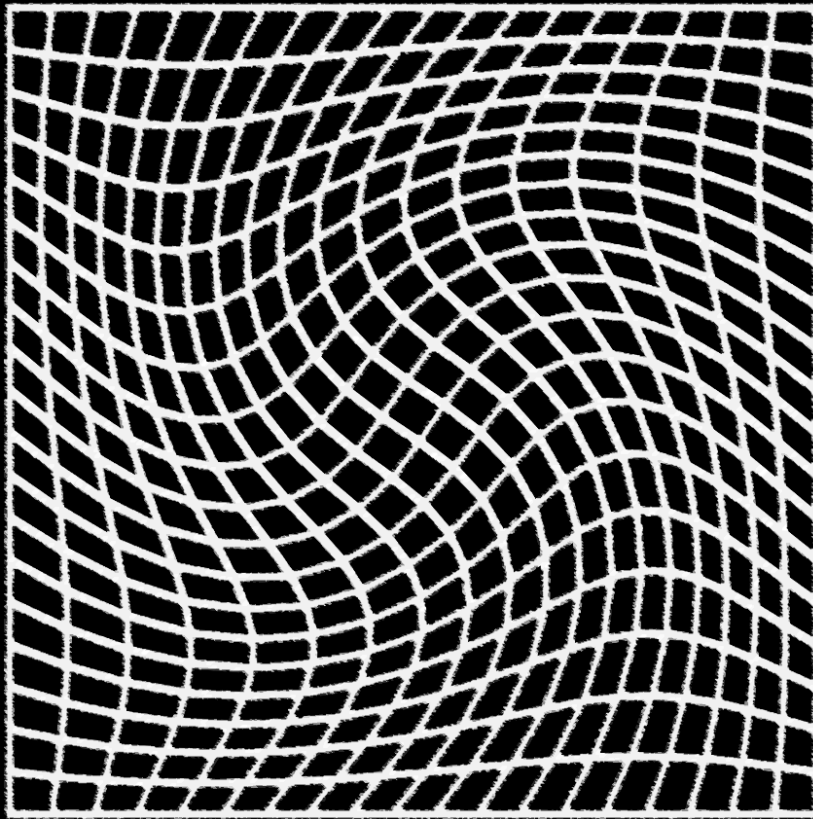


Geometry



of parametric
quantum
systems

Pavel Cejnar

@ipnp.troja.mff.cuni.cz

Closed bound quantum system

Hamiltonian

$$\hat{H}$$

Eigenvalues

$$E_i$$

Eigenvectors

$$|i\rangle$$

Closed bound quantum system with variable

$\lambda \equiv (\lambda_1, \lambda_2, \dots, \lambda_n) \in \mathbb{R}^n$ parameters

Hamiltonian

$$\hat{H}(\lambda)$$

Eigenvalues

$$E_i(\lambda)$$

Eigenvectors

$$|i(\lambda)\rangle$$

Closed bound quantum system with variable parameters

Hamiltonian

$$\lambda \equiv (\lambda_1, \lambda_2, \dots, \lambda_n) \in \mathbb{R}^n$$

parameters

$$\hat{H}(\lambda)$$

“controllable quantum systems”

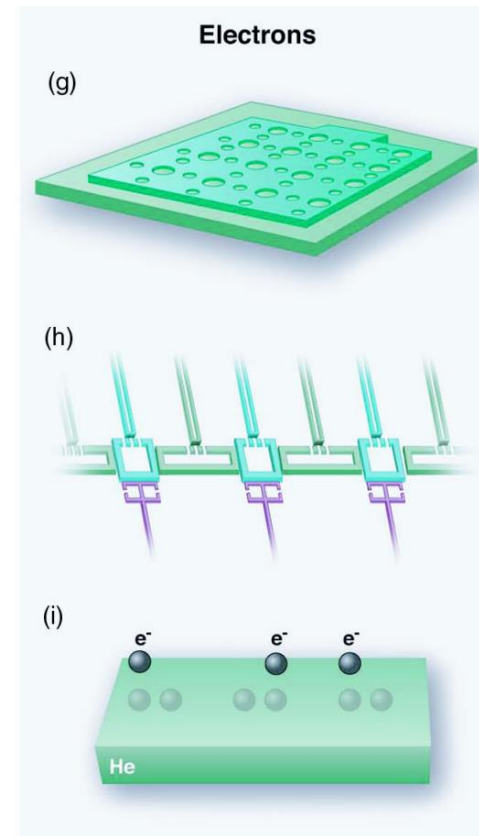
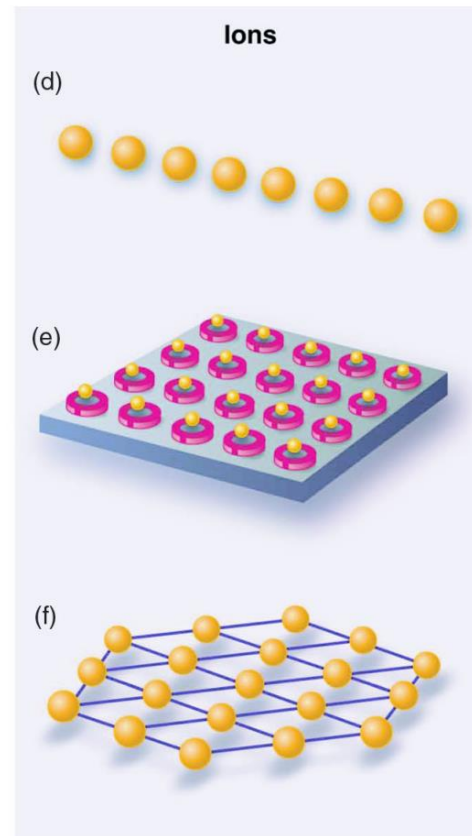
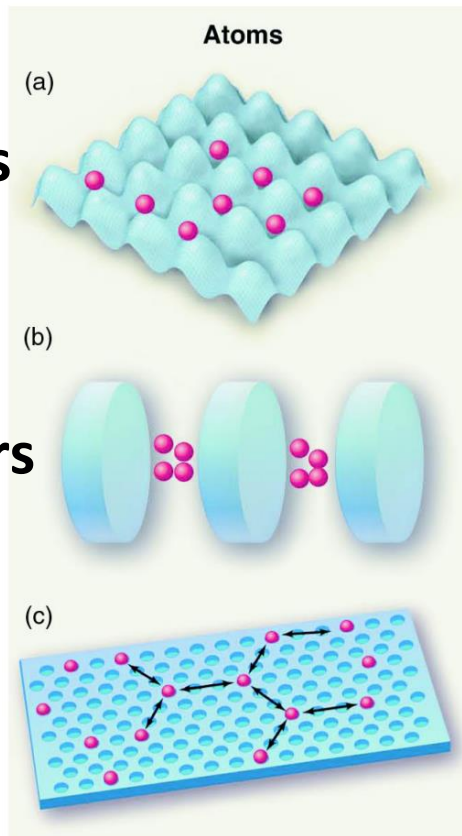
“quantum simulators”

Eigenvalues

$$E_i(\lambda)$$

Eigenvectors

$$|i(\lambda)\rangle$$



Closed bound quantum system with variable parameters

Hamiltonian

$$\hat{H}(\boldsymbol{\lambda})$$

Eigenvalues

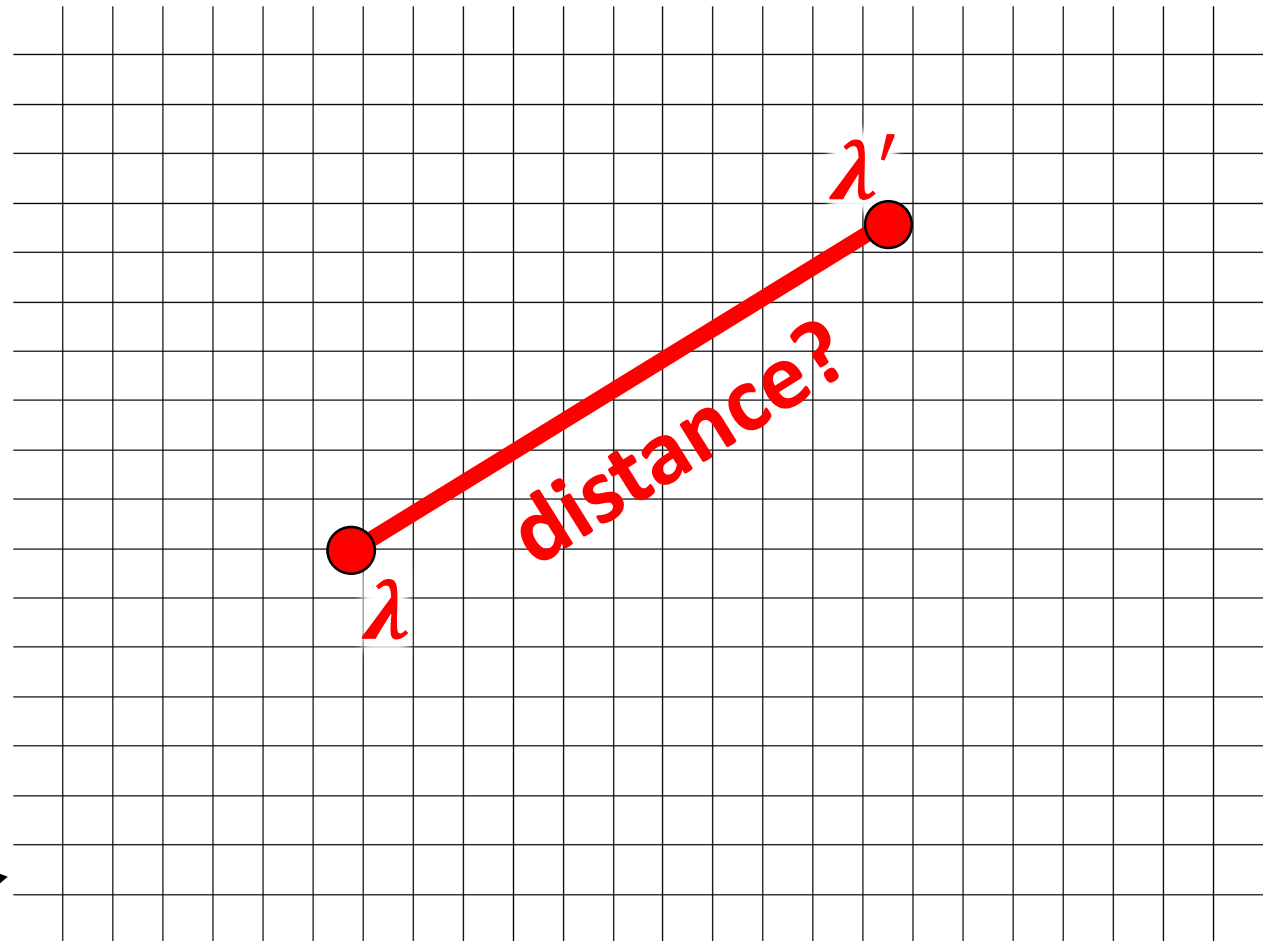
$$E_i(\boldsymbol{\lambda})$$

Eigenvectors

$$|i(\boldsymbol{\lambda})\rangle$$

$$\boldsymbol{\lambda} \equiv (\lambda_1, \lambda_2, \dots, \lambda_n) \in \mathbb{R}^n$$

