

The CMS Precision Proton Spectrometer (PPS): results, status and prospects

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on behalf of the CMS and TOTEM Collaborations

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Outline

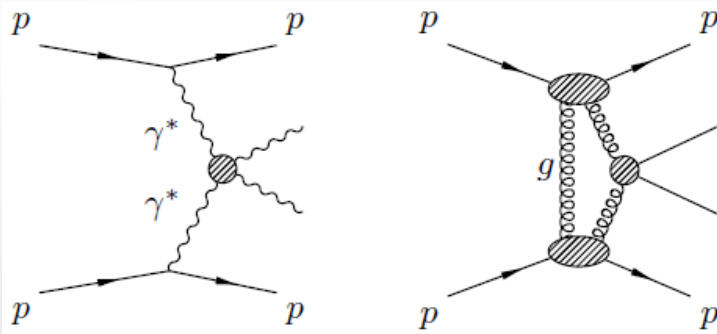


- PPS physics case
- Experimental apparatus
- Operation in LHC Run 2
- Tracking efficiency
- Detector alignment and proton reconstruction
- (Semi)Exclusive $\gamma\gamma \rightarrow \ell^+ \ell^-$ results with 2016 data
- Prospects for LHC Run 3

PPS physics case

- **Central exclusive production** ($J^{PC} = 0^{++}$ central final state)

- Colour-singlet exchanges with large rapidity gaps between the central system and the outgoing protons
- Two-photon, photon-pomeron or two-pomeron exchanges at LHC energies allow access to a large variety of processes
 - see e.g. *JHEP* 1608 (2016) 119, *Phys.Lett. B* 777 (2018) 303-323 (in pp), *Nature Phys.* 13 (2017) no.9, 852-858, *arXiv:1810.04602* (UPC in PbPb)



- **Electroweak physics:** diboson and dilepton production, anomalous coupling searches
- **QCD:** dijet, trijet, $t\bar{t}$ production
- **BSM direct searches:** new resonances, missing mass...

- **Advantages of the forward proton measurement**

- Strong background suppression by requiring kinematic match with the central system
- Reduced theory uncertainties related to proton dissociation



PPS physics
case

Experimental
apparatus

Operation in
LHC Run 2

Tracking
efficiency

Detector
alignment and
proton
reconstruction

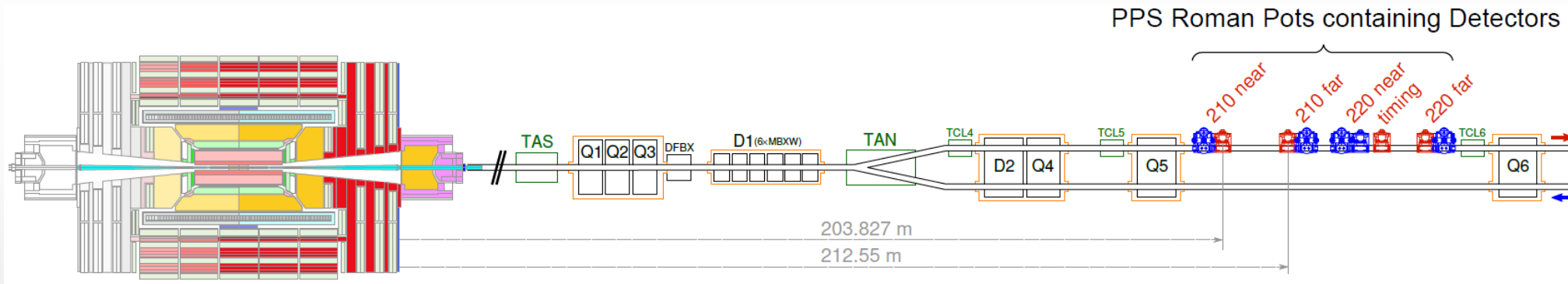
(Semi)Exclusive
 $\gamma\gamma \rightarrow \ell^+\ell^-$
2016 results

Prospects for
LHC Run 3

Introduction to PPS

- **Near beam magnetic spectrometer at IP5 of the LHC**

- Joint CMS+TOTEM project to include horizontal Roman Pots in the standard, high luminosity, CMS data taking, TDR in 2014 [*CERN-LHC-2014-021*]
- Started in 2016 (one year early) thanks to the availability of the legacy TOTEM silicon strips detectors
- Collected data throughout the whole LHC Run 2



- **Two complementary measurements**

- **Tracking detectors** measure the proton displacement with respect to the beam, which is translated into energy-momentum loss thanks to the knowledge of the beam optics
- **Timing detectors** measure the proton arrival time in both arms w.r.t. the reference clock that keeps the two stations synchronized, allowing the calculation of the longitudinal position of the interaction



PPS physics
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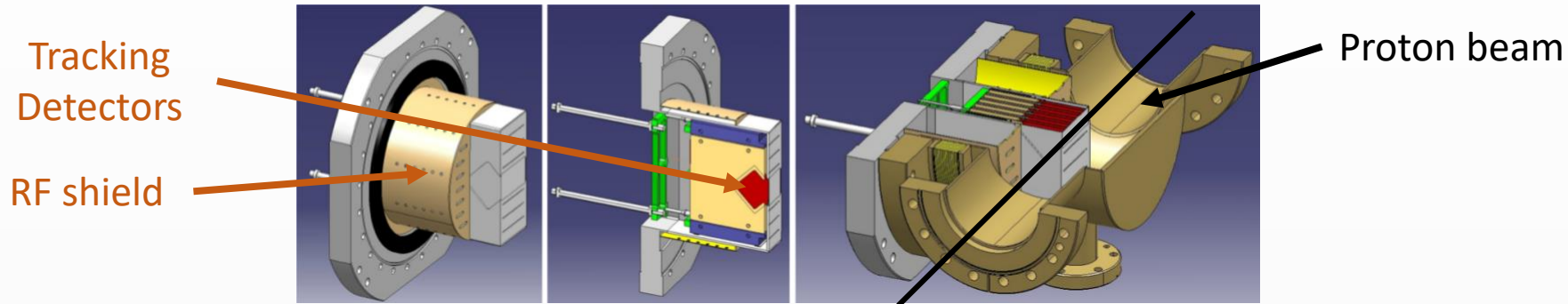
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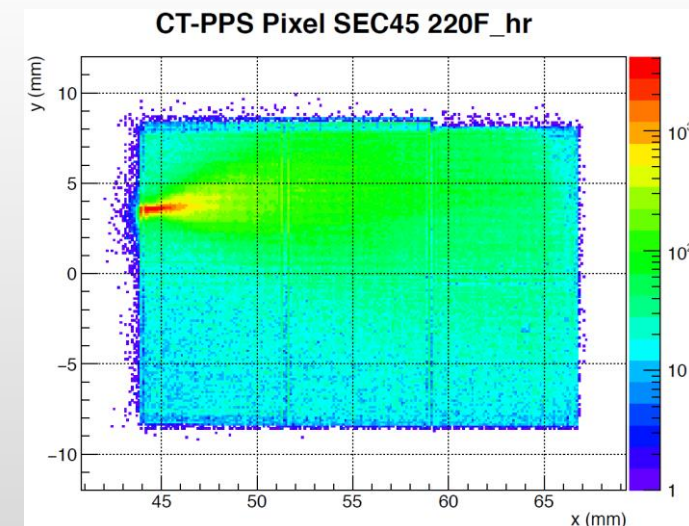
Prospects for
LHC Run 3

Experimental challenges

- **Roman Pots need to operate at few mm from the beam to maximize acceptance**
 - RF shielding installed to limit the impedance caused by the RP insertion



- **Detectors must tolerate high levels of non-uniform irradiation**
 - For $\sim 100 \text{ fb}^{-1}$ (Run 2): $\sim 5 \cdot 10^{15}$ protons/cm² in tracking detectors
 - Spatial resolution required: $\sim 10\text{-}30 \text{ }\mu\text{m}$
 - Timing resolution required: $\sim 20 \text{ ps}$, to reject high pileup



PPS physics case

Experimental apparatus

Operation in LHC Run 2

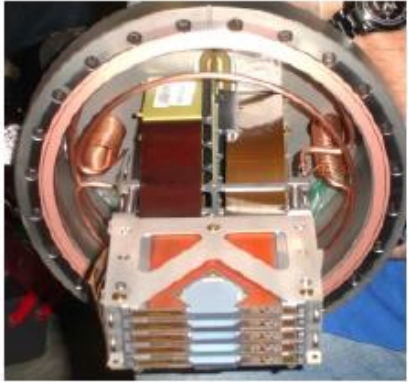
Tracking efficiency

Detector alignment and proton reconstruction

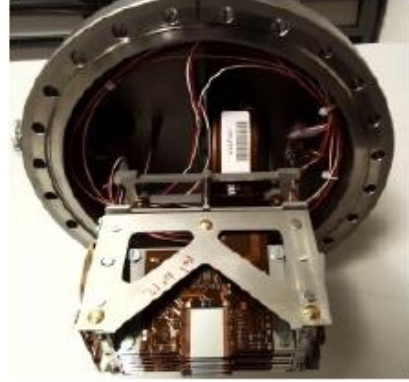
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Prospects for LHC Run 3

