## Shear and pressure

## distributions in the proton

## from lattice QCD

## Will Detmold

work wth Phiala Shanahan
PRL (2019) arXiv:I8I0.07589
PRD (2019) arXiv: 1810.04626

## Gluon structure

## Gluons offer a new window on nuclear structure

- Past 60+ years: detailed view of quark structure of nucleons
- Gluon structure also important - Unpolarised gluon PDF dominant at small longitudinal momentum fraction
- Other aspects of gluon structure relatively unexplored

Parton distributions in the proton


## Gluon angular momentum

- Gluon helicity much less well constrained
- Major focus of RHIC-spin program
- Asymmetries in polarised $p p \rightarrow \pi X, D X, B X$, jets
- Orbital angular momentum of gluons even less understood
- GluonTMDs
- Major motivation for EIC




## Gluon structure

How much do gluons contribute to the proton's

- Momentum
- Spin
- Mass
- D-term

What are the gluon distributions in a proton


- PDFs, GPDs,TMDs
- Pressure, Shear
- Gluon 'radius/
radii'

How is the gluon structure of a proton modified in a nucleus

\author{

- Gluon 'EMC' effect <br> - Exotic glue
}


## Gluon structure

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How is the gluon structure of a proton modified
in a nucleus


Exotic glue

## Energy-momentum tensor

Many of these properties derived from Energy-Momentum Tensor (conserved Noether current associated with Lorentz translations)

Matrix elements of traceless gluon EMT for spin-half nucleon:

$$
\begin{array}{r}
\left\langle p^{\prime}, s^{\prime}\right| G_{\{\mu \alpha}^{a} G_{\nu\}}^{a \alpha}|p, s\rangle=\bar{U}\left(p^{\prime}, s^{\prime}\right)( \\
\left.A_{g}(t) \gamma_{\{\mu} P_{\nu\}}+B_{g}(t) \frac{i P_{\{\mu} \sigma_{\nu\} \rho} \Delta^{\rho}}{2 M_{N}}+D_{g}(t) \frac{\Delta_{\{\mu} \Delta_{\nu\}}}{4 M_{N}}\right) U(p, s) \\
\text { Generalised gluon form factors } \quad \Delta_{\mu}=p_{\mu}^{\prime}-p_{\mu} P_{\mu}=\left(p_{\mu}+p_{\mu}^{\prime}\right) / 2 \quad t=\Delta^{2}
\end{array}
$$

- Three generalised gluon form factors $A_{g}(t), B_{g}(t), D_{g}(t)$
- Sum rules with quark pieces in forward limit
- Momentum fraction $A_{a}(0)=\langle x\rangle_{a} \rightarrow \sum_{a=q, g} A_{a}(0)=1$
- Spin $J_{a}(t)=\frac{1}{2}\left(A_{a}(t)+B_{a}(t)\right)$

$$
\sum_{a=q, g} J_{a}(0)=\frac{1}{2}
$$

- D-terms $D_{a}(0)$ unknown but equally fundamental!


## D-term

D-term GFF encodes the pressure and shear distributions in the nucleon (Breit frame)

$$
\begin{gathered}
s(r)=-\frac{r}{2} \frac{d}{d r} \frac{1}{r} \frac{d}{d r} \widetilde{D}(r), \quad p(r)=\frac{1}{3} \frac{1}{r^{2}} \frac{d}{d r} r^{2} \frac{d}{d r} \widetilde{D}(r), \\
\widetilde{D}(r)=\int \frac{d^{3} \vec{p}}{2 E(2 \pi)^{3}} e^{-i \vec{p} \cdot \vec{r}} D\left(-\vec{p}^{2}\right)
\end{gathered}
$$

- Quark and gluon shear forces individually well-defined (i.e., scaledependent partial contributions $s_{q, g}(r)$
- Pressure defined only for the total system (pieces depend also on GFFs related to the trace terms of the EMT that cancel in the sum)


