

Particle physics perspectives after LHC run 2

BSM
Model Building
High- P_T

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ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: March 2019

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 8, 13 \text{ TeV}$$

Model	ℓ, γ	Jets†	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit	Reference
Extra dimensions	ADD $G_{KK} + g/q$	$0 e, \mu$	$1-4 j$	Yes	36.1	M_D 7.7 TeV
	ADD non-resonant $\gamma\gamma$	2γ	-	-	36.7	M_S 8.6 TeV
	ADD OBH	-	$2 j$	-	37.0	M_{BH} 8.9 TeV
	ADD BH high $\sum p_T$	$\geq 1 e, \mu$	$\geq 2 j$	-	3.2	M_{BH} 8.2 TeV
	ADD BH multijet	-	$\geq 3 j$	-	3.6	M_{BH} 9.55 TeV
	RS1 $G_{KK} \rightarrow \gamma\gamma$	2γ	-	-	36.7	G_{KK} mass 4.1 TeV
	Bulk RS $G_{KK} \rightarrow WW/ZZ$	multi-channel	-	-	36.1	G_{KK} mass 2.3 TeV
	Bulk RS $G_{KK} \rightarrow WW/ZZ \rightarrow qqqq$	$0 e, \mu$	$2 j$	-	139	G_{KK} mass 2.8 TeV
	Bulk RS $g_{KK} \rightarrow tt$	$1 e, \mu$	$\geq 1 b, \geq 1J/2j$	Yes	36.1	g_{KK} mass 1.8 TeV
	ZUED / RPP	$1 e, \mu$	$\geq 2 b, \geq 3 j$	Yes	36.1	KK mass 1.8 TeV
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2 e, \mu$	-	-	139	Z' mass 5.1 TeV
	SSM $Z' \rightarrow \tau\tau$	2τ	-	-	36.1	Z' mass 2.42 TeV
	Leptophobic $Z' \rightarrow bb$	-	$2 b$	-	36.1	Z' mass 2.1 TeV
	Leptophobic $Z' \rightarrow tt$	$1 e, \mu$	$\geq 1 b, \geq 1J/2j$	Yes	36.1	Z' mass 3.0 TeV
	SSM $W' \rightarrow \ell\nu$	$1 e, \mu$	-	Yes	79.8	W' mass 5.6 TeV
	SSM $W' \rightarrow \tau\nu$	1τ	-	Yes	36.1	W' mass 3.7 TeV
	HVT $V' \rightarrow WW \rightarrow qqqq$ model B	$0 e, \mu$	$2 j$	-	139	V' mass 4.0 TeV
	HVT $V' \rightarrow WH/ZH$ model B	multi-channel	-	-	36.1	V' mass 2.93 TeV
	LRSM $W'_R \rightarrow tb$	multi-channel	-	-	36.1	W' mass 3.25 TeV
	CI	CI $qqqq$	-	$2 j$	-	37.0
CI $\ell\ell qq$		$2 e, \mu$	-	-	36.1	Λ 40.0 TeV η_{LL}
CI $tttt$		$\geq 1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Λ 2.57 TeV $ C_{\text{eff}} = 4\pi$
DM	Axial-vector mediator (Dirac DM)	$0 e, \mu$	$1-4 j$	Yes	36.1	m_{med} 1.55 TeV $g_s=0.25, g_t=1.0, m(\chi) = 1 \text{ GeV}$
	Colored scalar mediator (Dirac DM)	$0 e, \mu$	$1-4 j$	Yes	36.1	m_{med} 1.67 TeV $g=1.0, m(\chi) = 1 \text{ GeV}$
	$V\chi\chi$ EFT (Dirac DM)	$0 e, \mu$	$1 j, \leq 1 j$	Yes	3.2	$M_\chi < 150 \text{ GeV}$
	Scalar reson. $\phi \rightarrow t\bar{t}$ (Dirac DM)	$0-1 e, \mu$	$1 b, 0-1 j$	Yes	36.1	m_ϕ 700 GeV $y = 0.4, \lambda = 0.2, m(\chi) = 10 \text{ GeV}$
