

The inclusive Higgs width in the SMEFT

Ilaria Brivio

Institut für Theoretische Physik, Universität Heidelberg

based on 1906.06949
in collaboration with T. Corbett and M. Trott



The Niels Bohr
International Academy



SMEFT = Effective Field Theory with **SM fields** + **symmetries**

a Taylor expansion in canonical dimensions ($v, E/\Lambda$):

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \frac{1}{\Lambda^3} \mathcal{L}_7 + \frac{1}{\Lambda^4} \mathcal{L}_8 + \dots$$


$$\mathcal{L}_n = \sum_i C_i \mathcal{O}_i^{d=n}$$

C_i free parameters (**Wilson coefficients**)

\mathcal{O}_i invariant **operators** that form
a complete basis

 describes impact of **any** UV ▶ nearly decoupled: $\Lambda \gg v, E$

▶ converges to SM in the limit $E \ll \Lambda$

 not just a parameterization but **a consistent QFT** \rightarrow loops, RGE etc.

An important observable: the Higgs width

a crucial observable for the Higgs sector

SM: $\Gamma_H \simeq 4 \text{ MeV} \rightarrow \frac{\Gamma_H}{m_H} \simeq 3 \cdot 10^{-5} \ll 1$

\Rightarrow Higgs measurements can be factored into

$$\sigma(i \rightarrow H) \times Br(H \rightarrow f)$$


An important observable: the Higgs width

a crucial observable for the Higgs sector

SM: $\Gamma_H \simeq 4 \text{ MeV} \rightarrow \frac{\Gamma_H}{m_H} \simeq 3 \cdot 10^{-5} \ll 1$

⇒ Higgs measurements can be factored into

$$\sigma(i \rightarrow H) \times Br(H \rightarrow f)$$

SMEFT: probe separately production and decay 

$$Br(H \rightarrow f) = \frac{\Gamma(H \rightarrow f)}{\Gamma_H^{\text{tot}}}$$


An important observable: the Higgs width

a crucial observable for the Higgs sector

$$\text{SM:} \quad \Gamma_H \simeq 4 \text{ MeV} \quad \rightarrow \quad \frac{\Gamma_H}{m_H} \simeq 3 \cdot 10^{-5} \ll 1$$

⇒ Higgs measurements can be factored into

$$\sigma(i \rightarrow H) \times Br(H \rightarrow f)$$

SMEFT: probe separately production and decay 

$$Br_{\text{SMEFT}}(H \rightarrow f) = \left[\frac{\Gamma(H \rightarrow f)}{\Gamma_H^{\text{tot}}} \right]_{\text{SM}} \left[1 + \frac{\delta\Gamma(H \rightarrow f)}{\Gamma_{\text{SM}}(H \rightarrow f)} - \frac{\delta\Gamma_H^{\text{tot}}}{\Gamma_{H,\text{SM}}^{\text{tot}}} \right]$$

- ▶ both $\delta\Gamma(H \rightarrow f)$ and $\delta\Gamma_H^{\text{tot}}$ need to be determined
- ▶ $\delta\Gamma_H^{\text{tot}}$ enters all processes → **strong impact** on global SMEFT analyses!

The Higgs width in the SMEFT - setup

Focus on the **leading** contributions:

- ▶ LO in the EFT: up to SM - \mathcal{L}_6 **interference**.

$$\Gamma_H = \Gamma_{H,\text{SM}} \left[1 + \frac{\delta\Gamma_H}{\Gamma_{H,\text{SM}}} \right] \quad \frac{\delta\Gamma_H}{\Gamma_{H,\text{SM}}} = \sum_i a_i \bar{C}_i = \sum_i a_i \left(C_i \frac{v^2}{\Lambda^2} \right)$$

- ▶ **tree level**.

SM couplings $H\gamma\gamma$, $HZ\gamma$, Hgg included as effective vertices for $H \rightarrow f\bar{f}\gamma / \gamma\gamma / gg$.
Neglected in other channels.

- ▶ up to **4-body** decays.

The Higgs width in the SMEFT - setup

Focus on the **leading** contributions:

- ▶ LO in the EFT: up to SM - \mathcal{L}_6 **interference**.

$$\Gamma_H = \Gamma_{H,SM} \left[1 + \frac{\delta\Gamma_H}{\Gamma_{H,SM}} \right] \quad \frac{\delta\Gamma_H}{\Gamma_{H,SM}} = \sum_i a_i \bar{C}_i = \sum_i a_i \left(C_i \frac{v^2}{\Lambda^2} \right)$$

- ▶ **tree level**.

SM couplings $H\gamma\gamma$, $HZ\gamma$, Hgg included as effective vertices for $H \rightarrow f\bar{f}\gamma / \gamma\gamma / gg$.
Neglected in other channels.

- ▶ up to **4-body** decays.

