

### Energy Frontier of the Intensity Frontier: Millicharged Particle in Neutrino Exps. and FerMINI

#### Yu-Dai Tsai, Fermilab (WH674W) / U. Chicago

with Magill, Plestid, Maxim Pospelov (1806.03310, PRL '19)

with Kelly (1812.03998, submitted to PRL)

Email: <a href="mailto:ytsai@fnal.gov">ytsai@fnal.gov</a>, arXiv: <a href="mailto:https://arxiv.org/a/tsai">https://arxiv.org/a/tsai</a> y 1.html

#### **TEAM MCP!**



Maxim Pospelov
Minnesota / Perimeter



**Yu-Dai Tsai**(Presenting)
Fermilab / U.Chi



Arguelles, MIT



**Hostert,** Durham Team Dark Neutrino



**Gabriel Magill**McMaster



Ryan Plestid McMaster



**Kevin Kelly** Fermilab

### Outline

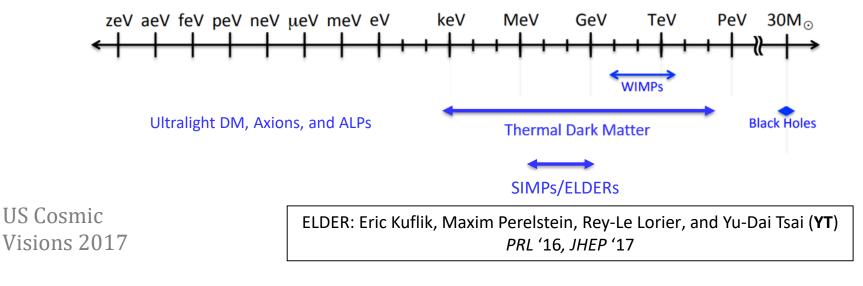
- Motivations
- Millicharged Particle (MCP) & Proton Fixed-Target Experiments
- Sensitivity Reach @ DUNE near detector (LAr-TPC ND)
   (and other neutrino detectors)
- II) Sensitivity Reach @ FerMINI setup (adding a low-cost detector in the ND complex, further extend MCP sensitivity)
- Discussion

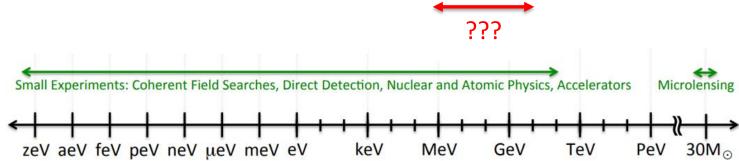
Yu-Dai Tsai, Fermilab

# Neutrino & Proton Fixed-Target (FT) Experiments: Natural habitats for signals of weakly interacting / long-lived / hidden particles But why? Why MeV - GeV+?

### Dark Matter/Hidden Particles Exploration

Dark Sector Candidates, Anomalies, and Search Techniques



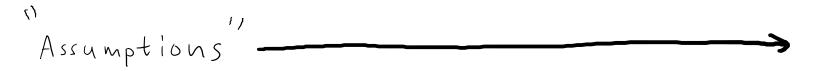


Proton fix-target/neutrino experiments are important for MeV ~ 10 GeV!

### Hidden Particles in Neutrino Experiments

- Neutrinos are weakly interacting particles. Just like
   Millicharged particles
- High statistics, e.g. DUNE plans  $\sim 10^{22}$  Protons on Target (POT)
- Shielded/underground: low background (e.g. solar v programs)
- Many of them existing and many to come: strength in numbers
- Produce hidden particles (from the beam!) without DMabundance or cosmological history assumptions:
   more "direct" than astrophysics/cosmological probes.
- Relatively high energy (LBNF/NuMI: 120 GeV; SPS: 400 GeV)

### Not all bounds are created with equal assumptions



Or, how likely is it that theorists would be able to argue our ways around them

Accelerator-based: Collider, Fixed-Target Experiments Some other ground based experiments

Astrophysical productions (not from ambient DM): energy loss/cooling, etc: Rely on modeling/observations of (extreme/complicated/rare) astro systems

Dark matter direct/indirect detection: abundance, velocity distribution, etc (reveal true story of DM)

2 different

Cosmology: assume cosmological history, species, etc

Signals of discoveries grow from anomalies Maybe nature is telling us something so we don't have to search in the dark? (systematics?)

### Some anomalies involving MeV-GeV+ Explanations

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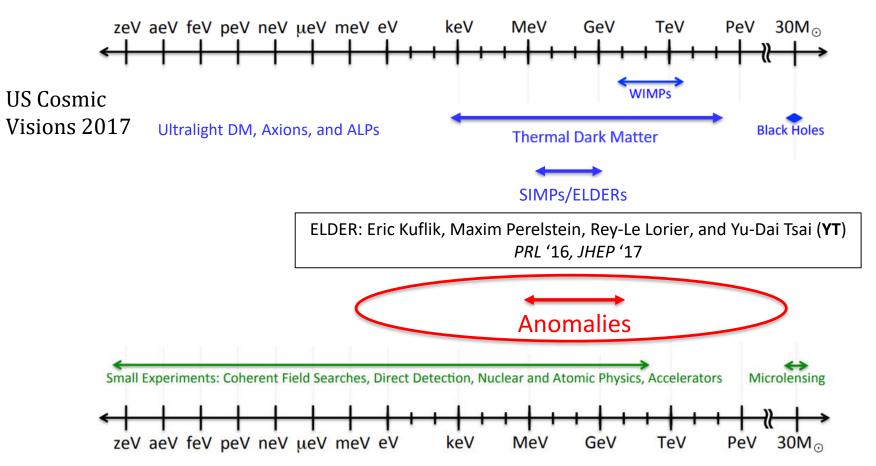
- Muon g-2
- Proton charge radius anomaly
- LSND & MiniBooNE anomaly
- EDGES result

•

Below ~ MeV there are also **strong astrophysical/cosmological bounds** 

### Dark Matter/Hidden Particles Exploration

Dark Sector Candidates, Anomalies, and Search Techniques



- Proton fix-target/neutrino experiments are important for MeV ~ 10 GeV!
- Many anomalies & anomaly explanations in this range!

