

# $\Upsilon$ production as a function of charged-particle multiplicity in pp collisions at $\sqrt{s} = 13$ TeV with ALICE

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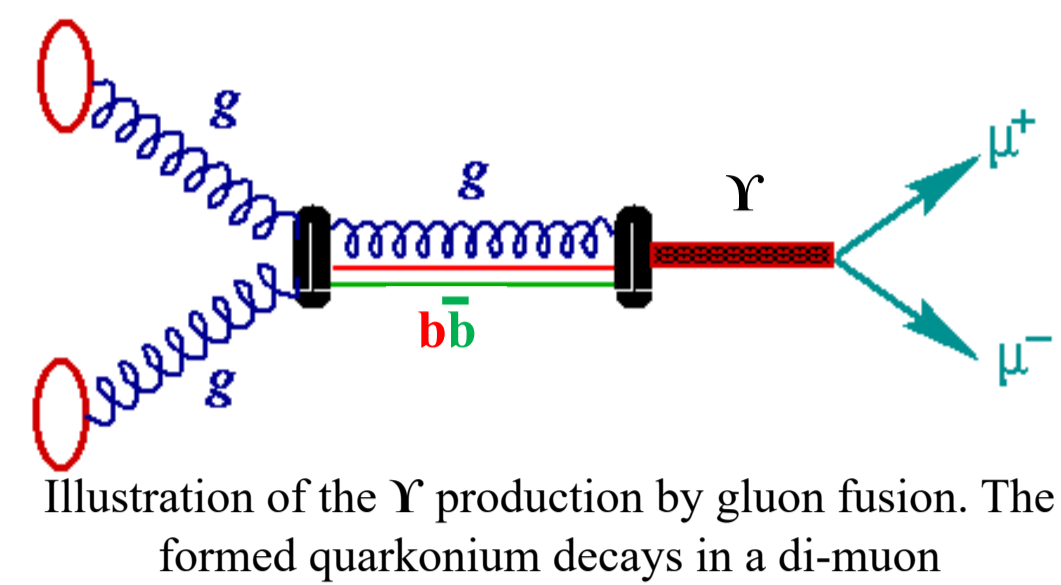
On behalf of the ALICE Collaboration