## Generalised parton distributions

## GFFs correspond to lowest moments of GPDs:

$$
\begin{aligned}
\int_{0}^{1} \mathrm{~d} x H_{g}(x, \xi, t)=A_{g}(t)+\xi^{2} D_{g}(t), & \int_{0}^{1} \mathrm{~d} x E_{g}(x, \xi, t)=B_{g}(t)-\xi^{2} D_{g}(t) \\
\int_{-1}^{1} \mathrm{~d} x x H_{q}(x, \xi, t)=A_{q}(t)+\xi^{2} D_{q}(t), & \int_{-1}^{1} \mathrm{~d} x x E_{q}(x, \xi, t)=B_{q}(t)-\xi^{2} D_{q}(t)
\end{aligned}
$$

Quark GPDs: constraints from JLab, HERA, COMPASS, by DVCS, DVMP, future improvements from JLab 12 GeV

- Gluon GPDs: almost unknown from experiment, future constraints are a central goal of EIC

Leading twist nucleon gluon GPDs:
Gluon field-
strength tensor

$$
\begin{gathered}
\int_{-\infty}^{\infty} \frac{d \lambda}{2 \pi} e^{i \lambda x}\left\langle p^{\prime}, s^{\prime}\right| G_{a}^{\{\mu \alpha}\left(-\frac{\lambda}{2} n\right)\left[\mathcal{U}_{\left[-\frac{\lambda}{2} n, \frac{\lambda}{2} n\right]}^{(A)}\right]_{a b}^{G_{b \alpha}^{\nu\}}\left(\frac{\lambda}{2} n\right)|p, s\rangle} \quad t=\Delta^{2} n^{2}=0 \quad \xi=-\frac{1}{2} n \cdot \Delta / n \cdot P \\
=\frac{1}{2}\left(H_{g}(x, \xi, t) \bar{U}\left(p^{\prime}, s^{\prime}\right) P^{\{\mu} \gamma^{\nu\}} U(p, s)+E_{g}(x, \xi, t) \bar{U}\left(p^{\prime}, s^{\prime}\right) \frac{P^{\{\mu} i \sigma^{\nu\} \alpha} \Delta_{\alpha}}{2 M} U(p, s)\right)+\ldots, \\
\text { GPDs(Bjorken x, skewness, mom transfer) }
\end{gathered}
$$

## D-term from JLab DVCS

Recent experimental determination of DVCS D-term and extraction of proton pressure distribution
V. D. Burkert, L. Elouadrhiri, and F. X. Girod, Nature 557, 396 (20I8)

$$
s(r)=-\frac{r}{2} \frac{d}{d r} \frac{1}{r} \frac{d}{d r} \widetilde{D}(r), \quad p(r)=\frac{1}{3} \frac{1}{r^{2}} \frac{d}{d r} r^{2} \frac{d}{d r} \widetilde{D}(r)
$$

- Strong repulsive pressure near the centre of the proton
- Binding pressure at greater distances.
- Peak pressure near the centre ~1035 Pascal, greater than pressure estimated for neutron stars
- Key assumptions: gluon D-term same as quark term, tripole form factor model, $D_{u}(t, \mu)=D_{d}(t, \mu)$


## Test assumptions in pressure extraction



Radial pressure distribution


## Gluon structure

## First-principles QCD calculations

$\rightarrow$ QCD benchmarks and predictions ahead of experiment


## Lattice QCD

## Numerical first-principles approach to non-perturbative QCD

- Discretise equations of QCD onto space-time grid
- Calculate physical quantities
- Take limit of vanishing discretisation afterwards



## Lattice QCD

## Numerical first-principles approach to non-perturbative QCD

QCD equations $\Longleftrightarrow$ integrals over the values of quark and gluon fields on each site/link (QCD path integral)

- $10^{12}$ variables (for state-of-the-art)

- Evaluate by importance sampling
- Paths near classical action dominate
- Calculate physics on a set (ensemble) of samples of the quark and gluon fields


## Lattice QCD

## Calculate the nucleon mass

- Create three quarks, annihilate them far from source

Tie together creation and annihilation operators in all possible ways (contractions)

- QCD path integral adds quark antiquark pairs and gluons automatically
- Measure exponentially decaying correlation to extract mass

$$
M(t)=\ln \left[\frac{C(t)}{C(t+1)}\right] \xrightarrow{t \rightarrow \infty} E_{0}
$$




