on the Higgs self-coupling modifier κ_λ with no systematic uncertainties (only statistical), 1% signal uncertainty, 1% signal uncertainty together with 1% uncertainty on the Higgs backgrounds (left) and assuming respectively $\times 1$, $\times 2$, $\times 0.5$ background yields (right).

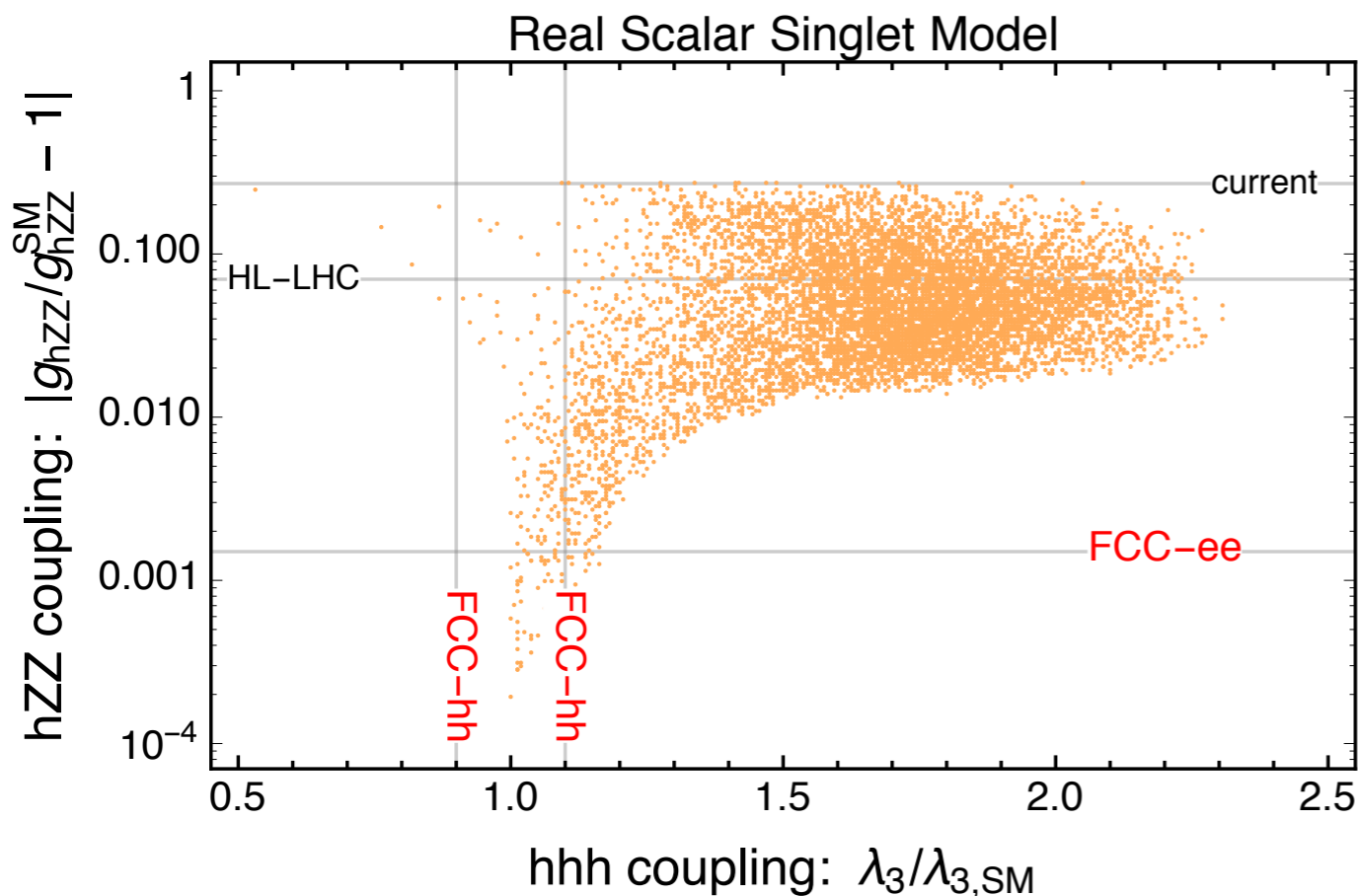
**Example of precision targets:
constraints on models with 1st order phase transition**

$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S \\ + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$

Example of precision targets: constraints on models with 1st order phase transition

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Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh

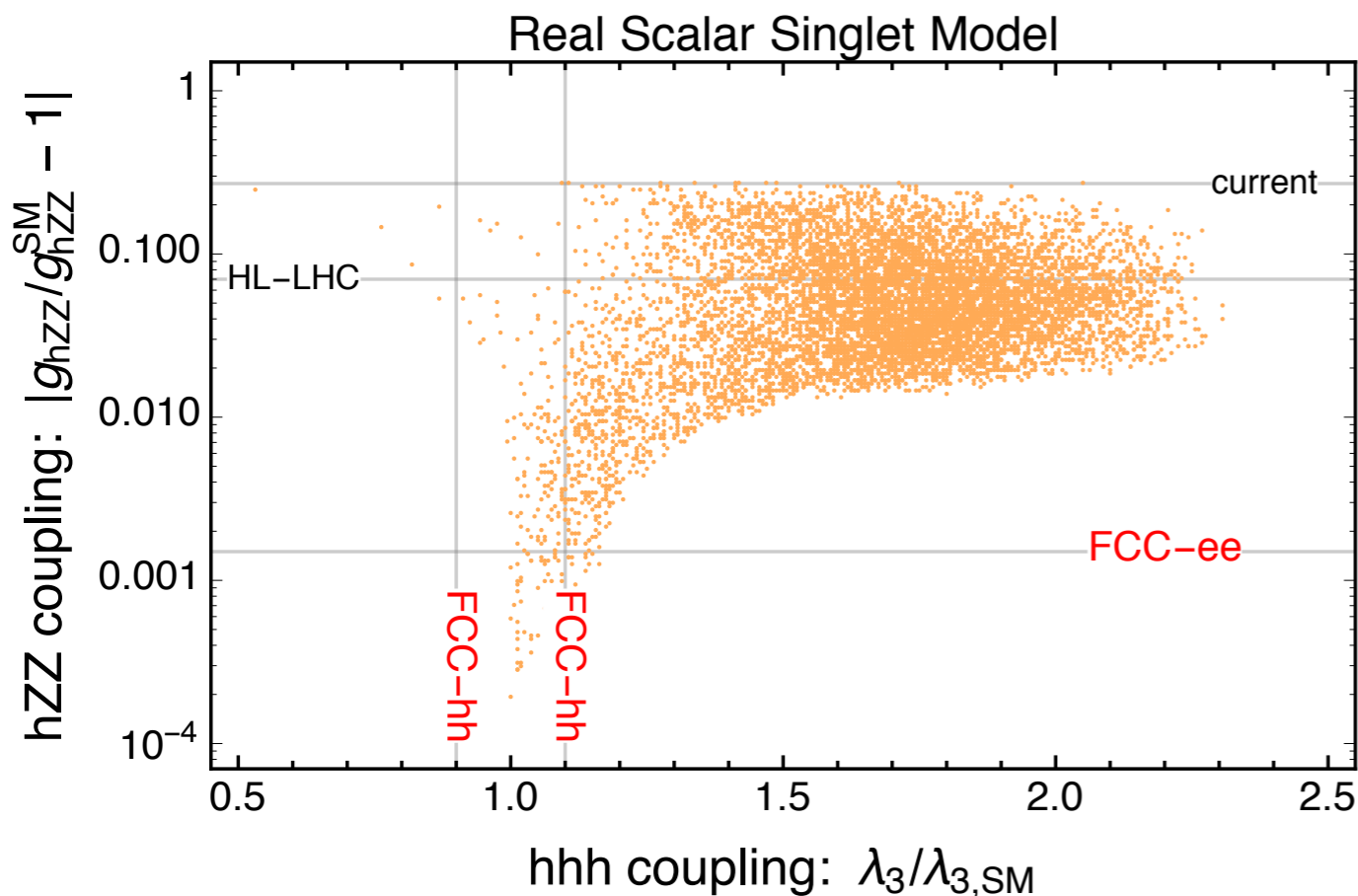


Parameter space scan for a singlet model extension of the Standard Model. The points indicate a first order phase transition.

Example of precision targets: constraints on models with 1st order phase transition

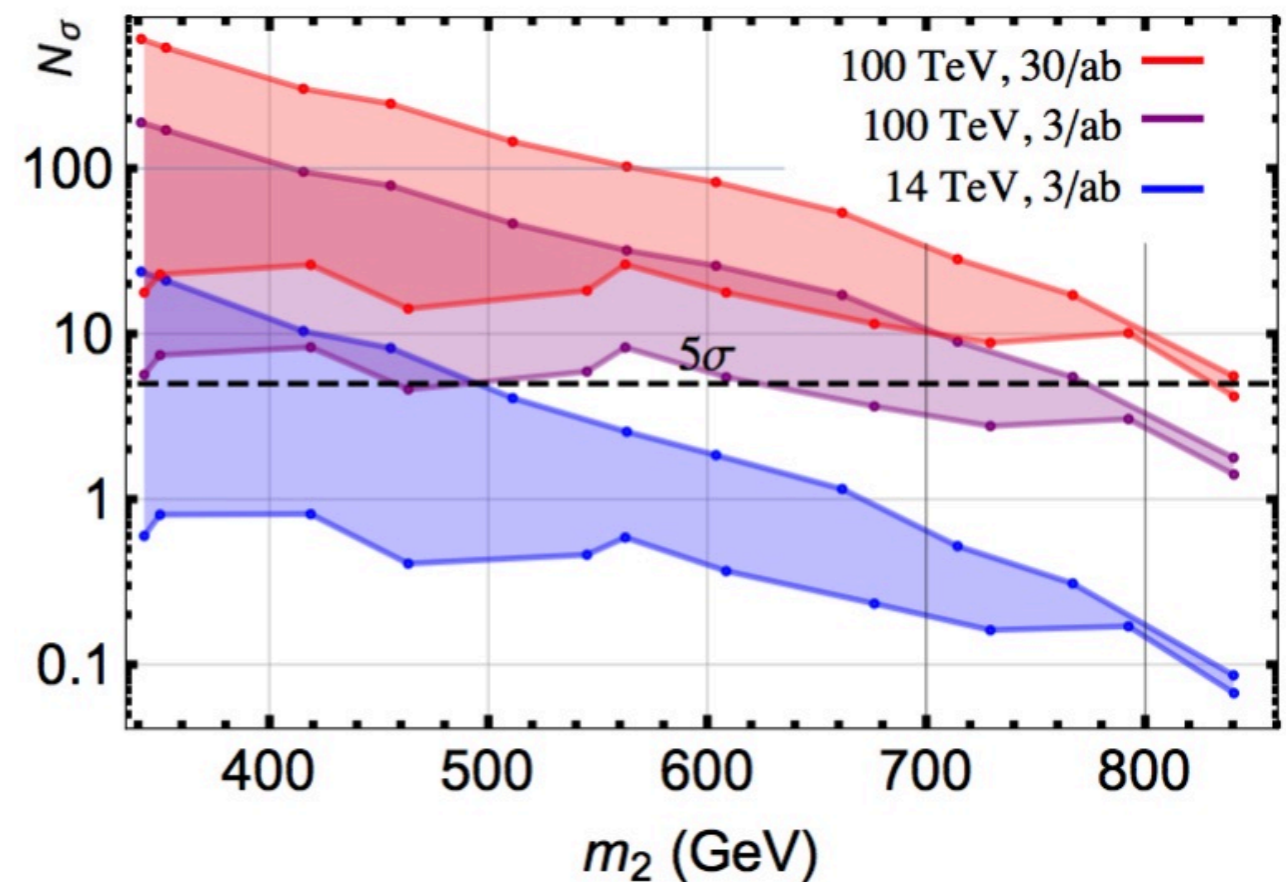
$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$

Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh



Parameter space scan for a singlet model extension of the Standard Model. The points indicate a first order phase transition.

Direct detection of extra Higgs states at FCC-hh



$h_2 \rightarrow h_1 h_1$ ($b\bar{b}\gamma\gamma + 4\tau$)
($h_2 \sim S, h_1 \sim H$)

Precision vs sensitivity

Precision vs sensitivity

- We often talk about “**precise**” Higgs measurements. What we actually aim at is “**sensitive**” tests of the Higgs properties, where *sensitive* refers to the ability to reveal BSM behaviours.

Precision vs sensitivity

- We often talk about “**precise**” Higgs measurements. What we actually aim at is “**sensitive**” tests of the Higgs properties, where *sensitive* refers to the ability to reveal BSM behaviours.
- **Sensitivity** may not require extreme precision
 - Going after “sensitivity”, rather than *just* precision, opens itself new opportunities ...

High- Q^2 observables : precision vs dynamic reach

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_k \mathcal{O}_k + \dots$$

$$O = | \langle f | L | i \rangle |^2 = O_{SM} [1 + O(\mu^2/\Lambda^2) + \dots]$$

High- Q^2 observables : precision vs dynamic reach

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_k \mathcal{O}_k + \dots$$

$$O = | \langle f | L | i \rangle |^2 = O_{SM} [1 + O(\mu^2 / \Lambda^2) + \dots]$$

For H decays, or inclusive production, $\mu \sim O(v, m_H)$

$$\delta O \sim \left(\frac{v}{\Lambda}\right)^2 \sim 6\% \left(\frac{\text{TeV}}{\Lambda}\right)^2 \Rightarrow \text{precision probes large } \Lambda$$

e.g. $\delta O = 1\% \Rightarrow \Lambda \sim 2.5 \text{ TeV}$

High- Q^2 observables : precision vs dynamic reach

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For H production off-shell or with large momentum transfer Q , $\mu \sim O(Q)$

$$\delta O \sim \left(\frac{Q}{\Lambda}\right)^2 \Rightarrow \text{kinematic reach probes large } \Lambda \text{ even if precision is "low"}$$

$$\text{e.g. } \delta O = 10\% \text{ at } Q = 1.5 \text{ TeV} \Rightarrow \Lambda \sim 5 \text{ TeV}$$

High- Q^2 observables : precision vs dynamic reach

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_k \mathcal{O}_k + \dots$$

$$O = | \langle f | L | i \rangle |^2 = O_{SM} [1 + O(\mu^2/\Lambda^2) + \dots]$$

For H decays, or inclusive production, $\mu \sim O(v, m_H)$

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$$\text{e.g. } \delta O = 1\% \Rightarrow \Lambda \sim 2.5 \text{ TeV}$$

For H production off-shell or with large momentum transfer Q , $\mu \sim O(Q)$

$$\delta O \sim \left(\frac{Q}{\Lambda}\right)^2 \Rightarrow \text{kinematic reach probes large } \Lambda$$