parameters



Eigenvector distance: Hilbert space distance

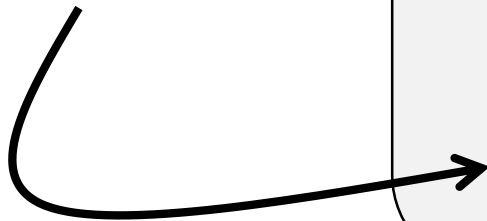
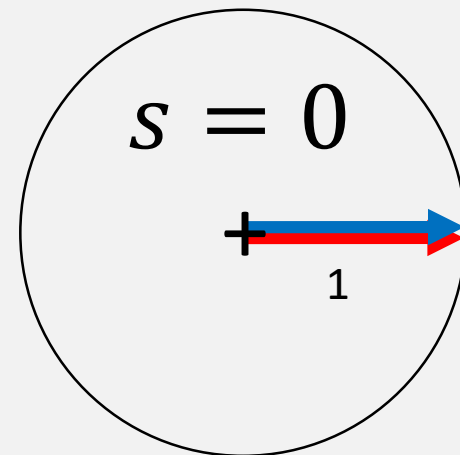
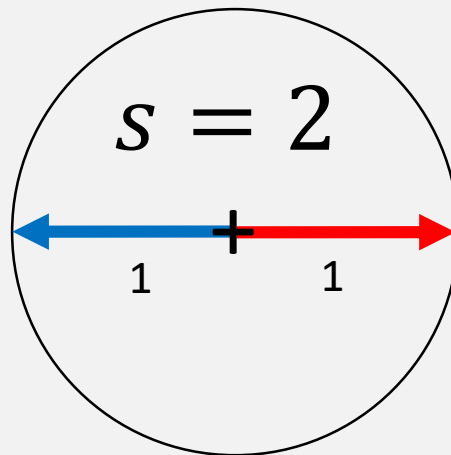
$$\lambda \equiv (\lambda_1, \lambda_2, \dots, \lambda_n) \in \mathbb{R}^n$$

$$(s^{(i)})^2 = \langle i(\lambda') - i(\lambda) | i(\lambda') - i(\lambda) \rangle$$

gauge dependent!

$$= 2 - 2\text{Re}\langle i(\lambda') | i(\lambda) \rangle$$

$|i(\lambda)\rangle$



Eigenvector distance: fidelity based distance

$$\boldsymbol{\lambda} \equiv (\lambda_1, \lambda_2, \dots, \lambda_n) \in \mathbb{R}^n$$

$$(ds^{(i)})^2 = 1 - \overbrace{|\langle i(\boldsymbol{\lambda} + d\boldsymbol{\lambda}) | i(\boldsymbol{\lambda}) \rangle|^2}^{\text{fidelity}}$$

gauge independent!

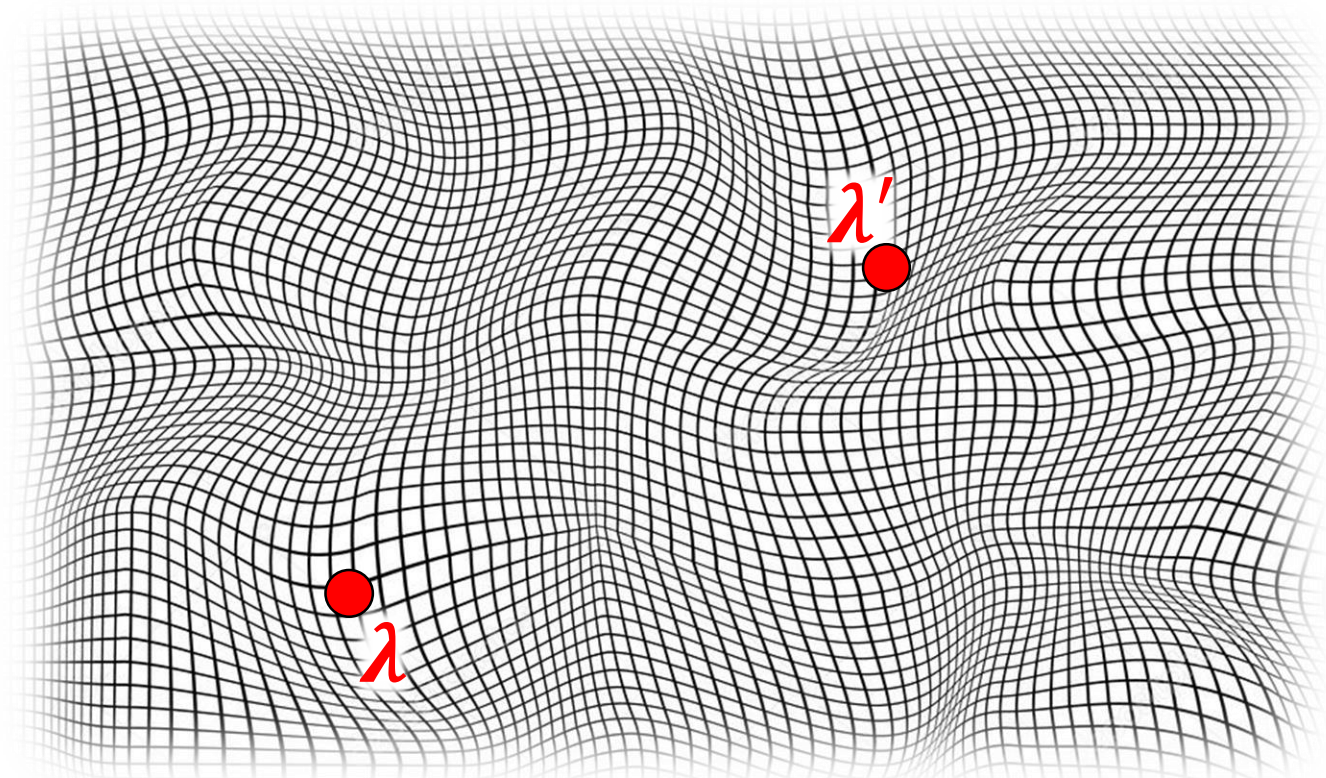
$$g_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) = \frac{1}{2} \left(\left\langle \frac{\partial}{\partial \lambda_\alpha} i(\boldsymbol{\lambda}) \middle| \frac{\partial}{\partial \lambda_\beta} i(\boldsymbol{\lambda}) \right\rangle - \left\langle \frac{\partial}{\partial \lambda_\alpha} i(\boldsymbol{\lambda}) \middle| i(\boldsymbol{\lambda}) \right\rangle \left\langle i(\boldsymbol{\lambda}) \middle| \frac{\partial}{\partial \lambda_\beta} i(\boldsymbol{\lambda}) \right\rangle + \text{c. c.} \right)$$

$|i(\boldsymbol{\lambda})\rangle$


$$(ds^{(i)})^2 = g_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) d\lambda_\alpha d\lambda_\beta$$

Eigenvector parameter manifold as a curved space

$$\boldsymbol{\lambda} \equiv (\lambda_1, \lambda_2, \dots, \lambda_n) \in \mathbb{R}^n$$



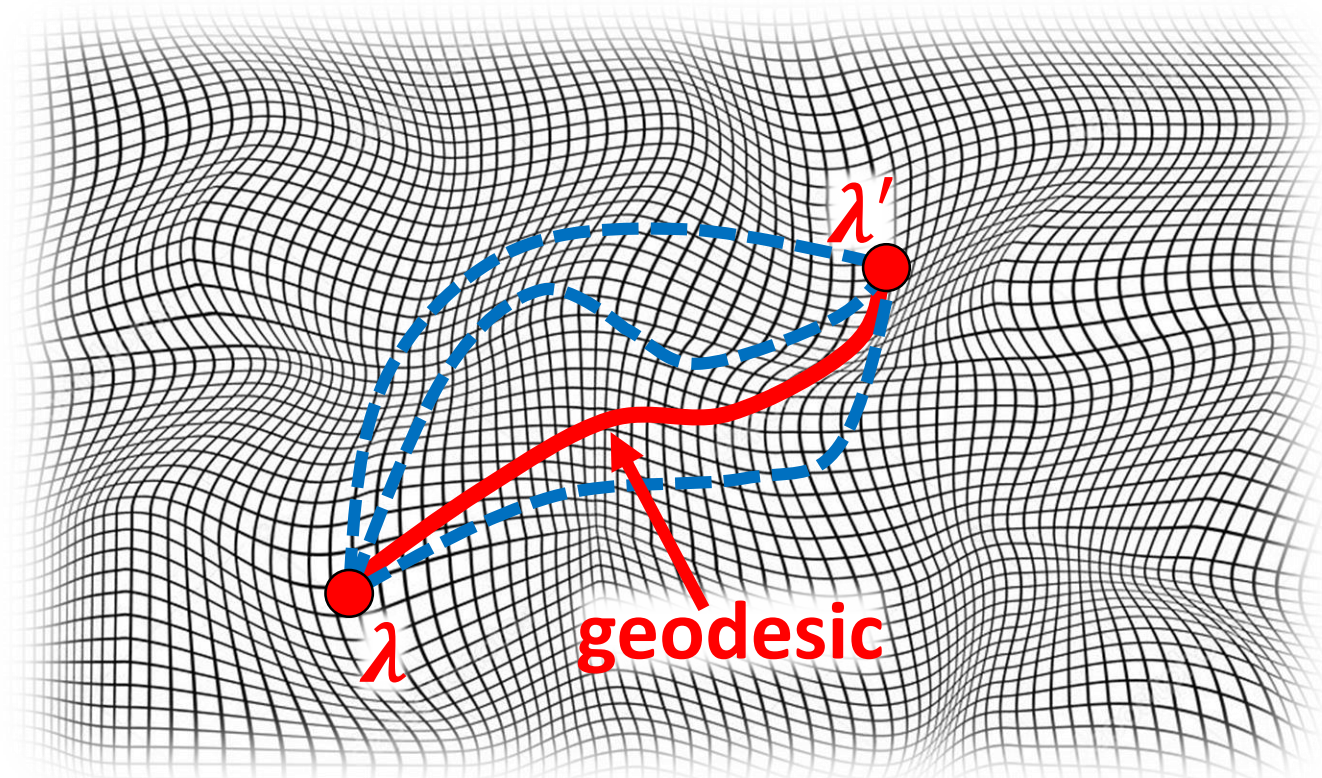
$|i(\boldsymbol{\lambda})\rangle$



$$(ds^{(i)})^2 = g_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) d\lambda_\alpha d\lambda_\beta$$