#### Anomaly & New Physics in Neutrino FT Experiments

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1) Light Scalar & Dark Photon at Borexino & LSND

Pospelov & YT, PLB '18, 1706.00424 (proton charge radius anomaly)

2) Dipole Portal Heavy Neutral Lepton

Magill, Plestid, Pospelov & **YT** , *PRD '18*, <u>1803.03262</u>

(Short-baseline LSND/MiniBoonE anomalies) See Ian's talk for more!

3) Millicharged Particles in Neutrino Experiments

Magill, Plestid, Pospelov & **YT**, *PRL '19*, <u>1806.03310</u>

(EDGES 21-cm measurement anomaly)

•

### MeV – GeV + anomalies: Not just search in the dark

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#### 4) Millicharged Particles in FerMINI Experiments

Kelly & **YT,** <u>1812.03998</u>

(EDGES Anomaly)

5) Dark Neutrino at Scattering Experiments: CHARM-II & MINERVA

**Argüelles**, Hostert, **YT**, <u>1812.08768</u>

(MiniBooNE Anomaly) Also see Pedro/lan's talk for more!

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Two new papers OUT! Happy to chat

### Millicharged Particles

Electric charge quantization?
Other implications (dark sector, etc)

### Finding Minicharge

- Is electric charge quantized? A long-standing question!
- U(1) allows arbitrarily small (any real number) charges.
   Why don't we see them in electric charges?
   Motivates Dirac quantization, Grand Unified Theory (GUT), etc, to explain such quantization
- A test to see if e/3 is the minimal charge
- MCP could have natural link to dark sector (dark photon, etc)
- Could account for dark matter (DM) (WIMP or Freeze-in scenarios)
- Used for the cooling of gas temperature to explain the EDGES result [EDGES collab., Nature, (2018), Barkana, Nature, (2018)].
   A small fraction of the ~ MeV-100 MeV DM as MCP to explain the EDGES anomalous 21-cm absorption spectrum

# Millicharged Particle: Models

### mCP Model

Small charged particles under U(1) hypercharge

$$\mathcal{L}_{\text{mCP}} = i\bar{\psi}(\partial \!\!\!/ - i\epsilon' e \!\!\!/ \!\!\!/ + M_{\text{mCP}})\psi$$

- Can just consider these Lagrangian terms by themselves (no extra mediator, i.e., dark photon), one can call this a "pure" MCP
- Or this could be from Kinetic Mixing
  - give a nice origin to this term
  - an example that gives rise to dark sectors
  - easily compatible with Grand Unification Theory
  - I will not spend too much time on the model

### Kinetic Mixing and MCP Phase

dark fermion

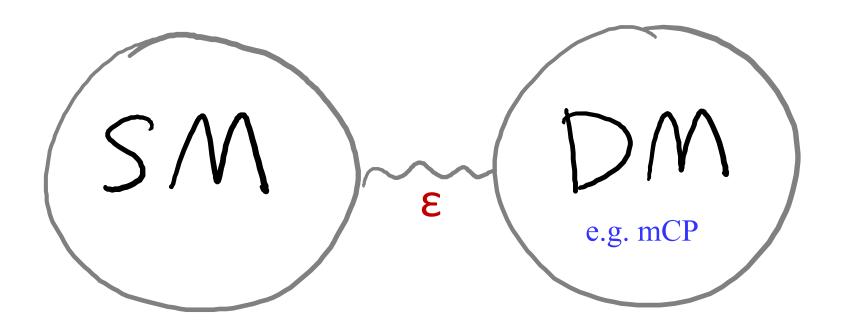


See, Holdom, 1985

$$\mathcal{L} = \mathcal{L}_{SM} - \frac{1}{4} B'_{\mu\nu} B^{\mu\nu'} - \frac{\kappa}{2} B'_{\mu\nu} B^{\mu\nu} + i \bar{\psi} (\partial \!\!\!/ + i e' B' + i M_{mCP}) \psi$$

- New Fermion ψ charged under U(1)'
- Field redefinition into a more convenient basis for massless B',  $B' \to B' + \kappa B$
- new fermion acquires an small EM charge Q (the charge of mCP  $\psi$ ):  $Q = \kappa e' \cos \theta_W \quad \epsilon \equiv \kappa e' \cos \theta_W / e.$

### The Rise of Dark Sector

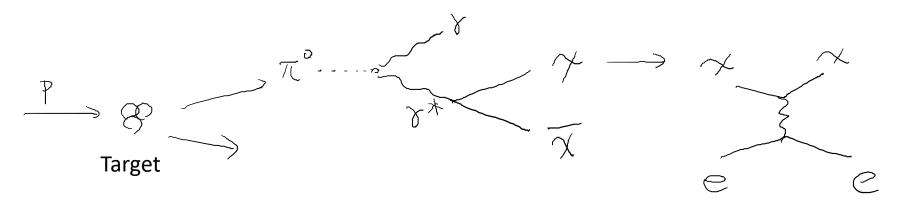


### **IMPORTANT NOTE**

- Our search is simply a search for particles (fermion  $\chi$ ) with  $\{$ mass, electric charge $\} = \{m_\chi, \epsilon e\}$
- Minimal theoretical inputs/parameters
- mCPs do not have to be DM in our searches
- The bounds we derive still put constraints on DM as well as dark sector scenarios.
- Not considering bounds on dark photon (not necessary for mCP particles)
- Similar bound/sensitivity applies to scalar mCPs
- There are additional motivations to search for "pure" MCP!

### Millicharged Particle: Signature

### MCP (or general light DM): production & detection



- production:meson decays
- - ☐ Heavy mesons are important for higher mass mCP's in high enough beam energy
  - Important and often neglected!