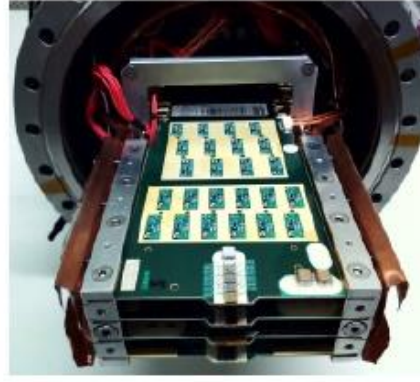
PPS detectors for Run 2



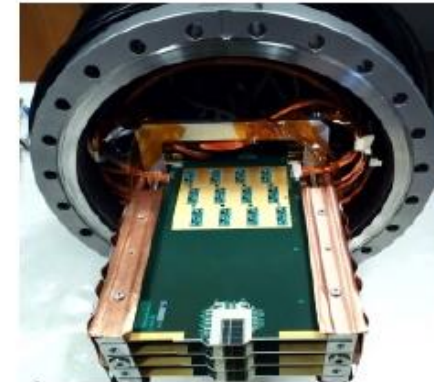
TOTEM si-strips



3D pixels



scCVD (diamond)



Ultra-fast silicon detectors

- **2016 Detectors**

- Tracking: 2 stations of TOTEM Si-strips detectors (10 planes), 20 μm resolution. Limited radiation resistance ($\Phi_{\text{max}} \sim 5 \cdot 10^{14} \text{p/cm}^2$), no multi-track capability.
- Timing: diamond detectors in cylindrical RP

- **2017 Detectors**

- Tracking: 1 station of TOTEM si-strips, 1 station of silicon 3D pixels (6 planes with CMS Phase 1 tracker readout chips), $\sigma_x \sim 15 \mu\text{m}$ and $\sigma_y \sim 30 \mu\text{m}$, $\Phi_{\text{max}} \sim 5 \cdot 10^{15} \text{p/cm}^2$
- Timing: 1 station with 3 planes of single-layer diamond with expected $\sigma_t = 80 \text{ps/plane}$ and 1 plane of UFSD with expected $\sigma_t = 30 \text{ps/plane}$ ($\Phi_{\text{max}} \sim 10^{14} \text{p/cm}^2$)

- **2018 Detectors**

- Tracking: two 3D pixels stations
- Timing: 1 station of diamond detectors (2 single-layer + 2 double-layer)

In both arms



PPS physics case

Experimental apparatus

Operation in LHC Run 2

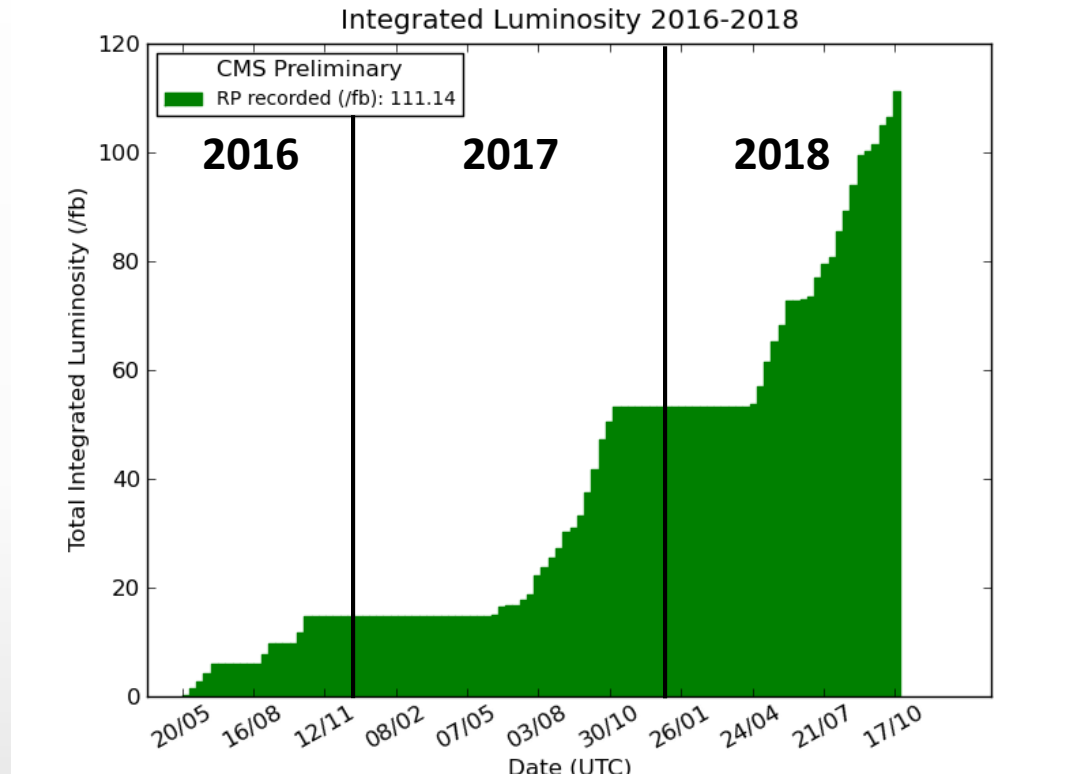
Tracking efficiency

Detector alignment and proton reconstruction

(Semi)Exclusive $\gamma\gamma \rightarrow \ell^+\ell^-$
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Prospects for LHC Run 3

LHC Run 2 operation



2016: 40% of the CMS statistics
2017: 88% of the CMS statistics
2018: 92% of the CMS statistics

- **$\sim 115 \text{ fb}^{-1}$ collected by PPS during Run 2**
 - Always inserted and taking data, fully integrated in CMS runs, in 2017 and 2018
 - Very high stability in both 2017 and 2018



PPS physics case

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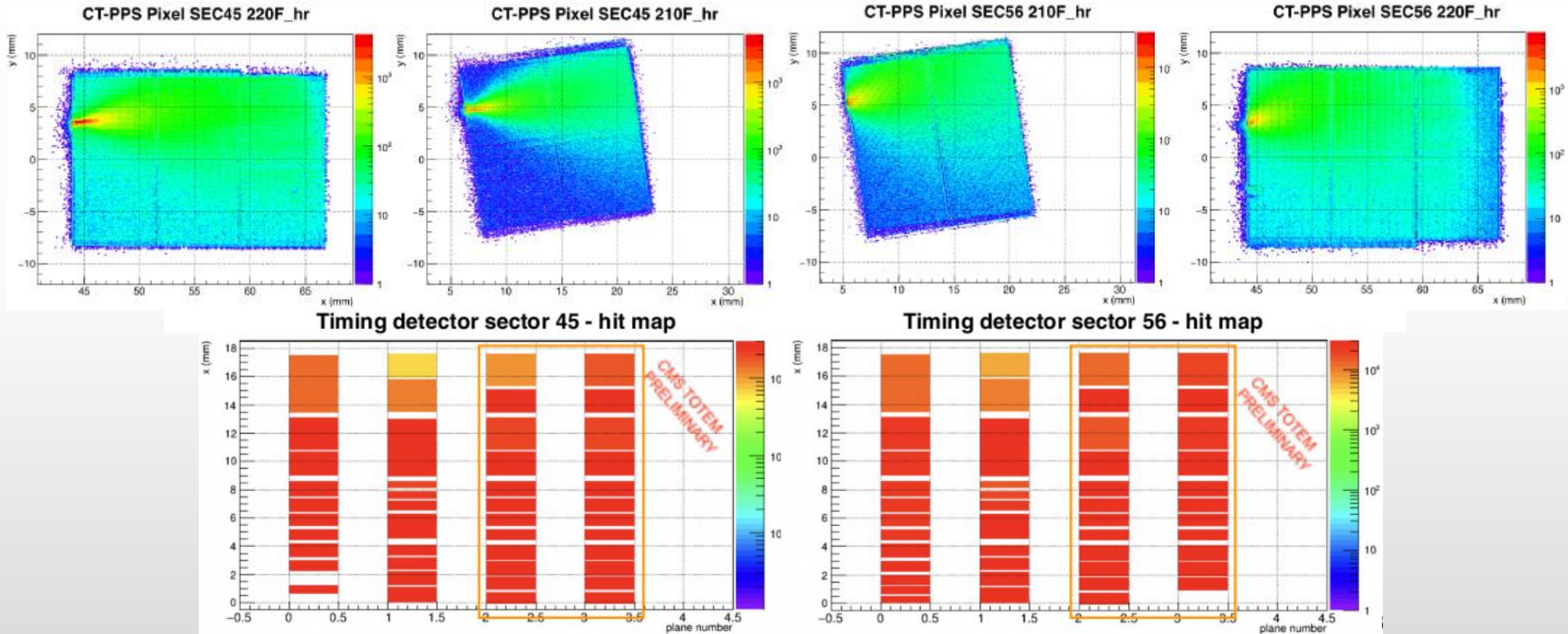
(Semi)Exclusive $\gamma\gamma \rightarrow \ell^+\ell^-$
2016 results