## Lattice QCD

- Calculations use world's largest computers
- Many millions of CPU/GPU hours
- Specifically designed processors for QCD (QCDOC precursor of BlueGene computers)



## Lattice QCD

- Ground state hadron spectrum reproduced
- p-n mass splitting reproduced

- Predictions for new states with controlled uncertainties



## Lattice QCD Matrix Elements

## Calculate matrix elements

- Create three quarks (correct quantum numbers) at a source and annihilate the three quarks at sink far from source
- Insert operator at intermediate timeslice

- Remove time-dependence by dividing out with two-point correlators:

$$
\frac{C_{3}\left(t, \tau, \overrightarrow{\left.p^{\prime}, \vec{q}\right)}\right.}{C_{2}\left(t-\tau, p^{\prime}\right) C_{2}(\tau, p)} \xrightarrow{t \rightarrow \infty}\left\langle N\left(p^{\prime}\right)\right| \mathcal{O}(q)|N(p)\rangle
$$

## Gluon GFFs from LQCD

## Construct system of equations for generalised gluon form factors

## Ratios of 3 pt and 2 pt correlation functions:

$$
\begin{gathered}
R_{s ;\{;, i}\left(\vec{p}, \vec{p}^{\prime}, t_{f}, \tau\right)=\frac{C_{s ; \mathcal{P}^{2}, i}^{3 \mathrm{p}}\left(\vec{p}, \vec{p}^{\prime}, t_{f}, \tau\right)}{C_{s}^{2 \mathrm{pt}}\left(\vec{p}^{\prime}, t_{f}\right)} \sqrt{\frac{C_{s}^{2 \mathrm{pt}}\left(\vec{p}, t_{f}-\tau\right) C_{s}^{2 \mathrm{pt}}\left(\vec{p}^{\prime}, t_{f}\right) C_{s}^{2 \mathrm{pt}}\left(\vec{p}^{\prime}, \tau\right)}{C_{s}^{\mathrm{pt}}\left(\vec{p}^{\prime}, t_{f}-\tau\right) C_{s}^{2 \mathrm{pt}}\left(\vec{p}, t_{f}\right) C_{s}^{2 \mathrm{pt}}(\vec{p}, \tau)}} \stackrel{t_{f} \gg \tau \gg 0}{\longrightarrow} \frac{\operatorname{Tr}\left[\Gamma_{s}\left(p^{\prime}+M_{N}\right) \mathcal{F}_{i}\left[A_{g}, B_{g}, D_{g}\right]\left(p+M_{N}\right)\right]}{8 \sqrt{E_{\vec{p}}^{(N)} E_{\overrightarrow{p^{\prime}}}^{(N)}\left(E_{\vec{p}}^{(N)}+M_{N}\right)\left(E_{\vec{p}^{\prime}}^{(N)}+M_{N}\right)}} \\
\mathcal{F}_{\mu \nu}\left[A_{g}, B_{g}, D_{g}\right]=A_{g}(t) \gamma_{\{\mu} P_{\nu\}}+B_{g}(t) \frac{i P_{\{\mu} \sigma_{\nu\} \rho} \Delta^{\rho}}{2 M_{N}}+D_{g}(t) \frac{\Delta_{\{\mu} \Delta_{\nu\}}}{4 M_{N}}
\end{gathered}
$$

## Generalised gluon form factors

$$
\Delta_{\mu}=p_{\mu}^{\prime}-p_{\mu} \quad P_{\mu}=\left(p_{\mu}+p_{\mu}^{\prime}\right) / 2 . \quad t=\Delta^{2}
$$

- Nucleon spin up/down: $\Gamma_{s= \pm 1}$
- Sink and operator momenta:

$$
\begin{aligned}
& \left|\overrightarrow{p^{\prime}}\right|^{2} \leq 5(2 \pi / L)^{2} \\
& |\vec{\Delta}|^{2} \leq 18(2 \pi / L)^{2}
\end{aligned}
$$