ALICE

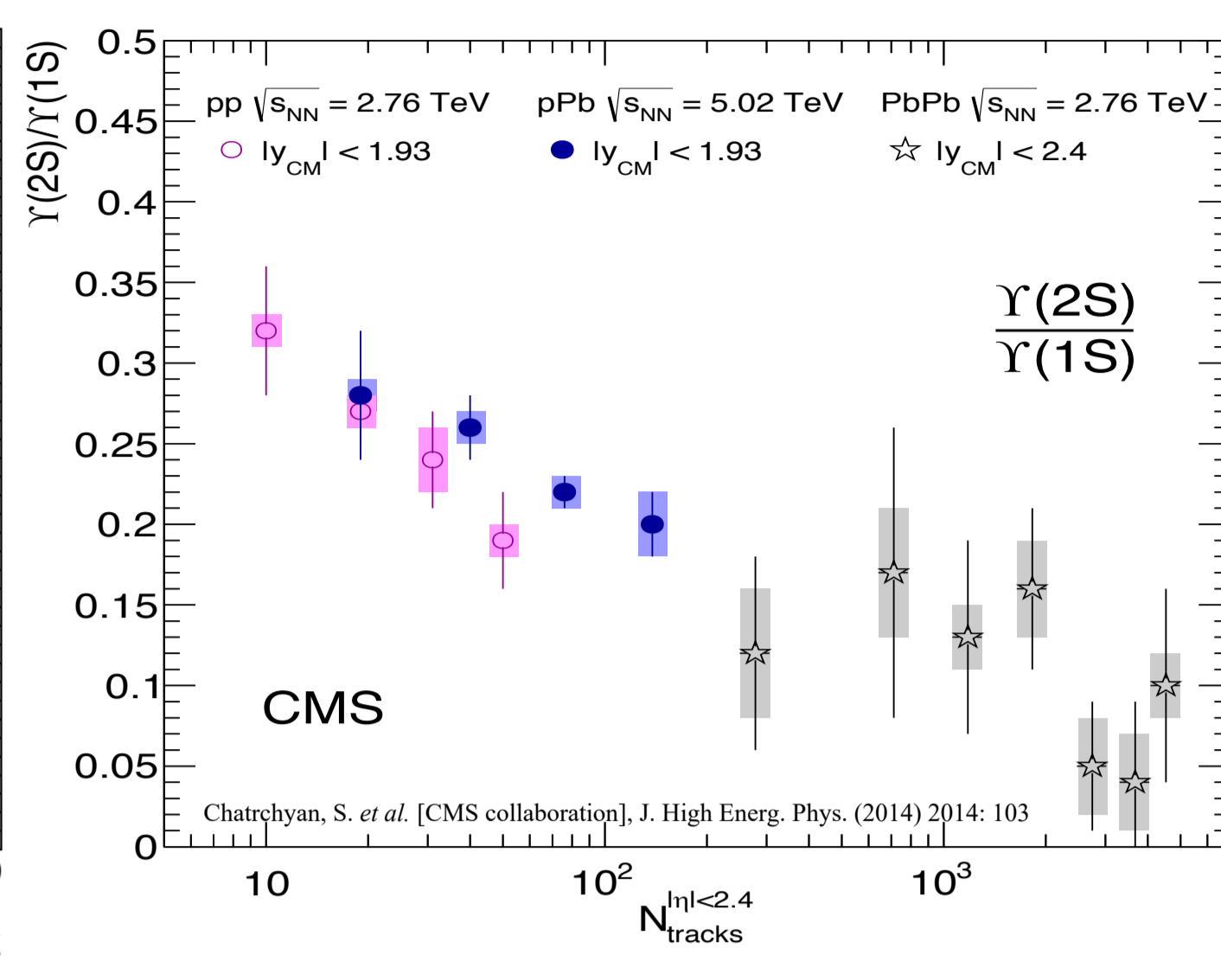
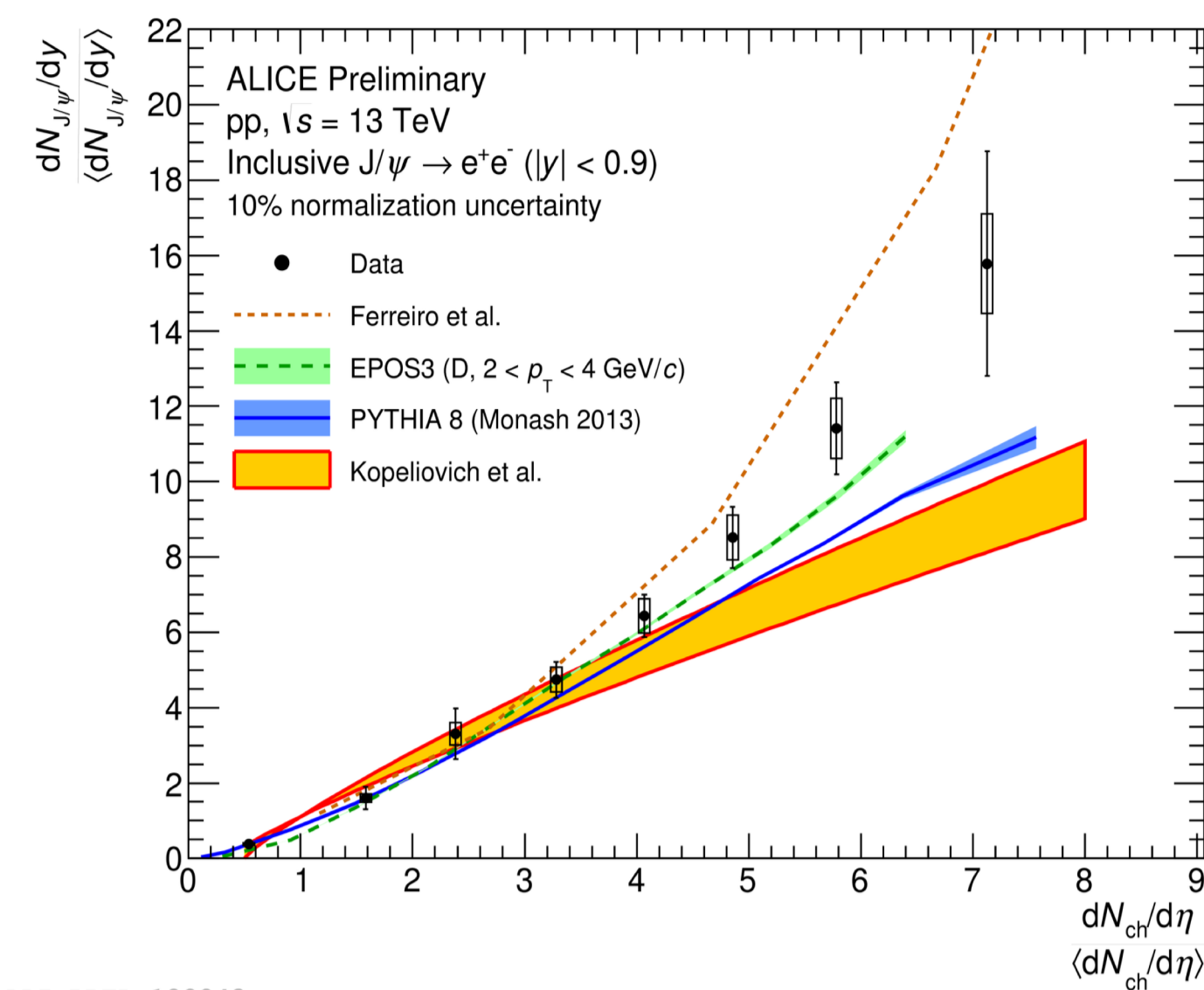
## 1. Quarkonium production as a function of charged-particle multiplicity

