

UNIVERSIDAD DE OVIEDO



WZ SM measurements and EWK SUSY searches

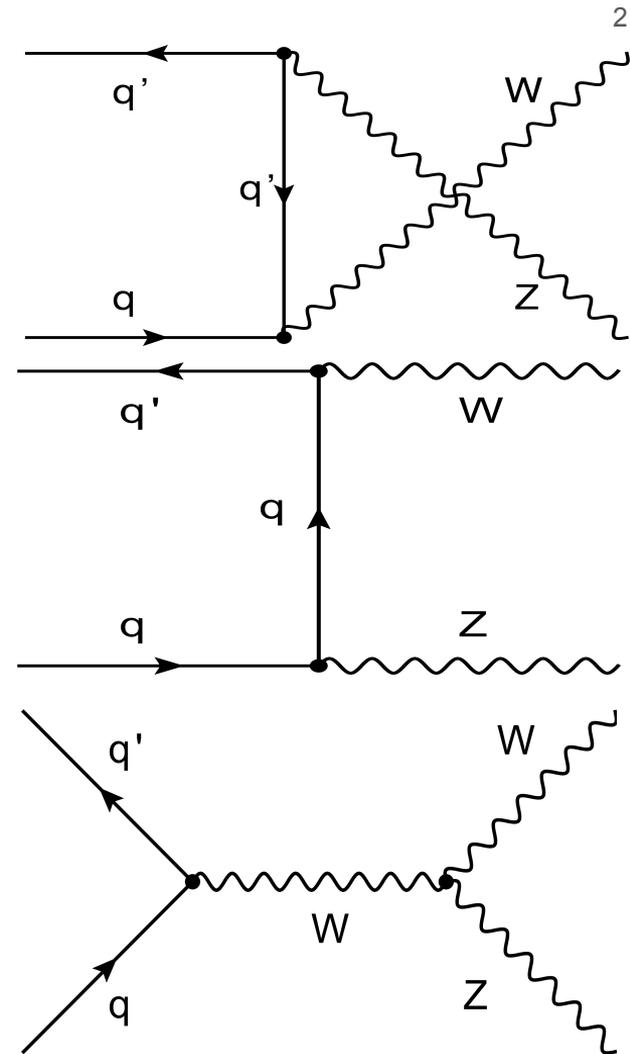
3rd Red LHC Workshop

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Why measure WZ?

- Dominant process at the **3-lepton final state**.
- Key process for our understanding of the SM:
 - Charged final state: charge asymmetry is induced in the final state by the ($\bar{u}d$ - $u\bar{d}$) pdf imbalance.
 - “Clean” final state: achieving very high ($\sim 85\%$) selection purities. Easier to check any SM deviation differentially (pointing to higher order corrections or BSM presence).
- And also for BSM searches:
 - As a signal: unique sensitivity to variations in the **WWZ** vertex coupling.
 - As a background: principal SM background in the 3-lepton final state. Different BSM models predict final states with very similar behavior (typically $WZ+p_T^{\text{miss}}$).



WZ Inclusive cross-section

→ Cross-section measured in the 3 lepton final state:

→ Use MVA-based lepton ID to reduce non-prompt contribution.

→ $p_T^{\text{miss}} > 30$ GeV (to decrease DY contamination).

