



# Damage mechanisms and limits of Nb<sub>3</sub>Sn due to direct high energy proton beam impact

Project overview and status

A. Will, J. Schubert, D. Wollmann, *et al.*



# Superconducting wires/ tapes

Nb-Ti strand (LHC)

Nb<sub>3</sub>Sn strand (HL-LHC)

HTS tapes  
(future acc. magnets..?)

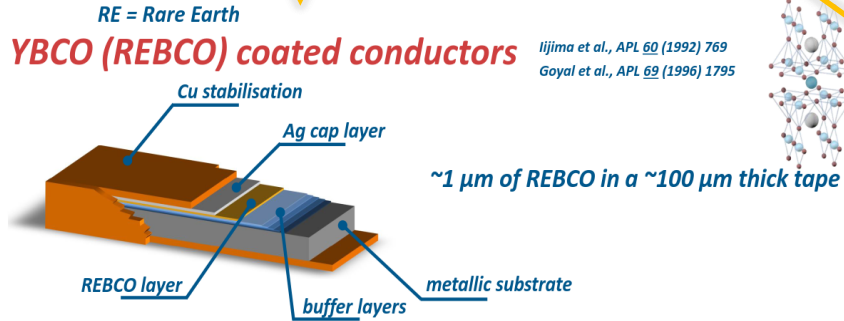


Image from C. Senatore, Lecture on Superconductivity, 2018

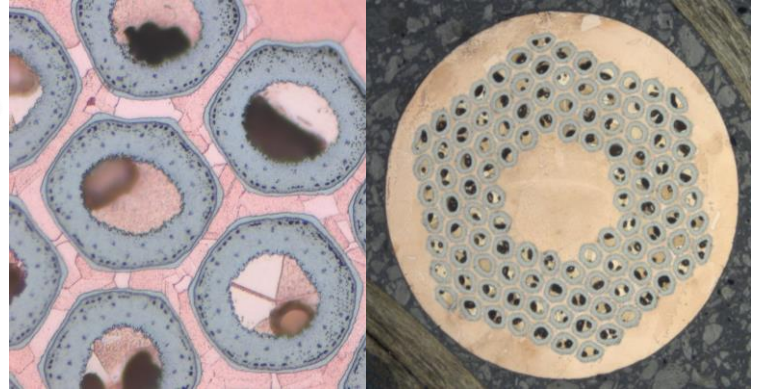
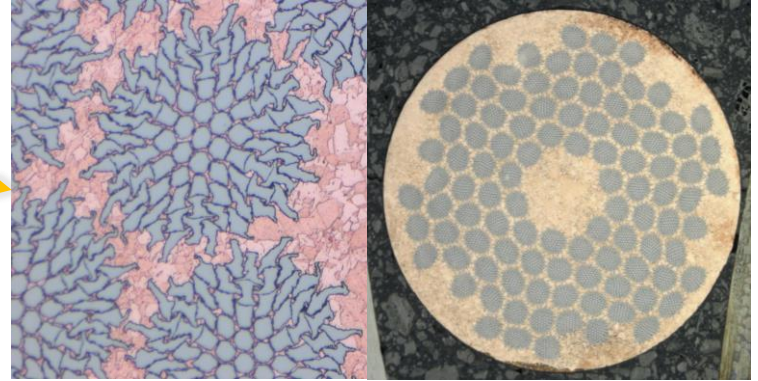


Image courtesy M. Meyer, CERN

# Beam impact experiment at cryogenic temperatures

- First **cryogenic beam impact experiment** was performed in August 2018
- Irradiation facility **HiRadMat** provides 440 GeV proton beams towards fixed targets
  - $1.2 \times 10^{11}$  ppb
  - 25 ns bunch spacing
  - 3.4 MJ max. per pulse
  - beam size  $\sim 1$ mm