Conventions and assumptions:

- ▶ Warsaw basis

Grzadkowski, Iskrzynski, Misiak, Rosiek 1008.4884

- ▶ **U(3)⁵ flavor symmetry** $\rightarrow g_{Hff} \sim y_f$ also in the SMEFT

- ▶ **inclusive** calculation, not differential \rightarrow CP odd terms do not contribute

- ▶ **2 input schemes**: $\{m_W, m_Z, G_F\}$, $\{\alpha_{em}, m_Z, G_F\}$

 [backup](#)

The Higgs width in the SMEFT

leading channels:

$$H \rightarrow \bar{f}f$$

$$H \rightarrow gg$$

$$H \rightarrow \gamma\gamma$$

$$H \rightarrow \bar{f}f\gamma$$

$$H \rightarrow 4f$$

- ▶ $4f, \gamma\gamma, \bar{b}b$ most relevant ones individually
- ▶ **all** need to be calculated for $\delta\Gamma_H^{\text{tot}} \rightarrow Br$

The Higgs width in the SMEFT

leading channels:

$$H \rightarrow \bar{f}f$$

$$H \rightarrow gg$$

$$H \rightarrow \gamma\gamma$$

$$H \rightarrow \bar{f}f\gamma$$

$$H \rightarrow 4f$$

$$\frac{\Gamma(H \rightarrow \bar{f}f)}{\Gamma_{SM}(H \rightarrow \bar{f}f)} \simeq 1 + 2\delta g_{Hff}$$

$$\delta g_{Hff} = \bar{C}_{H\Box} - \frac{\bar{C}_{HD}}{4} - \bar{C}_{HI}^{(3)} + \frac{\bar{C}'_{II}}{2} - \frac{5}{2}\bar{C}_{fH}$$

Alonso, Jenkins, Manohar, Trott 1312.2014

- ▶ renorm. of the H field
- ▶ contrib. to μ decay $\rightarrow G_F \rightarrow v$
- ▶ direct \mathcal{O}_{fH} contribution $\left(-\frac{3}{2}\bar{C}_{fH}\right)$
+
contrib. to $m_f \rightarrow y_f$ $\left(-\bar{C}_{fH}\right)$

- ▶ $4f, \gamma\gamma, \bar{b}b$ most relevant ones individually
- ▶ **all** need to be calculated for $\delta\Gamma_H^{\text{tot}} \rightarrow Br$

The Higgs width in the SMEFT

leading channels:

$$H \rightarrow \bar{f}f$$

$$H \rightarrow gg$$

$$H \rightarrow \gamma\gamma$$

$$H \rightarrow \bar{f}f\gamma$$

$$H \rightarrow 4f$$

$$\frac{\Gamma(h \rightarrow gg)}{\Gamma^{SM}(h \rightarrow gg)} \simeq 1 + \frac{16\pi^2}{g_s^2 I^g} \bar{C}_{HG}, \quad I^g \simeq 0.375$$

Manohar,Wise 0601212

- ▶ $4f, \gamma\gamma, \bar{b}b$ most relevant ones individually
- ▶ **all** need to be calculated for $\delta\Gamma_H^{\text{tot}} \rightarrow Br$

The Higgs width in the SMEFT

leading channels:

$$H \rightarrow \bar{f}f$$

$$H \rightarrow gg$$

$$H \rightarrow \gamma\gamma$$

$$H \rightarrow \bar{f}f\gamma$$

$$H \rightarrow 4f$$

$$\frac{\Gamma(h \rightarrow \gamma\gamma)}{\Gamma^{SM}(h \rightarrow \gamma\gamma)} \simeq 1 + \frac{16\pi^2}{e^2 I^\gamma} \mathcal{C}_{\gamma\gamma}, \quad I^\gamma \simeq -1.65$$

$$\mathcal{C}_{\gamma\gamma} = s_\theta^2 \bar{C}_{HW} + c_\theta^2 \bar{C}_{HB} - s_\theta c_\theta \bar{C}_{HWB}$$

Bergström, Hulth Nucl.Phys.B259(1985)137
Manohar, Wise 0601212

- ▶ $4f, \gamma\gamma, \bar{b}b$ most relevant ones individually
- ▶ **all** need to be calculated for $\delta\Gamma_H^{\text{tot}} \rightarrow Br$

The Higgs width in the SMEFT

leading channels:

$$H \rightarrow \bar{f}f$$

→ $H \rightarrow gg$

→ $H \rightarrow \gamma\gamma$

→ $H \rightarrow \bar{f}f\gamma$

$$H \rightarrow 4f$$

$$\frac{\Gamma(h \rightarrow \gamma\gamma)}{\Gamma^{SM}(h \rightarrow \gamma\gamma)} \simeq 1 + \frac{16\pi^2}{e^2 I_\gamma} \mathcal{C}_{\gamma\gamma}, \quad I_\gamma \simeq -1.65$$

$$\mathcal{C}_{\gamma\gamma} = s_\theta^2 \bar{C}_{HW} + c_\theta^2 \bar{C}_{HB} - s_\theta c_\theta \bar{C}_{HWB}$$

Bergström, Hulth Nucl.Phys.B259(1985)137
Manohar, Wise 0601212

loop-factor enhancement
in the relative correction:

tree-level SMEFT vs loop SM

- ▶ $4f, \gamma\gamma, \bar{b}b$ most relevant ones individually
- ▶ **all** need to be calculated for $\delta\Gamma_H^{\text{tot}} \rightarrow Br$