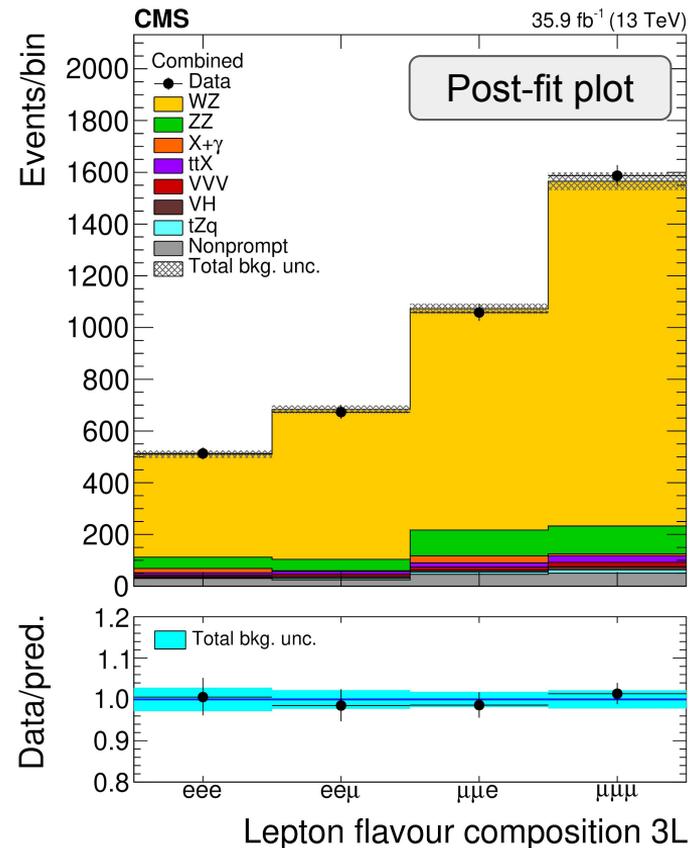
→ $N_{\text{B-tag}} = 0$ (to decrease tt and ttX contamination).

→ Final selection is extremely pure (~90% WZ).

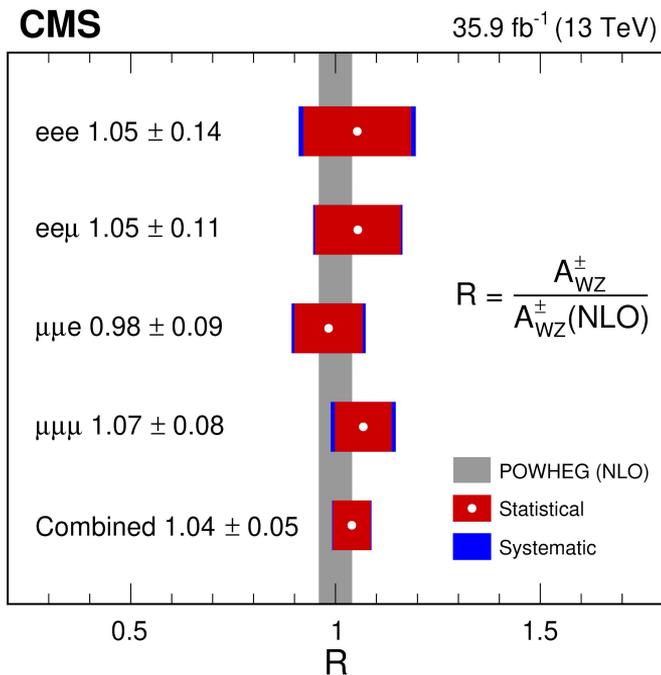
Category	$\sigma_{\text{tot}}(\text{pp} \rightarrow \text{WZ})$ [pb]
eee	$47.11^{+5.01}_{-4.63}$ (total) = $47.11^{+2.88}_{-2.79}$ (stat) $^{+0.46}_{-0.41}$ (theo) $^{+3.89}_{-3.47}$ (syst) ± 1.41 (lumi)
ee μ	$47.16^{+3.87}_{-3.61}$ (total) = $47.16^{+2.31}_{-2.29}$ (stat) $^{+0.45}_{-0.38}$ (theo) $^{+2.83}_{-2.52}$ (syst) ± 1.33 (lumi)
e $\mu\mu$	$47.70^{+3.58}_{-3.55}$ (total) = $47.70^{+2.00}_{-1.96}$ (stat) $^{+0.45}_{-0.39}$ (theo) $^{+2.66}_{-2.61}$ (syst) ± 1.42 (lumi)
$\mu\mu\mu$	$49.00^{+3.18}_{-3.03}$ (total) = $49.00^{+1.57}_{-1.53}$ (stat) $^{+0.41}_{-0.35}$ (theo) $^{+2.42}_{-2.22}$ (syst) ± 1.39 (lumi)

→ Biggest uncertainties arise from lepton and b-tag efficiencies (~2-3%).