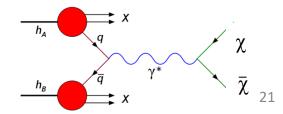
detection:scattering electron

$$BR(\pi^0 \rightarrow 2\gamma) = 0.99$$

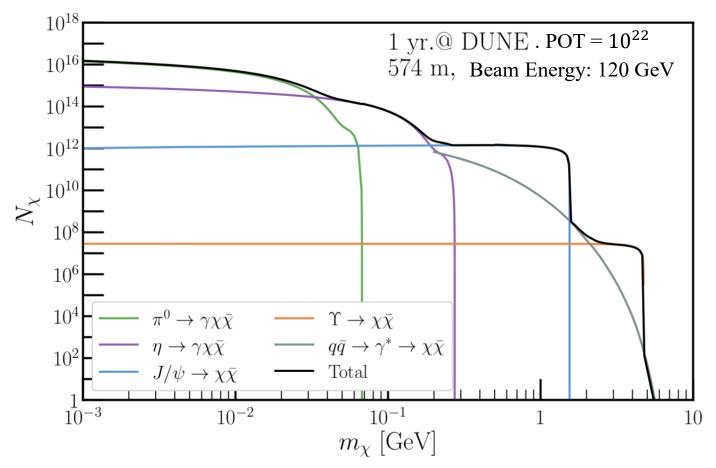
$$BR(\pi^0 \rightarrow \forall e^- e^+) = 0.01$$

BR(
$$\pi^0 \rightarrow e^- e^+$$
) = 6 \* 10<sup>-6</sup>

$$BR(J/\psi \to e^- e^+) = 0.06$$



### MCP Production/Flux



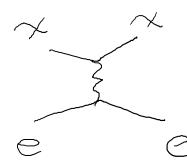
- We use PYTHIA to generate neutral meson Dalitz or direct decays from the pp collisions and rescale by considering,  $\text{BR}(\mathcal{M} \to \chi \bar{\chi}) \approx \epsilon^2 \times \text{BR}(\mathcal{M} \to X e^+ e^-) \times f\left(\frac{m_\chi}{M}\right)$ ,
- M: mass of the parent meson, X:additional particles,  $f(m_\chi/M)$ : phase space factor
- We also include Drell-Yan production for the high mass MCPs (see arXiv:1812.03998)

### MCP Detection: electron scattering

- Light mediator: the total cross section is dominated by the small  $Q^2$  contribution, we have  $\sigma_{\rm e\chi} = 4\pi \ \alpha^2 \epsilon^2/Q_{min}^2$ .
- lab frame:  $Q^2 = 2m_e (E_e m_e)$ ,  $E_e m_e$  is the electron recoil energy.
- Expressed in recoil energy threshold,  $E_e^{\,(min)}$ , we have

$$\sigma_{e\chi} = 2.6 \times 10^{-25} \text{cm}^2 \times \epsilon^2 \times \frac{1 \text{ MeV}}{E_e^{(\text{min})} - m_e}.$$

- Sensitivity greatly enhanced by accurately measuring low energy electron recoils for mCP's & light dark matter - electron scattering,
- See e.g., Magill, Plestid, Pospelov, YT, <u>1806.03310</u> & deNiverville, Frugiuele, <u>1807.06501</u> (for sub-GeV DM)



### MCP @ Neutrino Detectors

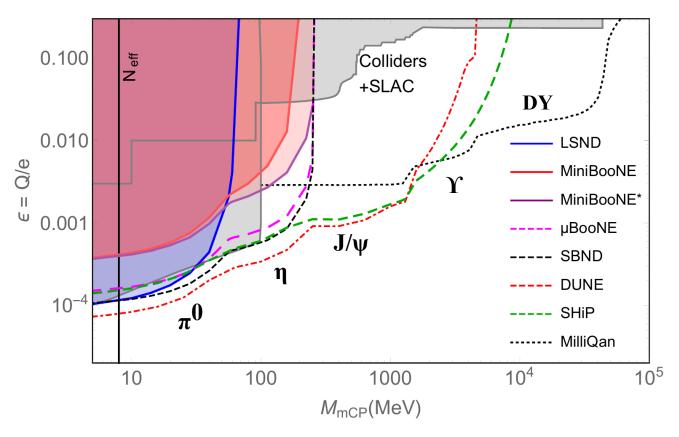
### MCP Signals in Neutrino Detector

• signal events  $n_{\text{event}}$ 

$$n_{
m event} \simeq \sum_{
m Energies} N_\chi(E_i) imes rac{N_e}{
m Area} imes \sigma_{e\chi}(E_i; \ m_\chi) imes {\cal E}.$$
 detection efficiency

- $N_{\chi}(E_i)$ : number of mCPs with energy  $E_i$  arriving at the detector.
- $N_e$ : total number of electrons inside the active volume of the detector
- Area: active volume divided by the average length traversed by particles inside the detector.
- $\sigma_{e\chi}(E_i)$ : detection cross section consistent with the angular and recoil cuts in the experiment
- Here,  $n_{\text{event}} \propto \varepsilon^4$ .  $\varepsilon^2$  from  $N_x$  and  $\varepsilon^2$  from  $\sigma_{\text{ex}}$
- Throughout this paper, we choose a credibility interval of  $1 \alpha = 95\%$  (~ 2 sigma)
- Roughly,  $\varepsilon_{sensitivity} \propto E_{e,R,min}^{1/4} Bg^{1/8}$

### **Preview: Sensitivity and Contributions**



- MilliQan: Haas, Hill, Izaguirre, Yavin, (2015), + (LOT arXiv:1607.04669)
- $N_{eff}$ : Boehm, Dolan, and McCabe (2013)
- Colliders/Accelerator: Davidson, Hannestad, Raffelt (2000) + refs within.
- SLAC mQ: Prinz el al, PRL (1998); Prinz, Thesis (2001).