The Higgs width in the SMEFT

leading channels:

$$H \rightarrow \bar{f}f$$

$$H \rightarrow gg$$

$$H \rightarrow \gamma\gamma$$

$$H \rightarrow \bar{f}f\gamma$$

$$H \rightarrow 4f$$

available as $H \rightarrow ZZ^*$, $H \rightarrow WW^*$ $\times Br(Z, W)$
relying on narrow width approx. for Z, W .

good in SM but **not sufficient** in the SMEFT!

main reason: tree $\gamma\gamma, Z\gamma$ mediated diagrams

also missing:

- ▶ CC - NC interference
- ▶ crossed-current interference in ZZ diagrams
- ▶ $\delta\Gamma_V, \delta m_V^2$ corrections for off-shell boson

- ▶ $4f, \gamma\gamma, \bar{b}b$ most relevant ones individually
- ▶ **all** need to be calculated for $\delta\Gamma_H^{\text{tot}} \rightarrow Br$

The Higgs width in the SMEFT

leading channels:

$$H \rightarrow \bar{f}f$$

$$H \rightarrow gg$$

$$H \rightarrow \gamma\gamma$$

$$H \rightarrow \bar{f}f\gamma$$

$$H \rightarrow 4f$$

available as $H \rightarrow ZZ^*$, $H \rightarrow WW^*$ $\times Br(Z, W)$
relying on narrow width approx. for Z, W .

good in SM but **not sufficient** in the SMEFT!

main reason: tree $\gamma\gamma, Z\gamma$ mediated diagrams

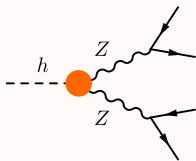
also missing:

- ▶ CC - NC interference
- ▶ crossed-current interference in ZZ diagrams
- ▶ $\delta\Gamma_V, \delta m_V^2$ corrections for off-shell boson

- ▶ $4f, \gamma\gamma, \bar{b}b$ most relevant ones individually
- ▶ **all** need to be calculated for $\delta\Gamma_H^{\text{tot}} \rightarrow Br$

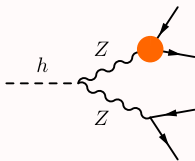
H → 4f in the SMEFT

① corrections to SM diagrams

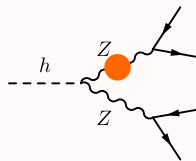


$$\propto g_{\mu\nu} \text{ (SM-like)}$$

$$\propto g_{\mu\nu} p \cdot q - p_\nu q_\mu \text{ (} Z_{\mu\nu} Z^{\mu\nu} h \text{)}$$



$$\delta g_L, \delta g_R$$



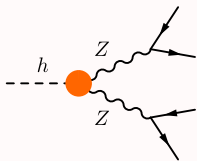
$$\frac{-im_Z \delta \Gamma_Z + (2m_Z - i\Gamma_Z) \delta m_Z}{p^2 - m_Z^2 + i\Gamma_Z m_Z}$$



hard to extract from
MC simulation!
full treatment requires
analytic calculation

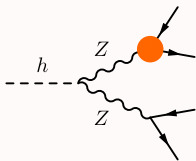
H \rightarrow 4f in the SMEFT

① corrections to SM diagrams

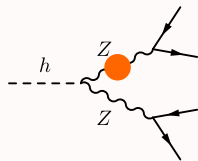


$$\propto g_{\mu\nu} \text{ (SM-like)}$$

$$\propto g_{\mu\nu} p \cdot q - p_\nu q_\mu \text{ (} Z_{\mu\nu} Z^{\mu\nu} h \text{)}$$

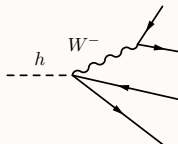
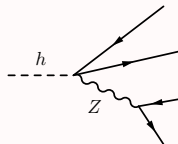
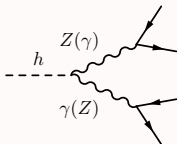
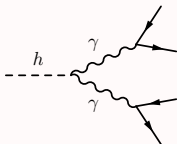


$$\delta g_L, \delta g_R$$



$$\frac{-im_Z \delta \Gamma_Z + (2m_Z - i\Gamma_Z) \delta m_Z}{p^2 - m_Z^2 + i\Gamma_Z m_Z}$$

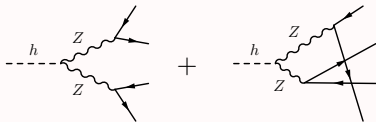
② genuine SMEFT diagrams



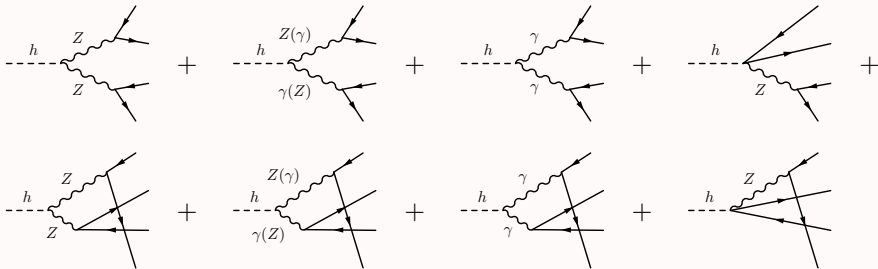
H \rightarrow 4f in the SMEFT - complexity