Prospects for LHC Run 3

2018 data: hit maps

Sector 45 ← IP5 → Sector 56

2018 Layout



- Less than 0.05% bad or noisy pixels



PPS physics case

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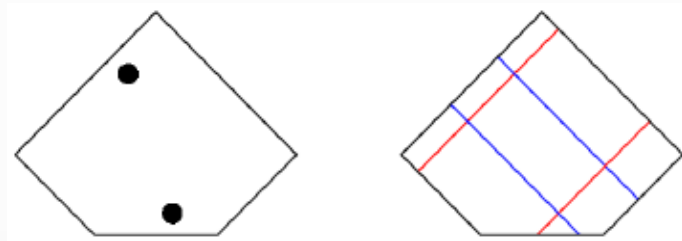
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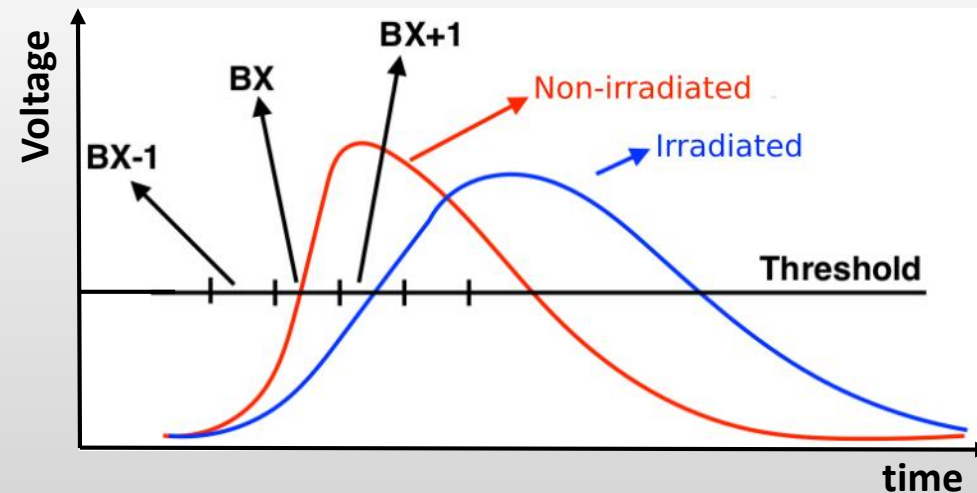
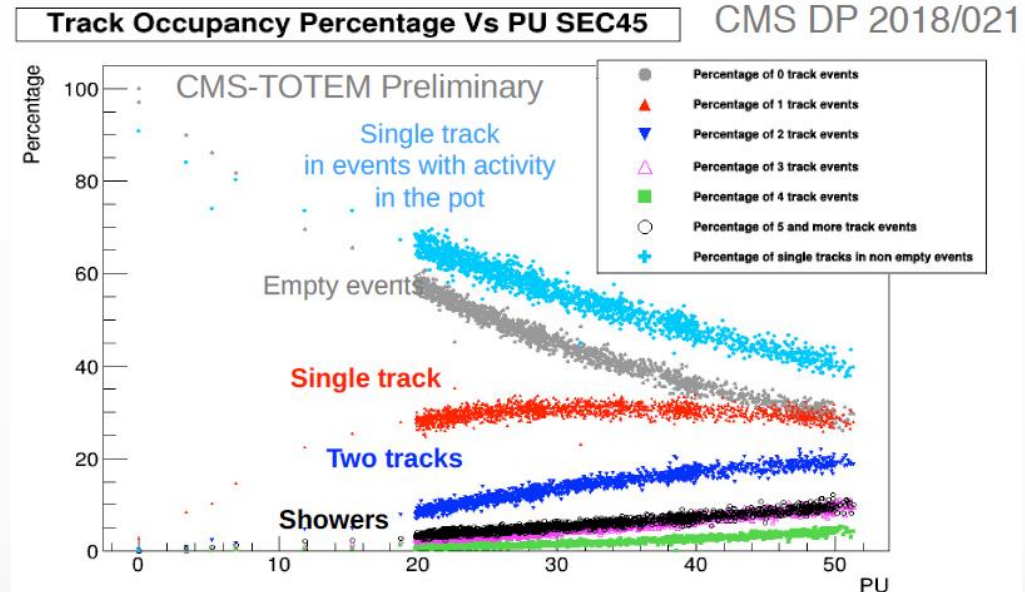
Prospects for LHC Run 3

Tracking efficiency studies

- **Main strips inefficiency due to no multi-tracking:**
 - 30% efficiency at pileup of 50
 - No problem for pixel detectors



- **Radiation damage effects on the sensors reduced when moving to pixels:**
 - However, non-uniform irradiation does affect the readout chip performance
 - Efficiency loss mitigated by moving the stations vertically during technical stops



PPS physics case

Experimental apparatus

Operation in LHC Run 2

Tracking efficiency

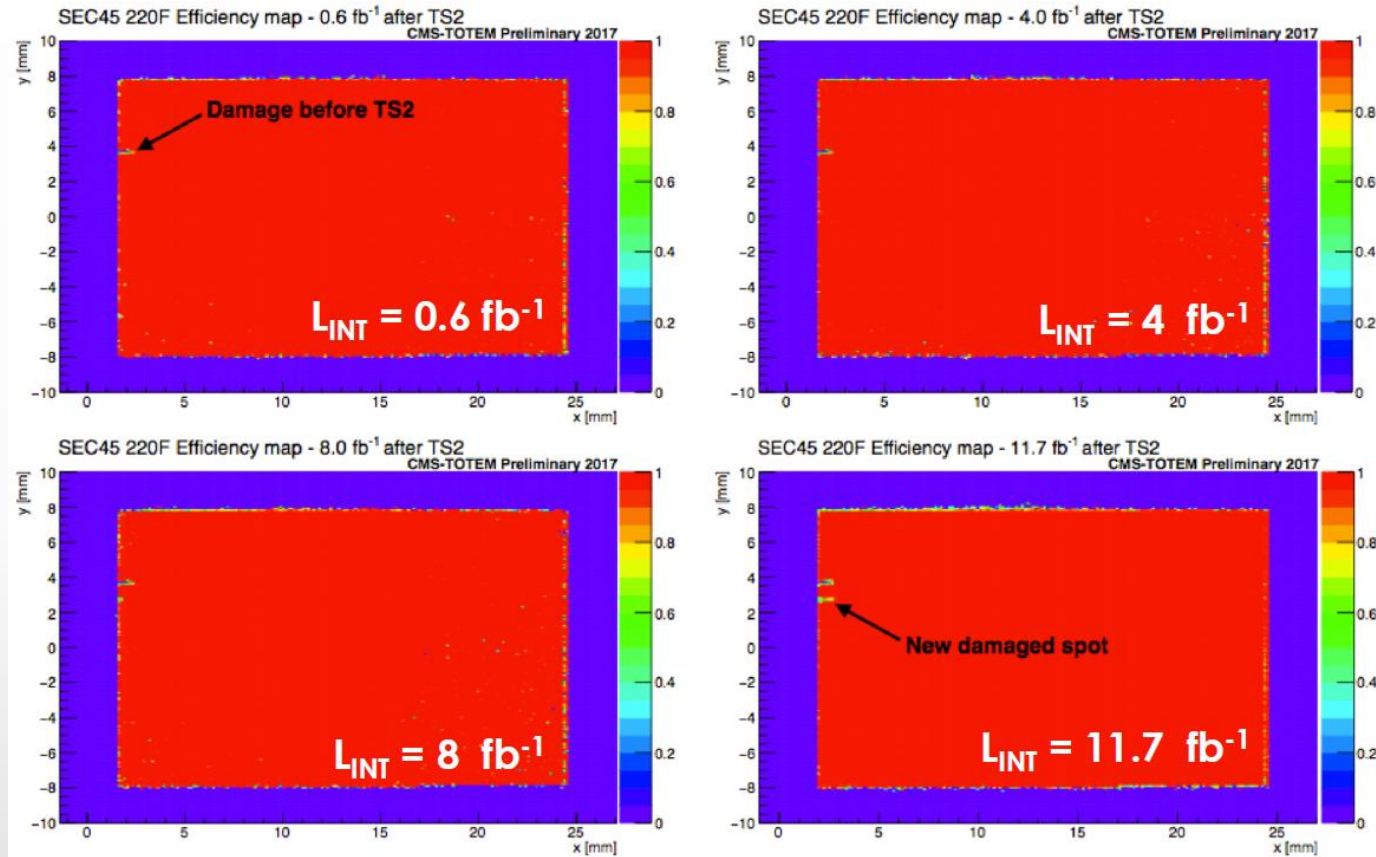
Detector alignment and proton reconstruction