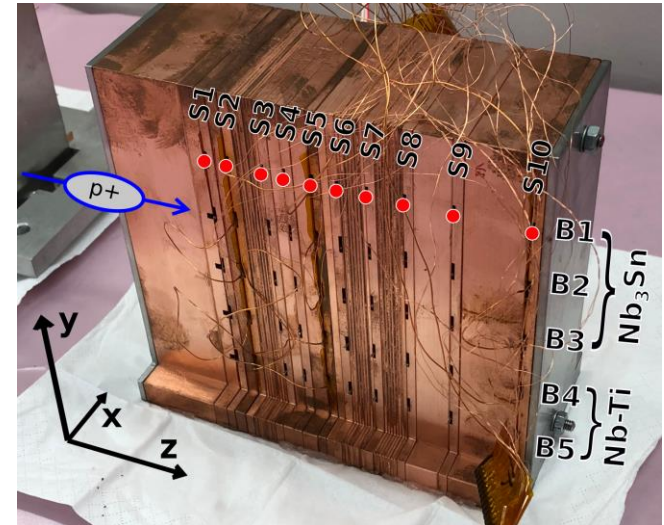
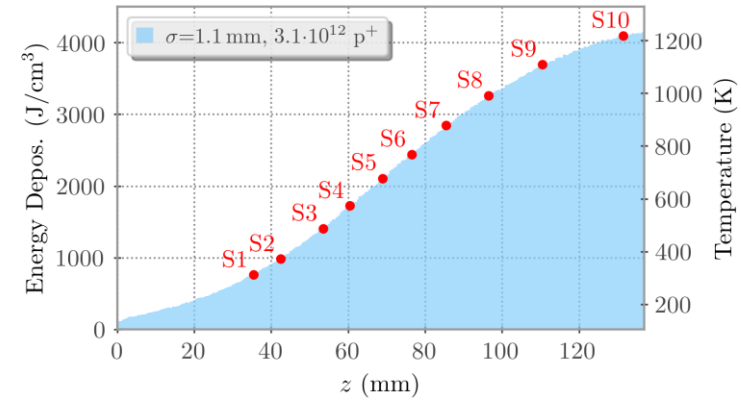
# Sample Holder

## Energy deposition

- **Temperature** reached along a **block of copper** for 440 GeV proton beam, simulated with FLUKA
- Embedded samples along beam axis reach various hot spot temperatures

## As in real failure case

- proton beam interacts with matter before impacting superconductors



# Nb<sub>3</sub>Sn Strain dependence

Mentink, Goedeke et al.

- Scaling law for Nb<sub>3</sub>Sn from literature
- normalized strain dependent function  $s(\varepsilon)$ 
  - $s(0) = 1$
  - monotonic decreasing towards zero for higher strains
  - simplest assumption  $s(\varepsilon) = 1 - \varepsilon^n$
- Kramer form linearizes the  $I_c$  curve
- Flux-pinning force