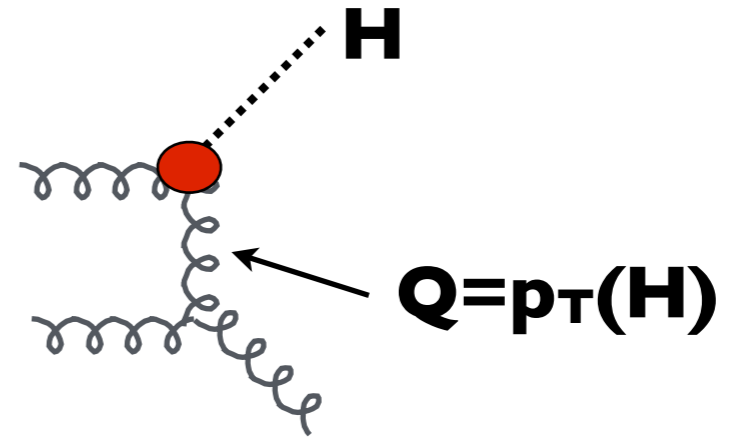
even if precision is “low”

$$\text{e.g. } \delta O = 10\% \text{ at } Q = 1.5 \text{ TeV} \Rightarrow \Lambda \sim 5 \text{ TeV}$$

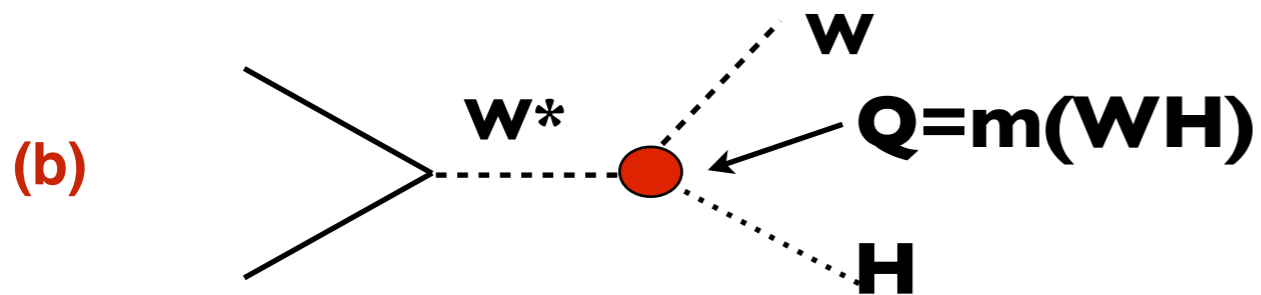
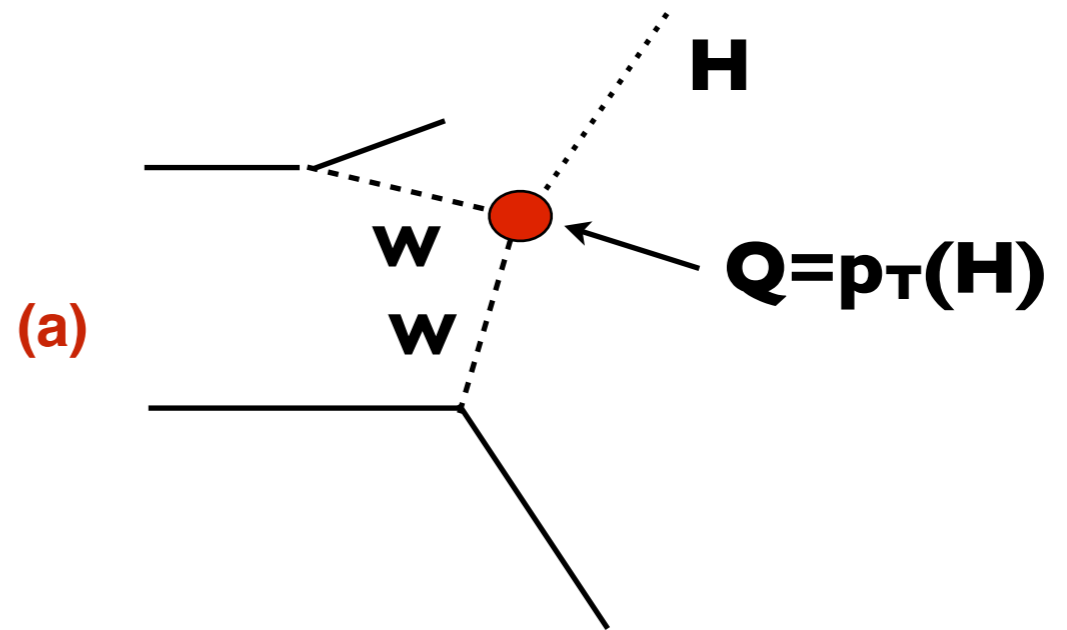
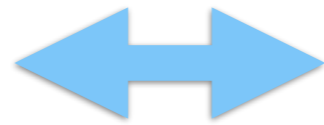
Complementarity between precise measurements at ee collider and large- Q studies at 100 TeV

Examples

$\delta BR(H \rightarrow gg)$



$\delta BR(H \rightarrow WW^*)$

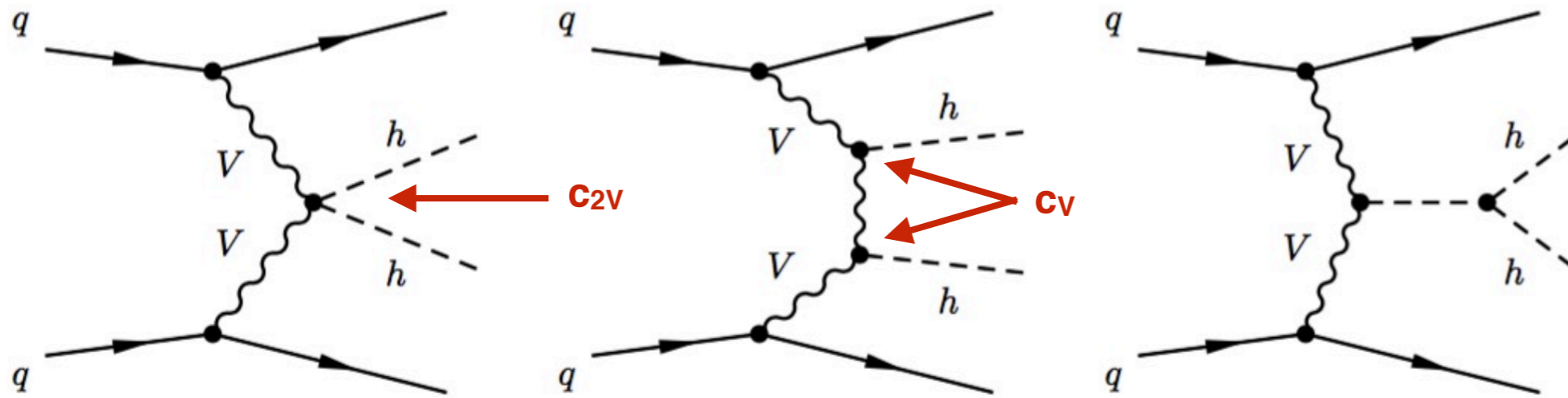


$$L_{D=6} = \frac{ig c_W}{2 \Lambda^2} (H^\dagger \sigma^a D^\mu H) D^\nu V_{\mu\nu}^a$$

➔

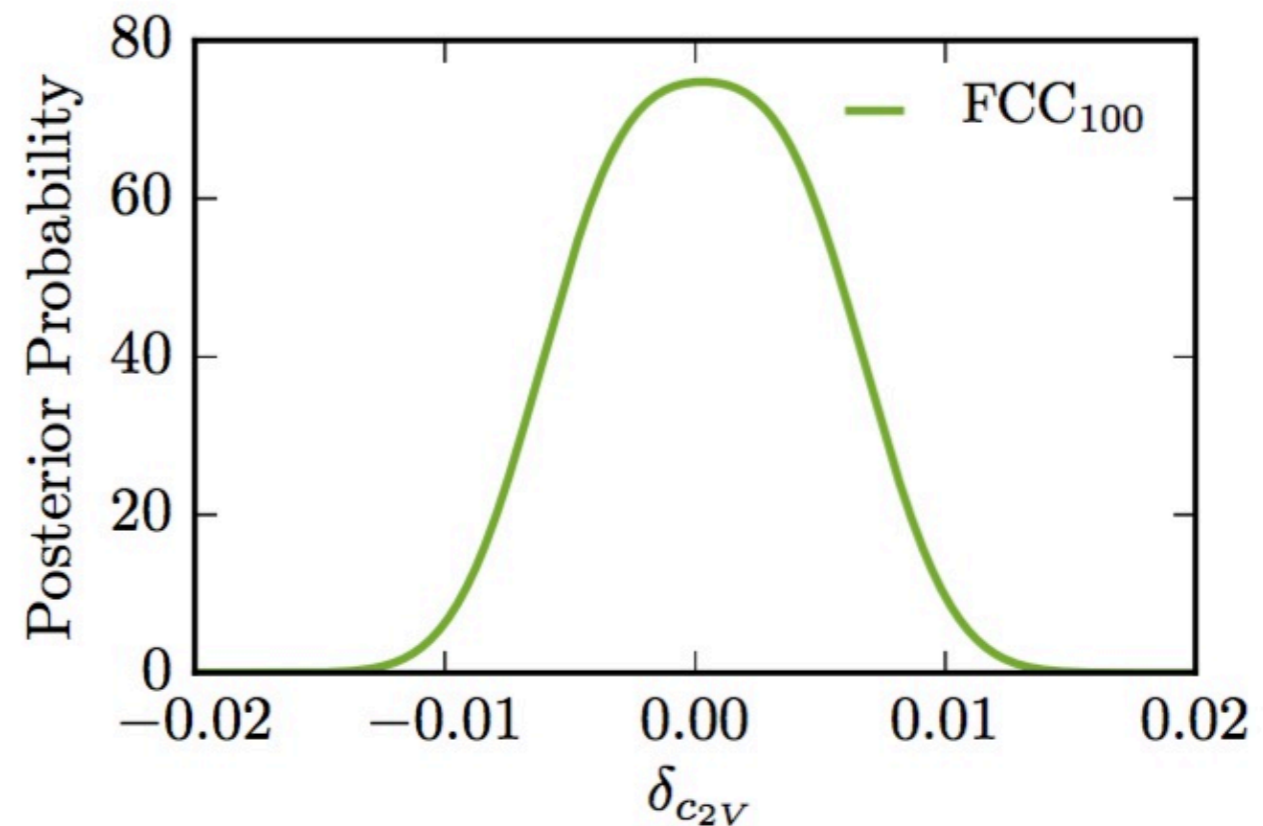
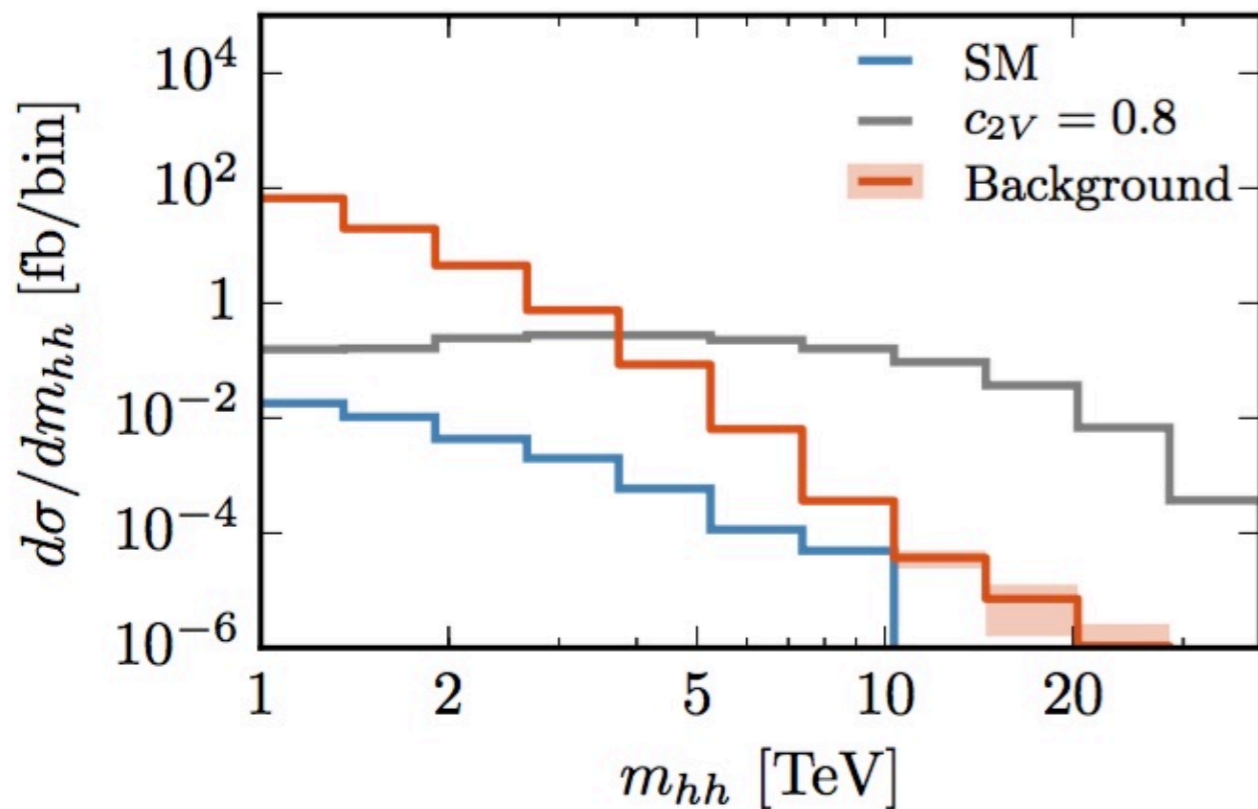
$$\frac{\sigma}{\sigma_{SM}} \sim \left(1 + c_W \frac{\hat{s}}{\Lambda^2} \right)^2$$

Example: high mass $VV \rightarrow HH$

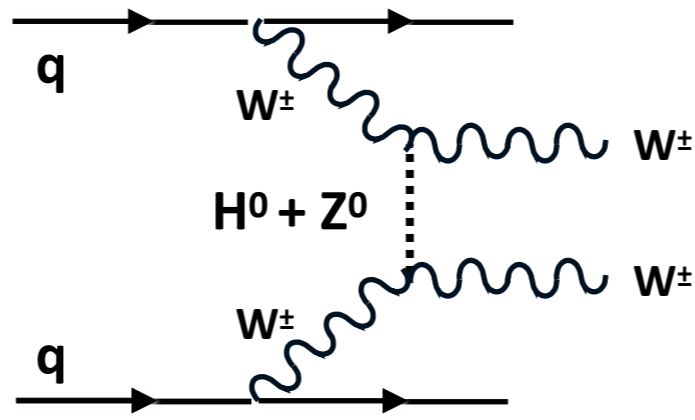


$$A(V_L V_L \rightarrow HH) \sim \frac{\hat{s}}{v^2} (c_{2V} - c_V^2) \cdot \text{where}$$

$$\begin{cases} c_V = g_{HVV}/g_{HVV}^{SM} \\ c_{2V} = g_{HHVV}/g_{HHVV}^{SM} \end{cases} \Rightarrow (c_{2V} - c_V^2)_{SM} = 0$$