- Operator index choices: two different irreducible representations of $\mathrm{H}(4)$

$$
\begin{aligned}
& \mathcal{O}_{i=\{11, \ldots, 6\}}^{\tau_{3}^{(6)}}=\left\{\frac{(-i)^{\delta_{\nu 0}}}{\sqrt{2}}\left(\mathcal{O}_{\mu \nu}+\mathcal{O}_{\nu \mu}\right), \quad 0 \leq \mu<\nu \leq 3\right\} \\
& \mathcal{O}_{1}^{\tau_{1}^{(3)}}=\frac{1}{2}\left(\mathcal{O}_{11}+\mathcal{O}_{22}-\mathcal{O}_{33}+\mathcal{O}_{00}\right), \cdots,
\end{aligned}
$$

## Gluon GFFs from LQCD

One ensemble, $\mathrm{m}_{\pi} \sim 450 \mathrm{MeV}$ (physical masses running now)

| $L / a$ | $T / a$ | $\beta$ | $a m_{l}$ | $a m_{s}$ | $a(\mathrm{fm})$ | $L(\mathrm{fm})$ | $T(\mathrm{fm})$ | $m_{\pi}(\mathrm{MeV})$ | $m_{K}(\mathrm{MeV})$ | $m_{\pi} L$ | $m_{\pi} T$ | $N_{\text {cfg }}$ | $N_{\text {meas }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 32 | 96 | 6.1 | -0.2800 | -0.2450 | $0.1167(16)$ | 3.7 | 11.2 | $450(5)$ | $596(6)$ | 8.5 | 25.6 | 2821 | 203 |

Clean plateaus in effective masses for $\left|\overrightarrow{p^{\prime}}\right|^{2} \leq 5(2 \pi / L)^{2}$


## Gluon GFFs from LQCD

## PION: Clean signals in 3pt/2pt ratios (examples)







Plateau fits
---- Solution of system of equations projected back
smeared-point smeared-smeared

## Gluon GFFs from LQCD

## PION: Clean signals in 3pt/2pt ratios (examples)






$t_{f}=13$ (dark points), $t_{f}=18$ (pale points)

Plateau fits
--- - Solution of system of equations projected back
smeared-point
smeared-smeared

## Gluon GFFs from LQCD

## NUCLEON: Clean signals in 3pt/2pt ratios (examples)




Plateau fits
-=- - Solution of system of equations projected back

$\tau=4$ (dark points), $\tau=6$ (pale points)
smeared-point smeared-smeared

## Gluon GFFs from LQCD

## NUCLEON: Clean signals in 3pt/2pt ratios (examples)







Plateau fits
-=- - Solution of system of equations projected back
$1 t_{f}=12$ (dark points), $t_{f}=14$ (pale points)

## smeared-point

 smeared-smeared
## Gluon GFFs from LQCD

Solve system of equations for GFFs in bins in $t=\left(p^{\prime}-p\right)^{2}$

Pion


Nucleon


- Colour coding: three momentum transfer $\vec{\Delta}^{2}=\left(\vec{p}^{\prime}-\vec{p}\right)^{2}$
- Point size $\propto$ number of three-momenta at that $\vec{\Delta}^{2}$
- Grey bands: bins in $t$


## Renormalisation

## Non-perturbative RI-MOM renormalisation of gluon operator

- Mixing with quark operator neglected

Found to be small in lattice PT e.g., Alexandrou et al., | $6 \mid 1.0690$ |

- One-loop perturbative matching to $\overline{\mathrm{MS}}$ scheme:Yang et al., $16 \mid 2.02855$

$$
\mathcal{O}^{\overline{\mathrm{MS}}}\left(\mu^{2}\right)=Z_{\mathcal{O}}^{\overline{\mathrm{MS}}}\left(\mu^{2}\right) \mathcal{O}^{\mathrm{latt}}=\mathcal{R}^{\overline{\mathrm{MS}}}\left(\mu^{2}, \mu_{R}^{2}\right) Z_{\mathcal{O}}^{\mathrm{RI}-\mathrm{MOM}}\left(\mu_{R}^{2}\right) \mathcal{O}^{\text {latt }}
$$