$$h \rightarrow e^+ e^- e^+ e^-$$

SM



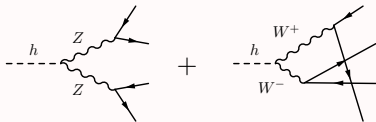
interfering with



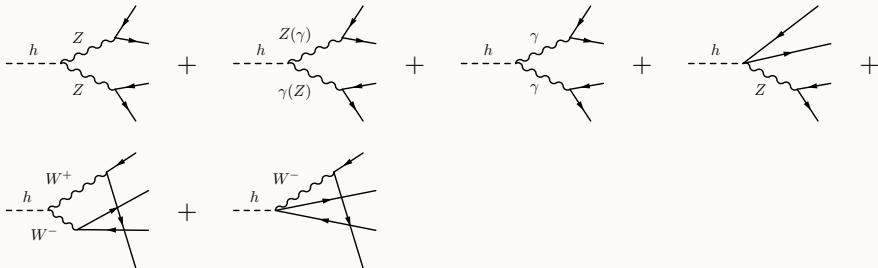
H \rightarrow 4f in the SMEFT - complexity

$$h \rightarrow \bar{u} u \bar{d} d$$

SM

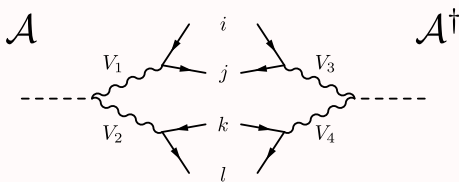


interfering with



H \rightarrow 4f - analytic calculation

fully analytical treatment. automated with general decomposition:



$$\mathcal{A}A^\dagger \sim g_{HV_1V_2} g_{HV_3V_4} \sum_n \mathcal{T}^{(n)}$$

$$\mathcal{T}^{(n)} = \mathcal{K}^{(n)} \left(g_{L,R}^{ij,V_1}, g_{L,R}^{ij,V_3}, g_{L,R}^{kl,V_2}, g_{L,R}^{kl,V_4} \right) \mathcal{F}_{V_1V_2V_3V_4}^{(n)}(p_a, m_a), \quad a = \{i, j, k, l\}$$

for $m_a \equiv 0$ there are only **8** independent $\mathcal{F}_{V_1V_2V_3V_4}$. For each $\{V\}$ set:

- ▶ numerical integration of phase space: **Vegas** in Mathematica T. Hahn 0404043
- ▶ cross-check: RAMBO + 2 independent parameterizations of phase space

Kleiss, Stirling, Ellis
Comput. Phys. Commun. 40(1986)359

H → 4f - results

Example: $H \rightarrow e^+ e^- \mu^+ \mu^-$ $m_f \equiv 0$, m_W scheme

$$\frac{\delta\Gamma(H \rightarrow e^+ e^- \mu^+ \mu^-)}{\Gamma_{\text{SM}}(H \rightarrow e^+ e^- \mu^+ \mu^-)} = \sum_i a_i \bar{C}_i = \sum_i a_i \left(C_i \frac{v^2}{\Lambda^2} \right)$$

	\bar{C}_{HW}	\bar{C}_{HB}	\bar{C}_{HWB}	$\bar{C}_{H\Box}$	\bar{C}_{HD}	$\bar{C}_{HI}^{(1)}$	$\bar{C}_{HI}^{(3)}$	\bar{C}_{He}	$\bar{C}_{Hq}^{(1)}$	$\bar{C}_{Hq}^{(3)}$	\bar{C}_{Hu}	\bar{C}_{Hd}	\bar{C}'_{ll}
Z	-0.78	-0.22	0.30	2	0.17	4.38	-1.62	-3.52					3.
A	1.04	-1.08	-0.68										
E						-2.23	-2.23	1.80					
G			-0.38		0.06	0.15	1.14	0.15	-0.39	-1.34	-0.20	0.15	-0.83
tot	0.26	-1.30	-0.76	2.	0.23	2.30	-2.71	-1.58	-0.39	-1.34	-0.20	0.15	2.17

Z	corrections to SM diagram
A	γ diagrams
E	contact diagrams ($HZee$)
G	$\delta\Gamma_Z^{\text{tot}}/\Gamma_{Z,SM}$ on + off-shell Z

H → 4f - results

Example: $H \rightarrow e^+ e^- \mu^+ \mu^-$ $m_f \equiv 0$, m_W scheme

$$\frac{\delta\Gamma(H \rightarrow e^+ e^- \mu^+ \mu^-)}{\Gamma_{\text{SM}}(H \rightarrow e^+ e^- \mu^+ \mu^-)} = \sum_i a_i \bar{C}_i = \sum_i a_i \left(C_i \frac{v^2}{\Lambda^2} \right)$$

	\bar{C}_{HW}	\bar{C}_{HB}	\bar{C}_{HWB}	$\bar{C}_{H\Box}$	\bar{C}_{HD}	$\bar{C}_{HI}^{(1)}$	$\bar{C}_{HI}^{(3)}$	\bar{C}_{He}	$\bar{C}_{Hq}^{(1)}$	$\bar{C}_{Hq}^{(3)}$	\bar{C}_{Hu}	\bar{C}_{Hd}	\bar{C}'_{ll}
Z	-0.78	-0.22	0.30	2	0.17	4.38	-1.62	-3.52					3.
A	1.04	-1.08	-0.68										
E						-2.23	-2.23	1.80					
G			-0.38		0.06	0.15	1.14	0.15	-0.39	-1.34	-0.20	0.15	-0.83
tot	0.26	-1.30	-0.76	2.	0.23	2.30	-2.71	-1.58	-0.39	-1.34	-0.20	0.15	2.17

