From collisions to analysis Lecture 1

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Lecture 1 outline:

- Introduction
- Trigger and DAQ basics
- Low-level trigger
- High-level trigger

Lecture 2 outline:

- Real-time analysis
- Interaction with analysis
- Upgrades



Acknowledgements to G. Raven and V. Gligorov for their slides from previous CERN/FNAL school lectures on these topics

Collisions to analysis



Collisions to analysis





Event reconstruction converts raw data hits to:

- Tracks, ECAL/HCAL clusters.
- e, γ, μ, τ
- Composite objects: missing E_T , H_T , vertices

Charged particle identification (π,K,p)



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Requires precise alignment and calibration of the detector.

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Full event reconstruction involved complicated algorithms, which are typically best suited to the high flexibility of CPUs.

Cost of full event reconstruction

A ballpark figure for LHC experiments is 1 second / CPU process, but compromises can be made in exchange for speed, e.g.:



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Note: not all CPU processes have the same speed, so better to talk in terms of *throughput* of a fully loaded multicore PC.



Even if we could afford to process/store all of the raw data, we couldn't power/cool the necessary electronics without spoiling the material budget.

Not such a constraint for LHCb



Its smaller events, and unique geometry, allow the LHCb upgrade to have trigger-less readout at a luminosity of 2x10³² cm⁻²s⁻¹. LHCB-TDR-016

Or for the relatively lower event rates in PbPb

The Alice upgrade will have continuous readout at 50 kHz in PbPb collisions.



Detector	Input to Online System (GByte/s)	Peak Output to Local Data Storage (GByte/s)	
TPC	1000	50.0	
TRD	81.5	10.0	
ITS	40	10.0	
Others	25	12.5	
Total	1146.5	82.5	

Collisions to analysis





What do the GHz of background interactions look like?

$$\frac{1}{N_{\rm ev}} \frac{dN_{\rm charged}}{d\eta} \sim 5$$

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@ pileup ~ 30, there are still ~10 charged particles with $p_T > 3$ GeV.



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Final state particles get softer with higher multiplicity decays and/or more complicated cascades.

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Signals without any leptons will always be difficult.

At an absolute minimum we must have single lepton triggers with p_T thresholds below ~25 GeV, without *prescales*.

Collisions to analysis



Collisions to analysis



Trigger and DAQ



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- The decision must be made and returned in a fixed *latency* of a few μ s, while the data are buffered in front-end *pipelines*.

Deadtime and de-randomisation



The *pipelines* buffer the data over the trigger latency.

They also de-randomise the data, reducing *deadtime* due to subsequent processing stages being *busy*.

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Trigger latency



What can we do in a few µs with part of the data?



- Simple pattern recognition and *locality*
- Low data rates

Versus what we can't yet do



- High data rates
- Complicated pattern recognition
- Need to link many sub-detectors

Real life example



Figure 1: The ATLAS TDAQ system in Run 2 with emphasis on the components relevant for triggering. L1Topo and FTK were being commissioned during 2015 and not used for the results shown here.

E.g., ATLAS thresholds

Table 1: Comparison of selected primary trigger thresholds (in GeV) at the end of Run 1 and during 2015 together with typical offline requirements applied in analyses (the 2012 offline thresholds are not listed but have a similar relationship to the 2012 HLT thresholds). Electron and tau identification are assumed to fulfil the 'medium' criteria unless otherwise stated. Photon and *b*-jet identification ('b') are assumed to fulfil the 'loose' criteria. Trigger isolation is denoted by 'i'. The details of these selections are described in Section 6.