- Quarkonium: bound state of heavy  $Q\bar{Q}$  e.g. charmonium ( $c\bar{c}$ ), bottomonium ( $b\bar{b}$ )



- Quarkonium production in vacuum: not fully understood

- Theoretical models:
  - Colour-Evaporation Model (CEM) [1,2]
  - Colour-Singlet Model (CSM) [3]
  - Non Relativistic QCD (NRQCD) [4]



- Relative quarkonium yields increase more than linearly with increasing charged-particle multiplicity in the central rapidity region [5]
  - Linear increase: compatible with Multi-Parton Interaction (MPI) [5]
  - Non linear increase: additional effects are needed (percolation model, collectivity)[6,7]
- $\Upsilon$  excited-to-ground state ratio in the central rapidity region: decreasing trend with the event activity measured by CMS not well understood (initial-state vs. final-state effects) [8]
- Comparison between  $J/\psi$  and  $\Upsilon(1S)$  at forward rapidity as a function of multiplicity at mid rapidity can help to understand this suppression

### Inner Tracking System

- 6 layers (2 drifts, 2 strips, 2 pixels)
- Vertex measurement

### Silicon Pixel Detector (SPD)

(2 innermost ITS layers)

- Vertex determination
- Multiplicity estimator** in  $|\eta| < 1$  & in  $|z_{\text{vertex}}| < 10$  cm

A tracklet is a pair of hits in the two SPD layers aligned with the primary vertex within a fiducial window ( $\theta, \phi$ )

## 2. The ALICE detector

Schematic view of the ALICE detector [9]. Inner tracking system for multiplicity and muon spectrometer for  $\Upsilon$  measurement

### Muon Spectrometer

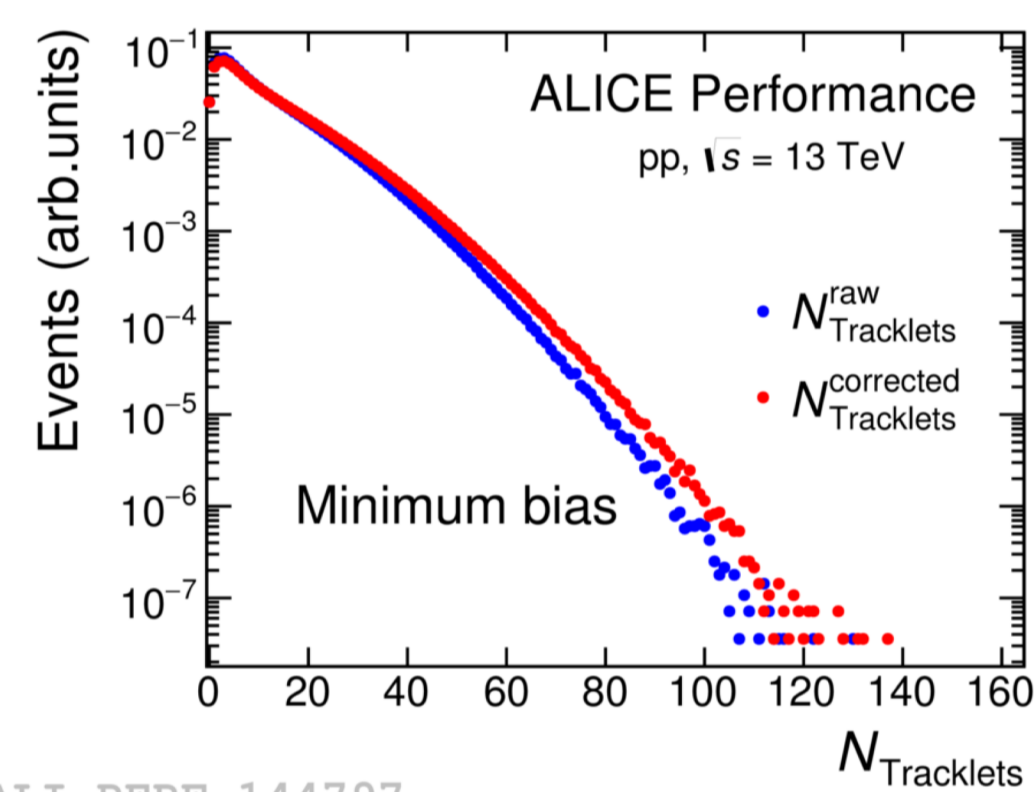
#### Di-muon ( $\mu^+\mu^-$ ) reconstruction with unlike-sign dimuon triggered events at forward rapidity