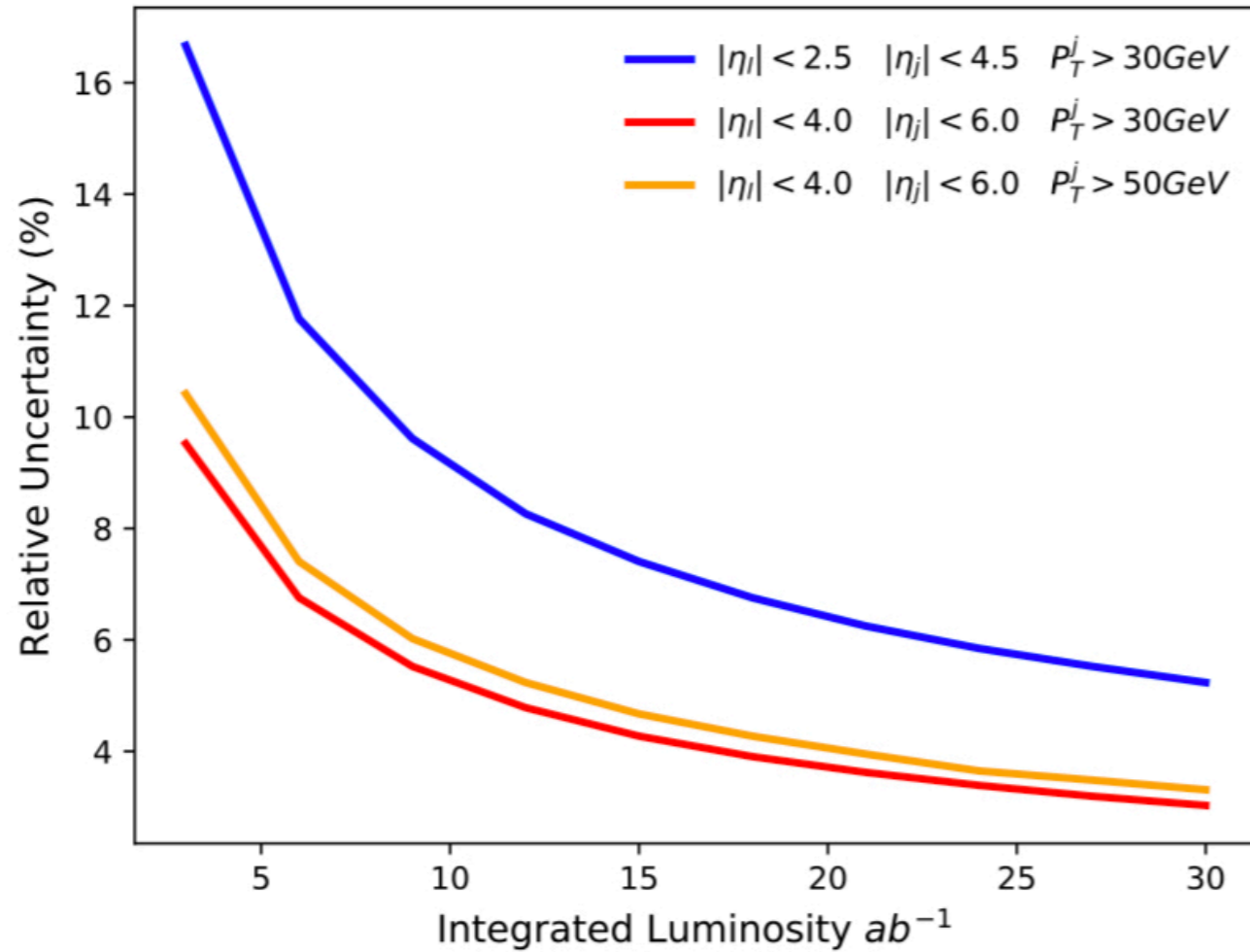


$W_L W_L$ scattering



large m_{WW}

VBS $W_L W_L$ Same Sign Cross Uncertainty



FCC-hh Simulation (Delphes)

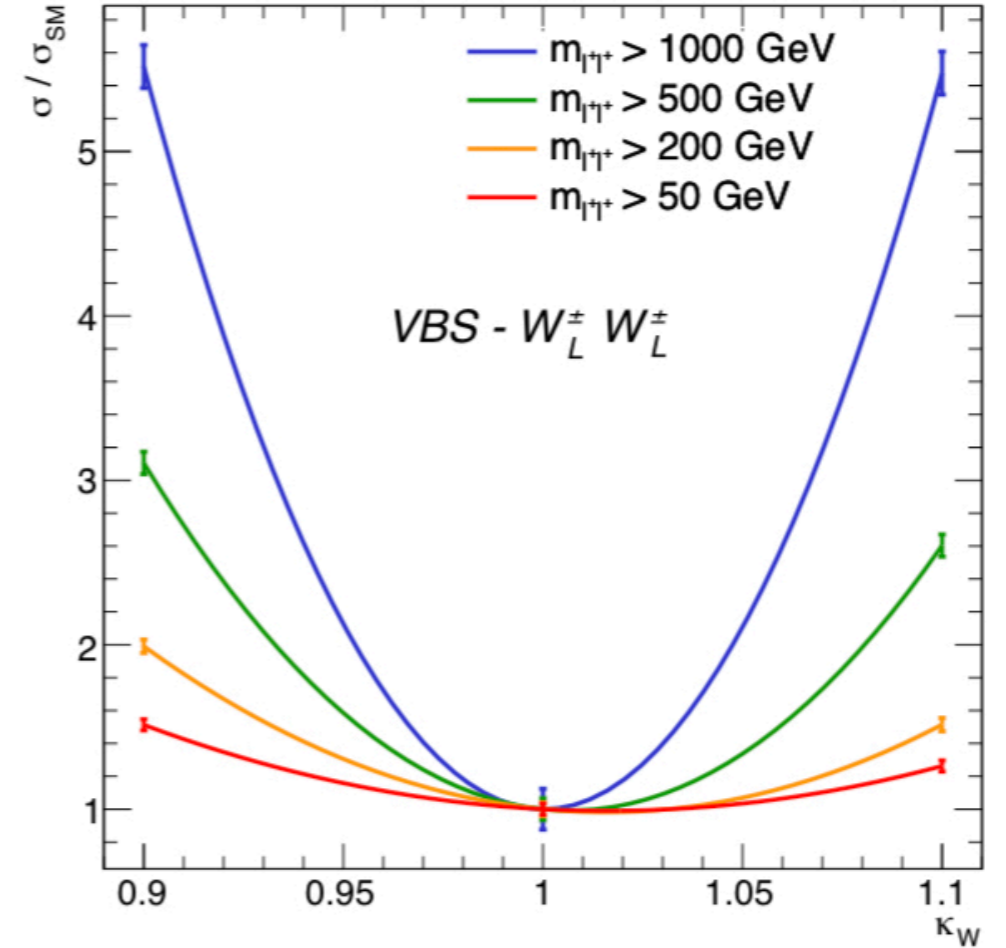


Table 4.5: Constraints on the HWW coupling modifier κ_W at 68% CL, obtained for various cuts on the di-lepton pair invariant mass in the $W_L W_L \rightarrow HH$ process.

m_{l+l^+} cut	$> 50\text{ GeV}$	$> 200\text{ GeV}$	$> 500\text{ GeV}$	$> 1000\text{ GeV}$
$\kappa_W \in$	[0.98, 1.05]	[0.99, 1.04]	[0.99, 1.03]	[0.98, 1.02]

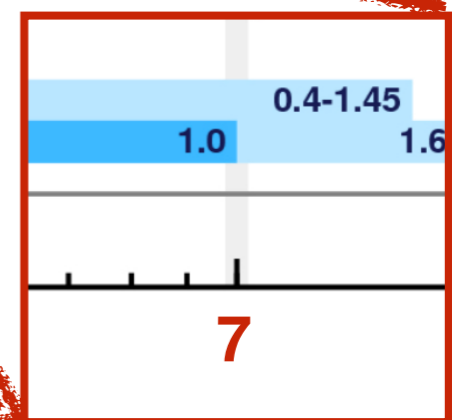
$$\kappa_W = \frac{g_{HWW}}{g_{HWW}^{SM}}$$

**Direct discovery reach:
the power of 100 TeV**

Model	Signature	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference	
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	2-6 jets 1-3 jets E_T^{miss} 36.1	\tilde{q} [2x, 8x Degen.] \tilde{q} [1x, 8x Degen.] 0.43 0.71 0.9 1.55 $m(\tilde{\chi}_1^0) < 100 \text{ GeV}$ $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1712.02332 1711.03301
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, μ	2-6 jets E_T^{miss} 36.1	\tilde{g} \tilde{g} Forbidden 0.95-1.6 2.0 $m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{\chi}_1^0) = 900 \text{ GeV}$	1712.02332 1712.02332
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	3 e, μ $ee, \mu\mu$	4 jets 2 jets E_T^{miss} 36.1	\tilde{g} \tilde{g} 1.2 1.85 $m(\tilde{\chi}_1^0) < 800 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$	1706.03731 1805.11381
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	0 e, μ 3 e, μ	7-11 jets 4 jets E_T^{miss} 36.1	\tilde{g} \tilde{g} 0.98 1.8 $m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200 \text{ GeV}$	1708.02794 1706.03731
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	0-1 e, μ 3 e, μ	3 b 4 jets E_T^{miss} 79.8 36.1	\tilde{g} \tilde{g} 1.25 2.25 $m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300 \text{ GeV}$	ATLAS-CONF-2018-041 1706.03731
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/t\tilde{\chi}_1^\pm$	Multiple Multiple Multiple	36.1 36.1 36.1	\tilde{b}_1 \tilde{b}_1 \tilde{b}_1 Forbidden 0.9 Forbidden 0.58-0.82 Forbidden 0.7 $m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(b\tilde{\chi}_1^0) = 1$ $m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(b\tilde{\chi}_1^0) = \text{BR}(t\tilde{\chi}_1^\pm) = 0.5$ $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, m(\tilde{\chi}_1^\pm) = 300 \text{ GeV}, \text{BR}(t\tilde{\chi}_1^\pm) = 1$	1708.09266, 1711.03301 1708.09266 1706.03731
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$	0 e, μ	6 b E_T^{miss} 139	\tilde{b}_1 \tilde{b}_1 Forbidden 0.23-0.48 0.23-1.35 $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0) = 100 \text{ GeV}$ $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0) = 0 \text{ GeV}$	SUSY-2018-31 SUSY-2018-31
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$	0-2 e, μ	0-2 jets/1-2 b E_T^{miss} 36.1	\tilde{t}_1 1.0 $m(\tilde{\chi}_1^0) = 1 \text{ GeV}$	1506.08616, 1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1$, Well-Tempered LSP	Multiple	36.1	\tilde{t}_1 0.48-0.84 $m(\tilde{\chi}_1^0) = 150 \text{ GeV}, m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}, \tilde{t}_1 \approx \tilde{t}_L$	1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$	1 $\tau + 1 e, \mu, \tau$	2 jets/1 b E_T^{miss} 36.1	\tilde{t}_1 1.16 $m(\tilde{\tau}_1) = 800 \text{ GeV}$	1803.10178
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0 e, μ	2 c E_T^{miss} 36.1	\tilde{t}_1 \tilde{t}_1 0.46 0.85 $m(\tilde{\chi}_1^0) = 0 \text{ GeV}$ $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$ $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1805.01649 1805.01649 1711.03301
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$	0 e, μ	mono-jet E_T^{miss} 36.1	\tilde{t}_1 0.43 $m(\tilde{\chi}_1^0) = 0 \text{ GeV}, m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 180 \text{ GeV}$	1706.03986
EW direct	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ	2-3 e, μ $ee, \mu\mu$	≥ 1 E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.17 0.6 $m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 10 \text{ GeV}$	1403.5294, 1806.02293 1712.08119
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via WW	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 0.42 $m(\tilde{\chi}_1^0) = 0$	ATLAS-CONF-2019-008
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh	0-1 e, μ	2 b E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.68 $m(\tilde{\chi}_1^0) = 0$	1812.09432
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via $\tilde{\ell}_L/\tilde{\nu}$	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 1.0 $m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	ATLAS-CONF-2019-008
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}_1\nu(\tau\tilde{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}_1\nu(\nu\tilde{\nu})$	2 τ	E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.76 0.22 $m(\tilde{\chi}_1^0) = 0, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 100 \text{ GeV}, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1708.07875 1708.07875
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ 2 e, μ	0 jets ≥ 1 E_T^{miss} 139 36.1	$\tilde{\ell}$ $\tilde{\ell}$ 0.7 0.18 $m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	ATLAS-CONF-2019-008 1712.08119
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ	$\geq 3 b$ 0 jets E_T^{miss} 36.1 36.1	\tilde{H} \tilde{H} 0.13-0.23 0.3 0.29-0.88 $\text{BR}(\tilde{\chi}_1^0 \rightarrow h\tilde{G}) = 1$ $\text{BR}(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$	1806.04030 1804.03602
Long-lived particles	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet E_T^{miss} 36.1	$\tilde{\chi}_1^\pm$ $\tilde{\chi}_1^\pm$ 0.15 0.46 Pure Wino Pure Higgsino	1712.02118 ATL-PHYS-PUB-2017-019
	Stable \tilde{g} R-hadron	Multiple	36.1	\tilde{g} 2.0 $m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1902.01636, 1808.04095
	Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$	Multiple	36.1	\tilde{g} [$\tau(\tilde{g}) = 10 \text{ ns}, 0.2 \text{ ns}$] 2.05 2.4	1710.04901, 1808.04095
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu/\mu\tau$	$e\mu, e\tau, \mu\tau$	3.2	$\tilde{\nu}_\tau$ 1.9 $\lambda'_{311} = 0.11, \lambda'_{132/133/233} = 0.07$	1607.08079
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\nu\nu$	4 e, μ	0 jets E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [$\lambda'_{333} \neq 0, \lambda'_{12k} \neq 0$] 0.82 1.33 $m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1804.03602
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$	4-5 large-R jets	36.1	\tilde{g} [$m(\tilde{\chi}_1^0) = 200 \text{ GeV}, 1100 \text{ GeV}$] \tilde{g} [$\lambda'_{112} = 2e-4, 2e-5$] 1.05 1.3 1.9 2.0 $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, \text{bino-like}$	1804.03568 ATLAS-CONF-2018-003
	$\tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$	Multiple	36.1	\tilde{g} [$\lambda'_{323} = 2e-4, 1e-2$] 0.55 1.05 $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, \text{bino-like}$	ATLAS-CONF-2018-003
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$	2 jets + 2 b	36.7	\tilde{t}_1 [qq, bs] 0.42 0.61 0.4-1.45 $\text{BR}(\tilde{t}_1 \rightarrow q\mu) = 100\%, \text{cos}\theta_{\tilde{t}_1} > 20\%$	1710.07171
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 e, μ 1 μ	2 b DV 36.1 136	\tilde{t}_1 \tilde{t}_1 1.0 1.6 $\text{BR}(\tilde{t}_1 \rightarrow q\mu) = 100\%, \text{cos}\theta_{\tilde{t}_1} > 20\%$	1710.05544 ATLAS-CONF-2019-006

10⁻¹ 1 Mass scale [TeV]