→ Measurement is closer to MATRIX $\sigma_{\text{NNLO}} = 49.9 \pm 2.5$ pb than POWHEG $\sigma_{\text{NLO}} = 42.5 \pm 1.7$ pb predictions.



Charge asymmetry measurements



→ All flavor channels are split into +/- total charge.

→ Best-likelihood fit to the asymmetry parameter:

$$A_{WZ}^{+-} = \frac{\sigma_{\text{tot}}(\text{pp} \rightarrow W^+Z)}{\sigma_{\text{tot}}(\text{pp} \rightarrow W^-Z)} = 1.48 \pm 0.06 (\text{stat}) \pm 0.02 (\text{syst}) \pm 0.01 (\text{theo})$$

→ Results are normalized to the expected value.

→ Most uncertainties “cancel out” in the + and - terms, result is driven by statistical uncertainties.

→ Expect uncertainties to drop significantly (~50%) with Run-II luminosity.

[10.1007/JHEP04\(2019\)122](https://arxiv.org/abs/10.1007/JHEP04(2019)122)

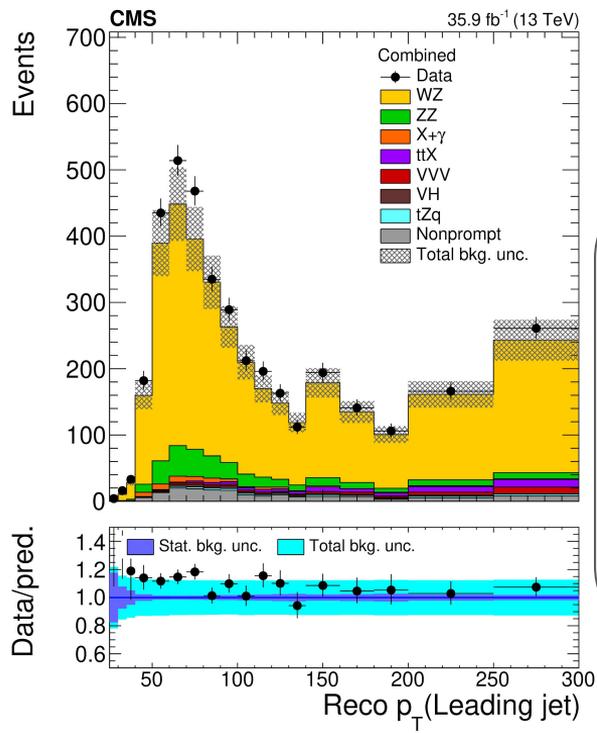
Differential cross-section(s) measurements

→ Also providing cross-section results in differential form for:

→ p_T (leading jet): probe to recoil of the WZ system.

→ p_T (W), p_T (Z), M (WZ): typically sensitive to high energy BSM models.

[10.1007/JHEP04\(2019\)122](https://arxiv.org/abs/10.1007/JHEP04(2019)122)