$$J_c(B, T, \varepsilon) = C_1 (1 - t^2) b^{-0.5} (1 - b)^2$$

$$t = \frac{T}{T_c \cdot s(\varepsilon)^{\frac{1}{3}}}$$

$$b = \frac{B}{B_{c2}(T, \varepsilon)}$$

$$B_{c2}(T, \varepsilon) = B_{c2,0} \cdot s(\varepsilon) \cdot (1 - t^{1.5})$$

$$f_K = J_c^{0.5} \cdot B^{0.25} \propto (1 - b) = \frac{B_{c2} - B}{B_{c2}}$$

$$\vec{F}_p = \vec{J}_c \times \vec{B}$$

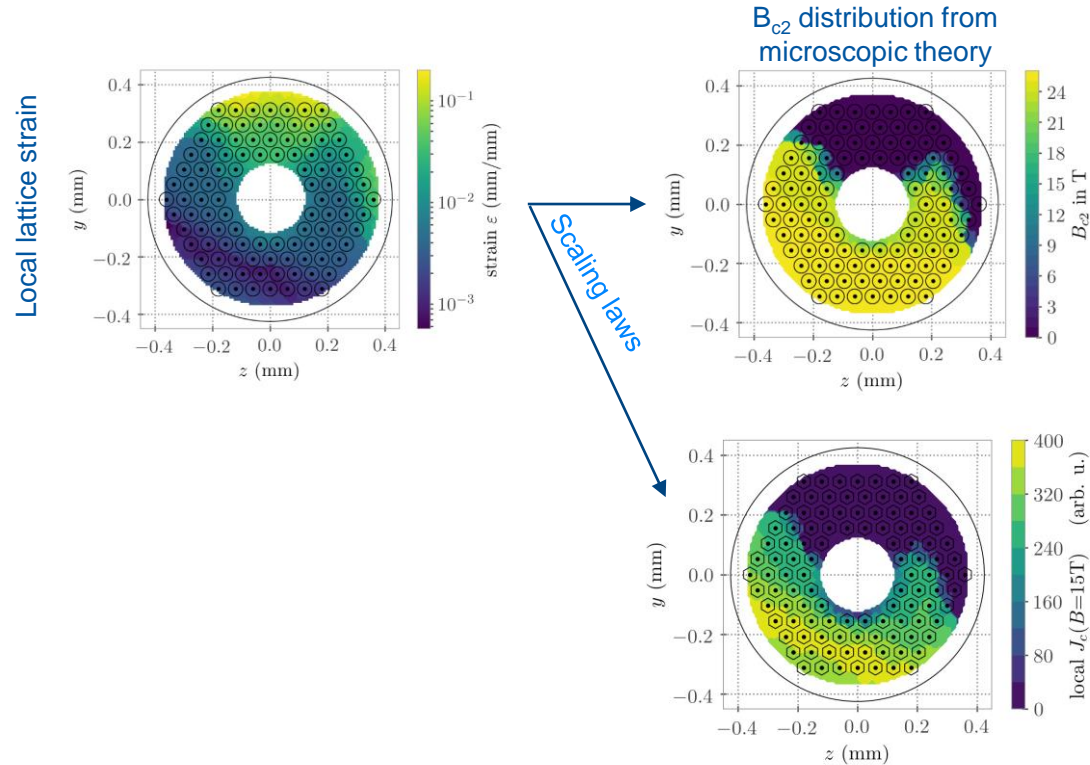
# Results of thermo-mechanical simulations

Thermo mechanical simulations performed\* with ANSYS

- study the stresses within the strand cross-section due to thermal gradients and expansion
- study the interaction of the strands with the sample holder

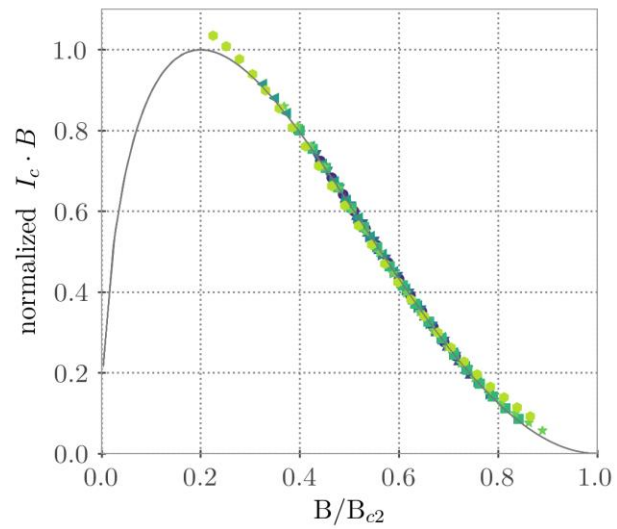
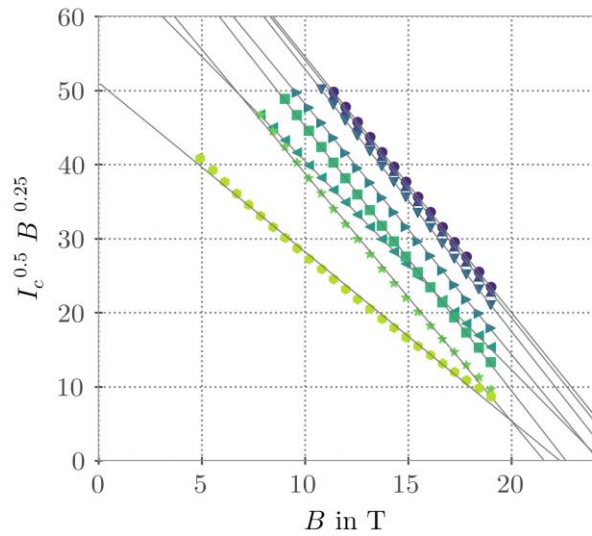
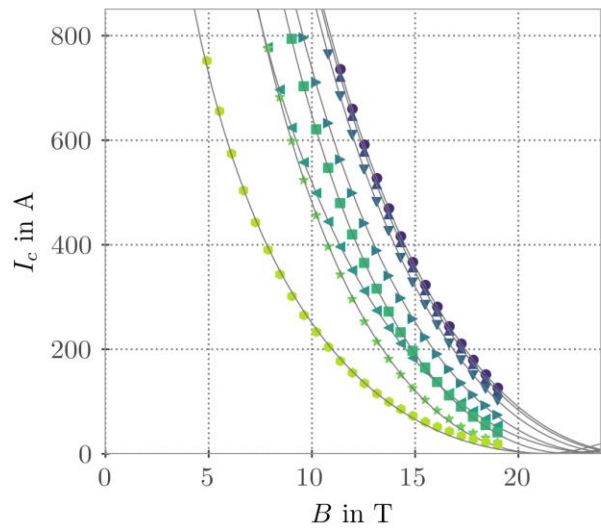
Use scaling laws from literature to predict  $I_c$  degradation of full strand

Filament cracking via “strain irreversibility cliff” model (Cheggour et al.)