Year	2012		2015			
\sqrt{s}	8 TeV		13 TeV			
Peak luminosity	$7.7 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$		$5.0 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$			
		p _T th	eshold [GeV], criteria			
Category	L1 HLT		L1	HLT	Offline	
Single electron	18	24i	20	24	25	
Single muon	15	24i	15	20i	21	
Single photon	20	120	22i	120	125	
Single tau	40	115	60	80	90	
Single jet	75	360	100	360	400	
Single <i>b</i> -jet	n/a	n/a	100	225	235	
$E_{\mathrm{T}}^{\mathrm{miss}}$	40	80	50	70	180	
Dielectron	2×10	2×12,loose	2×10	2×12,loose	15	
Dimuon	2×10	2×13	2×10	2×10	11	
Electron, muon	10, 6	12, 8	15, 10	17, 14	19, 15	
Diphoton	16, 12	35, 25	2×15	35, 25	40, 30	
Ditau	15i, 11i	27, 18	20i, 12i	35, 25	40, 30	
Tau, electron	11i, 14	28i, 18	12i(+jets), 15	25, 17i	30, 19	
Tau, muon	8, 10	20, 15	12i(+jets), 10	25, 14	30, 15	
Tau, E ^{miss}	20, 35	38, 40	20, 45(+jets)	35, 70	40, 180	
Four jets	4×15	4×80	3×40	4×85	95	
Six jets	4×15	6×45	4×15	6×45	55	
Two b-jets	75	35b,145b	100	50b,150b	60	
Four(Two) (b-)jets	4×15	2×35b, 2×35	3×25	2×35b, 2×35	45	
B-physics (Dimuon)	6, 4	6, 4	6, 4	6, 4	6,4	

ATLAS L1 bandwidth division



Wed, 4 Sep 2019 (showing first 14 of 22 entries)

[19] arXiv:1909.00761 [pdf, other]

Performance of electron and photon triggers in ATLAS during LHC Run 2

ATLAS Collaboration

Comments: 55 pages in total, author list starting page 39, 26 figures, 10 tables, submitted to EPJC. All figures including auxiliary figures are available at this https URL Subjects: High Energy Physics - Experiment (hep-ex)



Importance of calibration/resolution



LHCb



Not well suited to fast low-level triggers.

LHCb







- 1. Run at a *levelled* luminosity of 4x10³²cm⁻²s⁻¹
- 2. Profit from smaller events to take x10 higher L1 accept rate than ATLAS/CMS.









So far we have talked about *inclusive* selections, i.e. based on part of the signal.

What I just showed was an example of where an *exclusive* selection is required to classify the signals. This typically requires full offline quality event reconstruction.

LHCb L0 (1 MHz) trigger



L0 trigger	$E_{\rm T}/p_{\rm T}$ threshold			SPD threshold
	2015	2016	2017	
Hadron	$> 3.6 { m ~GeV}$	$> 3.7 { m ~GeV}$	$> 3.46 { m ~GeV}$	< 450
Photon	$> 2.7 { m ~GeV}$	$> 2.78 { m ~GeV}$	$> 2.47 { m ~GeV}$	< 450
Electron	$> 2.7 { m ~GeV}$	$> 2.4 { m ~GeV}$	$> 2.11 { m GeV}$	< 450
Muon	$> 2.8 { m ~GeV}$	$> 1.8 { m ~GeV}$	$> 1.35 { m ~GeV}$	< 450
Muon high $p_{\rm T}$	$> 6.0 { m ~GeV}$	$> 6.0 { m ~GeV}$	$> 6.0 { m ~GeV}$	none
Dimuon	$> 1.69 \ { m GeV^2}$	$> 2.25 \ { m GeV^2}$	$> 1.69 \ { m GeV^2}$	< 900

How are these thresholds decided?

LHCb L0 (≈1 MHz) bandwidth division

Why don't we just raise the luminosity?

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Going beyond the low-level triggers

LHCb software trigger farm in Run-II:~27000 physical cores running ~55000 processes.

The need for offline quality alignment and calibration

The LHCb RICH detectors

FIG. 6. Difference between the measured and expected Cherenkov angle, $\Delta \theta_C$ plotted as a function of the azimuthal angle ϕ and fitted with $\theta_x \cos(\phi) + \theta_y \sin(\phi)$, for one side of the RICH 2 detector [6]. The upper plot is prior to alignment, and shows a dependency of the angle θ_C on the angle ϕ . The bottom plot is after the alignment correction, and $\Delta \theta_C$ is uniform in ϕ .

E.g., offline quality RICH PID for HLT2

Performance slightly better than the offline version from Run-I.

We'll see tomorrow how RICH PID is a crucial requirement for the *Turbo stream*.

JINST 14 (2019) P0401353

Partial event reconstruction in HLT1

HLT1 selections

Most (≈100 kHZ) of the rate is taken by inclusive one- and two-track heavy-flavour lines.