(Semi)Exclusive $\gamma\gamma \rightarrow \ell^+\ell^-$ 2016 results

Prospects for LHC Run 3

2017 Pixel detector efficiency

- **Efficiency computed with tracks reconstructed within the same station**
 - Evolution of the radiation damage vs. integrated luminosity after LHC second Technical Stop ($\sim 18 \text{ fb}^{-1}$ taken before TS2)
 - Inefficiency spot caused by radiation damage is moved away from the high-occupancy region when the station is lifted
 - The radiation effect starts to be visible at $\sim 8 \text{ fb}^{-1}$
- **Very high performance overall: average efficiency $\sim 98\%$**
 - Few damaged pixels: $\sim 1.5 \times 0.3 \text{ mm}^2$, caused by non-uniform irradiation of the readout chip



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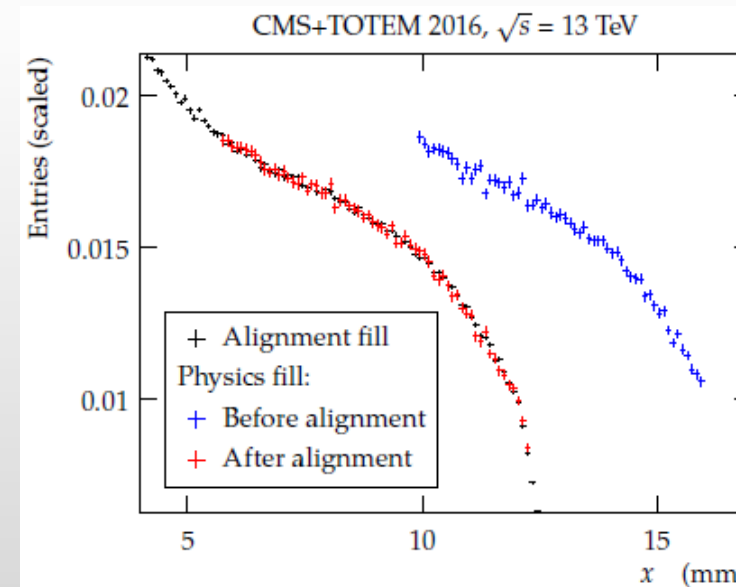
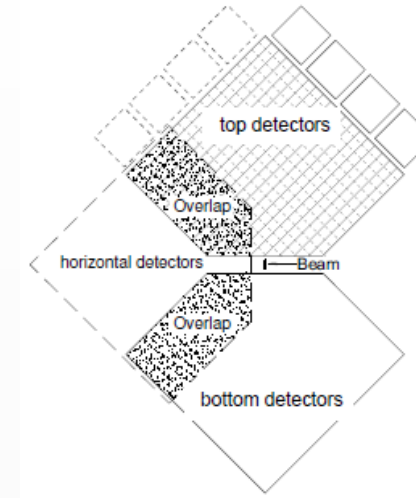
Prospects for
LHC Run 3

Detector alignment

Two step detector alignment procedure:

- Absolute RP alignment in dedicated low-luminosity runs:
 - Beam-based alignment, to establish the position of the RPs w.r.t. the LHC collimators and the beam
 - Exploit the geometrical properties of elastic scattering events ($pp \rightarrow pp$) to extract positions of all the RPs, with both horizontal and vertical pots inserted very close to the beam
- Fill-by-fill alignment:
 - Need to redetermine the RPs positions w.r.t. the beam: the RP position during standard data taking differs from that of the alignment run; the beam position may also change due to the RP or beam movement
 - Use inclusive sample of protons and match the proton tracks distribution with those of the alignment run

See CERN-LHC-2014-021 and CERN-TOTEM-NOTE-2017-001 for further details



PPS physics case

Experimental apparatus

Operation in LHC Run 2

Tracking efficiency

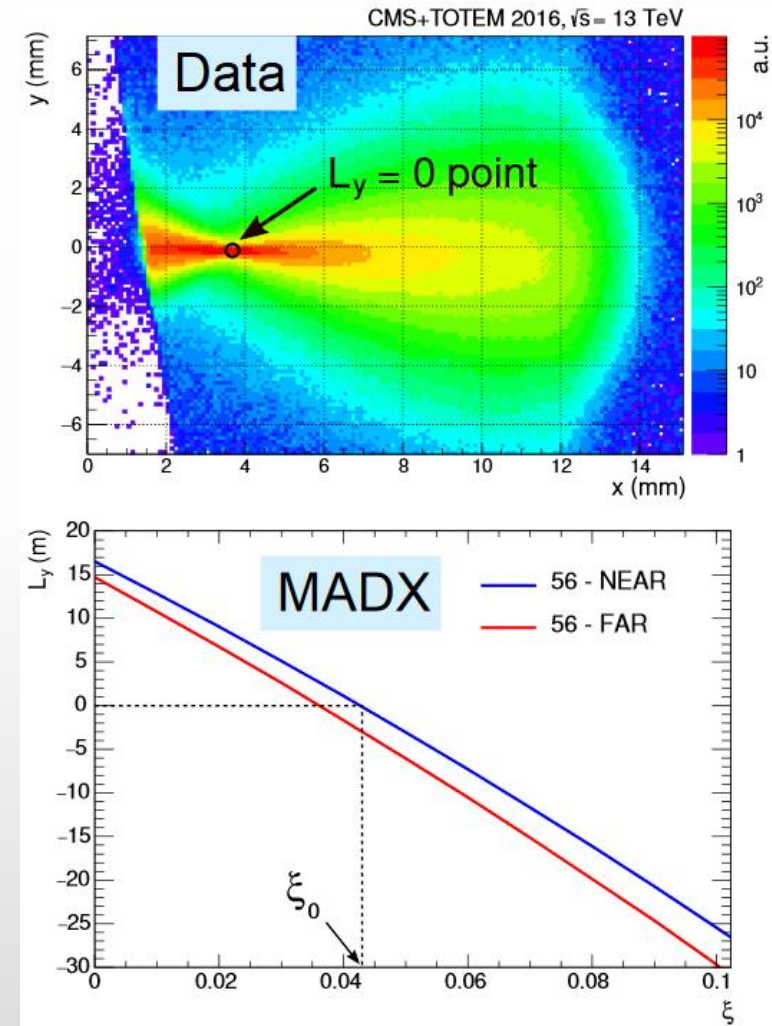
Detector alignment and proton reconstruction

(Semi)Exclusive $\gamma\gamma \rightarrow l^+l^-$ 2016 results

Prospects for LHC Run 3

LHC optics and dispersion corrections

- **Very good knowledge of the LHC beam optics is needed in order to correctly reconstruct the proton fractional momentum loss ξ**
 - Significant data-driven corrections need to be made to the nominal optics
 - MADX software is used to simulate LHC optics
 - The model parameters are tuned to measurements performed with RPs and beam-position monitors
 - The dispersion is calibrated by using the effective length pinch point ($L_y(x) = 0$)
- In the end: **non-linear calibration of ξ vs. the measured track position, $\xi = x/D_x(\xi)$**
- *For the full documentation, see:
New J. Phys. 16 (2014) 103041, CERN-TOTEM-NOTE-2017-002*



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Prospects for LHC Run 3

Proton reconstruction

- **Knowledge of beam optics allows the proton fractional momentum loss ξ to be computed:**

- From ξ the invariant mass and rapidity of the centrally produced state X is determined