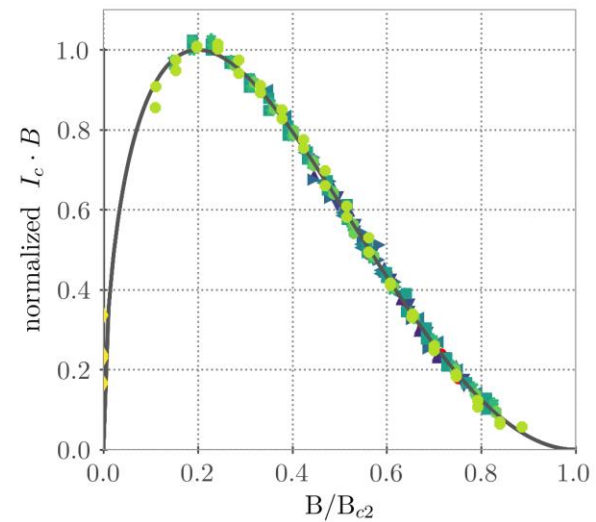
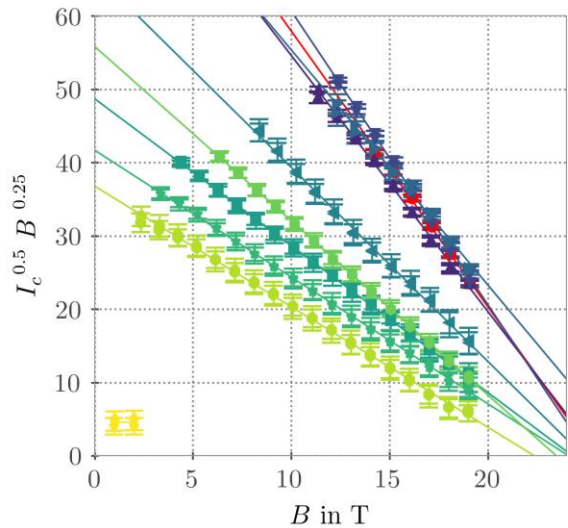
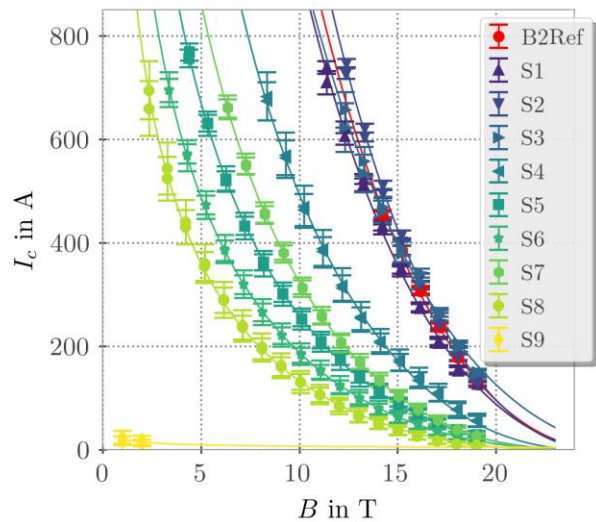