JINST 14 (2019) P0401355

Inclusive HLT1 track line(s) performance

- $D^+ \rightarrow K^- \pi^+ \pi^+$ \checkmark D⁰ \rightarrow K⁻ $\pi^+\pi^-\pi^+$
 - $\Lambda_{c}^{+} \rightarrow pK^{-}\pi^{+}$

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There is an interesting interplay with the reconstruction because, e.g., the two-track line is only really useful if we can reconstruct tracks down to a p_T threshold of ~500 MeV.

The 10 PB disk buffer in, e.g., 2017

- HLT1 output rate ~150 kHz
- HLT2 throughput ~80 kHz out-of-fill (and ~30 kHz in fill).
- Average machine efficiency 30-50%.

The 10 PB disk buffer in, e.g., 2017

- HLT1 output rate ~150 kHz
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HLT2 selections

About 500 HLT2 "lines" by the end of Run-II

The zoo of b hadron decay modes

http://pdglive.lbl.gov/

Indentation is used to indicate a subchannel of a previous reaction. All resonant subchannels have been corrected for resonance branching fractions to the final state so the sum of the subchannel branching fractions can exceed that of the final state. For inclusive branching fractions, e. g., $B \rightarrow D^{\pm}X$, the values usually are multiplicities, not branching fractions. They can be greater than one.

Mode			Eraction (Γ_i / Γ)	Scale Factor/ Conf. Level	P (MeV/c)
Semileptor	nic and leptonic modes				
Γ ₁	$\ell^+ \nu_{\ell} X$	[1]	$(10.99 \pm 0.28)\%$		
Γ ₂	$e^+ v_e X_c$		$(10.8 \pm 0.4)\%$		
Γ ₃	$\mathrm{D}\ell^+ v_\ell \mathrm{X}$		(9.7 ± 0.7)%		
Γ_4	-0 D $\ell^+ \nu_{\ell}$	[1]	$(2.35 \pm 0.09)\%$		2310
Γ ₅	$D_0^{\tau} v_{\tau}$ D $\tau^+ v_{\tau}$		$(7.7 \pm 2.5) \times 10^{-3}$		1911
Г ₆	$D^{+}(2007)^{0}\ell^{+}\nu_{\ell}$	[1]	(5.66 ± 0.22)%		2258
Γ ₇	$D^{*}(2007)^{0}\tau^{+}\nu_{\tau}$		(1.88 ± 0.20)%	S=1.0	1839
Γ ₈	$D^{-}\pi^{+}\ell^{+}\nu_{\ell}$		$(4.4 \pm 0.4) \times 10^{-3}$		2306
Γ ₉	$-* - * 0$ $D_{\alpha}(2420)^{0}\ell^{+}\nu D_{\alpha} \rightarrow D^{-}\pi^{+}$		$(2.5 \pm 0.5) \times 10^{-3}$		
Γ_{10}	$(D_{\mathfrak{g}}(2460))^{0}\ell^{+}\nu_{\mathfrak{g}}, D_{\mathfrak{g}} \longrightarrow D^{-}\pi^{+}$		$(1.53 \pm 0.16) \times 10^{-3}$	S=1.0	2065
Г ₁₁	$D^{(*)2}n\pi\ell^{\nu}v_{\ell}(n \ge 1)^{2}$		(1.86 ± 0.26)%		
Γ ₁₂	$D^* \pi^+ \ell^+ v_\ell$		$(6.0 \pm 0.4) \times 10^{-3}$		2254
Γ ₁₃	$\underline{D}_{1'}(2420)^{0}\ell^{+}\nu_{\ell}, \underline{D}_{1'0} \rightarrow D^{*}\pi^{+}$		$(3.03 \pm 0.20) \times 10^{-3}$		2084

Used as key variable in with a "Bonzai" BDT [JINST 8 (2013) P02013] to provide a few kHz of ~pure bbbar.

Few kHz of inclusive B lines gives:

• $B^+ \rightarrow \overline{D}{}^0 \pi^+ \square B^0 \rightarrow D^- \pi^+ \land B^+ \rightarrow J/\psi(e^+e^-)K^+ \land B^+ \rightarrow J/\psi(\mu^+\mu^-)K^+$

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Backup slides follow from here...