#### Analysis cuts:

- Single muons
  - Tracking-trigger matching ( $p_T$  trigger threshold  $\sim 0.5$  GeV/c)
  - $17.6 < R_{\text{DCA}} < 89.5$  cm
  - $p_{\text{DCA}}$  within  $6\sigma$
  - $-4 < \eta < 2.5$
- Di-muons
  - $2.5 < y < 4$
  - Total di-muon charge = 0

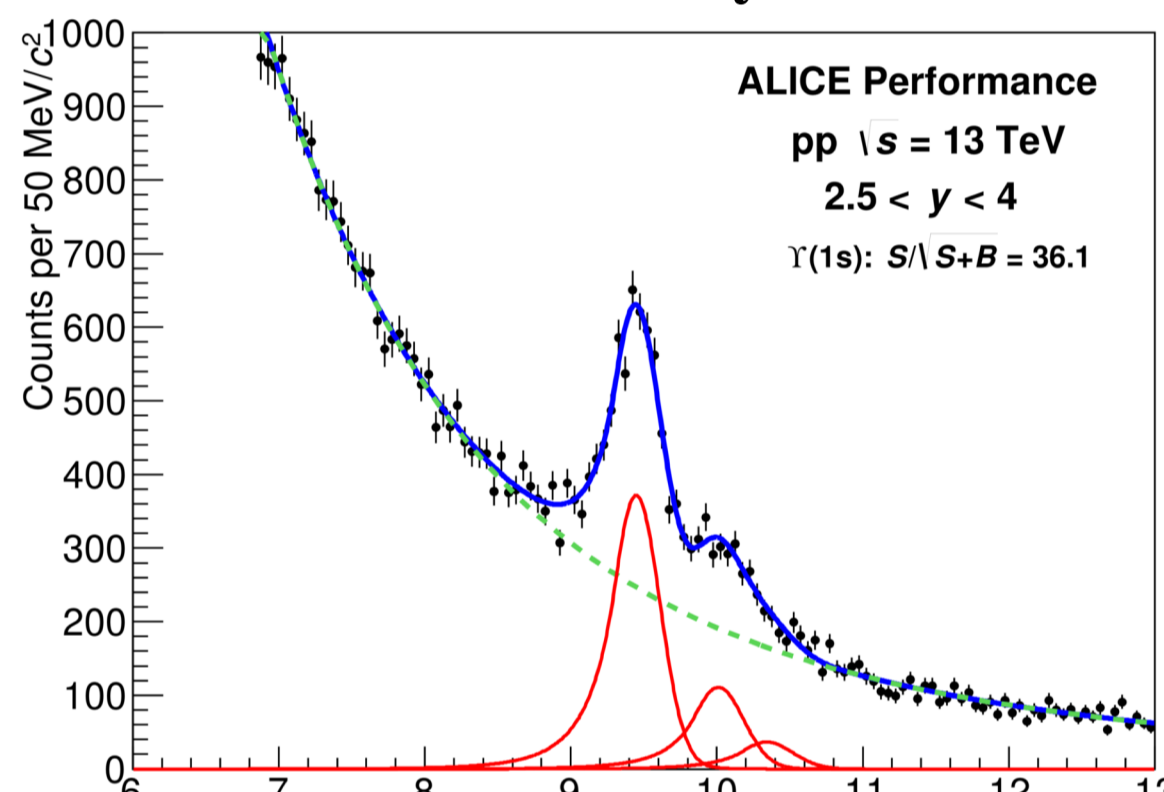
## 3. Analysis strategy

### Self-normalized charged-particle density:



- Multiplicity estimation ( $|\eta| < 1$ ):**
  - Multiplicity is estimated from reconstructed SPD tracklets
  - The raw tracklet ( $N_{\text{tracklet}}^{\text{raw}}$ ) distribution (blue data points) is corrected for the SPD inefficiency ( $N_{\text{tracklet}}^{\text{corrected}}$ , red data points)
  - Number of **tracklets** is approximately proportional to the **charged-particle multiplicity**: estimated from Monte Carlo simulations
  - Only INEL>0 (having at least one charged particle in  $|\eta| < 1$ ) events are selected for multiplicity measurement