Eigenvector parameter manifold as a curved space

$$\boldsymbol{\lambda} \equiv (\lambda_1, \lambda_2, \dots, \lambda_n) \in \mathbb{R}^n$$



$|i(\boldsymbol{\lambda})\rangle$



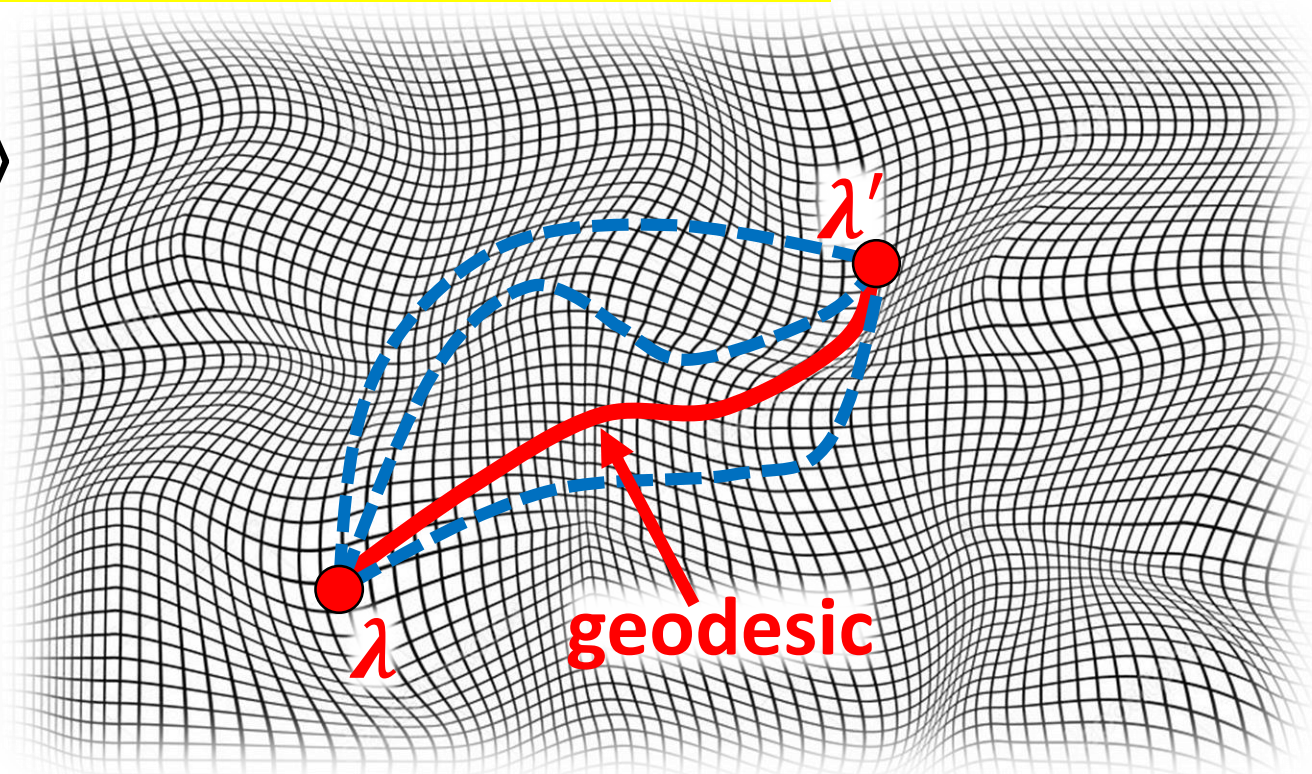
$$(ds^{(i)})^2 = g_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) d\lambda_\alpha d\lambda_\beta$$

Driven dynamics: adiabatic driving

$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = \hat{H}(\lambda(t)) |\psi(t)\rangle$$

$$|\psi(0)\rangle = |i(\lambda)\rangle$$

$$|\psi(T)\rangle = |i(\lambda')\rangle$$



$$|i(\lambda)\rangle$$

$$(ds^{(i)})^2 = g_{\alpha\beta}^{(i)}(\lambda) d\lambda_\alpha d\lambda_\beta$$

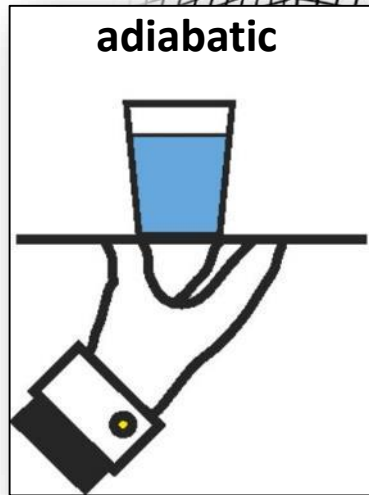
Driven dynamics: adiabatic driving

$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = \hat{H}(\lambda(t)) |\psi(t)\rangle$$

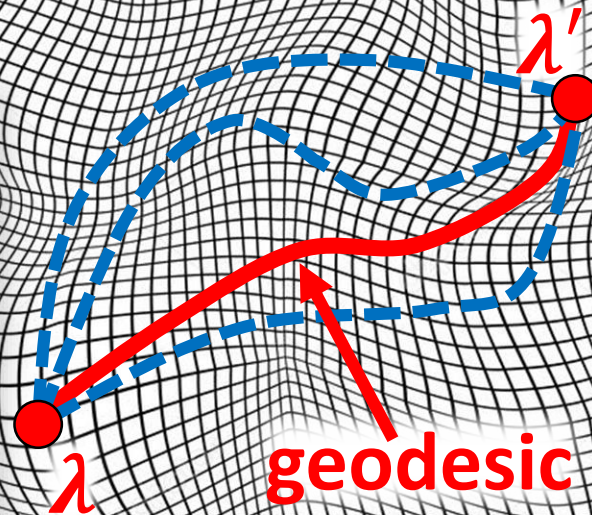
$$|\psi(0)\rangle = |i(\lambda)\rangle$$

$$|\psi(T)\rangle = |i(\lambda')\rangle$$

D.Sels, A.Polkovnikov
PNAS 114, E3909 (2017)



$$|i(\lambda)\rangle$$



$$\left(ds^{(i)} \right)^2 = g_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) d\lambda_{\alpha} d\lambda_{\beta}$$

Driven dynamics: adiabatic driving

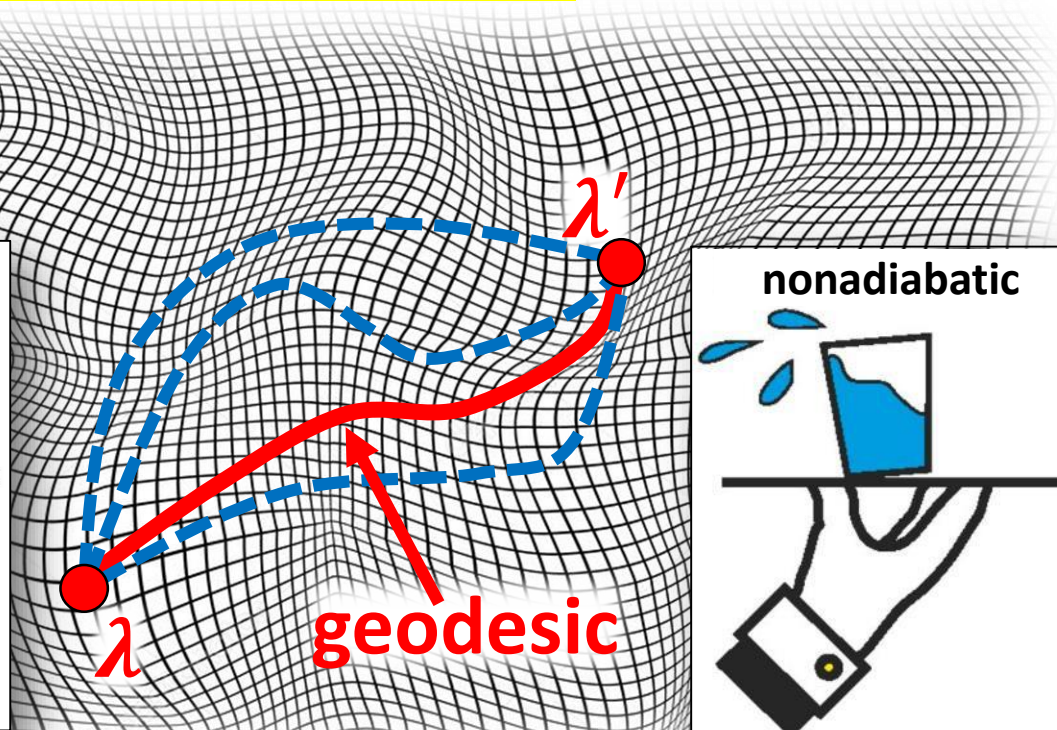
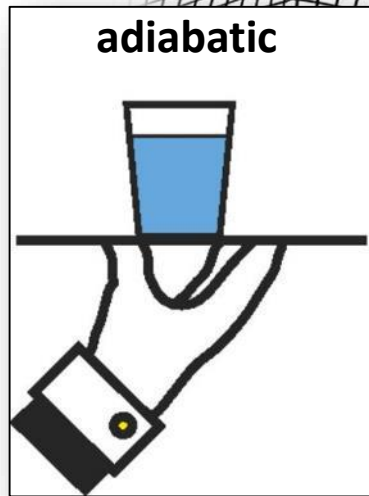
$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = \hat{H}(\lambda(t)) |\psi(t)\rangle$$

$$|\psi(0)\rangle = |i(\lambda)\rangle$$

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D.Sels, A.Polkovnikov
PNAS 114, E3909 (2017)

$$|i(\lambda)\rangle$$



$$\left(ds^{(i)} \right)^2 = g_{\alpha\beta}^{(i)}(\lambda) d\lambda_{\alpha} d\lambda_{\beta}$$

Driven dynamics: counterdiabatic driving

$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = \left[\hat{H}(\lambda(t)) + \dot{\lambda}(t) \cdot \hat{A}(\lambda(t)) \right] |\psi(t)\rangle$$