# Simulations



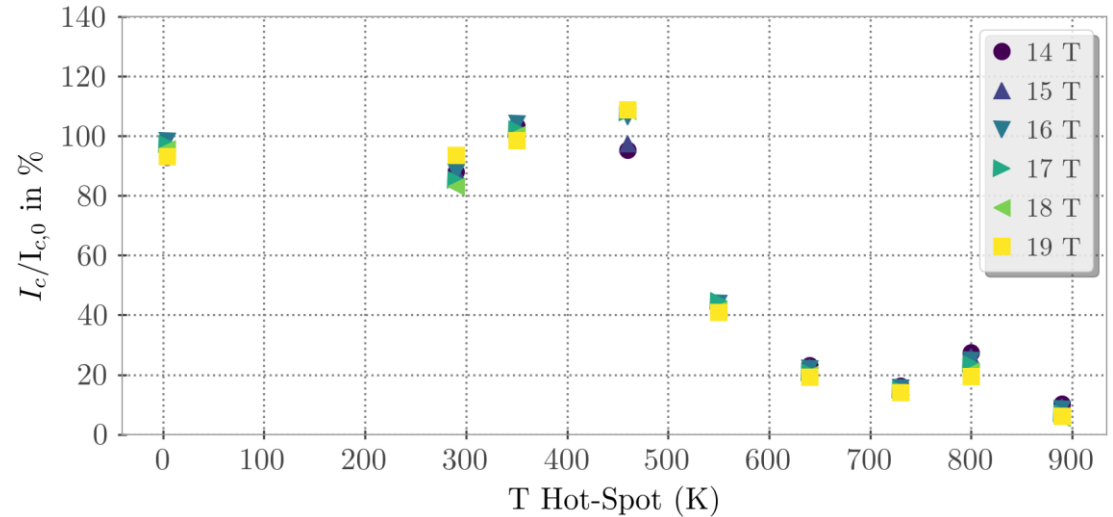
# Measurements





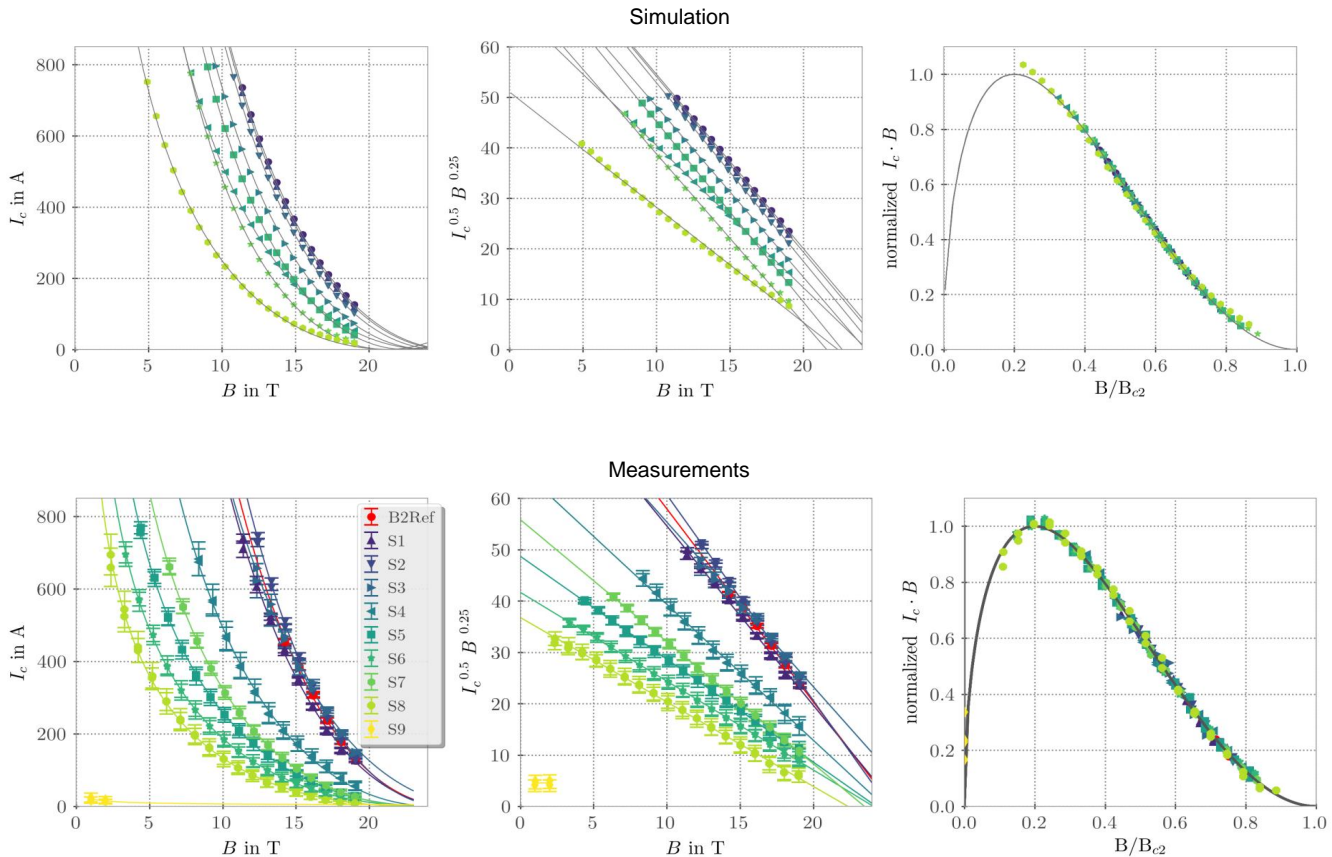
# Absolute $I_c$ degradation

- Absolute  $I_c$  degrades strongly above hot-spot temperature of  $\sim 450\text{K}$
- Almost independent of magnetic field



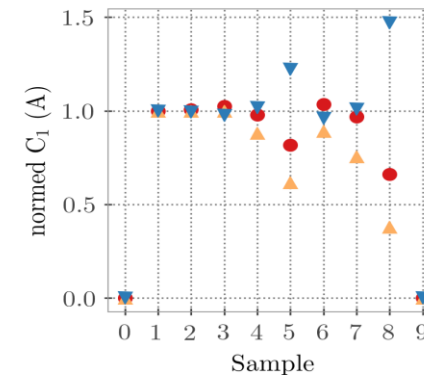
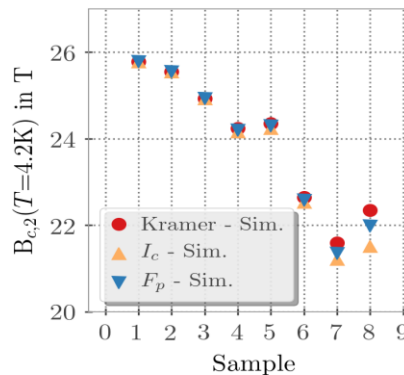
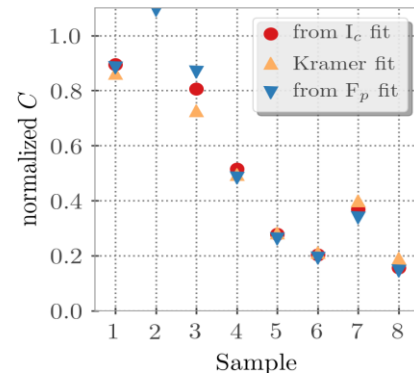
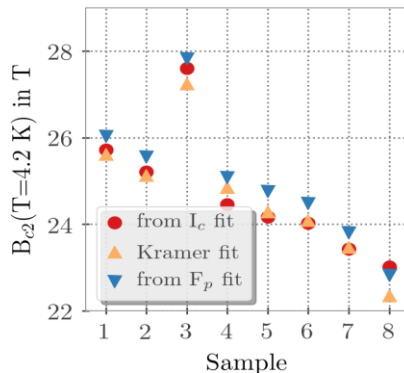