@14 TeV

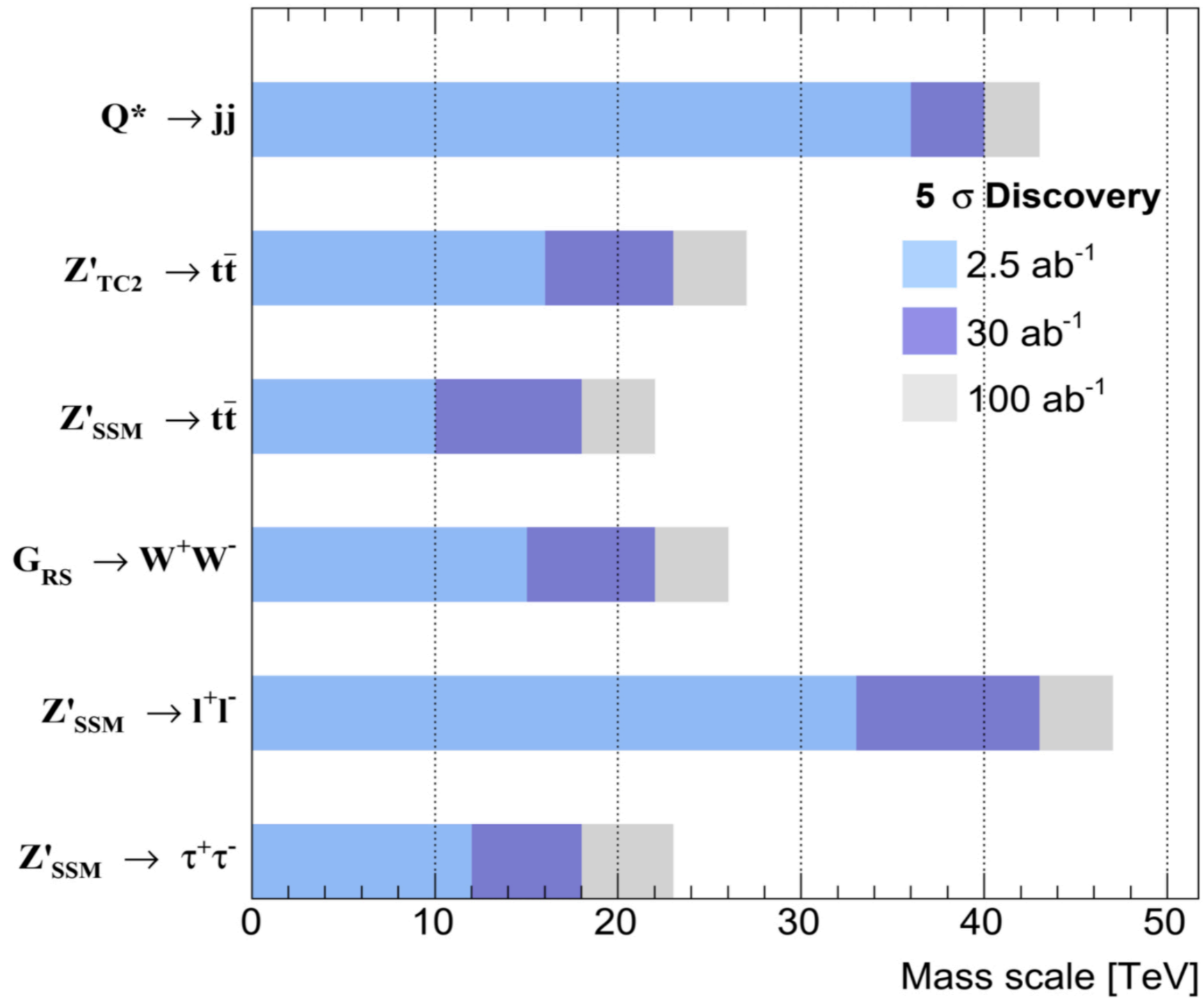


@100 TeV

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

s-channel resonances

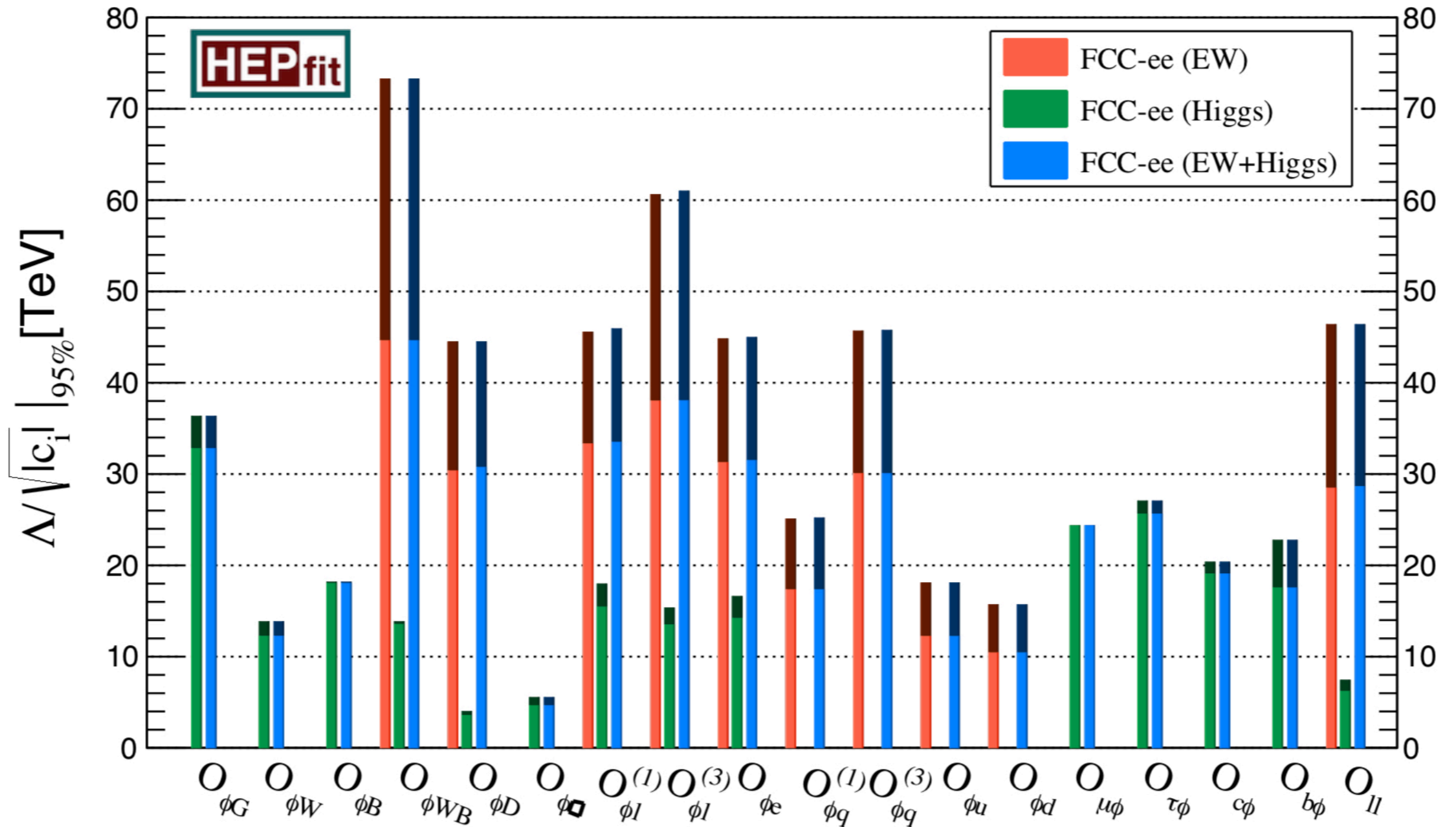
FCC-hh Simulation (Delphes), $\sqrt{s} = 100$ TeV



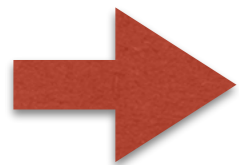
FCC-hh reach ~ 6 x HL-LHC reach

Global EFT fits to EW and H observables at FCC-ee

=> see L. Reina lecture 3



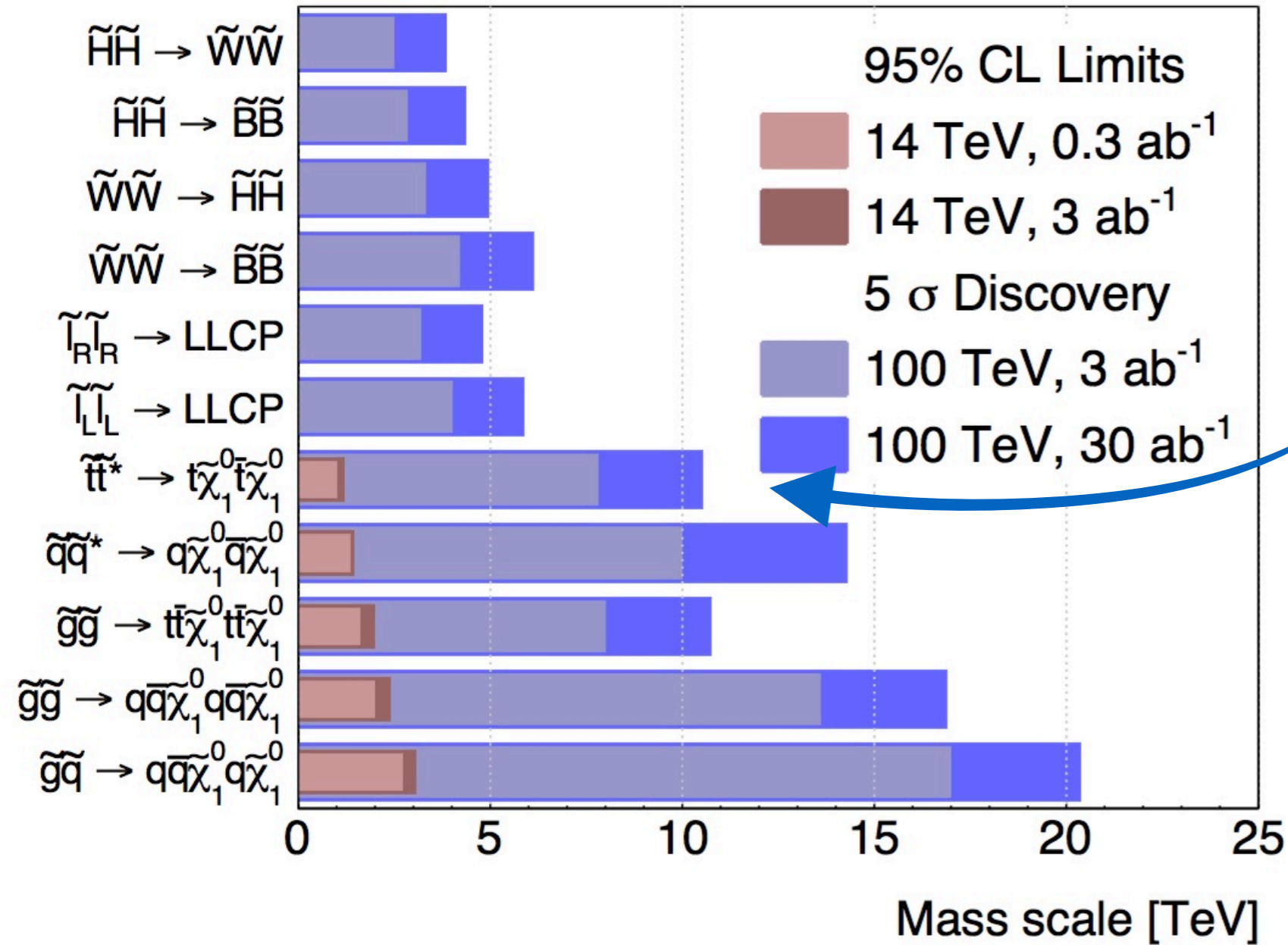
Constraints on the coefficients of various EFT op's from a global fit of (i) EW observables, (ii) Higgs couplings and (iii) EW+Higgs combined. Darker shades of each color indicate the results neglecting all SM theory uncertainties.



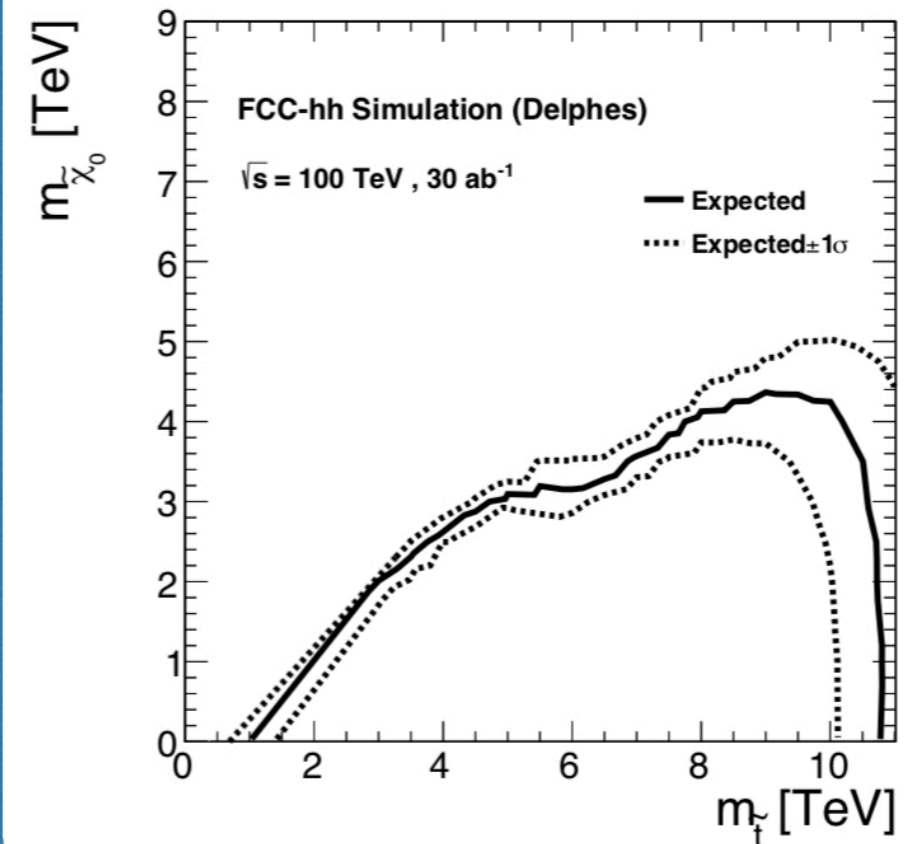
100 TeV is the appropriate CoM energy to directly search for new physics appearing indirectly through precision EW and H measurements at the future ee collider

SUSY reach at 100 TeV

Early phenomenology studies



New detector performance studies



WIMP DM theoretical constraints

See lecture 1 by M. Mc Cullough

For particles held in equilibrium by pair creation and annihilation processes, ($\chi \chi \leftrightarrow \text{SM}$)

$$\Omega_{\text{DM}} h^2 \sim \frac{10^9 \text{GeV}^{-1}}{M_{\text{pl}}} \frac{1}{\langle \sigma v \rangle}$$

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$$\Omega_{\text{DM}} h^2 \sim 0.12 \times \left(\frac{M_{\text{DM}}}{2 \text{TeV}} \right)^2 \left(\frac{0.3}{g_{\text{eff}}} \right)^4$$

