# Primordial Gravitational Waves in Non-standard Cosmologies

Based on NB & Fazlollah Hajkaim – arXiv:1905.10410 [astro-ph.CO]

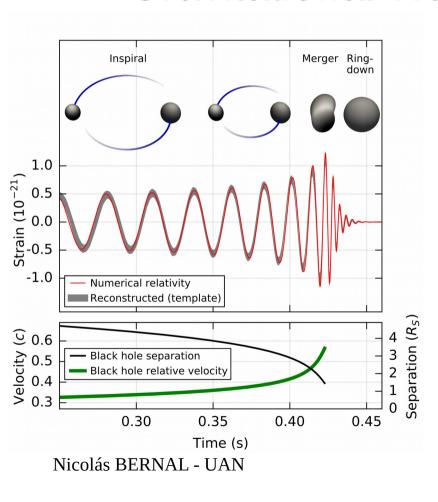
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CoCo 2019 May 31<sup>st</sup>, 2019

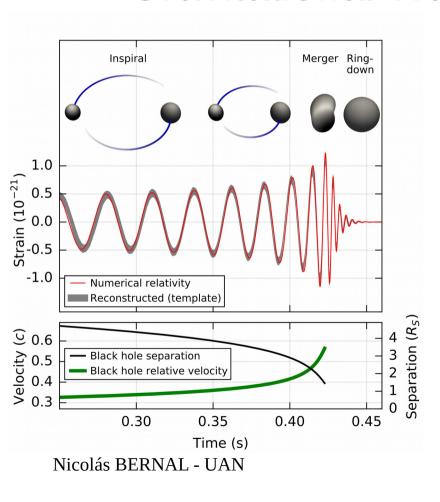


#### Gravitational Waves have been detected!



- O(10) solar mass black hole exists
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- NS-NS merger test NS paradigm
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- We can observe the early Universe through GWs!

#### GWs as probes of the early Universe

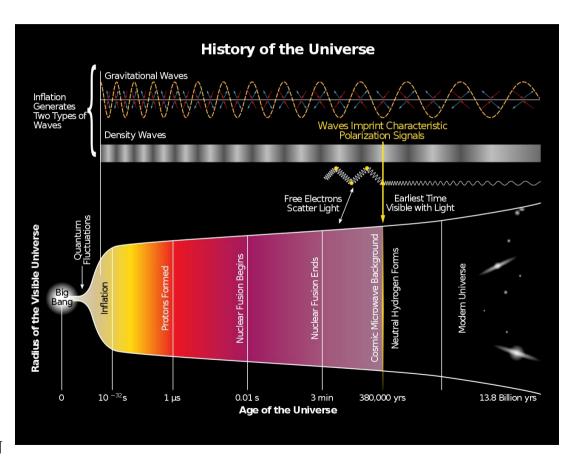
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## The Early Universe





#### Primordial Gravitational Waves

- GW are represented by spatial metric perturbations (transverse & traceless).
- The evolution of GWs is described by the linearized Einstein eq.

$$\ddot{h}_{ij} + 3H \, \dot{h}_{ij} - \frac{\nabla^2}{a^2} h_{ij} = 16\pi \, G \, \Pi_{ij}^{TT}$$

 $\Pi^{TT}$  is the transverse-traceless part of the anisotropic stress tensor  $\Pi_{ij} = \frac{T_{ij} - p g_{ij}}{T_{ij}}$ 

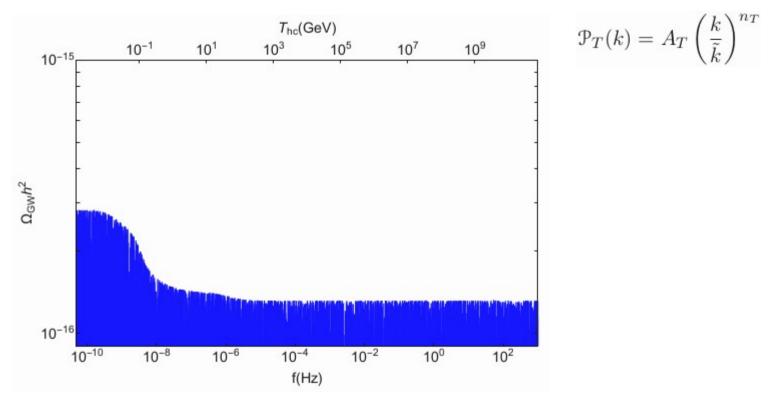
$$\Pi_{ij} = \frac{T_{ij} - p \, g_{ij}}{a^2}$$