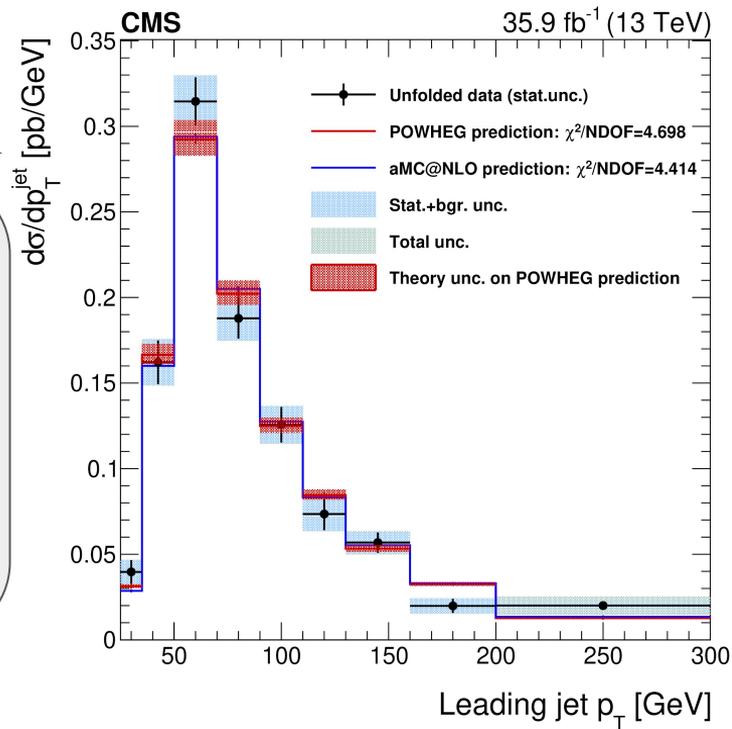


Unfolding with:

- Yields area constrain.

- Bias scale (towards measured total x-sec).

-No additional (i.e. Tikhonov) regularization terms.



Anomalous Couplings searches (I)

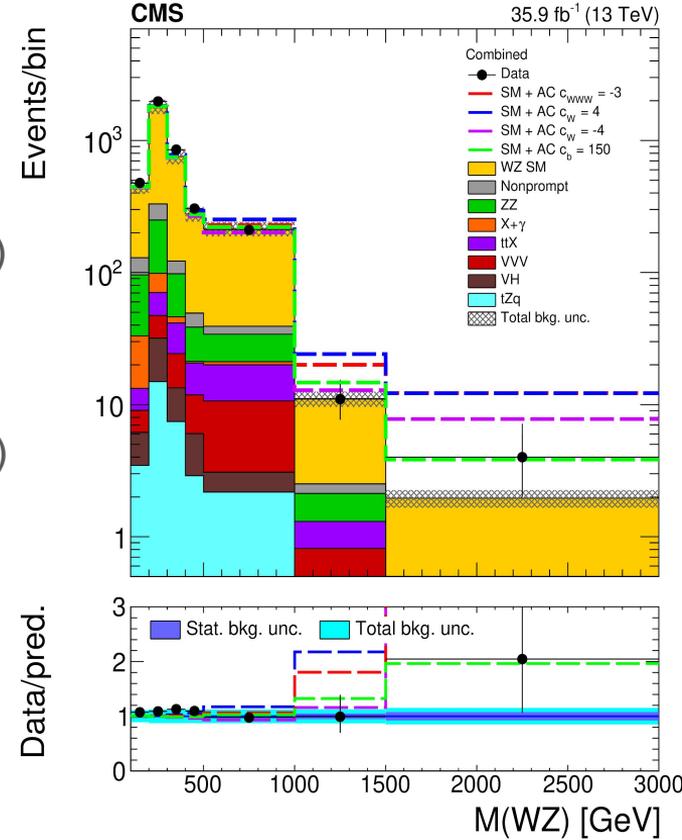
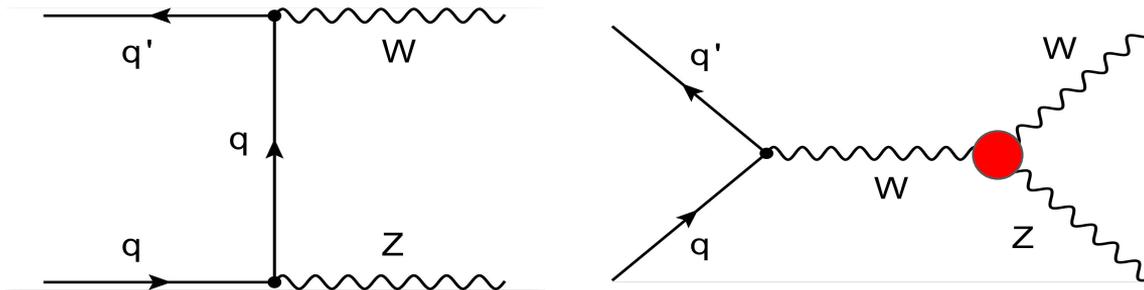
[10.1007/JHEP04\(2019\)122](https://arxiv.org/abs/10.1007/JHEP04(2019)122)