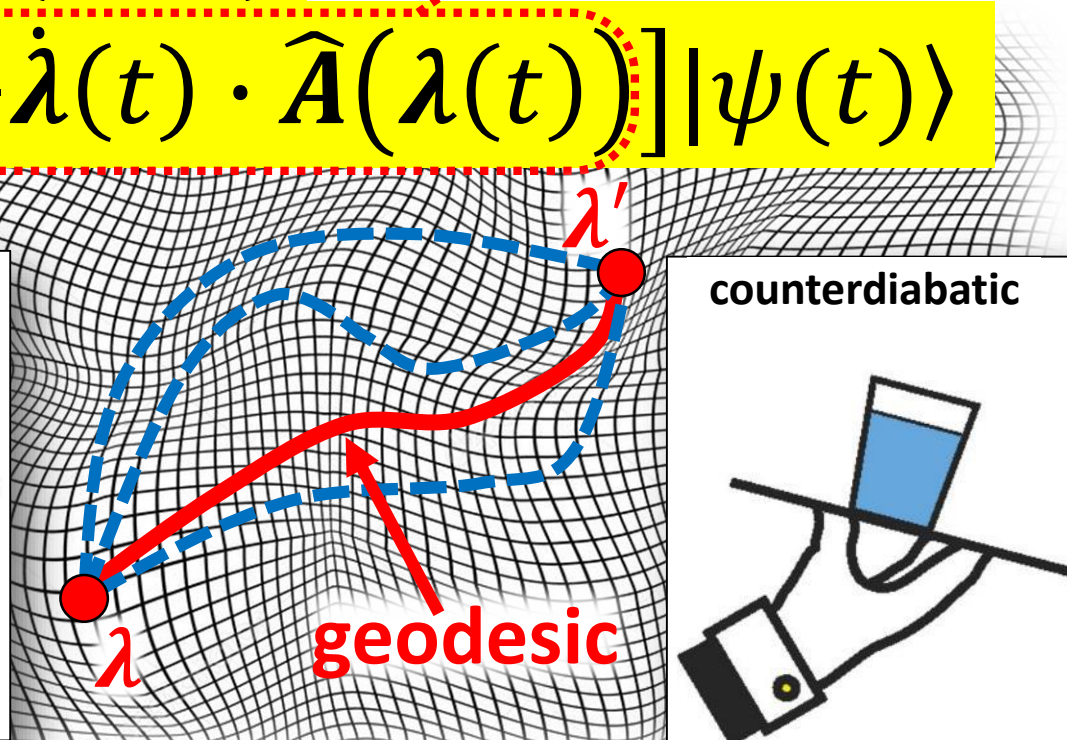
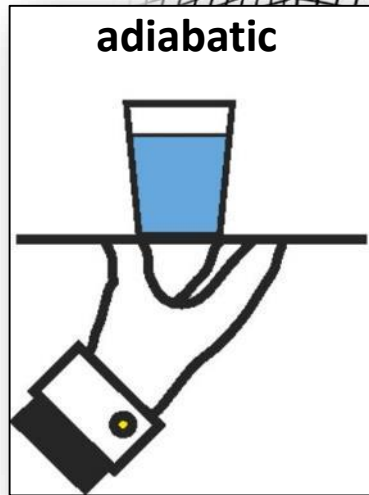
balancing term

$$|\psi(0)\rangle = |i(\lambda)\rangle$$

$$|\psi(T)\rangle = |i(\lambda')\rangle$$

D.Sels, A.Polkovnikov
PNAS 114, E3909 (2017)

$$|i(\lambda)\rangle$$



$$(ds^{(i)})^2 = g_{\alpha\beta}^{(i)}(\lambda) d\lambda_{\alpha} d\lambda_{\beta}$$

Piš, barde, středej ...
(Write, bard, collect ...)

\$15 million

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Metric tensor and curvature

$$g_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) = \mathbf{Re} \left(\left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle - \left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| i \right\rangle \left\langle i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle \right)$$
$$= \mathbf{Re} \sum_{j(\neq i)} \frac{\left\langle i \middle| \frac{\partial \hat{H}}{\partial \lambda_\alpha} \middle| j \right\rangle \left\langle j \middle| \frac{\partial \hat{H}}{\partial \lambda_\beta} \middle| i \right\rangle}{(E_i - E_j)^2}$$

Metric tensor

Metric tensor and curvature

$$g_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) = \mathbf{Re} \left(\left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle - \left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| i \right\rangle \left\langle i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle \right)$$
$$= \mathbf{Re} \sum_{j(\neq i)} \frac{\left\langle i \middle| \frac{\partial \hat{H}}{\partial \lambda_\alpha} \middle| j \right\rangle \left\langle j \middle| \frac{\partial \hat{H}}{\partial \lambda_\beta} \middle| i \right\rangle}{(E_i - E_j)^2}$$

Metric tensor

$$v_{\alpha\beta}^{(i)}(\boldsymbol{\lambda}) = -2 \mathbf{Im} \sum_{j(\neq i)} \frac{\left\langle i \middle| \frac{\partial \hat{H}}{\partial \lambda_\alpha} \middle| j \right\rangle \left\langle j \middle| \frac{\partial \hat{H}}{\partial \lambda_\beta} \middle| i \right\rangle}{(E_i - E_j)^2}$$

Curvature tensor

$$= -2 \mathbf{Im} \left(\left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle - \left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| i \right\rangle \left\langle i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle \right)$$

Metric tensor and curvature

$$g_{\alpha\beta}^{(i)}(\lambda) = \text{Re} \left(\left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle - \left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| i \right\rangle \left\langle i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle \right)$$

$$= \text{Re} \sum_{j(\neq i)} \frac{\langle i | \frac{\partial \hat{H}}{\partial \lambda_\alpha} | j \rangle \langle j | \frac{\partial \hat{H}}{\partial \lambda_\beta} | i \rangle}{(E_i - E_j)^2}$$

Metric tensor

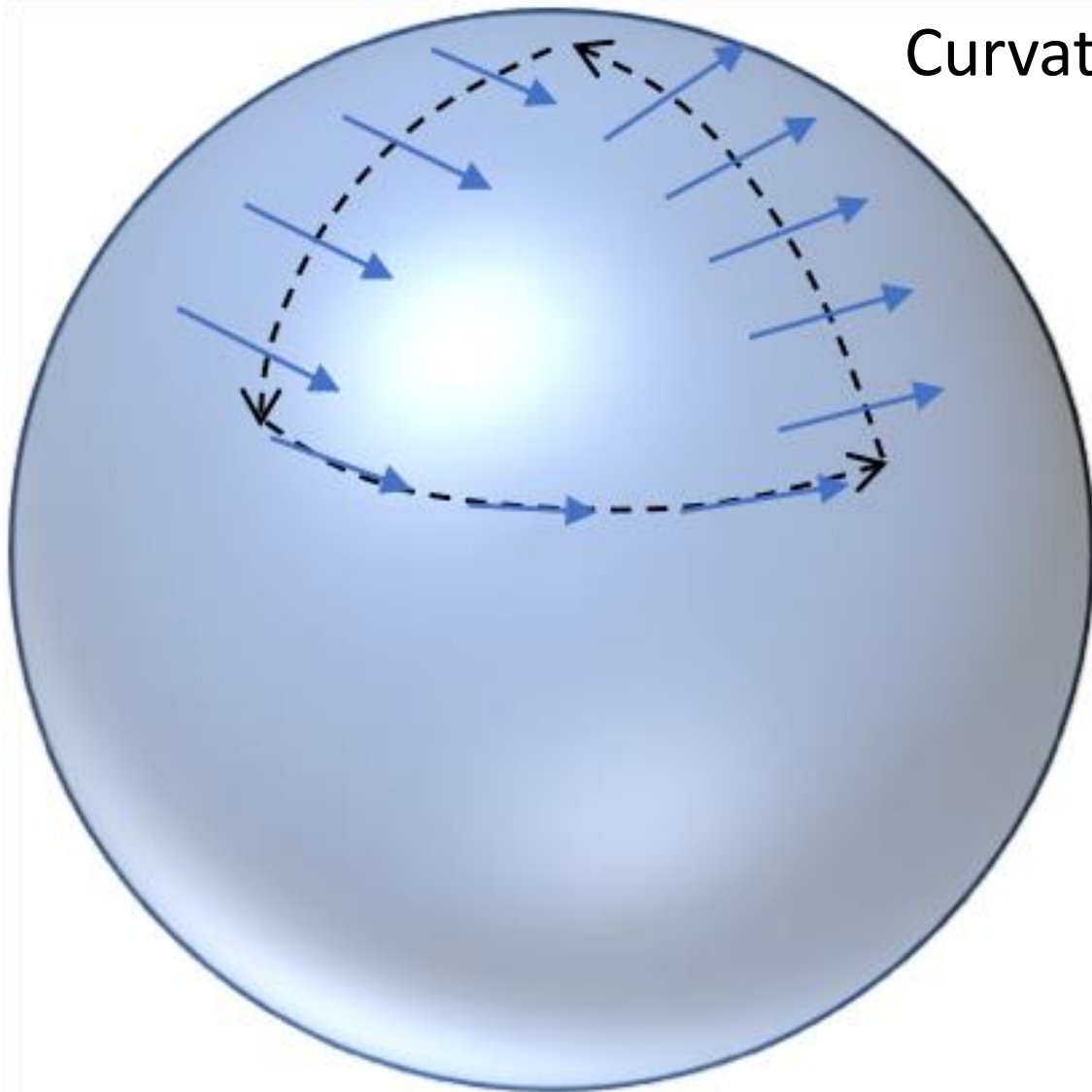
Geometric tensor $\chi_{\alpha\beta}^{(i)}(\lambda) = g_{\alpha\beta}^{(i)}(\lambda) - i 2v_{\alpha\beta}^{(i)}(\lambda)$

$$v_{\alpha\beta}^{(i)}(\lambda) = -2 \text{Im} \sum_{j(\neq i)} \frac{\langle i | \frac{\partial \hat{H}}{\partial \lambda_\alpha} | j \rangle \langle j | \frac{\partial \hat{H}}{\partial \lambda_\beta} | i \rangle}{(E_i - E_j)^2}$$