### **Background for Future Measurements**

- Single-electron background for ongoing/future experiments for MicroBooNE, SBND, DUNE, and SHiP?
- Two classes of backgrounds:
  - From neutrino fluxes (calculable),
     [i.e. ve → ve and vn → ep], greatly reduced by
     maximum electron recoil energy cuts E<sub>e</sub>(max), because no low Q² enhancement (through W/Z, not γ)
  - 2) Other sources such as beam related: dirt related events, mis-id particles external: cosmics, multiply a factor of the neutrino-caused background to account for these background
  - 3) More on background control: **Harnik**, **Liu**, **Ornella**,**1902.03246**MeV-Scale Physics in Lar-TPC: **ArgoNeuT**, **1810.06502** (**Ivan Lepetic +**)

### **Summary Table**

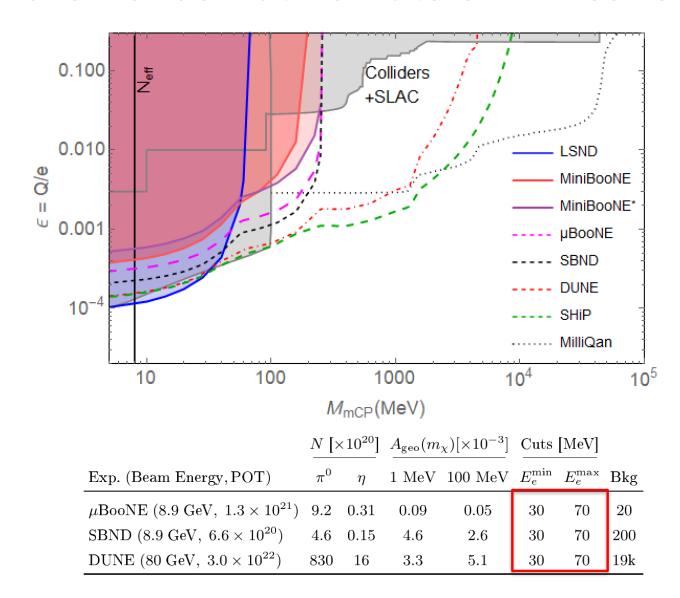
	$N \times 10^{20}$		$A_{\rm geo}(m_\chi)[\times 10^{-3}]$		$\operatorname{Cuts} \left[ \operatorname{MeV} \right]$		
Exp. (Beam Energy, POT)	$\pi^0$	$\eta$	$1~{ m MeV}$	$100~{ m MeV}$	$E_e^{\min}$	$E_e^{\max}$	Bkg
LSND (0.8 GeV, $1.7 \times 10^{23}$ )	130		20		18	52	300
mBooNE (8.9 GeV, $2.4 \times 10^{21}$ )	17	0.56	1.2	0.68	130	530	2k
mBooNE* (8.9 GeV, $1.9 \times 10^{20}$ )	1.3	0.04	1.2	0.68	75	850	0
$\mu \text{BooNE} (8.9 \text{ GeV}, 1.3 \times 10^{21})$	9.2	0.31	0.09	0.05	2	40	16
SBND (8.9 GeV, $6.6 \times 10^{20}$ )	4.6	0.15	4.6	2.6	2	40	230
DUNE (80 GeV, $3.0 \times 10^{22}$ )	830	16	3.3	5.1	2	40	19k
SHiP (400 GeV, $2.0 \times 10^{20}$ )	4.7	0.11	130	220	100	300	140

- $\varepsilon \propto E_{e.\,R,min}^{1/4} \, Bg^{1/8}$
- At LArTPC, the wire/pixel spacing is assumed to be around 3 mm, the ionization stopping power is approximately 2.5 MeV/cm: electrons with total energy larger than at least 2 MeV produce tracks long enough to be reconstructed across two wires/pixels. DUNE LArTPC ND, Using CDR config. Efficiency of 0.2 for Cherenkov detectors, 0.5 for nuclear emulsion detectors, and 0.8 for liquid argon time projection chambers.

## Recasting Existing Analysis: LSND, MiniBooNE, and MiniBooNE\* (DM Run)

- LSND: hep-ex/0101039. Measurement of electron-neutrino electron elastic scattering
- MiniBooNE: arXiv:1805.12028.
   Electron-Like Events in the MiniBooNE Short-Baseline Neutrino Experiment, combines data from both neutrino and antineutrino runs and consider a sample of 2.4 × 10<sup>21</sup>POT for which we take the single electron background to be 2.0 × 10<sup>3</sup> events and the measured rate to be 2.4 × 10<sup>3</sup>
- MiniBooNE\* (DM run): arXiv:1807.06137 (see Bishai's talk). Electron recoil analysis  $\cos \theta > 0$  is imposed (\*except for at MiniBooNE's dark matter run where a cut of  $\cos \theta > 0.99$  effectively reduces backgrounds to zero [Dharmapalan, MiniBooNE, (2012)]).
- We did not include their timing cuts in our calculations, since they were optimized by the MiniBooNE collaboration

### More Conservative Cuts on Threshold



#### Remarks

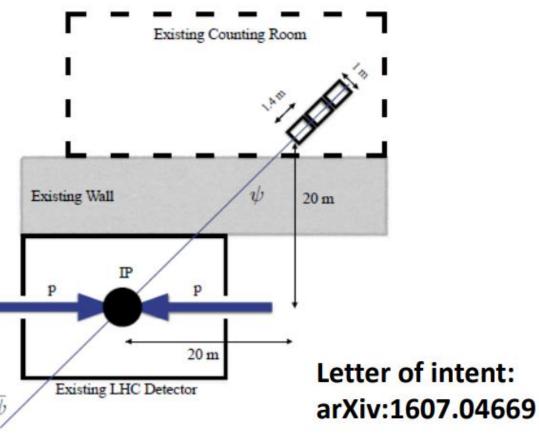
- Our technique can be applied to more generic light dark matter and other weakling interacting particles
- For mCP, or generically light dark matters with lighter mediator
  - Production from heavy neutral mesons are important (sometimes neglected in literature)
  - Signature favor low electron-recoil energy threshold
  - v background reduced by max electron-recoil energy cuts
- For more involved analysis (with your help): including
   realistic background, E<sub>e, R,min</sub> cut, etc, with Animesh, Albert,
   Jae, Yun-Tse on realistic analysis on different experiments.

### FerMINI Proposal:

Putting milliQan-Type Minicharged Particle Detector in the Fermilab Beamlines: NuMI or LBNF Independent from the DUNE ND MCP probe

#### MilliQan Detector: General Idea

- Require triple incidence in small time window (15 nanoseconds)
- With Q down to 10<sup>-3</sup> e, each
   MCP produce averagely ~ 1
   photo-electron observed per ~ 1
   meter long scintillator



**Andrew Haas, Fermilab (2017)** 

- A. Haas, C. S. Hill, E. Izaguirre, I. Yavin, arXiv:1410.6816, PRD '15
- ArXiv:1607.04669, Letter of Intent (LOT)

- Total: 1 m × 1 m (transverse plane) × 3 m (longitudinal) plastic scintillator array.
- Array oriented such that the long axis points at the CMS Interaction Point.
- The array is subdivided into 3 sections each containing 400 5 cm × 5 cm × 80 cm scintillator bars optically coupled to highgain photomultiplier (PMT).
- A triple-incidence within a 15 ns time window along longitudinally contiguous bars in each of the 3 sections will be required in order to reduce the darkcurrent noise (the dominant background).