→ Anomalous gauge couplings based on an EFT approach:

$$\delta\mathcal{L}_{AC} = c_{WWW} \text{Tr}[W_{\mu\nu} W^{\nu\rho} W_{\rho}^{\mu}] + c_w (D_{\mu}H)^{\dagger} W^{\mu\nu} (D_{\nu}H) + c_b (D_{\mu}H)^{\dagger} B^{\mu\nu} (D_{\nu}H)$$

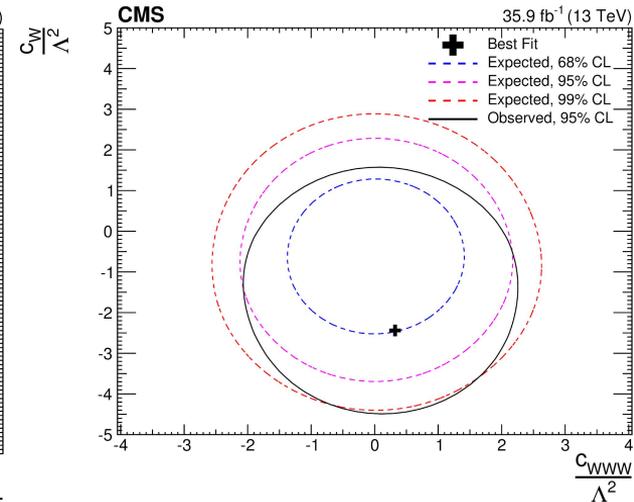
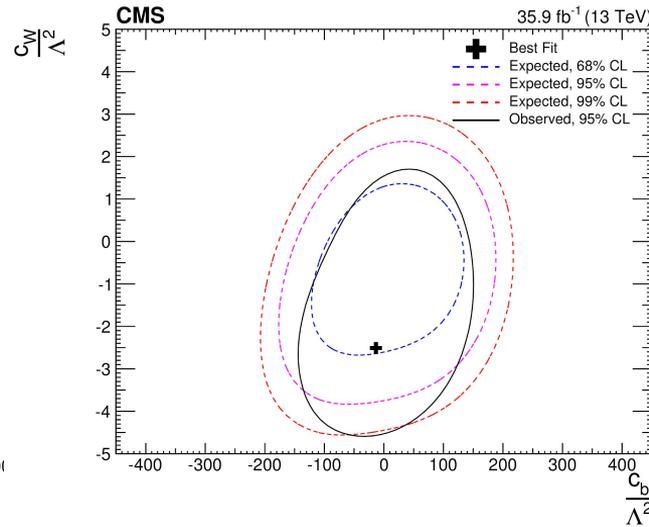
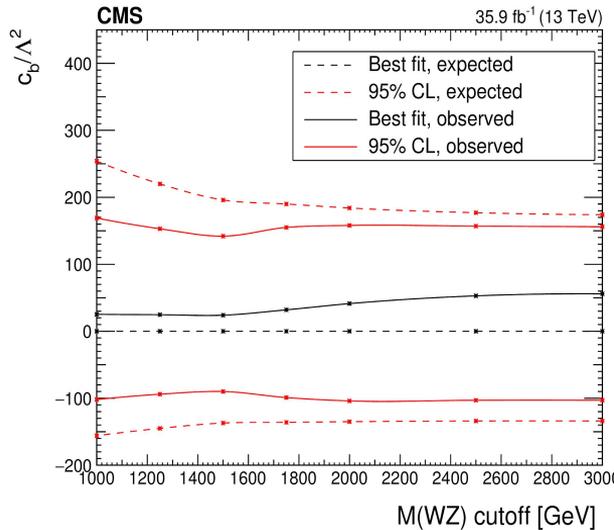
→ Signals modelled as linear (dim. 6) and quadratic (dim. 8) multiplicative terms to the cross-section (differentially).

→ Results extracted from likelihood fit to the shape of $M(WZ)$ (reconstructed assuming $p_z^{\nu} = 0$).