Calculate RI-MOM coefficient using Landau-gauge fixed gluon 2pt function

$$
\begin{aligned}
& \left.\left(Z_{\hat{\mathcal{O}}}^{\mathrm{RI}-\mathrm{MOM}}\left(\mu_{R}^{2}\right)\right)^{-1}=\frac{\left.4 p^{2}\left\langle\hat{\mathcal{O}}_{\alpha \beta} \operatorname{Tr}\left[A_{\tau}(p) A_{\tau}(-p)\right]\right\rangle\right\rangle}{\Lambda_{\hat{\mathcal{O}}}^{\text {tre }}(p)\left\langle\operatorname{Tr}\left[A_{\tau}(p) A_{\tau}(-p)\right]\right\rangle} \right\rvert\, \begin{array}{c}
p^{2}=\mu_{R}^{2} \\
\tau \neq \alpha \neq \beta \\
, p_{\tau}=0
\end{array} \\
& \Lambda_{\mathcal{O}}^{\text {tree }}(p)=\left\langle\hat{\mathcal{O}}_{\alpha \beta}^{\mathfrak{R}} \operatorname{Tr}\left[A_{\tau}(p) A_{\tau}(-p)\right]\right\rangle_{\text {amp }}^{\mathrm{tree}^{\text {tre }}} \text {. }
\end{aligned}
$$

(1) Wilson-flowed gluon 2pts

这 No flow in 2pts


## Gluon GFFs from LQCD






Uncertainties from renormalisation not shown

## Gluon GFFs from LQCD











Cross-sections: GFF not shown in each projection taken to its central value

## LQCD Pion GFFs

## Pion gluon GFFs $\mathrm{m}_{\boldsymbol{\pi}} \sim 450 \mathrm{MeV}$

Solve system of equations simultaneously for both hypercubic irreps for each binned fourmomentum transfer

Dipole-like fall-off with momentum transfer

- Momentum fraction $A_{a}(0)=\langle x\rangle_{a}$

$$
\rightarrow \sum_{a=q, g} A_{a}(0)=1
$$

- D-terms $D_{a}(0)$ related to pressure and shear distributions




## LQCD Pion GFFs

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Dipole-like fall-off with momentum transfer

- Momentum fraction $A_{a}(0)=\langle x\rangle_{a}$ $\longrightarrow \sum_{a=q, g} A_{a}(0)=1$
- D-terms $D_{a}(0)$ related to pressure and shear distributions
gluon: Shanahan, Detmold, PRD (2019)
quark: Brommel Ph.D. thesis (2007) $\mathrm{m}_{\pi} \sim 840 \mathrm{MeV}$




## LQCD pion pressure



## LQCD Nucleon GFFs

## Nucleon gluon GFFs, $\mathrm{m}_{\boldsymbol{\pi}} \sim 450 \mathrm{MeV}$

Dipole-like fall-off with momentum transfer




## Gluon momentum fraction

Gluon momentum fraction $A_{a}(0)=\langle x\rangle_{a}$


## LQCD Nucleon GFFs

## Nucleon gluon GFFs, $\mathrm{m}_{\boldsymbol{\pi}} \sim 450 \mathrm{MeV}$

Tripole-like fall-off with momentum transfer


Gluon GFFs: Shanahan, Detmold, PRD (2019) PRL (2019)
Quark GFFs: P. Hägler et al. (LHPC), PRD77, 094502 (2008)
Expt quark GFFs (BEG): Burkert et al, Nature 557, 396 (2018)


## Nucleon D-term GFFs

## Nucleon gluon GFFs, $\mathrm{m}_{\boldsymbol{\pi}} \sim 450 \mathrm{MeV}$

Tripole-like fall-off with momentum transfer


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Key assumptions in pressure extraction from DVCS

- Gluon D-term same as quark term in magnitude and shape
Factor of $\sim 2$ difference in magnitude, somewhat different tdependence
- Tripole form factor model LQCD results consistent with ansatz, but more general form is less well constrained
- Isovector quark D-term vanishes $D_{u-d}(t) \sim 0$ from other LQCD studies


## LQCD proton pressure

Nucleon pressure using LQCD results for quark and gluon GFFs, $\mathrm{m}_{\boldsymbol{\pi}} \sim 450 \mathrm{MeV}$