# Comparison

- $I_c$  of S1-3 barely affected by strain dependent degradation (both Simulation and Measurement)
- From S4 much stronger degradation is observed than predicted by strain scaling only
- Shift in  $B_{c2}$  in all type of fits
- Pinning force maximum unaltered



# Fit-results comparison

- $I_c$  of S1-3 barely affected by strain dependent degradation (both Simulation and Measurement)
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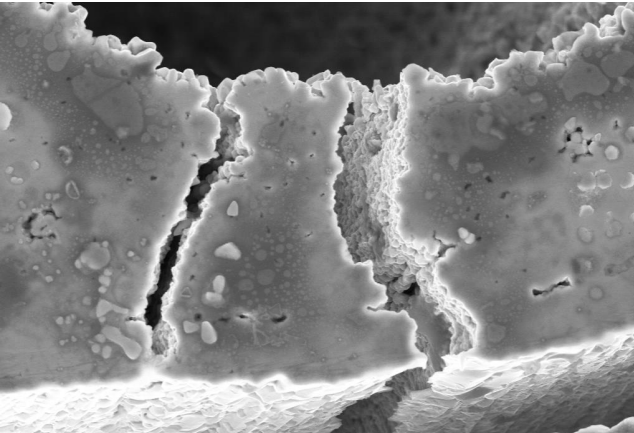


