

Institut de recherche en mathématique et physique Centre de Cosmologie, Physique des Particules et Phénoménologie







Plan

Yesterday

- BSM Model
- Matrix-element and decay

Today

- Phase-Space integration at LO
- Merging at LO (MLM and CKKW(L))

Friday

aMC@NLO versus NLO computation

Phase-Space

Olivier Mattelaer CP3/UCLouvain

Matrix-Element



Matrix-Element



Matrix-Element



Example: QCD 2 \rightarrow 2



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Three very different pole structures contributing to the same matrix element.

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ToolBox of yesterday: Multi-Channel

*Method used in MadGraph

Trick in MadEvent: Split the complexity

$$\int |M_{tot}|^2 = \int \frac{\sum_i |M_i|^2}{\sum_j |M_j|^2} |M_{tot}|^2 = \sum_i \int \frac{|M_i|^2}{\sum_j |M_j|^2} |M_{tot}|^2$$

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Key Idea

- Any single diagram is "easy" to integrate (pole structures/ suitable integration variables known from the propagators)
- Divide integration into pieces, based on diagrams
- All other peaks taken care of by denominator sum

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N Integral

- Errors add in quadrature so no extra cost
- "Weight" functions already calculated during $|\mathcal{M}|^2$ calculation
- Parallel in nature

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P1 qq wpwm

 $s = 725.73 \pm 2.07 \text{ (pb)}$

<u>Graph</u>	Cross-Section ↓	<u>Error</u>	<u>Events (K)</u>	<u>Unwgt</u>	<u>Luminosity</u>
G2.2	<u>377.6</u>	1.67	142.285	7941.0	21
G3	<u>239</u>	1.16	220.04	10856.0	45.5
G 1	<u>109.1</u>	0.378	70.88	3793.0	34.8

P1 gg wpwm

s= 20.714 ± 0.332 (pb)

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G1.2	<u>20.71</u>	0.332	7.01	373.0	18

term of the above sum.

each term might not be gauge invariant

Cut Impact

- Events are generated according to our best knowledge of the function
 - →Basic cut include in this "best knowledge"
 - →Custom cut are ignored



Cut Impact



Might miss the contribution and think it is just zero.

Matching/Merging

Olivier Mattelaer CP3/UCLouvain

PS alone vs matched samples

Parton Shower are using collinear approximation how accurate does it works





- I. Fixed order calculation
- 2. Computationally expensive
- 3. Limited number of particles
- 4. Valid when partons are hard and well separated
- 5. Quantum interference correct
- 6. Needed for multi-jet description



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Approaches are complementary: merge them!



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Difficulty: avoid double counting, ensure smooth distributions



2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia

• Regularization of matrix element divergence



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- Regularization of matrix element divergence
- Correction of the parton shower for large momenta



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Double counting between ME and PS easily avoided using phase space cut between the two: PS below cutoff, ME above cutoff.

- So double counting problem easily solved, but what about getting smooth distributions that are independent of the precise value of Q^c?
- Below cutoff, distribution is given by PS
 need to make ME look like PS near cutoff
- Let's take another look at the PS!





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Corresponds to the matrix element BUT with α_s evaluated at the scale of each splitting



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Sudakov suppression due to not allowing additional radiation above the scale t_{cut}

 $|\mathcal{M}|^2(\hat{s}, p_3, p_4, \ldots)$

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$$\mathcal{M}|^2 \to |\mathcal{M}|^2 \frac{\alpha_s(t_1)}{\alpha_s(t_0)} \frac{\alpha_s(t_2)}{\alpha_s(t_0)}$$

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 - 3. Use some algorithm to apply the equivalent Sudakov suppression $(\Delta_q(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{\text{cut}}, t_2))^2$



We are of course not interested in e⁺e⁻ but p-p(bar)
 what happens for initial state radiation?



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 what happens for initial state radiation?
- Let's do the same exercise as before:

$$\mathcal{P} = (\Delta_{Iq}(t_{cut}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{cut}, t_2))^2 \frac{\alpha_s(t_1)}{2\pi} \frac{P_{gq}(z)}{z} \frac{f_q(x_1, t_1)}{f_q(x_1', t_1)} \frac{\alpha_s(t_2)}{2\pi} P_{qg}(z') \\ \times \hat{\sigma}_{q\bar{q} \to e\nu}(\hat{s}, ...) f_q(x_1', t_0) f_{\bar{q}}(x_2, t_0) \\ t_{cut_1} t_1 t_1 t_2 t_{cut_1} t_{$$

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Mattelaer Olívíer

Monte-Carlo Lecture: 2019

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ME (convoluted with PDF) with α_s evaluated at the scale of each splitting PDF reweighting



 $(\Delta_{Iq}(t_{\rm cut},t_0))^2 \Delta_g(t_2,t_1) (\Delta_q(t_{\rm cut},t_2))^2 \frac{\alpha_s(t_1)}{2\pi} \frac{P_{gq}(z)}{z} \frac{f_q(x_1,t_1)}{f_q(x_1',t_1)} \frac{\alpha_s(t_2)}{2\pi} P_{qg}(z')$

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Sudakov suppression due to non-branching above scale t_{cut}





• Again, use a clustering scheme to get a parton shower history



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- Now, reweight both due to α_s and PDF

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• Use some algorithm to apply the equivalent Sudakov suppression



Matching schemes

- We still haven't specified how to apply the Sudakov reweighting to the matrix element
- Three general schemes available in the literature:
 - CKKW scheme [Catani,Krauss,Kuhn,Webber 2001; Krauss 2002]
 - Lönnblad scheme (or CKKW-L) [Lönnblad 2002]
 - MLM scheme [Mangano unpublished 2002; Mangano et al. 2007]
[Catani, Krauss, Kuhn, Webber 2001] [Krauss 2002]

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• Apply the required Sudakov suppression

 $(\Delta_{Iq}(t_{\rm cut},t_0))^2 \Delta_g(t_2,t_1) (\Delta_q(t_{\rm cut},t_2))^2$

analytically, using the best available (NLL) Sudakovs.