Z	corrections to SM diagram
A	γ diagrams
E	contact diagrams ($HZee$)
G	$\delta\Gamma_Z^{\text{tot}}/\Gamma_{Z,SM}$ on + off-shell Z

Impact of previously neglected contributions

(1) photon-mediated diagrams

$\mathcal{O}(1 - 250)\%$ effect

	with γ			without γ		
	\bar{C}_{HW}	\bar{C}_{HB}	\bar{C}_{HWB}	\bar{C}_{HW}	\bar{C}_{HB}	\bar{C}_{HWB}
$h \rightarrow e^+ e^- \mu^+ \mu^-$	0.26	-1.30	-0.38	-0.77	-0.22	0.30
$h \rightarrow \bar{u} u \bar{c} c$	1.45	-2.63	-0.29	-0.77	-0.22	1.33
$h \rightarrow e^+ e^- \bar{d} d$	0.50	-1.55	-0.37	-0.77	-0.22	0.47
$h \rightarrow e^+ e^- e^+ e^-$	0.02	-2.28	0.27	-0.76	-0.21	0.44
$h \rightarrow \bar{u} u \bar{u} u$	1.39	-2.72	-0.14	-0.76	-0.21	1.19
$h \rightarrow e^+ e^- \bar{\nu}_e \nu_e$	-1.49	0.01	-0.06	-1.48	-0.007	-0.07

Impact of previously neglected contributions

(2) $Z - W$ interference terms

$\mathcal{O}(1 - 200)\%$ effect

$\delta\Gamma(H \rightarrow e^+ e^- \bar{\nu}_e \nu_e)/\Gamma_{\text{SM}}$ omitting γ and $\delta\Gamma_Z, \delta\Gamma_W$ contrib.

	\bar{C}_{HW}	\bar{C}_{HB}	\bar{C}_{HWB}	$\bar{C}_{H\Box}$	\bar{C}_{HD}	$\bar{C}_{HI}^{(1)}$	$\bar{C}_{HI}^{(3)}$	\bar{C}_{He}	$\bar{C}_{Hg}^{(1)}$	$\bar{C}_{Hg}^{(3)}$	\bar{C}_{Hu}	\bar{C}_{Hd}	\bar{C}'_{ll}
ZZ	-0.04	-0.01	-0.003	0.09	-0.008	0.004	-0.19	-0.04	0	0	0	0	0.14
WW	-1.49	0	0	2.0	-0.50	0	-3.77	0	0	0	0	0	3.00
WZ	0.04	0.004	-0.06	-0.10	-0.04	-0.01	0.21	0	0	0	0	0	-0.14
full	-1.49	-0.007	-0.07	2.	-0.55	-0.008	-3.74	-0.04	0	0	0	0	3.
NW	-1.46	-0.01	-0.003	2.	-0.49	0.004	-3.77	-0.04	0.	0.	0.	0.	3.

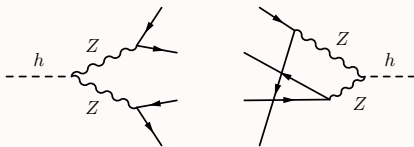
full | ZZ + WW + WZ

NW | ZZ + WW

Impact of previously neglected contributions

(3) NC crossed - interference terms

$\mathcal{O}(\text{few} - 40)\%$ effect



$\delta\Gamma(H \rightarrow e^+e^-e^+e^-)/\Gamma_{\text{SM}}$ incl. only ZZ and HZee diagrams

	\bar{C}_{HW}	\bar{C}_{HB}	\bar{C}_{HWB}	$\bar{C}_{H\Box}$	\bar{C}_{HD}	$\bar{C}_{HI}^{(1)}$	$\bar{C}_{HI}^{(3)}$	\bar{C}_{He}	$\bar{C}_{Hq}^{(1)}$	$\bar{C}_{Hq}^{(3)}$	\bar{C}_{Hu}	\bar{C}_{Hd}	\bar{C}'_{ll}
full	-0.75	-0.22	0.43	2.	0.28	2.09	-3.91	-1.64	0	0	0	0	3.
NW	-0.78	-0.221	0.30	2.	0.17	2.15	-3.85	-1.73	0.	0.	0.	0.	3.

$$\begin{array}{l} \text{full} \\ \text{NW} \end{array} \left| \begin{array}{l} |\mathcal{A}_{ijkl}|^2 + |\mathcal{A}_{ilkj}|^2 + 2\text{Re}\mathcal{A}_{ijkl}\mathcal{A}_{ilkj}^\dagger \\ |\mathcal{A}_{ijkl}|^2 + |\mathcal{A}_{ilkj}|^2 \end{array} \right.$$