LQ	Scalar LQ 1 st gen	$1, 2 e$	$\geq 2 j$	Yes	36.1	LQ mass 1.4 TeV $\beta = 1$
	Scalar LQ 2 nd gen	$1, 2 \mu$	$\geq 2 j$	Yes	36.1	LQ mass 1.56 TeV $\beta = 1$
	Scalar LQ 3 rd gen	2τ	$2 b$	-	36.1	LQ mass 1.0 TeV $\mathcal{B}(LQ_s^+ \rightarrow b\tau) = 1$
	Scalar LQ 3 rd gen	$0-1 e, \mu$	$2 b$	Yes	36.1	LQ mass 970 GeV $\mathcal{B}(LQ_s^+ \rightarrow t\tau) = 0$
	Scalar LQ 3 rd gen	$0-1 e, \mu$	$2 b$	Yes	36.1	LQ mass 970 GeV $\mathcal{B}(LQ_s^+ \rightarrow t\tau) = 0$
Heavy quarks	VLQ $TT \rightarrow Ht/Zt/Wb + X$	multi-channel	-	-	36.1	T mass 1.37 TeV SU(2) doublet
	VLQ $BB \rightarrow Wt/Zb + X$	multi-channel	-	-	36.1	B mass 1.34 TeV SU(2) doublet
	VLQ $T_{5/3} T_{5/3} T_{5/3} \rightarrow Wt + X$	$2(SS)/\geq 3 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	$T_{5/3}$ mass 1.64 TeV $\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$
	VLQ $Y \rightarrow Wb + X$	$1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Y mass 1.85 TeV $\mathcal{B}(Y \rightarrow Wb) = 1, c_Y(Wb) = 1$
	VLQ $B \rightarrow Hb + X$	$0 e, \mu, 2 \gamma$	$\geq 1 b, \geq 1 j$	Yes	79.8	B mass 1.21 TeV $\kappa_B = 0.5$
VLQ $QQ \rightarrow WqWq$	$1 e, \mu$	$\geq 4 j$	Yes	20.3	Q mass 690 GeV	
Excited fermions	Excited quark $q^* \rightarrow qg$	-	$2 j$	-	139	q^* mass 6.7 TeV only u^* and $d^*, \Lambda = m(q^*)$
	Excited quark $q^* \rightarrow q\gamma$	1γ	$1 j$	-	36.7	q^* mass 5.3 TeV only u^* and $d^*, \Lambda = m(q^*)$
	Excited quark $b^* \rightarrow b\gamma$	-	$1 b, 1 j$	-	36.1	b^* mass 2.6 TeV
	Excited lepton ℓ^*	$3 e, \mu$	-	-	20.3	ℓ^* mass 3.0 TeV $\Lambda = 3.0 \text{ TeV}$
	Excited lepton ν^*	$3 e, \mu, \tau$	-	-	20.3	ν^* mass 1.6 TeV $\Lambda = 1.6 \text{ TeV}$
	Other	Type III Seesaw	$1 e, \mu$	$\geq 2 j$	Yes	79.8
LRSM Majorana ν		2μ	$2 j$	-	36.1	N_R mass 3.2 TeV $m(W_R) = 4.1 \text{ TeV}, g_L = g_R$
Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$		$2, 3, 4 e, \mu$ (SS)	-	-	36.1	$H^{\pm\pm}$ mass 870 GeV DY production
Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$		$3 e, \mu, \tau$	-	-	20.3	$H^{\pm\pm}$ mass 400 GeV DY production, $\mathcal{B}(H^{\pm\pm} \rightarrow \ell\tau) = 1$
Multi-charged particles		-	-	-	36.1	multi-charged particle mass 1.2 TeV DY production, $ q = 5e$
Magnetic monopoles		-	-	-	7.0	monopole mass 1.4 TeV DY production, $ g = 1g_D, \text{spin } 1/2$