Gluon GFFs: Shanahan, Detmold, PRD (2019), PRL (2019)
Quark GFFs: P. Hägler et al. (LHPC), PRD77, 094502 (2008)

## LQCD + Expt proton pressure

Nucleon pressure using LQCD results for gluon GFF, JLab results for quark GFF


Gluon GFFs: Shanahan, Detmold, PRD (2019), PRL (2019)
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## LQCD proton shear



Gluon GFFs: Shanahan, Detmold, PRD (2019), PRL (2019) Quark GFFs: P. Hägler et al. (LHPC), PRD77, 094502 (2008)
Expt quark GFFs (BEG): Burkert et al, Nature 557, 396 (2018)

Tangential shear
vector field $4 \pi r^{2} T_{i j} \boldsymbol{e}_{j}^{\phi}$

| -0.02 | -0.01 | 0 | 0.01 | 0.02 | 0.03 |
| :--- | :--- | :--- | :--- | :--- | :--- |

## Gluon structure from LQCD

## LQCD calculations of proton and pion energy momentum tensor

- Gluon and quark gravitational form factors
- Shear and pressure distributions
- New physical mass calculations are ongoing
- Complements recent experimental studies
- Support analysis assumptions
- Suggest target kinematics for future model independent extractions at JLabl2 and EIC


## Next: pressure in nuclei

## Pressure in light nuclei c.f. pressure in the nucleon?



[^0]Quark GFFs: P. Hägler et al. (LHPC), PRD77, 094502 (2008)

## Next: pressure in nuclei

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## Gluon structure of nuclei

## How does the gluon structure of a nucleon change in a nucleus?

## European Muon

 Collaboration (1983): "EMC effect"Modification of per-nucleon cross section of nucleons bound in nuclei

Gluon analogue?

Ratio of structure function $F_{2}$ per nucleon for iron and deuterium

$$
F_{2}\left(x, Q^{2}\right)=\sum_{q=u, d, s \ldots} x e_{q}^{2}[\underbrace{\left.q\left(x, Q^{2}\right)+\bar{q}\left(x, Q^{2}\right)\right]}_{\substack{\text { Number density of } \\ \text { partons of flavour q }}}
$$



Longitudinal momentum fraction

## Nuclear glue, $\mathrm{m}_{\pi} \sim 450 \mathrm{MeV}$

Deuteron gluon momentum fraction
Look for nuclear (EMC-type) effects in the first moments of the spin-independent gluon structure function

## Doubly challenging

- Nuclear matrix element
- Gluon observable (suffer from poor signal-to-noise)

Ratio $\propto$ matrix element for $0 \ll \tau \ll t$


## Gluon momentum fraction

NPLQCD Collaboration PRD96 0945 I2 (20I7)

- Matrix elements of the spin-independent gluon operator in nucleon and light nuclei
- Present statistics: can't distinguish from no-EMC effect scenario
- Small additional uncertainty from mixing with quark operators

Ratio of gluon momentum fraction in nucleus to nucleon


## Gluon structure of nuclei

## Exotic Glue

Contributions to nuclear structure from gluons not associated with individual nucleons in nucleus

Exotic glue operator:
nucleon $\quad\langle p| \mathcal{O}|p\rangle=0$
nucleus $\langle N, Z| \mathcal{O}|N, Z\rangle \neq 0$
Jaffe and Manohar,"Nuclear Gluonometry" Phys. Lett. B223 (I989) 218

## Non-nucleonic glue in deuteron

NPLQCD Collaboration PRD96 0945 I2 (20I7)

First moment of gluon transversity distribution in the deuteron, $\mathrm{m}_{\boldsymbol{\pi}} \sim 800 \mathrm{MeV}$

First evidence for non-nucleonic gluon contributions to nuclear structure

- Hypothesis of no signal ruled out to better than one part in 107
- Magnitude relative to momentum fraction as expected from large- $\mathrm{N}_{c}$


Ratio $\propto$ matrix element for $0 \ll \tau \ll t$

Ratio of 3pt and 2pt functions



[^0]:    Gluon GFFs: Shanahan, Detmold, PRD (2019), PRL (2019)