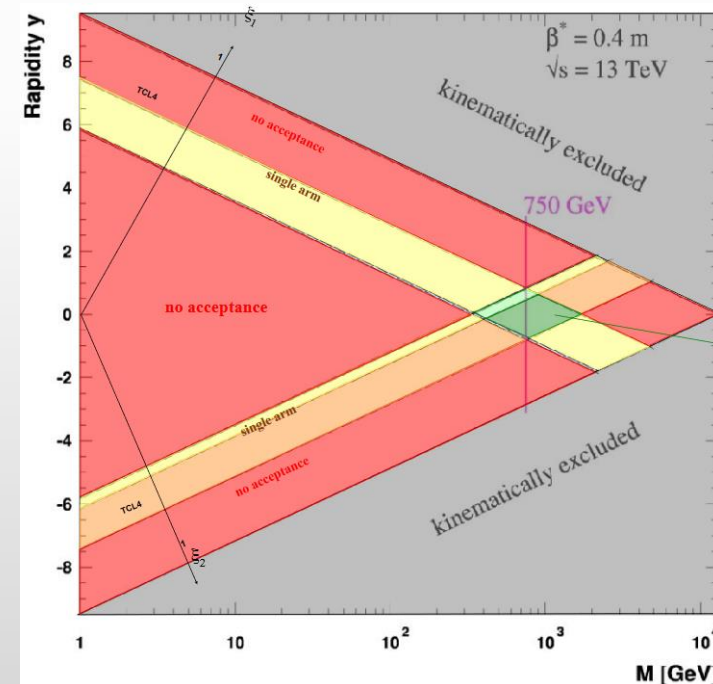
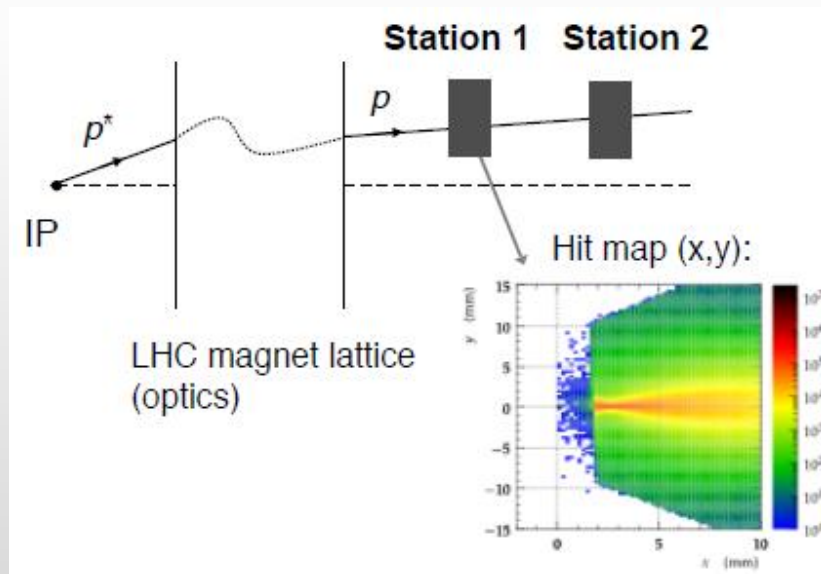
- **Double arm mass acceptance in the $\sim 400 - 2000$ GeV range**

- Lower limit mainly due to the minimum distance from the beam (can vary depending on beam conditions, e.g. crossing angle)
- Upper limit due to collimators

$$\xi = 1 - \frac{|\mathbf{p}_f|}{|\mathbf{p}_i|}$$

$$M_X = \sqrt{s \xi_1 \xi_2}$$

$$y_X = \frac{1}{2} \log\left(\frac{\xi_1}{\xi_2}\right)$$



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(Semi)Exclusive $\gamma\gamma \rightarrow l^+l^-$
2016 results

Prospects for LHC Run 3

(Semi)exclusive $\gamma\gamma \rightarrow \ell^+ \ell^-$

- Search for **opposite-charge lepton pairs produced by two photons, with forward proton tagging**

- See (JHEP 07 (2018) 153 (arXiv:1803.04496 [hep-ex]))

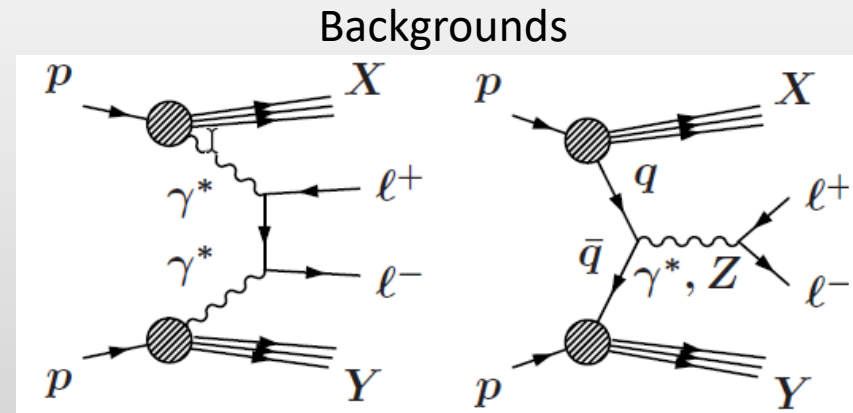
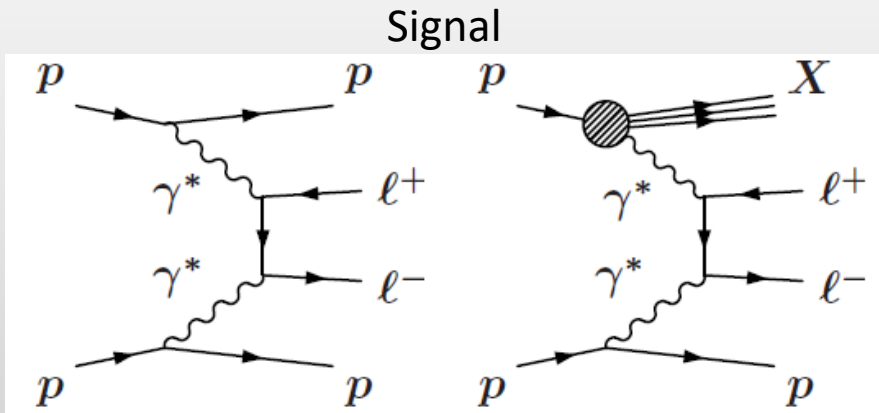
- Known EWK process:**

- Proton ξ related to \mathbf{p}_T and η of the leptons:** $\xi(\ell^+ \ell^-) = \frac{1}{\sqrt{s}} [\mathbf{p}_T(\ell^+) e^{\pm \eta(\ell^+)} + \mathbf{p}_T(\ell^-) e^{\pm \eta(\ell^-)}]$
 - Elastic contribution:** low theoretical uncertainty (E-M proton form factors, ...)
 - Single dissociation component:**
 - Wider photon virtuality spectrum than for exclusive production
 - Sensitive to rapidity gap survival probability
 - Provides acceptance towards lower masses

- Backgrounds:**

- Double dissociation** contribution
 - Inclusive** (Drell-Yan, VBF) contribution

} + pileup proton



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 $\gamma\gamma \rightarrow \ell^+ \ell^-$
2016 results

Prospects for
LHC Run 3

(Semi)Exclusive $\gamma\gamma \rightarrow \ell^+ \ell^-$ data selection

- **9.4 fb⁻¹ of 2016 data used**
- **Pre-selection:**
 - Trigger: ≥ 2 leptons with $p_T(\mu^\pm) > 38$ GeV and $p_T(e^\pm) > 33$ GeV
 - Offline selection: $p_T(\ell^\pm) > 50$ GeV, $m(\ell^+ \ell^-) > 110$ GeV (above Z mass peak), well-reconstructed protons with $\xi(\ell^+ \ell^-)$ in PPS coverage ≥ 1
 - Refitted dilepton vertex ($\chi^2 < 10$, $z < 15$ cm) separated from neighbouring tracks (0.5 mm veto)
 - Leptons produced back-to-back in transverse plane: $a \equiv 1 - \left| \frac{\Delta\Phi}{\pi} \right| < \begin{cases} 0.009 (\mu^+ \mu^-) \\ 0.006 (e^+ e^-) \end{cases}$
- Strong background suppression by **requiring 2σ matching between ξ measurements** (central tracker vs. PPS)
- **Data-driven background estimation**, using inclusive $DY \rightarrow \ell^+ \ell^-$ data and double dissociative simulated events together with randomly selected protons from the data passing the ξ match requirement
- **Expected background:**

$$\begin{cases} 1.49 \pm 0.07 \text{ (stat.)} \pm 0.53 \text{ (syst.)} (\mu^+ \mu^-) \\ 2.36 \pm 0.09 \text{ (stat.)} \pm 0.47 \text{ (syst.)} (e^+ e^-) \end{cases}$$