ATLAS SUSY Searches* - 95% CL Lower Limits

March 2019

ATLAS Preliminary

$$\sqrt{s} = 13 \text{ TeV}$$

Model	Signature	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference
Inclusive Searches	$q\bar{q}, \bar{q} \rightarrow q\bar{q}^0$	$0 e, \mu$ mono-jet	E_T^{miss} 36.1	q [2x, 8x Degen] $m(q) < 100 \text{ GeV}$
	$\bar{q}\bar{q}, \bar{q} \rightarrow q\bar{q}^0$	$0 e, \mu$ 2-6 jets	E_T^{miss} 36.1	q [1x, 8x Degen] 0.43 0.9 1.55 $m(q) - m(\bar{q}) = 5 \text{ GeV}$
	$\bar{q}\bar{q}, \bar{q} \rightarrow q\bar{q}^0$	$0 e, \mu$ 2-6 jets	E_T^{miss} 36.1	\bar{q} 2.0 $m(\bar{q}) = 200 \text{ GeV}$
	$\bar{q}\bar{q}, \bar{q} \rightarrow q\bar{q}(\ell\ell)\bar{q}^0$	$3 e, \mu$ 4 jets	E_T^{miss} 36.1	\bar{q} 1.85 $m(\bar{q}) = 900 \text{ GeV}$
	$\bar{q}\bar{q}, \bar{q} \rightarrow q\bar{q}WZ\bar{q}^0$	$0 e, \mu$ 7-11 jets	E_T^{miss} 36.1	\bar{q} 1.2 $m(\bar{q}) - m(\bar{q}^0) = 50 \text{ GeV}$
	$\bar{q}\bar{q}, \bar{q} \rightarrow q\bar{q}WZ\bar{q}^0$	$3 e, \mu$ 4 jets	E_T^{miss} 36.1	\bar{q} 1.8 $m(\bar{q}) < 400 \text{ GeV}$
	$\bar{q}\bar{q}, \bar{q} \rightarrow t\bar{t}\bar{q}^0$	$0-1 e, \mu$ 3 b	E_T^{miss} 79.8	\bar{q} 0.98 2.25 $m(\bar{q}) - m(\bar{q}^0) = 200 \text{ GeV}$
3 rd gen. squarks direct production	$\bar{b}_1\bar{b}_1, \bar{b}_1 \rightarrow b\bar{b}_1^0/\bar{t}\bar{t}_1^0$	Multiple	E_T^{miss} 36.1	\bar{b}_1 Forbidden $m(\bar{b}_1) = 300 \text{ GeV}, \text{BR}(\bar{b}_1^0) = 1$
	$\bar{b}_1\bar{b}_1, \bar{b}_1 \rightarrow b\bar{b}_1^0/\bar{t}\bar{t}_1^0$	Multiple	E_T^{miss} 36.1	\bar{b}_1 Forbidden 0.58-0.82 0.7 $m(\bar{b}_1) = 300 \text{ GeV}, \text{BR}(\bar{b}_1^0) = 0.5$
	$\bar{b}_1\bar{b}_1, \bar{b}_1 \rightarrow b\bar{b}_1^0/\bar{t}\bar{t}_1^0$	Multiple	E_T^{miss} 36.1	\bar{b}_1 Forbidden 0.23-0.48 0.23-1.35 $m(\bar{b}_1) = 200 \text{ GeV}, m(\bar{b}_1^0) = 300 \text{ GeV}, \text{BR}(\bar{b}_1^0) = 1$
	$\bar{t}_1\bar{t}_1, \bar{t}_1 \rightarrow W\bar{t}_1^0/\bar{b}_1^0$	$0-2 e, \mu$ 0-2 jets/1-2 b	E_T^{miss} 36.1	\bar{t}_1 1.0 $\Delta m(\bar{t}_1^0, \bar{t}_1^0) = 130 \text{ GeV}, m(\bar{t}_1^0) = 100 \text{ GeV}$
	$\bar{t}_1\bar{t}_1, \bar{t}_1 \rightarrow W\bar{t}_1^0/\bar{b}_1^0$	Multiple	E_T^{miss} 36.1	\bar{t}_1 0.48-0.84 $m(\bar{t}_1^0) = 150 \text{ GeV}, m(\bar{t}_1^0) - m(\bar{t}_1^0) = 5 \text{ GeV}, \bar{t}_1 = \bar{t}_1$
	$\bar{t}_1\bar{t}_1, \bar{t}_1 \rightarrow \tau\bar{t}_1/\bar{b}_1, \bar{t}_1 \rightarrow \tau\bar{t}_1/\bar{b}_1$	$1 \tau + 1 e, \mu, \tau$ 2 jets/1 b	E_T^{miss} 36.1	\bar{t}_1 1.16 $m(\bar{t}_1) = 800 \text{ GeV}$
	$\bar{t}_1\bar{t}_1, \bar{t}_1 \rightarrow c\bar{t}_1^0/\bar{c}\bar{c}, \bar{t}_1 \rightarrow c\bar{t}_1^0/\bar{c}\bar{c}$	$0 e, \mu$ 2 c	E_T^{miss} 36.1	\bar{t}_1 0.85 $m(\bar{t}_1) = 0 \text{ GeV}$
	$\bar{t}_1\bar{t}_1, \bar{t}_1 \rightarrow c\bar{t}_1^0/\bar{c}\bar{c}, \bar{t}_1 \rightarrow c\bar{t}_1^0/\bar{c}\bar{c}$	$0 e, \mu$ mono-jet	E_T^{miss} 36.1	\bar{t}_1 0.46 0.43 $m(\bar{t}_1, \bar{t}_1) - m(\bar{t}_1^0) = 50 \text{ GeV}$