Curvature tensor

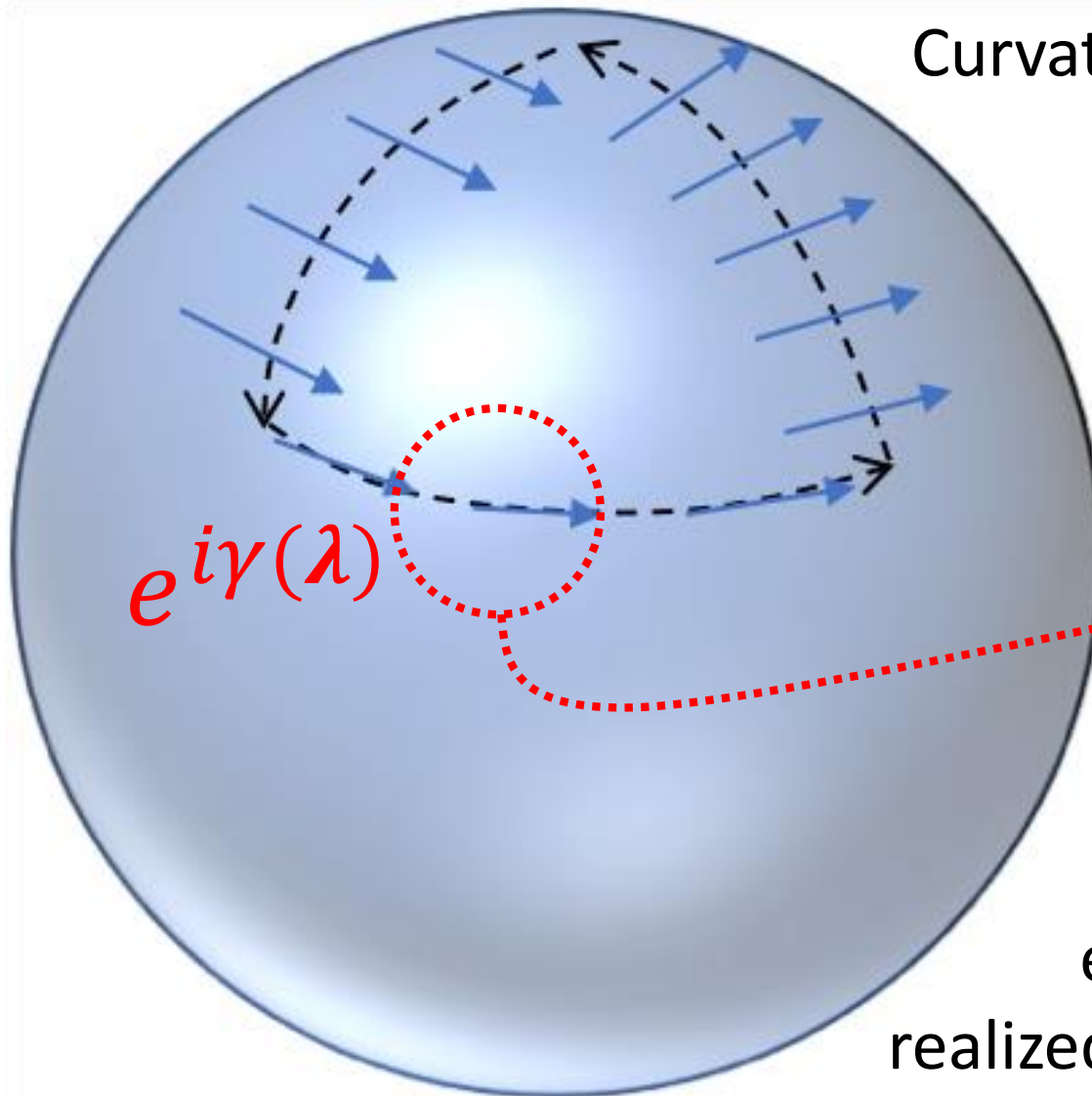
$$= -2 \text{Im} \left(\left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle - \left\langle \frac{\partial}{\partial \lambda_\alpha} i \middle| i \right\rangle \left\langle i \middle| \frac{\partial}{\partial \lambda_\beta} i \right\rangle \right)$$

Berry's phase



Curvature can be measured
by **parallel transport**

Berry's phase



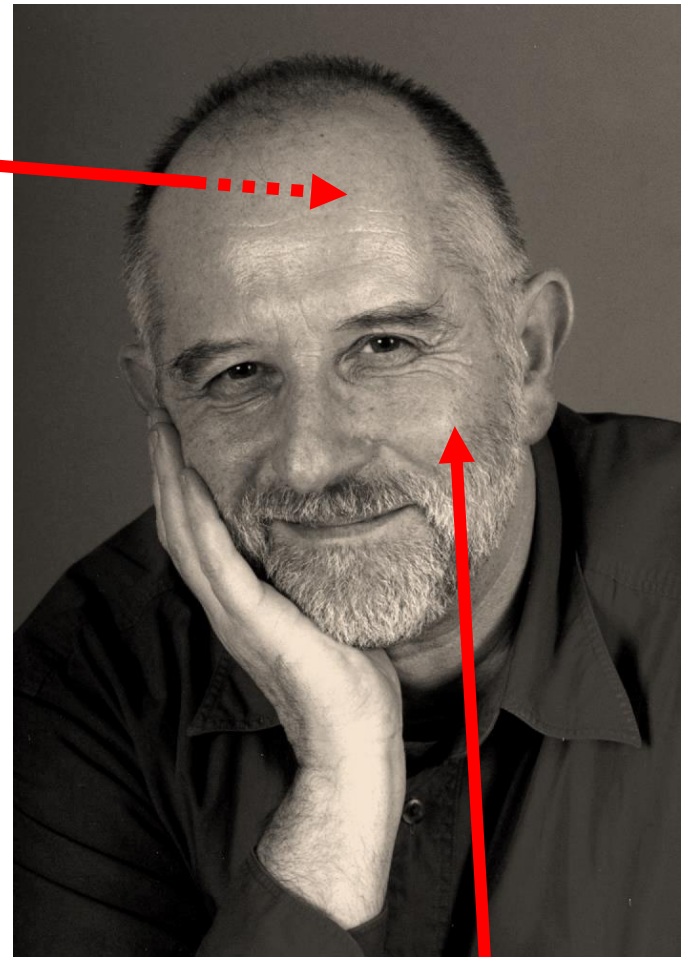
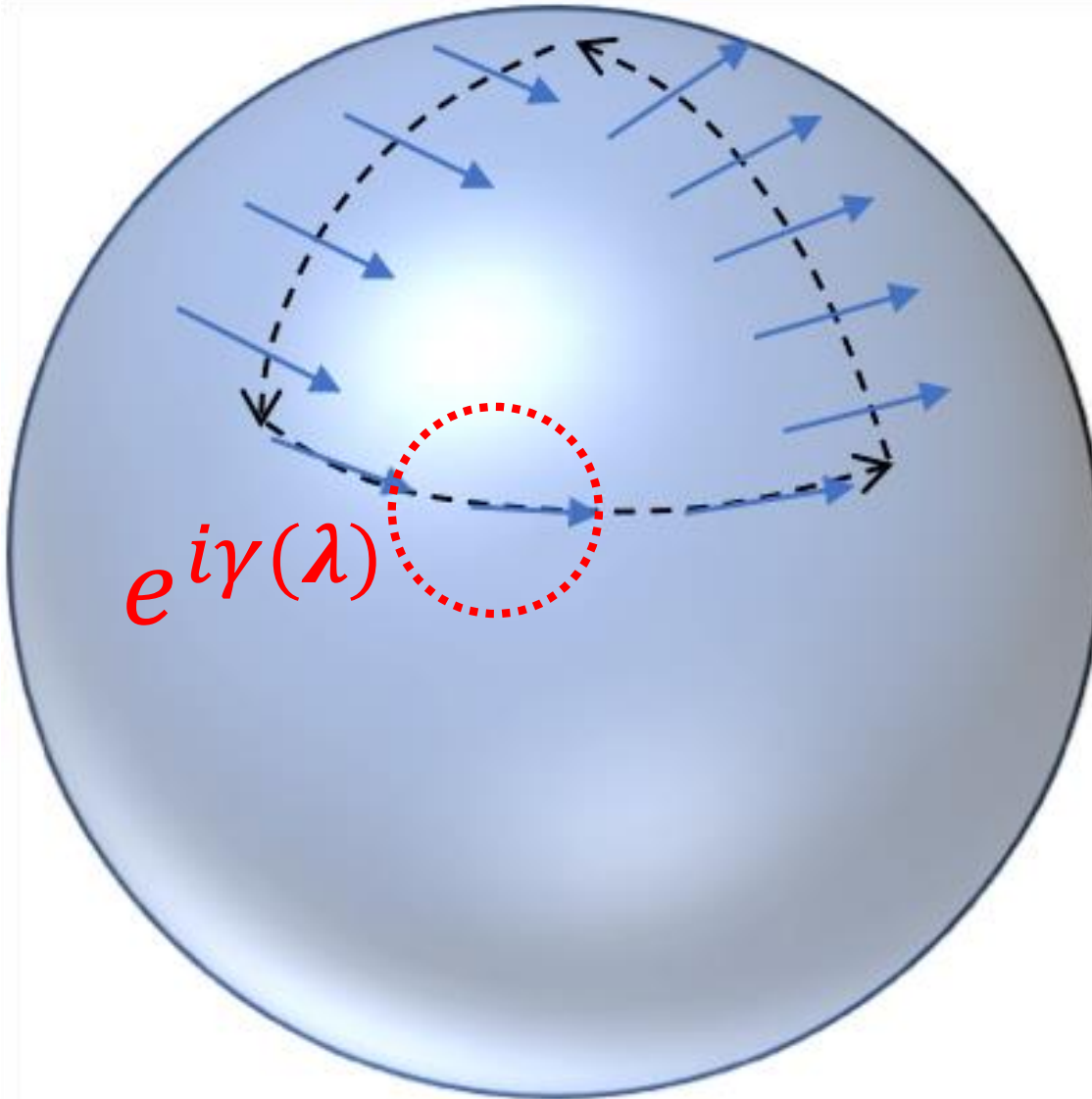
Curvature can be measured
by **parallel transport**

Local phase (gauge)
of the eigenvector

Parallel transport
in the Hamiltonian
eigenstate manifold is
realized by **adiabatic driving**

Berry's phase

*1983



Michael Berry

*1941

Berry's face

Berry's phase

Ig Nobel Prize 2000

Levitation without Meditation

Michael Berry & Andrey Geim



Degeneracies & avoided crossings

$$\chi_{\alpha\beta}^{(i)}(\lambda) = \sum_{j(\neq i)} \frac{\langle i | \frac{\partial \hat{H}}{\partial \lambda_\alpha} | j \rangle \langle j | \frac{\partial \hat{H}}{\partial \lambda_\beta} | i \rangle}{(E_i - E_j)^2} \quad \text{!!!}$$

Singularities

$$E_i = E_j$$

Degeneracies & avoided crossings

$$\chi_{\alpha\beta}^{(i)}(\lambda) = \sum_{j(\neq i)} \frac{\langle i | \frac{\partial \hat{H}}{\partial \lambda_\alpha} | j \rangle \langle j | \frac{\partial \hat{H}}{\partial \lambda_\beta} | i \rangle}{(E_i - E_j)^2} \quad \text{!!!}$$

Singularities $E_i = E_j$

Hamiltonian 2 x 2

$$\hat{H}(\lambda) = \begin{pmatrix} H_{11}(\lambda) & H_{12}(\lambda) \\ H_{12}^*(\lambda) & H_{22}(\lambda) \end{pmatrix}$$

solutions

$$E_{\pm}(\lambda) = \frac{H_{11}(\lambda) + H_{22}(\lambda)}{2} \pm \frac{1}{2} \sqrt{\left(\frac{H_{11}(\lambda) - H_{22}(\lambda)}{2} \right)^2 + (\operatorname{Re} H_{12}(\lambda))^2 + (\operatorname{Im} H_{12}(\lambda))^2}$$

eigenvalue equation

$$\operatorname{Det}[\hat{H}(\lambda) - E \hat{\mathbb{1}}] = 0$$

Degeneracies & avoided crossings

$$\chi_{\alpha\beta}^{(i)}(\lambda) = \sum_{j(\neq i)} \frac{\langle i | \frac{\partial \hat{H}}{\partial \lambda_\alpha} | j \rangle \langle j | \frac{\partial \hat{H}}{\partial \lambda_\beta} | i \rangle}{(E_i - E_j)^2} \quad \text{!!!}$$