Impact of previously neglected contributions

(4) $\delta\Gamma_V$, δm_V from off-shell boson

$\mathcal{O}(\text{few})\%$ effect

narrow width approx.:
$$\frac{\delta\Gamma(H \rightarrow VV^* \rightarrow 4f)}{\Gamma_{SM}(H \rightarrow VV^* \rightarrow 4f)} = (-1) \frac{\delta\Gamma_V}{\Gamma_{V,SM}} + \dots$$

full calculation:

$h \rightarrow e^+ e^- \mu^+ \mu^-$	-0.820	$\delta\Gamma_Z/\Gamma_{Z,SM}$	
$h \rightarrow e^+ e^- e^+ e^-$	-0.748	$\delta\Gamma_Z/\Gamma_{Z,SM}$	
$h \rightarrow e^+ \nu_e \bar{\nu}_\mu \mu^-$	-0.915	$\delta\Gamma_W/\Gamma_{W,SM}$	
$h \rightarrow e^+ \nu_e \bar{\nu}_e e^-$	-0.914	$\delta\Gamma_W/\Gamma_{W,SM}$	$- 0.038 \delta\Gamma_Z/\Gamma_{Z,SM}$

H → 4f summary

- ▶ we did a **fully analytic** calculation, with numerical integration of phase space
- ▶ also generated all channels with MG5_aMC@NLO using **SMEFTsim**
→ agreement to 1% or better ✓
Brivio, Jiang, Trott 1709.06492
- ▶ analytic treatment has a few advantages:
 - ▶ allows to separate contributions
 - ▶ easier to **linearize in $\delta\Gamma_V, \delta m_V$**
 - ▶ more stable for the massless fermions case with γ diagrams
 - ▶ cancellations are reproduced exactly
 - ▶ calculation can be **automated** in a dedicated package
↪ a new reweighting tool coming soon
- ▶ some previously neglected contributions turn out to be relevant:
 γ diagrams and **$Z - W$ interference**

The total Higgs width in the SMEFT

putting together all the main contributions* we obtain

$$\Gamma_H^{\text{tot}} = \Gamma_{H,SM}^{\text{tot}} \left[1 + \frac{\delta\Gamma_H^{\text{tot}}}{\Gamma_{H,SM}^{\text{tot}}} \right]$$

$$\Gamma_{H,SM}^{\text{tot}} = 4.100 \text{ MeV}$$

$$\begin{aligned} \frac{\delta\Gamma_H^{\text{tot}}}{\Gamma_{H,SM}^{\text{tot}}} = & -1.50 \tilde{C}_{HB} - 1.21 \tilde{C}_{HW} + 1.21 \tilde{C}_{HWB} + 50.6 \tilde{C}_{HG} \\ & + 1.83 \tilde{C}_{H\Box} - 0.43 \tilde{C}_{HD} + 1.17 \tilde{C}'_{II} \\ & - 7.85 Y_c \Re \tilde{C}_{uH} - 48.5 Y_b \Re \tilde{C}_{dH} - 12.3 Y_\tau \Re \tilde{C}_{eH} \\ & + 0.002 \tilde{C}_{Hq}^{(1)} + 0.06 \tilde{C}_{Hq}^{(3)} + 0.001 \tilde{C}_{Hu} - 0.0007 \tilde{C}_{Hd} \\ & - 0.0009 \tilde{C}_{HI}^{(1)} - 2.32 \tilde{C}_{HI}^{(3)} - 0.0006 \tilde{C}_{He}, \end{aligned}$$

* $gg + \gamma\gamma + \bar{b}b + \bar{c}c + \tau^+\tau^- + 4f + \bar{f}f\gamma$

- ▶ the inclusive Higgs width is a crucial observable for the Higgs sector
- ▶ improved calculation of $H \rightarrow 4f, H \rightarrow \bar{f}f\gamma$ without relying on the narrow width approx. for Z, W
→ important for LHC measurements
- ▶ joined all the main channels into $\delta\Gamma_H^{\text{tot}}$
- ▶ Br's and Γ_H^{tot} are computed **once and for all!**
→ can be used directly without further MC generation
- ▶ an automated package for the reweighting to appear soon!
- ▶ possible refinements:
full massive fermions treatment
phase space integration with cuts
...

Backup slides

X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_{\mu}^{A\nu} G_{\nu}^{B\rho} G_{\rho}^{C\mu}$	Q_{φ}	$(\varphi^{\dagger}\varphi)^3$	$Q_{e\varphi}$	$(\varphi^{\dagger}\varphi)(\bar{l}_p e_r \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu}^{A\nu} G_{\nu}^{B\rho} G_{\rho}^{C\mu}$	$Q_{\varphi\Box}$	$(\varphi^{\dagger}\varphi)\Box(\varphi^{\dagger}\varphi)$	$Q_{u\varphi}$	$(\varphi^{\dagger}\varphi)(\bar{q}_p u_r \tilde{\varphi})$
Q_W	$\varepsilon^{IJK} W_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$	$Q_{\varphi D}$	$(\varphi^{\dagger} D^{\mu}\varphi)^{\star} (\varphi^{\dagger} D_{\mu}\varphi)$	$Q_{d\varphi}$	$(\varphi^{\dagger}\varphi)(\bar{q}_p d_r \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^{\dagger}\varphi G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{l}_p \gamma^{\mu} l_r)$
$Q_{\varphi \tilde{G}}$	$\varphi^{\dagger}\varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu}^I \varphi) (\bar{l}_p \tau^I \gamma^{\mu} l_r)$
$Q_{\varphi W}$	$\varphi^{\dagger}\varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{e}_p \gamma^{\mu} e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^{\dagger}\varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{q}_p \gamma^{\mu} q_r)$
$Q_{\varphi B}$	$\varphi^{\dagger}\varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu}^I \varphi) (\bar{q}_p \tau^I \gamma^{\mu} q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^{\dagger}\varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{u}_p \gamma^{\mu} u_r)$
$Q_{\varphi WB}$	$\varphi^{\dagger} \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{d}_p \gamma^{\mu} d_r)$
$Q_{\varphi \tilde{W}B}$	$\varphi^{\dagger} \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^{\dagger} D_{\mu} \varphi) (\bar{u}_p \gamma^{\mu} d_r)$