Primordial GW spectrum

$$\Omega_{\rm GW}(t,k) = \frac{1}{\rho_c(t)} \frac{d\rho_{\rm GW}(t,k)}{d\ln k} = \frac{1}{12 a^2(\eta) H^2(\eta)} \mathcal{P}_T(k) \left[ X'(\eta,k) \right]^2$$

with the primordial tensor power spectrum  $\mathfrak{P}_T(k) = A_T \left(\frac{k}{\tilde{\iota}}\right)^{NT}$ and the transfer function  $\frac{d^2X(u)}{du^2} + \frac{2}{a(u)} \frac{da(u)}{du} \frac{dX(u)}{du} + X(u) = 0.$ 

#### Primordial Gravitational Waves



For a primordial scale invariant spectrum and standard cosmology

### Standard Cosmology

 Total energy density dominated by SM radiation from end of inflation until matter-radiation equality



→ SM energy density:  $\rho_R \sim a^{-4}$ 

→ Photon temperature:  $T \sim a^{-1}$ 

Hubble expansion rate

$$\rightarrow$$
 H  $\sim$  T<sup>2</sup>/M<sub>P</sub>

$$\rightarrow$$
 H ~ a<sup>-2</sup>

## Non-standard Cosmologies

- The energy density of the Universe was dominated by  $\phi$  with an equation of state  $\omega_\phi \equiv p_\phi$  /  $\rho_\phi$

$$\begin{split} \frac{d\rho_{\phi}}{dt} + 3(1+\omega_{\phi}) H \, \rho_{\phi} &= -\Gamma_{\phi} \, \rho_{\phi} \,, \\ \frac{ds_R}{dt} + 3 \, H \, s_R &= + \frac{\Gamma_{\phi} \, \rho_{\phi}}{T} \end{split}$$

$$\rho_R(T) = \frac{\pi^2}{30} g_{\star}(T) T^4$$

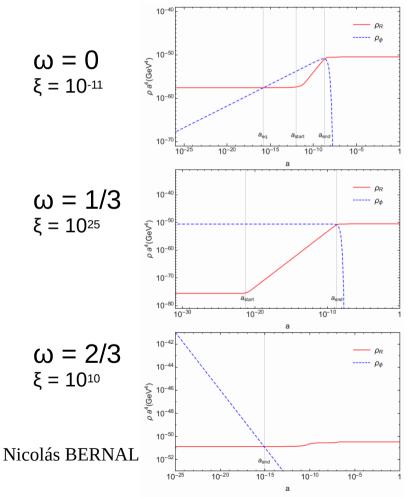
$$s_R(T) = \frac{\rho_R + p_R}{T} = \frac{2\pi^2}{45} h_{\star}(T) T^3$$

$$H^2 = \frac{\rho_{\phi} + \rho_R + \rho_m + \rho_{\Lambda}}{3 M_{Pl}^2}$$

• Initial conditions: 
$$T_{end}$$
 and  $\xi \equiv \frac{\rho_{\phi}}{\rho_{R}}\Big|_{T=T_{max}}$ 

with  $T_{max}=10^{14}$  GeV (pivot scale)