	$\bar{t}_1\bar{t}_1, \bar{t}_1 \rightarrow h\bar{t}_1 + h$	$1-2 e, \mu$ 4 b	E_T^{miss} 36.1	\bar{t}_1 0.32-0.88 $m(\bar{t}_1^0) = 0 \text{ GeV}, m(\bar{t}_1) - m(\bar{t}_1^0) = 180 \text{ GeV}$
	EW direct	$\bar{\chi}_1^+ \bar{\chi}_1^0$ via WZ	$2-3 e, \mu$ ee, $\mu\mu$	E_T^{miss} 36.1
$\bar{\chi}_1^+ \bar{\chi}_1^0$ via WW		$2 e, \mu$	E_T^{miss} 139	$\bar{\chi}_1^+/\bar{\chi}_1^0$ 0.42 $m(\bar{\chi}_1^+) = 10 \text{ GeV}$
$\bar{\chi}_1^+ \bar{\chi}_1^0$ via Wh		$0-1 e, \mu$ 2 b	E_T^{miss} 36.1	$\bar{\chi}_1^+/\bar{\chi}_1^0$ 0.68 $m(\bar{\chi}_1^+) = 0$
$\bar{\chi}_1^+ \bar{\chi}_1^0$ via $\ell\ell/\nu\nu$		$2 e, \mu$	E_T^{miss} 139	$\bar{\chi}_1^+/\bar{\chi}_1^0$ 1.0 $m(\bar{\chi}_1^+) = 0$
$\bar{\chi}_1^+ \bar{\chi}_1^0/\bar{\chi}_2^0, \bar{\chi}_1^+ \rightarrow \tau_1\nu(\tau\nu), \bar{\chi}_2^0 \rightarrow \tau_1\nu(\tau\nu)$		2τ	E_T^{miss} 36.1	$\bar{\chi}_1^+/\bar{\chi}_2^0$ 0.76 $m(\bar{\chi}_1^+) = 0, m(\bar{\chi}_2^0) = 100 \text{ GeV}, m(\bar{\chi}_1^+) - m(\bar{\chi}_2^0) = 5 \text{ GeV}$
$\bar{\chi}_1^+ \bar{\chi}_1^0, \bar{\chi}_1^+ \rightarrow \tau_1\nu(\tau\nu), \bar{\chi}_2^0 \rightarrow \tau_1\nu(\tau\nu)$		$2 e, \mu$ 0 jets	E_T^{miss} 139	$\bar{\chi}_1^+$ 0.7 $m(\bar{\chi}_1^+) = 0$
$\bar{\chi}_1^+ \bar{\chi}_1^0, \bar{\chi}_1^+ \rightarrow \tau_1\nu(\tau\nu), \bar{\chi}_2^0 \rightarrow \tau_1\nu(\tau\nu)$		$2 e, \mu$ ≥ 1	E_T^{miss} 36.1	$\bar{\chi}_1^+$ 0.18 $m(\bar{\chi}_1^+) - m(\bar{\chi}_2^0) = 5 \text{ GeV}$
Long-lived particles	$\bar{H}\bar{H}, \bar{H} \rightarrow hG/ZG$	$0 e, \mu$ $\geq 3 b$	E_T^{miss} 36.1	\bar{H} 0.13-0.23 0.29-0.88 $\text{BR}(\bar{H}^0 \rightarrow hG) = 1$
	Direct $\bar{\chi}_1^+ \bar{\chi}_1^0$ prod., long-lived $\bar{\chi}_1^+$	Disapp. trk 1 jet	E_T^{miss} 36.1	$\bar{\chi}_1^+$ 0.15 0.46 Pure Wino
	Stable \bar{g} R-hadron	Multiple	E_T^{miss} 36.1	\bar{g} 2.0 Pure Higgsino
RPV	LFV $pp \rightarrow \bar{\nu}_e + X, \bar{\nu}_e \rightarrow e\mu/\tau\mu/\tau\tau$	$e\mu, \tau\mu, \tau\tau$	E_T^{miss} 3.2	$\bar{\nu}_e$ 1.9 $A'_{111} = 0.11, A_{1212(13213)} = 0.07$
	$\bar{\chi}_1^+ \bar{\chi}_1^0/\bar{\chi}_2^0 \rightarrow W/Z\ell\ell/\nu\nu$	$4 e, \mu$ 0 jets	E_T^{miss} 36.1	$\bar{\chi}_1^+/\bar{\chi}_2^0$ 0.82 1.33 $m(\bar{\chi}_1^+) = 100 \text{ GeV}$
	$\bar{\chi}_1^+ \bar{\chi}_1^0/\bar{\chi}_2^0 \rightarrow W/Z\ell\ell/\nu\nu$	4-5 large-R jets	E_T^{miss} 36.1	$\bar{\chi}_1^+/\bar{\chi}_2^0$ 1.05 1.3 $m(\bar{\chi}_1^+) = 200 \text{ GeV}, 1100 \text{ GeV}$
	$\bar{\chi}_1^+ \bar{\chi}_1^0/\bar{\chi}_2^0 \rightarrow W/Z\ell\ell/\nu\nu$	Multiple	E_T^{miss} 36.1	$\bar{\chi}_1^+/\bar{\chi}_2^0$ 1.05 2.0 $m(\bar{\chi}_1^+) = 200 \text{ GeV}, \text{bino-like}$
	$\bar{\chi}_1^+ \bar{\chi}_1^0/\bar{\chi}_2^0 \rightarrow W/Z\ell\ell/\nu\nu$	Multiple	E_T^{miss} 36.1	$\bar{\chi}_1^+/\bar{\chi}_2^0$ 0.55 1.05 $m(\bar{\chi}_1^+) = 200 \text{ GeV}, \text{bino-like}$
	$\bar{\chi}_1^+ \bar{\chi}_1^0/\bar{\chi}_2^0 \rightarrow W/Z\ell\ell/\nu\nu$	2 jets + 2 b	E_T^{miss} 36.1	$\bar{\chi}_1^+/\bar{\chi}_2^0$ 0.42 0.6 $m(\bar{\chi}_1^+) = 200 \text{ GeV}, \text{bino-like}$