### Self-normalized $\Upsilon$ yields:



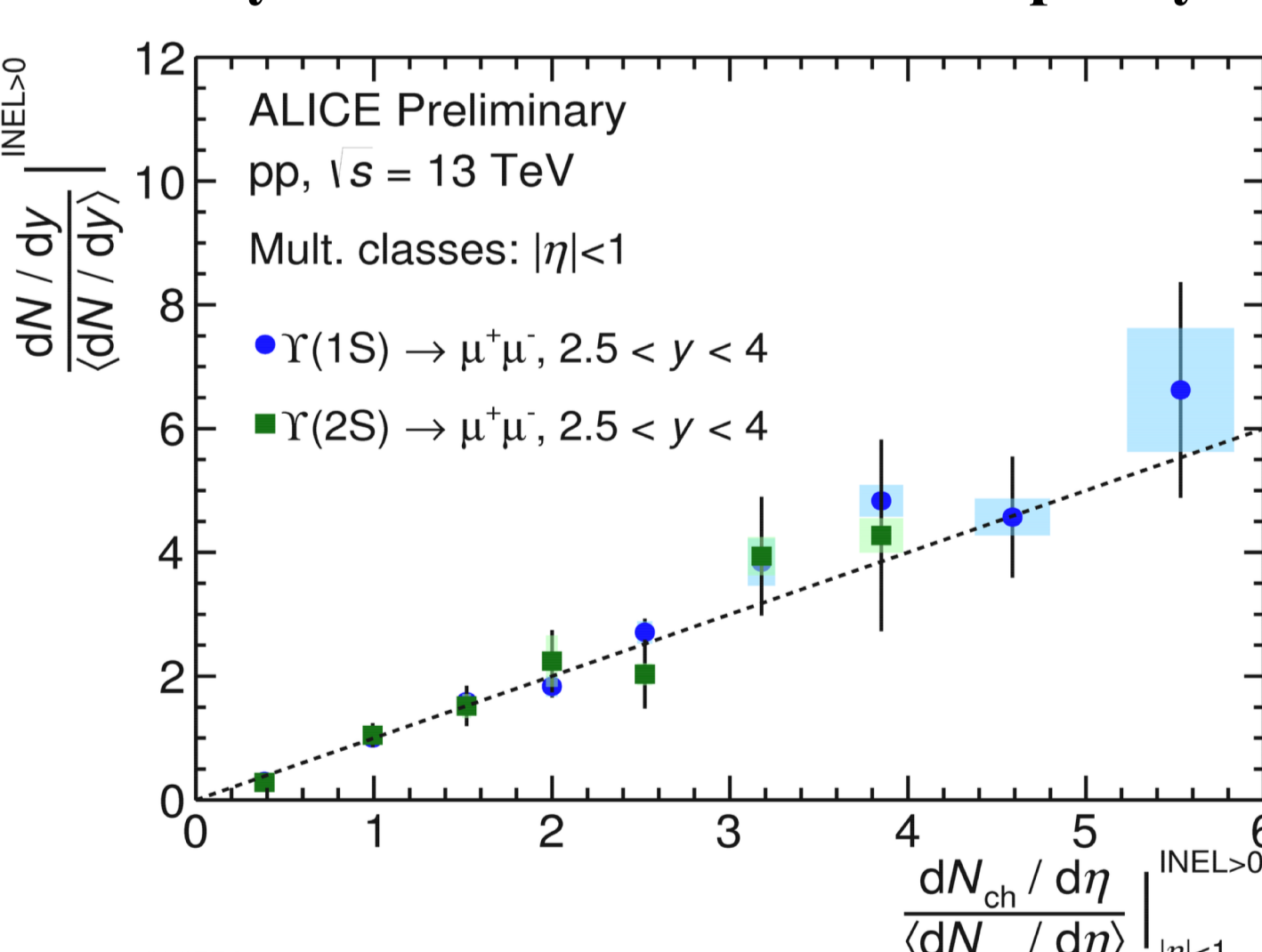
- $\Upsilon$  Signal extraction (in  $2.5 < y < 4$ ):**  $\Upsilon$  yield extracted from a fit to the di-muon invariant mass spectrum combining:
  - a signal function (Double Crystal Ball) for each  $\Upsilon$  state
  - a background function (variable width Gaussian/ product of two exponentials/ product of an exponential & a power law)
  - Systematic uncertainties are computed by varying fit conditions

- Self-normalized charged-particle density:** The average charged-particle density ( $dN_{\text{ch}}/d\eta$ ), in multiplicity bin  $i$  is normalized by the average charged-particle density from the integrated case ( $\langle dN_{\text{ch}}/d\eta \rangle_{\text{total}}$ )
- Systematic uncertainties for charged-particle multiplicity are computed by varying different MC inputs and vertex ranges
- Efficiency factors ( $\epsilon$ ) and corresponding systematic uncertainties are applied for trigger selection and INEL>0 selection