Singularities $E_i = E_j$

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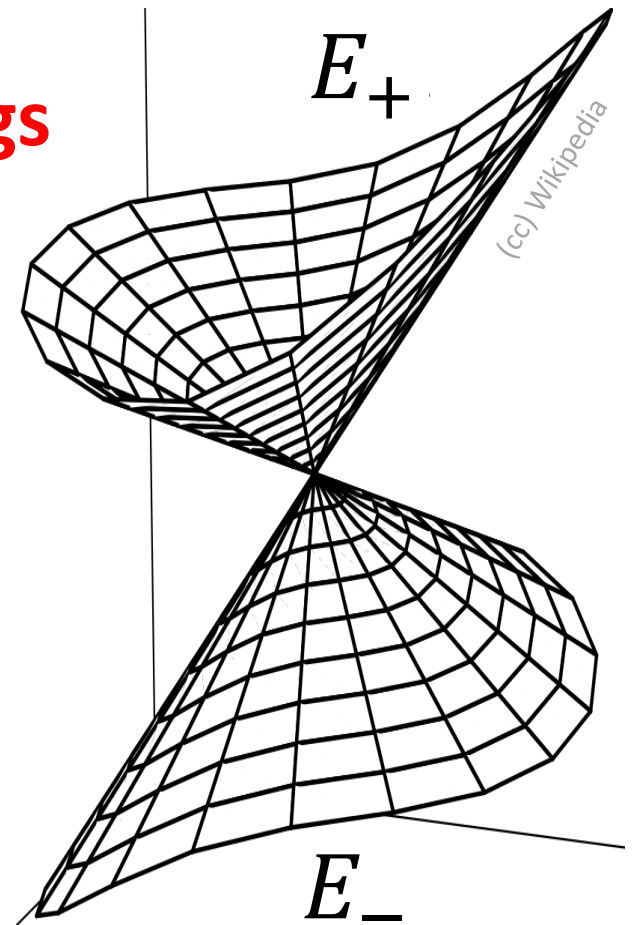
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0 $\stackrel{!}{=}$

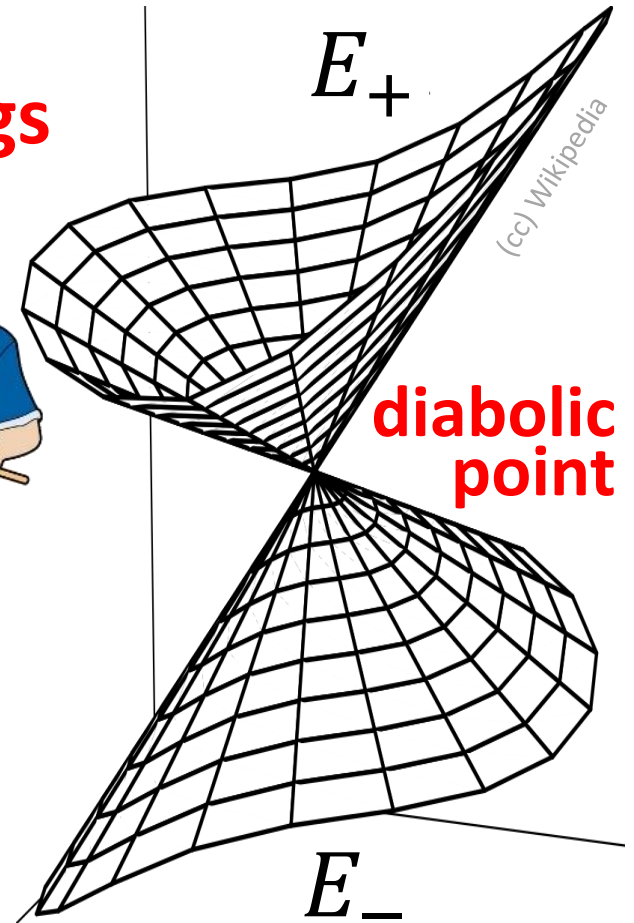
conical intersection

$$\propto |\lambda - \lambda_0|$$



Degeneracies & avoided crossings

$$\chi_{\alpha\beta}^{(i)}(\lambda) = \sum_{j(\neq i)} \frac{\langle i | \frac{\partial \hat{H}}{\partial \lambda_\alpha} | j \rangle \langle j | \frac{\partial \hat{H}}{\partial \lambda_\beta} | i \rangle}{(E_i - E_j)^2} \quad !!!$$



Singularities $E_i = E_j$

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0 $\stackrel{!}{=}$

conical intersection

$$\propto |\lambda - \lambda_0|$$

$$\propto |\lambda - \lambda_0|^2$$

$$\propto |\lambda - \lambda_0|^2$$

$$\propto |\lambda - \lambda_0|^2$$

Degeneracies & avoided crossings

$$\chi_{\alpha\beta}^{(i)}(\lambda) = \sum_{j(\neq i)} \frac{\langle i | \frac{\partial \hat{H}}{\partial \lambda_\alpha} | j \rangle \langle j | \frac{\partial \hat{H}}{\partial \lambda_\beta} | i \rangle}{(E_i - E_j)^2} \quad !!!$$

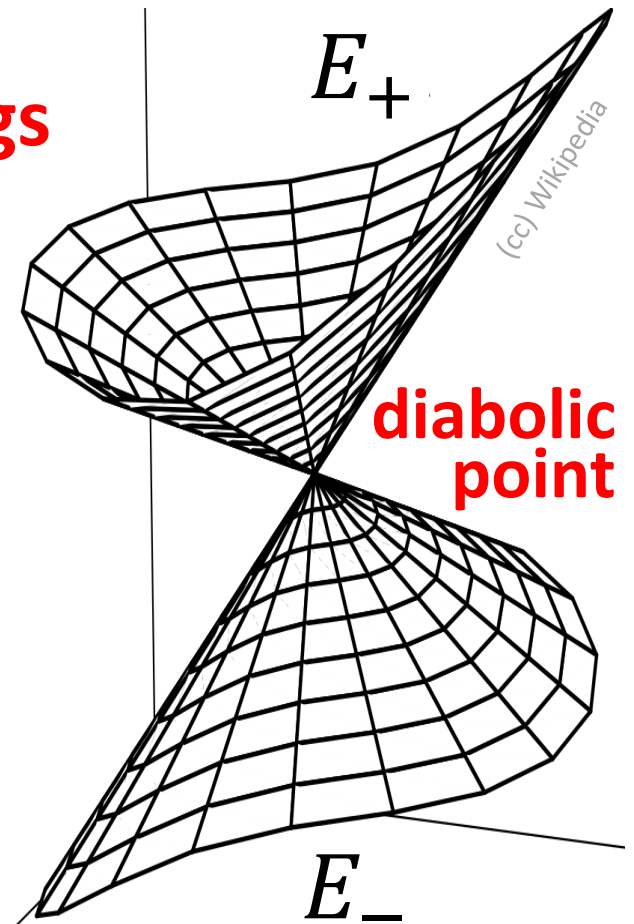
Singularities $E_i = E_j$

Hamiltonian 2 x 2

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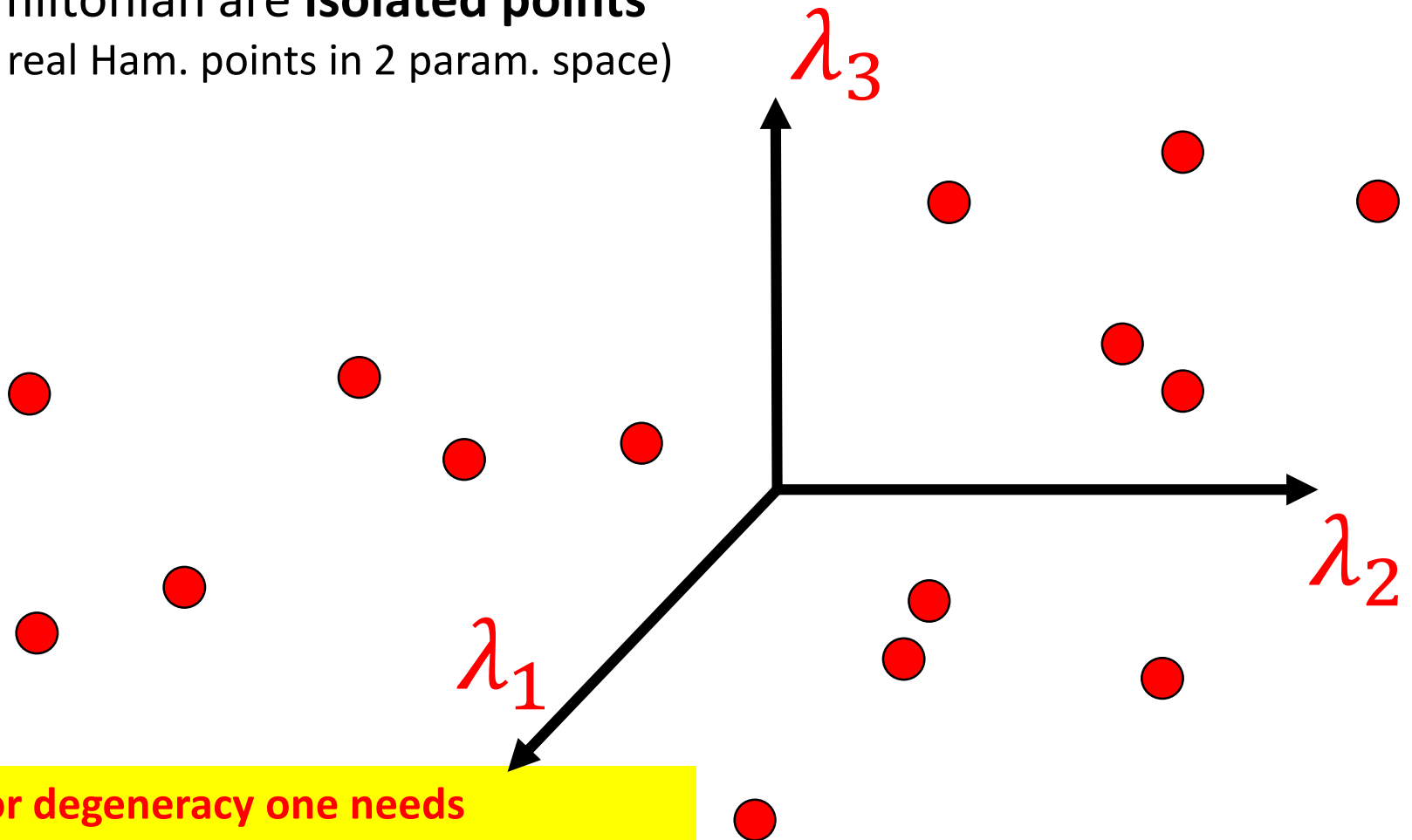


For degeneracy one needs
3 independent functions to vanish!

Degeneracies & avoided crossings

For systems with 3 parameters degeneracies of a complex Hamiltonian are **isolated points**

(for real Ham. points in 2 param. space)



For degeneracy one needs
3 independent functions to vanish!