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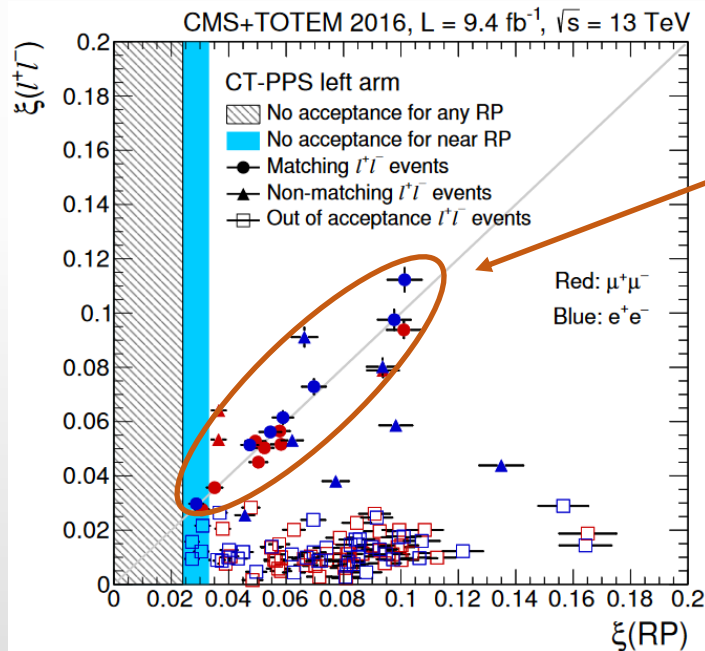
Detector alignment and proton reconstruction

(Semi)Exclusive $\gamma\gamma \rightarrow \ell^+ \ell^-$ 2016 results

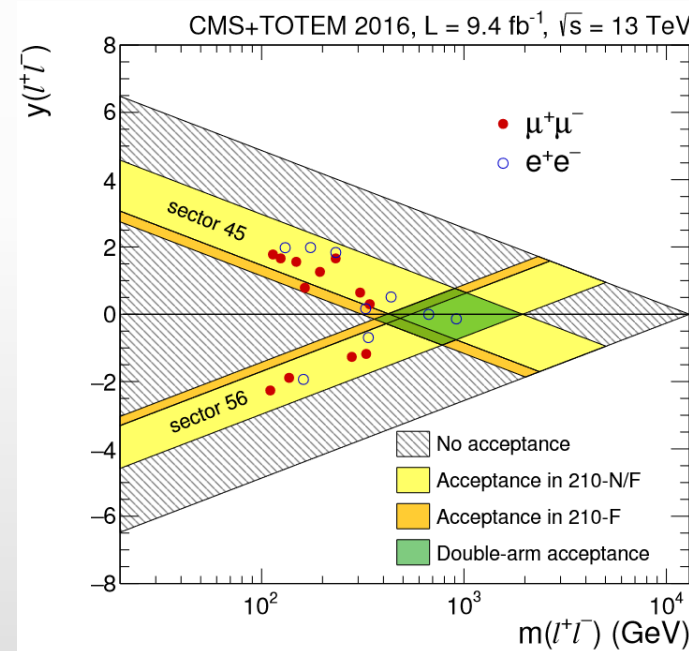
Prospects for LHC Run 3

(Semi)Exclusive $\gamma\gamma \rightarrow \ell^+ \ell^-$ 2016 results

- First observation of (semi)exclusive (two-photon) production of dileptons with tagged protons
 - 20 matching events, with a total estimated background of 3.85 events
 - Combined significance of 5.1σ over background only hypothesis



Events matched
in Sector 45
(Left arm)



This measurement proves that PPS performs according to expectations



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Prospects for
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PPS analysis prospects

- **Search for two-photon production of a boson pair:**

- Study neutral quartic gauge coupling, suppressed in SM
- Sensitive to resonances (axion-like particles, new particle exchanges)
- Provides model-independent bounds on massive charged particles, only parameterized by spin, mass and "effective charge"

- See *arXiv: 1411.6629*

- **Search for anomalous quartic gauge couplings:**

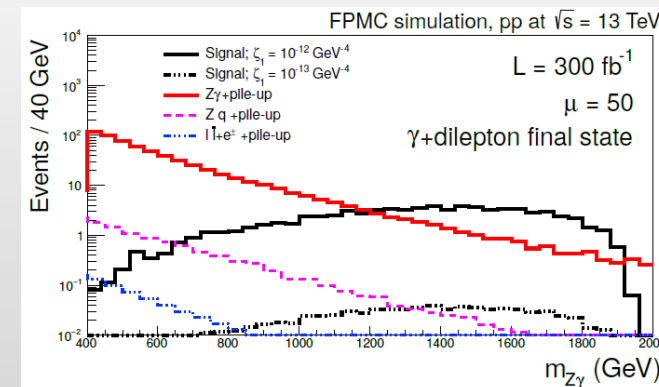
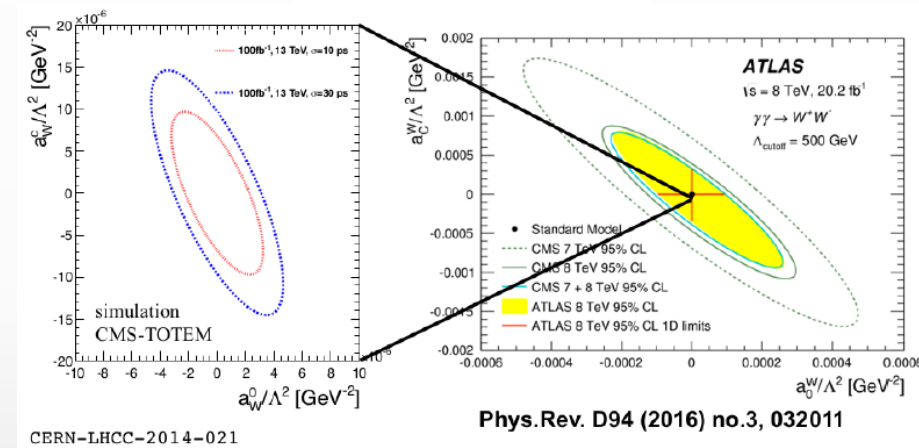
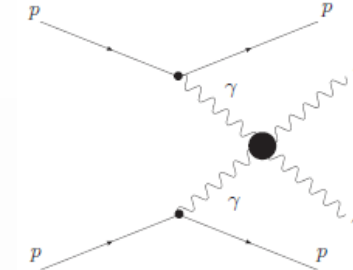
$$\gamma\gamma \rightarrow W^+W^-, \gamma\gamma \rightarrow \gamma Z, \gamma\gamma \rightarrow ZZ...$$

- $\gamma\gamma \rightarrow W^+W^-$: PPS TDR expectations with two order of magnitude improvement w.r.t. Run 1 searches

- See *arXiv: 1604.04464*, *arXiv: 1607.03745*

- $\gamma\gamma \rightarrow \gamma Z$: the combined sensitivity in the γZ channel, for 300 fb⁻¹, goes beyond the one expected for $Z \rightarrow \gamma\gamma\gamma$ decay searches by ~3 orders of magnitude

- See *arXiv: 1703.10600*



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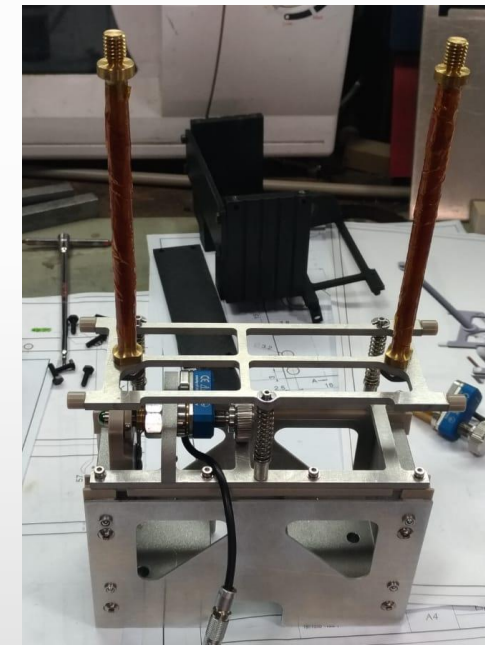
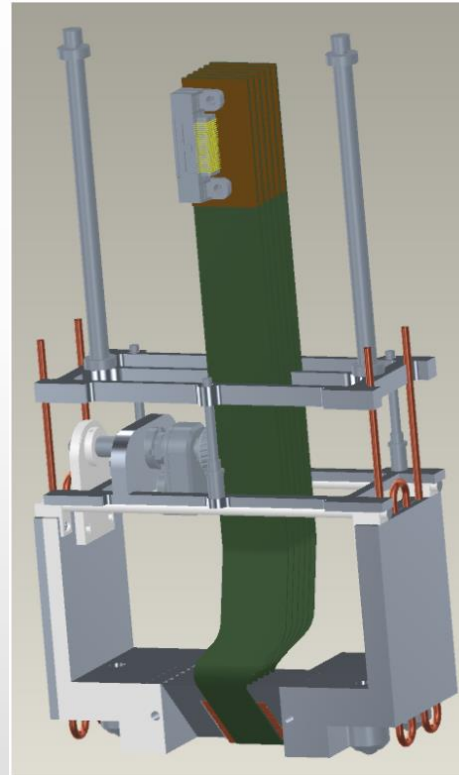
PPS Prospects