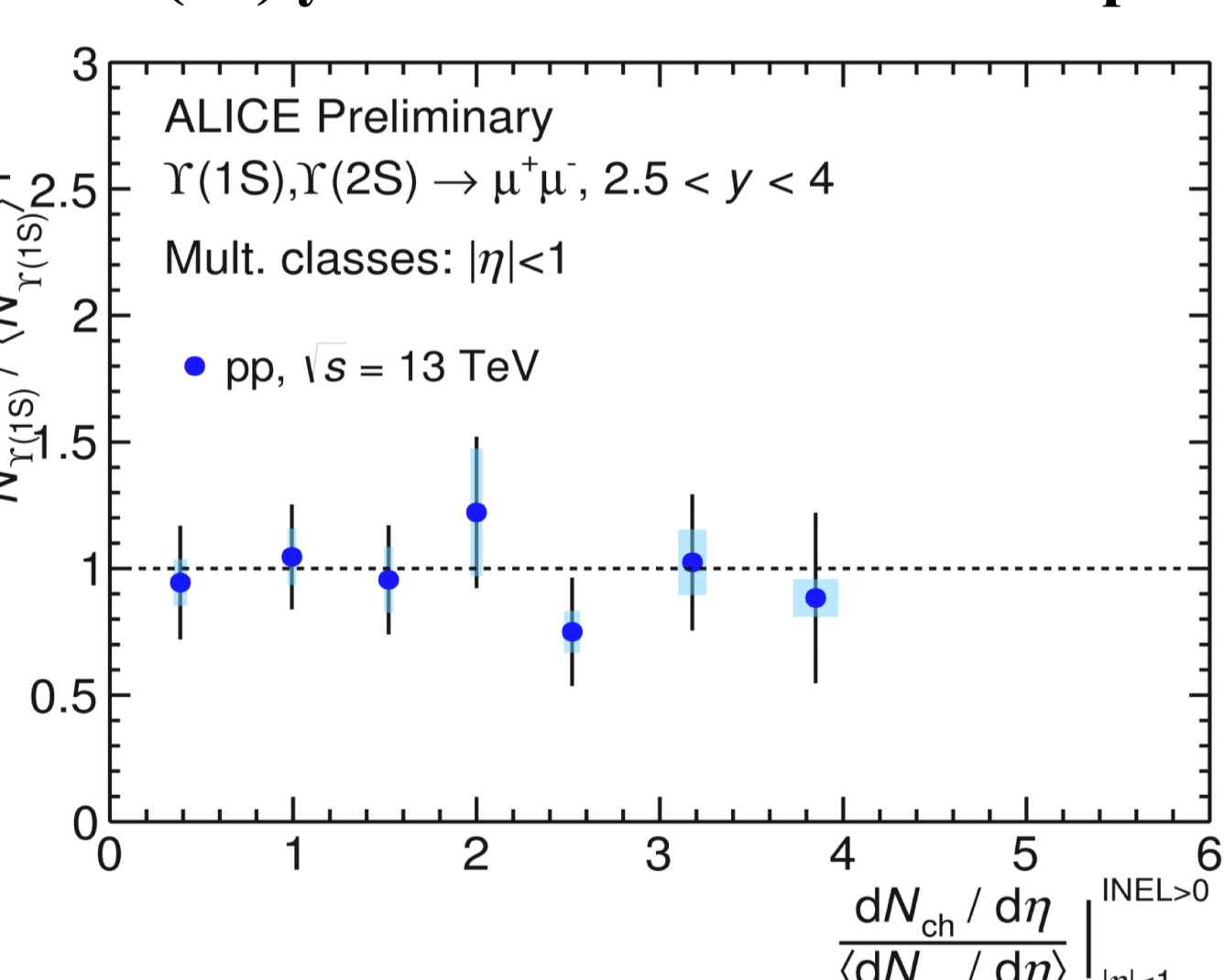
- The  $\Upsilon$  yield obtained from signal extraction in each multiplicity interval is normalized by INEL>0 events in the corresponding multiplicity interval. Afterwards it is divided by multiplicity integrated  $\Upsilon$  yield (which is also normalized by total INEL>0 events).
- Efficiency factors ( $\epsilon$ ) and corresponding systematic uncertainties are applied to account for trigger and INEL>0 selection