### MilliQan: Design

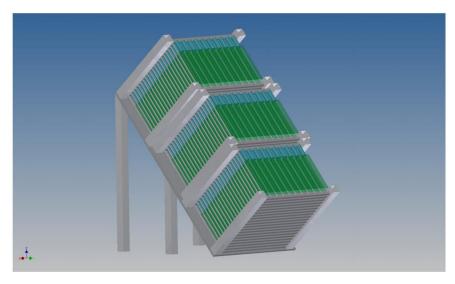
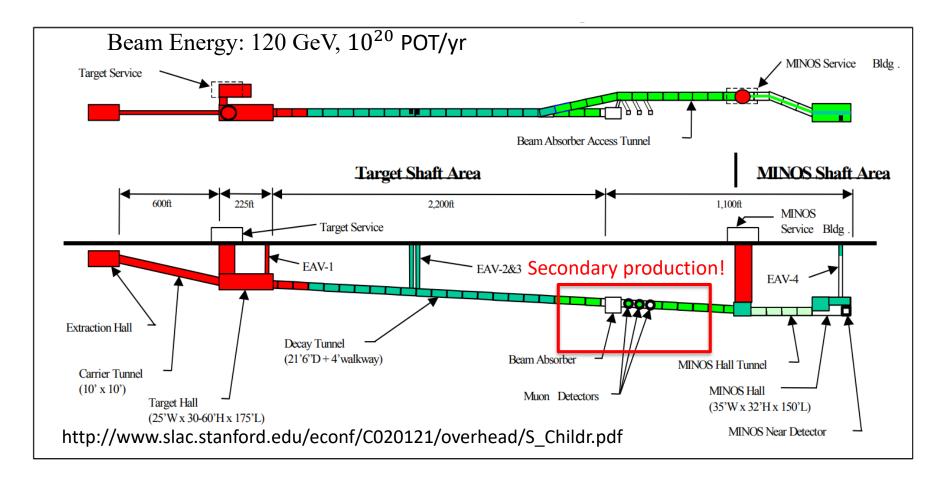


Figure from 1607.04669 (milliQan LOT)

### FerMINI:

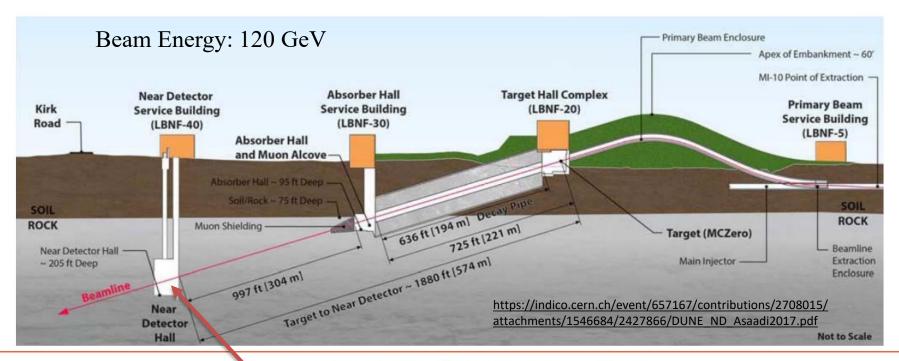
A Fermilab Search for Minicharged Particle

### Site 1: NuMI Beam & MINOS ND Hall

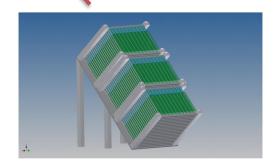


NuMI: Neutrinos at the Main Injector (See Todd's talk)
MINOS: Main Injector Neutrino Oscillation Search, ND: Near Detector
(MINERvA: Main Injector Experiment for v-A is also here)

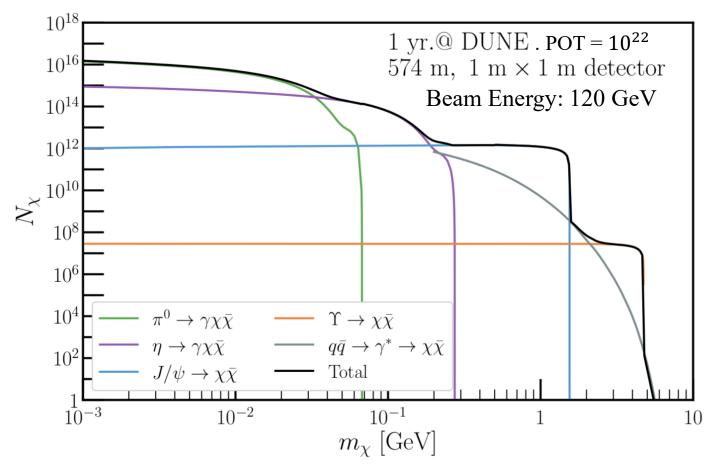
#### Site2: LBNF Beam & DUNE ND Hall



Jonathan Asaadi - University of Texas Arlington



#### **MCP Production/Flux**



- We use PYTHIA to generate neutral meson Dalitz or direct decays from the pp collisions and rescale by considering,  $\text{BR}(\mathcal{M} \to \chi \bar{\chi}) \approx \epsilon^2 \times \text{BR}(\mathcal{M} \to X e^+ e^-) \times f\left(\frac{m_\chi}{M}\right)$ ,
- M: mass of the parent meson, X:additional particles,  $f(m_\chi/M)$ : phase space factor
- We also include Drell-Yan production for the high mass MCPs.