Anomalous Couplings searches (II)

- Results for 1D and 2D scans over AC parameters (both dim. 6 and dim.6 + dim.8).
- Also providing results with varying cut-off energy scale.
- Measurements show very small deviations (within 1σ) of SM predictions. Driven by a fluctuation at high $M(WZ)$ (Slide 6). Constraints are stronger than previous WZ results.
- Strictly limited by statistical power at the high energy tails: improvements at Run II.



Prospects - boson polarization

→ Increasingly higher statistics allow for new high precision measurements.

→ Sensitivity is enough to study the polarization of W/Z bosons in WZ production. Typically sensitive to angles between final state leptons and its mother boson:

$$\frac{d\sigma}{\sigma d\cos\theta_3} \equiv \frac{3}{8} \left[(1 \mp \cos\theta_3)^2 f_L^{W^\pm} + (1 \pm \cos\theta_3)^2 f_R^{W^\pm} + 2\sin^2\theta_3 f_0^{W^\pm} \right]$$

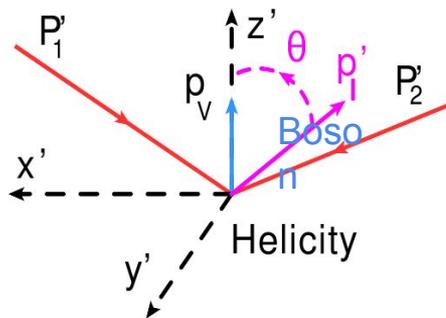
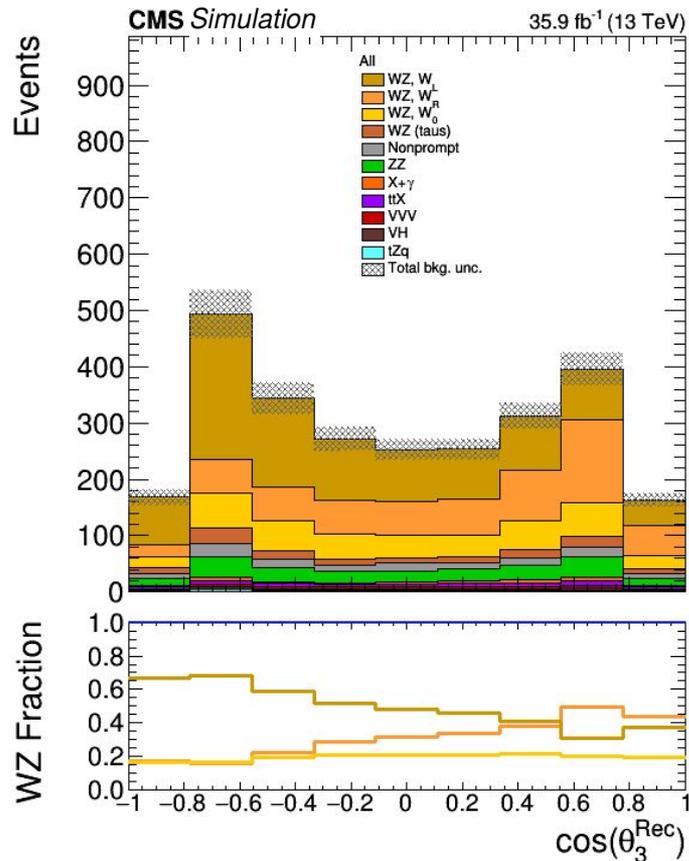
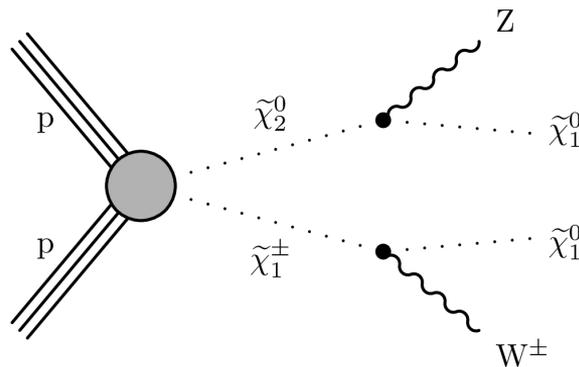
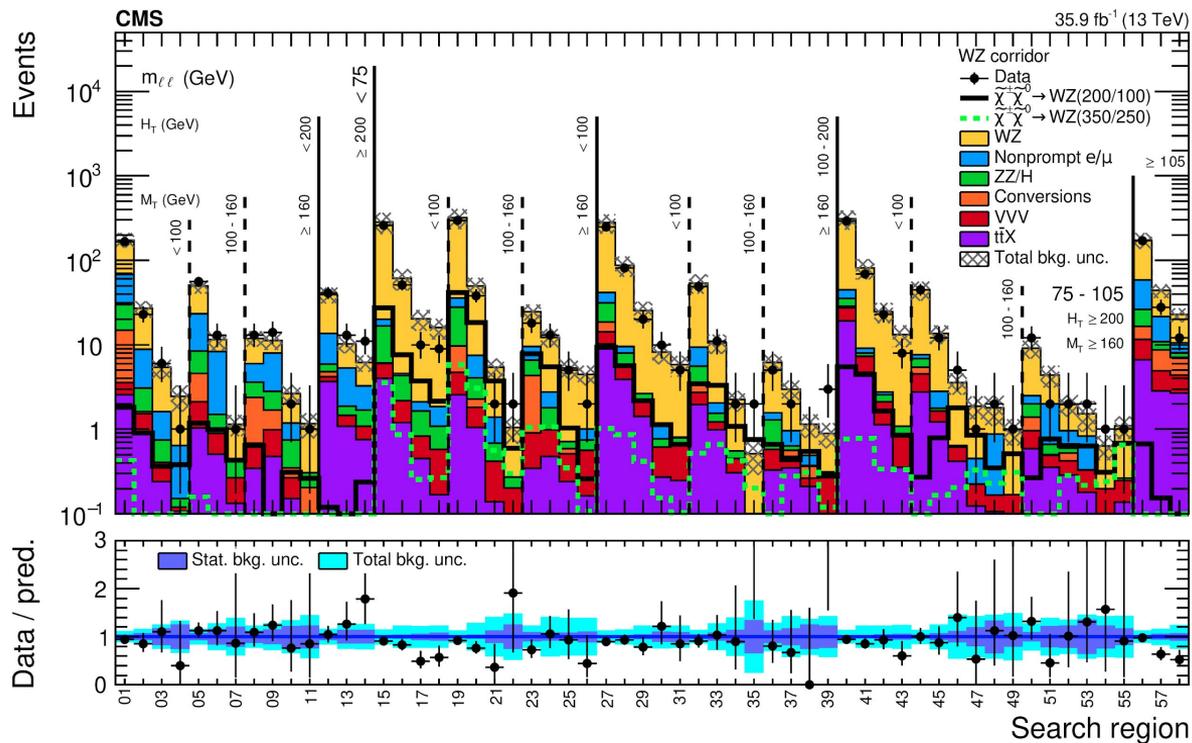


Figure from Baglio, Ninh *JHEP* 1904 (2019) 065



WZ-like Supersymmetric searches

[10.1007/JHEP03\(2018\)160](https://arxiv.org/abs/1803.09817)