## 4. Results

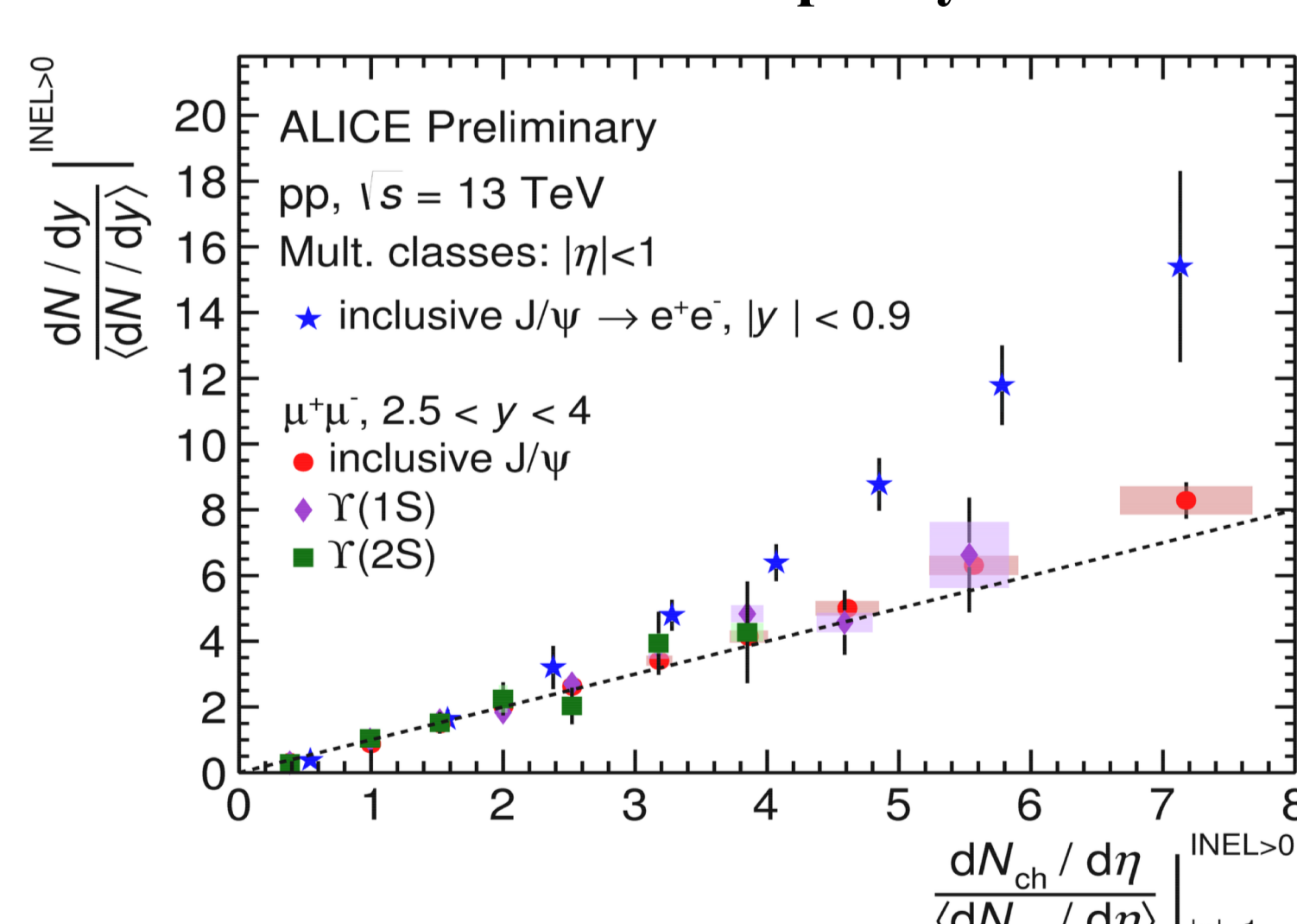
### Self-normalized $\Upsilon(1S)$ & $\Upsilon(2S)$ yields as a function of multiplicity:



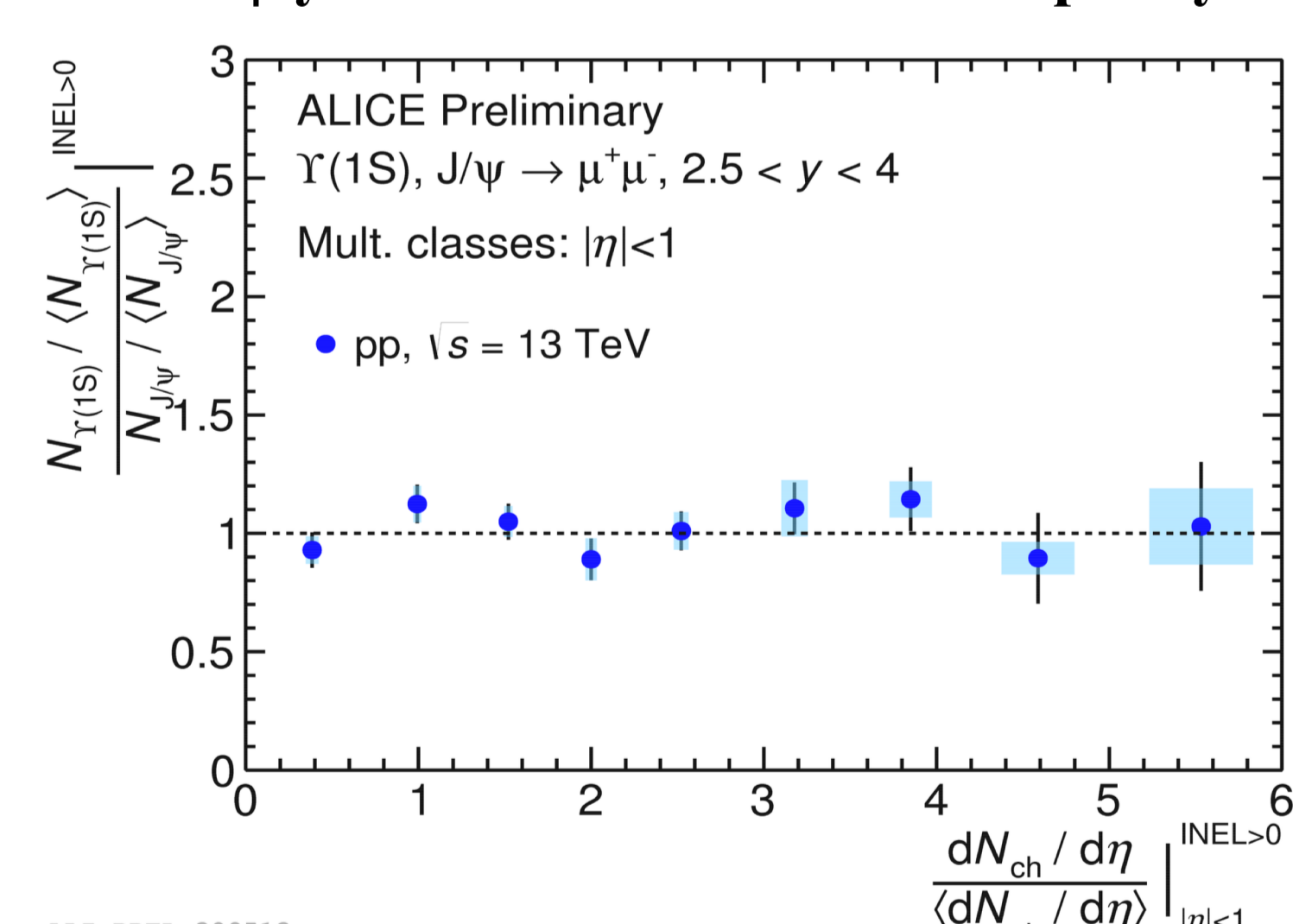
### Ratio between self-normalized $\Upsilon(2S)$ & $\Upsilon(1S)$ yields as a function of multiplicity:



### Self-normalized quarkonium yields as a function of multiplicity:



### Ratio between self-normalized $\Upsilon(1S)$ & $J/\psi$ yields as a function of multiplicity:



- Statistical errors are shown in bars and systematic uncertainties are represented by the boxes
- A linear increase is observed for both of  $\Upsilon$  states and  $J/\psi$  measured at forward rapidity
- The self-normalized ratio of  $\Upsilon(2S)$  over  $\Upsilon(1S)$  as a function of multiplicity is constant, compatible with unity within the uncertainties, up to a relative multiplicity of about 4
- The self-normalized ratio of  $\Upsilon(1S)$  over  $J/\psi$  as a function of multiplicity is constant, compatible with unity within the uncertainties, up to a relative multiplicity of about 6
- A faster than linear increase was observed for  $J/\psi$  measured at mid-rapidity [5]
- A 20% decrease in  $\Upsilon(2S) / \Upsilon(1S)$  as measured in the central rapidity region by CMS [8] is still compatible within current uncertainties

## 5. Conclusions

- First ALICE measurement for  $\Upsilon(1S)$  &  $\Upsilon(2S)$  as a function of charged-particle multiplicity in pp collisions at  $\sqrt{s} = 13$  TeV
- Quarkonia ( $J/\psi$ ,  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ ) measured at forward rapidity as a function of charged-particle multiplicity measured at mid rapidity show a linear increase with increasing multiplicity
- The self-normalized ratio of the  $\Upsilon(2S)$  over  $\Upsilon(1S)$  and  $\Upsilon(1S)$  over  $J/\psi$  are compatible with unity as a function of multiplicity
- No effects are seen w.r.t the resonance mass and quark content at forward rapidity
- Introducing a rapidity gap between the hard probe and the multiplicity estimator reveals a linear behavior which is different from previous observations without rapidity gap

## References

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