#### Signature: Triple Incidence

 The averaged number of photoelectron (PE) seen by the detector from single MCP is:

$$N_{PE} \simeq \rho_{scint} \times \left\langle -\frac{dE}{dx} \right\rangle \times l_{scint} \times LY \times e_{det}.$$
 • LY: light yield

•  $e_{det}$ : detection efficiency

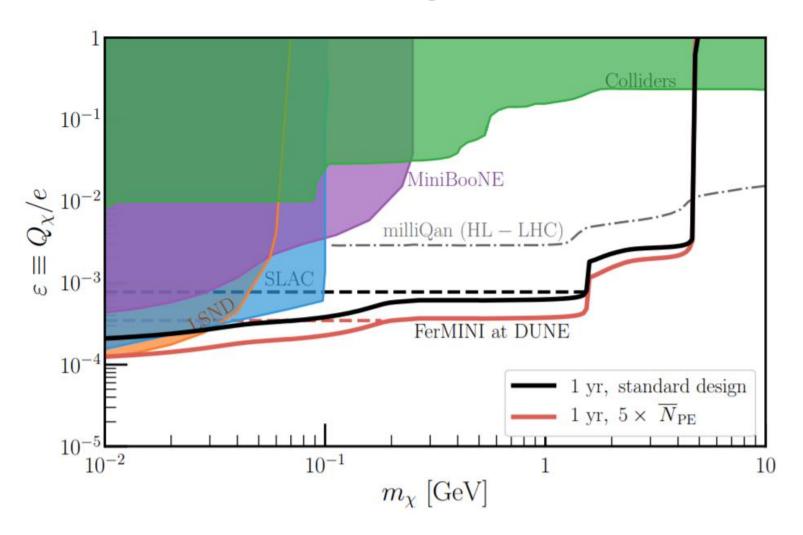
 $N_{PE} \sim \epsilon^2 \times 10^6$ , so  $\epsilon \sim 10^{-3}$  roughly gives one PE in 1 meter scintillation bar

- Based on Poisson distribution, zero event in each bar correspond to  $P_0 = e^{-N_{PE}}, \text{ so the probability of seeing triple incident of one or more}$  photoelectron is:  $P = \left(1 e^{-N_{PE}}\right)^3,$
- $N_{x,detector} = N_x \times P$ .

## Background

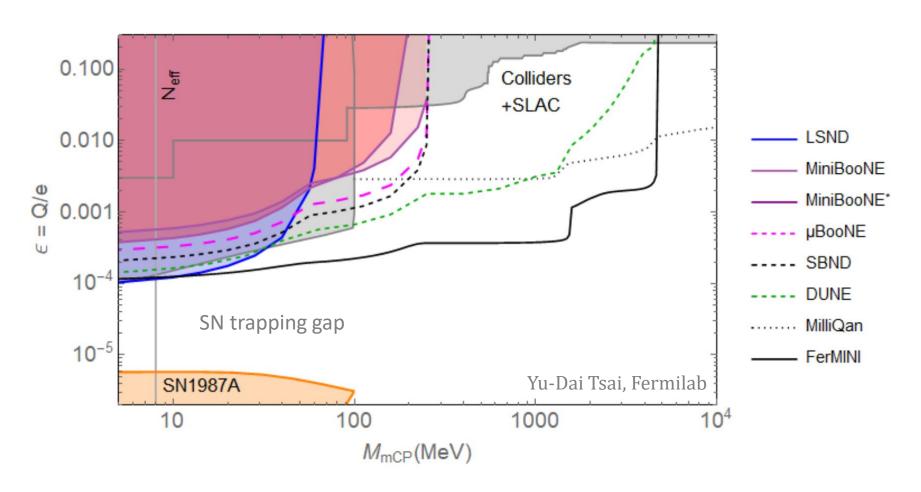
- We will discuss two major detector
   backgrounds and the reduction technique
- SM charged particles from background radiation (e.g., cosmic muons):
  - Offline veto of events with > 10 PEs
  - Offset middle detector
- Dark current: triple coincidence
  - ~ 300 events in one year of trigger-live time

## FerMINI @ DUNE



NuMI/MINOS Hall is a viable alternative site

## Compilation



One can combine the MCP detector with neutrino detector to improve sensitivity or reduce background

## Advantages: Timeliness, Low-cost, Movable, Tested, Easy to Implement, ...

- 1. LHC entering long shutdown
- 2. Can develop at NuMI/MINOS and then move to DUNE
- 3. NuMI operating, shutting down in 5 years (DO IT NOW!)
- 4. Sensitivity better than milliQan for MCP below 5 GeV and don't have to wait for HL-LHC

JOIN THE PROPOSAL! ytsai@fnal.gov

#### **Alternative Detector Setup & New Ideas**

- Combine with neutrino detector: behind, in front, or sandwich them: mixed signature
- Combine with DUNE PRISM: moving up and down
- FerMINI+DUNE 3DST
- Better scintillator material
- Can search for millicharged quarks, fermions with small electric dipole
- New ideas from you are welcomed!

#### **Looking Ahead**

- Exploring Energy Frontier of the Intensity Frontier (complementary to and before HL-LHC upgrade)
- Near-future (and almost free) opportunity
   (NuMI Facility, SBN program, DUNE Near Detector, etc.)
- Other new low-cost alternatives/proposals (~\$1M) to probe hidden particles and new forces (FerMINI is just a beginning!)
- Dark sectors in neutrino telescopes
- Cosmology-driven models!

# Thank You Thanks for the invitation again!

## Backup Slides

#### Dark Current Background @ PMT

#### Major Background Source!

- dark-current frequency to be  $v_B$ = **500 Hz** for estimation. (from 1607.04669, milliQan L.O.T.)
- For each tri-PMT set (each connect to the three connected scintillation bar), the background rate for triple incidence is  $v_B^3 \Delta t^2 = 2.8 \times 10^{-8} \text{ Hz}$ , for  $\Delta t = 15 \text{ ns}$ .
- There are 400 such set in the nominal design.
- The total background rate is 400 x 2.8 x  $10^{-8}$  ~  $10^{-5}$  Hz
- ~ 300 events in one year of trigger-live time

#### (detail) Meson Production Details

- At LSND, the  $\pi$ 0 (135 MeV) spectrum is modeled using a Burman-Smith distribution
- Fermilab's Booster Neutrino Beam (BNB):  $\pi 0$  and  $\eta$  (548 MeV) mesons.  $\pi 0$ 's angular and energy spectra are modeled by the **Sanford-Wang distribution**.  $\eta$  mesons by the Feynman Scaling hypothesis.
- SHiP/DUNE: pseudoscalar meson production using the BMPT distribution, as before, but use a beam energy of 80 GeV
- J/ $\psi$  (3.1 GeV), we assume that their energy production spectra are described by the distribution from **Gale**, **Jeon**, **Kapusta**, **PLB '99**, nucl-th/9812056.
- Upsilon, Y (9.4 GeV): Same dist., normalized by data from HERA-B, I. Abt et al., PLB (2006), hep-ex/0603015.
- Calibrated with existing data [e.g. NA50, EPJ '06, nucl-ex/0612012, Herb et al., PRL '77]. and simulations from other groups [e.g. deNiverville, Chen, Pospelov, and Ritz, Phys. Rev. D95, 035006 (2017), arXiv:1609.01770 [hepph].]