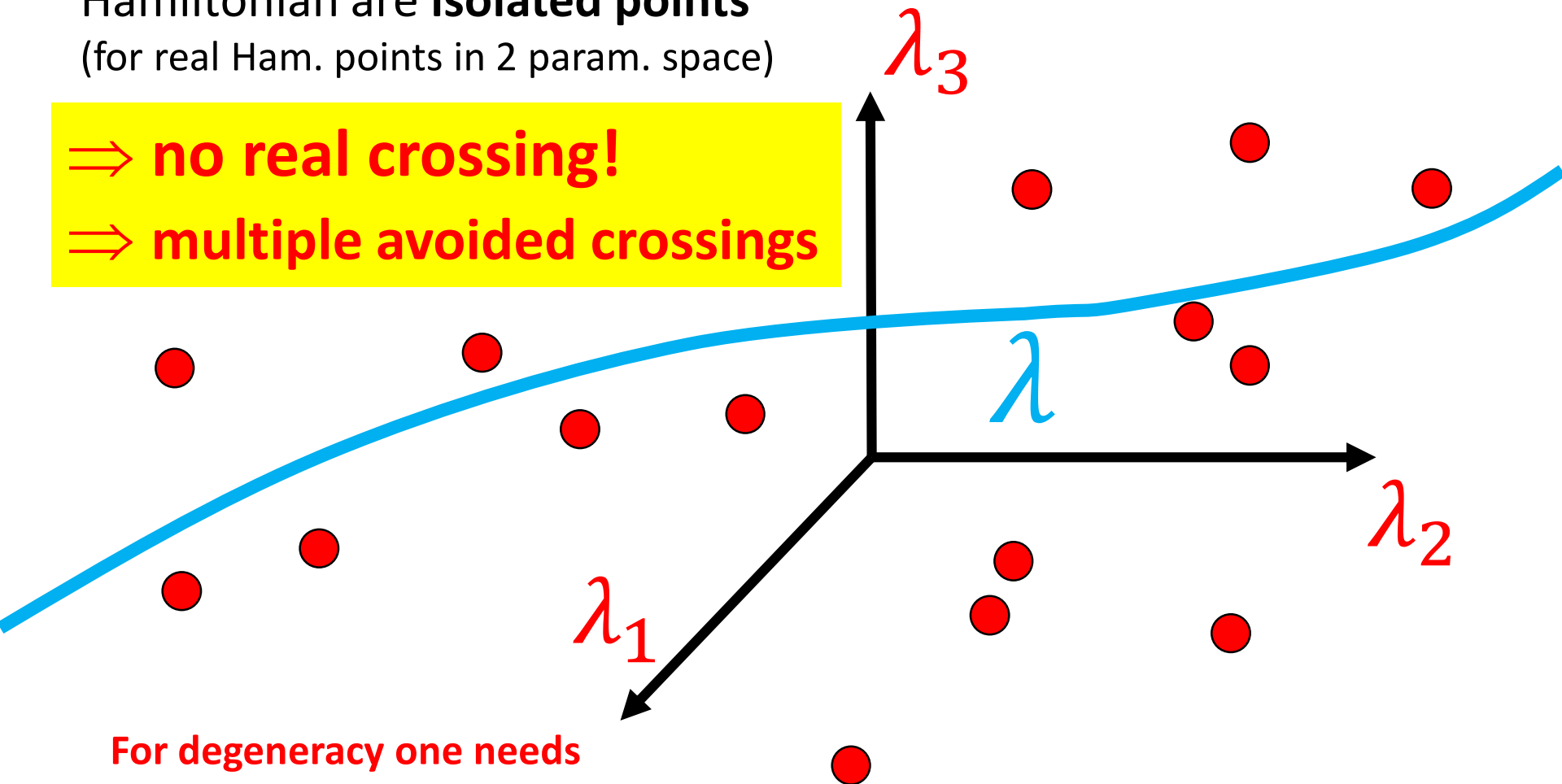
Degeneracies & avoided crossings

For systems with 3 parameters degeneracies of a complex Hamiltonian are **isolated points**

(for real Ham. points in 2 param. space)

⇒ **no real crossing!**

⇒ **multiple avoided crossings**



**For degeneracy one needs
3 independent functions to vanish!**

Degeneracies & avoided crossings



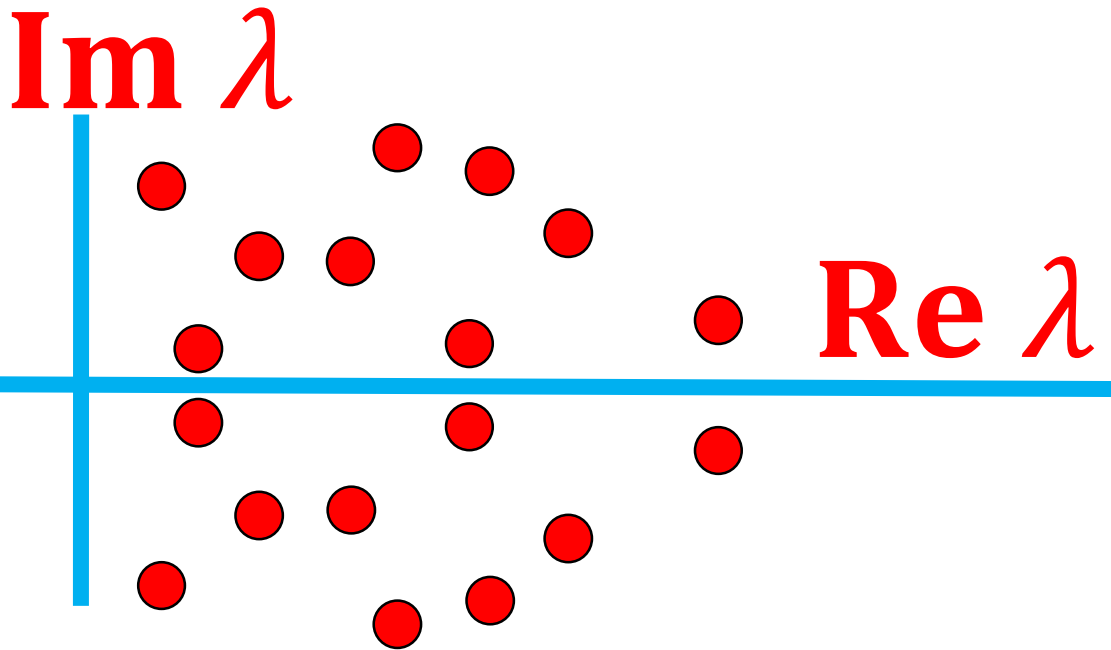
A diagram illustrating energy levels E as a function of a parameter λ . The vertical axis is labeled E and the horizontal axis is labeled λ . The plot shows a dense set of black lines representing energy levels. As λ increases, the energy levels generally trend downwards. Many levels cross each other, but at certain points, they do not cross; instead, they form avoided crossings, where the two levels repel each other and then cross again. This behavior is characteristic of systems with discrete energy levels that interact with each other.

E

λ

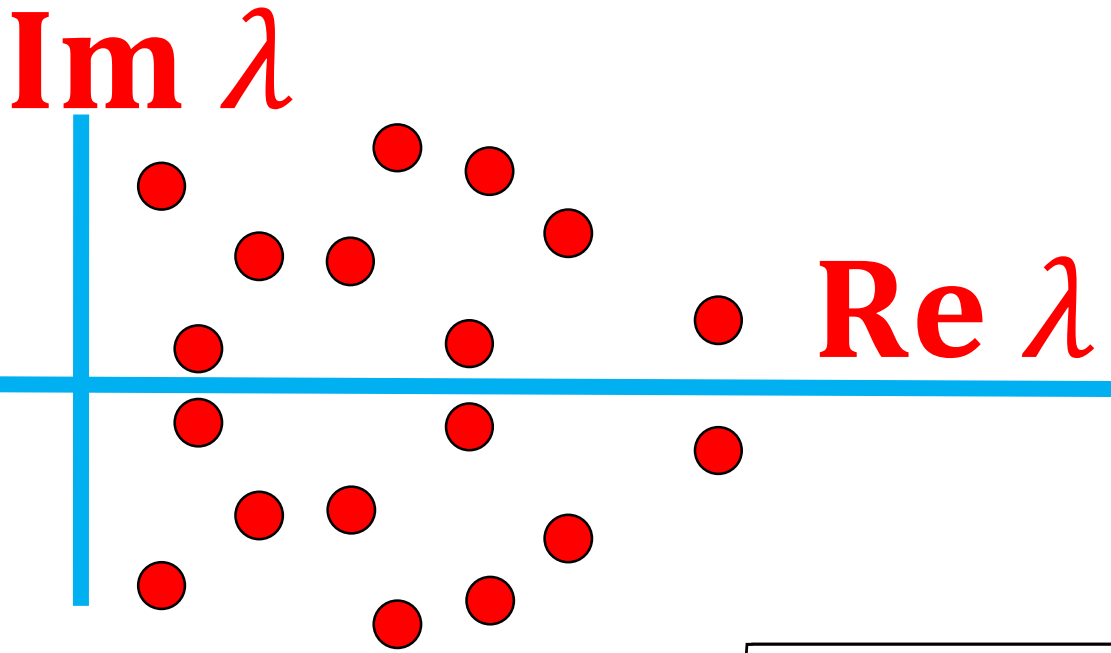
Degeneracies & avoided crossings } & complex extensions

For systems with 1 parameter we seek for degeneracies in $\lambda \in \mathbb{C}$



Degeneracies & avoided crossings } & complex extensions

For systems with 1 parameter we seek for degeneracies in $\lambda \in \mathbb{C}$