PPS prospects for Run 3

- PPS approved for operation as standard CMS subsystem in LHC Run 3 (2021-2023)

- **Experimental apparatus:**

- Tracking:
 - 2 horizontal stations with silicon 3D pixels sensors (similar to the 3D pixels designed for the CMS Phase2 Tracker R&D)
 - PROC600 readout chip (same as CMS pixel detector layer 1)
 - New internally motorized detector package, to distribute the radiation damage and reduce its impact
- Timing:
 - 2 horizontal stations equipped with double-layer diamond sensors
 - Optimized readout electronics



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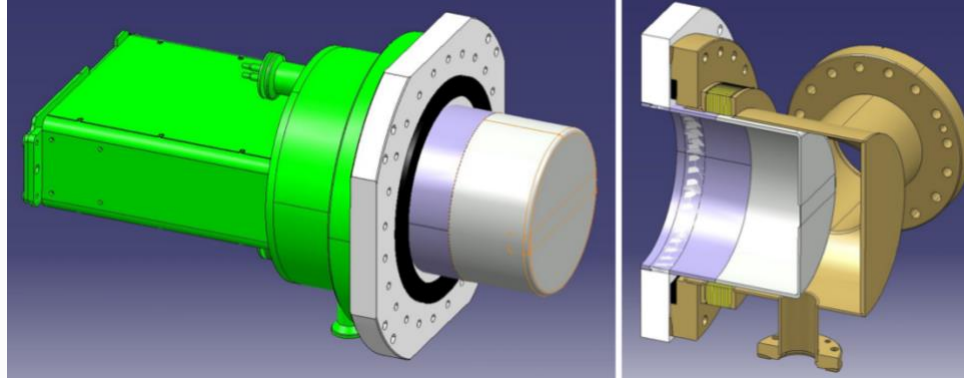
Summary



- PPS has proven the feasibility of continuously operating a near-beam proton spectrometer at a high-luminosity hadron collider
- **PPS has successfully collected $\sim 115 \text{ fb}^{-1}$ worth of data** during LHC Run 2, with very good overall performance
- **PPS has for the first time observed at more than 5σ significance the proton-tagged two-photon production of lepton pairs** with $\sim 10 \text{ fb}^{-1}$ of 2016 data
- 2017 and 2018 detector performance studies are being finalized
- **The preparation for LHC Run 3 is ongoing:**
 - **New detectors are getting ready** to be installed for the future data taking
 - **A rich physics programme** lies ahead, with many final states to be studied

Backup Slides

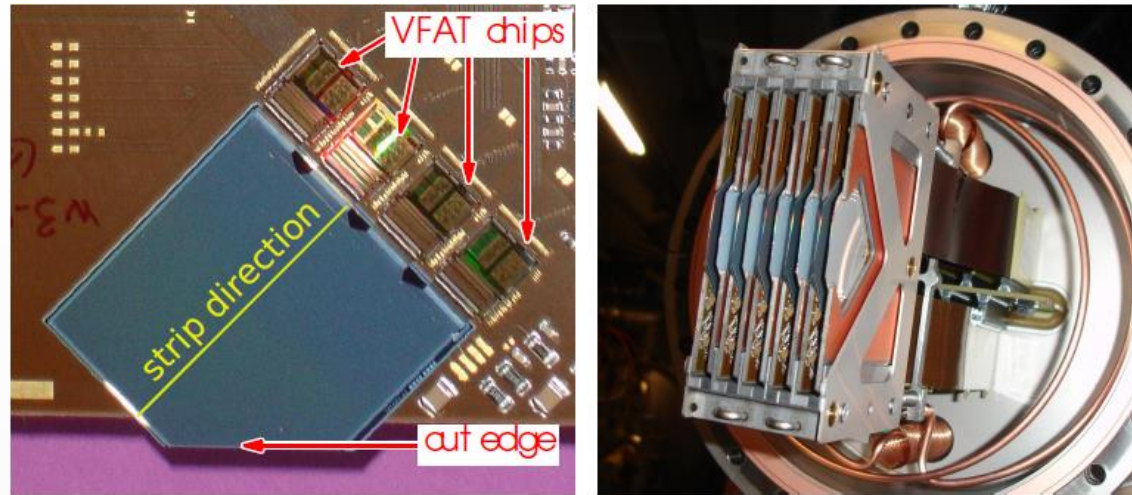
RPs for timing detectors



- **New cylindrical RP design to host larger detectors and reduce the impedance**
- Timing RPs are **equipped with a 300 μm thick window towards the beam**
 - The thickness is required to compensate the pressure gradient on the larger window
 - No vertical stations needed because the alignment is done by propagating tracks from the tracking stations



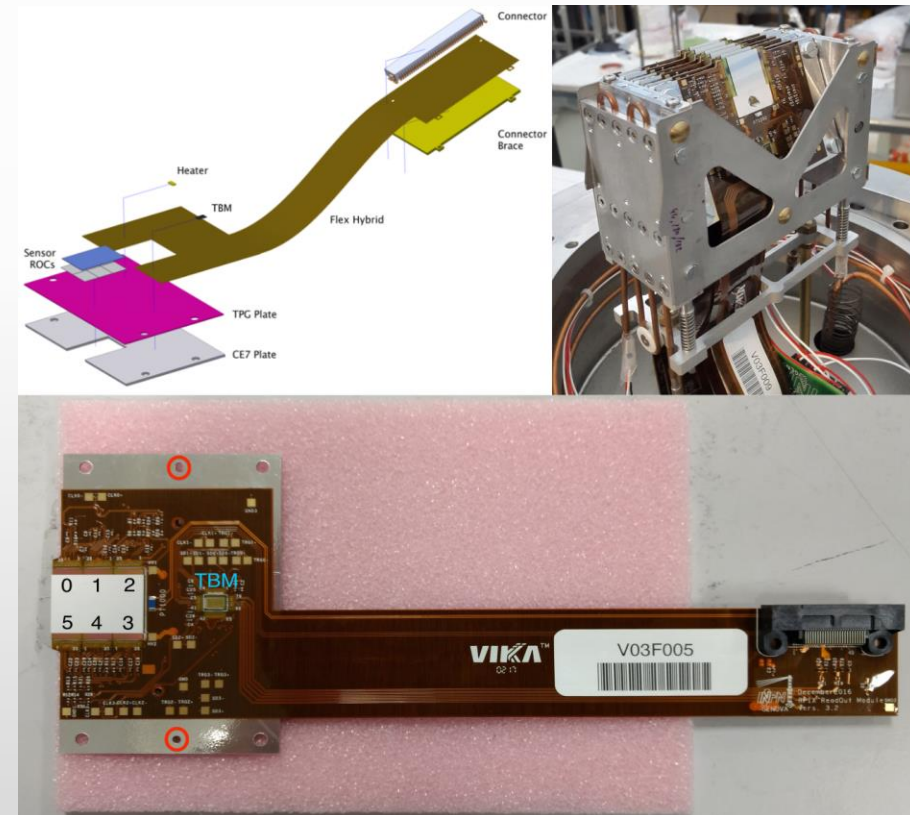
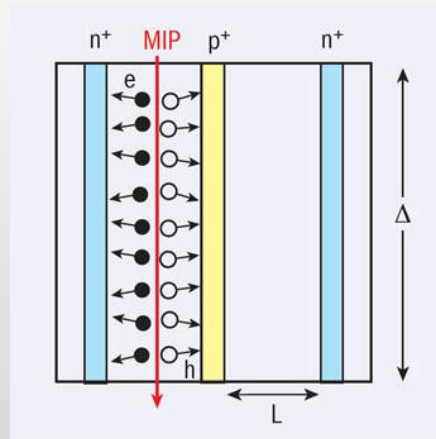
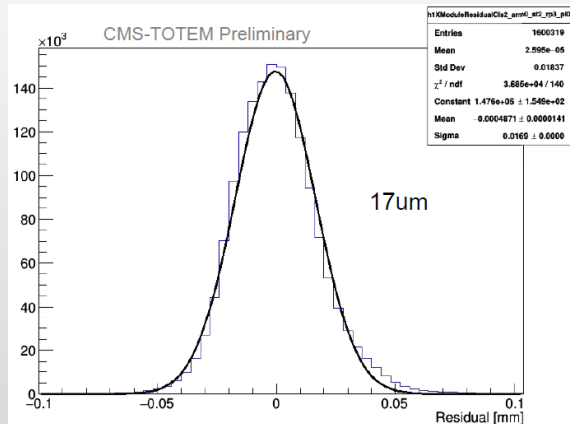
Silicon strips detectors