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

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†Small-radius (large-radius) jets are denoted by the letter j (J).

$\sqrt{s} = 8 \text{ TeV}$ $\sqrt{s} = 13 \text{ TeV}$ partial data full data

10⁻¹ 1 Mass scale [TeV]

- Large number of analyses trying to cover all corners of model and parameter space
- Very stringent bounds banning new physics beyond the TeV frontier.
- Is it over? NO!
 - Take the gap seriously: EFTs
 - Critically assess reach of experimental analyses

NP gap: EFT

- If there is a mass gap there won't be impressive 5σ signals but small (consistent) deviations
- In this case effective field theory is the way to go
- Global fits will be crucial to achieve a consistent picture
- Top-down approach (matching to specific models) needed to characterize the new physics

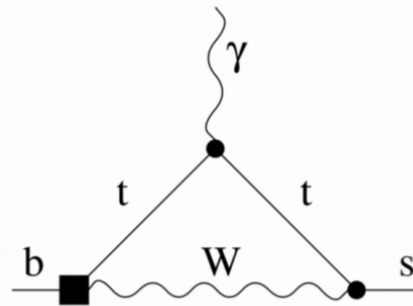
Why Global Fits (+ top-down)

- Example: top anomalous couplings and $b \rightarrow s\gamma$

$$R_{b \rightarrow s\gamma} \equiv \frac{B(b \rightarrow s\gamma)}{B(b \rightarrow s\gamma)_{SM}} - 1 = -2.45 \Delta C_7 - 0.58 \Delta C_8 \quad \text{Grzadkowski, Misiak ('08)}$$

$$= -2.51 \delta v_L + 131 v_R - 256 g_L + 0.58 g_R - 2.45 C_7^{(p)} - 0.58 C_8^{(p)}$$

$$\frac{-ig_w}{\sqrt{2}} [\gamma_\mu (v_L P_L + v_R P_R) + \frac{i\sigma_{\mu\nu} q^\nu}{M_W} (g_L P_L + g_R P_R)]$$



$$\mathcal{O}_7 = \frac{em_b}{16\pi^2} \bar{s}_L \sigma^{\mu\nu} b_R F_{\mu\nu},$$

$$\mathcal{O}_8 = \frac{g_s m_b}{16\pi^2} \bar{s}_L \sigma^{\mu\nu} T^A b_R G_{\mu\nu}^A,$$