[Catani, Krauss, Kuhn, Webber 2001] [Krauss 2002]

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• Perform "truncated showering": Run the parton shower starting at t_0 , but forbid any showers above the cutoff scale t_{cut} .



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- Perform "truncated showering": Run the parton shower starting at t_0 , but forbid any showers above the cutoff scale t_{cut} .
 - ✓ Best theoretical treatment of matrix element
 - Requires dedicated PS implementation
 - Mismatch between analytical Sudakov and (non-NLL) shower
 - Implemented in Sherpa (v. I.I)



[Lönnblad 2002]

[Hoeche et al. 2009]



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- \checkmark Automatic agreement between Sudakov and shower
- Requires dedicated PS implementation
 - Need multiple implementations to compare between showers
- Implemented in Adriane, Sherpa (v. 1.2), and Pythia 8

[M.L. Mangano, ~2002, 2007] [J.A. et al 2007, 2008]



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- Perform jet clustéring after PS if hardest jet $k_{TI} > t_{cut}$ or there are jets not matched to partons, reject the event
- The resulting Sudakov suppression from the procedure is $(\Delta_{Iq}(t_{\rm cut}, t_0))^2 (\Delta_q(t_{\rm cut}, t_0))^2$

which turns out to be a good enough approximation of the correct expression $(\Delta_{Iq}(t_{cut}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{cut}, t_2))^2$

Mattelaer Olívíer

[M.L. Mangano, ~2002, 2007] [J.A. et al 2007, 2008]

• The simplest way to do the Sudakov suppression is to run the shower on the event, starting from t₀!



- \checkmark Simplest available scheme
- ✓ Allows matching with any shower, without modification
- Sudakov suppression not exact, minor mismatch with shower
- Implemented in AlpGen, HELAC, MadGraph+Pythia 6

Highest multiplicity sample

- In the previous, assumed we can simulate all parton multiplicities by the ME
- In practice, we can only do limited number of final-state partons with matrix element (up to 4-5 or so)
- For the highest jet multiplicity that we generate with the matrix element, we need to allow additional jets above the matching scale *t*_{cut}, since we will otherwise not get a jet-inclusive description but still can't allow PS radiation harder than the ME partons
- Need to replace t_{cut} by the clustering scale for the softest ME parton for the highest multiplicity

matching schemes

- We have a number of choices to make in the above procedure. The most important are:
 - I. The clustering scheme used to determine the parton shower history of the ME event
 - 2. What to use for the scale of hard emission
 - 3. How to divide the phase space between parton showers and matrix elements

Cluster schemes

1. The clustering scheme used inside MadGraph and Sherpa to determine the parton shower history is the Durham k_T scheme. For e⁺e⁻:

$$k_{Tij}^2 = 2\min(E_i^2, E_j^2)(1 - \cos\theta_{ij})$$

and for hadron collisions, the minimum of: $k_{Tibeam} = m_i^2 + p_{Ti}^2 = (E_i + p_{zi})(E_i - p_{zi})$ and $k_{Tij}^2 = \min(p_{Ti}^2, p_{Tj}^2)R_{ij}$ with $R_{ij} = 2[\cosh(y_i - y_j) - \cos(\phi_i - \phi_j)] \simeq (\Delta y)^2 + (\Delta \phi)^2$

Find the smallest k_{Tij} (or k_{Tibeam}), combine partons *i* and *j* (or *i* and the beam), and continue until you reach a $2 \rightarrow 2$ (or $2 \rightarrow 1$) scattering.

2. In AlpGen a more naive cone algorithm is used.

MLM in MadGraph

- Cannot use the standard k_{T} clustering:
 - MadGraph and Sherpa only allow clustering according to valid diagrams in the process. This means that, e.g., two quarks or quark-antiquark of different flavor are never clustered, and the clustering always gives a physically allowed parton shower history.
 - If there is an on-shell propagator in the diagram (e.g. a top quark), only clustering according to diagrams with this propagator is allowed.

Hard scale

2. The clustering provides a convenient choice for factorization scale Q²:



Cluster back to the 2 \rightarrow 2 (here qq \rightarrow W-g) system, and use the W boson transverse mass in that system.

Phase-space division

3. How to divide the phase space between PS and ME: This is where the schemes really differ:

AlpGen: MLM Cone MadGraph: MLM Cone, k_T or shower-k_T Sherpa: CKKW

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- Regularization of matrix element divergence
- Correction of the parton shower for large momenta
- Smooth jet distributions



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2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia

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2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia In the soft-collinear approximation of Parton Shower MCs, parameters are used to tune the result \Rightarrow Large variation in results (small prediction power)



PS alone vs ME matching

In a matched sample these differences are irrelevant since the behavior at high pt is dominated by the matrix element.



Summary of Matching Procedure

- I. Generate ME events (with different parton multiplicities) using parton-level cuts ($p_T^{ME}/\Delta R$ or k_T^{ME})
- 2. Cluster each event and reweight α_s and PDFs based on the scales in the clustering vertices
- 3. Apply Sudakov factors to account for the required nonradiation above clustering cutoff scale and generate parton shower emissions below clustering cutoff:
 - a. (CKKW) Analytical Sudakovs + truncated showers
 - b. (CKKW-L) Sudakovs from truncated showers
 - c. (MLM) Sudakovs from reclustered shower emissions