$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		B -violating			
Q_{ledq}	$(\bar{l}_p^j e_r)(\bar{d}_s^j q_t^j)$	Q_{duq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^\gamma)^T C l_t^k]$		
$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	Q_{qqqu}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	Q_{qqq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jkmn} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^m)^T C l_t^n]$		
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	Q_{duu}	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				

Shifts from input parameters

SM case.

Parameters in the canonically normalized Lagrangian : $\bar{v}, \bar{g}_1, \bar{g}_2, s_{\bar{\theta}}$

The values can be inferred from the measurements e.g. of $\{\alpha_{\text{em}}, m_Z, G_f\}$:

$$\alpha_{\text{em}} = \frac{\bar{g}_1 \bar{g}_2}{\bar{g}_1^2 + \bar{g}_2^2}$$

$$m_Z = \frac{\bar{g}_2 \bar{v}}{2c_{\bar{\theta}}}$$

$$G_f = \frac{1}{\sqrt{2}\bar{v}^2}$$

→

$$\hat{v}^2 = \frac{1}{\sqrt{2}G_f}$$

$$\sin \hat{\theta}^2 = \frac{1}{2} \left(1 - \sqrt{1 - \frac{4\pi\alpha_{\text{em}}}{\sqrt{2}G_f m_Z^2}} \right)$$

$$\hat{g}_1 = \frac{\sqrt{4\pi\alpha_{\text{em}}}}{\cos \hat{\theta}}$$

$$\hat{g}_2 = \frac{\sqrt{4\pi\alpha_{\text{em}}}}{\sin \hat{\theta}}$$

in the SM at tree-level $\bar{\kappa} = \hat{\kappa}$

Shifts from input parameters

SMEFT case.

Parameters in the canonically normalized Lagrangian : $\bar{v}, \bar{g}_1, \bar{g}_2, s_{\bar{\theta}}$

The values can be inferred from the measurements e.g. of $\{\alpha_{\text{em}}, m_Z, G_f\}$:

$$\begin{aligned}\alpha_{\text{em}} &= \frac{\bar{g}_1 \bar{g}_2}{\bar{g}_1^2 + \bar{g}_2^2} \left[1 + \bar{v}^2 C_{HWB} \frac{\bar{g}_2^3 / \bar{g}_1}{\bar{g}_1^2 + \bar{g}_2^2} \right] & \hat{v}^2 &= \frac{1}{\sqrt{2} G_f} \\ m_Z &= \frac{\bar{g}_2 \bar{v}}{2 c_{\bar{\theta}}} + \delta m_Z(C_i) & \sin \hat{\theta}^2 &= \frac{1}{2} \left(1 - \sqrt{1 - \frac{4\pi \alpha_{\text{em}}}{\sqrt{2} G_f m_Z^2}} \right) \\ G_f &= \frac{1}{\sqrt{2} \bar{v}^2} + \delta G_f(C_i) & \hat{g}_1 &= \frac{\sqrt{4\pi \alpha_{\text{em}}}}{\cos \hat{\theta}} \\ & & \hat{g}_2 &= \frac{\sqrt{4\pi \alpha_{\text{em}}}}{\sin \hat{\theta}}\end{aligned}$$

in the SM at tree-level $\bar{\kappa} = \hat{\kappa}$

in the SMEFT $\bar{\kappa} = \hat{\kappa} + \delta \kappa(C_i)$

Shifts from input parameters

To have numerical predictions it is necessary to replace $\bar{\kappa} \rightarrow \hat{\kappa} + \delta\kappa(C_i)$ for all the parameters in the Lagrangian.

$\{\alpha_{\text{em}}, m_Z, G_f\}$ scheme

$$\delta m_Z^2 = m_Z^2 \hat{v}^2 \left(\frac{c_{HD}}{2} + 2c_{\hat{\theta}} s_{\hat{\theta}} c_{HWB} \right)$$

$$\delta G_f = \frac{\hat{v}^2}{\sqrt{2}} \left((c_{HI}^{(3)})_{11} + (c_{HI}^{(3)})_{22} - (c_{II})_{1221} \right)$$

$$\delta g_1 = \frac{s_{\hat{\theta}}^2}{2(1-2s_{\hat{\theta}}^2)} \left(\sqrt{2}\delta G_f + \delta m_Z^2/m_Z^2 + 2\frac{c_{\hat{\theta}}^3}{s_{\hat{\theta}}} c_{HWB} \hat{v}^2 \right)$$