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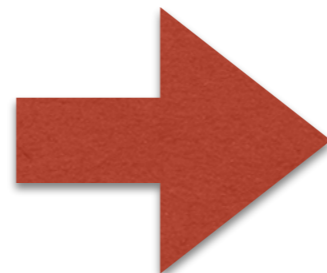
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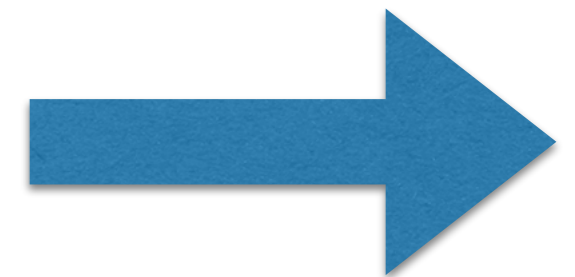
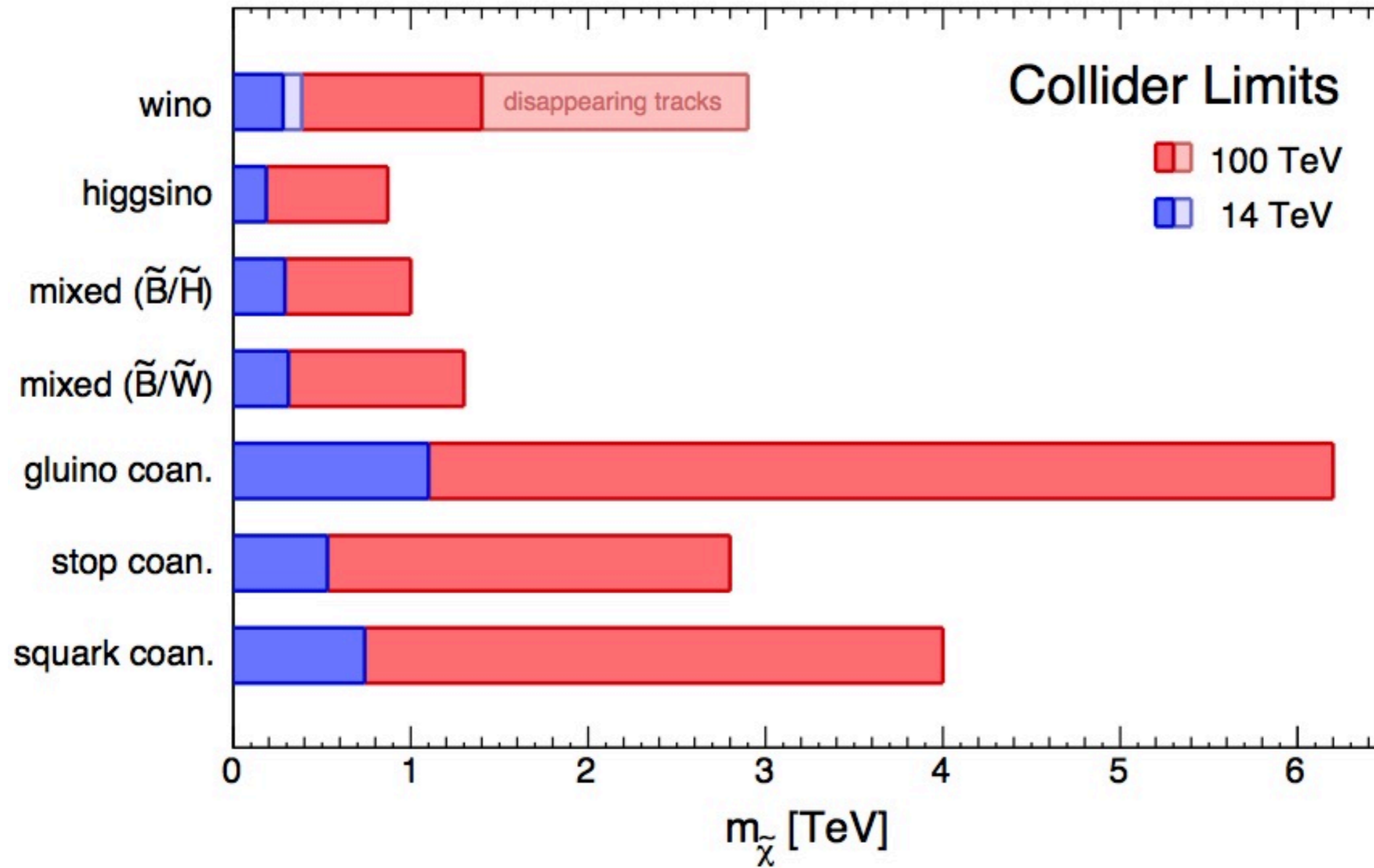
$$\Omega_{\text{wimp}} h^2 \lesssim 0.12$$



$$M_{\text{wimp}} \lesssim 2 \text{TeV} \left(\frac{g}{0.3} \right)^2$$

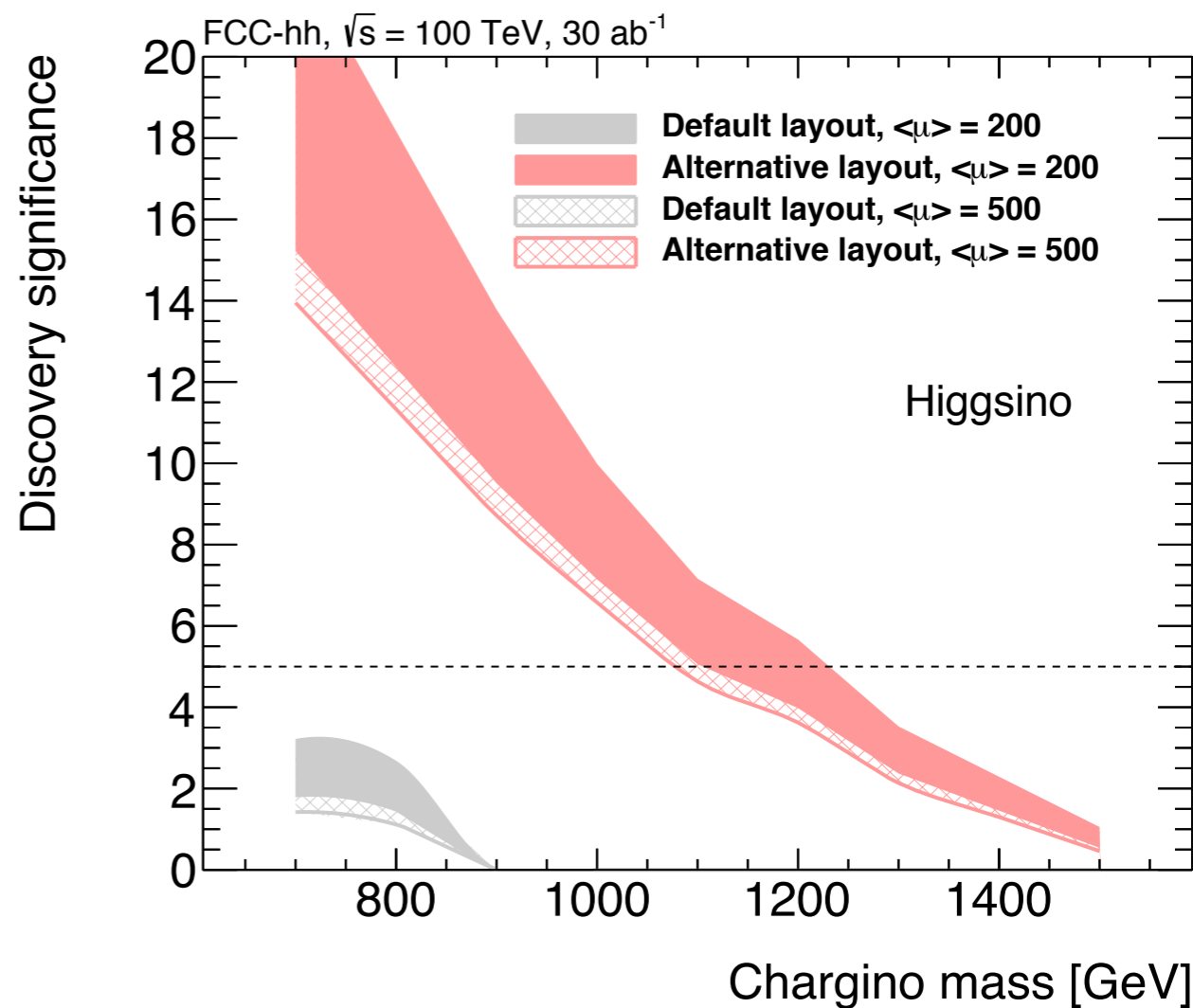
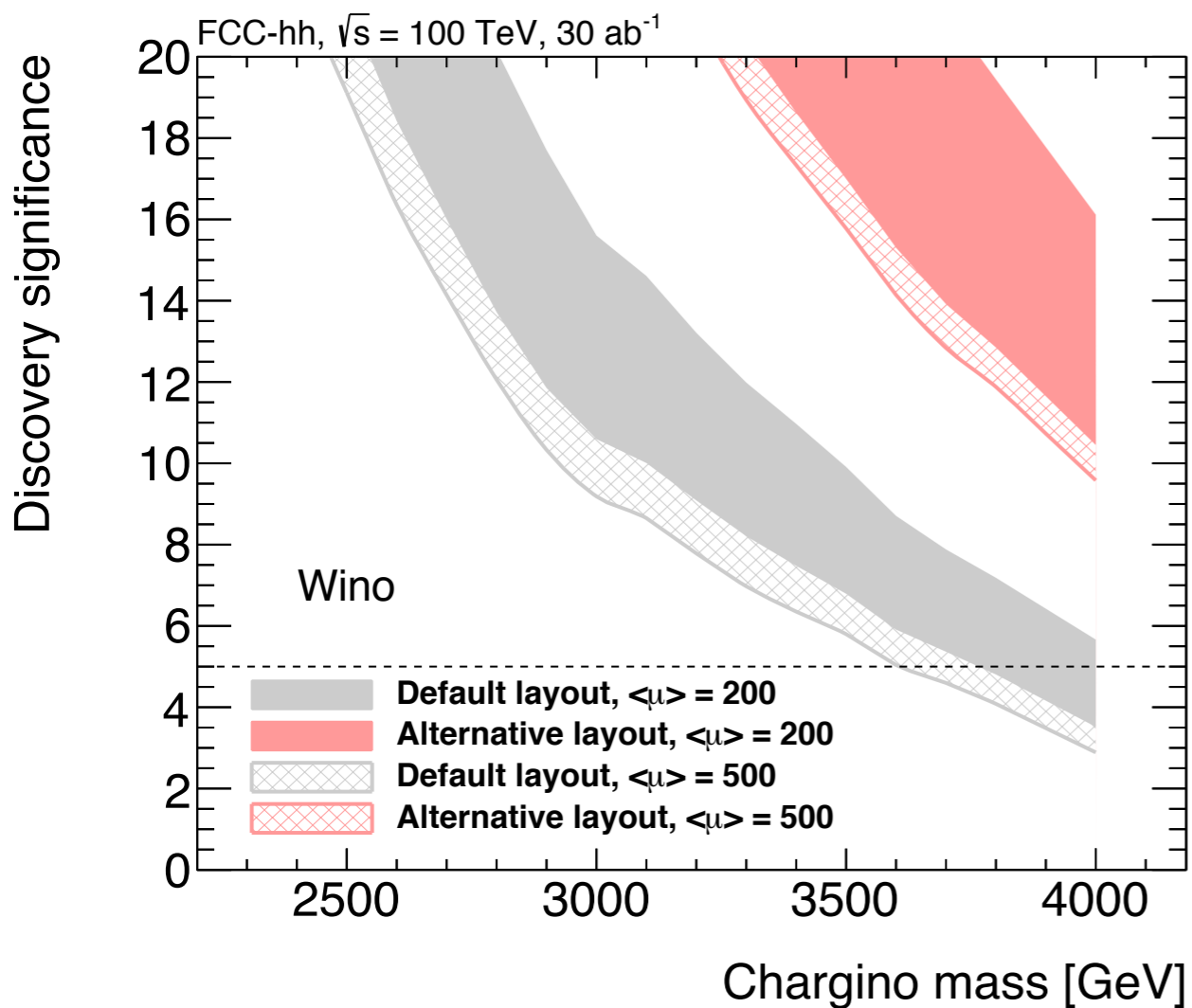
DM reach at 100 TeV

Early phenomenology studies



New detector performance studies

Disappearing charged track analyses (at ~full pileup)