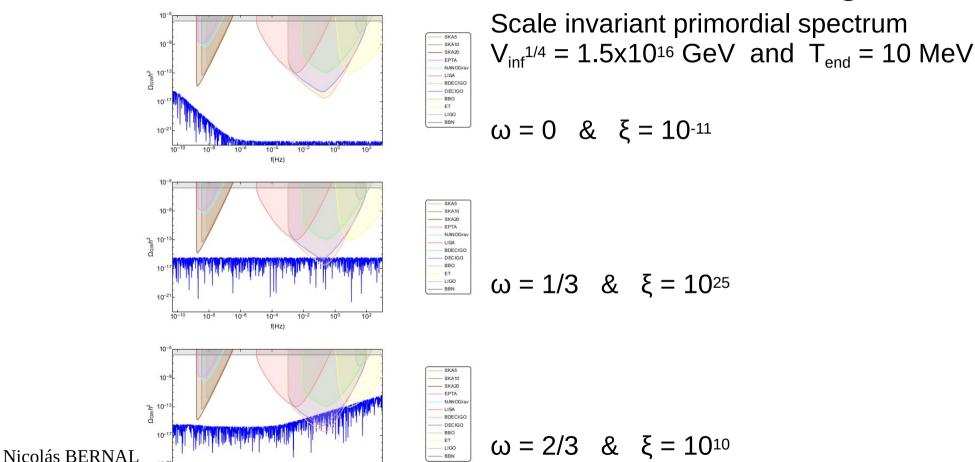
## Non-standard Cosmologies

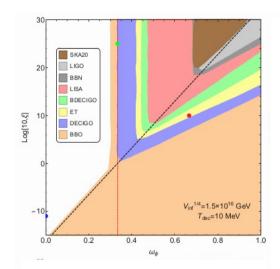


$$\rho_{\phi}(a) \propto a^{-3(1+\omega_{\phi})}$$

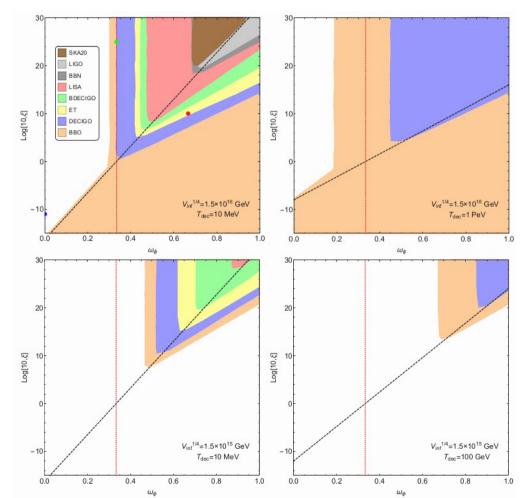
$$\rho_R(a) \propto \begin{cases} a^{-4} & \text{for} \quad a \ll a_{\text{start}}, \\ a^{-\frac{3}{2}(1+\omega_{\phi})} & \text{for } a_{\text{start}} \ll a \ll a_{\text{dec}}, \\ a^{-4} & \text{for } a_{\text{dec}} \ll a, \end{cases}$$

$$T(a) \propto \begin{cases} a^{-1} & \text{for} \quad a \ll a_{\text{start}}, \\ a^{-\frac{3}{8}(1+\omega_{\phi})} & \text{for } a_{\text{start}} \ll a \ll a_{\text{dec}}, \\ a^{-1} & \text{for } a_{\text{dec}} \ll a. \end{cases}$$

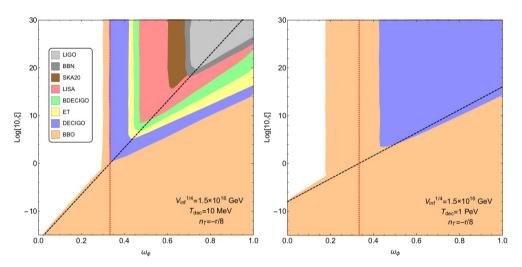




Scale invariant primordial spectrum



Scale invariant primordial spectrum



Consistency relation in the single-field slow-roll scenario:

$$r = -8 \, n_T$$
.

where

$$r \equiv \frac{A_T}{A_S}$$

and  $V_{inf}^{1/4} = 1.5 \times 10^{16} \text{ GeV}$ 

#### Conclusions & Outlook

- GWs offer a new messenger to explore the early Universe, beyond EM waves.
- The existence of a PGW background is one of the most crucial predictions of the inflationary scenario of the early Universe.
- The spectrum of the inflationary GWs depends on:
  - \* the power spectrum of primordial tensor perturbations
  - \* the expansion rate of the Universe.
- GWs can probe the early Universe!
- Upcoming GW observatories could shed light on the cosmological properties of the early Universe.

## Muchas gracias!