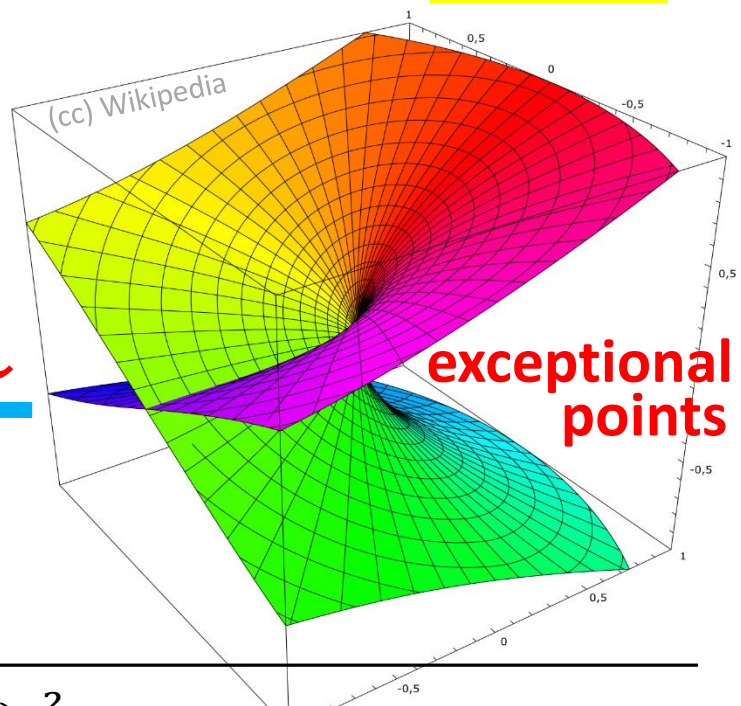
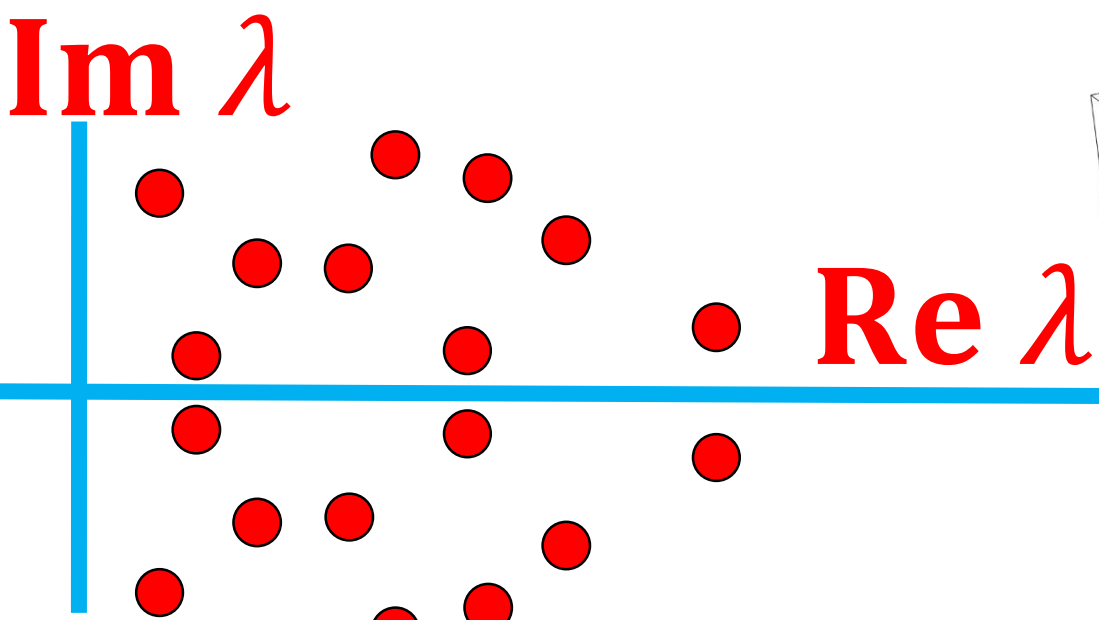


$$E_{\pm}(\lambda) = \frac{H_{11}(\lambda) + H_{22}(\lambda)}{2} \pm \frac{1}{2} \sqrt{\underbrace{\left(\frac{H_{11}(\lambda) - H_{22}(\lambda)}{2} \right)^2 + (\operatorname{Re} H_{12}(\lambda))^2 + (\operatorname{Im} H_{12}(\lambda))^2}_{\propto (\lambda - \lambda_0)}}}$$

$0 \stackrel{!}{=} \underbrace{\hspace{15em}}_{\propto \sqrt{\lambda - \lambda_0}}$

Degeneracies & avoided crossings & complex extensions

For systems with 1 parameter we seek for degeneracies in $\lambda \in \mathbb{C}$



$$E_{\pm}(\lambda) = \frac{H_{11}(\lambda) + H_{22}(\lambda)}{2} \pm \frac{1}{2} \sqrt{\left(\frac{H_{11}(\lambda) - H_{22}(\lambda)}{2}\right)^2 + (\text{Re } H_{12}(\lambda))^2 + (\text{Im } H_{12}(\lambda))^2}$$

0 $\stackrel{!}{=}$ → $\underbrace{\left(\frac{H_{11}(\lambda) - H_{22}(\lambda)}{2}\right)^2 + (\text{Re } H_{12}(\lambda))^2 + (\text{Im } H_{12}(\lambda))^2}_{\propto (\lambda - \lambda_0)}$

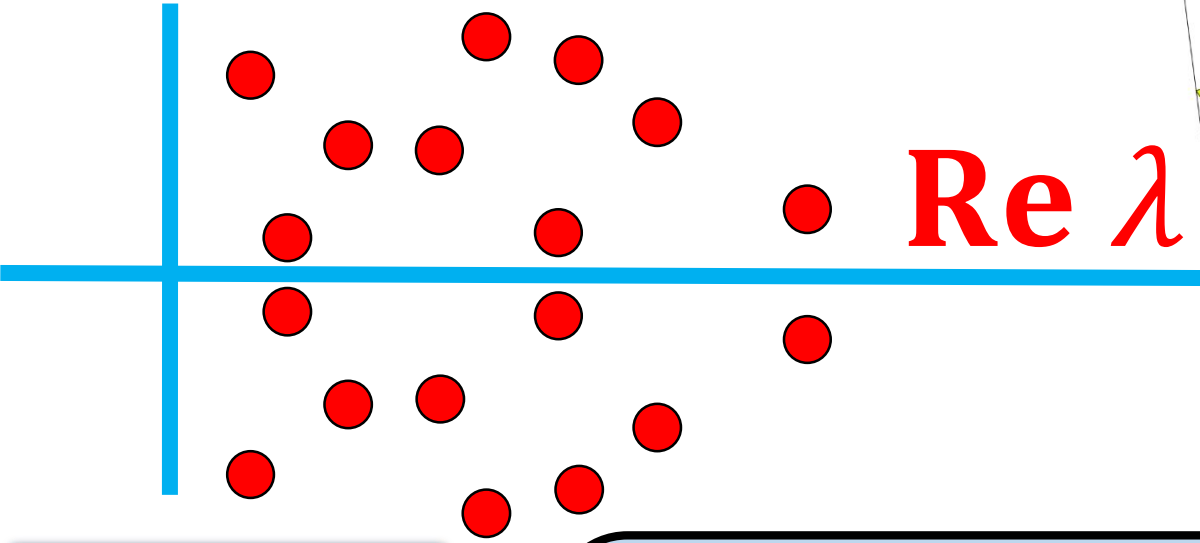
square root branching point $\propto \sqrt{\lambda - \lambda_0}$

Degeneracies & avoided crossings

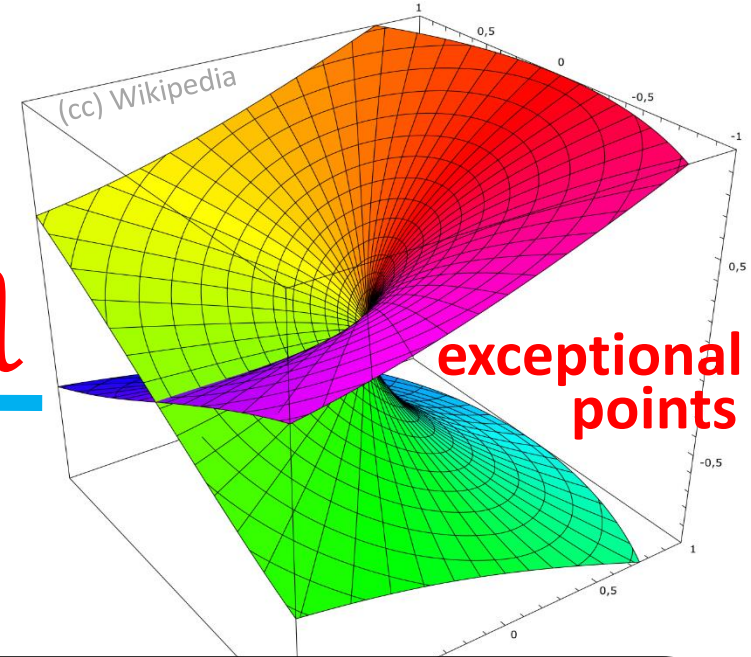
& complex extensions

For systems with 1 parameter we seek for degeneracies in $\lambda \in \mathbb{C}$

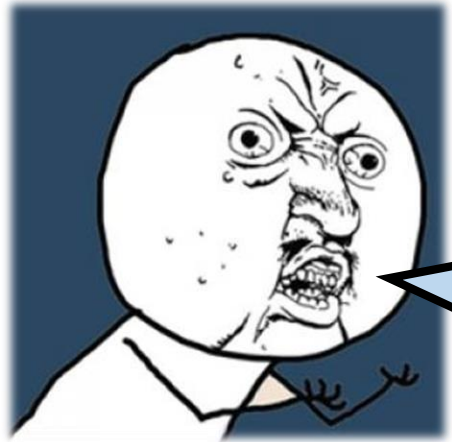
$\text{Im } \lambda$



$\text{Re } \lambda$

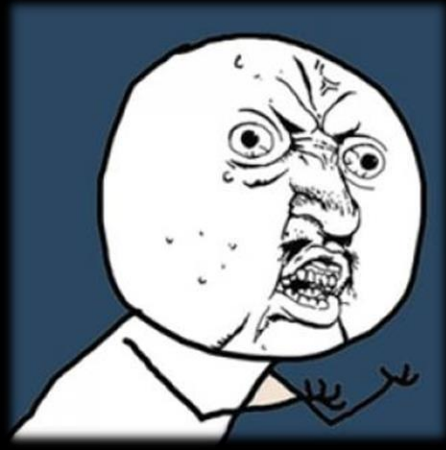


exceptional points



Accumulation of exceptional points near $(\text{Re}\lambda, \text{Im}\lambda) = (\lambda_c, 0)$ indicates a **quantum critical point** ...

This is where our work starts.

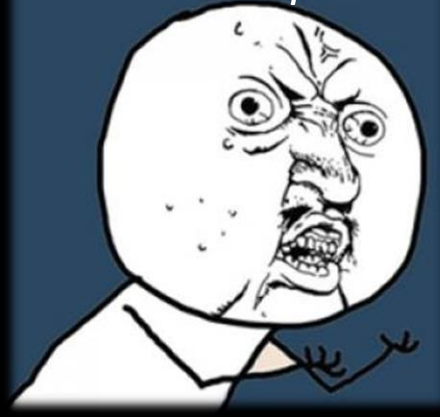


Thank you!

This is where our work starts.

Some home reading:

- M.V. Berry, in: *Geometric Phases in Physics*, edited by A. Shapere, F. Wilczek (World Scientific, Singapore, 1989)
“Quantal phase factors accompanying adiabatic changes” (1984)
“The quantum phase, five years after” (1988)
- M. Kolodrubetz, D. Sels, P. Mehta, A. Polkovnikov, *Physics Reports* 697 (2017) 1
“Geometry and non-adiabatic response in quantum and classical systems”
- M. Tomka, T. Souza, S. Rosenberg, A. Polkovnikov, arXiv:1606.05890
“Geodesic paths for quantum many-body systems”
- P. Stránský, M. Dvořák, P. Cejnar, *Physical Review E* 97 (2018) 012112
“Exceptional points near first- and second-order quantum phase transitions”



Thank you!