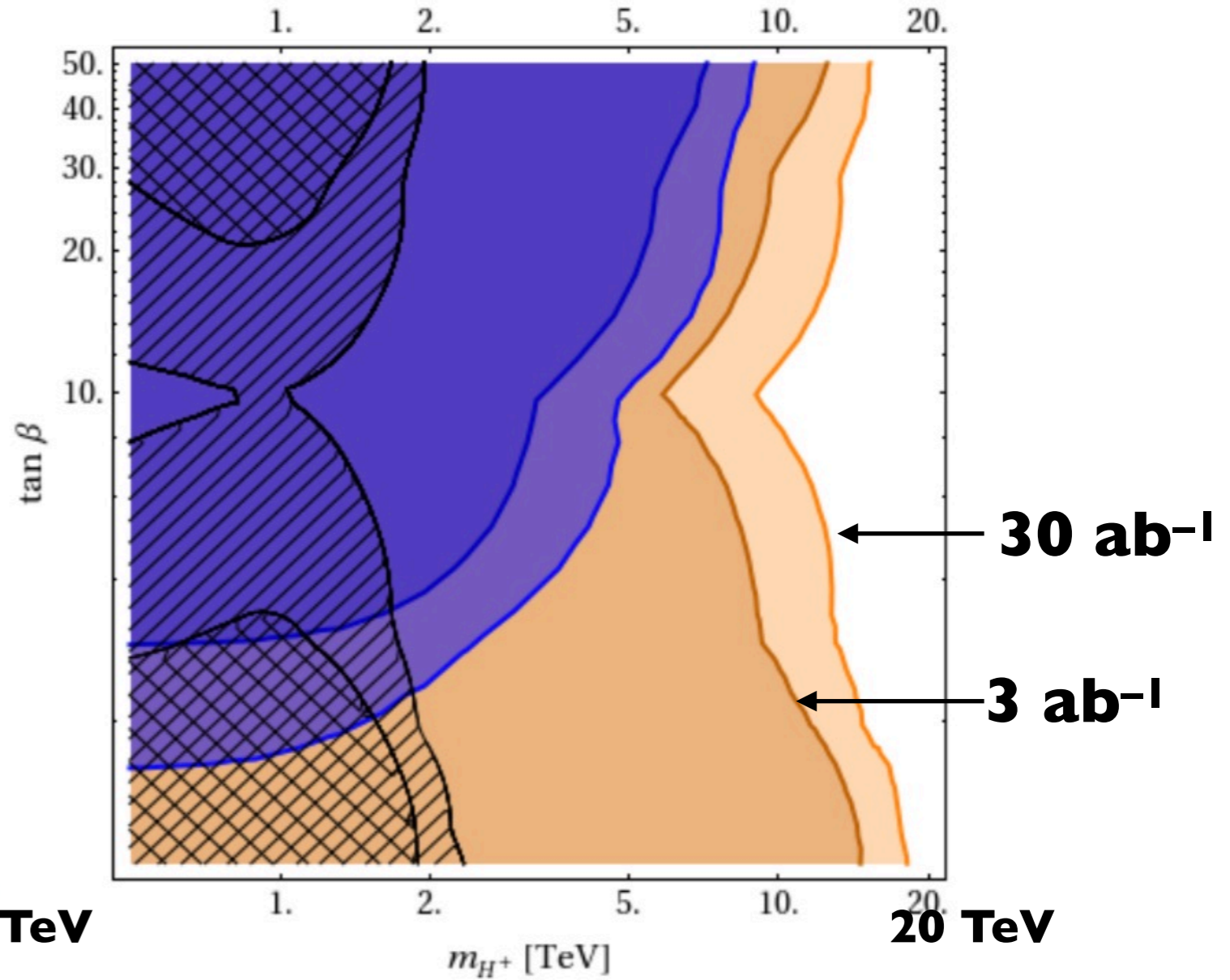
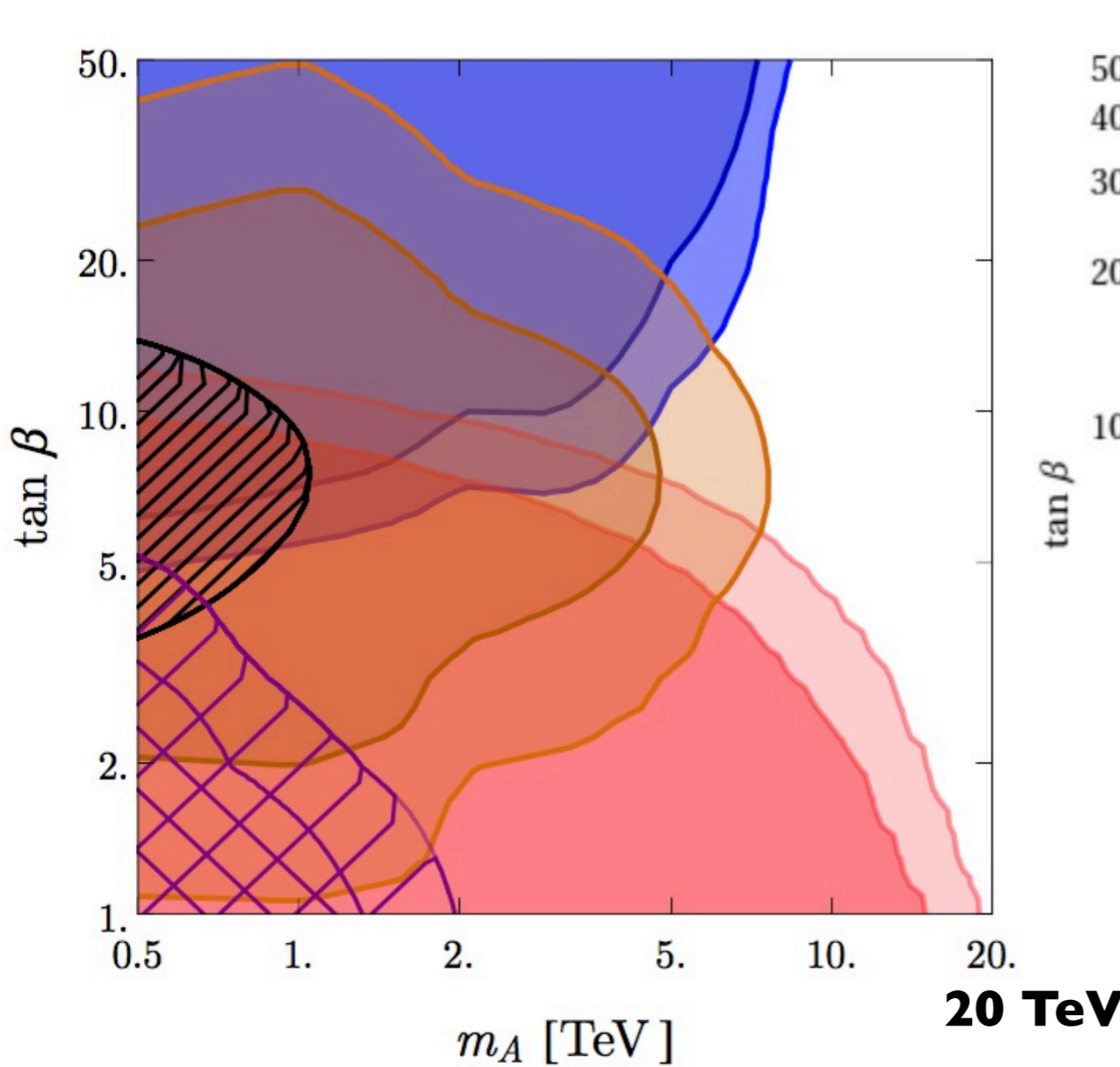
=> coverage beyond the upper limit of the thermal WIMP mass range for both higgsinos and winos !!

$$M_{wimp} \lesssim 2 \text{ TeV} \left(\frac{g}{0.3} \right)^2$$

MSSM Higgs @ 100 TeV

- $bbH^0/A^0 \rightarrow bb\tau\tau$
- $bbH^0/A^0 \rightarrow bbt\bar{t}$
- $t(t)H^0/A^0 \rightarrow t(t)t\bar{t}$

- $tbH^+ \rightarrow tbT\bar{V}$
- $tbH^+ \rightarrow tbt\bar{b}$
- LHC 3 ab^{-1}**
- LHC 0.3 ab^{-1}**



N. Craig, J. Hajer, Y.-Y. Li, T. Liu, H. Zhang,
arXiv:1605.08744

J. Hajer, Y.-Y. Li, T. Liu, and J. F. H. Shiu,
arXiv:1504.07617

FCC-ee + FCC-hh, project timeline

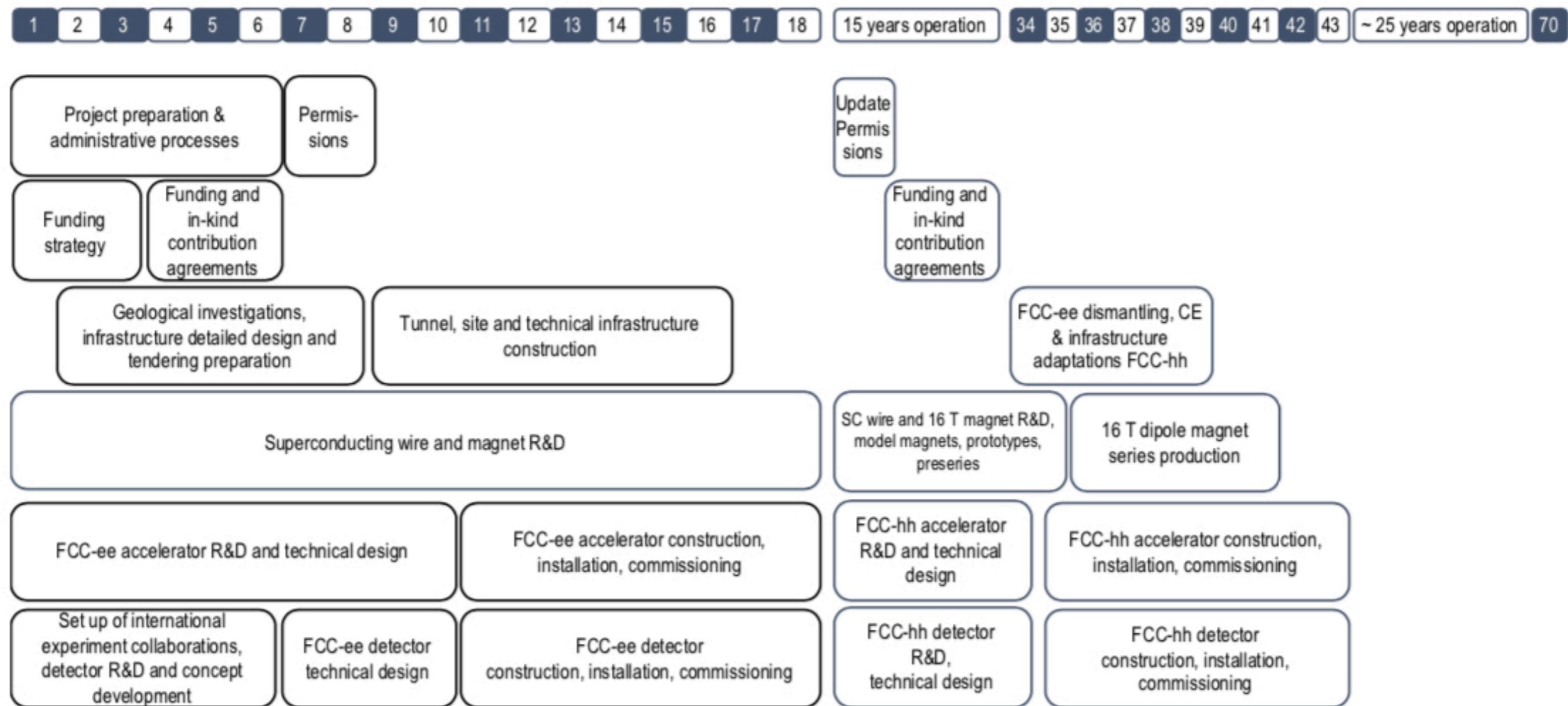
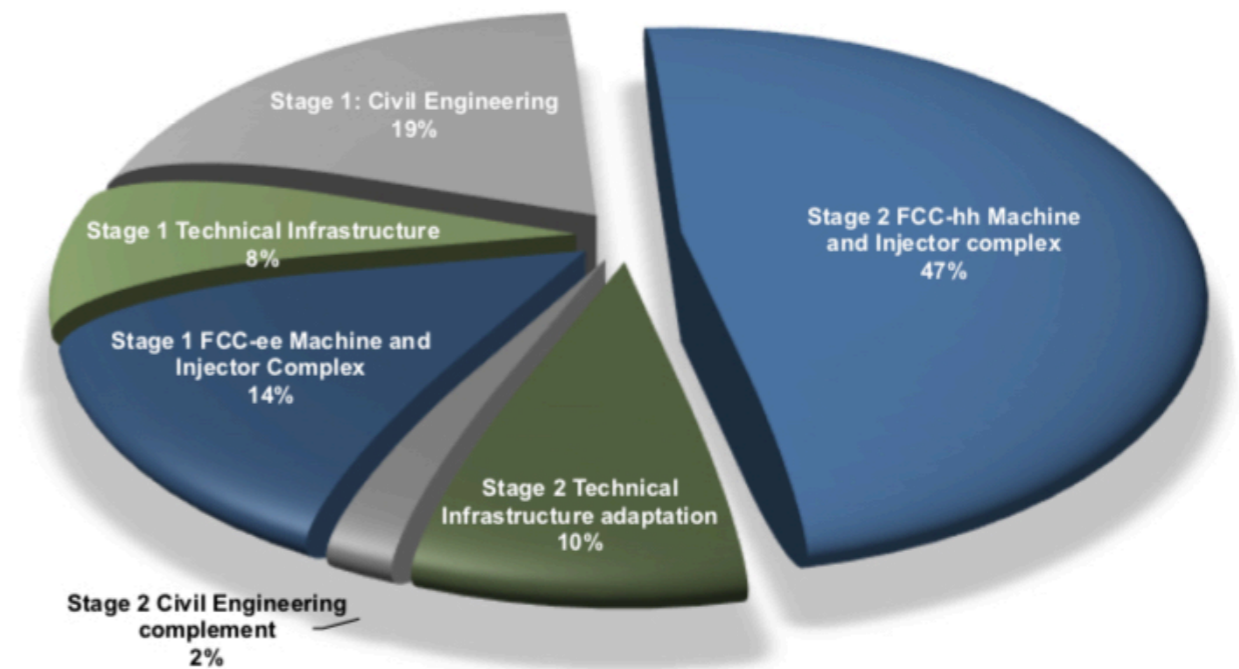


Table 5: Summary of capital cost to implement the integral FCC programme (FCC-ee followed by FCC-hh).

Domain	Cost in MCHF
Stage 1 - Civil Engineering	5,400
Stage 1 - Technical Infrastructure	2,200
Stage 1 - FCC-ee Machine and Injector Complex	4,000
Stage 2 - Civil Engineering complement	600
Stage 2 - Technical Infrastructure adaptation	2,800
Stage 2 - FCC-hh Machine and Injector complex	13,600
TOTAL construction cost for integral FCC project	28,600



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 - *complementary and synergetic precision studies of EW, Higgs and top properties*
 - *energy reach to allow direct discoveries at the mass scales possibly revealed by the precision measurements*

Additional material on physics at HE-LHC

For details see

P. Azzi, S. Farry, P. Nason, A. Tricoli, and D. Zeppenfeld, (conveners), et al, *Standard Model Physics at the HL-LHC and HE-LHC*, CERN-LPCC-2018-03, CERN, Geneva, 2018. <https://cds.cern.ch/record/2650160>.

M. Cepeda, S. Gori, P. J. Ilten, M. Kado, and F. Riva, (conveners), et al, *Higgs Physics at the HL-LHC and HE-LHC*, CERN-LPCC-2018-04, CERN, Geneva, 2018. <https://cds.cern.ch/record/2650162>.