One-by-one limits

bound	δv_L	v_R	g_L	g_R	$C_7^{(p)}$	$C_8^{(p)}$
upper	0.03	0.0025	0.0004	0.57	0.04	0.15
lower	-0.13	-0.0007	-0.0013	-0.15	-0.14	-0.56

But contributions are rarely one by one

Why Global Fits (+ top-down)

- Example: top anomalous couplings and $b \rightarrow s\gamma$

$$\Delta\mathcal{L} = \bar{Q}(i\not{D} - M_Q)Q - \left[\lambda_{Qt} \bar{Q}_L \tilde{\phi} t_R + \lambda_{Qb} \bar{Q}_L \phi b_R + \text{h.c.} \right] \quad Q_{L,R} \sim (3, 2)_{1/6}$$

$$R_{b \rightarrow s\gamma} = 131 v_R - 2.45 C_7^{(p)} - 0.58 C_8^{(p)} \approx 171 v_R$$

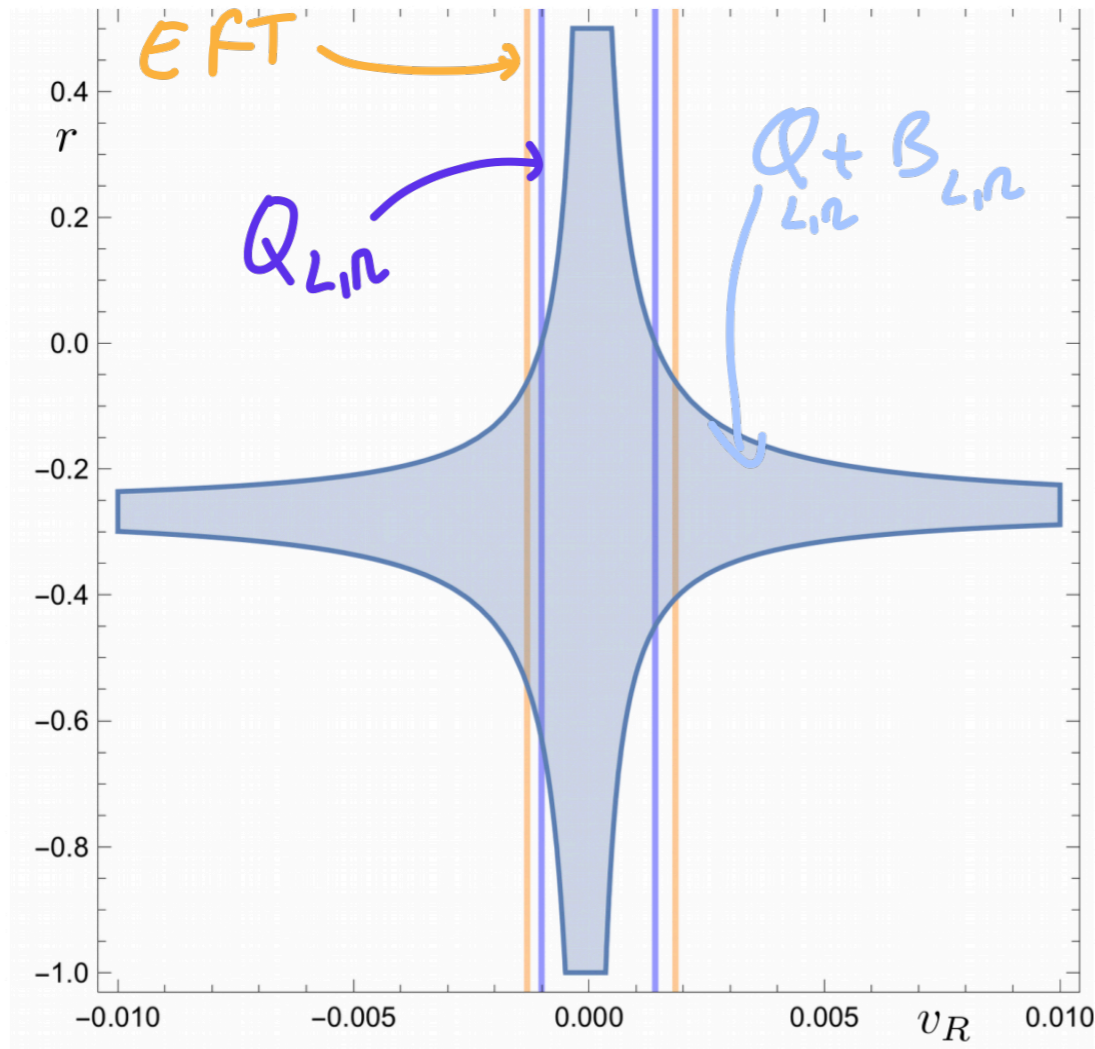
$$\Delta\mathcal{L}' = \bar{B}(i\not{D} - M_B)B - \left[\lambda_{q_2 B} \bar{q}_{2L} \phi B_R + \tilde{\lambda} \bar{Q}_R \phi B_L + \text{h.c.} \right] \quad B_{L,R} \sim (3, 1)_{-1/3}$$