# Filament Cracking

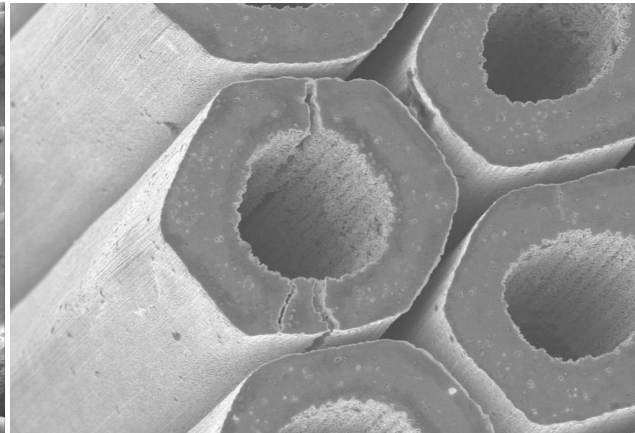
- strain alone can not explain the large reduction in absolute  $I_c$  (independent of T and B)
- Reduction in superconducting cross section could explain
  - Filament cracking due to too large strain?

$$J_c(B, T, \varepsilon) = C_1 (1 - t^2) b^{-0.5} (1 - b)^2$$

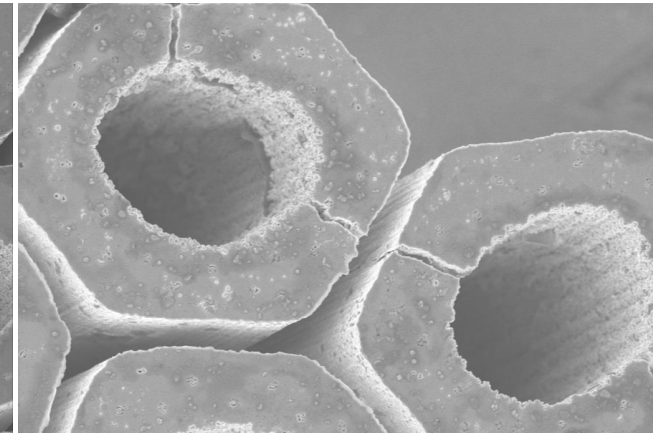
Nb3Sn samples from HRMT31 Experiment, Thanks to M. Meyer from the SEM Lab



2  $\mu\text{m}$  | EHT = 10.00 kV  
WD = 11.2 mm  
Signal A = InLens  
Meyer Mickael  
Date :9 Jul 2019  
Mag = 5.00 K X



20  $\mu\text{m}$  | EHT = 10.00 kV  
WD = 11.3 mm  
Signal A = InLens  
Meyer Mickael  
Date :9 Jul 2019  
Mag = 903 X



20  $\mu\text{m}$  | EHT = 10.00 kV  
WD = 11.2 mm  
Signal A = InLens  
Meyer Mickael  
Date :9 Jul 2019  
Mag = 1.26 K X