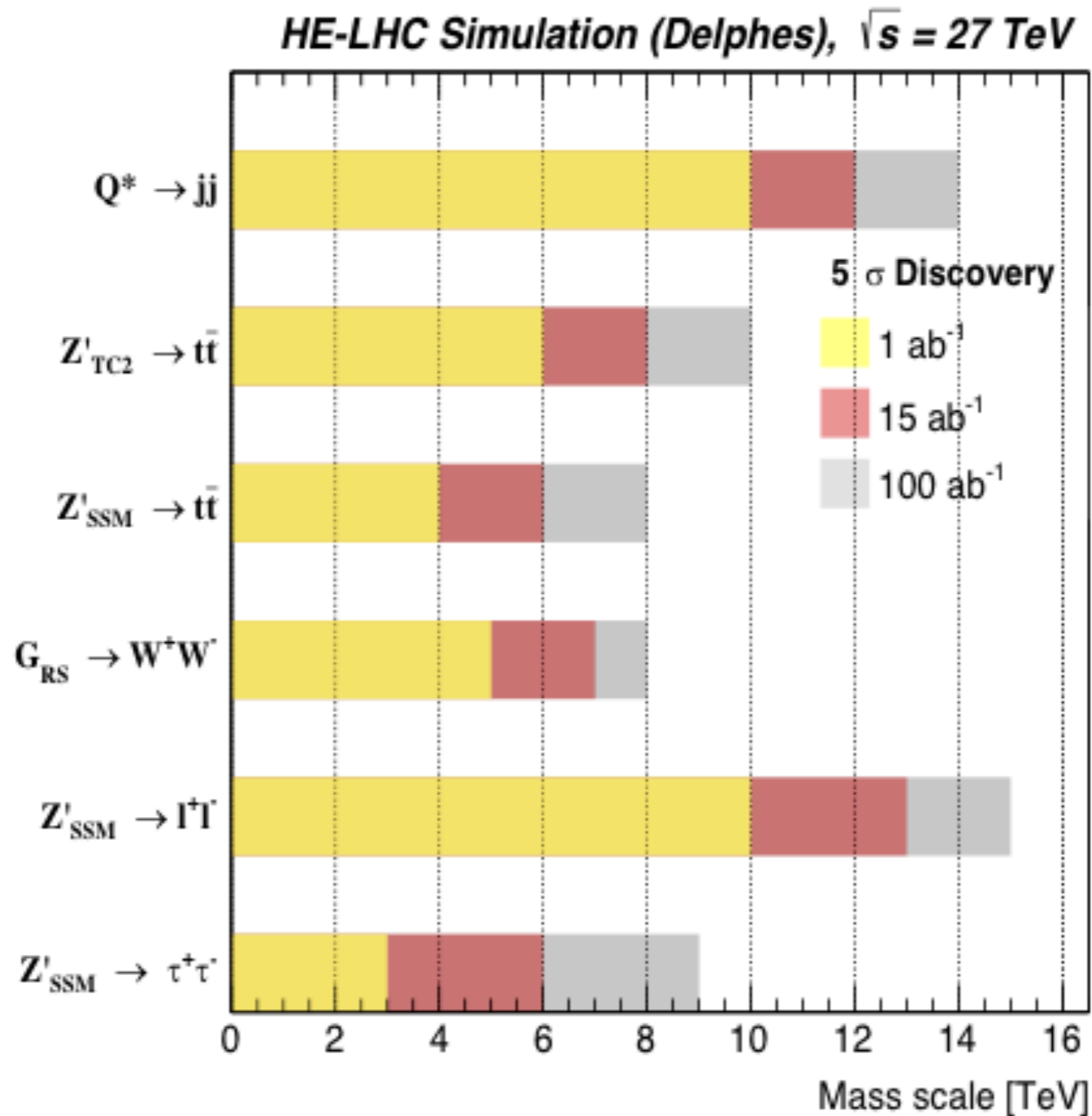
X. Cid-Vidal, M. D'Onofrio, P. J. Fox, R. Torre, and K. Ulmer, (conveners), et al, *Beyond the Standard Model Physics at the HL-LHC and HE-LHC*, CERN-LPCC-2018-05, CERN, Geneva, 2018. <https://cds.cern.ch/record/2650173>.

A. Cerri, V. V. Gligorov, S. Malvezzi, J. Martin Camalich, and J. Zupan, (conveners), et al, *Flavour Physics at the HL-LHC and HE-LHC*, CERN-LPCC-2018-06, CERN, Geneva, 2018. <https://cds.cern.ch/record/2650175>.

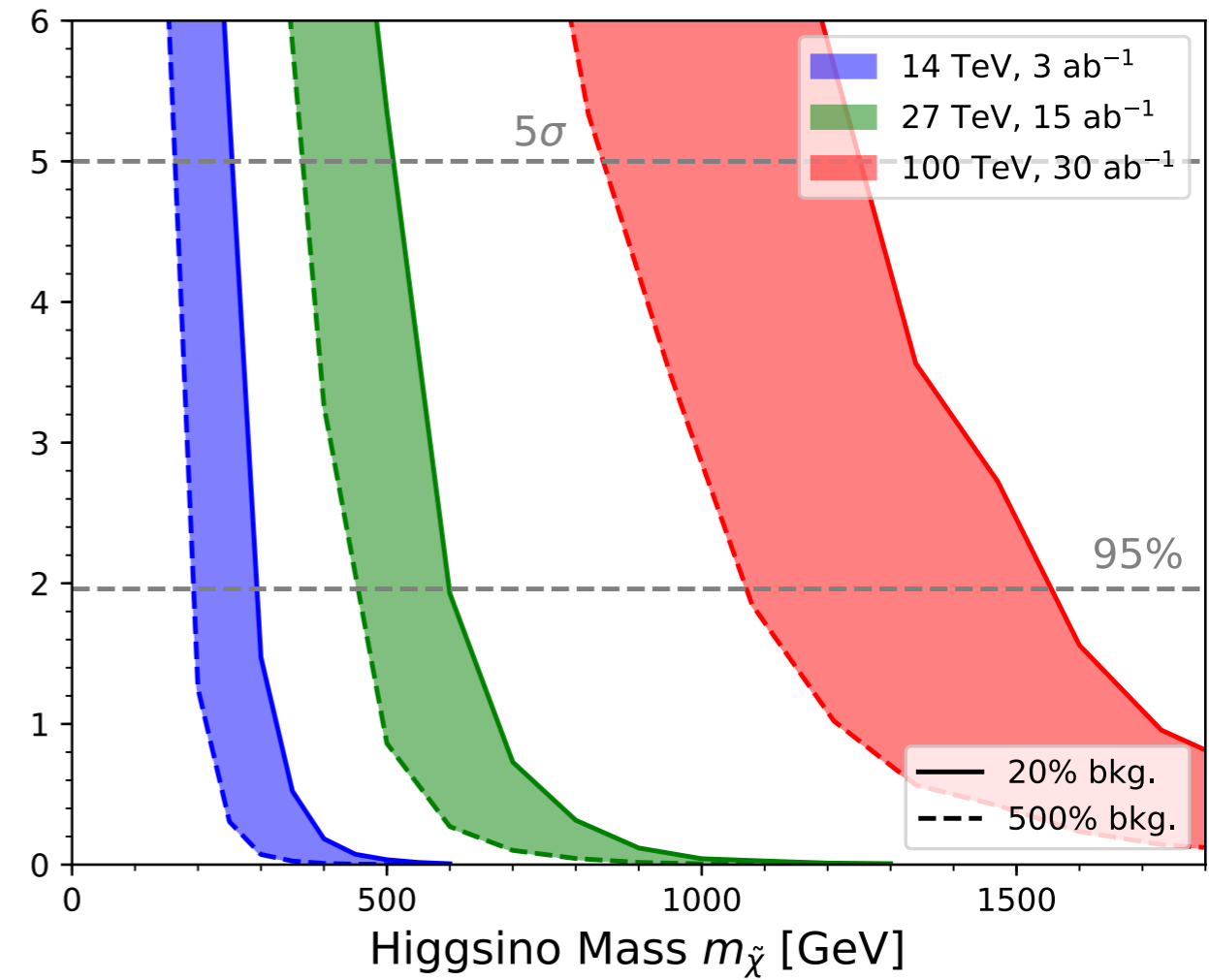
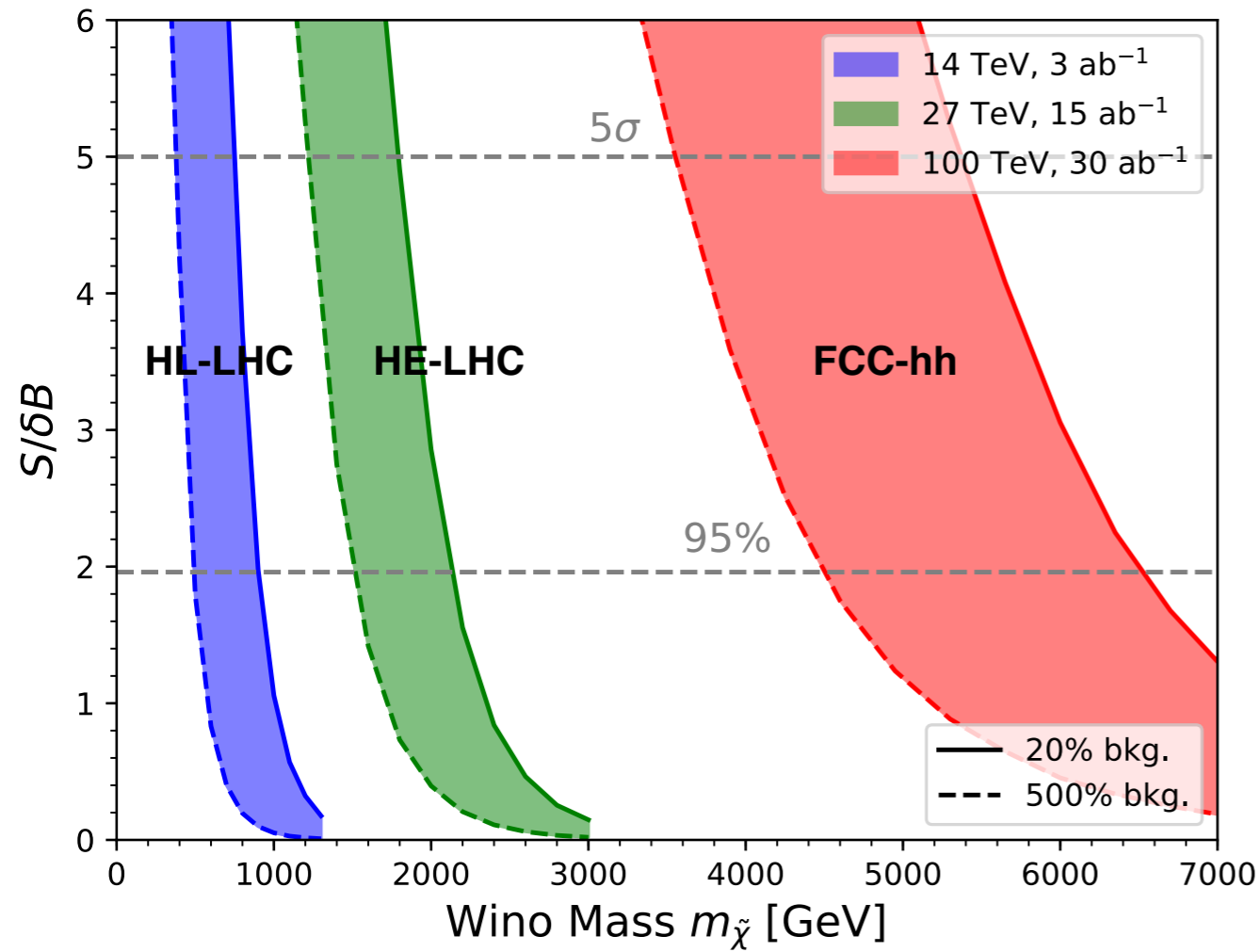
Z. Citron, A. Dainese, J. F. Grosse-Oetringhaus, J. M. Jowett, Y.-J. Lee, U. Wiedemann, and M. A. Winn, (conveners), et al, *Future physics opportunities for high-density QCD at the LHC with heavy-ion and proton beams*, CERN-LPCC-2018-07, CERN, Geneva, 2018. [arXiv:1812.06772](https://arxiv.org/abs/1812.06772) [hep-ph]. <https://cds.cern.ch/record/2650176>.

The ATLAS and CMS Collaborations, *Report on the Physics at the HL-LHC and Perspectives for the HE-LHC*, CERN-LPCC-2019-01, CERN, Geneva, 2019. <https://cds.cern.ch/record/2651134>.

(I) extension of mass reach for discovery: s-channel resonances



(I) EW-ino DM searches



(I+II) precision measurements and EWSB probes: Higgs observables

Examples of goals in the Higgs sector:

- (a) improve the sensitivity to the Higgs self-coupling
- (b) reduce to the few percent level all major Higgs couplings
- (c) improve the sensitivity to possible invisible Higgs decays
- (d) measure the charm Yukawa coupling

	$gg \rightarrow H$	WH	ZH	ttH	HH
N_{27}	2.2×10^9	5.4×10^7	3.7×10^7	4×10^7	2.1×10^6
N_{27}/N_{14}	13	12	13	23	19

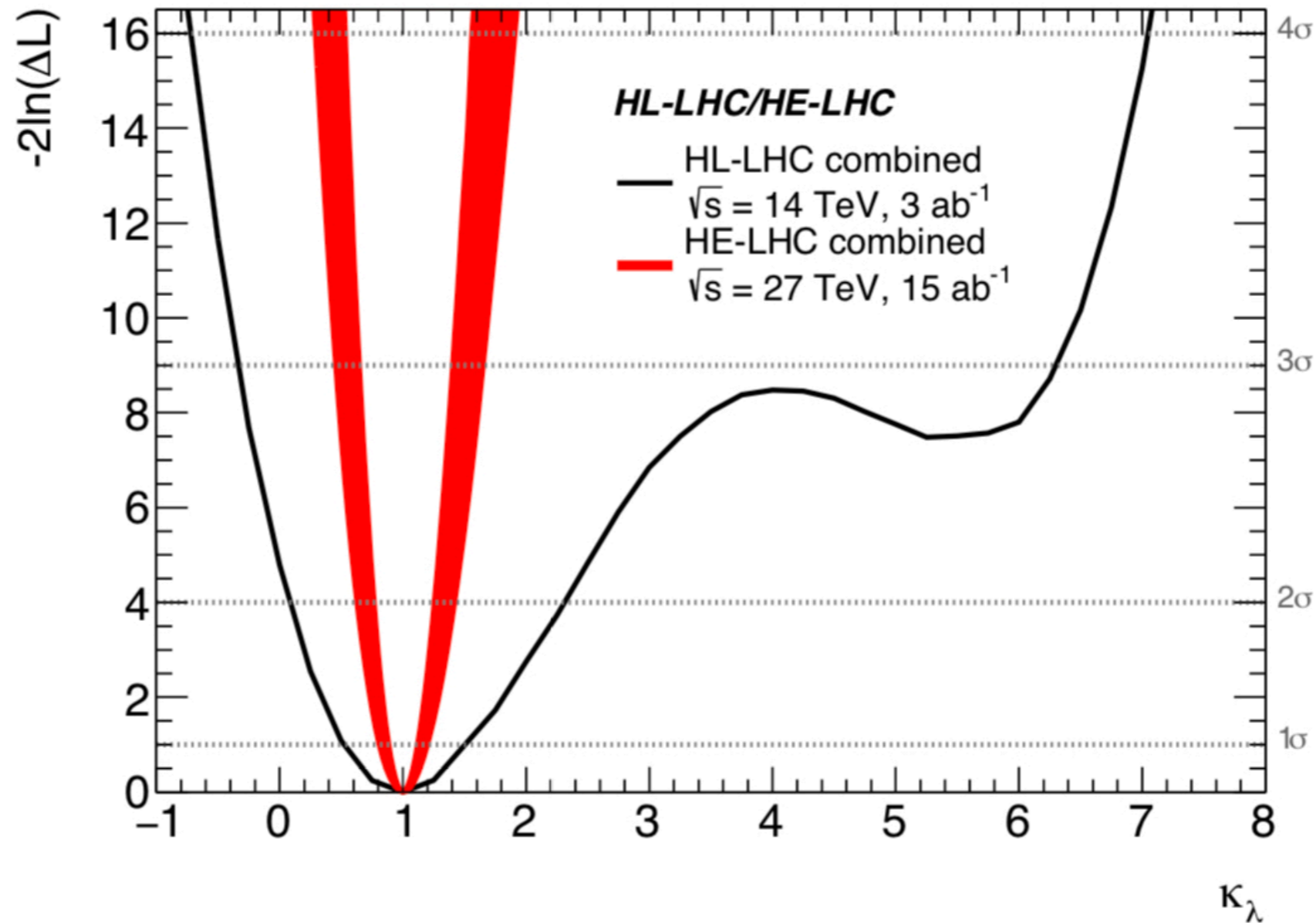
$$N_{27} = \sigma(27 \text{ TeV}) * 15 \text{ ab}^{-1}$$

$$N_{14} = \sigma(14 \text{ TeV}) * 3 \text{ ab}^{-1}$$

Higgs self-coupling at HE-LHC vs HL-LHC

HL-LHC: $\lambda/\lambda_{\text{SM}} \sim 1 \pm 0.5$ (68%CL)