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- MiniBooNE\* (DM run): arXiv:1807.06137 (came out after our v1). Electron recoil analysis

We did not include their timing cuts in our calculations, since they were optimized by the MiniBooNE collaboration to the signal's timing profile.

#### Dark Current Background @ PMT

#### Major Background!

- We take the dark-current frequency to be  $v_B$ = 500 Hz for estimation. (from 1607.04669, milliQan L.O.T.)
- For each tri-PMT set (each connect to the three connected scintillation bar), the background rate for triple incidence is  $v_B^3 \Delta t^2 = 2.8 \times 10^{-8} \text{ Hz}$ , for  $\Delta t = 15 \text{ ns}$ .
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- The total background rate is 400 x 2.8 x  $10^{-8}$  ~  $10^{-5}$  Hz
- ~ 300 events in one year of trigger-live time

#### **Number of photoelectrons (PEs)**

For moderately small epsilon and heavy enough MCP (>> electron mass),
 one can use Bethe equation to estimate average energy loss.

$$\left\langle -\frac{dE}{dx} \right\rangle = Kz^2 \frac{Z}{A} \frac{1}{\beta^2} \left[ \frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\text{max}}}{I^2} - \beta^2 - \frac{\delta(\beta \gamma)}{2} \right] .$$

z charge number of incident particle

Z atomic number of absorber

A atomic mass of absorber  $g \text{ mol}^{-1}$ 

 $K = 4\pi N_A r_e^2 m_e c^2$  0.307 075 MeV mol<sup>-1</sup> cm<sup>2</sup>

(Coefficient for dE/dx)

 $I \quad \text{mean excitation energy} \qquad \text{eV (Nota bene!)} \qquad W_{\text{max}} = \frac{2m_e c^2 \beta^2 \gamma^2}{1 + 2\gamma m_e / M + (m_e / M)^2} .$ 

 $\delta(\beta\gamma)$  density effect correction to ionization energy loss

- M: charged particle mass
- For very small epsilon (related to the finite length effect), one have to consider most probable energy deposition & consider landau distribution for the energy transfer.

## Background: Detector & Beam Related

## **Detector Background**

- We will discuss two major detector
   backgrounds and the reduction technique
- SM charged particles from background radiation (e.g., cosmic muons):
   offline veto + offset middle detector
- Dark current: triple incidence
- See 1607.04669 (milliQan LOT)

## Reduce background from SM

- Reduce background from SM charged particles
- Offline-vetoes of large-PE events: Offline veto of events with > 10 PEs
   Background events (such as cosmic muons) produce a large number of PEs
   (typically > 10<sup>3</sup> PEs ) would be vetoed.
- Offset the middle detector array: Charged standard model particles
  skimming the edge, producing only a low number of PE. Offsetting the
  middle detector, or making it slightly smaller/larger, would prevent these
  types of events from producing signature in all three arrays.
- These would also reduce the charged particles directly or indirectly from the beam, e.g., vµ from the beam striking the detector (or nearby rock) and producing a muon

#### Dark Current Background @ PMT

#### Major Background!

- We take the dark-current frequency to be  $v_B$ = 500 Hz for estimation. (from 1607.04669, milliQan L.O.T.)
- For each tri-PMT set (each connect to the three connected scintillation bar), the background rate for triple incidence is  $v_B^3 \Delta t^2 = 2.8 \times 10^{-8} \text{ Hz}$ , for  $\Delta t = 15 \text{ ns}$ .
- There are 400 such set in the nominal design.
- The total background rate is 400 x 2.8 x  $10^{-8}$  ~  $10^{-5}$  Hz
- ~ 300 events in one year of trigger-live time

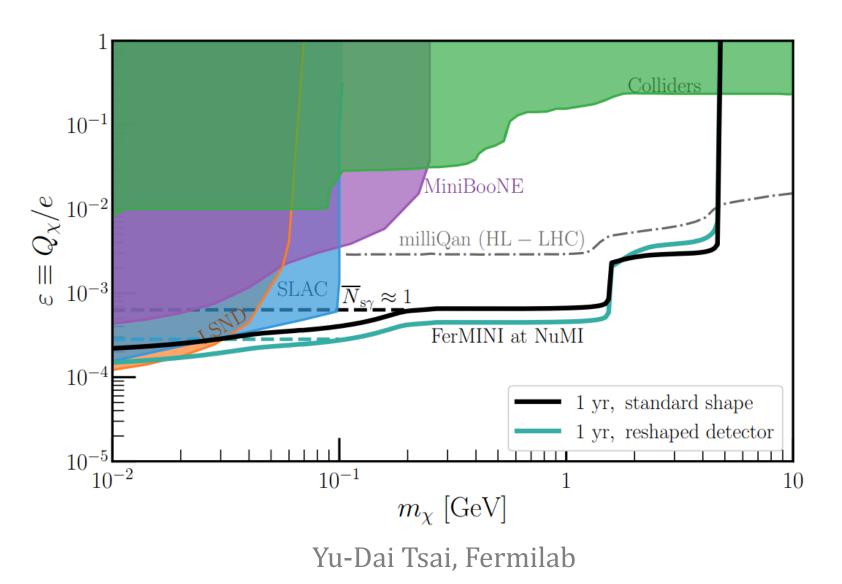
#### Beam Related Background:

- Beam produced charged particle went through several shielding already, including absorber and rocks.
- Each beams have muon monitors.
- Determine the SM charged particle rate on site
- Remaining beam / dirt / rock produced charged particle: vetoed similar to the previous veto of cosmic muons.
- Neutrino produced background:  $O(10^{-19})$ , negligible.
- To be conservative, we assume the beam related background ≈ dark current background for our sensitivity determination.