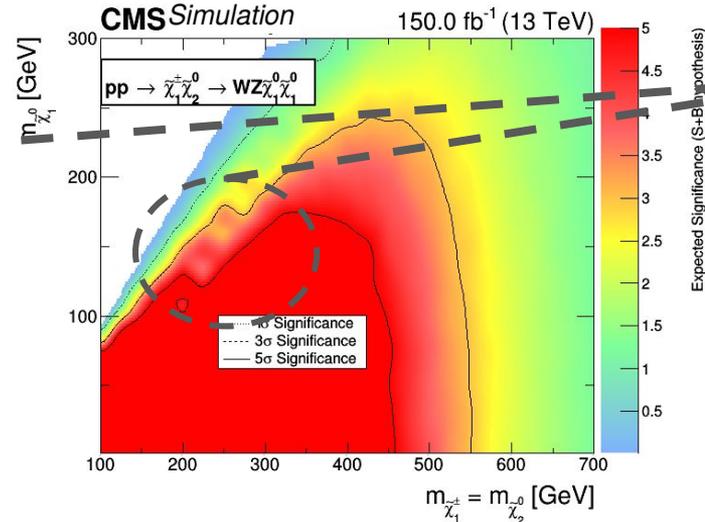
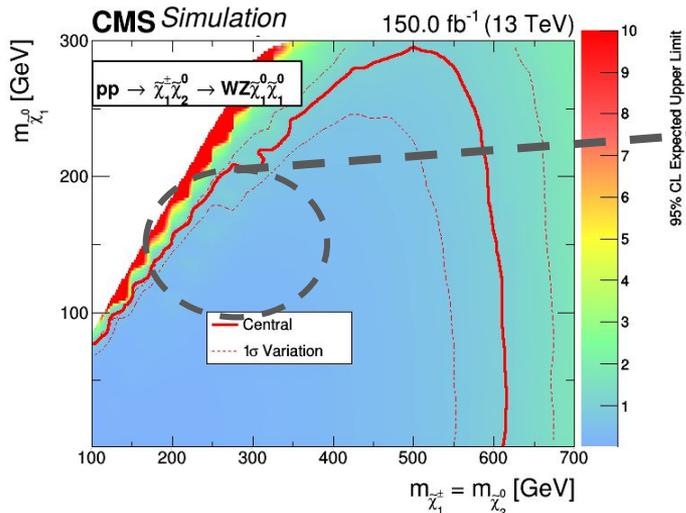
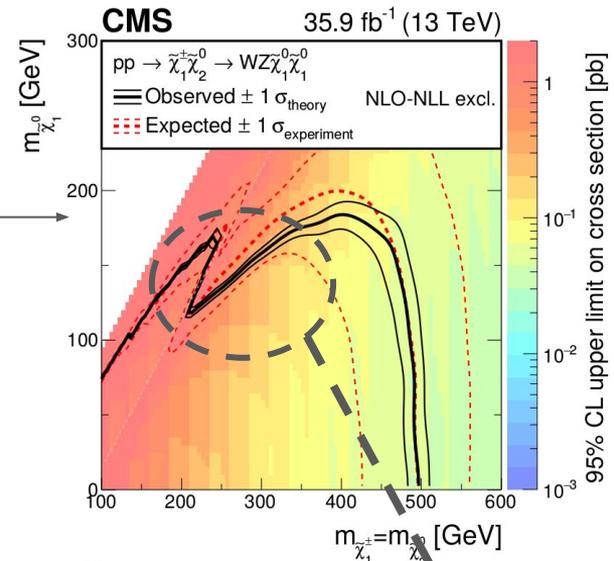
→ $T_{\tilde{\chi}}WZ$ model predicting final states very similar to SM WZ production.

→ Slightly higher p_T^{miss} , m_T^W , etc.

→ Search is limited by our knowledge of WZ differential cross-sections (enhanced p_T^{miss} tails).

Supersymmetry Prospects

- So far excluded NLSP/LSP up to $\sim 500/200$ GeV.
- Run II analysis driven by techniques to improve reach:
 - Increased knowledge of WZ cross-section.
 - Enhanced non-prompt background discriminance.
 - New multivariate techniques.
- Greater exclusion, but also increased discovery reach:



WZ Gap

Conclusions

- New results on WZ production show overall agreement with SM expectations.
- WZ Cross-sections entering the high precision era, systematics require careful assessment in Run-II analysis. Minor tensions in anomalous couplings ($\sim 1\sigma$) to be clarified with new data.
- Additional prospects: expanding into boson spin measurements, analyze boson polarization observables.
- Multileptonic SUSY searches will greatly benefit by enhanced knowledge of WZ cross-section. Prospects: clarifying the status of the WZ “gap” with reach for both exclusion and discovery.

Back-Up