$$\delta g_2 = -\frac{c_{\hat{\theta}}^2}{2(1-2s_{\hat{\theta}}^2)} \left(\sqrt{2}\delta G_f + \delta m_Z^2/m_Z^2 + 2\frac{s_{\hat{\theta}}^3}{c_{\hat{\theta}}} c_{HWB} \hat{v}^2 \right)$$

$$\delta s_{\hat{\theta}}^2 = 2c_{\hat{\theta}}^2 s_{\hat{\theta}}^2 (\delta g_1 - \delta g_2) + c_{\hat{\theta}} s_{\hat{\theta}} (1 - 2s_{\hat{\theta}}^2) c_{HWB} \hat{v}^2$$

$$\delta m_h^2 = m_h^2 \hat{v}^2 \left(2c_{H\Box} - \frac{c_{HD}}{2} - \frac{3c_H}{2\lambda m} \right)$$

Shifts from input parameters

To have numerical predictions it is necessary to replace $\bar{\kappa} \rightarrow \hat{\kappa} + \delta\kappa(C_i)$
for all the parameters in the Lagrangian.

$\{m_W, m_Z, G_f\}$ scheme

$$\delta m_Z^2 = m_Z^2 \hat{v}^2 \left(\frac{c_{HD}}{2} + 2c_{\hat{\theta}} s_{\hat{\theta}} c_{HWB} \right)$$

$$\delta G_f = \frac{\hat{v}^2}{\sqrt{2}} \left((c_{HI}^{(3)})_{11} + (c_{HI}^{(3)})_{22} - (c_{II})_{1221} \right)$$

$$\delta g_1 = -\frac{1}{2} \left(\sqrt{2} \delta G_f + \frac{1}{s_{\hat{\theta}}^2} \frac{\delta m_Z^2}{m_Z^2} \right)$$

$$\delta g_2 = -\frac{1}{\sqrt{2}} \delta G_f$$

$$\delta s_{\hat{\theta}}^2 = 2c_{\hat{\theta}}^2 s_{\hat{\theta}}^2 (\delta g_1 - \delta g_2) + c_{\hat{\theta}} s_{\hat{\theta}} (1 - 2s_{\hat{\theta}}^2) c_{HWB} \hat{v}^2$$

$$\delta m_h^2 = m_h^2 \hat{v}^2 \left(2c_{CH\Box} - \frac{c_{HD}}{2} - \frac{3c_{CH}}{2\lambda m} \right)$$

The SMEFTsim package

an UFO & FeynRules model with*:

Brivio, Jiang, Trott 1709.06492
feynrules.irmp.ucl.ac.be/wiki/SMEFT

1. the complete B-conserving Warsaw basis for 3 generations , including all complex phases and ~~CP~~ terms
2. automatic field redefinitions to have **canonical kinetic terms**
3. automatic **parameter shifts** due to the choice of an input parameters set

↪ backup

Main scope:

estimate **tree-level** $|\mathcal{A}_{\text{SM}}\mathcal{A}_{\text{d=6}}^*|$ **interference** terms \rightarrow theo. accuracy $\gtrsim 1\%$

* at the moment only LO, unitary gauge implementation

The SMEFTsim package

6 different implementations available

Brivio, Jiang, Trott 1709.06492

$$\textcircled{3} \text{ flavor structures } \left\{ \begin{array}{l} \text{general} \\ U(3)^5 \text{ symmetric} \\ \text{linear MFV} \end{array} \right. \times \textcircled{2} \text{ input schemes } \left\{ \begin{array}{l} \hat{\alpha}_{em}, \hat{m}_Z, \hat{G}_f \\ \hat{m}_W, \hat{m}_Z, \hat{G}_f \end{array} \right.$$

feynrules.irmp.ucl.ac.be/wiki/SMEFT

[viki: SMEFT](#)
Standard Model Effective Field Theory -- The SMEFTsim package
Authors
 Ilaria Brivio, Yun Jiang and Michael Trott
ilaria.brivio@nbi.ku.dk, yunjiang@nbi.ku.dk, michael.trott@cern.ch
 NBIA and Discovery Center, Niels Bohr Institute, University of Copenhagen

Pre-exported UFO files (include restriction cards)

	Set A		Set B	
	α scheme	m_W scheme	α scheme	m_W scheme
Flavor general SMEFT	SMEFTsim_A_general_alphaScheme_UFO.tar.gz ↓	SMEFTsim_A_general_MwScheme_UFO.tar.gz ↓	SMEFT_alpha_UFO.zip ↓	SMEFT_mW_UFO.zip ↓
MFV SMEFT	SMEFTsim_A_MFV_alphaScheme_UFO.tar.gz ↓	SMEFTsim_A_MFV_MwScheme_UFO.tar.gz ↓	SMEFT_alpha_MFV_UFO.zip ↓	SMEFT_mW_MFV_UFO.zip ↓
$U(3)^5$ SMEFT	SMEFTsim_A_U35_alphaScheme_UFO.tar.gz ↓	SMEFTsim_A_U35_MwScheme_UFO.tar.gz ↓	SMEFT_alpha_FLU_UFO.zip ↓	SMEFT_mW_FLU_UFO.zip ↓