$$R_{b \rightarrow s\gamma} = 131 v_R - 2.45 C_7^{(p)} - 0.58 C_8^{(p)} \approx 171 v_R \left(1 + 3.82 \frac{\lambda_{q_2 B} \tilde{\lambda} M_Q}{\lambda_{Qt}^* M_B} \right)$$

Matching done with MatchMaker: C. Anastasiou, A. Carmona,
A. Lazopoulos, J.S.

Why Global Fits (+ top-down)

$$R_{b \rightarrow s\gamma} = 131 v_R - 2.45 C_7^{(p)} - 0.58 C_8^{(p)} \approx 171 v_R \left(1 + \underbrace{3.82 \frac{\lambda_{q_2 B} \tilde{\lambda} M_Q}{\lambda_{Qt}^* M_B}}_r \right)$$



$$\text{Br}^{\text{exp}}[b \rightarrow s\gamma] = (3.43 \pm 0.21 \pm 0.07) \times 10^{-4},$$

$$\text{Br}^{\text{SM}}[b \rightarrow s\gamma] = (3.36 \pm 0.23) \times 10^{-4}.$$

$$-0.17 \leq R_{b \rightarrow s\gamma} \leq 0.24$$

$$-0.0013 \leq v_R \leq 0.0018$$

$$-0.0010 \leq v_R \leq 0.0014$$

v_R unconstrained for $r \approx -0.26$

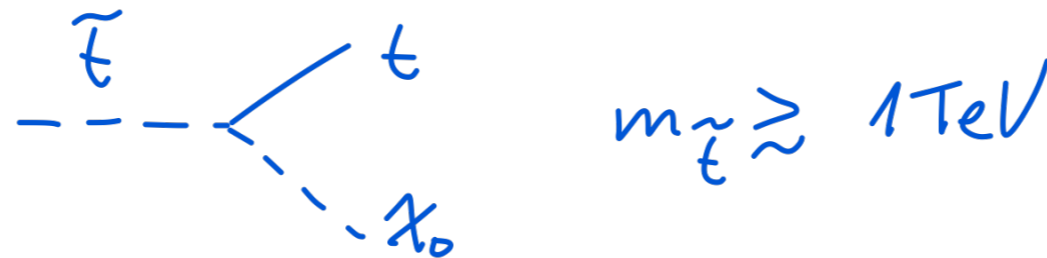
- But B contributes to other observables $\sigma_{Z_{SL}}, V_{is}, \gamma_{sb}, \dots$

Critically assess exp. analyses

- Experimental analyses largely inspired by minimal/simplified models
- Most models are not minimal and their signatures can be vastly different from the ones we are searching for.
 - R-parity violating SUSY
 - VL quarks with non-standard decays