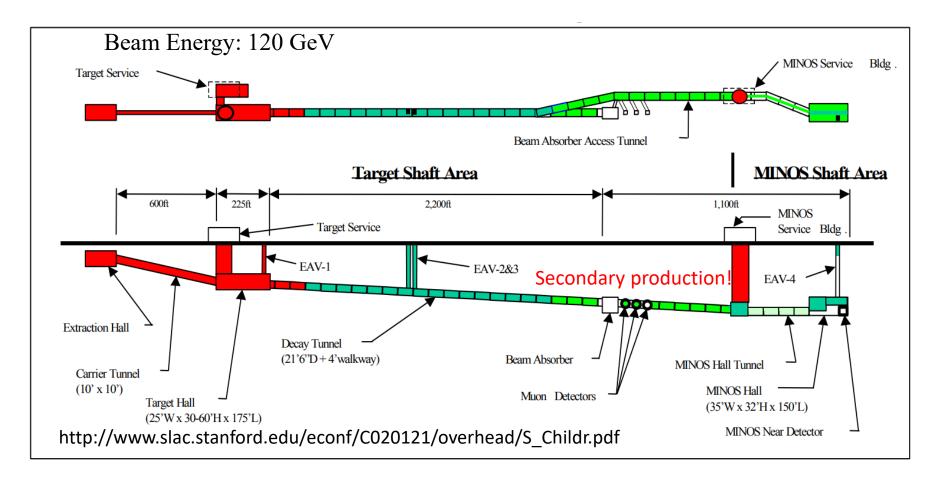
### FerMINI: Increasing scintillation photons

- Elongating the scintillator bar does not affect the background from dark current (basically determined by the number of PMTs)
- So we estimate the sensitivity of FerMINI at DUNE for five times larger scintillation capability
- And estimate the sensitivity of FerMINI at NuMI for five time more scintillation capability but five times less scintillator bar-PMT sets (actually reduce dark current background!)

## FerMINI @ MINOS



### NuMI Beam & MINOS ND Hall

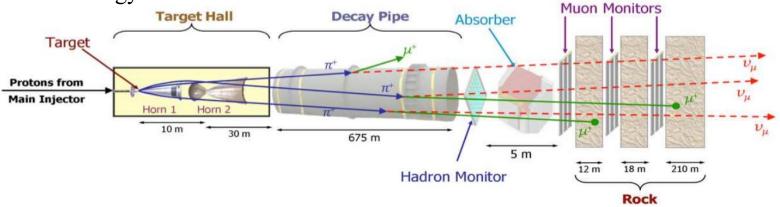


NuMI: Neutrinos at the Main Injector

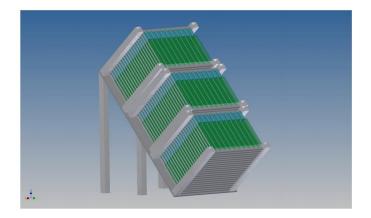
MINOS: Main Injector Neutrino Oscillation Search

## FerMINI @ NuMI-MINOS Hall

Beam Energy: 120 GeV







# FerMINI: Discussions & Alternative Designs

## Detection Limitation: $N_{photon} \leq 1$

- Define:  $\varepsilon_{low}$  as  $N_{sintilator\ photon} = 1$
- Roughly around or below this, one really have to worry about scintillator performance
- One can elongate the scintillator or consider alternative materials to help.

Material	Photons/keV	Density (g/cm³)	* Length needed (cm)	Speed (ns)	Cost for 5x5 cm (\$)	Notes
Plastic BC408	10	1.03	145	~2	~200	Current choice
Nal	38	3.67	11	~230	~800	Slow, fragile
LaBr3(Ce)	63	5.08	5	~16	~3000	Radioactive
Liquid Xe	62	2.95	8	~2 / ~34	~1000?	Cryogenic, ultraviolet

Andy Haas, Fermilab, <u>2017</u>

<sup>\*</sup> Length needed to get 3 photons for charge 1/1000 e

#### **Alternatives (Straightforward)**

- 1. Quadruple incidence: further background reduction, sacrifice event rate but potentially gain better control of background, reduce the background naively by **10**-5
- 1. Basically zero background experiment?
- 2. Different lengths for each detectors
- 3. Different materials:

Material	Photons/keV	Density (g/cm³)	* Length needed (cm)	Speed (ns)	Cost for 5x5 cm (\$)	Notes
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Andy Haas, Fermilab, <u>2017</u>

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#### **Summary Table**

	$N \times 10^{20}$		$A_{\rm geo}(m_\chi)[\times 10^{-3}]$		$\underline{\mathrm{Cuts}\;[\mathrm{MeV}]}$		
Exp. (Beam Energy, POT)	$\pi^0$	$\eta$	$1~{ m MeV}$	$100~{\rm MeV}$	$E_e^{\min}$	$E_e^{\max}$	Bkg
LSND $(0.8 \text{ GeV}, 1.7 \times 10^{23})$	130		20		18	52	300
mBooNE (8.9 GeV, $2.4 \times 10^{21}$ )	17	0.56	1.2	0.68	130	530	2k
mBooNE* (8.9 GeV, $1.9 \times 10^{20}$ )	1.3	0.04	1.2	0.68	75	850	0
$\mu \text{BooNE } (8.9 \text{ GeV}, \ 1.3 \times 10^{21})$	9.2	0.31	0.09	0.05	2	40	16
SBND $(8.9 \text{ GeV}, 6.6 \times 10^{20})$	4.6	0.15	4.6	2.6	2	40	230
DUNE (80 GeV, $3.0 \times 10^{22}$ )	830	16	3.3	5.1	2	40	19k
SHiP (400 GeV, $2.0 \times 10^{20}$ )	4.7	0.11	130	220	100	300	140

- $\varepsilon \propto E_{e,R,min}^{1/4} Bg^{1/8}$
- DUNE LArTPC ND, Using CDR configuration Efficiency of 0.2 for Cherenkov detectors, 0.5 for nuclear emulsion

detectors, and 0.8 for liquid argon time projection chambers.