HE-LHC: $\lambda/\lambda_{\text{SM}} \sim 1 \pm 0.15$ (68%CL)

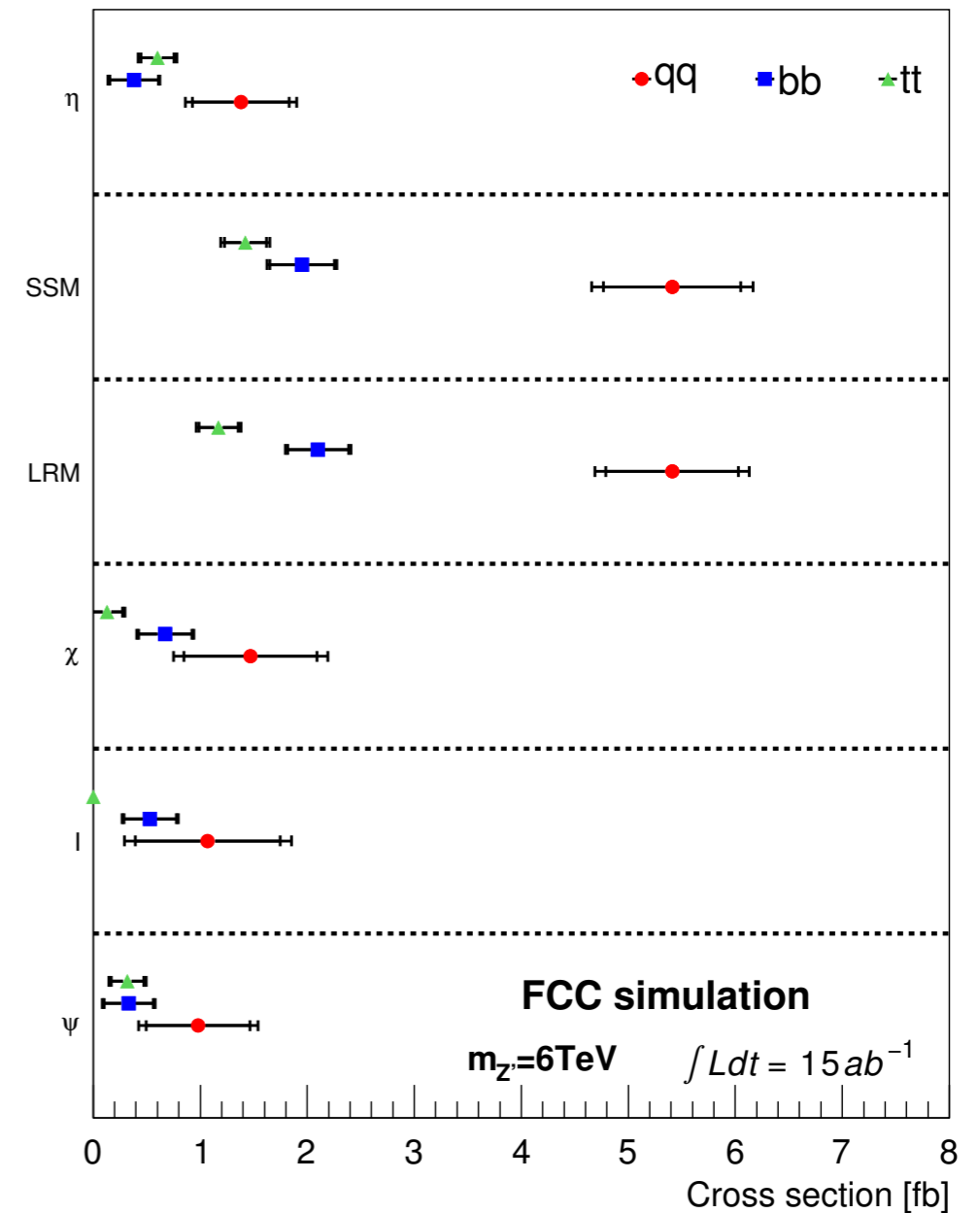
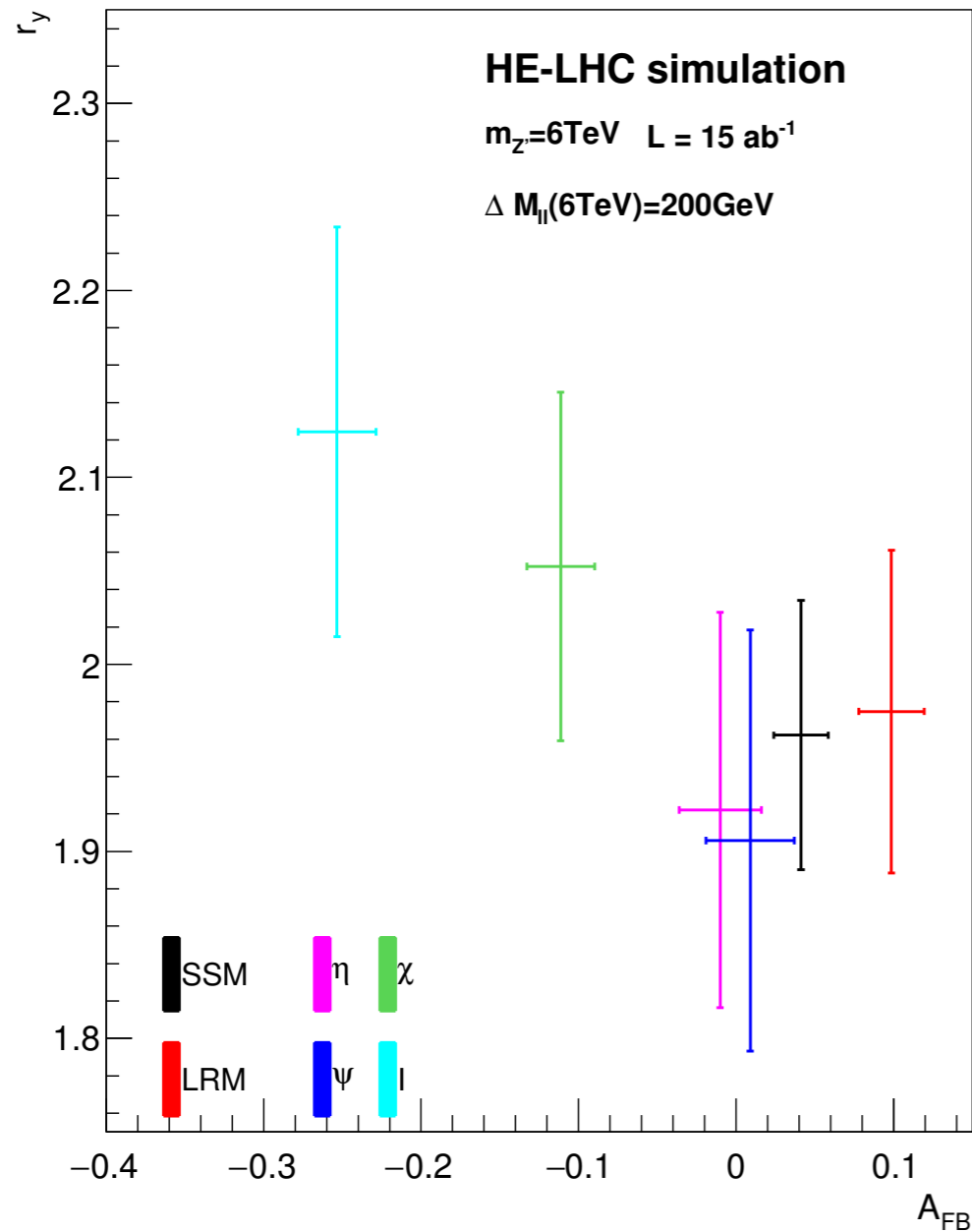


See also:

D. Gonçalves, T. Han, F. Kling, T. Plehn, and M. Takeuchi, *Higgs Pair Production at Future Hadron Colliders: From Kinematics to Dynamics*, arXiv:1802.04319 [hep-ph].

(IV) Exploration at 27 TeV of LHC discoveries: characterization of Z' models within reach of LHC observation

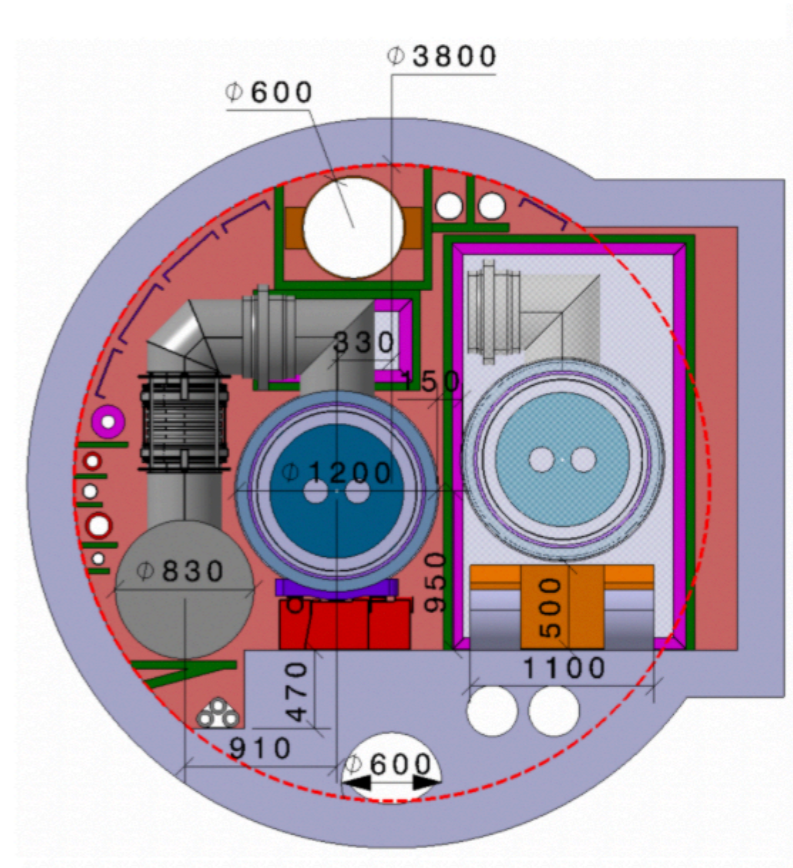
NB: uncertainty bars reflect very conservative syst assumptions



Colours: different Z' models, leading to observation at HL-LHC in $Z' \rightarrow$ dilepton decay for $m(Z')=6 \text{ TeV}$

HE-LHC: the challenges

- 16T Nb₃Sn magnets: more challenging than for FCC-hh, due to reduced space in the tunnel (requires dedicated R&D)



- SPS upgrade, to SC technology, to allow injection at 0.9-1.3 TeV
- Full replacement and strengthening of all infrastructure on the surface and underground cryogenics
- Significant civil engineering work both on the surface and in the tunnel (new SPS transfer lines, new caverns for cryogenics, 2 new shafts, ...)
- Overhaul/full replacement of detectors (radiation damage after HL-LHC, limited lifetime of key systems like magnets, use of new technologies, ...)
- ...

HE-LHC, project timeline/cost

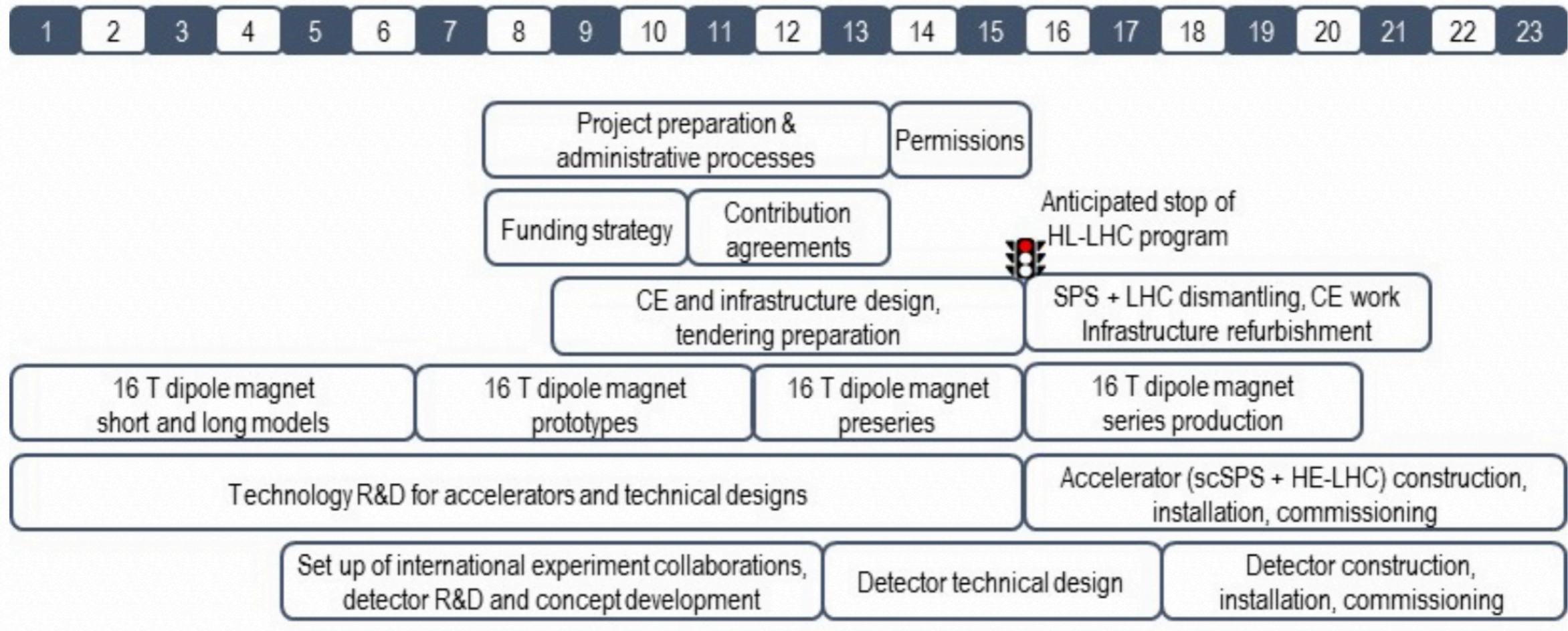


Figure 7: Overview of implementation timeline for the HE-LHC project starting in 2020. Numbers in the top row indicate the year. Physics operation would start in the mid 2040ies.

Domain	Cost in MCHF
Collider	5,000
Injector complex	1,100
Technical infrastructure	800
Civil Engineering	300
TOTAL cost	7,200

Table 2: Summary of capital cost for implementation of the HE-LHC project.