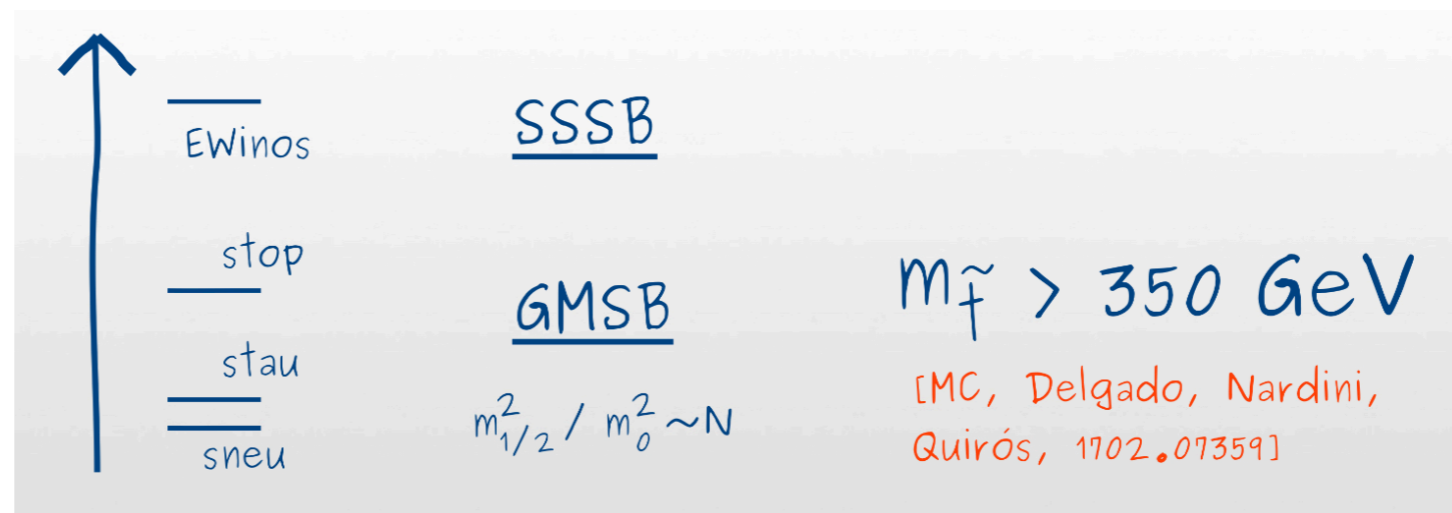
Critically assess exp. analyses

- R-parity violating SUSY
- Most analyses assume R-parity with lightest neutralino as DM candidate



- Things change if R-parity violated and different spectrum

Borrowed from
M. Chala

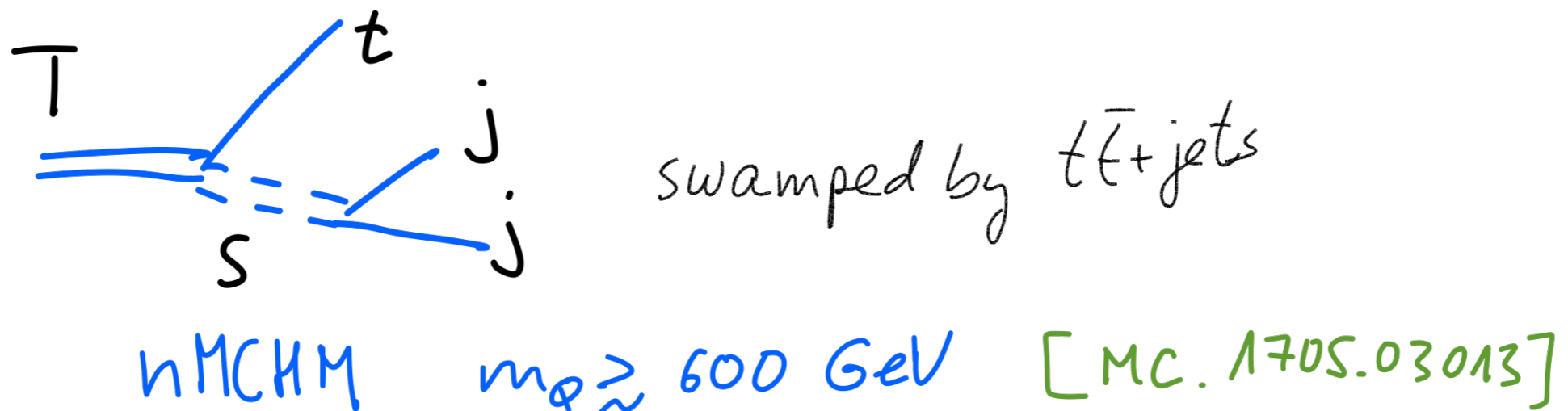


Critically assess exp. analyses

- VL quarks with non standard decays
- Searches for VL quarks assume SM decays



- Things change if there are new light particles. Very common in non-minimal models.



Critically assess exp. analyses

- Experimental analyses largely inspired by minimal/simplified models result in very stringent bounds on new physics
- Realistic, natural models are usually far from minimal
- Non-optimized searches result in much milder bounds, leaving plenty of room for new physics to show up in future LHC runs and other experiments.