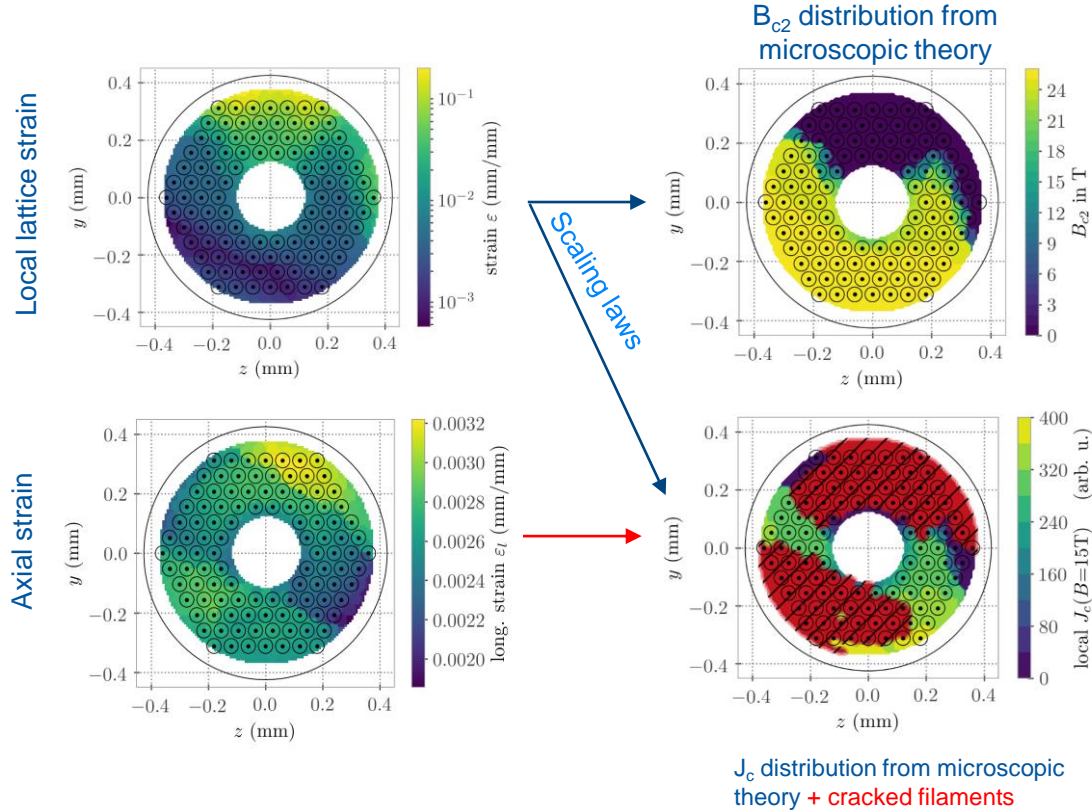
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# Conclusion and further actions

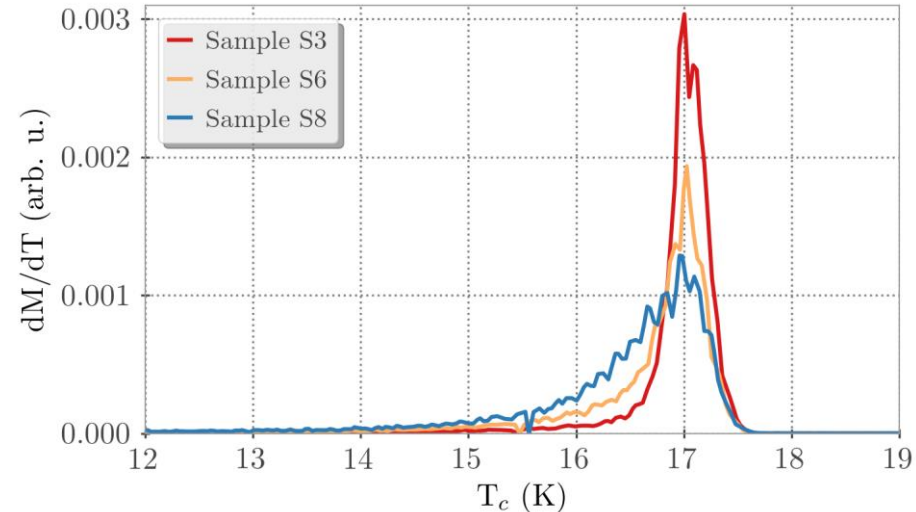
$$T_c(\varepsilon) = T_{c,0} \cdot s(\varepsilon)^{\frac{1}{3}}$$
$$B_{c2}(T, \varepsilon) = B_{c2}(T) \cdot s(\varepsilon)$$

$B_{c2}$  and  $T_c$  degradation dominated by lattice strain

- Measured  $T_c$  distribution:
  - Broadening of  $T_c$  distribution towards lower values as predicted: with higher exposed energy deposition, higher strains are expected.
  - Due to strain gradient, distribution is widened and no sharp peak
- Well reproduced with Simulations!

$I_c$  degradation is a combination of lattice & axial strain, major part due to filament cracking

- Filament Cracking is difficult to predict due to nature of  $Nb_3Sn$
- Continue optical analysis with SEM!  
(as seen before)





# Thank you for your attention



A. Will, D. Wollmann, J. Schubert, A. Oslandsbotn, A. Monteuis, K. Stachon, A.S. Mueller, A. Bernhard, M. Bonura, C. Senatore, M. Mentink, A. Verweij, B. Bordini

## Thanks to all Collaborators

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# Backup Slide I

## Experimental challenges

$I_c$  measurements on short wires are challenging for highly inhomogeneous  $J_c$  wires  
very low  $n$ -values  
Sample holders add a lot of copper  
Current sharing