- 300 μm thickness
- 66 μm pitch
- strips oriented at 45° w.r.t. edge facing beam
- Dead area close to the cut edge: only $\sim 50 \mu\text{m}$
- operated at -20°C , bias voltage $\sim 100 \text{ V}$
- 5+5 planes per RP (2 strip orientations for 2D reconstruction)

3D pixel detectors

- 3D technology grants high radiation hardness (up to $\sim 5 \cdot 10^{15}$ p/cm²)
- Pixel size 100×150 μm , 230 μm thickness
- Low dead region close to the edge (~ 50 μm)
- 6 planes per RP
 - 18° planes tilt to improve hit cluster size
- Operation at -20 °C and in vacuum ($P < 20$ mbar)
- Pixel tracker works as expected
 - Residuals according with test beam measurements



Timing detectors

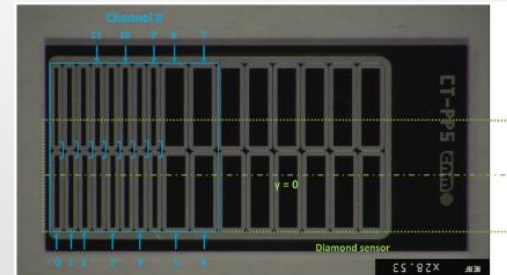
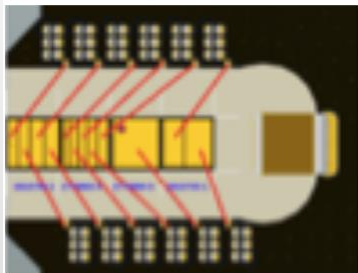
Diamond Sensors

- CVD diamond sensors
- Radiation hard
- Macro-pixels of varying size
- Single-plane resolution in test-beam:
 $\sigma_t \sim 80$ ps

Ultra-Fast Silicon Detectors

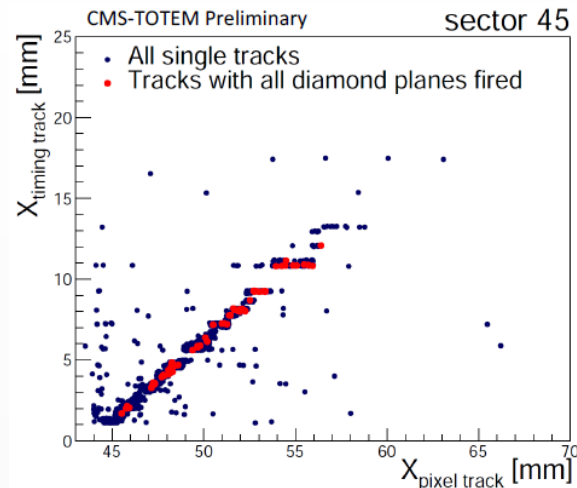
- Based on LGAD technology
- Limited radiation hardness
- Macro-pixels of varying size
- Single-plane resolution in test-beam:
 $\sigma_t \sim 30$ ps

Common readout electronics



Timing detectors performance

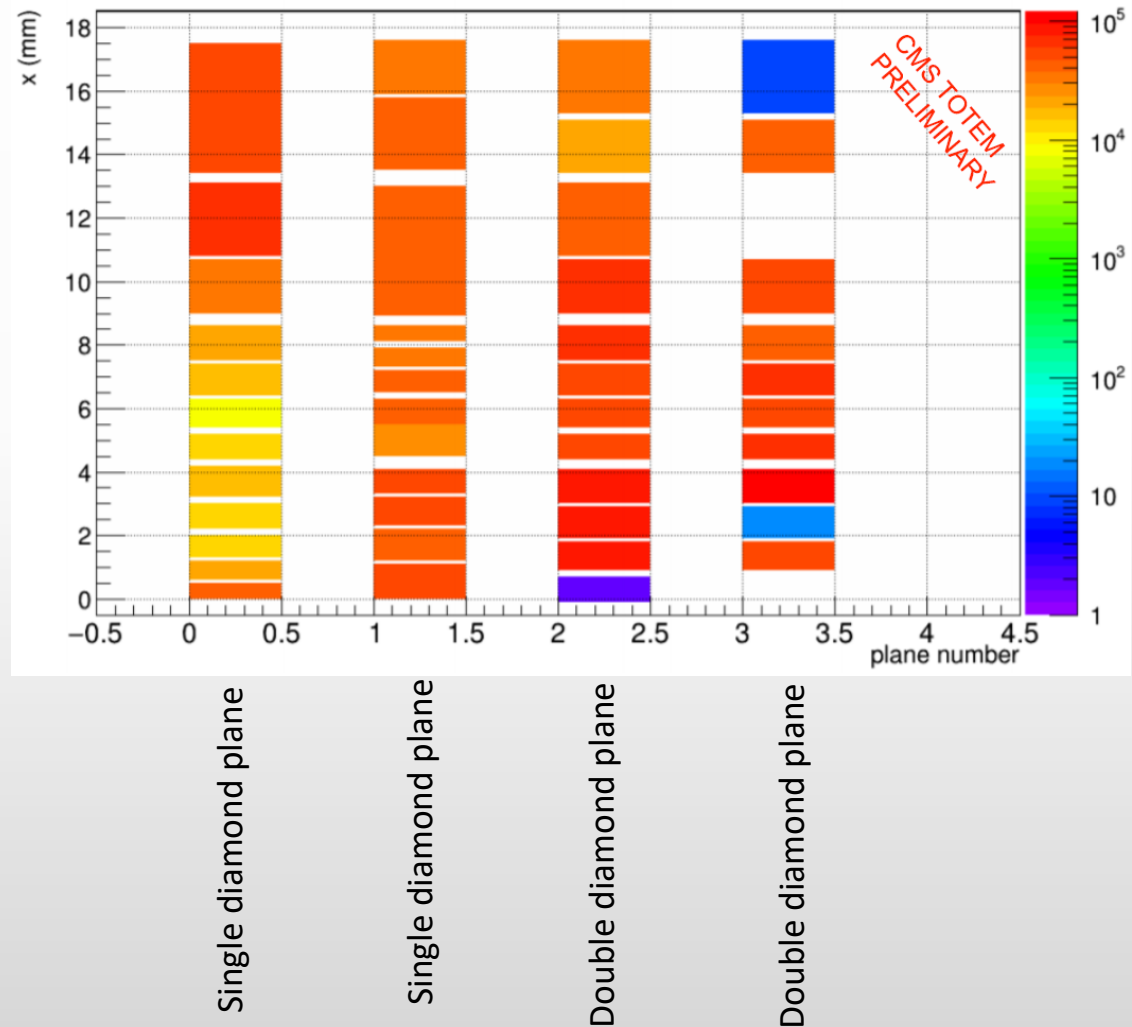
2017



- Horizontal position of the reconstructed track in the PPS pixels vs. horizontal position of in the PPS timing in low pile-up data ($\langle \text{PU} \rangle \sim 0.8$)
- The data sample is selected requiring a single vertex reconstructed in CMS, a single track reconstructed in the PPS pixels in each of the arms and a single track reconstructed in the PPS timing in each of the arms
- **Blue points** show all the events passing the double arm selection
- **Red points** represent the subsample with a single hit per diamond detector plane and all the diamond planes firing

2018

CT-PPS Timing SEC56 220C



Proton transport model

- The propagation of the proton from the IP to the RPs is approximately modeled with this linear formula:

$$\mathbf{d}(s) = T(s) \cdot \mathbf{d}^*$$

where $\mathbf{d} = \left(x, \theta_x, y, \theta_y, \frac{\Delta p}{p}\right)^T$, the * symbol refers to quantities at the IP,

$$T = \begin{pmatrix} v_x & L_x & m_{13} & m_{14} & D_x \\ \frac{dv_x}{ds} & \frac{dL_x}{ds} & m_{23} & m_{24} & \frac{dD_x}{ds} \\ m_{31} & m_{32} & v_y & L_y & D_y \\ m_{41} & m_{42} & \frac{dv_y}{ds} & \frac{dL_y}{ds} & \frac{dD_y}{ds} \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix}, \begin{matrix} v_{x,y} = \sqrt{\beta_{x,y}/\beta^*} \cos(\Delta\mu_{x,y}) \\ L_{x,y} = \sqrt{\beta_{x,y} \cdot \beta^*} \sin(\Delta\mu_{x,y}) \end{matrix}$$

β is the betatron amplitude and $\Delta\mu_{x,y}$ is the relative phase advance ($\Delta\mu_{x,y} = \int_{IP}^{RP} \frac{ds}{\beta_{x,y}}$)

The leading terms are:

- $x \approx D_x(\xi) \cdot \xi$
- $y \approx L_y(\xi) \cdot \theta_y^*$

